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Nb-doped TiO₂ Nanofibers for Lithium Ion Batteries

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Abstract

Niobium doped nanofibers elaborated by facile, single-step electrospinning present higher rate capability in electrochemical cycling experiments than non-doped materials. This is attributed to the reduction of Li⁺ diffusion path lengths and enhanced intimate inter-particle contact, in combination with improved intra-particle conductivity. Niobium doping has a significant effect on the electronic structure and provokes a substantial decrease in particle size.

Introduction

Since their introduction into the market in the early 90's by *Sony*, lithium ion batteries (LIB) have come a long way.¹ Originally intended to serve only for small portable electronic devices, they have conquered numerous other fields of domestic and military applications and have become an

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indispensable component of present day technologised society. Today LIB are at the dawn of entering not only the transportation market on the large scale but also to backbone the transition from fossil to renewable energy sources. In this development graphite, which is today's commercially established anode material, has limited high cycling rate properties due to necessary solid electrolyte interphase (SEI) formation and the risk of lithium plating when overcharged linked to its low working potential. An elevated performance is however crucial for application in the transportation sector, not only for rapid acceleration but also for reduced recharging times and recovery of braking energy. In the search of alternative anode materials TiO_2 has been widely recognised as a viable candidate due to its cost effectiveness, environmental benignancy and lithium storage capability. Of particular interest is its tetragonal polymorph, the anatase phase. It has a high theoretical capacity (335 mAh/g), a flat operating potential and low volume expansion (ca. 3 %) during Li^+ charge/discharge leading to long cycle life and durability.²

However, the poor electronic conductivity (ca. 10^{-12} to 10^{-7} S/cm), the frequent aggregation of titania nanoparticles, and the low Li^+ diffusion rate (ca. 10^{-15} to $10^{-9} \text{ cm}^2/\text{s}$) reduce the electrochemical performance of anatase- TiO_2 . Considerable efforts have been made to overcome the ionic and electronic transport limitation of titania anodes.³⁻⁶ One approach is the tailoring of morphology and particle size of TiO_2 -based materials. In this respect, one dimensional (1D) nanostructured materials such as nanofibers and nanotubes are particularly interesting for LIBs, due to their large surface/volume ratio, their vectorial electron (along the long dimension) and Li^+ (along the lateral direction) transport properties, and their ability to accommodate lithiation-induced stresses.^{7,8}

Soft-chemistry and template-based syntheses are usually employed to synthesise fibrous nanostructured electrodes, e.g.⁹⁻¹² Among the methods for generating 1D nanostructures, electrospinning is a simple and versatile technique for preparing ultrathin nanofibers and nanotubes of polymers,¹³ composites¹⁴ and ceramics¹⁵ with controlled and reproducible diameters. Recent years have witnessed a growing development of electrospun TiO_2 -based electrodes for lithium-ion batteries, fuel cells and other conversion and energy storage devices.¹⁶⁻²⁰

Various methods are pursued to further enhance the performance of electrospun titania when

used as an active insertion material in LIB. One of them is the preparation of composite nanostructured electrodes interconnecting titania with a conducting additive nanophase (based on carbon), which improves the electron transfer. With this aim, porous carbon nanofibers loaded with TiO₂ nanoparticles,²¹ composite TiO₂/C²² and TiO₂/graphene nanofibers²³ with improved reversible capacity and high rate behaviour have been recently reported. Another approach is surface modification to promote faster Li⁺ diffusion and electron transport, but also to suppress particle agglomeration.²⁴ For instance, the layer-by-layer self-assembly approach was combined with electrospinning for the fabrication of MoO₂-modified TiO₂ nanohybrids with improved lithium-storage properties, compared to the uncoated materials.¹⁹ A further approach to enhance the electrical conductivity of TiO₂ nanofibers is by loading with metal nanoparticles. LIB anodes based on electrospun titania nanofibers embedding 10% Au or Ag nanoparticles showed a 20% enhancement of the specific capacity and a 2-fold rate performance compared to the bare TiO₂.²⁵ TiO₂/Ag composite nanotubes fabricated by coaxial electrospinning also showed increased rate performance and cycling stability.²⁶

Doping with aliovalent ions is also a facile strategy to modify the electronic properties of electrospun titanium oxide.^{17,18,27–30} Nevertheless, only few papers report this procedure for application of the materials in LIB.³¹ This is astonishing as, for example, *Wang et al.* show that the conductivity, which is the crucial parameter for high performance cycling, can be raised two orders of magnitude by doping mesoporous TiO₂ with Nb.³² These findings are in agreement with the results of *Archana et al.* stating an enhanced Li⁺ mobility and diffusion coefficient for Nb-doped TiO₂ electrospun fibers³³ as well as findings by *Sheppard et al.* describing metallic-type conduction behaviour of Nb-doped TiO₂.³⁴

In the present paper we report the synergistic effects of nanostructuring by electrospinning and doping with niobium on TiO₂ anode materials for lithium-storage applications. Nanofibrous TiO₂ with and without Nb doping are synthesised via electrospinning and characterised via XRD, SEM and XAS. Furthermore we provide results from DFT calculations to explain in detail the effect of Nb doping on electronic band structure, explaining the observed increase of the electronic

conductivity. The potency of the nanofibers as lithium insertion material was tested by preparing film electrodes and applying them in half-cells vs. lithium. Simple, standardized synthesis and electrode preparation method were chosen to elucidate only the effect of Nb-doping on the electrochemical cycling properties, avoiding possible pseudocapacitive storage contribution from morphology tailoring or enhanced conductivity by elaborated electrode formulation, as such might superimpose to the effects of the dopant and hence disguise its benefits.

Experimental

Synthesis of 10% Nb-doped and non-doped TiO₂ nanofibers

The detailed synthesis of titania nanofibers has already been reported elsewhere.¹⁷ Briefly, the carrier polymer solution based on polyvinyl pyrrolidone (average Mw \approx 1,300,000 g/mol, Aldrich) in absolute ethanol (puriss., Sigma-Aldrich) was mixed with a solution containing 0.52 ml of titanium(IV)isopropoxide (97%, Aldrich), (and 0.040 ml of niobium ethoxide, 99.95%, Aldrich) and 1 ml of acetic acid (Sigma-Aldrich), corresponding to a dopant content of 10 at%. Electrospinning of the final solution was carried out in air at room temperature with a standard syringe and a grounded collector plate configuration. The distance between the needle tip and the collector plate was 10 cm, the applied voltage 15 kV and the flow rate 0.5 ml/h. The as-prepared fibers were calcined in air at 500 °C at a heating rate of 5 °C/min for 6 hours in order to remove the carrier polymer.

Electrode and cell preparation

For electrochemical measurements, the titania powder which serves as active insertion material was mixed with carbon black and PVdF with a weight composition of 80:8:12, respectively. An NMP-based slurry of this mixture was tape casted on a copper foil using a doctor blade system leading to homogeneous \approx 25 μ m thick films which were cut and used as working electrodes.

These were built in *Swagelok* cells together with lithium metal as reference and counter electrode as well as *Whatman* fiber glass separators. A standard mixture of organic solvents containing ethylene carbonate, propylene carbonate and dimethyl carbonate in composition 1:1:3 and 1 M LiPF₆ was used as the electrolyte. Since both electrolyte and lithium metal are very sensitive to air and moisture, the assembly of the electrochemical cells was carried out in a glove box under Ar atmosphere.

Characterization

The nanofiber morphology was verified by using a scanning electron microscope (SEM) *Hitachi S-4800 FEG* equipped with energy-dispersive spectroscopy analysis (EDS) and a *JEOL 1200 EXII* transmission electron microscope (TEM) operating at 120 kV equipped with a CCD camera *SIS Olympus Quemesa* (11 million pixels). The crystal structures were identified by powder X-ray diffraction (XRD) using a *Philips X'pert* diffractometer with Cu K_α radiation at room temperature, in the range from 10-70° 2Θ using a step size of 0,026° and dwell time of 3500 sec. *Fullprof* was used for determination of unit cell parameters (Rietveld refinement). Nitrogen physisorption was carried out on a *Micromeritics ASAP2020* and the adsorption and desorption isotherms were analyzed using *Brunauer-Emmett-Teller* (BET) as well *Density Functional Theory* (DFT) methods to calculate specific surface area and pore size distribution, respectively.

The surface composition of the non-doped and Nb-doped TiO₂ fibers was monitored by X-ray photoelectron spectroscopy (XPS) on an *ESCALAB 250 (Thermo Electron)*. The X-ray excitation was provided by a monochromatic Al K_α (1486.6 eV) source with a beam diameter of 400 μm. A constant analyzer energy mode was used for the electron detection (20 eV pass energy). The detection of the emitted photoelectrons was performed perpendicularly to the sample surface. Data quantification was performed on the *Avantage* program. The background signal was removed using the Shirley method. The surface atomic concentrations were determined from photoelectron peaks areas using the atomic sensitivity factors reported by Scofield. Binding energies (BE) of all core levels were referred to the C-C bond of C 1s at 284.8 eV.

Raman spectroscopy was carried using *LabRam ARAMIS IR²*, *Horiba Jobin Yvon* with blue diode laser $\lambda = 473\text{ nm}$ or helium neon laser $\lambda = 633\text{ nm}$.

XAS measurements were carried out at ambient temperature on the beamlines A (HASYLAB @ DESY, Hamburg, Germany) and XAFS (ELETTRA, Trieste, Italy). The Ti K-edge spectra were recorded in transmission mode on homogeneous pressed pellets of adequate thickness. Fourier transform were performed using k^2 weighting and the structural parameters were determined by curve-fitting procedures using *Artemis* data analysis software with embedded *FEFF* tool for calculating scattering paths on the basis of the anatase structure. To fit Nb-doped anatase in the R-space, FEFF paths based on Nb₂O₅ crystal structure in which the central Nb atom was exchanged for Ti were added, in order to include Ti-Nb bonds in the calculations.

Electronic structure calculations based on DFT and a *generalized gradient approximation* (GGA) using an exchange-correlation potential by Perdew, Burke and Ernzerhof³⁵ were performed for anatase TiO₂ and three supercells of 24 atoms (NbTi₇O₁₆), 48 atoms (NbTi₁₅O₃₂) and 108 atoms (NbTi₃₅O₇₂) which correspond to Nb doping of about 12 %, 6 % and 3 %, respectively. The augmented plane-wave method + local orbital (APW+lo) as implemented in the WIEN2k code³⁶ was used with the muffin-tin radii $R_{\text{mt}}(\text{Ti}) = 2.0\text{ a.u.}$, $R_{\text{mt}}(\text{O}) = 1.4\text{ a.u.}$ and $R_{\text{mt}}(\text{Nb}) = 2.0\text{ a.u.}$, the semicore states: Ti 3s, Ti 3p, O 2s, Nb 4s and Nb 4p, the plane-wave cutoff: $\min(R_{\text{mt}}) \cdot \max(K) = 7$ (where K is a reciprocal lattice vector) and the magnitude of the largest vector in the charge-density Fourier expansion: $G_{\text{max}} = 15\text{ Ry}^{-1/2}$. The experimental values of the lattice constants were used for the calculations but the internal atomic positions were moved in order to minimize the internal atomic forces. A similar approach was successfully used for the electronic structure calculations of Sn³⁷ and Zr³⁸ doped anatase TiO₂. Self-consistency was achieved with an energy tolerance of 10^{-4} Ry and a force tolerance of $10^{-3}\text{ Ry Bohr}^{-1}$. The XAS spectra were calculated from the dipole transition strengths and the partial density of states (PDOS) in the same way as previous calculations for Li₄Ti₅O₁₂.³⁹ The calculated Ti K-edge spectra were averaged over all the Ti atoms of the supercells. For comparison with experiments the calculated spectra were convoluted with a Lorentzian function (FWHM=2 eV).

The electrochemical cycling was carried out in galvanostatic mode on a *Biologic* multichannel potentiostat. Different cycling rates were applied ranging from C/20 to 5 C whereas 1 C, corresponds to the insertion of 1 Li per hour into TiO₂ which is equal to a current of 336 mA/g.

Results

Electrode properties and electrochemical characterization

Film electrodes were prepared via tape casting of slurry made of Nb-doped or non-doped TiO₂/C/PVdF mixture which resulted in homogeneous, uniform and smooth film electrodes with a loading of Nb-doped or non-doped TiO₂ of roughly 1.5 mg/cm². Energy dispersive X-ray spectroscopy (EDS) analysis of such electrodes reveals a homogeneous distribution of active insertion material TiO₂ and dopant Nb, see Fig. 1(b) and (c) respectively. Furthermore a Nb content of 8.1±0.2 at% is determined from the ratio of Ti/Nb. This is close to the aspired value of 10 at%, while at the same time no unexpected elements were detected. EDS has a penetration depth of ≈ 10 μm, assuring that the here stated value can be considered as the bulk concentration.

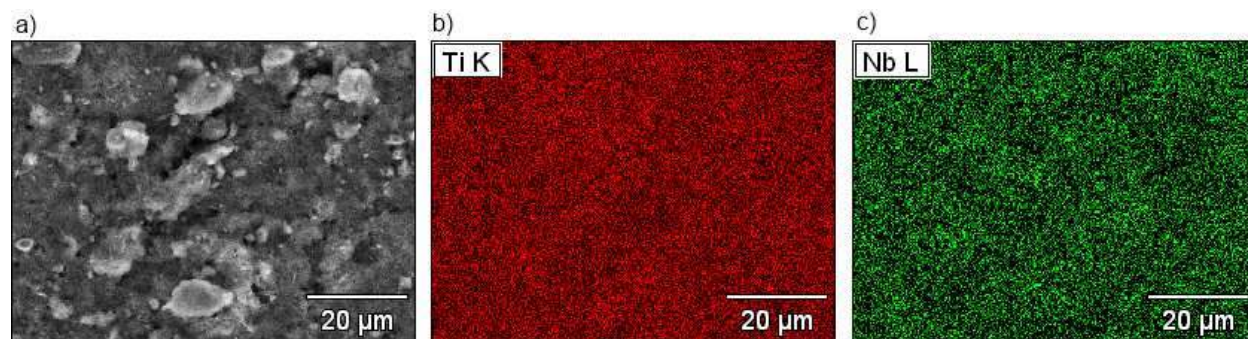


Figure 1: a) EDX image and b) elemental mapping of Ti and c) of Nb of tape casted film electrode.

The galvanostatic cycling curve for the first cycle with doped and non-doped TiO₂ electrodes are presented in Fig. 2(a). The curves are quite similar, revealing typical insertion behaviour of nanoscaled anatase with a plateau at 1.75 V which is followed by a slope which correspond to the biphasic and subsequent monophasic lithium insertion mechanism, respectively. It is noteworthy that the irreversible voltage plateau is less pronounced for the Nb-doped sample. In the first cycle a

high irreversibility is observed which suggests a SEI formation and trapping of Li^+ in irreversible sites. The samples were cycled 20 cycles at subsequently increased cycling rate and the mean values of each cycling rate are presented in 2(b). Similar average capacity values at low cycling rates $C/20$ and $C/5$ are obtained for doped and non-doped TiO_2 , see Fig. 2(b) 140 vs. 128 mAh/g and 97 vs. 92 mAh/g , respectively. However, as the cycling rate is increased, a growing difference between the capacities obtained for doped and non-doped TiO_2 electrodes appears, while no general difference between the galvanostatic cycling curves is observed (not shown here). Although the overall capacity values are well below those found by Wang *et al.* (142 and 108 mAh/g @ $C/4$ and 1 C , respectively for mesoporous Nb-doped TiO_2 with $128 \text{ m}^2/\text{g}$)³² the direct comparison of non-doped and Nb-doped material reveals a superior high rate capability of the Nb-doped sample. At highest cycling rate 5 C the doped samples achieve double the capacity of non-doped TiO_2 , 23 vs. 10 mAh/g . This phenomena of increased rate capability while maintaining absolute capacities was also recently reported for N-doped TiO_2 fibers.³¹

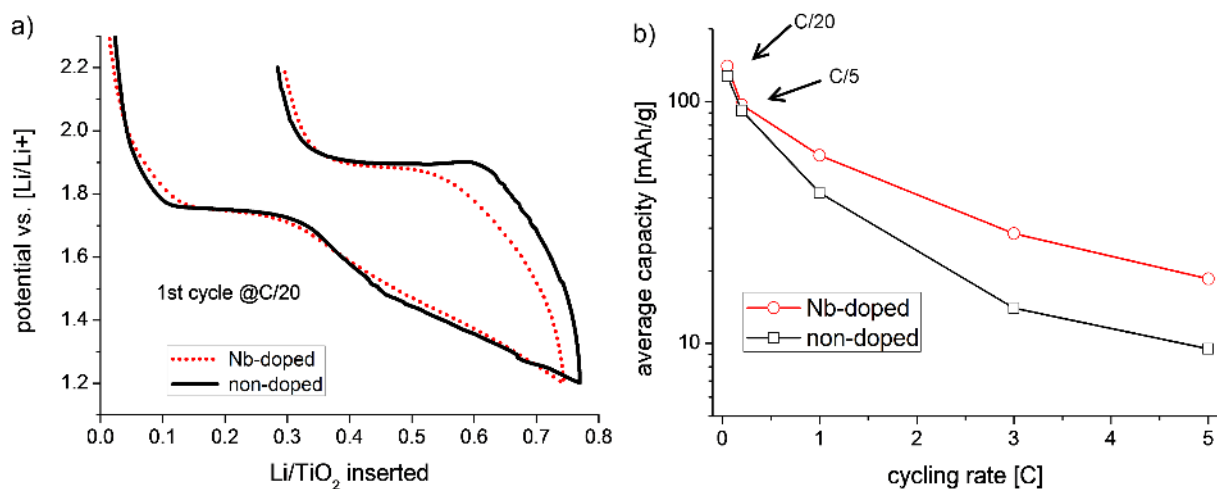


Figure 2: Comparison of a) galvanostatic cycling curve and b) mean rate capability of non-doped and Nb-doped electrospun TiO_2 nanofiber electrodes.

In order to elucidate the reason for this significant improvement in rate capability a series of characterization techniques and calculations were performed which illuminate the diverse influence of Nb doping on electrospun TiO_2 fibers.

Morphology

In Fig. 3(a) the SEM and (c) TEM micrographs show the morphology of the electrospun Nb-doped TiO₂ nanofibers. Uniform polycrystalline fibers with average diameter of $51 \pm 19 \text{ nm}$ are obtained whereas the length can reach several micrometers. Note that the TEM contrast variations reflect the polycrystalline microstructure of the fibers, consisting of well grown nanocrystals with different orientations. The graph in Fig. 3(b) shows the diameter distribution revealing the conformity of the nanofiber morphology, with $\geq 95 \%$ of all values falling into the narrow range 20-80 nm. SEM and TEM of non-doped TiO₂ show an equally narrow fiber distribution of diameter with a slightly higher average diameter of about 62 nm (not presented here). This small difference agrees with our previous observations¹⁷ and with the findings by Archana *et al.*³³ but is in opposition to the results of An and Ahn,⁴⁰ who showed rather a slight increase in the diameter of the fibers on Nb doping. It should be pointed out though that in any of these cases the magnitude of changes in diameter is small.

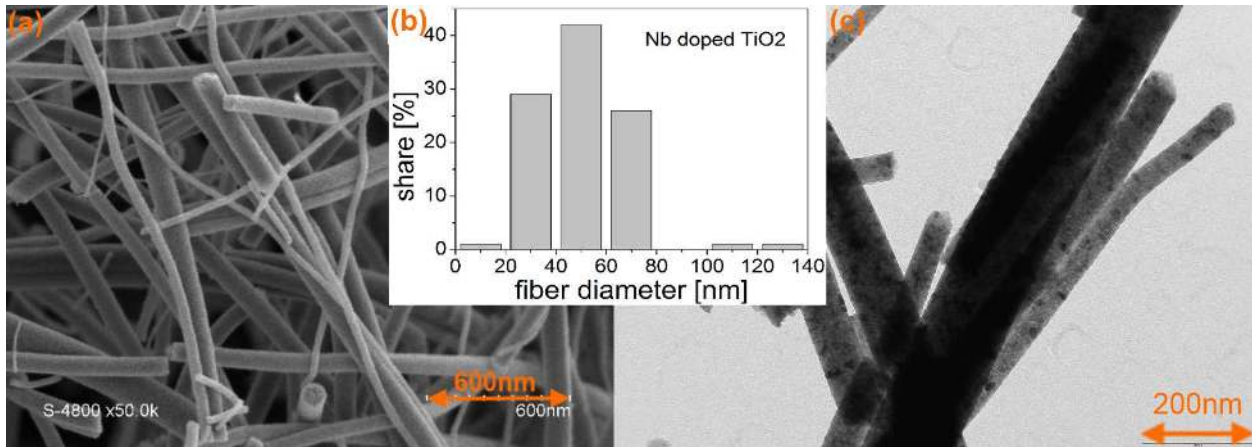


Figure 3: (a) SEM and (c) TEM micrographs of Nb-doped electrospun TiO₂ nanofibers and (b) distribution of fiber diameter.

The results of nitrogen physisorption reveal also a great conformity of doped and non-doped samples. A type IV physisorption isotherm⁴¹ is observed which indicates mesoporosity, and BET specific surface area of 66 and 44 m²/g for doped and non-doped sample, respectively was determined (not shown). We attribute this gain in specific surface area mainly to the reduced grain size,

as pointed out in Sec. . The similarity of the results is in agreement with the assumption that Nb doping has no significant influence on the global morphology of TiO₂ electrospun nanofibers, as shown in previous work.¹⁷

Crystal structure

When comparing crystal and local structure of doped and non-doped samples by XRD and Raman spectroscopy, (Fig. 4) we find a clear signature of anatase phase of both samples and no visible segregation of possible Nb-containing phases. While Raman spectra Fig. 4(b) reveals a single anatase phase, the XRD patterns Fig. 4(a) shows very minor traces of rutile as well as brookite in both samples, which is coherent with previous findings.¹⁷ The discrepancy between these two results is due to the lower sensitivity of Raman spectroscopy compared to XRD. In the XRD patterns we find besides a peak broadening, corresponding to a crystallite size reduction from 15 to 7 nm determined using the *Scherrer* equation, a slight but significant peak shift to lower diffraction angles for the Nb-doped samples. According to *Ruiz et al.* Nb doping obstructs grain coarsening, which *Guidi et al.* attribute to the reduced ionic oxygen mobility.^{42,43} The XRD peak shift corresponds to an increase of lattice parameter a (3.789 Å, ICSD009853) of $\approx 0,014$ Å equivalent to $\approx 3.7\%$. Whether this expansion is due to larger ionic radius of Nb⁵⁺ (0.64 Å) compared to that of Ti⁴⁺ (0.61 Å), or a strengthening in Coulombic repulsion force between Nb and Ti cannot be answered at this point with certainty. Our findings are in agreement with the general idea of lattice expansion through substitutional doping of Ti by Nb, except that the magnitude of expansion is somewhat smaller than those found in the literature.^{40,42,44,45} This is likely to be related to the influence of the nanostructuring on the lattice parameter which already accounts for an expansion of $\approx 2.4\%$ compared to reference value.⁴⁶

The Raman spectra show exclusively the typical bands of anatase, Fig. 4(b). A closer look reveals however small differences in the Nb-doped TiO₂ spectra. Not only a blue shift of the characteristic E_g band by ≈ 1 cm⁻¹ from the reference value of 144 cm⁻¹⁴⁷ is observed, but also a deformation of the line shape. This asymmetry might result from the sum of a distribution of

peaks reflecting a slight inhomogeneous Nb distribution. In addition the creation of defects through doping is contributing to the observed peak broadening. The E_g band corresponds to the O-Ti-O bending vibration. If applying the harmonic oscillator model, the exchange of Ti for heavier Nb atom should lead to a shift to lower wavenumber unless this change in mass is compensated by increased binding energy. Using the example of substitutional doping of Ti by Zr in TiO_2 *Lejon and Österlund* attributed the observed Raman blue shift to surface stresses as well as spatial phonon confinement effects linked to Zr-doping induced decrease of particle size, while at the same time an expansion of the lattice parameter is observed.⁴⁸ *Ruiz et al.* have described a direct relation between Raman line shift and expansion of lattice parameters, which substantiates the insertion of Nb dopant in TiO_2 lattice;⁴² analogous results are published by *Babu et al.* for substitution of oxygen by N dopant in TiO_2 .²⁸

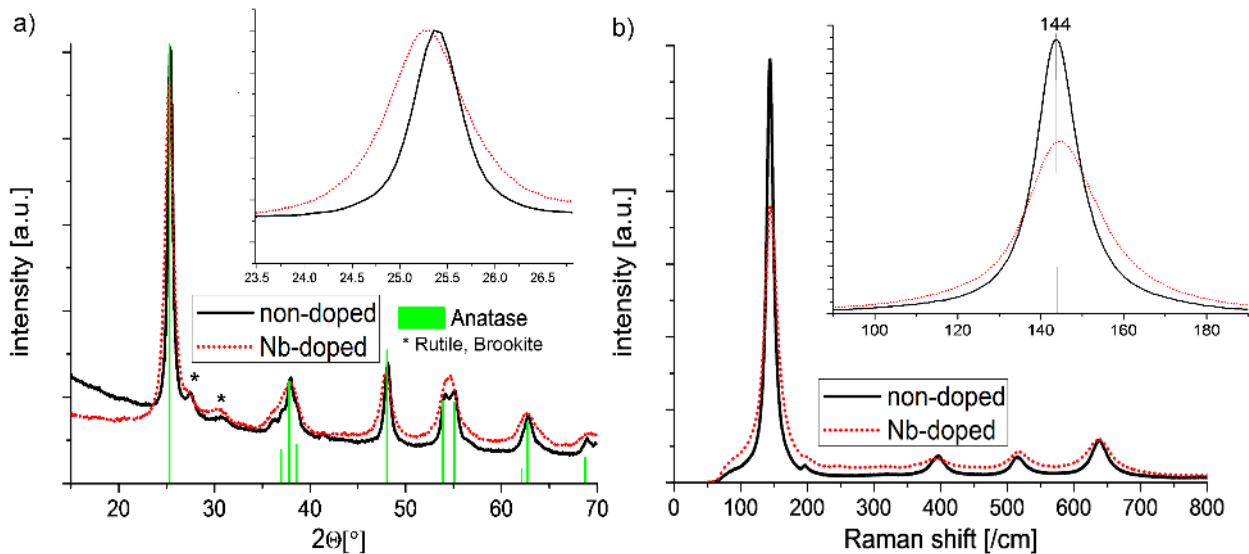


Figure 4: a) XRD diffractogram with anatase reference [ICSD 009853] and b) Raman spectra comparison of non-doped and Nb-doped TiO_2 nanofibers, showing the complete survey and an enlarged cut out of most intense signal.

Surface composition analysis

Since the structural characterization does not provide sufficient proof of Nb presence in the TiO_2 lattice XPS measurements were carried out. XPS spectra of the nanofibers are presented in Fig.

5(a) showing characteristic Nb 3d twin peaks for the Nb-doped sample and in Fig. 5(b) comparison of Ti 2p signals for non-doped and Nb-doped TiO₂ nanofibers. Ti 2p_{1/2} and Nb 3d_{3/2} peaks appear at 458.7 and 207.2 eV, respectively, indicating that titanium is present as Ti^{IV+} and niobium as Nb^{V+}. No sign of Ti^{III+} (455 eV⁴⁹) or Nb^{IV+} (205.9 eV⁵⁰) are visible in the spectra. These results foster the assumption that an effective aliovalent doping is accomplished. Furthermore we can deduct that charge compensation of Nb^{V+} is not achieved by reducing Ti^{IV+} to Ti^{III+} but more likely producing oxygen vacancies in the anatase lattice. It is nevertheless noteworthy that when interpreting XPS data of doped, high specific surface area, nanostructured samples one has to take due care as surface state effects, change in Fermi level position and band bending are influencing the peak position.^{5,42,51}

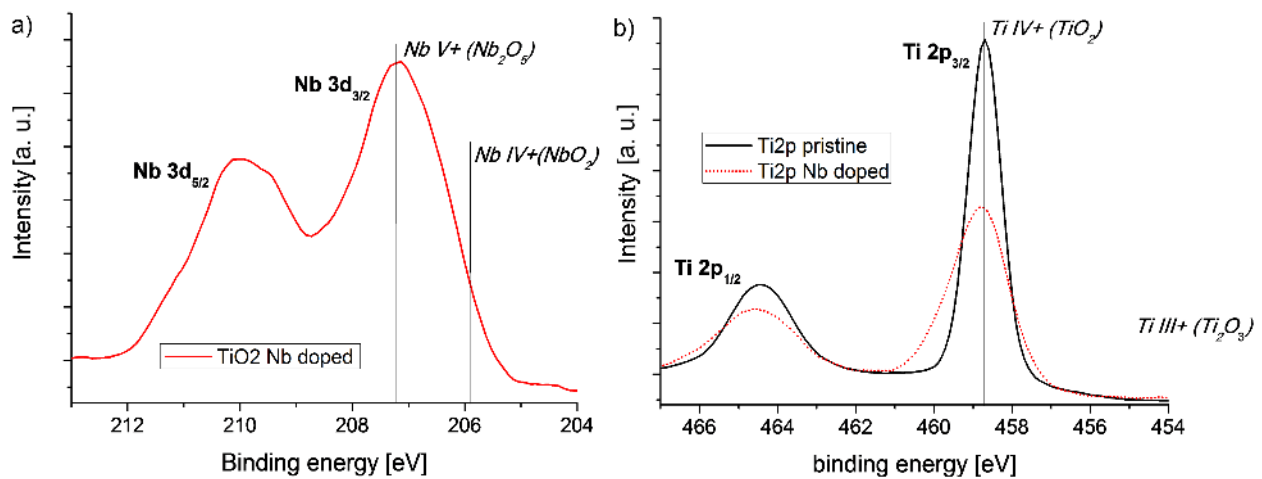


Figure 5: XPS spectra a) Nb 3d peaks and b) Ti 2p peaks comparison of doped and non-doped electrospun TiO₂ nanofibers.

From quantification measurements we obtained a Nb-to-Ti ratio of $\approx 16\%$ which denotes a slight agglomeration of Nb on the surface compared to the from weighed portion expected 10% bulk concentration. This suggests the segregation of Nb at the surface of electrospun TiO₂ fibers. According to the literature it is mainly the oxygen activity and oxygen affinity of the dopant that determines whether or not surface segregation occurs.^{52,53} The possibility of inhomogeneous Nb dopant distribution is not considered since the high solubility of Nb in TiO₂ is widely known.^{44,54,55}

Theoretical and experimental XAS

The Ti K-edge X-ray absorption near edge spectroscopy (XANES) spectra (4966 eV) of non-doped and Nb-doped TiO₂ are shown in Fig. 6. The absorption spectra reveal great resemblance, having the same number of pre-edge and edge features peaks and absorption edge position, merely small differences in intensity are observed. This is in good agreement with results from DFT calculations of the XAS spectra which illustrate that no significant changes of the Ti K-edge are expected under the influence of Nb doping, see Fig. 7. The observed differences between the relative amplitudes of the experimental and theoretical edge features are mainly due to energy dependent broadening effects caused by finite lifetime of the final electronic state which was not taken into account in the calculations for simplicity.

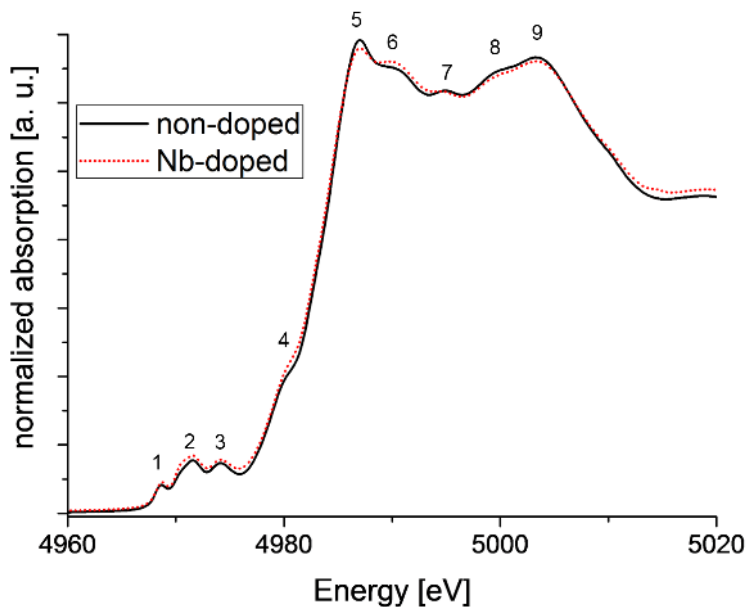


Figure 6: Measured Ti K-edge of XANES spectra of non-doped and Nb-doped TiO₂.

All the peaks of the experimental XANES spectra are reproduced by the calculations and can be assigned to main peaks of the Ti p DOS of the conduction band except for the pre-edge peak 1 which is purely quadrupolar as in the case of rutile TiO₂.⁵⁶ The pre-edge peaks 2 and 3 contain the dipolar components and reflect the t_{2g} - e_g splitting of Ti in TiO₆ octahedral environment.

Since the comparison of XANES at the Ti K-edge of non-doped and Nb-doped anatase provides little information on the incorporation of Nb within the anatase lattice, we turned our attention to

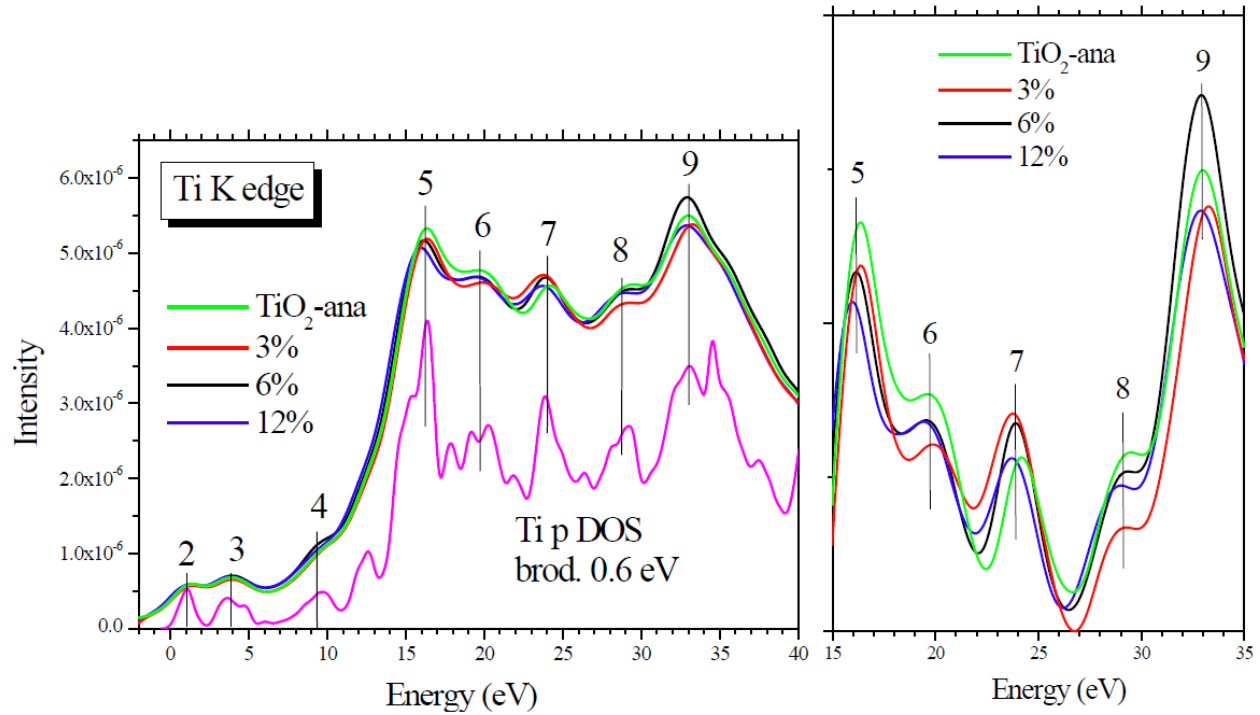


Figure 7: Theoretical XANES spectra at K-edge of non-doped and Nb-doped anatase TiO_2 (about 3, 6 and 12 at. % of Nb). The broadened Ti p DOS is shown for comparison. The origin of energy is taken at the Fermi level.

the EXAFS signal which bears information on local structure. The k^2 weighted EXAFS signal of doped and non-doped TiO_2 is shown in Fig. 8(a), and the derived Fourier transforms in Fig. 8(b). The spectra of non-doped and Nb-doped TiO_2 are rather similar, indicating that local structure is largely preserved. Only a slight decrease of the amplitude of the EXAFS spectrum of the Nb-doped sample is observed, corresponding to a slight decrease of the maxima in the Fourier transform, cf. Fig. 8(b).

According to the crystal structure of anatase, these spectra were fitted in the R-space using one oxygen coordination shell for the closest peak, and two titanium and one oxygen shell for the following ones. For the Nb-doped sample, Nb incorporation in the anatase lattice was taken into account through a combination of the FEFF Ti-Ti paths of anatase and corresponding Ti-Nb paths derived from Nb_2O_5 in which central Nb atom was replaced by Ti. A slight variation of interatomic distances is observed upon Nb insertion (cf. Tab. 1) which suggests that Nb predominately occupies Ti sites. Furthermore, slightly higher Debye-Waller factors are observed for Nb-doped than

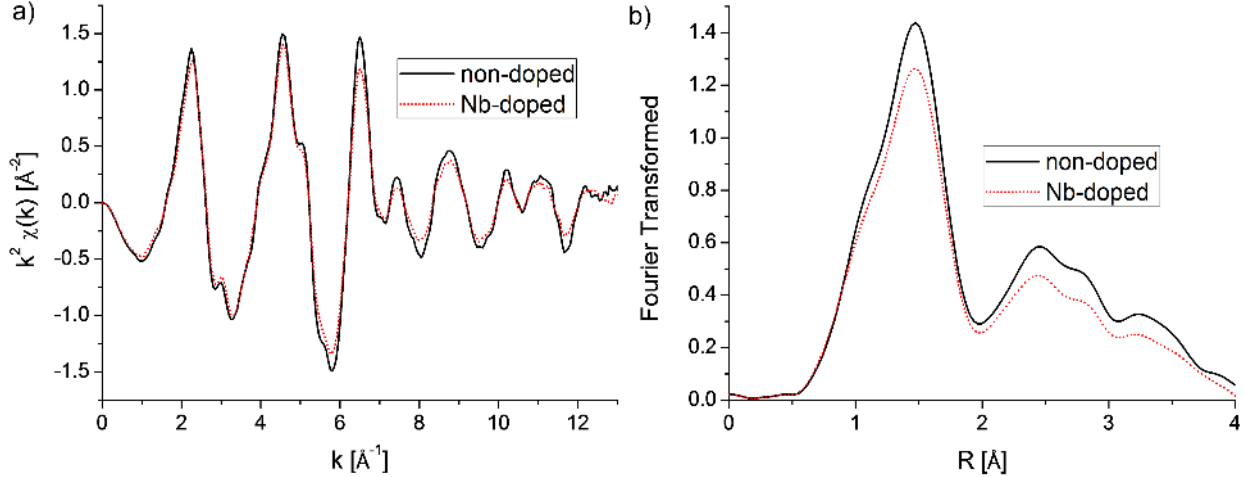


Figure 8: K^2 weighted a) EXAFS signal and b) Fourier transformed of non-doped and doped TiO_2 nanofibers. The latter revealing position of next neighbor shells of central Ti atom for Nb-doped and non-doped TiO_2 .

for non-doped TiO_2 , which reflects an increase in disorder of the structure due to incorporation of Nb in the anatase lattice. This explains the reduced intensity of the FT spectra of the doped sample, and confirms the substitutional doping of Nb.

Table 1: EXAFS Fitting Parameters for non-doped and 10% Nb-doped anatase

	atom	CN	interatomic distance [Å]		Debye-Waller factor [Å] ²	
			non-doped	Nb-doped	non-doped	Nb-doped
1	O	6	1.948(±0.005)	1.949(±0.004)	0.0055(±0.0006)	0.0069(±0.0006)
2	Ti/Nb	4	3.068(±0.007)	3.068(±0.007)	0.0066(±0.0009)	0.0067(±0.0009)
3	Ti/Nb	4	3.88(±0.10)	3.85(±0.08)	0.023(±0.022)	0.020(±0.024)
4	O	8	3.844(±0.034)	3.86(±0.06)	0.004(±0.003)	0.007(±0.004)

The R-factor reflects the absolute misfit between experimental data and the fitted curve. Its variation as a function of the degree of Nb doping imposed during the fitting procedure follow a stretched parabolic slope with a minimum at 22 at% Nb dopant, indicating the best agreement between measured and fitted EXAFS data, see Fig. 9. This result is in line with the assumption that Nb partially substitutes Ti atoms in the anatase lattice whereas the deviation of the peak minimum (≈ 22 at%) from the expected Nb-content (≈ 10 at%) is considered acceptable regarding the width of the minimum and assumptions and simplifications made in the computational approach.

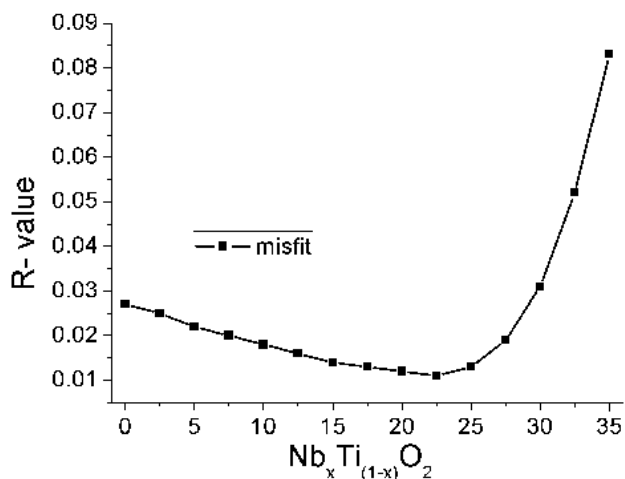


Figure 9: Misfit of EXAFS fit as function of extent of Nb doping of TiO₂

Electronic structure

The partial density of states (PDOS) of Nb 4d and Ti 3d are similar, but the additional Nb 4d electron (compared to the Ti 3d electron number) pushes the Fermi level from the top of the valence band for non-doped TiO₂ to the bottom of the conduction band for Nb-doped TiO₂, Fig.10. The lowest Nb 4d empty states are not in the band gap but overlap with Ti 3d empty states to form a band in the range 0-2 eV. The Fermi level is located within this band, which indicates that Nb-doped anatase TiO₂ has a metallic conductivity due to the delocalization of this additional Nb 4d electron as previously discussed.^{33,34,44,46,57}

Summing up the above stated findings we can assume that the main reason for the improved rate capability of Nb-doped TiO₂ lies in the enhancement of charge transfer in TiO₂. Substituting Ti⁺⁴ by aliovalent Nb⁺⁵ introduces additional charge carries which improve the low bulk conductivity of TiO₂. This effect is especially prominent at elevated cycling rates as kinetic limitations become more salient. Our results are very much in agreement with findings by *Han et al.*³¹ reporting a doubling of rate capability at increased cycling rate while absolute capacities are only slightly improved for aliovalent doped TiO₂ electrospun fibers. The possibilities to compare our results to those of other groups are however limited as few groups focus on studying the isolated effect of the dopant like we did. In most existing studies side contributions deriving from complex syn-

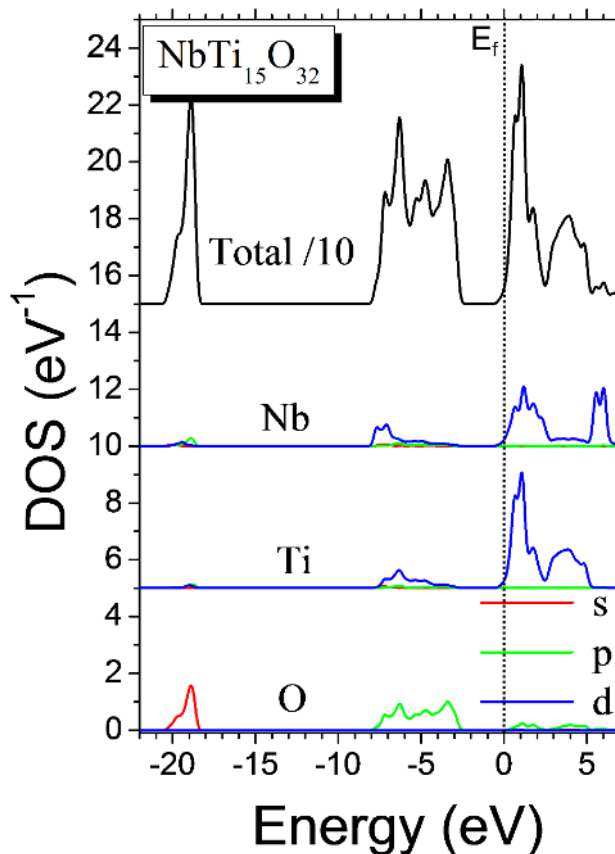


Figure 10: Calculated partial densities of states of NbTi₁₅O₃₂ ($\approx 6at\%$). The origin of energy is taken at the Fermi level (dotted line)

thesis method to achieve specific morphology^{26,32,58,59} or elaborate electrode formulation^{20,60–62} are superimposing to the effect of dopant and are therefore hardly comparable to our results. It is needless to say that a specific work on electrode formulation will probably bring rate capability and capacity retention improvements.

Conclusions

We showed that the Nb dopant is homogeneously distributed within the electrospun anatase fibers, and it is embedded in the lattice by substituting Ti. The statistical substitution of every tenth Ti atom by Nb has no influence on the global morphology and only slight changes of the lattice were observed via XRD and Raman. DFT calculations and XAS experimental results go hand in hand in showing that no relevant changes of Ti K absorption edge occur. EXAFS indicates an increased

disorder in doped samples which underlines the embedding of Nb by substitutional doping in the bulk of TiO₂. Galvanostatic cycling shows that the Li insertion mechanism remains largely unaffected by niobium doping of TiO₂ anatase. However when comparing the rate capability, superior performance of Nb-doped sample is evident. We ascribe this improvement predominantly to two of the observed Nb-doping induced effects; Firstly the global increase in the electronic conductivity which is related to the change in Fermi level position and secondly the decrease in crystal size that is beneficial for interparticle contact and reduces diffusion paths, which is crucial for high performance cycling. The ancillary effect of increasing specific surface area is likely to also contribute to the observed improvement.

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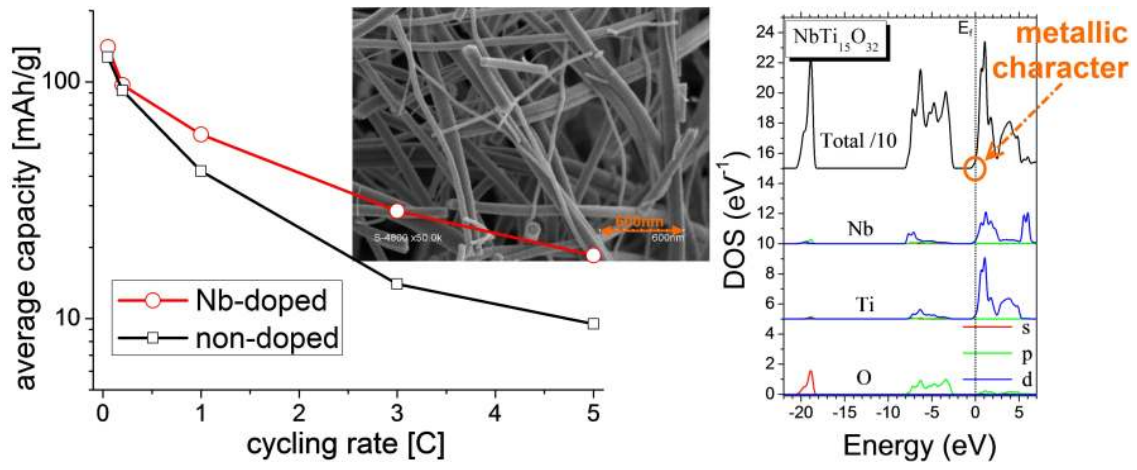


Figure 11: Table of Content figure