

NEGATIVE INDEX JACOBI FORMS AND QUANTUM MODULAR FORMS

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ABSTRACT. In this paper, we consider the Fourier coefficients of a special class of meromorphic Jacobi forms of negative index. Much recent work has been done on such coefficients in the case of Jacobi forms of positive index, but almost nothing is known for Jacobi forms of negative index. In this paper we show, from two different perspectives, that their Fourier coefficients have a simple decomposition in terms of partial theta functions. The first perspective uses the language of Lie super algebras, and the second applies the theory of elliptic functions. In particular, we find a new infinite family of rank-crank type PDEs generalizing the famous example of Atkin and Garvan. We then describe the modularity properties of these coefficients, showing that they are “mixed partial theta functions”, along the way determining a new class of quantum modular partial theta functions which is of independent interest.

1. INTRODUCTION AND STATEMENT OF RESULTS

Since the introduction of the theory of Jacobi forms by Eichler and Zagier [20], connections between Jacobi forms and modular-type objects have been a question of central interest, with applications to many areas including Siegel modular forms, mock modular forms, and Lie theory. In this paper we study the Fourier coefficients of a special family of negative index Jacobi forms. In particular, consider for $N \in \mathbb{N}$, $M \in \mathbb{N}_0$ the functions

$$\phi_{M,N}(z; \tau) := \frac{\vartheta(z + \frac{1}{2}; \tau)^M}{\vartheta(z; \tau)^N},$$

where $\vartheta(z; \tau)$ is the usual Jacobi theta function

$$(1.1) \quad \vartheta(z; \tau) := -i\zeta^{-\frac{1}{2}}q^{\frac{1}{8}}(q)_\infty(\zeta)_\infty(\zeta^{-1}q)_\infty.$$

Here, $q := e^{2\pi i\tau}$ ($\tau \in \mathbb{H}$), $\zeta := e^{2\pi iz}$ ($z \in \mathbb{C}$), and for $n \in \mathbb{N}_0 \cup \{\infty\}$, $(a)_n := (a; q)_n := \prod_{j=0}^{n-1} (1 - aq^j)$ is the q -Pochhammer symbol. The relationship between Jacobi forms and modular forms has appeared in many guises and stems back to important work on holomorphic Jacobi forms, which states that they have theta decompositions relating them to half-integral weight modular forms [20]. The situation for meromorphic positive index Jacobi forms well-understood as well; a meromorphic Jacobi form of positive index has Fourier coefficients which are *almost mock modular forms*, which in turn are holomorphic parts of *almost harmonic Maass forms* [8, 17, 21, 34, 46]. Loosely speaking, almost harmonic weak Maass forms are sums of harmonic weak Maass functions under iterates of the raising operator multiplied by almost holomorphic modular forms. In this paper, we describe new decompositions of the Jacobi forms $\phi_{M,N}(z; \tau)$ which complement this long history of previous work on positive index Jacobi forms in the much more mysterious case of negative index. In addition to being of interest in the subject of general Jacobi forms, here we give further applications of such decompositions, focusing on the special subfamily $\phi_N(\tau) := \phi_{0,N}(\tau)$ as they are of great interest in various areas such as number theory, representation theory, combinatorics, and physics. Here we outline just a few such occurrences.

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Firstly, for various choices of N , the functions ϕ_N are of combinatorial interest. In particular, the function ϕ_1 is related to the famous Andrews-Dyson-Garvan crank generating function (see (1.7) and (1.8)), which was used by Andrews and Garvan to provide a combinatorial explanation for the Ramanujan congruences for the partition function [4], as postulated by Dyson [18]. In this paper we describe relations between powers of the crank generating function with certain Appell-Lerch series, giving a new family of PDEs generalizing the “rank-crank PDE” of Atkin and Garvan [6] (see Theorem 1.3). This beautiful identity of Atkin and Garvan gives a surprising connection between the rank and crank generating functions which can be used to show various congruences relating ranks and cranks, as well as useful relations between the rank and crank moments [6]. We also note the other examples of similar PDEs related to combinatorics have shown up in, for example Section 3.2 of [9], where the function $\phi_{1,3}(z; \tau)$ is studied in relation to overpartitions.

Secondly, the functions ϕ_N contain information about certain affine vertex algebras and their associated affine Lie algebras studied by Kac and Wakimoto [24]. More precisely, let $\mathcal{S}(1)$ be the graded vector space

$$\mathcal{S}(1) := \text{Sym} \bigoplus_{n=0}^{\infty} \left(q^{n+\frac{1}{2}} (\zeta \mathbb{C} \oplus \zeta^{-1} \mathbb{C}) \right),$$

and $\mathcal{S}(N) = \mathcal{S}(1)^{\otimes N}$. The vector space $\mathcal{S}(N)$ can be given the structure of a vertex algebra, the bosonic $\beta\gamma$ -ghost vertex algebra of rank N and central charge $-N$. The graded character of this vertex algebra has a nice product form in the domain $|q|^{\frac{1}{2}} < |\zeta_n| < |q|^{-\frac{1}{2}}$, where $\zeta_n := e^{2\pi i z_n}$,

$$(1.2) \quad \text{ch}[\mathcal{S}(N)](z_1, \dots, z_N; \tau) = \prod_{n=1}^N q^{\frac{1}{24}} \text{ch}[\mathcal{S}(1)](z_n; \tau) = q^{\frac{N}{24}} \prod_{n=1}^N \frac{1}{\left(\zeta_n^{-1} q^{\frac{1}{2}} \right)_{\infty} \left(\zeta_n q^{\frac{1}{2}} \right)_{\infty}}.$$

It specializes to $\phi_N(z; \tau)$ for the following choices:

$$i^N \zeta^{\frac{N}{2}} q^{-\frac{N}{6}} (q)_{\infty}^{-N} \text{ch}[\mathcal{S}(N)] \left(z - \frac{\tau}{2}, \dots, z - \frac{\tau}{2}; \tau \right) = \phi_N(z; \tau).$$

The algebra $\mathcal{S}(N)$ contains as commuting subalgebras the rank one Heisenberg vertex algebra $\mathcal{H}(1)$ and the simple affine vertex algebra of $sl(N)$ at level minus one, $L_{-1}(sl(N))$. Note that minus one is not an admissible level in the case of $N = 2$ and that these algebras do not form a mutually commuting pair inside $\mathcal{S}(2)$. However for $N > 2$ it was shown in [1] that $L_{-1}(sl(N))$ and $\mathcal{H}(1)$ form such a mutually commuting pair inside $\mathcal{S}(N)$. The character of the highest-weight module \mathcal{F}_{μ} , $\mu \in \mathbb{R}$, of $\mathcal{H}(1)$ takes the form

$$(1.3) \quad \text{ch}[\mathcal{F}_{\mu}(z; \tau)] = \frac{\zeta^{\sqrt{N}\mu} q^{-\frac{\mu^2}{2}}}{q^{\frac{1}{24}}(q)_{\infty}},$$

so that the Fourier coefficients in ζ of $\phi_N(z; \tau)$ immediately allow one to compute the multiplicity with which the character of $\mathcal{F}_{r/\sqrt{N}}$ appears. In physics language, such a multiplicity is called the *branching function* of the coset $\mathcal{S}(N)/\mathcal{H}(1)$.

This leads to the second conformal field theory and vertex algebra importance of decomposing a meromorphic Jacobi form. One of the most interesting classes of vertex algebras is given by $V_k(\mathfrak{g})$, the universal affine vertex algebra of the simple Lie algebra \mathfrak{g} at level $k \in \mathbb{C}$. For certain rational admissible levels, $V_k(\mathfrak{g})$ is not simple and one instead prefers to study its simple quotient $L_k(\mathfrak{g})$. The characters of irreducible highest-weight modules at admissible level $L_k(\mathfrak{g})$ are the sum expansions in special domains of meromorphic Jacobi forms [26]. Understanding these sum expansions is crucial in studying the modular data of the corresponding conformal field theory [14, 15].

Fourthly, the functions ϕ_N appear in the denominator identities of affine Lie super algebras [27]. In [14, 15] the denominator identity of $\widehat{sl}(2|1)$ was an essential ingredient to study the relations

between characters of admissible level $L_k(\mathfrak{sl}(2))$, while we use the identities for the family $\widehat{\mathfrak{sl}}(N|1)$ to prove one of our central theorems.

Finally, the functions ϕ_N also occur in string theory; we only expound upon one example. The reciprocal of the Igusa cusp form $\Phi_{10}(Z)$ ($Z \in \mathbb{H}_2$, the Siegel upper half plane of genus 2) arises as the partition function of quarter-BPS dyons in the type II compactification on the product of a K_3 surface and an elliptic curve. Write

$$(1.4) \quad \frac{1}{\Phi_{10}(Z)} = \sum_{m=-1}^{\infty} \psi_m(z; \tau) \rho^m \quad \left(Z = \begin{pmatrix} \tau & z \\ z & w \end{pmatrix}, \rho = e^{2\pi i w} \right).$$

For $m > 0$, the Fourier coefficients of the functions ψ_m are the degeneracies of single-centered black holes and two-centered black holes with total magnetic charge invariant equal to m . This case is studied in pathbreaking work of Dabholkar, Murthy, and Zagier [17] (see also [32] for the appearance of mock modular forms in the context of quantum gravity partition functions and AdS3/CFT2, as well as [31] for a relation between multi-centered black holes and mock Siegel-Narain theta functions). The coefficient of $m = -1$ equals

$$\frac{-1}{\eta^{18}(\tau) \vartheta(z; \tau)^2}$$

(note that our theta function $\vartheta(z; \tau)$ differs from the theta function $\theta_1(z; \tau)$ in the notation of [17] by a factor of i). Analogously to the case of Jacobi forms of positive index, one may view Theorem 1.3 below as a decomposition giving a “polar part” but no “finite part” as described in [8, 17] and stated in more detail in (3.6). This is consistent with a string theoretic interpretation of ψ_{-1} in (1.4) in that there are no single-centered black holes and the degeneracies are all interpreted as accounting for two-centered black holes (see [16, 36]). In contrast, in the case $m > 0$, the mock part of ψ_m corresponds to single-centered black holes and the Appel Lerch sum corresponds to two-centered black hole bound states [17].

Returning to the problem of studying $\phi_{M,N}$, we define its Fourier coefficients by

$$\phi_{M,N}(z; \tau) =: \sum_{r \in \frac{M-N}{2} + \mathbb{Z}} \chi(M, N, r; \tau) \zeta^r$$

and in particular we set $\chi(N, r; \tau) := \chi(0, N, r; \tau)$. Note that wallcrossing occurs; the coefficients $\chi(M, N, r; \tau)$ are only well-defined if we fix a range for z . We show that the Fourier coefficients $\chi(M, N, r; \tau)$ can be described using *partial theta functions* (i.e., sums over half a lattice which when summed over a full lattices becomes a theta function), whose modularity properties near the real line we also describe using quantum modular forms. Quantum modular forms were recently defined by Zagier in [44] (see also [10, 12, 22, 30]). Although the definition is not rigorous, Zagier gave a number of motivating examples. Roughly speaking, a *weight k quantum modular form* is a function $f: \mathcal{Q} \rightarrow \mathbb{C}$ for some subset $\mathcal{Q} \subseteq \mathbb{P}_1(\mathbb{Q})$ such that for any γ in a congruence subgroup Γ , the cocycle $f|_k(1 - \gamma)$ extends to an open set of \mathbb{R} and is “nice” (e.g. continuously differentiable, smooth, etc). In fact, our study of the modularity of the partial theta functions shows that they are what Zagier refers to as strong quantum modular forms, namely that they have a near-modular property for asymptotic expansions defined at every point in a subset of $\mathbb{P}_1(\mathbb{Q})$. Moreover, this behavior comes from the “leaking” of modularity properties of a non-holomorphic Eichler integral defined on the lower half plane (see (6.7)).

Returning to the Fourier coefficients of $\phi_{M,N}$, we define a *mixed partial theta function* to be a linear combination of quasi-modular forms multiplied with partial theta functions. These functions have known connections to many interesting combinatorial functions, such as concave and convex compositions [3], unimodal sequences [25, 38], and stacks [42]. Throughout, we abuse notation to say that any function is a modular form, partial theta function, mixed partial theta function, etc.

if it is equal to such a function up to multiplication by a rational power of q . Our main result is the following.

Theorem 1.1. *For any $N \in \mathbb{N}$, $M \in 2\mathbb{N}_0$ and $r \in \mathbb{Z}$ with $M < N$, the functions $\chi(M, N, r; \tau)$ are mixed partial theta functions.*

Remarks. 1. *The quasi-modular forms appearing in the decomposition of the mixed partial theta functions are canonically determined by the Laurent expansion of the Jacobi form $\phi_{M,N}$ (see Theorem 1.4).*

2. *Using the techniques of this paper it is easy to relax the condition on M to allow any natural number less than N , however we restrict to even M for notational convenience.*

We first consider the case of $\phi_N(z; \tau)$, which we study from two perspectives. Our first viewpoint describes the Fourier coefficients as derivatives of partial theta functions of a rescaled version of the root lattice

$$A_{N-1} := \mathbb{Z}\alpha_1 \oplus \cdots \oplus \mathbb{Z}\alpha_{N-1}$$

of $\mathfrak{sl}(N)$. Here, the α_n are the simple roots of $\mathfrak{sl}(N)$, which are linear functionals on the Cartan subalgebra $\mathfrak{h} \cong \mathbb{C}^{N-1}$ of $\mathfrak{sl}(N)$. The Gram matrix of A_{N-1} is the Cartan matrix of $\mathfrak{sl}(N)$. We denote the bilinear form by $(\cdot | \cdot)$ and abbreviate $t^2 := (t|t)$ for t in A_{N-1} . For r in \mathbb{Z} , we define the subset of $\frac{1}{N}A_{N-1}$

$$(1.5) \quad T_r := \left\{ \sum_{n=1}^{N-1} a_n \alpha_n \mid a_n = \frac{(N-n)n}{2} - \frac{rn}{N} + (N-1)m_n, m_n \in \mathbb{Z}, \left(m_{N-1} - \frac{1}{2}\right) \left(r - \frac{1}{2}\right) \geq 0 \right\}$$

and its *partial theta function*

$$P_r(\tau) := \sum_{t \in T_r} e^t q^{\frac{t^2}{2(N-1)}}.$$

The e^t are functions on the Cartan subalgebra \mathfrak{h} defined by $e^t : u \mapsto e^{t(u)}$ for u in \mathfrak{h} . Its evaluation for $u \in \mathfrak{h}$ is then denoted by $P_r(u; \tau)$. We call P_r a partial theta function, because the theta function obtained by summing over the complete lattice $\frac{1}{N}A_{N-1}$,

$$\theta_N(\tau) := \sum_{t \in \frac{1}{N}A_{N-1}} e^{t(0)} q^{\frac{t^2}{2(N-1)}},$$

is a modular form of weight $(N-1)/2$ for $\Gamma(M)$ with $M = N^2(N-1)/2$. This statement is true, since θ_N is the theta function of the lattice $\frac{1}{\sqrt{2M}}A_{N-1}$. The level of this lattice is M , and the modularity of theta functions of lattices is discussed for example in [19].

Further let ∂ be the differential operator

$$(1.6) \quad \partial_\alpha e^\lambda := (\lambda|\alpha) e^\lambda \quad \text{and} \quad \partial := \prod_{\alpha \in \Delta_0^+} \partial_\alpha,$$

where Δ_0^+ is the set of positive roots of $\mathfrak{sl}(N)$. Finally, set $d_N := \prod_{j=1}^N j!$ and let $\text{sign}(r) = 1$ if $r \geq 0$ and -1 otherwise. Then we have the following.

Theorem 1.2. *For $N \geq 2$, the r -th Fourier coefficient of $\phi_N(z; \tau)$ is given by*

$$\begin{aligned} \chi(N, r; \tau) &= \frac{i^N \text{sign}\left(r - \frac{N}{2}\right) q^{\frac{-r^2}{2N}}}{\eta(\tau)^{N(N+1)} d_{N-1}} \sum_{t \in T_{r - \frac{N}{2}}} \prod_{\alpha \in \Delta_0^+} (t|\alpha) q^{\frac{t^2}{2(N-1)}} \\ &= \frac{i^N \text{sign}\left(r - \frac{N}{2}\right) q^{\frac{-r^2}{2N}}}{\eta(\tau)^{N(N+1)} d_{N-1}} \partial P_{r - \frac{N}{2}}(\tau) \Big|_{\{e^t = 1 \mid t \in \frac{1}{N}A_{N-1}\}}. \end{aligned}$$

Remarks. 1. Using (1.2) and (1.3), Theorem 1.2 implies the character decomposition

$$ch[\mathcal{S}(N)](z, \dots, z; \tau) = \sum_{r \in \mathbb{Z}} ch \left[\mathcal{F}_{\frac{r}{\sqrt{N}}}(z; \tau) \right] ch[\mathcal{B}_r](\tau),$$

where

$$ch[\mathcal{B}_r](\tau) := \frac{\text{sign}(r)}{\eta(\tau)^{N^2-1} d_{N-1}} \sum_{t \in T_r} \prod_{\alpha \in \Delta_0^+} (t|\alpha) q^{\frac{t^2}{2(N-1)}}.$$

The $ch[\mathcal{B}_r](\tau)$ are then characters of $L_{-1}(\mathfrak{sl}(N))$.

2. The proof of the theorem uses the denominator identity of both $\mathfrak{sl}(N|1)$ and $\widehat{\mathfrak{sl}}(N|1)$ as well as Weyl's character formula for $\mathfrak{sl}(N)$.

3. The case $N = 1$ follows from the denominator identity of $\widehat{\mathfrak{gl}}(1|1)$ (see Proposition 4.4). In this case, the Fourier coefficients relate to the characters of a well-known logarithmic conformal field theory, the $\mathcal{W}(2,3)$ -algebra of central charge -2 . The modularity of the coefficients has been studied from a different perspective in [13].

The second approach is based on a generalization of a deep identity of Atkin and Garvan. To state it, we first recall the rank and crank generating functions (whose combinatorial meanings are not needed in this paper), which arise in many contexts and in particular give combinatorial explanations of Ramanujan's congruences (for example see [4, 7, 18]). Specifically, the generating functions are given as follows:

$$(1.7) \quad \begin{aligned} \mathcal{R}(\zeta; q) &:= \frac{(1-z)}{(q)_\infty} \sum_{n \in \mathbb{Z}} \frac{(-1)^n q^{\frac{n}{2}(3n+1)}}{1 - \zeta q^n}, \\ \mathcal{C}(\zeta; q) &:= \frac{(q)_\infty}{(\zeta q)_\infty (\zeta^{-1} q)_\infty}. \end{aligned}$$

We also need the normalized versions

$$\mathcal{R}^*(\zeta; q) := \frac{\zeta^{\frac{1}{2}} q^{-\frac{1}{24}} \mathcal{R}(\zeta; q)}{1 - \zeta}, \quad \mathcal{C}^*(\zeta; q) := \frac{\zeta^{\frac{1}{2}} q^{-\frac{1}{24}} \mathcal{C}(\zeta; q)}{1 - \zeta}.$$

Note that ϕ_N is essentially the N -th power of \mathcal{C}^* as for $N \in \mathbb{N}$ we have

$$(1.8) \quad \phi_N(z; \tau) = i^N \eta(\tau)^{-2N} \mathcal{C}^*(\zeta; q)^N.$$

The simplest case of our decomposition relies on the fact that \mathcal{C}^* , and thus ϕ_1 , is essentially an Appell-Lerch sum thanks to the following classical partial fractions expansion (for example, see Theorem 1.4 of [47]).

$$(1.9) \quad \mathcal{C}^*(\zeta; q) = \frac{\zeta^{\frac{1}{2}}}{\eta(\tau)} \sum_{n \in \mathbb{Z}} \frac{(-1)^n q^{\frac{n(n+1)}{2}}}{1 - \zeta q^n}.$$

For the cube of the crank generating function, Atkin and Garvan [6] proved the following rank-crank PDE which is very useful in establishing congruences and relations between the moments of the rank and crank generating functions:

$$(1.10) \quad \zeta(q)_\infty^2 \mathcal{C}^*(\zeta; q)^3 = \left(3\mathcal{D}_q + \frac{1}{2}\mathcal{D}_\zeta + \frac{1}{2}\mathcal{D}_\zeta^2 \right) \mathcal{R}^*(\zeta; q).$$

Here and throughout $\mathcal{D}_x := x \frac{\partial}{\partial x}$. Note that this gives a description of ϕ_3 in terms of Appell-Lerch sums by (1.8).

Zwegers [47] nicely generalized (1.10) for arbitrary odd powers of the crank generating function using the theory of elliptic forms. For similar results using another clever proof, see also the paper of Chan, Dixit, and Garvan [11].

In this paper, we prove a new family of analogous PDEs which are of independent interest. Moreover, we package Zwegers' family of PDEs in a way which illuminates their structure coming from negative index Jacobi forms. To describe this, we need the Appell-Lerch sums

$$(1.11) \quad F_N(z, u; \tau) := \zeta^{\frac{N}{2}} w^{\frac{N}{2}} \sum_{n \in \mathbb{Z}} \frac{(-w)^{Nn} q^{\frac{N}{2}n(n+1)}}{1 - \zeta w q^n},$$

where $w := e^{2\pi i u}$. We also require the Laurent coefficients of $\phi_{M,N}(z; \tau)$ at $z = 0$:

$$(1.12) \quad \phi_{M,N}(z; \tau) = \frac{D_N(\tau)}{(2\pi i z)^N} + \frac{D_{N-2}(\tau)}{(2\pi i z)^{N-2}} + \dots + O(1).$$

Note that only even or odd Laurent coefficients occur, depending on the parity of N , since $\vartheta(-z; \tau) = -\vartheta(z; \tau)$. It is not hard to see that the coefficients D_j are quasimodular forms. Explicitly, they can be computed quickly in terms of the usual Eisenstein series

$$G_k(\tau) := -\frac{B_k}{2k} + \sum_{n \geq 1} \sigma_{k-1}(n) q^n,$$

where $\sigma_k(n) := \sum_{d|n} d^k$ and B_k is the usual k -th Bernoulli number. Specifically, it easily follows from the Jacobi triple product formula that

$$\vartheta(z; \tau) = -2\pi z \eta^3(\tau) \exp \left(-2 \sum_{k \geq 1} G_{2k}(\tau) \frac{(2\pi i z)^{2k}}{(2k)!} \right).$$

The following result puts Zwegers family of PDEs as well as our new family of PDEs into a common framework. Setting

$$\delta_e := \begin{cases} 0 & \text{if } N \in 2\mathbb{N} - 1, \\ 1 & \text{if } N \in 2\mathbb{N}, \end{cases}$$

we find:

Theorem 1.3. *For any $N \in \mathbb{N}$, $M \in \mathbb{N}_0$, we have*

$$\phi_{2M, N+2M}(z; \tau) = (-1)^{1+\delta_e} \sum_{j=0}^{\frac{N-1-\delta_e}{2} + M} \frac{D_{2j+\delta_e+1}(\tau)}{(2j+\delta_e)!} \mathcal{D}_w^{2j+\delta_e} (F_N(z, u; \tau)) \Big|_{w=1}.$$

Remark. *Note that Theorem 1.3 is more explicit than Zwegers' rank-crank type PDEs as it gives the modular coefficients of the PDEs directly from the structure of the Jacobi form $\phi_{M,N}$. Chan, Dixit, and Garvan also remarked that it would be interesting to find such an explicit expression for the quasimodular forms in the decomposition in that case.*

Armed with the decomposition in Theorem 1.3 into Appell-Lerch sums, we can easily pick off the Fourier coefficients $\chi(M, N, r; \tau)$ and write them in terms of the Laurent coefficients of $\phi_{M,N}$ and certain partial theta functions

$$\Theta_{\frac{1}{2}+\delta_e}(N, r; \tau) := \sum_{n \geq 0} (-1)^{Nn} \left(n + \frac{r}{N} \right)^{\delta_e} q^{\frac{N}{2} \left(n + \frac{r}{N} \right)^2}.$$

Specifically, if we let

$$\rho(r) := \begin{cases} r & \text{if } r \geq \frac{N}{2}, \\ N - r & \text{if } r < \frac{N}{2}, \end{cases}$$

then the Fourier coefficients of $\phi_{M,N}$ are as follows.

Theorem 1.4. *For any $N \in \mathbb{N}$, $M \in \mathbb{N}_0$, $r \in \frac{N}{2} + \mathbb{Z}$, $0 \leq \text{Im}(z) < \text{Im}(\tau)$, we have*

$$\chi(2M, N + 2M, r; \tau) = (-1)^{1+\delta_e} q^{-\frac{r^2}{2N}} \sum_{j=0}^{\frac{N-1-\delta_e}{2}+M} \frac{D_{2j+\delta_e+1}(\tau)}{(2j+\delta_e)!} N^{j+\delta_e} 2^j \mathcal{D}_q^j \left(\Theta_{\frac{1}{2}+\delta_e}(N, \rho(r); \tau) \right).$$

If $N > 1$ is odd, these partial theta functions fit into the pioneering work of Folsom, Ono, and Rhoades [22] which gives startling relations between the asymptotic expansions of the rank and crank generating functions, generalizing and proving beautiful formulas of Ramanujan. Their work shows that $\Theta_{\frac{1}{2}}(N, r; \tau)$ is a strong quantum modular form for odd $N > 1$. Although their theorem does not directly apply for $N = 1$, in this case we essentially obtain an eta quotient which is trivially a quantum modular form at cusps where it vanishes.

For even N , both the hypergeometric representations used to determine quantum sets and the proof of quantum modularity are not applicable. Here we use the innovative approach of Lawrence and Zagier [29] to study quantum modularity properties (see also [45]). A key ingredient in our investigation is a beautiful identity of Warnaar [39] which relates certain partial and false theta functions (see (6.4)). Our main result for studying quantum modularity for even N is the following, which gives a new family of quantum modular forms.

Theorem 1.5. *For any $N \in 2\mathbb{N}$ and $r \in \mathbb{Z}$, $\Theta_{\frac{3}{2}}(N, r; \tau)$ is a strong quantum modular form with quantum set $\widehat{Q}_{N,r}$ (defined in (6.3)) on $\Gamma_1(2N)$, multiplier system χ_r (defined in (3.5)), and weight $\frac{3}{2}$.*

Remarks. 1. *More details about the specific quantum modular properties can be found in the proof of Theorem 1.5 in Section 6.2.*

2. *More generally, using Proposition 3 of [43], our proof of Theorem 1.5 shows that $\Theta_{\frac{3}{2}}(N, r; \tau)$ has modularity properties on all of \mathbb{Q} . For this, we note that although the function is not defined on all of \mathbb{Q} , it has a well-defined asymptotic expansion at all points in \mathbb{Q} . This expansion still agrees with the non-holomorphic Eichler integral on the lower half plane (see Section 6), so one could say that $\Theta_{\frac{3}{2}}(N, r; \tau)$ is a quantum modular form on \mathbb{Q} if we allow “poles” at certain points in \mathbb{Q} .*

The paper is organized as follows. In Sections 2 and 3, we review the necessary notation and basic objects from Lie theory, Jacobi forms, and quantum modular forms. We give our first proof of the decomposition using Lie theory in Section 4 and our second proof using an analogue of the rank-crank PDE in Section 5. We conclude by describing the quantum modular properties of $\Theta_{\frac{1}{2}+\nu}(N, r; \tau)$ in Section 6.

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Acknowledgments to be added later.

2. PRELIMINARIES ON LIE SUPER ALGEBRAS AND CHARACTER IDENTITIES

In this section, we recall some known facts of the affine Lie superalgebra $\widehat{\mathfrak{sl}}(N|1)$, following [27], as well as the finite dimensional Lie algebra $\mathfrak{sl}(N)$ using [23].

2.1. The Lie super algebra $sl(N+1|1)$. In this subsection the Lie super algebra $sl(N+1|1)$ and its root system are defined.

The even subalgebra of the Lie super algebra $sl(N+1|1)$ is $gl(N+1)$ and the odd part decomposes into the standard representation of the even subalgebra and its conjugate. In order to define the Lie super algebra, it is convenient to first introduce its root system. It lies in the lattice

$$L_N := \mathbb{Z}\varepsilon_1 \oplus \cdots \oplus \mathbb{Z}\varepsilon_{N+1} \oplus \mathbb{Z}\delta$$

with bilinear form

$$(\varepsilon_j|\varepsilon_k) := \delta_{j,k}, \quad (\varepsilon_j|\delta) := 0, \quad (\delta|\delta) := -1.$$

Thus, its signature is $(N+1, 1)$. The set of roots is $\Delta = \Delta_0 \cup \Delta_1 \subset L_N$, where the set of even roots (resp. odd roots) is denoted by Δ_0 (resp. Δ_1). They are

$$(2.1) \quad \Delta_0 := \{\varepsilon_j - \varepsilon_k | 1 \leq j, k \leq N+1, j \neq k\}, \quad \Delta_1 := \{\varepsilon_j - \delta, \delta - \varepsilon_j | 1 \leq j \leq N+1\}.$$

It is useful to split these sets into positive and negative subroot spaces, where

$$\begin{aligned} \Delta_0^+ &:= \{\varepsilon_j - \varepsilon_k | 1 \leq j < k \leq N+1\}, & \Delta_1^+ &:= \{\varepsilon_j - \delta | 1 \leq j \leq N+1\}, \\ \Delta^+ &:= \Delta_0^+ \cup \Delta_1^+, & \Delta &:= \Delta^+ \cup \Delta^-. \end{aligned}$$

A distinguished system of simple positive roots is then chosen to be

$$\Pi := \{\alpha_j = \varepsilon_j - \varepsilon_{j+1}, \beta = \varepsilon_{N+1} - \delta | 1 \leq j \leq N\}.$$

The α_j are even roots and β is the only distinguished odd simple root. The inner products of simple positive roots are

$$(\alpha_j|\alpha_k) = \begin{cases} 2 & \text{if } j = k, \\ -1 & \text{if } j = k \pm 1, \\ 0 & \text{otherwise,} \end{cases} \quad (\beta|\alpha_j) = -\delta_{j,N}, \quad (\beta|\beta) = 0.$$

Hence β is an isotropic root. Simple even roots generate the even root lattice

$$A_N := \mathbb{Z}\alpha_1 \oplus \cdots \oplus \mathbb{Z}\alpha_N.$$

Its dual lattice is

$$A'_N := \mathbb{Z}\lambda_1 \oplus \cdots \oplus \mathbb{Z}\lambda_N,$$

where the inner product of the fundamental weights λ_j with simple roots is $(\lambda_j|\alpha_k) = \delta_{j,k}$ and $(\lambda_j|\delta) = 0$. Roots and weights act on the Cartan subalgebra, which is

$$\mathfrak{h} := \bigoplus_{\alpha \in \Pi} \mathbb{C}h_\alpha = \mathfrak{h}_0 \oplus \mathbb{C}h_\beta, \quad \mathfrak{h}_0 := \mathbb{C}h_{\alpha_1} \oplus \cdots \oplus \mathbb{C}h_{\alpha_N},$$

with basis $\{h_\alpha\}$ parameterized by simple positive roots, and \mathfrak{h}_0 the Cartan subalgebra of $sl(N+1)$. The fundamental weights λ_j are identified with elements of the dual \mathfrak{h}_0^* of \mathfrak{h}_0 via $\lambda_j(h_{\alpha_k}) = \delta_{j,k}$. A bilinear form $(\ , \)$ on \mathfrak{h} is induced from the form on its dual space via

$$(h_\alpha, h_{\alpha'}) := (\alpha|\alpha').$$

We remark that the Lie superalgebra $sl(N+1|1)$ is then the $\mathbb{Z}/2\mathbb{Z}$ -graded algebra generated by $\{h_\alpha, e_\alpha^\pm | \alpha \in \Pi\}$ subject to the Serre-Chevalley relations (2.2) and the graded Jacobi identity. The parity of h_α and $e_{\alpha_j}^\pm$ is even, while the e_β^\pm are odd. We denote the graded anti-symmetric bracket by $[\ , \]: sl(N+1|1) \times sl(N+1|1) \rightarrow sl(N+1|1)$. Then the Serre-Chevalley relations of the algebra are

$$(2.2) \quad [h_\alpha, h_{\alpha'}] = 0, \quad [h_\alpha, e_{\alpha'}^\pm] = \pm(\alpha|\alpha')e_{\alpha'}^\pm, \quad [e_\alpha^+, e_{\alpha'}^-] = \delta_{\alpha, \alpha'}h_\alpha, \quad (\text{ad } e_\alpha^\pm)^{1-(\alpha|\alpha')}e_{\alpha'}^\pm = 0$$

for all $\alpha, \alpha' \in \Pi$ and $\alpha \neq \alpha'$ in the last equation. The bilinear form $(\ , \)$ on \mathfrak{h} can be extended to an invariant non-degenerate graded symmetric form on $sl(N+1|1)$, which we also denote by $(\ , \)$.

2.2. The even Weyl group and denominator identity of $sl(N+1|1)$. We now introduce the even Weyl group and the denominator identity of the Lie super algebra $sl(N+1|1)$.

For this, we first need to define the Weyl vector ρ . It is the difference of the even Weyl vector ρ_0 and the odd one ρ_1 ; namely

$$\rho_0 := \frac{1}{2} \sum_{\alpha \in \Delta_0^+} \alpha = \frac{1}{2} \sum_{j=1}^{N+1} (N+2-2j)\varepsilon_j, \quad \rho_1 := \frac{1}{2} \sum_{\alpha \in \Delta_1^+} \alpha = \frac{1}{2} \left(-(N+1)\delta + \sum_{j=1}^{N+1} \varepsilon_j \right),$$

$$\rho := \rho_0 - \rho_1 = \frac{1}{2} \left((N+1)\delta + \sum_{j=1}^{N+1} (N+1-2j)\varepsilon_j \right).$$

The group of even Weyl reflections W^\sharp acts on the dual of the even root lattice, A'_N , and is generated by $\sigma_j, j = 1, \dots, N$ defined by

$$\sigma_j: A'_N \rightarrow A'_N, \quad \sigma_j(\lambda) := \lambda - (\alpha_j | \lambda) \alpha_j.$$

This action naturally extends to the lattice L_N via $\sigma_j(\delta) = 0$ and

$$\sigma_j(\varepsilon_k) := \varepsilon_k - (\varepsilon_k | \alpha_j) \alpha_j = \varepsilon_k - (\delta_{j,k} - \delta_{j+1,k})(\varepsilon_j - \varepsilon_{j+1}) = \begin{cases} \varepsilon_{j+1} & j = k, \\ \varepsilon_j & j = k - 1, \\ \varepsilon_k & \text{otherwise.} \end{cases}$$

Hence the even Weyl group W^\sharp is just the group S_{N+1} permuting the ε_j . Orthonormality of the ε_j implies that the even Weyl group preserves the bilinear form $(\cdot | \cdot)$. Following [27] we define

Definition. A **regular exponential function** on A'_N is a finite linear combination of exponentials of the form e^λ for $\lambda \in A'_N$. A **rational exponential function** is the quotient A/B of two regular exponential functions A and $B \neq 0$. The even Weyl group W^\sharp acts on the field of these functions as $e^\lambda \mapsto e^{w(\lambda)}$ for any $w \in W^\sharp$. The **Weyl denominator** of $sl(N+1|1)$ is the rational exponential function

$$R = \frac{\prod_{\alpha \in \Delta_0^+} (1 - e^{-\alpha})}{\prod_{\alpha \in \Delta_1^+} (1 + e^{-\alpha})}.$$

We saw that the even Weyl group W^\sharp is just S_{N+1} , the signum of an element w in W^\sharp is $\sigma(w) := (-1)^n$ if w can be written as a composition of n transpositions. Theorem 2.1 of [27] applied to our situation gives

Lemma 2.1. *The Weyl denominator of the Lie super algebra $sl(N+1|1)$ is*

$$e^\rho R = \sum_{w \in W^\sharp} \sigma(w) w \left(\frac{e^\rho}{(1 - e^{\delta - \varepsilon_{N+1}})} \right).$$

2.3. The denominator identity of the affine Lie super algebra $\widehat{sl}(N+1|1)$. We turn our focus to the affinization of $sl(N+1|1)$, that is

$$\widehat{sl}(N+1|1) := \mathbb{C}[t, t^{-1}] \otimes sl(N+1|1) \oplus \mathbb{C}C \oplus \mathbb{C}d$$

with bracket

$$[t^n \otimes x, t^m \otimes y] := t^{n+m} \otimes [x, y] + n\delta_{n+m,0}(x, y)C, \quad [d, t^n \otimes x] := nt^n \otimes x, \quad [C, t^n \otimes x] := [C, d] := 0$$

for all $x, y \in \mathfrak{sl}(N+1|1)$ and $n, m \in \mathbb{Z}$. The Cartan subalgebra extends to its affine counterpart

$$\widehat{\mathfrak{h}} := \mathfrak{h} \oplus \mathbb{C}d \oplus \mathbb{C}C$$

and the bilinear form extends as

$$(t^n \otimes x, t^m \otimes y) := (x, y)\delta_{n+m,0}, \quad (C, d) := 1, \quad (C, C) := (d, d) := (C, t^n \otimes x) := (d, t^n \otimes x) := 0.$$

We identify C and d with linear functionals on \mathfrak{h} using the bilinear form $(\ , \)$ and extend A'_N to

$$\widehat{A}'_N := A'_N \oplus \mathbb{Z}d \oplus \mathbb{Z}C.$$

The bilinear form extends as

$$(C|d) := 1 \quad \text{and} \quad (C|C) := (d|d) = (C|\lambda) = (d|\lambda) = 0 \quad \forall \lambda \in A'_N.$$

The lattice $A_N \subset A'_N$ is then also a sublattice of \widehat{A}'_N . The affine Weyl vector is

$$\widehat{\rho} := \rho + Nd.$$

Note that N is the dual Coxeter number of $\mathfrak{sl}(N+1|1)$. For $\alpha \in A_N$, we define

$$(2.3) \quad t_\alpha: \widehat{A}'_N \rightarrow \widehat{A}'_N, \quad \lambda \mapsto \lambda + (\lambda|C)\alpha - \left((\lambda|\alpha) + \frac{1}{2}(\alpha|\alpha)(\lambda|C) \right) C.$$

The group of even Weyl translations is $\{t_\alpha | \alpha \in A_N\}$. Conjugation by a Weyl rotation gives for any $w \in W^\sharp$, $\alpha \in A_N$

$$(2.4) \quad w \circ t_\alpha \circ w^{-1} = t_{w(\alpha)}.$$

Let

$$Y := \left\{ h \in \widehat{\mathfrak{h}} \mid \operatorname{Re}(C(h)) > 0 \right\}$$

be the domain of all elements in $\widehat{\mathfrak{h}}$ on which the action of C has positive real part. Let $\widehat{\mathcal{F}}$ be the field of meromorphic functions on Y and define $\mathfrak{q} := e^{-C}$. Thus $|\mathfrak{q}(y)| < 1$ for all y in Y . Any element λ of L' extends to a linear function on $\widehat{\mathfrak{h}}^*$ by defining $\lambda(C) = \lambda(d) = 0$. In this way rational exponential functions on L' embed in $\widehat{\mathcal{F}}$.

Definition. *The denominator of $\widehat{\mathfrak{sl}}(N+1|1)$ is*

$$\widehat{R} = R \prod_{j=1}^{\infty} (1 - \mathfrak{q}^j)^{n+1} \prod_{\alpha \in \Delta_0} (1 - \mathfrak{q}^j e^\alpha) \prod_{\alpha \in \Delta_1} (1 + \mathfrak{q}^j e^\alpha)^{-1}.$$

We need Theorem 4.1 of [27], which states

Lemma 2.2. *The denominator of $\widehat{\mathfrak{sl}}(N+1|1)$ satisfies*

$$e^{\widehat{\rho}} \widehat{R} = \sum_{\alpha \in A_N} t_\alpha \left(e^{\widehat{\rho}} R \right).$$

2.4. The Weyl character formula of $\mathfrak{sl}(N+1)$. We also require a well-known variant of the dimension formula, which itself is a corollary of the famous character formula of Weyl [41]. Let $\lambda = m_1 \lambda_1 + \cdots + m_N \lambda_N$ be a dominant weight of $\mathfrak{sl}(N+1)$; that is, all m_j are natural numbers. Letting V_λ be the corresponding irreducible highest-weight module, then the character formula is in our notation

$$(2.5) \quad \operatorname{ch}[V_\lambda] = \frac{\sum_{w \in W^\sharp} \sigma(w) e^{w(\lambda + \rho_0)}}{e^{\rho_0} \prod_{\alpha \in \Delta_0^+} (1 - e^{-\alpha})}.$$

Since ρ_1 is W^\sharp invariant we can replace ρ_0 by ρ in this formula. Let $m \in \mathbb{N}$ and let v be the linear map from the regular exponential functions on $\frac{1}{m}A'_N$ to the complex numbers defined by $v(e^\lambda) = 1$ for every $\lambda \in \frac{1}{m}A'_N$. Let V_λ be the irreducible finite dimensional highest-weight module of highest-weight λ . Hence $v(\text{ch}[V_\lambda])$ is just the dimension of this module. The application of v to both nominator and denominator of the character formula (2.5) vanishes, but the quotient is finite. Using (1.6), we find [23]

$$v(\text{ch}[V_\lambda]) = \frac{v\left(\partial \sum_{w \in W^\sharp} \sigma(w) e^{w(\lambda + \rho_0)}\right)}{v\left(\partial e^{\rho_0} \prod_{\alpha \in \Delta_0^+} (1 - e^{-\alpha})\right)} = \frac{\prod_{\alpha \in \Delta_0^+} (\lambda + \rho_0 | \alpha)}{\prod_{\alpha \in \Delta_0^+} (\rho_0 | \alpha)}.$$

Note that this is Weyl's character formula for irreducible finite-dimensional highest-weight modules. The second equality also holds if we replace $\lambda + \rho_0$ by $zw(\lambda + \rho_0)$ for any complex number z and any w in W^\sharp .

Definition. If $m \in \mathbb{N}$ and μ in $\frac{1}{m}A'_N$, then v_μ is the rational exponential function

$$(2.6) \quad v_\mu = \frac{\sum_{w \in W^\sharp} \sigma(w) e^{w(\mu)}}{e^{\rho_0} \prod_{\alpha \in \Delta_0^+} (1 - e^{-\alpha})}.$$

Note that if $\mu - \rho_0$ is dominant, then this is just the character of the irreducible highest-weight module of highest-weight $\mu - \rho_0$. We now closely follow the argument of the proof of the dimension formula of [23].

Lemma 2.3. If $m \in \mathbb{N}$ and μ in $\frac{1}{m}A'_N$, then

$$v(v_\mu) = \frac{1}{d_N} \prod_{\alpha \in \Delta_0^+} (\mu | \alpha) = \frac{v(\partial e^\mu)}{d_N}.$$

Proof: Using the explicit description of the positive even roots in (2.1), it is easy to compute

$$\prod_{\alpha \in \Delta_0^+} (\rho_0 | \alpha) = \prod_{n=1}^N n! = d_N.$$

For an arbitrary weight $\mu \in A'_N$, there exists a unique $w \in W^\sharp$ such that $w(\mu + \rho_0) - \rho_0$ is dominant. Letting $\ell(w)$ be the number of positive roots that are mapped to negative ones by w , then $(-1)^{\ell(w)} = \sigma(w)$ (see [23]). Then using that the even Weyl group respects the bilinear form.

$$\begin{aligned} v\left(v_{\frac{\mu}{M}}\right) &= v\left(\frac{\sum_{w \in W^\sharp} \sigma(w) e^{w(\frac{\mu}{M})}}{e^{\rho_0} \prod_{\alpha \in \Delta_0^+} (1 - e^{-\alpha})}\right) = (-1)^{\ell(w)} \frac{\prod_{\alpha \in \Delta_0^+} \left(w\left(\frac{\mu}{M}\right) | \alpha\right)}{\prod_{\alpha \in \Delta_0^+} (\rho_0 | \alpha)} = \frac{(-1)^{\ell(w)}}{d_N} \prod_{\alpha \in \Delta_0^+} \left(\frac{\mu}{M} | w^{-1}(\alpha)\right) \\ &= \frac{(-1)^{\ell(w) + \ell(w^{-1})}}{d_N} \prod_{\alpha \in \Delta_0^+} \left(\frac{\mu}{M} | \alpha\right) = \frac{1}{d_N} \prod_{\alpha \in \Delta_0^+} \left(\frac{\mu}{M} | \alpha\right) = \frac{v\left(\partial e^{\frac{\mu}{M}}\right)}{d_N}. \end{aligned}$$

□

3. BASIC FACTS ON JACOBI FORMS AND QUANTUM MODULAR FORMS

3.1. Jacobi Forms. Here we recall some special Jacobi forms and previous work on Fourier coefficients of Jacobi forms. Jacobi forms are functions from $\mathbb{C} \times \mathbb{H} \rightarrow \mathbb{C}$ which satisfy both an elliptic and a modular transformation law. For the precise definition and basic facts on Jacobi forms we refer the reader to [20]. In this paper, we are particularly interested in the classical Jacobi theta function, defined in (1.1). The following transformation laws are well-known (for example, see [35] (80.31) and (80.8)).

Lemma 3.1. *For $\lambda, \mu \in \mathbb{Z}$ and $\gamma = \begin{pmatrix} a & b \\ c & d \end{pmatrix} \in \mathrm{SL}_2(\mathbb{Z})$, we have that*

$$\vartheta \left(\frac{z}{c\tau + d}; \frac{a\tau + b}{c\tau + d} \right) = \psi^3(\gamma)(c\tau + d)^{\frac{1}{2}} e^{\frac{\pi icz^2}{c\tau + d}} \vartheta(z; \tau),$$

$$\vartheta(z + \lambda\tau + \mu; \tau) = (-1)^{\lambda + \mu} q^{-\frac{\lambda^2}{2}} e^{-2\pi i \lambda z} \vartheta(z; \tau),$$

where $\psi(\gamma)$ is the multiplier arising in the transformation law of Dedekind's eta function.

We also require the following theta functions of weight $\frac{1}{2} + \nu$ defined for $r \in \mathbb{Z}$, $\nu \in \{0, 1\}$

$$\vartheta_{\frac{1}{2} + \nu}(N, r; \tau) := \sum_{n \in \mathbb{Z}} (-1)^{nN} \left(n + \frac{r}{N} - \frac{1}{2} \right)^\nu q^{\frac{N}{2} \left(n + \frac{r}{N} - \frac{1}{2} \right)^2}.$$

We define for convenience the following shifted versions when $N \in 2\mathbb{N}$

$$(3.1) \quad \tilde{\vartheta}_{\frac{1}{2} + \nu}(N, r; \tau) := \vartheta_{\frac{1}{2} + \nu} \left(N, r + \frac{N}{2}; \tau \right).$$

It is trivial to show the following identities:

$$(3.2) \quad \tilde{\vartheta}_{\frac{1}{2} + \nu}(N, r + N; \tau) = \tilde{\vartheta}_{\frac{1}{2} + \nu}(N, r; \tau),$$

$$(3.3) \quad \tilde{\vartheta}_{\frac{1}{2} + \nu}(N, -r; \tau) = (-1)^\nu \tilde{\vartheta}_{\frac{1}{2} + \nu}(N, r; \tau),$$

$$(3.4) \quad \vartheta_{\frac{1}{2}}(1, 0; \tau) = \tilde{\vartheta}_{\frac{3}{2}}(2, r; \tau) = 0.$$

In Sections 5 and 6, we need the following modular transformations, which can be derived as special cases of the transformation formulas for the theta functions of Shimura [37].

Proposition 3.2. *If $N \in 2\mathbb{N}$ and $r \in \mathbb{Z}$, then we have:*

$$\tilde{\vartheta}_{\frac{3}{2}}(N, r; \tau + 1) = e \left(\frac{r^2}{2N} \right) \tilde{\vartheta}_{\frac{3}{2}}(N, r; \tau),$$

$$\tilde{\vartheta}_{\frac{3}{2}} \left(N, r; -\frac{1}{\tau} \right) = \frac{2}{\sqrt{N}} (-i\tau)^{\frac{3}{2}} \sum_{k=1}^{\frac{N}{2}-1} \sin \left(\frac{2\pi kr}{N} \right) \tilde{\vartheta}_{\frac{3}{2}}(N, k; \tau).$$

Moreover, for $\nu \in \{0, 1\}$, $\tilde{\vartheta}_{\frac{1}{2} + \nu}(N, r; \tau)$ is a modular form of weight $\frac{1}{2} + \nu$ on $\Gamma_1(2N)$ with multiplier

$$(3.5) \quad \chi_r \begin{pmatrix} a & b \\ c & d \end{pmatrix} := \begin{cases} e \left(\frac{r^2}{2N} \right) & \text{if } c = 0, \\ e \left(\frac{br^2}{2N} \right) \left(\frac{2Nc}{d} \right) & \text{if } c \neq 0. \end{cases}$$

We remark that in Proposition 3.2, $\vartheta_{\frac{1}{2}+\nu}(N, r; \tau)$ are actually modular forms on a slightly larger congruence subgroup, but we have chosen to use $\Gamma_1(2N)$ for ease of exposition.

We next recall the structure of Fourier coefficients of positive index Jacobi forms for comparison. It is well known that holomorphic Jacobi forms have a theta decomposition involving the functions

$$\vartheta_{m,b}(z; \tau) := \sum_{\substack{\lambda \in \mathbb{Z} \\ \lambda \equiv b \pmod{2m}}} e^{\frac{\pi i \lambda^2 \tau}{2m} + 2\pi i \lambda z}.$$

The components of this decomposition are classical (vector-valued) modular forms [20]. The Fourier coefficients of meromorphic Jacobi forms of positive index are also understood. Specifically, in [8], Folsom and the first author, building on illuminating work of Dabholkar, Murthy, and Zagier [17] and Zwegers [46], considered the Kac-Wakimoto character of level (M, N) with $M > N$, $M, N \in 2\mathbb{N}$, which essentially corresponds to the meromorphic Jacobi form $\phi_{M,N}$ (the general case with $M > N$ is considered in [34]). These Kac-Wakimoto characters have a decomposition into a finite and a polar part, where the finite part has a theta decomposition similar to that of holomorphic Jacobi forms (but involving mock modular forms), and where the polar part is

(3.6)

$$\varphi^P(z; \tau) := - \sum_{j=1}^{\frac{N}{2}} \frac{D_{2j}(\tau)}{(2\pi i)^{2j-1} (2j-1)!} \frac{\partial^{2j-1}}{\partial u^{2j-1}} \left(e^{\pi i (M-N)u} \zeta^{-\frac{M-N}{2}} F_{M-N} \left(\frac{\tau}{2} - u, z - \frac{\tau}{2}; \tau \right) \right) \Big|_{u=0}.$$

Here D_j is the j -th Laurent coefficient of the level (M, N) Kac-Wakimoto character. Thus we see that our functions $\phi_{M,N}$ have decompositions which are strikingly similar to the decompositions of positive index Jacobi forms, although in our case there are no associated “finite parts”. As mentioned in Remark 1 following Theorem 1.3, this has an interesting interpretation in physics.

3.2. Quantum modular forms. In this section, we recall some definitions and examples of quantum modular forms and describe the quantum sets in Theorem 1.5. We begin with a few definitions (see [28] for additional background on quasimodular forms).

Definition. A function $f: \mathbb{H} \rightarrow \mathbb{C}$ is an **almost holomorphic modular form of weight k** on a congruence subgroup Γ if it transforms as a modular form of weight k for Γ and is a polynomial in $\frac{1}{\text{Im}(\tau)}$ with coefficients which are holomorphic on $\mathbb{H} \cup \mathbb{P}_1(\mathbb{Q})$. Moreover, f is a **quasimodular form of weight k** if it is the constant term of an almost holomorphic modular form of weight k .

Quantum modular are then defined as follows (see [44] for background on quantum modular forms).

Definition. For any infinite “quantum set” $\mathcal{Q} \subseteq \mathbb{Q}$, we say a function $f: \mathcal{Q} \rightarrow \mathbb{C}$ is a **quantum modular form of weight k** on a congruence subgroup Γ if for all $\gamma \in \Gamma$, the cocycle

$$r_\gamma(\tau) := f|_k(1 - \gamma)(\tau)$$

extends to an open subset of \mathbb{R} and is analytically “nice”. Here “nice” could mean continuous, smooth, real-analytic etc. We say that f is a **strong quantum modular form** if there is a formal power series over \mathbb{C} attached to each point in \mathcal{Q} with a stronger modularity requirement (see [44]).

Remark. All of the quantum modular forms occurring in this paper have cocycles defined on \mathbb{R} which are real-analytic except at one point. Moreover, they have full asymptotic expansions towards rational points in their quantum sets which agree with the asymptotic expansions of mock modular forms defined on the lower half-plane.

Especially relevant for us are certain partial theta functions which were shown to be quantum modular forms in recent work of Folsom, Ono, and Rhoades [22]; namely

$$G(a, b; \tau) := \sum_{n \geq 0} (-1)^n q^{(n + \frac{a}{b})^2}.$$

For any $a, b \in \mathbb{Z}$ with $(a, b) = 1$, $a > 0$, define the following quantum set, where all fractions are assumed to have coprime denominator and numerator throughout

$$\mathcal{Q}_{a,b} := \left\{ \frac{h}{k} \in \mathbb{Q} : h > 0, b|2h, b \nmid h, k \equiv a \pmod{b}, k \geq a \right\}.$$

Since for $r = \frac{j}{2} \in \frac{1}{2} + \mathbb{Z}$

$$\Theta_{\frac{1}{2}}(N, r; \tau) = G\left(j, 2N; \frac{N\tau}{2}\right),$$

it suffices to study the quantum modular properties of $G(a, b; \tau)$. Although $a = 0$ is excluded, it is easy to handle this case directly. Note that $G(0, 1; \frac{\tau}{2})$ is essentially a modular form as $G(0, 1; \tau) = \frac{\eta(\tau)^2}{2\eta(2\tau)} + \frac{1}{2}$ and also that $G(0, 2N; \tau) = G(0, 1; 2N\tau)$. It is clear that $G(0, 1; \frac{\tau}{2})$ is quantum modular at any cusps where the eta quotient vanishes, namely for $\tau \in \{\frac{h}{k} \in \mathbb{Q} : k \equiv 1 \pmod{2}\}$. For $a > 0$, the situation is more subtle. Folsom, Ono and Rhoades proved that $G(a, b; \tau)$ have the following quantum properties:

Theorem 3.3 ([22]). *For b even, $G(a, b; \tau)$ is a strong quantum modular form of weight $1/2$ with quantum set $\mathcal{Q}_{a,b}$.*

Remark. *Although [22] only states the theorem for $0 < a < b$, an inspection of the proof shows that it is true for general integers $(a, b) = 1$ with $a > 0$ and b even*

When N is even, we also have the analogous weight $\frac{3}{2}$ partial theta functions $\Theta_{\frac{3}{2}}(N, r; \tau)$ (see Theorem 1.5).

4. THE FOURIER COEFFICIENTS AND PARTIAL THETA FUNCTIONS OF A_N

In this section, we prove Theorem 1.2.

Proof of Theorem 1.2: Define a subdomain of Y

$$X := \{h \in Y | \operatorname{Re}(C(h)) > \operatorname{Re}(\alpha(h)) > 0 \forall \alpha \in \Delta^+\}$$

so that in particular $|q(x)| < |e^{(\delta - \varepsilon_{N+1})(x)}| < 1$ for all x in X . We begin with the following crucial lemma.

Lemma 4.1. *As a function in X , we have*

$$e^{\hat{\rho}} \hat{R} = e^{Nd - \rho_1} q^{-\frac{(N+1)(N+2)}{24}} \sum_{r \in \mathbb{Z}} \operatorname{sign}(-r) (-1)^r \mathfrak{q}^{-\frac{r^2}{2(N+1)}} \mathfrak{q}^{\frac{r}{2}} e^r \sum_{t \in T_r} \sum_{w \in W^\#} \sigma(w) w \left(e^t \mathfrak{q}^{\frac{t^2}{2N}} \right).$$

Proof: Inserting the statement of Lemma 2.1 into the one of Lemma 2.2 gives

$$e^{\hat{\rho}} \hat{R} = \sum_{\alpha \in A_N} t_\alpha \left(e^{\hat{\rho}} R \right) = \sum_{\alpha \in A_N} t_\alpha \left(\sum_{w \in W^\#} \sigma(w) w \left(\frac{e^{\hat{\rho}}}{(1 + e^{\delta - \varepsilon_{N+1}})} \right) \right).$$

Using (2.4) and the bijectivity of the map $w : A_N \rightarrow A_N$ for every $w \in W^\#$, we get

$$e^{\hat{\rho}} \hat{R} = \sum_{\alpha \in A_N} \sum_{w \in W^\#} \sigma(w) w \left(t_{w^{-1}(\alpha)} \left(\frac{e^{\hat{\rho}}}{(1 + e^{\delta - \varepsilon_{N+1}})} \right) \right) = \sum_{\alpha \in A_N} \sum_{w \in W^\#} \sigma(w) w \left(t_\alpha \left(\frac{e^{\hat{\rho}}}{(1 + e^{\delta - \varepsilon_{N+1}})} \right) \right).$$

Let $\alpha = m_1\alpha_1 + \cdots + m_N\alpha_N$ be an element of A_N , and set $m_0 := m_{N+1} := 0$. By (2.3), we have

$$t_\alpha(\widehat{\rho}) = \widehat{\rho} + N\alpha - \left(\sum_{n=1}^{N+1} \left(m_n + \frac{N}{2}(m_n - m_{n-1})^2 \right) \right) C \quad \text{and} \quad t_\alpha(\delta - \varepsilon_{N+1}) = \delta - \varepsilon_{N+1} - m_N C.$$

Hence

$$e^{\widehat{\rho}} \widehat{R} = \sum_{m \in \mathbb{Z}^N} \sum_{w \in W^\#} \sigma(w) w \left(e^{\widehat{\rho}} \prod_{n=1}^{N+1} e^{N m_n \alpha_n} \mathbf{q}^{m_n + \frac{N}{2}(m_n - m_{n-1})^2} \left(1 + e^{\delta - \varepsilon_{N+1}} \mathbf{q}^{m_N} \right)^{-1} \right),$$

where we used the short-hand notation $m = (m_1, \dots, m_N)$ and kept as before $m_0 = m_{N+1} = 0$. Recall that $\alpha_n = \varepsilon_n - \varepsilon_{n+1}$. We split the exponential of the affine Weyl vector as

$$e^{\widehat{\rho}} = e^{Nd - \rho_1} \prod_{n=1}^{N+1} e^{\frac{(N+2-2n)}{2} \varepsilon_n}.$$

Note that $Nd - \rho_1$ is invariant under $W^\#$. Letting $q_n := m_n - m_{n-1}$, we then find the identity

$$\sum_{n=1}^{N+1} \left(\frac{N}{2} q_n^2 + m_n \right) = \frac{1}{2N} \sum_{n=1}^{N+1} \left(N \left(q_n + \frac{1}{2} \right) + (1-n) \right)^2 - \frac{(N+1)(N+2)}{24}.$$

Defining the set

$$S := \left\{ (s_1, \dots, s_{N+1}) \in \frac{1}{2} \mathbb{Z}^{N+1} \mid s_n = N \left(q_n + \frac{1}{2} \right) + (1-n), q_n \in \mathbb{Z}, \sum_{n=1}^{N+1} q_n = 0 \right\},$$

we obtain

$$e^{\widehat{\rho}} \widehat{R} = e^{Nd - \rho_1} \mathbf{q}^{-\frac{(N+1)(N+2)}{24}} \sum_{s \in S} \sum_{w \in W^\#} \sigma(w) w \left(\prod_{n=1}^{N+1} e^{s_n \varepsilon_n} \mathbf{q}^{\frac{s_n^2}{2N}} \left(1 + e^{\delta - \varepsilon_{N+1}} \mathbf{q}^{-\frac{s_{N+1}}{N} - \frac{1}{2}} \right)^{-1} \right).$$

In the domain X we can expand in a geometric series to find that $e^{\widehat{\rho}} \widehat{R}$ equals

$$e^{Nd - \rho_1} \mathbf{q}^{-\frac{(N+1)(N+2)}{24}} \left(\sum_{\substack{s \in S \\ s_{N+1} \leq -\frac{1}{2}}} \sum_{r=0}^{\infty} (-1)^r \sum_{w \in W^\#} \sigma(w) w \left(e^{r(\delta - \varepsilon_{N+1})} \mathbf{q}^{-r(\frac{s_{N+1}}{N} + \frac{1}{2})} \prod_{n=1}^{N+1} e^{s_n \varepsilon_n} \mathbf{q}^{\frac{s_n^2}{2N}} \right) \right. \\ \left. - \sum_{\substack{s \in S \\ s_{N+1} > -\frac{1}{2}}} \sum_{r=1}^{\infty} (-1)^r \sum_{w \in W^\#} \sigma(w) w \left(e^{-r(\delta - \varepsilon_{N+1})} \mathbf{q}^{r(\frac{s_{N+1}}{N} + \frac{1}{2})} \prod_{n=1}^{N+1} e^{s_n \varepsilon_n} \mathbf{q}^{\frac{s_n^2}{2N}} \right) \right).$$

Since the double sum converges absolutely in the domain X , we can interchange summations. Define

$$g_r := \begin{cases} \sum_{\substack{s \in S \\ s_{N+1} \leq -\frac{1}{2}}} \sum_{w \in W^\#} \sigma(w) w \left(e^{-r(\delta - \varepsilon_{N+1})} \mathbf{q}^{r(\frac{s_{N+1}}{N} + \frac{1}{2})} \prod_{n=1}^{N+1} e^{s_n \varepsilon_n} \mathbf{q}^{\frac{s_n^2}{2N}} \right) & \text{if } r \leq 0, \\ \sum_{\substack{s \in S \\ s_{N+1} > -\frac{1}{2}}} \sum_{w \in W^\#} \sigma(w) w \left(e^{-r(\delta - \varepsilon_{N+1})} \mathbf{q}^{r(\frac{s_{N+1}}{N} + \frac{1}{2})} \prod_{n=1}^{N+1} e^{s_n \varepsilon_n} \mathbf{q}^{\frac{s_n^2}{2N}} \right) & \text{if } r > 0. \end{cases}$$

Then

$$e^{\widehat{\rho}} \widehat{R} = e^{Nd - \rho_1} \mathbf{q}^{-\frac{(N+1)(N+2)}{24}} \sum_{r \in \mathbb{Z}} (-1)^r \text{sign}(-r) g_r.$$

We can express $\varepsilon_{N+1} - \delta$ in terms of the odd Weyl vector and positive even simple roots:

$$(4.1) \quad \varepsilon_{N+1} - \delta = -\delta + \frac{1}{N+1} \left(\sum_{n=1}^{N+1} \varepsilon_n + \sum_{n=1}^N n(\varepsilon_{n+1} - \varepsilon_n) \right) = \frac{2}{N+1} \rho_1 + \varepsilon_{N+1} - \frac{1}{N+1} \sum_{n=1}^{N+1} \varepsilon_n.$$

We see that $\varepsilon_{N+1} - \delta - \frac{2}{N+1} \rho_1$ is in $\frac{1}{N+1} A_N$. For $(s_1, \dots, s_{N+1}) \in S$, we find that

$$(4.2) \quad \left(s_{N+1} - \frac{r}{N+1} + r \right)^2 + \sum_{n=1}^N \left(s_n - \frac{r}{N+1} \right)^2 = \frac{Nr^2}{N+1} + 2rs_{N+1} + \sum_{n=1}^{N+1} s_n^2.$$

Combining (4.1) and (4.2), we can rewrite

$$e^{-r(\delta - \varepsilon_{N+1})} \mathbf{q}^{r\left(\frac{s_{N+1}}{N+1} + \frac{1}{2}\right)} \prod_{n=1}^{N+1} e^{s_n \varepsilon_n} \mathbf{q}^{\frac{s_n^2}{2N}} = e^{r\frac{2\rho_1}{N+1}} \mathbf{q}^{-\frac{r^2}{2(N+1)}} \mathbf{q}^{\frac{r}{2}} \prod_{n=1}^{N+1} e^{t_n \varepsilon_n} \mathbf{q}^{\frac{t_n^2}{2N}}$$

with $t_n := s_n - \frac{r}{N+1} + r\delta_{n,N+1}$. Then

$$\sum_{n=1}^{N+1} t_n \varepsilon_n = \sum_{n=1}^{N+1} \sum_{j=1}^n t_j (\varepsilon_j - \varepsilon_{j+1}) = \sum_{n=1}^N \sum_{j=1}^n t_j \alpha_n.$$

Here we used that $t_1 + \dots + t_{N+1} = 0$, which follows from the same property for the s_n . Let q_j be as in the definition of the set S ; in particular we can write $q_j = m_j - m_{j-1}$ with integers m_j for $1 \leq j \leq N$, and $m_{N+1} = 0$. Then

$$\sum_{j=1}^n t_j = -\frac{rn}{N+1} + \sum_{j=1}^n s_j = \frac{n(N-n+1)}{2} - \frac{rn}{N+1} + \sum_{j=1}^n Nq_j = \frac{(N-n+1)n}{2} - \frac{rn}{N+1} + Nm_n.$$

Using the sets T_r (1.5), we finally get

$$g_r = e^{r\frac{2\rho_1}{N+1}} \mathbf{q}^{-\frac{r^2}{2(N+1)}} \mathbf{q}^{\frac{r}{2}} \sum_{t \in T_r} \sum_{w \in W^\sharp} \sigma(w) w \left(e^t \mathbf{q}^{\frac{t^2}{2N}} \right).$$

□

Letting, $\mathfrak{z} = -e^{\frac{2\rho_1}{N+1}}$, we deduce the following.

Corollary 4.2. *The identity $A = BC$ holds as functions on X , where*

$$A := \prod_{j=1}^{N+1} \frac{\mathbf{q}^{\frac{1}{24}}}{(-e^{\delta - \varepsilon_j}; \mathbf{q})_\infty (-e^{\varepsilon_j - \delta}; \mathbf{q}; \mathbf{q})_\infty}, \quad B := \frac{\mathbf{q}^{-\frac{(N+1)^2}{24}}}{(\mathbf{q}; \mathbf{q})_\infty^{N+1} \prod_{\alpha \in \Delta_0} (e^\alpha \mathbf{q}; \mathbf{q})_\infty (e^{-\alpha} \mathbf{q}; \mathbf{q})_\infty},$$

$$C := \sum_{r \in \mathbb{Z}} \text{sign}(-r) \mathfrak{z}^r \mathbf{q}^{-\frac{r^2}{2(N+1)}} \mathbf{q}^{\frac{r}{2}} \sum_{t \in T_r} v_t \mathbf{q}^{\frac{t^2}{2N}}.$$

Proof: The corollary follows immediately from Lemma 4.1 by inserting v_t in the definition of C in (2.6). □

Evaluating the expressions in this equality provides a nice expansion of $\phi_{N+1}(z; \tau)$.

Corollary 4.3. *Inside the range $|q| < |\zeta| < 1$, we have*

$$\phi_N(z; \tau) = \frac{i^N}{d_{N-1}\eta(\tau)^{N^2+N}} \sum_{r \in \mathbb{Z}} \zeta^{r+\frac{N}{2}} \text{sign}(r) q^{-\frac{(r+\frac{N}{2})^2}{2N}} \sum_{t \in T_r} \prod_{\alpha \in \Delta_0^+} (t|\alpha) q^{\frac{t^2}{2(N-1)}}.$$

Proof: The evaluation v maps every regular exponential e^λ for $\lambda \in \frac{1}{N}A_{N-1}$ to 1. The application of v to A and B is finite for $|q(x)| < |\mathfrak{z}^{-1}(x)| < 1$ and $x \in X$, and the same is true for C by Lemma 2.3. The identity (4.1) implies that $v(e^{\delta-\varepsilon_j}) = e^{-\frac{2\rho_1}{N}} = -\zeta$ for all $j = 1, \dots, N$, so that

$$v(A) = \mathfrak{q}^{\frac{N}{24}} (\mathfrak{z}^{-1}; \mathfrak{q})_\infty^{-N} (\mathfrak{z}\mathfrak{q}; \mathfrak{q})_\infty^{-N} \quad \text{and} \quad v(B) = \mathfrak{q}^{-\frac{N^2}{24}} (\mathfrak{q}; \mathfrak{q})_\infty^{-N^2}.$$

By Lemma 2.3,

$$v(v_t) = \frac{1}{d_{N-1}} \prod_{\alpha \in \Delta_0^+} (t|\alpha)$$

and the evaluation $v(C)$ follows. All three evaluations $v(A), v(B), v(C)$ are meromorphic functions on $\left\{x = -2\pi i\tau d + \frac{4\pi i z h_{\rho_1}}{N-1} : \text{Im}(\tau) > \text{Im}(z) > 0\right\}$, so that the result follows with $\zeta = \mathfrak{z}^{-1}(x)$ and $q = \mathfrak{q}(x)$. \square

This completes the proof as Corollary 4.3 and Lemma 2.3 imply Theorem 1.2. \square

The case $N = 1$ can be proven in a very similar manner using (1.9), which is the denominator identity of $\widehat{\mathfrak{gl}}(1|1)$ (see Example 4.1 of [27]).

Example 4.4. *The Fourier coefficients of $\phi_1(z; \tau)$ are given by*

$$\chi(1, r; \tau) = \frac{i q^{-\frac{r^2}{2}}}{\eta(\tau)^3} \sum_{m=0}^{\infty} (-1)^m q^{\frac{(m+|r-\frac{1}{2}|+\frac{1}{2})^2}{2}}.$$

Proof: Suppose $|q| < |\zeta| < 1$. Expanding (1.9) in a geometric series and rewriting easily gives the statement. \square

5. SECOND VIEWPOINT ON THE DECOMPOSITION INTO PARTIAL THETA FUNCTIONS

In this section, we prove Theorem 1.3 and use it to extract the Fourier coefficients of $\phi_{M,N}$ in Theorem 1.4. A key ingredient for the proof of Theorem 1.3 is the following result whose proof is deferred to Section 5.2.

Lemma 5.1. *For $N \in \mathbb{N}$, there exist meromorphic functions $f_j^*(\tau)$ for $0 \leq j \leq \frac{N-1-\delta_e}{2}$ with $f_{\frac{N-1-\delta_e}{2}}^*(\tau) \neq 0$ such that for all $r \in \mathbb{Z}$*

$$\sum_{j=0}^{\frac{N-1-\delta_e}{2}} f_j^*(\tau) \mathcal{D}_q^j \left(\vartheta_{\frac{1}{2}+\delta_e}(N, r; \tau) \right) = 0.$$

5.1. Proof of Theorem 1.3 for $M = 0$. The first step in the proof of Theorem 1.3 is to show the following decomposition for the case when $M = 0$:

Proposition 5.2. *For $N \in \mathbb{N}$ there exist meromorphic functions $g_j(\tau)$ such that*

$$\phi_N(z; \tau) = \sum_{j=0}^{\frac{N-1-\delta_e}{2}} g_j(\tau) \mathcal{D}_w^{2j+\delta_e} (F_N(z, u; \tau)) \Big|_{w=1}.$$

Proof: We first determine the elliptic transformations of $F_N(z; \tau)$ and prove that, although this function does not in general transform as a negative index Jacobi form, we can “correct” the elliptic transformations to match those of $\phi_N(z; \tau)$. The following periodicity property is evident:

$$F_N(z + 1, u; \tau) = (-1)^N F_N(z, u; \tau).$$

For the elliptic transformation $z \mapsto z + \tau$, a direct calculation gives

$$\begin{aligned} & (-1)^N \zeta^{-N} q^{-\frac{N}{2}} F_N(z + \tau, u; \tau) - F_N(z, u; \tau) \\ &= \sum_{r=0}^{N-1} \zeta^{r-\frac{N}{2}} q^{-\frac{1}{2N}(r-\frac{N}{2})^2} \sum_{n \in \mathbb{Z}} (-1)^{Nn} w^{Nn+r-\frac{N}{2}} q^{\frac{N}{2}(n-\frac{1}{2}+\frac{r}{N})^2}. \end{aligned}$$

Thus, we have the following elliptic transformation formula for the iterated derivative of $F_N(z; \tau)$:

$$\begin{aligned} & (-1)^N \zeta^{-N} q^{-\frac{N}{2}} \mathcal{D}_w^{2j+\delta_e}(F_N(z + \tau, u; \tau))|_{w=1} - \mathcal{D}_w^{2j+\delta_e}(F_N(z, u; \tau))|_{w=1} \\ &= 2^j N^{j+\delta_e} \sum_{r=0}^{N-1} \zeta^{r-\frac{N}{2}} q^{-\frac{1}{2N}(r-\frac{N}{2})^2} \mathcal{D}_q^j \left(\vartheta_{\frac{1}{2}+\delta_e}(N, r; \tau) \right). \end{aligned}$$

We now use the functions f_j^* from Lemma 5.1 to “correct” the elliptic transformation by defining

$$P_N(z; \tau) := \sum_{j=0}^{\frac{N-1-\delta_e}{2}} \frac{f_j^*(\tau)}{2^j N^{j+\delta_e}} \mathcal{D}_w^{2j+\delta_e}(F_N(z, u; \tau))|_{w=1},$$

so that

$$P_N(z + 1; \tau) = \zeta^{-N} q^{-\frac{N}{2}} P_N(z + \tau; \tau) = (-1)^N P_N(z; \tau).$$

Thus, $P_N(z; \tau)$ satisfies the same elliptic transformations as $\phi_N(z; \tau)$. It also has poles in the same locations and of the same order; namely poles in $\mathbb{Z}\tau + \mathbb{Z}$ of order N . Hence the product

$$p_N(z; \tau) := \vartheta(z; \tau)^N P_N(z; \tau)$$

is an entire elliptic function and therefore constant in z . It remains to show that $P_N(z; \tau) \neq 0$, which we prove by looking at the behavior as $z \rightarrow 0$. The principal part as $z \rightarrow 0$ of $\mathcal{D}_w^j(F_N(z, u; \tau))$ only comes from the $n = 0$ term in (1.11), which contributes

$$\frac{(\zeta w)^{\frac{N}{2}}}{1 - \zeta w} = - \sum_{m \geq 0} \frac{B_m\left(\frac{N}{2}\right) (2\pi i(u+z))^{m-1}}{m!} = -\frac{1}{2\pi i(u+z)} + O(1),$$

where $B_m(x)$ is the usual m -th Bernoulli polynomial. Thus, as $z \rightarrow 0$,

$$(5.1) \quad \mathcal{D}_w^j(F_N(z, u; \tau))|_{w=1} = \frac{(-1)^{j+1} j!}{(2\pi i z)^{j+1}} + O(1)$$

and so

$$P_N(z; \tau) = \frac{(-1)^N (N-1)! f_{\frac{N-1-\delta_e}{2}}^*(\tau)}{N^{\delta_e} (2N)^{\frac{N-1-\delta_e}{2}} (2\pi i z)^N} + O(z^{-N+1}).$$

We can then use the well-known formula

$$\vartheta'(0; \tau) = -2\pi\eta(\tau)^3,$$

and compare the coefficients of z^{-N} to give

$$p_N(z; \tau) = \frac{(N-1)!(-i)^N f_{\frac{N-1-\delta_e}{2}}^*(\tau)}{N^{\delta_e}(2N)^{\frac{N-1-\delta_e}{2}}} \neq 0,$$

as by assumption $f_{\frac{N-1-\delta_e}{2}}^* \neq 0$. By absorbing the constants into the f_j^* , Proposition 5.2 follows. \square

Proof of Theorem 1.3 for $M = 0$: To finish the proof for $M = 0$, we connect the functions g_j in Proposition 5.2 to the Laurent coefficients of ϕ_N given in (1.12) by comparing principal parts. Namely, using (5.1), we easily read off:

$$g_j(\tau) = \frac{(-1)^{\delta_e+1} D_{2j+\delta_e+1}(\tau)}{(2j+\delta_e)!}.$$

\square

5.2. Proof of Lemma 5.1. For N odd, Lemma 2.1 of [47] easily gives Lemma 5.1 by rearranging terms. The condition $f_0 \neq 0$ (in the notation of [47]) is not stated explicitly in the statement; however the proof shows that one can choose $f_0 = 1$ in Lemma 2.1 of [47]. Now suppose that N is even. For $k \in \mathbb{N}$, consider the Ramanujan-Serre derivative, which raises the weight of a modular form by 2:

$$\mathcal{E}_k := \mathcal{D}_q - \frac{k}{12} E_2(\tau),$$

and its iterated version starting at weight $\frac{3}{2}$ given by $\mathcal{E}^n := \mathcal{E}_{2n-\frac{1}{2}} \circ \mathcal{E}_{2n-\frac{5}{2}} \circ \dots \circ \mathcal{E}_{\frac{7}{2}} \circ \mathcal{E}_{\frac{3}{2}}$. By rearranging, it is enough to show that there are holomorphic functions f_j such that for all $r \in \mathbb{Z}$

$$\sum_{j=0}^{\frac{N}{2}-1} f_j(\tau) \mathcal{E}^j \left(\vartheta_{\frac{3}{2}}(N, r; \tau) \right) = 0.$$

This is clearly equivalent to the following, where $\tilde{\vartheta}_{\frac{1}{2}+\nu}(N, r; \tau)$ is defined in (3.1).

Lemma 5.3. *If $N \in 2\mathbb{N}$, then there exist meromorphic functions $f_j(\tau)$ with $f_{\frac{N}{2}-1}(\tau) \neq 0$ such that for all $r \in \mathbb{Z}$*

$$(5.2) \quad \sum_{j=0}^{\frac{N}{2}-1} f_j(\tau) \mathcal{E}^j \left(\tilde{\vartheta}_{\frac{3}{2}}(N, r; \tau) \right) = 0.$$

Proof: The approach taken here is similar to Zwegers' proof of Lemma 2.1 in [47], although we give details for the reader's convenience. Using (3.2) and (3.3), it suffices to prove the lemma for $1 \leq r \leq \frac{N}{2} - 1$. By (3.4), we may simply choose $f_0 = 1$ for $N = 2$. Thus, we assume for the remainder of the proof that $N \geq 4$.

Consider the vector-valued form

$$\vec{\vartheta}_N(\tau) := \left(\tilde{\vartheta}_{\frac{3}{2}}(N, 1; \tau), \dots, \tilde{\vartheta}_{\frac{3}{2}}\left(N, \frac{N}{2} - 1; \tau\right) \right)^T$$

and the matrix-valued form

$$T_N := \left(\vec{\vartheta}_N, \mathcal{E} \left(\vec{\vartheta}_N \right), \dots, \mathcal{E}^{\frac{N}{2}-2} \left(\vec{\vartheta}_N \right) \right).$$

Using Proposition 3.2, we see that

$$T_N(\tau + 1) = \text{diag} \left(e \left(\frac{j^2}{2N} \right) \right)_{1 \leq j \leq \frac{N}{2}-1} T_N(\tau),$$

$$T_N \left(-\frac{1}{\tau} \right) = \frac{2}{\sqrt{N}} (-i\tau)^{\frac{3}{2}} \left(\sin \left(\frac{2\pi k\ell}{N} \right) \right)_{1 \leq k, \ell \leq \frac{N}{2}-1} T_N(\tau) \text{diag} \left(\tau^{2j-2} \right)_{1 \leq j \leq \frac{N}{2}-1}.$$

Hence

$$\det(T_N)(\tau + 1) = e \left(\frac{(N-1)(N-2)}{48} \right) \det(T_N)(\tau),$$

$$\det(T_N) \left(-\frac{1}{\tau} \right) = (-i\tau)^{\frac{(N-1)(N-2)}{4}} \det(T_N)(\tau).$$

Here we used the following elementary determinant formula:

$$\det \left(\frac{2(-i)^{\frac{3}{2}}}{\sqrt{N}} \sin \left(\frac{2\pi \ell k}{N} \right) \right)_{1 \leq k, \ell \leq \frac{N}{2}-1} = (-i)^{\frac{(N-1)(N-2)}{4}}.$$

From these transformations we see that $\det(T_N)$ is a modular form on $\text{SL}_2(\mathbb{Z})$ with the same multiplier system as $\eta^{\frac{(N-1)(N-2)}{2}}$. It is then easy to see that the ratio

$$\frac{\det(T_N)(\tau)}{\eta(\tau)^{\frac{(N-1)(N-2)}{2}}} =: \alpha$$

is a constant. Assuming that α is non-zero, it follows that $T_N(\tau)$ is invertible for all $\tau \in \mathbb{H}$.

To see that α is non-zero, observe that by elementary row operations we can write

$$\det(T_N) = \det \left(\vec{\vartheta}_N, \mathcal{D}_q \left(\vec{\vartheta}_N \right), \dots, \mathcal{D}_q^{\frac{N}{2}-2} \left(\vec{\vartheta}_N \right) \right).$$

For $1 \leq r \leq \frac{N}{2} - 1$, note that $\tilde{\vartheta}_{\frac{3}{2}}(N, j; \tau) = \frac{j}{N} q^{\frac{j^2}{2N}} (1 + O(q))$. But then we have

$$(5.3) \quad \det(T_N)(\tau) = \det \left(\text{diag} \left(\frac{j}{N} q^{\frac{j^2}{2N}} \right)_{1 \leq j \leq \frac{N}{2}-1} \left(\left(\frac{m^2}{2N} \right)^{k-1} + O(q) \right)_{1 \leq m, k \leq \frac{N}{2}-1} \right).$$

The right hand side is clearly non-zero as the first matrix on the right hand side of (5.3) is diagonal with non-zero entries and the second is a Vandermonde matrix with distinct generators. We can now solve (5.2) by writing it as

$$T_N \left(f_0, \dots, f_{\frac{N}{2}-2} \right)^T + f_{\frac{N}{2}-1} \mathcal{E}^{\frac{N}{2}-1} \left(\vec{\vartheta}_N \right) = 0.$$

Choosing $f_{\frac{N}{2}-1} = 1$, it is easy to see that we obtain holomorphic functions f_j satisfying (5.2) as follows:

$$\left(f_0, \dots, f_{\frac{N}{2}-2} \right)^T = -T_N^{-1} f_{\frac{N}{2}-1} \mathcal{E}^{\frac{N}{2}-1} \left(\vec{\vartheta}_N \right).$$

□

5.3. Proof of Theorem 1.3. We now use the above decomposition of $\phi_N(z; \tau)$ to deduce Theorem 1.3 for any M satisfying the conditions of the theorem. We may define a *heat operator*, $\mathcal{H} := \mathcal{H}_{-\frac{N}{2}} := 2N\mathcal{D}_q + \mathcal{D}_\zeta^2$ which has the property that it preserves the elliptic transformations of a function satisfying the elliptic transformation properties of an index $-\frac{N}{2}$ Jacobi form (note that this differs by a constant from the heat operator defined in [20]). The main idea is to show the following:

Proposition 5.4. *For any $N \in \mathbb{N}$ and $M \in \mathbb{N}_0$, there exist meromorphic functions $f_j(\tau)$ for $j = 0, 1, \dots, M$ with $f_M(\tau) \neq 0$ such that*

$$(5.4) \quad \phi_{2M, N+2M}(z; \tau) = \sum_{j=0}^M f_j(\tau) \mathcal{H}^j(\phi_N(z; \tau)).$$

For the proof we need a few preliminaries on special determinants and orthogonal polynomials. We recall that a matrix H is called a *Hankel matrix* if it is constant on each skew diagonal. It is well known that the determinants of Hankel matrices are connected to orthogonal polynomials and continued fractions. Such functions have a long history which is explained in many places; we refer the reader to Chapter 11 of [40] for more details. Given a sequence of numbers c_0, c_1, c_2, \dots which are the moments of a sequence of orthogonal polynomials $p_n(x)$, we say that $p_n(x)$ is a sequence of orthogonal polynomials *relative to c_n* . For any sequence $(c_n)_{n \in \mathbb{N}_0}$, we define the following sequence of Hankel determinants:

$$\Delta_n := \begin{vmatrix} c_0 & c_1 & \cdots & c_n \\ c_1 & c_2 & \cdots & c_{n+1} \\ \vdots & \vdots & \vdots & \vdots \\ c_n & c_{n+1} & \cdots & c_{2n} \end{vmatrix}.$$

Then we have the following:

Theorem 5.5 ([40] Theorem 50.1). *Let c_n be a sequence of numbers. Then $\Delta_n \neq 0$ for all n if and only if there exists a sequence of orthogonal polynomials relative to c_n .*

We now proceed to the proof of Proposition 5.4.

Proof of Proposition 5.4: First note that ϕ_N and $\phi_{2M, N+2M}$ have the same elliptic transformation properties. Moreover, the right hand side of (5.4) has poles of order exactly $N + 2M$ for $z \in \mathbb{Z} + \mathbb{Z}\tau$, as does $\phi_{2M, N+2M}$. It suffices to choose f_j such that the right hand side has zeros of order at least $2M$ for $z \in \frac{1}{2} + \mathbb{Z} + \mathbb{Z}\tau$, as dividing the right hand side of (5.4) by the left hand side gives a holomorphic elliptic function, and hence a constant.

Thus, it suffices to choose f_j to cancel out the first $2M$ Taylor coefficients of (5.4) at $z = -\frac{1}{2}$. Note, that as $\vartheta(z + \frac{1}{2})$ is even, we only have even order Taylor coefficients to cancel out. We expand

$$\mathcal{H}^j(\phi_N(z; \tau)) =: \sum_{k \geq 0} T_{2k, j}(\tau) \left(z + \frac{1}{2}\right)^{2k}.$$

Thus we need to find a solution of the equation

$$T \begin{pmatrix} f_0 \\ \vdots \\ f_{M-1} \end{pmatrix} + f_M \begin{pmatrix} T_{0, M} \\ \vdots \\ T_{2M, M} \end{pmatrix} = 0,$$

where

$$T := \begin{pmatrix} T_{0,0} & \cdots & T_{0, M-1} \\ \vdots & \vdots & \vdots \\ T_{2M-2,0} & \cdots & T_{2M-2, M-1} \end{pmatrix}.$$

Choosing $f_M(\tau) = 1$ it suffices to prove that $\det(T)(\tau) \neq 0$. To show this, it is enough to show that the q -series for $\det(T)$ has at least one non-vanishing coefficient; we look at the lowest order term. By the definition of $\vartheta(z; \tau)$, we have that

$$\vartheta(z; \tau) = q^{\frac{1}{8}} \left(e^{\pi i(z + \frac{1}{2})} + e^{-\pi i(z + \frac{1}{2})} \right) + O\left(q^{\frac{9}{8}}\right).$$

By rescaling, it suffices to study the matrix of Taylor coefficients of $\Psi(v; \tau) := q^{-\frac{N}{8}} \sec^N(v)$ and its iterated heat operators. By row reduction, we can replace the action of the operator \mathcal{H}^j with $\frac{\partial^{2j}}{\partial v^{2j}}$ and study the resulting determinant. By rescaling the k -th row of the resulting matrix by multiplying by $(2k)!$, we need to show nonvanishing of

$$S := \left| \left(\frac{\partial^{2(j+k)}}{\partial v^{2(j+k)}} (\sec^N(v)) \Big|_{v=0} \right)_{0 \leq j, k \leq M-1} \right|.$$

These numbers are known as the (signless) higher order Euler numbers, and we define

$$E_{2j}^{(N)} := \frac{\partial^{2j}}{\partial v^{2j}} (\sec^N(v)) \Big|_{v=0}.$$

By Theorem 5.5, it suffices to show that the higher order Euler numbers are for each N a moment sequence for a sequence of orthogonal polynomials. But Lemma 1.3 of [33] gives the desired sequence of orthogonal polynomials for any N . □

Finally, we compare Laurent expansions to determine a more convenient decomposition. Using the following geometric series expansion:

$$(5.5) \quad F_N(z, u; \tau) = \sum_{\substack{n \geq 0 \\ r \geq \frac{N}{2}}} (-1)^{Nn} q^{\frac{N}{2}n^2 + nr} \zeta^r w^{r+Nn} - \sum_{\substack{n \geq 1 \\ r < \frac{N}{2}}} (-1)^{Nn} q^{\frac{N}{2}n^2 - nr} \zeta^r w^{r-Nn},$$

where r runs through $\frac{N}{2} + \mathbb{Z}$, we directly find that

$$\mathcal{H}_{-\frac{N}{2}}^j (F_N(z, u; \tau)) = \mathcal{D}_w^{2j} (F_N(z, u; \tau)).$$

Since we have already shown Theorem 1.3 for $M = 0$ in Section 5.1, we see that for any $N \in \mathbb{N}$, $M \in \mathbb{N}_0$ that there exist meromorphic functions $g_j(\tau)$ such that

$$\phi_{2M, N+2M}(z; \tau) = \sum_{j=0}^{\frac{N-1-\delta_e}{2} + M} g_j(\tau) \mathcal{D}_w^{2j+\delta_e} (F_N(z, u; \tau)) \Big|_{w=1}.$$

Using (5.1) gives Theorem 1.3.

5.4. Extracting the Fourier coefficients of $\phi_{2M, N+2M}$. In this section, we use Theorem 1.3 to deduce Theorem 1.4 by expanding the Appell-Lerch series. Recall that throughout, we fix the range $0 \leq \text{Im}(z) < \text{Im}(\tau)$. We also assume that $0 < \text{Im}(u) < \text{Im}(\tau) - \text{Im}(z)$.

Proof: We use the geometric series expansion given in (5.5) to pick off the ζ^r coefficient of $F_N(z, u; \tau)$ as

$$[\zeta^r] (F_N(z, u; \tau)) = \begin{cases} q^{-\frac{r^2}{2N}} \sum_{n \geq 0} (-1)^{Nn} q^{\frac{N}{2}(n+\frac{r}{N})^2} w^{N(n+\frac{r}{N})} & \text{if } r \geq \frac{N}{2}, \\ -q^{-\frac{r^2}{2N}} \sum_{n \geq 1} (-1)^{Nn} w^{-N(n-\frac{r}{N})} q^{\frac{N}{2}(n-\frac{r}{N})^2} & \text{if } r < \frac{N}{2}. \end{cases}$$

Differentiating and applying Theorem 1.3 completes the proof. □

6. PROOF OF THEOREM 1.5

In this section, we complete the proof of Theorem 1.5 by showing strong quantum modularity of $\Theta_{\frac{3}{2}}(N, r; \tau)$ for any $N \in 2\mathbb{N}$, $r \in \mathbb{Z}$. We find it convenient to take sums and differences of two Fourier coefficients (which makes describing the asymptotic expansions in Proposition 6.5 much cleaner). Specifically, we define

$$\Theta_{\frac{3}{2}}^{\pm}(N, r; \tau) := \Theta_{\frac{3}{2}}(N, r; \tau) \pm \Theta_{\frac{3}{2}}(N, -r; \tau).$$

One easily sees that

$$(6.1) \quad \Theta_{\frac{3}{2}}^{-}(N, r; \tau) = \tilde{\vartheta}_{\frac{3}{2}}(N, r; \tau) + \frac{r}{N} q^{\frac{r^2}{2N}},$$

and for $0 \leq r \leq N - 1$ we compute:

$$(6.2) \quad \Theta_{\frac{3}{2}}^{+}(N, r; \tau) = \frac{1}{N} \sum_{\substack{n \geq -r \\ n \equiv \pm r \pmod{N}}} n q^{\frac{n^2}{2N}},$$

where for $r = 0$, we consider each summand with multiplicity 2.

6.1. Quantum sets. In this section, we show that the following is a quantum set for $\Theta_{\frac{3}{2}}(N, r; \tau)$:

$$(6.3) \quad \widehat{\mathcal{Q}}_{N,r} := \begin{cases} \left\{ \frac{h}{k} \in \mathbb{Q} : \frac{N}{2} | k, \text{ ord}_2(k) = \text{ord}_2(N) - 1 \right\} & \text{if } \frac{N}{2} \nmid r, \\ \left\{ \frac{h}{k} \in \mathbb{Q} : \text{ord}_2(k) > \text{ord}_2(N) \right\} & \text{if } r \equiv \frac{N}{2} \pmod{N}, \\ \left\{ \frac{h}{k} \in \mathbb{Q} : \text{ord}_2(k) = \text{ord}_2(N) \right\} & \text{if } r \equiv 0 \pmod{N}. \end{cases}$$

Here $\text{ord}_2(n)$ is the 2-adic valuation of an integer n . We will later see that the congruence subgroup for $\Theta_{\frac{3}{2}}(N, r; \tau)$ is $\Gamma_1(2N)$, which acts on $\widehat{\mathcal{Q}}_{N,r}$. Note that it is enough to determine a quantum set for each $0 \leq r \leq N - 1$, as if we shift $r \mapsto r \pm N$ in $\Theta_{\frac{3}{2}}(N, r; \tau)$, we only change the function by a finite polynomial in rational powers of q .

Lemma 6.1. *For any $N \in 2\mathbb{N}$, $0 \leq r \leq N - 1$, $\Theta_{\frac{3}{2}}(N, r; \tau)$ is well-defined on $\widehat{\mathcal{Q}}_{N,r}$.*

Proof: Throughout, we abuse terminology to say that a partial theta function is *convergent* at a root of unity if it has a q -hypergeometric representation which is.

We start with $\frac{N}{2} \nmid r$. Recall the following beautiful identity of Warnaar (p. 17 of [39])

$$(6.4) \quad \sum_{n \geq 0} \frac{(q; q^2)_n (Aq; q^2)_n (Aq)^n}{(-Aq)_{2n+1}} = \sum_{n \geq 0} (-A)^n q^{n(n+1)}.$$

Multiplying both sides of (6.4) by $A^{\frac{r}{N}}$, applying \mathcal{D}_A , and then setting $A = -q^{\frac{2r}{N}-1}$, we see that

$$\sum_{n \geq 0} \left(n + \frac{r}{N} \right) q^{\left(n + \frac{r}{N} \right)^2}$$

is convergent at any root of unity q such that $(q; q^2)_{\infty} (-q^{\frac{2r}{N}}; q^2)_{\infty} = 0$ and $(q^{\frac{2r}{N}})_{\infty} \neq 0$. Making the change of variables $q \rightarrow q^{\frac{N}{2}}$, shows that $\widehat{\mathcal{Q}}_{N,r}$ is a quantum set for $\Theta_{\frac{3}{2}}(N, r; \tau)$. We note that other hypergeometric representations in the literature could also be used to derive the same quantum set; see for example (1.2) discussed by Alladi in [2].

If $r = \frac{N}{2}$, note that $\Theta_{\frac{3}{2}}(N, \frac{N}{2}; \tau) = \Theta_{\frac{3}{2}}(1, \frac{1}{2}; N\tau)$, so we first study $\Theta_{\frac{3}{2}}(2, 1; \tau)$. This function can be split into two pieces, where the first is

$$\sum_{n \geq 0} q^{\frac{(n+\frac{1}{2})^2}{2}} = \frac{\eta(2\tau)^2}{\eta(\tau)}.$$

The eta-quotient on the right hand side is easily seen to vanish at any cusp $\frac{h}{k}$ with $2|k$. The second piece is related to a “sum of tails” by (3.11) of [5], which states that

$$(6.5) \quad \sum_{n \geq 1} nq^{\frac{n^2+n}{2}} = \frac{(q^2; q^2)_{\infty}}{(q; q^2)_{\infty}} \sum_{n \geq 1} \frac{(-1)^n q^n}{1 - q^n} + \sum_{n \geq 0} \left(\frac{(q^2; q^2)_{\infty}}{(q; q^2)_{\infty}} - \frac{(q^2; q^2)_n}{(q; q^2)_{n+1}} \right).$$

The right hand side of (6.5) terminates for even order roots of unity. By combining, the series $\Theta_{\frac{3}{2}}(2, 1; \tau)$ converges at even order roots of unity, and rescaling $q \mapsto q^N$ gives the quantum set in the definition of $\widehat{\mathcal{Q}}_{N,r}$.

If $r = 0$, we use (3.8) in [5], which states that

$$(6.6) \quad 4 \sum_{n \geq 1} (-1)^n nq^{n^2} = -2 \frac{(q; q)_{\infty}}{(-q; q)_{\infty}} \sum_{n \geq 1} \frac{q^n}{1 - q^{2n}} + \sum_{n \geq 0} \left(\frac{(q; q)_{\infty}}{(-q; q)_{\infty}} - \frac{(q; q)_n}{(-q; q)_n} \right).$$

The right hand side of (6.6) terminates for q an odd order root of unity. Letting $\tau \mapsto \tau + \frac{1}{2}$ shows that $\sum_{n \geq 0} nq^{n^2}$ converges for $\tau \in \{\frac{h}{k} \in \mathbb{Q} : k \equiv 2 \pmod{4}\}$. Finally, letting $\tau \mapsto \frac{N}{2}\tau$ shows that $\Theta_{\frac{3}{2}}(N, 0; \tau)$ converges at the claimed rational points. \square

6.2. Quantum modularity of $\Theta_{\frac{3}{2}}(N, r; \tau)$. In this subsection we use asymptotic expansions to study the quantum modularity of $\Theta_{\frac{3}{2}}^+(N, r; \tau)$. For $\tau \in \mathbb{H}_- := \{\tau \in \mathbb{C} : \text{Im}(\tau) < 0\}$, consider the nonholomorphic Eichler integral

$$(6.7) \quad \Theta_{\frac{1}{2}}^*(N, r; \tau) := -\frac{1}{2\pi\sqrt{iN}} \int_{\bar{\tau}}^{i\infty} \frac{\tilde{\vartheta}_{\frac{1}{2}}(N, r; z)}{(z - \tau)^{\frac{3}{2}}} dz.$$

We prove that $\Theta_{\frac{1}{2}}^*(N, r; \tau)$ agrees (up to one elementary term) with $\Theta_{\frac{3}{2}}^+(N, r; \tau)$ for $\tau = \frac{h}{k} \in \widehat{\mathcal{Q}}_{N,r}$ to infinite order and that $\Theta_{\frac{1}{2}}^*(N, r; \tau)$ satisfies a nice transformation law for $\tau \in \mathbb{H}_-$, which demonstrates that $\Theta_{\frac{3}{2}}(N, r; \tau)$ is a strong quantum modular form. Moreover, the proof clearly shows that the asymptotic expansions attached to rational numbers in the quantum set also arise from a mock modular form on the lower half-plane.

In order to give the asymptotic expansions of these functions, we require the following periodic sequence defined for any fixed $\frac{h}{k} \in \mathbb{Q}$:

$$\gamma_{N,r}(n) := \begin{cases} e\left(\frac{hn^2}{2kN}\right) & \text{if } n \equiv \pm r \pmod{N} \text{ and } r \neq 0, \\ 2e\left(\frac{hn^2}{2kN}\right) & \text{if } n \equiv 0 \pmod{N} \text{ and } r = 0, \\ 0 & \text{otherwise.} \end{cases}$$

We also need the following property of $\gamma_{N,r}(n)$:

Lemma 6.2. *If $N \in 2\mathbb{N}$, $0 \leq r \leq N - 1$, and $\frac{h}{k} \in \widehat{\mathcal{Q}}_{N,r}$, then $\gamma_{N,r}(n)$ is periodic of mean value zero.*

The periodicity property is clear. To prove the mean value zero property, it suffices to show that the quadratic Gauss sum $\mathcal{G}(\frac{N}{2}h, hr, k)$ is zero, where for $a, b \in \mathbb{Z}$, $c \in \mathbb{N}$ we define

$$\mathcal{G}(a, b, c) := \sum_{n=0}^{c-1} e\left(\frac{an^2 + bn}{c}\right).$$

We can evaluate these sums using the following easily-verified facts about Gauss sums.

Lemma 6.3. *Let a, b, c be integers. Then $\mathcal{G}(a, b, c) = 0$ if any of the following are satisfied.*

- (1) *We have that $(a, c) > 1$ and $(a, c) \nmid b$.*
- (2) *We have that $c \equiv 0 \pmod{4}$ and b is odd.*
- (3) *We have that $b = 0$ and $c \equiv 2 \pmod{4}$.*

We are now in position to prove Lemma 6.2.

Proof of Lemma 6.2: For $r \not\equiv 0 \pmod{\frac{N}{2}}$, by definition of the quantum set we have $\frac{N}{2} \mid k$, but hr is not divisible by $\frac{N}{2}$. Thus, the Gauss sum $\mathcal{G}(\frac{N}{2}h, hr, k)$ is zero by (1) of Lemma 6.3.

For $r = \frac{N}{2}$, note that $\mathcal{G}(\frac{N}{2}h, \frac{N}{2}h, k) = \mathcal{G}(m, m, n)$, where $m := \frac{\frac{N}{2}h}{(\frac{N}{2}, k)}$, $n := \frac{k}{(\frac{N}{2}, k)}$. By assumption on the quantum set, we see that m is odd and $n \equiv 0 \pmod{4}$ and so the Gauss sum vanishes by (2) of Lemma 6.3.

Finally, if $r = 0$, we reduce the Gauss sum to $\mathcal{G}(m, 0, n)$ where $m := \frac{\frac{N}{2}h}{(\frac{N}{2}, k)}$ and $n := \frac{k}{(\frac{N}{2}, k)}$. By assumption on the quantum set, we have $n \equiv 2 \pmod{4}$. The result follows by (3) of Lemma 6.3. \square

We now return to our discussion of asymptotic expansions. First we recall that for any arithmetic function $\chi: \mathbb{Z} \rightarrow \mathbb{C}$ and $s \in \mathbb{C}$, we may formally define an associated L -function

$$L(s, \chi) := \sum_{n \geq 1} \frac{\chi(n)}{n^s}.$$

Then the asymptotic expansions for $\Theta_{\frac{3}{2}}^+(N, r; \tau)$ and $\Theta_{\frac{1}{2}}^*(N, r; \tau)$ are given as follows.

Proposition 6.4. *For any $N \in 2\mathbb{N}$, $0 \leq r \leq N - 1$ and $\frac{h}{k} \in \widehat{\mathcal{Q}}_{N, r}$, we have as $t \rightarrow 0^+$*

$$\Theta_{\frac{3}{2}}^+\left(N, r; \frac{h}{k} + \frac{it}{2\pi}\right) + \frac{r}{N} q^{\frac{r^2}{2N}} \sim \Theta_{\frac{1}{2}}^*\left(N, r; \frac{h}{k} - \frac{it}{2\pi}\right) \sim \frac{1}{N} \sum_{n \geq 0} \frac{L(-2n - 1, \gamma_{N, r})}{n!} \left(-\frac{t}{2N}\right)^n.$$

Here $L(-2n - 1, \gamma_{N, r})$ is defined by the analytic continuation of $L(s, \chi)$ to \mathbb{C} .

The key tool for proving these asymptotic expansions is the following.

Lemma 6.5. *Let $\chi: \mathbb{Z} \rightarrow \mathbb{C}$ be a periodic function with mean value 0. Then $L(s, \chi)$ extends holomorphically to all of \mathbb{C} and we have as $t \rightarrow 0^+$*

$$\sum_{n \geq 1} n \chi(n) e^{-n^2 t} \sim \sum_{n \geq 0} L(-2n - 1, \chi) \frac{(-t)^n}{n!}.$$

If χ is even, then we also have as $t \rightarrow 0^+$

$$(6.8) \quad \sum_{n \geq 1} n \chi(n) \Gamma\left(-\frac{1}{2}; 2n^2 t\right) e^{n^2 t} \sim -2\sqrt{\pi} \sum_{n \geq 0} L(-2n - 1, \chi) \frac{t^n}{n!},$$

where $\Gamma(\ell; t) := \int_t^\infty e^{-u} u^{\ell-1} du$.

Proof: The first claim can be proven using standard Mellin transform techniques, for example see the proof of the proposition on page 99 of [29]. For the second claim we proceed similarly as in [29]. The idea is to compute the Mellin transform of (6.8) in two different ways. Making a simple change of variables and interchanging summation and integration, we see that for $\text{Re}(s) > 1$

$$(6.9) \quad \int_0^\infty \sum_{n \geq 1} \chi(n)n\Gamma\left(-\frac{1}{2}; 2n^2t\right) e^{n^2t} t^{s-1} dt = L(2s-1, \chi) \int_0^\infty \Gamma\left(-\frac{1}{2}; 2t\right) e^{t} t^{s-1} dt.$$

Expanding the left hand side of (6.9), we find for any $m \in \mathbb{N}_0$:

$$(6.10) \quad \sum_{n \geq 1} \chi(n)n\Gamma\left(-\frac{1}{2}; 2n^2t\right) e^{n^2t} = \sum_{r=-1}^m \frac{b_r}{\left(\frac{r}{2}\right)!} t^{\frac{r}{2}} + O\left(t^{\frac{m+1}{2}}\right)$$

for certain coefficients $b_{\frac{r}{2}}$. The existence of this expansion can be seen using a shifted version of Proposition 3 of [43] (see Remark 1 following Proposition 3). We omit the details, however we remark that the proof requires the assumption that $\chi(n)$ has mean value zero together with the following behavior near 0 and ∞ :

$$(6.11) \quad \Gamma\left(-\frac{1}{2}; t\right) \sim t^{-\frac{3}{2}} e^{-t} \quad \text{as } t \rightarrow \infty,$$

$$\Gamma\left(-\frac{1}{2}; 2t\right) \sim \sqrt{2}t^{-\frac{1}{2}} - 2\sqrt{\pi} + t^{\frac{1}{2}} \sum_{j \geq 0} \alpha_j t^j \quad \text{as } t \rightarrow 0.$$

In fact, we will soon show that only even r occur in (6.10). Using (6.11), we see that the left hand side of (6.9) equals

$$\int_0^1 \left(\sum_{r=-1}^m \frac{b_r}{\left(\frac{r}{2}\right)!} t^{\frac{r}{2}} + O\left(t^{\frac{m+1}{2}}\right) \right) t^{s-1} dt + \int_1^\infty O\left(e^{-t}\right) t^{s-1} dt = \sum_{r=-1}^m \frac{b_r}{\left(\frac{r}{2}\right)! \left(\frac{r}{2} + s\right)} + f_m(s),$$

where f_m is analytic for $\text{Re}(s) > -\frac{m}{2}$. This gives that the residue of (6.9) at a half integer $-\frac{r}{2}$ is $\frac{b_r}{\left(\frac{r}{2}\right)!}$. We can determine the coefficients $b_{\frac{r}{2}}$ by computing the residues of (6.9) using the right hand side. First note that, as $t \rightarrow 0$,

$$(6.12) \quad \left(\Gamma\left(-\frac{1}{2}; 2t\right) + 2\sqrt{\pi} \right) e^t \sim \sqrt{2}t^{-\frac{1}{2}} + t^{\frac{1}{2}} \sum_{n \geq 0} \beta_n t^n$$

is an odd Laurent series in $t^{\frac{1}{2}}$. To see this, we use that

$$\Gamma\left(-\frac{1}{2}; x\right) = \frac{2e^{-x}}{\sqrt{x}} + 2\sqrt{\pi} (\text{erf}(\sqrt{x}) - 1),$$

where $\text{erf}(z) := \frac{2}{\sqrt{\pi}} \int_0^z e^{-x^2} dx$ is the usual error function and that $\text{erf}(-z) = -\text{erf}(z)$. We then compute

$$\int_0^\infty \Gamma\left(-\frac{1}{2}; 2t\right) e^{t} t^{s-1} dt = I_0 + I_1 + I_2$$

where

$$I_0 := -2\sqrt{\pi} \int_0^1 e^{t} t^{s-1} dt,$$

$$I_1 := \int_0^1 \left(\Gamma\left(-\frac{1}{2}; 2t\right) + 2\sqrt{\pi} \right) e^{t} t^{s-1} dt,$$

$$I_2 := \int_1^\infty \Gamma\left(-\frac{1}{2}; 2t\right) e^{tt^{s-1}} dt.$$

Using the formula

$$(6.13) \quad \int_0^1 e^{tt^{s-1}} dt = \sum_{j \geq 0} \frac{1}{j!(s+j)},$$

we find that I_0 has simple poles precisely at $s = -j$ for $j \in \mathbb{N}_0$ with residue $\frac{-2\sqrt{\pi}}{j!}$. By (6.12), another application of (6.13) shows that I_1 has no poles at negative integers. Finally, it is easy to prove that I_2 is entire. Thus, the residue of (6.9) at any negative integer $-j$ is $\frac{-2\sqrt{\pi}L(-2j-1, \chi)}{j!}$. Note that by the remark on page 99 of [29], $L(-2n, \chi) = 0$ for all $n \in \mathbb{N}$ as $\chi(n)$ is even. As I_1 only has simple poles, the zeros from the L -function cancels any pole of I_1 , so that the residue of (6.9) is zero at any half integer. Thus, b_j is only non-zero for integers j , in which case $b_j = -2\sqrt{\pi}L(-j-1, \chi)$. \square

We now have the necessary tools to determine the asymptotic expansions of $\Theta_{\frac{3}{2}}^+(N, r; \tau)$ and $\Theta_{\frac{1}{2}}^*(N, r; \tau)$.

Proof of Proposition 6.4: For $t > 0$, we use (6.2) to write

$$\Theta_{\frac{3}{2}}^+\left(N, r; \frac{h}{k} + \frac{it}{2\pi}\right) + \frac{r}{N} q^{\frac{r^2}{2N}} = \frac{1}{N} \sum_{\substack{n > 0 \\ n \equiv \pm r \pmod{N}}} ne\left(\frac{hn^2}{2kN}\right) e^{-\frac{tn^2}{2N}} = \frac{1}{N} \sum_{n > 0} n\gamma_{N,r}(n) e^{-\frac{tn^2}{2N}}.$$

This expansion, combined with Lemma 6.5 and Lemma 6.2, gives the asymptotic expansion for $\Theta_{\frac{3}{2}}(N, r; \tau)$ as $\tau \rightarrow \frac{h}{k}$. We next turn to $\Theta_{\frac{1}{2}}^*$. By a simple change of variables and term-by-term integration, we find for $\tau \in \mathbb{H}_-$

$$\begin{aligned} \Theta_{\frac{1}{2}}^*(N, r; \tau) &= -\frac{1}{2N\sqrt{\pi}} \sum_{\substack{n > 0 \\ n \equiv \pm r \pmod{N}}} n\Gamma\left(-\frac{1}{2}; -4\pi n^2 \frac{\text{Im}(\tau)}{2N}\right) q^{\frac{n^2}{2N}}, \\ \Theta_{\frac{1}{2}}^*\left(N, r; \frac{h}{k} - \frac{it}{2\pi}\right) &= -\frac{1}{2N\sqrt{\pi}} \sum_{n > 0} n\gamma_{N,r}(n) \Gamma\left(-\frac{1}{2}; -\frac{2n^2 t}{2N}\right) e^{-\frac{n^2 t}{2N}}. \end{aligned}$$

Noting that $\gamma_{N,r}(n)$ is even, another application of Lemma 6.5 gives the result. \square

Finally, we describe the cocycles for $\Theta_{\frac{1}{2}}^*(N, r; \tau)$. Since the proof is standard we have omitted it here; the necessary modular transformations of $\tilde{\vartheta}_{\frac{1}{2}}(N, r; \tau)$ are given in Proposition 3.2.

Lemma 6.6. For any $\gamma = \begin{pmatrix} a & b \\ c & d \end{pmatrix} \in \Gamma_1(2N)$, we have

$$\Theta_{\frac{1}{2}}^*(N, r; \gamma\tau) \chi_r(\gamma)^{-1} (c\tau + d)^{-\frac{3}{2}} - \Theta_{\frac{1}{2}}^*(N, r; \tau) = r_{-\frac{d}{c}}(\tau),$$

where

$$r_x(\tau) := \frac{1}{2\pi\sqrt{iN}} \int_x^{i\infty} \frac{\tilde{\vartheta}_{\frac{1}{2}}(N, r; z)}{(z - \tau)^{\frac{3}{2}}} dz.$$

We are now in a position to prove Theorem 1.5.

Proof of Theorem 1.5: By combining Lemma 6.6 and Proposition 6.3, we find that $\Theta_{\frac{3}{2}}^+(N, r; \tau)$ is a strong quantum modular form. More specifically, for any $\gamma \in \Gamma_1(2N)$, we easily see that the cocycle $r_{-\frac{a}{c}}(\tau)$ extends to a real-analytic function on $\mathbb{R} \setminus \{\gamma^{-1}\infty\}$. As $2\Theta_{\frac{3}{2}}(N, r; \tau) = \Theta_{\frac{3}{2}}^+(N, r; \tau) + \Theta_{\frac{3}{2}}^-(N, r; \tau)$, it suffices to prove that $\Theta_{\frac{3}{2}}^-(N, r; \tau)$ is a strong quantum modular form on the same congruence subgroup and with the same multiplier. First note that any rational power of q is a strong quantum modular form as its cocycle is real-analytic on \mathbb{R} . Thus, by Proposition 3.2 and (6.1), $\Theta_{\frac{3}{2}}^-(N, r; \tau)$ is strongly quantum modular with the same multiplier and congruence subgroup as $\tilde{\vartheta}_{\frac{3}{2}}(N, r; \tau)$ is, which is the same as for $\Theta_{\frac{3}{2}}^+(N, r; \tau)$ by Lemma 6.6. \square

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