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# Multiscale 3D analysis of creep cavities in AISI type 316 stainless steel

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A sample of AISI type 316 stainless steel from a power station steam header, showing reheat cracking, was removed from service and has been examined by a combination of microscale Xray computational tomography, nanoscale serial section focused ion beam-scanning electron microscopy, elemental energy dispersive X-ray mapping and transmission electron microscopy. Multiscale three-dimensional analysis using correlative tomography allowed key regions to be found and analysed with high resolution techniques. The grain boundary analysed was decorated with micrometre sized, facetted cavities,  $M_{23}C_6$  carbides, ferrite and G phase but no  $\sigma$  phase. Smaller intragranular M<sub>23</sub>C<sub>6</sub> particles were also observed, close to the grain boundaries. This intimate coexistence suggests that the secondary phases will control the nucleation and growth of the cavities. Current models of cavitation, based on isolated idealised grain boundary cavities, are oversimplified.

Keywords: Stainless steels, Creep cavitation, 3D imaging, Correlative tomography,

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# Introduction

Nickel based alloys and heat resistant steels are generally the materials of choice for the construction of engineered components and plant operating at high temperatures. The creep lifetime of such materials is dependent on the applied load history, temperature and creep properties.<sup>1</sup> Various methods for estimating the remaining creep life of components have been developed including ductility exhaustion,<sup>2</sup> life fraction and phenomenological approaches incorporating cavitation or microstructural damage.3

Welding is widely used in fabricating systems for high temperature operation in the nuclear, aeroengine and oil and gas sectors. However, creep lifetimes can be reduced by the associated residual stresses.<sup>4</sup> Furthermore, the transient thermal excursions arising from welding may produce suboptimal and inhomogeneous microstructures<sup>3</sup> as well as interactions with the environment, for example, with carbon dioxide.<sup>6</sup> Welds joining many heat resistant alloys, such as type 300 series stainless steels, alloy 800H, alloy 52 and alloy 617, are susceptible to reheat cracking. This creep failure mechanism is variously described in the literature as ductility dip cracking, relaxation cracking,

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stress relief cracking, post-weld heat treatment cracking and strain age cracking.

In the 1990s, widespread reheat cracking was found in non-stress relieved welded AISI type 316H austenitic stainless steel components in several advanced gas cooled reactors (AGRs) in the UK.<sup>7</sup> Interestingly, some of these were driven solely by the presence of residual stresses. Cracking in the heat affected zone (HAZ), a few millimetres from the weld fusion boundary, was caused by a combination of thermal relaxation of residual stress combined with very low material creep ductility (a fraction of 1%) at the operating temperature. The creep ductility of the HAZ of these welds was reduced owing to three effects: the thermomechanical plastic strains associated with welding, slow creep deformation rates such as those occurring under plant operational conditions<sup>8</sup> and a highly triaxial residual stress state associated with the weld geometry.9 A ductility exhaustion life assessment model has been developed by the industry to predict the onset of reheat crack initiation in welded type 316H stainless steel components.<sup>2,8</sup> Within this model, an empirical approach incorporating cavity nucleation and growth terms was adopted to describe the substantial effect of stress triaxiality on ductility.<sup>9</sup> The model has been shown to overpredict damage for welded feature tests and laboratory specimens<sup>4</sup> and continues to be applied to support operation of UK nuclear plants. However, the model lacks any direct functional relationship with the underlying damage development mechanisms. For example, it does not explain why reheat cracking appears to occur on a plane parallel to the fusion boundary, at initiation, rather than normal to the maximum principal macroscopic stress

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direction (which acts along the welding direction and parallel to the crack plane), nor why the ductility of the HAZ can be an order of magnitude lower than the parent material.

The length scales of creep damage increase with time as the governing process changes from being diffusion driven, through cavity nucleation at grain boundaries and vacancy condensation around second phase particles, to cavity growth and linkage at the grain scale, followed ultimately by multigrain fissure and failure. For example, in AISI type 316 austenitic stainless steel, vacancy diffusion and dislocation mechanisms occur at subnanometre dimensions, and the minimum stable cavity diameter (i.e. nucleation size) is  $\sim 30$  nm assuming a stress of 300 MPa (at 500°C) and a surface tension of 2.41  $\text{N-m}^{-1}$  for stainless steel.<sup>10</sup> Evidence from reheat cracking studies using conventional optical microscopy suggests that such cavities often form around grain boundary carbides.<sup>11</sup> As cavities grow, they tend to coalesce to form microcracks having lengths of the order of the material's grain size (typically 50-100 µm), which then link up to form macrocracks a few millimetres long that can be detected and sized by non-destructive testing.

Methods used to observe and characterise creep damage in metals can be classified as either destructive or non-destructive and further subdivided into surface [two-dimensional (2D)] and volumetric [three-dimensional (3D)] methods. The capability of any technique is determined by its spatial resolution, the contrast mechanism, its sampling volume and by practical issues. Traditional procedures for direct optical microscopic examination of the surface of metallographic specimens can be applied to provide information about the location, size distribution and orientation of cavities and cracks greater than  $\sim$ 500 nm. At finer scales, scanning electron microscopy can, in theory, reveal details to a resolution of a few nanometres, but in practice, the size of cavities that can be observed is limited by the efficacy of traditional metallographic preparation techniques. At the highest resolutions, transmission electron microscopy (TEM) can be used to characterise very small cavities but the volume sampled is very small and specimens are difficult to prepare.

Small angle scattering (SANS) is a semidestructive technique, which can give population distributions of cavities and carbides in the size range of  $10-500 \text{ nm}^{11-13}$  and thereby provide indirect 3D observations typically in 20 mm<sup>3</sup> sample volumes. However, observations have to be validated by other techniques to confirm the nature of the features being characterised.

New techniques for imaging the microstructure of material in three dimensions have emerged in recent years, most notably X-ray computational tomography (CT)<sup>14</sup> and dual beam focused ion beam (FIB)–scanning electron microscopy (SEM) serial sectioning. For example, Pyzalla *et al.*,<sup>15</sup> Gupta *et al.*,<sup>16</sup> Sket *et al.*<sup>17</sup> and Isaac *et al.*<sup>18</sup> have applied X-ray CT to map creep cavitation damage at the micrometre scale. To cover a larger range of scales, region of interest (RoI) (also known as zoom or local tomography) can be performed.<sup>14,19</sup> While conventional X-ray CT can map creep damage across relatively large volumes, it is insensitive to the grain boundaries and crystallographic relationships. Serial

backscattered diffraction as well as chemical mapping in 3D by energy dispersive X-ray (EDX). By incorporating both within a correlative tomography framework,<sup>21</sup> we can survey damage at the larger scale and then select RoIs for higher resolution study by FIB tomography, ensuring that the information thus obtained is statistically and spatially relevant in terms of the macrostructure of the creep damage (e.g. the final creep crack).

In the present paper, we focus on a particular joint from a pressure vessel (a boiler steam header) removed from an operational AGR plant owing to the development of extensive reheat cracking.<sup>7</sup> This component and cracking mode have been the focus of previous detailed examination by conventional optical microscopy, small angle neutron scattering  $^{11-13}$  and synchrotron tomography.<sup>22</sup> Using FIB serial sectioning of RoIs located by X-ray microtomography, we explore in detail the local environment of the creep cavities, in an effort to understand how this may have affected their nucleation, growth and eventual coalescence, which led to macroscopic reheat cracking. Previous work<sup>12</sup> has suggested that there are significant variations in carbide and cavity populations local to the reheat crack (see Fig. 1), and in the present study, we look at cavitation at just one RoI. Ultimately, we hope to expand our study to look at the whole region because it is our aim to understand the coevolution of the secondary phases and cavitation damage and develop functional relationships between the nature of internal damage and macroscopic quantities such as stress, strain, temperature and time in order to improve methods of creep life prediction.

# Experimental

#### Material composition and history

An AGR power station steam header (1C1/1 cast 55919) operating under internal pressure loading of 16 MPa at ~525°C was removed from service, after 65 000 h, following the detection of extensive creep cracking by non-destructive surface and ultrasonic volumetric inspections. A cross-sectional slice of material containing the cracked region was excised from the header to nozzle weld region, as indicated in Fig. 1. The header was made from AISI type 316H austenitic stainless steel with the composition shown in Table 1 and had a grain size of  $85 \pm 10 \ \mu\text{m}$ . A ~ 10 mm<sup>2</sup> and ~ 8 mm thick coupon was extracted from a location before the crack mouth and the outer surface of the header for the present studies. The coupons were cut into a series of nominally  $0.5 \text{ mm}^2$ cross-section matchsticks using a fine SiC circular saw. Multiple matchsticks were cut from each coupon, and their position within the coupon was recorded. The HAZ material associated with the nozzle attachment weld in this coupon has been previously identified as a region prone to high creep cavitation according to both numerical modelling and SANS experiments<sup>12</sup> (referred to as region number 12 in their work).

#### Correlative tomography strategy

The experimental method employed here has been termed 'correlative tomography'.<sup>21</sup> It describes a multi-scale 3D method of characterisation where the same RoI



1 Schematic diagram showing where sample was extracted from ex-service steam header. Coupon was originally extracted close to crack origin from which smaller matchstick was extracted, which is subject of studies here. Coordinate system shows relative orientation of matchstick with respect to original component and crack. Schematic of matchstick sample shows region of X-ray CT analysis, 3D visualisation of which is shown in Fig. 2a

is interrogated at different scales and using different imaging modalities. Here, X-ray CT data are used to survey a large volume, in this case, a millimetre sized matchstick of material, but there is the potential to image the whole region local to the reheat crack at a lower magnification first before continuing to progressively higher resolution analyses. The role of the large volume X-ray CT scan on the matchstick was threefold: to provide an overview of the sample and the features of interest visible at this scale, to provide an intermediate scale linking the higher resolution analysis to the component scale and to provide a means for identifying and locating RoIs for higher resolution investigation in 3D using FIB-SEM serial sectioning. Once an RoI was identified (in this case, two RoIs were identified) by nondestructive X-ray CT (Fig. 2a and b) for higher resolution examination, features on the sample surface that were visible in the X-ray CT data were located in the FIB-SEM and used as a means of spatially correlating the two datasets (for further details, see Burnett et al.).<sup>21</sup> Once correlated, it was possible to import knowledge about the location of the subsurface

Table 1 Composition of AISI type 316H austenitic stainless steel header (while there is no direct confirmation of Nb content for this header, other headers from same cast have had their Nb content measured as 0.06 wt-%)

Composition/wt-%	Header	Weld	
С	0.058	0.035	
Si	0.42	0.46	
Mn	1.61	0.94	
Р	0.009	0.011	
S	0.020	0.010	
Cr	17.33	18·36	
Ni	11.29	11.36	
Мо	2.31	2.91	
Cu	0.12	0.04	
V	0.04	0.06	
Ti	<0.01	0.02	
AI	<0.01	0.01	
Со	<0.01	0.03	
Ν	0.041	0.056	

RoI from the X-ray CT data into the coordinate system of the FIB-SEM. With the location of the RoIs defined in the FIB-SEM, the RoIs could be progressively revealed by serial sectioning using the FIB (Fig. 2c). In this manner, the same RoI can be imaged by X-ray CT and SEM, providing higher resolution insights and enabling crystallographic and potentially chemical information to be added to the picture. Once serial sectioning was complete, a thin slice of material was removed for analysis in the scanning transmission electron microscope (STEM) at an even higher resolution analysis (Fig. 2d) and crucially providing information on the chemistry and crystallography of the RoIs.

#### Microscale X-ray computed tomography

X-ray CT was carried out on the 'TOMCAT' synchrotron beamline at the Swiss Light Source in order to characterise the defect structure within a matchstick sample (of cross-section  $\sim 0.6 \times 0.6$  mm) cut from a coupon located close to the reheat crack from the side nearest the crack (see Fig. 1). Projections (1501) over  $360^{\circ}$  were collected with energy of 38 keV and a 2 s exposure time per frame. The source detector distance was  $\sim 25$  m, and the sample detector distance was as small as possible (5 mm). The 3D volume was reconstructed, using a Feldkamp-Davis-Kress reconstruction at a  $0.37 \ \mu\text{m}^3$  voxel size (the smallest achievable for this set-up), before visualisation in 3D using Avizo software (see Fig. 2a). This 3D map was used to visualise the defects throughout the sample at the micrometre scale before more detailed analysis by electron microscopy.

# Nanoscale serial section tomography

Two RoIs near the surface were identified from the Xray CT dataset for further investigation: one containing prolate spheroidal feature that was ultimately identified as a stringer and one containing a 'dark line' feature (*see* Fig. 2a and b). After spatial correlation of the X-ray CT data to the sample surface imaged with an FEI Helios 650 Dualbeam FIB-SEM, the locations of the RoIs could be identified and a physical cross-section of each RoI was prepared using the Ga ion FIB. The crosssections were then imaged using the electron beam



2 *a* left shows three unprocessed orthogonal virtual slices from reconstructed X-ray CT dataset. Total size of this volume was  $0.74 \times 0.6 \times 0.6$  mm. Subvolume is indicated, which has been segmented for 3D visualisation of visible features. Two principal features are that of 'stringers' and 'dark lines'. Width of dark lines has been exaggerated in segmented volume to make them clearly visible at this scale. Approximate location of two Rols chosen for further analysis is indicated. *b* Rol<sub>1</sub> shows large stringer, which was chosen for further analysis (image processed with edge preserving smoothing and non-local means filters). Rol<sub>2</sub> shows dark line, which was chosen for further analysis; other stringers are visible in this virtual slice and indicated by blue arrows (image processed with adaptive histogram equalisation and median filter). For both images, region for FIB-SEM analysis is indicated by dashed red outline. *c* image (SEM) (using mixed signal of secondary and backscatter electrons via through lens detector) of FIB prepared cross-section of two same Rols imaged via X-ray CT in *b*. Already for Rol<sub>13</sub> feature can be identified as solid multiphase feature sitting on grain boundary and Rol<sub>2</sub> can be identified as a string of cavities. *d* image (SEM) showing detail of two-same Rols once TEM lamella of the feature had been extracted. Rol<sub>1</sub> can now be clearly seen solid multiphase feature with some parts of structure charging under electron beam. Rol<sub>2</sub> shows detailed morphology some individual cavities that lie along boundary as well as wealth of secondary phases, which also exist there

and the through lens detector, collecting a mixture of secondary and backscatter electrons. A degree of caution was applied to the positioning of the first cross-section excised so as to just fall short of the feature of interest. From this point, material was successively removed from the cross-section to uncover the features of interest. Any mislocation of the feature either laterally or with depth could be addressed at this stage to position

1

2



3 *a* post-processed (edge preserving smoothing and non-local means filtering) virtual slice from reconstructed X-ray CT data showing stringer in Rol<sub>1</sub> and *b* SEM image of same stringer in Rol<sub>1</sub> taken with through lens detector on extracted TEM lamella

the feature of interest in the centre of the cross-section where possible (*see* Fig. 2c). Once this initial crosssection had been prepared, this became the starting point for destructive serial sectioning of the feature using the Auto Slice and View software, whereby a sequence of cross-sections is prepared one after another using the FIB with each cross-section imaged by the SEM,<sup>23</sup> creating a high resolution 3D image volume (Fig. 4). A nominal slice thickness of 20 nm and a pixel size of x=8.8 nm and y=11.2 nm was used for the present work.

After the serial sectioning, but before completely slicing the entire feature, a thin ( $<5 \mu$ m) lamella of material from the same region was extracted for analysis in the STEM (*see* Fig. 2*d*). Once extracted, further thinning, using the FIB, was required to produce an electron transparent lamella for STEM imaging.

We have been careful to assess the damage introduced by the FIB-SEM serial sectioning. The fact that each slice in the series smoothly blends to the next in the series shows that not only have we chosen an appropriate slice thickness ( $\sim 20$  nm) but that there is no significant damage (see videos S1 and S2 in the supplementary information) or misalignment between slices. The facetted nature of the cavities is further evidence that the original morphology of these features is preserved. It is clear that limited damage by the way of curtaining is present (see Fig. 4a) and is particularly hard to avoid when sectioning materials that offer vastly different resistance to the ion beam (i.e. between the cavities and the steel matrix); however, the continuity of the features through the slice series has not significantly affected the shapes of the cavities.

#### Scanning transmission electron microscopy

As mentioned above, a thin ( $<5 \mu$ m) lamella of material from each RoI was extracted for analysis in the STEM after FIB-SEM serial sectioning. Once extracted and mounted onto a copper TEM grid, further thinning, using the FIB, was required to produce an electron transparent lamella for STEM imaging. Scanning transmission electron microscopy imaging and EDX spectrum imaging were performed using a probe side aberration corrected FEI Titan G2 80-200 S/TEM operated at 200 kV (Fig. 5). Images (STEM) were collected using a convergence angle of 18 mrad and a high angle annular dark field (HAADF) detector with an inner angle of 54 mrad. EDX spectroscopy and compositional analysis were performed using the Super-X detector configuration (4 × 30 mm silicon drift detectors) with a solid angle of  $\sim 0.7$  srad and a beam current of  $\sim 0.6$  nA. Spectrum images were acquired with a dwell time of 30 µs per pixel. For display purposes, all elemental maps were processed using a three-pixel smoothing window in the Bruker Esprit software.

Energy dispersive X-ray spectrum images were also recorded in the FEI Helios 650 FIB-SEM as the sample extracted using the FIB was thicker than optimal for STEM analysis. All SEM-EDX spectroscopy measurements were made at 20 kV and high current using an Oxford Instruments Xmax80 EDX detector and Aztec software (*see* Fig. 6).

In all cases, the quantification of the EDX spectroscopy data extended only to the identification of the major, minor and trace elements.

### Results

#### Overview

Figure 1 shows the orientation of the matchstick with respect to the original welded component. As can be seen in Fig. 2, there are two types of feature visible (what we have termed stringers and dark lines), which have then been segmented in the subvolume (Fig. 2a). Both of these features appear dark with respect to the steel matrix and lighter with respect to the air background in the virtual slices of the reconstructed dataset, indicating an electron density (i.e. function of density and/or atomic number) somewhere intermediate to the steel and the air (Fig. 2b). These stringers have a range of sizes, from  $\sim 20 \ \mu m$  down to the resolution limit of the X-ray CT imaging ( $\sim 1 \mu m$ ). It is evident from the X-ray CT data that these features are generally aligned along the length of the matchstick (i.e. along the axis of the header). We have termed these cigar shaped (prolate spheroids) features 'stringers'. Their appearance, as shown in Fig. 2a, is typical of what is observed in all locations so far surveyed across the header. Figure 2 shows the success of the 'correlative tomography' approach as we have been able to locate and section one of the stringers ( $RoI_1$  in Fig. 2) identified from the X-ray CT data ready for higher resolution analysis using SEM and TEM. It should be noted that the total volume of the stringers in the volume measured is <0.1%, highlighting the need for guidance in finding these features, which look densely populated in 3D are in fact quite disperse. Figure 3 confirms the stringer as a solid multiphase feature and provides a comparison of the feature as observed by X-ray and SEM. As can be seen,

3



4 *a* unprocessed SEM image (using mixed signal of secondary and backscatter electrons via through lens detector) of representative region of cross-section created by FIB, one of series of images acquired during serial sectioning using Auto Slice and View. *b* 3D visualisation of cavities created after segmentation of FIB-SEM serial section volume. *c* 3D visualisation of cavities and M<sub>23</sub>C<sub>6</sub> phase prepared in same way as *b*. *d* 3D visualisation of cavities, M<sub>23</sub>C<sub>6</sub> and ferrite prepared in same way as *c* viewed from different perspective to show clearly the morphology of ferrite phase. Cavities are represented in blue, M<sub>23</sub>C<sub>6</sub> phase in red and the ferrite in green as indicated by key in *b* 

the shape orientation compares very favourably between the two. Multiple measurements from each image give an average length and height of the stringer of 17-5 and 4-1  $\mu$ m (for the maximum height) as measured by X-ray CT and 18.662 and 3.650  $\mu$ m as measured from the SEM image. This gives an agreement within 7% of the length and within 12% of the height which looks reasonable. While we have chosen comparative slices, we have not serially sectioned the entire feature so a degree of uncertainty remains in the position of the slice we have compared, so caution needs to be applied, and further measurements made, before, for example, altering the segmentation routine of the X-ray CT data.

Figure 2 shows again the success of the correlative tomography approach, this time in locating and sectioning in the FIB-SEM an example of what we have termed 'dark lines' from the virtual slices of the reconstructed X-ray CT data (RoI<sub>2</sub>). These fine features are very difficult to segment in the X-ray CT data; they are shown in the visualisation of the subvolume in Fig. 2*a* (their width has been slightly exaggerated for easy viewing). As can be seen in the segmented subvolume of the X-ray CT data in Fig. 2, these features are always aligned at ~45° to the length of the matchstick. Looking at the XY plane presented in

Fig. 1, this would result in a 45° plane coming out of the surface in such a way that a slice through it, following the XY plane, would reveal a line parallel to the Xdirection, i.e. as would be observed from a physical cross-section or metallographically prepared sample. This orientation was consistently observed for a total of 8-10 'dark lines' found in this sample, although there is some uncertainty as to whether some of the dark lines are separate or connected. The physical cross-sections prepared using the FIB allowed identification of the 'dark lines' as boundaries decorated with cavities and multiple additional phases (see  $RoI_2$  in Fig. 2d). The orientation of this boundary can be seen in the X-ray CT data in Fig. 2 from the YZ perspective. The contrast observed in these virtual slices is likely to be due to the presence of the cavities on the grain boundary with some influence from the other phases; this non-optimum contrast made it impossible to distinguish individual cavities from the X-ray CT alone. SAED in the TEM and STEM EDX spectrum imaging of a lamella prepared from this same region confirmed the 'dark line' as a grain boundary hosting a total of three phases identified in addition to the cavities (Fig. 5). The phases were M<sub>23</sub>C<sub>6</sub>, ferrite and G phase (Fig. 5k). High resolution 3D volumes obtained by serial sectioning



5 *a* Image (HAADF STEM) from grain boundary showing faceted cavity in contact with different second phase particles and matrix. *b-j* elemental maps recorded using STEM EDX spectrum imaging of area shown in *a. k* schematic diagram overlaid onto HAADF image showing key to different phases identified by composition and selected area electron diffraction (SAED), *V* is cavity,  $\alpha$  is ferrite,  $\gamma$  is austenite matrix and 's', 'm' and 'l' indicated small, medium and large M<sub>23</sub>C<sub>6</sub> phase particles respectively. *I* SAED pattern from area labelled  $\alpha$  in *k*; yellow grid is [5–23]<sub>bcc</sub> pattern and red grid is [2–11]<sub>bcc</sub> pattern

with the FIB containing the 'dark line' feature show the 3D arrangement of cavities and carbides around the grain boundary (*see* Fig. 4). This demonstrates very clearly the intimate association of the cavities with the second phases in the boundary.

# $M_{23}C_6$

Our study has highlighted three distinct morphologies of the  $M_{23}C_6$  carbides, namely, large irregular grain boundary particles that grow into the grains on both sides of the boundary, intermediate particles that also sit on the grain boundary but appear only grow into one grain, and small intragranular carbides (see Figs. 4c and d and 5k). Despite the three distinct morphologies, all appear to have the same composition (see Fig. 5). They are composed predominantly of Cr and C, as expected, but also contain Mo and Mn. As demonstrated in Fig. 5d, very little, if any, Fe is present in this phase, but its presence cannot be discounted completely. The large (500–2000 nm)  $M_{23}C_6$  irregular carbides on the grain boundary have their maximum dimension typically along the grain boundary essentially forming discontinuous plates (see regions marked 'l' in Fig. 5k). The intermediate (100-300 nm) M<sub>23</sub>C<sub>6</sub> carbide (see regions marked 'm' in Fig. 5k) particles also sit on the grain boundaries. Finally, there is a huge population of small intragranular carbides (45–90 nm) (see regions marked 's' in Fig. 5k, and in red in Fig. 4c and d). Close inspection in the TEM reveals that these intragranular carbides are cuboidal in nature, although their small size means that they often appear spheroidal when viewed by SEM.

 $M_{23}C_6$  is generally the first precipitate to form on aging 316 steels.<sup>24–27</sup> At 525°C,  $M_{23}C_6$  starts to form at grain boundaries after only a few tens of hours,<sup>25</sup> much earlier than is indicated on time-temperature-transformation curves generated by X-ray methods, which can detect precipitates only when present at overall levels >0.3 wt-%.<sup>26</sup> The transformation rate to M<sub>23</sub>C<sub>6</sub> increases significantly with the concentrations of both C and Mo. Weiss and Stickler suggested that initially the  $M_{23}C_6$  is a predominantly Cr/Fe carbide but that, over time, the molybdenum content increases to a composition of (Fe, Cr)<sub>21</sub>Mo<sub>3</sub>C<sub>6</sub>. They also identified the presence of Mo in the M<sub>23</sub>C<sub>6</sub>. Mn is also seen to be present in the M<sub>23</sub>C<sub>6</sub> here, consistent with the finding of Lai<sup>24</sup> that the solubility of Mn in  $M_{23}C_6$  is 8%. Nucleation of  $M_{23}C_6$  will occur first at high energy sites on grain boundaries or phase interfaces. Within the HAZ, nucleation will be most favoured on any ferriteaustenite boundaries because the  $\delta$  ferrite is Cr and Mo rich and the austenite is C rich.  $^{28}$  Austenite–austenite

grain boundaries are the next preferred nucleation site because grain edges and corners offer energetically favourable nucleation sites and diffusion is quickest along the grain boundaries. The precipitation of intragranular carbides also occurs, albeit more slowly, on twins, low angle boundaries and dislocations. Hence, the focus of these carbides on and around the grain boundary is seen in Fig. 4*c* and *d* and also shown by Jazaeri *et al.*<sup>11</sup> and Minami *et al.*<sup>29</sup> Previous SANS results<sup>30</sup> have detected a distribution of particles ~30 nm in size. Only a few particles close to this size, likely to be carbides, have been detected in the current study, but imaging was not optimised to observe such particles.

#### G phase

Scanning transmission electron microscopy, EDX spectrum imaging revealed small regions of what we have identified as G phase (see Fig. 5), which is a Ni, Nb and Si rich phase with small amounts of Fe and  $Cr^{2/2}$ previously reported with a composition of 51Ni-29Nb-14Si-4·5Fe-1·5Cr (wt-%).<sup>31</sup> The G phase also contains a surprisingly high level of Mo, which has not previously been reported as a constituent of this phase. The small size (smallest dimension typically <100 nm) of these particles has made examination using diffraction difficult. This appears to be the first time Nb rich G phase has been reported in an 18Cr-12Ni steel,<sup>27</sup> but both the diminutive size and the close proximity to multiple other phases would have made it quite easy to overlook in previous work. These particles were identified exclusively along the grain boundary and were intimately related to the  $M_{23}C_6$  carbides.

Although the Nb content of this header is unknown, others from the same cast have had their Nb content measured as 0.06 wt-%. This would explain the presence of G phase, which has previously been observed to form in preference to  $\sigma$  phase at grain boundaries in Nb stabilised austenitic stainless steels. Aging (8000 h) of such material at 500°C produced G phase and M<sub>23</sub>C<sub>6</sub> associated with each other at the grain boundaries,<sup>31</sup> as found in the present study.

#### Ferrite

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We have identified the presence of ferrite via STEM EDX and SAED in the TEM (see Fig. 5). Figure 4d shows that there are significant amounts of this ferrite growing away from the grain boundaries and the carbides, while none was found within the grains (rendered transparent here). The morphology of the ferrite is very striking; it is strongly facetted and grows away from the grain boundary (or more specifically it appears to initiate from the  $M_{23}C_6$ ) into the matrix with a roughly triangular cross-section. The maximum observed length of the ferrite was  $\sim 2 \,\mu\text{m}$ . This 3D visualisation also reveals the abundance of ferrite in the region of the boundary. The ferrite's composition is Fe rich, and it appears to contain significantly lower concentrations of all the same constituents as the matrix. However, the detection of these elements will be influenced by the surrounding matrix, and it is therefore very difficult to completely exclude their presence in the ferrite phase. Of particular note is the segregation of nickel into the austenite matrix immediately around the ferrite (see the edge of the ferrite phase in Fig. 5e).

#### Cavities

All the cavities have very irregular shapes, e.g. see Fig. 7. This can be missed in 2D slices but is quickly revealed in the 3D analysis. As well as the irregular shape of the cavities seen in Figs. 2, 4 and 5, close inspection reveals that almost all of the interfaces are strongly facetted, as seen in Fig. 5k (marked as V) and Fig. 7. The cavities have boundaries with all of the other grain boundary phases identified, and all of them appear more or less facetted. Facetted creep cavities have previously been observed in pure iron,<sup>32</sup> a Ni base superalloy,<sup>33</sup> copper<sup>34</sup> and  $\alpha/\beta$  brass<sup>35</sup> but have not previously been reported in 316 steel. The cavities appear from the FIB-SEM and TEM images to be fairly uniformly spaced along the grain boundary, and for the  $\sim 40$  cavities measured, we found a maximum single dimension of 2500 nm. Further work segmenting and quantifying the 3D data is required to properly describe the size distribution, although many cavities  $>1 \ \mu m$  can be clearly seen in Fig. 4.

This facetted cavity shape contrasts with the general assumption made when modelling the development of creep cavities at grain boundaries that isotropic surface tension interacting with a planar boundary will lead to their being lenticular. This assumption requires not only that void growth is diffusion controlled but that 'diffusion is sufficiently fast on the surface of the cavities to maintain their equilibrium lens shape'.<sup>36</sup>

#### Stringers

The most abundant features identified by X-ray CT are what have been termed 'stringers'. As can be seen in Fig. 2*a*, the stringers are aligned roughly along the length of the matchstick, which is parallel to the forming direction in the original component (see Figs. 1 and 2a). As seen in Fig. 2c and d for RoI<sub>1</sub>, these are cigar shaped particles that have an internal structure not apparent from the X-ray CT data. As described before, they have been detected with a range of sizes, from  $\sim 20 \ \mu m$  down to the resolution limit of the X-ray CT imaging ( $\sim 1 \mu m$ ). The detailed results of EDX microanalysis of the stringer are shown in Fig. 6. The chemistry of the stringers is complex, and three phases were identified: first, as expected, manganese sulphide; second, a Ca, Al, Si and O rich phase, which must be a silicate inclusion originating from the original casting; and, last, a Cr and Ti rich oxide phase, containing a small amount of Mn. The insulating nature of the silicate phase is apparent from its charging under the electron beam (see Fig. 2d, RoI<sub>1</sub>). None of the phases contained more than a trace of Fe. The composition and morphology probably originate from the initial casting and forging of the component. This particular stringer sits exactly upon a grain boundary, but one at which there appears to be no secondary phases or cavities. (Later in the preparation of the cross-section, a few cavities were detected at the interface between the stringer and the matrix.) All the stringers examined were in contact with a grain boundary and contained manganese sulphide and the silicate phase.

Generally, the presence of what are usually described just as 'manganese sulphide' inclusions is not seen as a cause for concern,<sup>37</sup> although their size suggests that they might act as stress raisers in highly stressed locations, acting as points of initiation for more extensive. The fact that these stringers are not always



6 a image\_(SEM) (using a mixed signal of secondary and backscatter electrons via 'through lens detector') of stringer imaged from extracted TEM lamella, b-I elemental maps of area shown in a recorded in SEM using EDX spectrum imaging, m backscatter electron image recorded in SEM showing more clearly internal phases as effect of charging is reduced and n schematic diagram showing key to phases identified by EDX spectroscopy

associated with cavitated boundaries suggests that, as previously thought, their presence is not critical to the creep behaviour. However, once a crack has started, it is possible that their large size may assist crack propagation, as they may provide an easy path to other nearby cavitated boundaries.

# Discussion

Techniques like X-ray CT and SANS are invaluable in providing a statistically significant overview of particles and cavities over large volumes. X-ray CT has the advantage of accurate spatial mapping in 3D, although, as a rule of thumb, the resolution (which is somewhat bigger than the pixel size) scales with the analysed volume with a multiplication factor of  $\sim 1/1000$ .<sup>14</sup> In other words 1 µm resolution is possible on a 1 mm diameter specimen and 10 µm resolution attainable on a 10 mm sample. However, in practice, the ability to distinguish different phases is also determined by the contrast between the phases. The high absorption of the X-rays by our steel sample results in noise that reduces the effective spatial resolution. The fact that we can see the 'dark lines' confirms that we have a resolution of the



7 Two SEM images recorded with through lens detector providing mix of backscatter and secondary electrons showing facetted nature of cavities

order of 1-2 µm. The absence of a sharp change in contrast between phases associated with the dark line, the fact that there are multiple different phases and the noise on top of everything result in an uncertainty in identifying the defects in this sample equivalent to about three to four voxels. SANS can, depending on the instrumental set-up, provide volumetric measurement of submicrometre features as reviewed by Rustichelli,<sup>38</sup> and is especially powerful in detecting the size and distribu-tion of nanometre sized defects.<sup>12,30</sup> However, neither of these techniques can necessarily distinguish between secondary phases and cavities, so these measurements must be supported by microscopic analysis of the microstructure.<sup>11</sup> Nor are traditional metallographic preparation techniques ideally suited to distinguishing between cavities and secondary phases because the second phases can be inadvertently removed during polishing or etching and the edges of cavities may be rounded. Crosssections created by FIB sectioning have previously been shown<sup>20</sup> to provide a superior preparation method to unambiguously identify cavities separately from secondary phases. In addition, the composition can be determined with submicrometre spatial resolution only by microanalysis of thin foils, and the combination of STEM EDX and SAED provides simultaneous phase identification. This demonstrates that a complete understanding of these complex microstructures requires study of the same area using a combination of techniques.

#### Microstructure

The intimate association of the cavities with the second phases in the boundary (clearly seen in Figs. 2, 4 and 5) shows that we need to understand the evolution of all the phases in this sample to understand how the local environment around the grain boundaries led to the nucleation of the cavities and affected their growth. At the service temperature of 525°C, volume diffusion of substitutional solutes in 316 steel will be very slow so the nucleation and growth of second phases involving them will be governed by solute supply through grain boundaries and interfaces and along dislocations.

A plastically strained zone would have been created during fabrication by the thermomechanical cycles associated with the original multipass weld. The dislocations in this strain affected zone, extending several millimetres from the weld fusion line, would have provided nucleation sites and rapid diffusion paths for solute supply to create the array of intragranular carbides seen in Fig. 4c and d. Orderly lines of precipitates have been observed (see Fig. 4a and c) close to the grain boundary in this sample similar to those previously found to be nucleated on dislocations.<sup>39,40</sup> Other intragranular carbides, which do not appear to be aligned, would still be expected to have formed on dislocations. Because the rate of solute supply from the bulk is so limited, the carbides connected by higher diffusivity paths coarsen during aging, leaving a reduced density of intragranular carbides and the carbides along the grain boundary coarsen to micrometre sized particles.29

Ferrite has previously been identified in type 316H stainless steel exposed to creep conditions in different forms, and its presence has been attributed to a range of different sources.<sup>41</sup> The starting material may contain  $\delta$  ferrite at levels of 1–2%, although a magnaprobe investigation of an unaged header (cast 53579) gave a ferrite content of <0.1%. In addition, welds in this

material are designed to contain an amount of  $\delta$  ferrite to alleviate hot cracking,<sup>42,43</sup> so regions close to the fusion boundary may have contained a few wt-% of  $\delta$  ferrite after welding.

Relatively short exposure (<1000 h) to temperatures of 600°C and above has been shown to be sufficient to completely transform the  $\delta$  ferrite,<sup>44</sup> usually forming  $\sigma$ phase and austenite, and sometimes also Laves and  $\chi$ phase. Tests at 550°C for up to 6200 h<sup>45,46</sup> have shown stress to accelerate the transformation of  $\delta$  ferrite in 316 steel, but extrapolating these results suggests that after 65 000 h, <50% of the  $\delta$  ferrite would have transformed. In another study,<sup>47</sup> 50% of the delta ferrite had converted to austenite, sigma ( $\sigma$ ), chi ( $\chi$ ) and M<sub>23</sub>C<sub>6</sub> after only 5000 h at 500°C. For the sample studied here, which has been exposed to 525°C for 65 000 h, it is therefore possible that some  $\delta$  ferrite might still be present. If this were the case, it would be expected to have become enriched in Cr and Mo as it transformed to austenite.<sup>28</sup> It is this availability of Cr and Mo that would promote the precipitation of  $\sigma$ ,  $\chi$  and Laves phase. However, all the ferrite analysed in the present study was predominantly composed of Fe and was depleted in Cr, Ni and Mo compared to the austenitic matrix, and we have not detected any  $\sigma$ ,  $\chi$  or Laves phase, so this seems to rule out the possibility that the ferrite seen around the grain boundary carbides remains from  $\delta$  ferrite formed during the manufacture of the header.

Thus, we look to other possible routes for the formation of this ferrite phase. Stewart and Spellward<sup>48</sup> observed similar triangular ferrite particles up to 1  $\mu$ m long and of measured composition 91Fe–4Cr–4Ni (wt-%) close to grain boundaries in a similar 316 steel with a similar service history and suggested that this might be the result of 'supersensitisation' where, once all the Cr had been incorporated into the carbides, the austenite transforms back to ferrite. No explanation was suggested as to why the removal of Cr (a ferrite former) should result in the FCC to BCC transformation.

The shape of the ferrite seen in the present study and in Ref. 48 is reminiscent of the saw tooth morphology of Widmanstätten ferrite. This is more usually associated with high temperature fcc to bcc transformations in ferritic steels, but Singhal and Martin<sup>49</sup> found that Widmanstätten ferrite nucleated on M<sub>23</sub>C<sub>6</sub> precipitates in a 24Ni-25Cr austenitic steel aged for 25 h at 750°C. All the ferrite observed in the present study was in contact with  $M_{23}C_6$  at a grain boundary and was depleted in Cr, Mo and Ni, but apparently, only Ni was rejected into the surrounding matrix during growth. This suggests that at the comparatively low temperature of 525°C, the diffusion of C was slow enough that, in forming  $M_{23}C_6$ , the strong carbide formers Cr and Mo might have reduced the local activity of C enough to stabilise ferrite, which, in forming, rejected the austenite forming Ni into the surrounding matrix. Thus, it is unlikely that the ferrite observed around the grain boundaries in this sample originates from the  $\delta$  ferrite present immediately after the manufacture of the header; the morphology, the composition and the chemistry of its surroundings all point to its having developed during service.

The detrimental influence of extended aging of 316 steel at high temperatures on its creep resistance has, in part, been attributed<sup>50</sup> to the formation of brittle intermetallic  $\sigma$  phase, which has often been found to replace M<sub>23</sub>C<sub>6</sub> during extended service and appears to be favoured in steels stabilised by Ti or Nb additions.<sup>51,52</sup> In particular, Morris *et al.* suggest that it is not just that  $\sigma$  phase is brittle but that small M<sub>23</sub>C<sub>6</sub> carbides are replaced by larger  $\sigma$  phase particles, allowing easier grain boundary sliding and higher local stress concentrations. No  $\sigma$  phase was found in the present study. This is in common with a TEM study of 316 steel with a similar composition and service history to that of the present study,<sup>26,48</sup> which also failed to detect  $\sigma$  phase by X-ray methods after 94 000 h at 600°C, so it appears that at temperatures as low as 525°C, other phases such as  $\sigma$  and  $\chi$  may not form even after tens of 1000 s of hours.<sup>28</sup>

#### Development of cavitation

Here, we consider the nucleation and growth of the cavities and how this occurred in association with the continuing evolution of the phases around the grain boundary.

The residual stress associated with the weld is the main driving force for cavity nucleation in the material surrounding it,<sup>4</sup> and it is very noticeable that the heavily cavitated boundaries in our sample, seen as 'dark lines' in virtual slices of the X-ray CT data, all lie at close to 45° to the plane of the crack (as described above). This suggests that the shear components of the residual stress field dominated the cavitation behaviour in the HAZ. Voids formed during creep of hydrogen reformer tubes have been observed to form mostly at grain boundaries that were neither perpendicular nor parallel to the hoop stress with those close to the outer wall, in particular, being most likely to lie at  $45^{\circ}$  to the hoop stress.<sup>53</sup> M<sub>23</sub>C<sub>6</sub> is known to form at grain boundaries of similar steels early in the service life of the header, so it is most likely that the  $M_{23}C_6$ carbides were present from an early stage and contributed to the formation of the cavities<sup>47</sup> (as well as to the ferrite as described above). The grain boundaries at 45° to the crack would experience the highest driving force for grain boundary sliding, which would be impeded by the growing M23C6; these and other second phase particles at the boundary would then act as stress raisers. Vacancies would be attracted to relieve the stress built up by the resulting dislocation pile-ups in the boundary around the particles, reducing the nucleation barrier to cavity formation at the carbide/boundary interface. Incoherent interfaces between  $M_{23}C_6$  and austenite have been suggested as preferred sites for creep voids.<sup>54,55</sup> The orientation of the cavitated boundaries is crucial when considering the mechanisms driving the cavity formation. We have observed in our sample a consistent set of findings with respect to the orientation of several cavitated boundaries (see the segmented subvolume of the X-ray CT data in Fig. 2a) examined in the region of the weld, but we fully acknowledge the need to extend the volumes analysed as in light of develop greater confidence in the generality of these

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local microstructural inhomogeneities and the complex stress state in this component (*see* Fig. 8 in Ref. 12) to develop greater confidence in the generality of these observations. Superimposed on the effects of the orientation of the grain boundaries planes relative to the principle stresses will be the effects of the crystallographic orientations of the adjoining grains themselves, relative to the principle stresses, relative to the local grain boundary plane and relative to each other. Dzieciol *et al.*<sup>34</sup> found a large scatter in the growth rates of individual creep voids in copper and suggested that local dislocation glide made a significant contribution to void growth. They pointed out that since dislocation glide occurs on particular crystallographic planes, void growth might be expected to depend on the crystallographic orientation of the adjoining grains.

The facetting of creep cavities has been previously been attributed either to the anisotropy of their surface energy or to kinetic effects by which slip on favoured planes will produce factetted voids. Two-dimensional dislocation simulations for example<sup>56</sup> showed that in single crystals slip could produce facets on originally round voids.

It is not clear whether the ferrite or the cavities nucleated first, but the irregular shape of the cavities and their facetted morphology suggest two alternative mechanisms for their formation. The irregular cavities could have formed either by coalescence of smaller cavities after separate nucleation events or, perhaps, by growth of one initial cavity along the interfaces between the  $M_{23}C_6$  carbides and either the matrix or the ferrite as well as along the grain boundary. In either case, continued stress assisted vacancy absorption might be expected to modify normal diffusional growth. The close association between the grain boundary phases and the cavities and the complexity of their interfaces is clearly visible in Figs. 2, 4 and 5 and, while there is still a lot of work to do to understand their evolution, it is clear that the role of the grain boundary phases cannot be ignored.<sup>57,58</sup> However they are formed, it is clear that for a fixed total volume of cavities (i.e. number of vacancies) the convoluted cavities would cause decohesion over a far wider area of grain surface than the formation of spherical cavities would, leading to a correspondingly higher embrittling effect. Moreover, in the later stages of creep, complex shaped cavities may grow faster than spherical, rod-like or ellipsoidal cavities as observed for a leaded  $\alpha/\beta$  brass by Isaac *et al.*<sup>35</sup>

If the presence of the grain boundary carbides is critical to the formation of the cavities, it is possible that failure occurred where it did because the dislocations in the HAZ (formed by plastic deformation resulting from the thermal stresses during and after welding) provided additional rapid pathways for solute and/or vacancy supply to the grain boundary region and the intergranular phases, accelerating the formation of cavities in this region. It is also likely that the presence of the ferrite, seen in Fig. 4d, reduced the ductility in the region of the cavitated boundaries, assisting in final fracture.

It is only by exploring and understanding the arrangement of the secondary phases along the grain boundary and their 3D relationship to the dislocation structure that we can hope to then create an accurate model describing the nucleation and growth of the cavities. Additional regions around the crack will be the subject of further investigation.

# Conclusions

Multiscale 3D analysis using the correlative tomography approach has allowed us to find key RoIs for high resolution SEM and TEM analysis. In particular X-ray CT analysis has been successfully applied to create a map of the defects buried inside a service aged stainless steel sample and to identify creep damaged zones. Preparation of cross-sections and lamellae using the FIB allowed site specific phase identification and chemical analysis in the regions identified as having the highest propensity for creep cavitation. Scanning electron microscopy, TEM and STEM revealed that the following.

1. The most heavily cavitated boundaries observed in this sample lie parallel to the crack in 2D and at an angle  $\sim 45^{\circ}$  to the plane of the crack in 3D.

2. The surfaces of all cavities appeared facetted.

3. The cavities at the grain boundary were of irregular shape and always connected to secondary phases including  $M_{23}C_6$  carbides, ferrite and G phase.

4. The ferrite appeared to have formed during service operation by transformation of the solute depleted austenite in contact with the growing  $M_{23}C_6$ .

5. Three morphologies of  $M_{23}C_6$  carbides were found: large irregular particles up to ~2 µm long sitting on the grain boundary and growing into the grains either side forming a discontinuous plate, intermediate particles (100–300 nm) sitting on the grain boundaries but only growing into one grain, and a large population of smaller (45–90 nm) cuboidal intragranular carbides.

6. Because of the intimacy of the contact between the cavities and grain boundary phases, these secondary phases are likely to control where, when and how the cavities themselves form. Cavity growth rates will depend on the local stress and diffusion conditions, which may be very different from those modelled assuming a homogeneous matrix on either side of the grain boundary.

7. 'Stringers' were found and identified as predominantly composed of Si, Ca, Al and O associated with manganese sulphide and are silicate inclusions from the original cast.

The complexity and facetted morphology of cavitated grain boundaries decorated with carbides of varying size, ferrite and G phase suggest that current models that are used to predict the lifetime of components, based on the nucleation and growth of spherical cavities, are oversimplified.

Further work will focus on a 3D nanoscale analysis of the decorated grain boundaries and on mapping the cavity sizes and morphologies and the microstructural variations in further volumes and chosen from both regions of low to high creep damage, with the aim of understanding the evolution of creep damage from initiation to rupture.

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