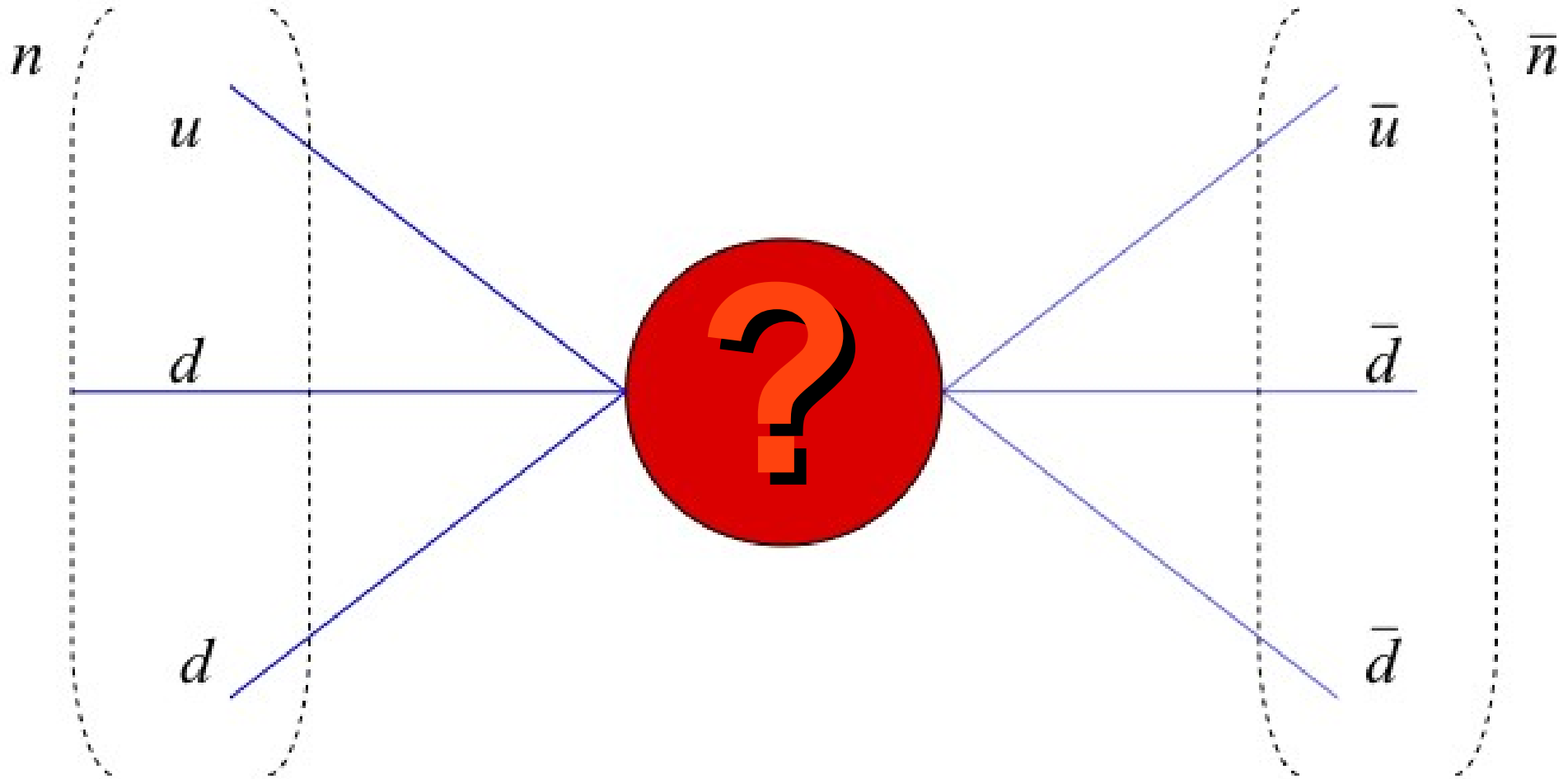


# Neutron Majorana mass from exotic instantons



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Max Born Symposium 2015*

# References

- 1) A. Addazi and M. Bianchi, arXiv:1407.2897
- 2) A. Addazi, arXiv:1501.04660
- 3) A. Addazi and M. Bianchi, arXiv:1502.01531
- 4) A. Addazi and M. Bianchi, arXiv:1502.08041
- 5) A. Addazi, arXiv:1504.06799 [hep-ph];
- 6) A. Addazi, arXiv:1505.00625 [hep-ph].
- 7) A. Addazi, arXiv:1505.02080 [hep-ph].
- 8) A. Addazi, arXiv:1506.06351.

# Majorana mass for a neutron?

A crazy idea suggested by  
Ettore Majorana

*Nuovo Cimento '37'*

(Idea reconsidered some years later by Bruno Pontecorvo..)

$$(udd)^2 / \mathcal{M}_{n\bar{n}}^5$$

$$\Delta B = 2$$

$$\propto n n + \text{h.c.}$$



**Physics Beyond Standard Model!  
Baryogenesis, Leptogenesis...**

# Why New Physics

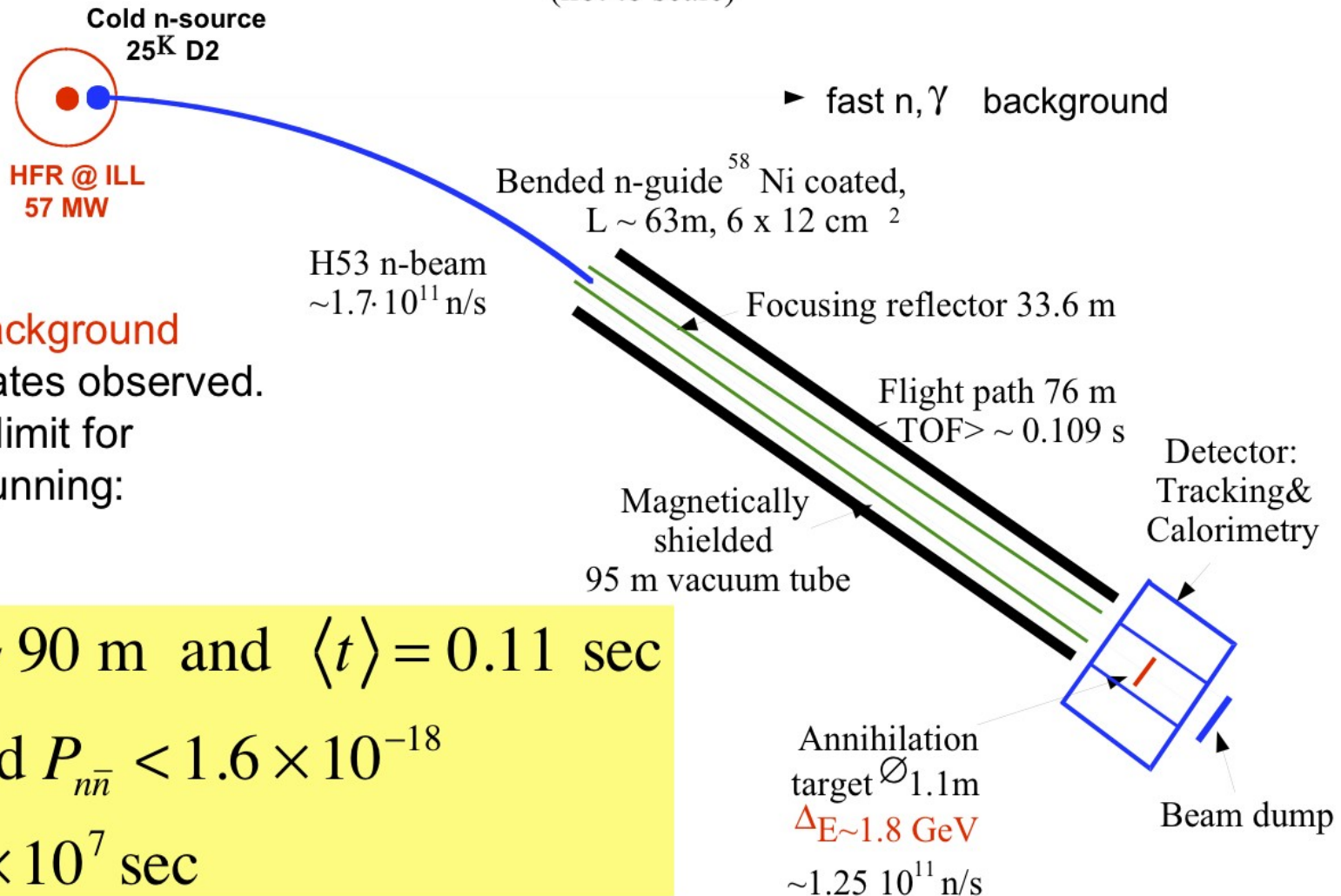
- **B,L** are accidental symmetry of SM lagrangian
- **Vafa-Witten theorem** protects non-perturbative QCD against dynamical violations of vector-like global symmetries, like B-number.
- **Sphalerons** can only change B-violations into L-one and viceversa. **They are suppressed** as  $10^{(-120)}$  in laboratory. They are relevant in early thermal bath.

# $n$ - $n\bar{n}$ experiments\*

See talks of Snow, Young and Kamyshev at BLV 2011

Heidelberg - ILL - Padova - Pavia  $n\bar{n}$  search experiment  
at Grenoble 89-91

(not to scale)



No GeV background

No candidates observed.

Measured limit for  
a year of running:

with  $L \sim 90 \text{ m}$  and  $\langle t \rangle = 0.11 \text{ sec}$

measured  $P_{n\bar{n}} < 1.6 \times 10^{-18}$

$\tau > 8.6 \times 10^7 \text{ sec}$

Baldo-Ceolin M. et al., Z. Phys. C63,409 (1994).

# **NNBar, why not??? ....KKbar oscillations observed...**

Very interesting if we compare this one  
with **Proton decay limit** (33<sup>th</sup> order higher)  
and **Neutrinoless Double Beta decay** (23<sup>th</sup> order)

*Theoretical side?*

**1) R-breaking MSSM connects nnbar with p-decays...Extra protecting symmetries?  
Alternatively?**

**2) Babu-Mohapatra GUT SO(10) without susy, 126  
Multiplets**

**3) Bary-majoron: Mohapatra, Barbieri, Dvali,  
Bereziani  
Bary-majoron and RH-neutron**

**4) unesthetic toy-models, 'ad hoc' with colored sextets,  
triplets of n-plets**

# Don't worry be happy: no matter destabilization during my PHD

$$\mathcal{H}_{\text{eff}} = \begin{pmatrix} m_n - V_n & \alpha \\ \alpha^* & m_n - V_{\bar{n}} \end{pmatrix} \quad (1)$$

where  $V_n$  and  $V_{\bar{n}}$  are the binding energies in the nucleus for a neutron and an antineutron.  $V_{\bar{n}} \ll V_n$ ,  $|V_{\bar{n}} - V_n| = \Delta V \sim V_n \sim -10\text{MeV}$ . The neutron in the nucleus is 'free' for a time that can be estimated from the indetermination principle

$$\Delta E \Delta t \sim 1 \longrightarrow t_{\text{free}} \sim \frac{1}{E_{\text{leg}}} \sim 10^{-23}\text{s}$$

The oscillation probability is

$$P_{n\bar{n}} = \frac{\alpha^2}{\alpha^2 + \Delta V^2} \sin^2 \sqrt{\alpha^2 + \Delta V^2} t \simeq \frac{4\alpha^2}{(\Delta V)^2} \longrightarrow \tau_a = \frac{1}{p_A} \sim 10^{30}\text{yr}$$

where  $\tau_A$  is the internuclear transition lifetime, and  $p_A$  the transition probability per second.

# Why the suppression of the Magnetic Field (with $t \ll \text{Beta decay}$ )

$$\mathcal{H}_{eff} = \begin{pmatrix} m + 2\omega & \alpha \\ \alpha^* & m - 2\omega \end{pmatrix}$$

with  $2\omega = |\mu B|$ . Rotated, with mixing angles

$$\tan 2\theta = \frac{\alpha}{\omega}; \quad \sin^2 2\theta = \frac{\alpha^2}{\alpha^2 + \omega^2}$$

With  $\Omega = \sqrt{\alpha^2 + \omega^2}$ , the transition probability has the form

$$\mathcal{P}_{n\bar{n}} = \sin^2 2\theta \sin^2(\Omega t) = \frac{\alpha^2}{\alpha^2 + \omega^2} \sin^2(\sqrt{\alpha^2 + \omega^2} t)$$

$B_{\text{Earth}} \sim 0.5 \text{ Gauss} \longrightarrow \omega \gg \alpha \sim 10^{-23} \text{ eV} \longrightarrow \sqrt{\alpha^2 + \omega^2} t \gg 1$ ,

from (4)  $\longrightarrow \mathcal{P}_{n\bar{n}} = \frac{\alpha^2}{2\omega^2} \simeq 0$ , there aren't any chances to try  $n - \bar{n}$



# Challenge accepted!!!

*We propose the possibility that such a  $NN\bar{b}$  transition, competitive with future experiments, can be generated by Non-perturbative effects of quantum gravity.*

**Exotic Instantons** in String theories!  
*Not existing at all in gauge theories  
(out of ADHM classifications)*

# Instantons in String theory

A simple geometric interpretation:  
nothing but “special D-branes”,  
Eucliden D-branes (E-branes) wrapping an  
internal cycle, that could intersect the  
`physical' D-branes.

**Exotic instantons are represented by E-  
branes not wrapping the same cycles of the  
ordinary D-branes!**

They are not in ADHM construction!

On the other hand, gauge instantons are  
wrapping the same cycles.

## **Examples:**

**I) in (un)oriented type IIA, instantons are E2  
branes wrapping 3 cycle**

**II) type I, E5 in internal space, with same  
magnetization of D9 (wrapping the entire CY3)**

**III) in (un)oriented IIB E(-1) or E3 wrapping  
wrapping the same holomorphic divisor as a  
stack of physical D7...etc..**

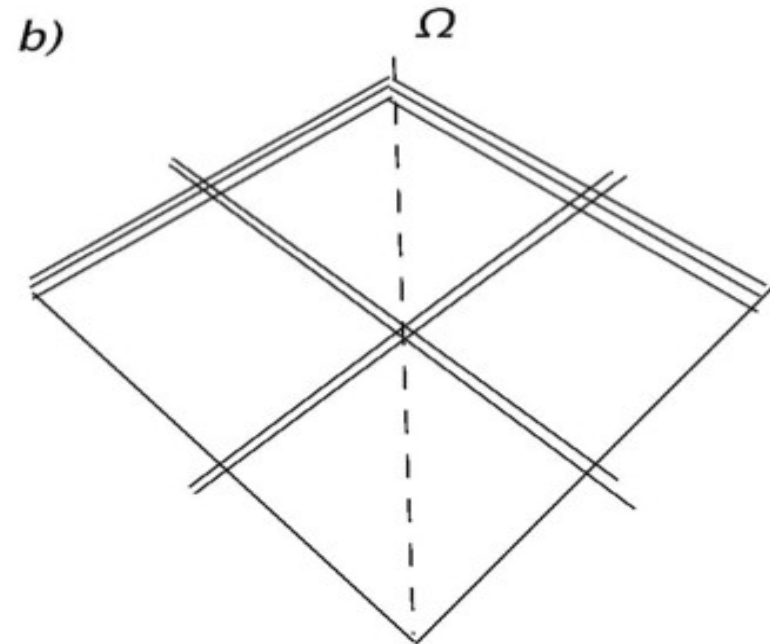
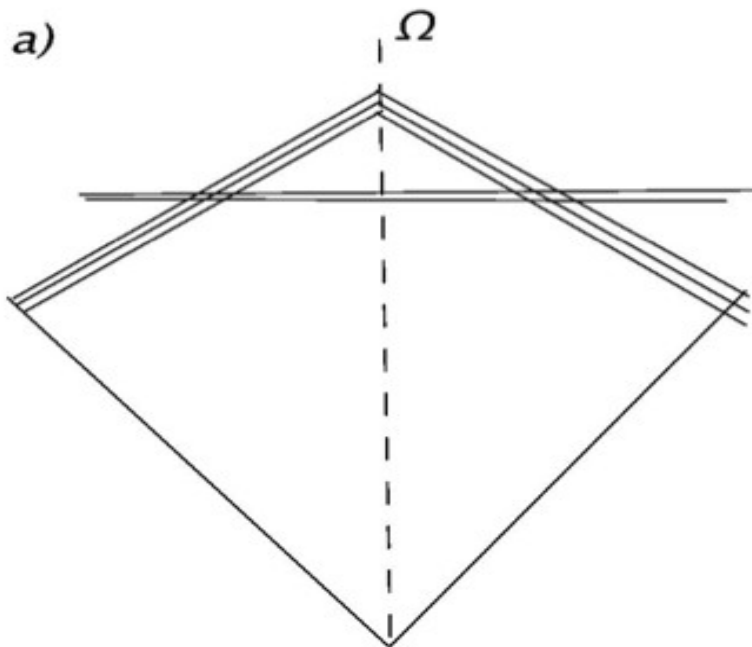
# We consider a simple class of models:

**Instruments:** unoriented string theory, D6-branes stacks  
Wrapping 3-cycles in CY3, Antisymmetric Mirror plane,  
E2-branes (gauge and exotic instantons).

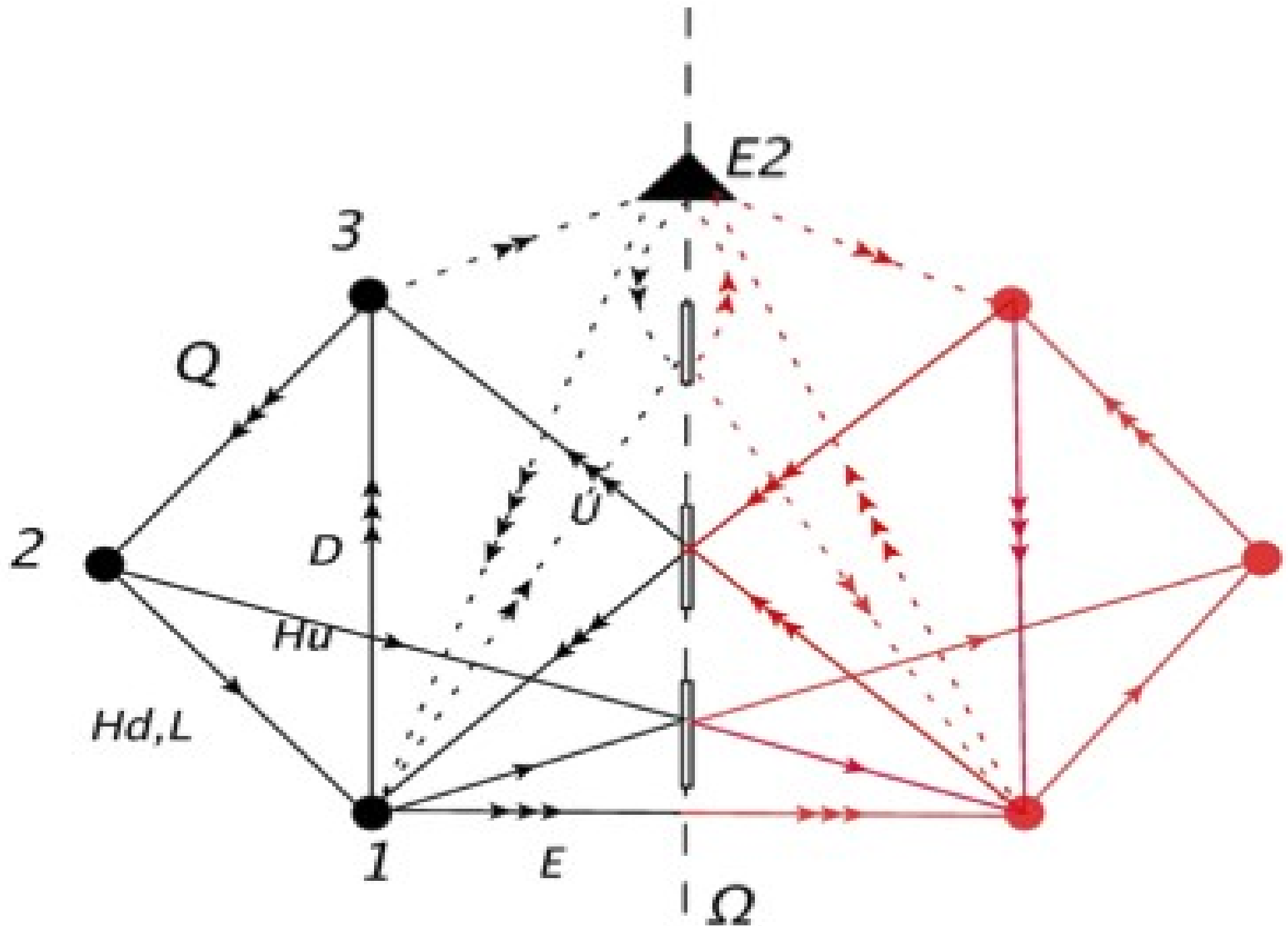
**“Music”:** (MS)SM + 4 extra U(1)

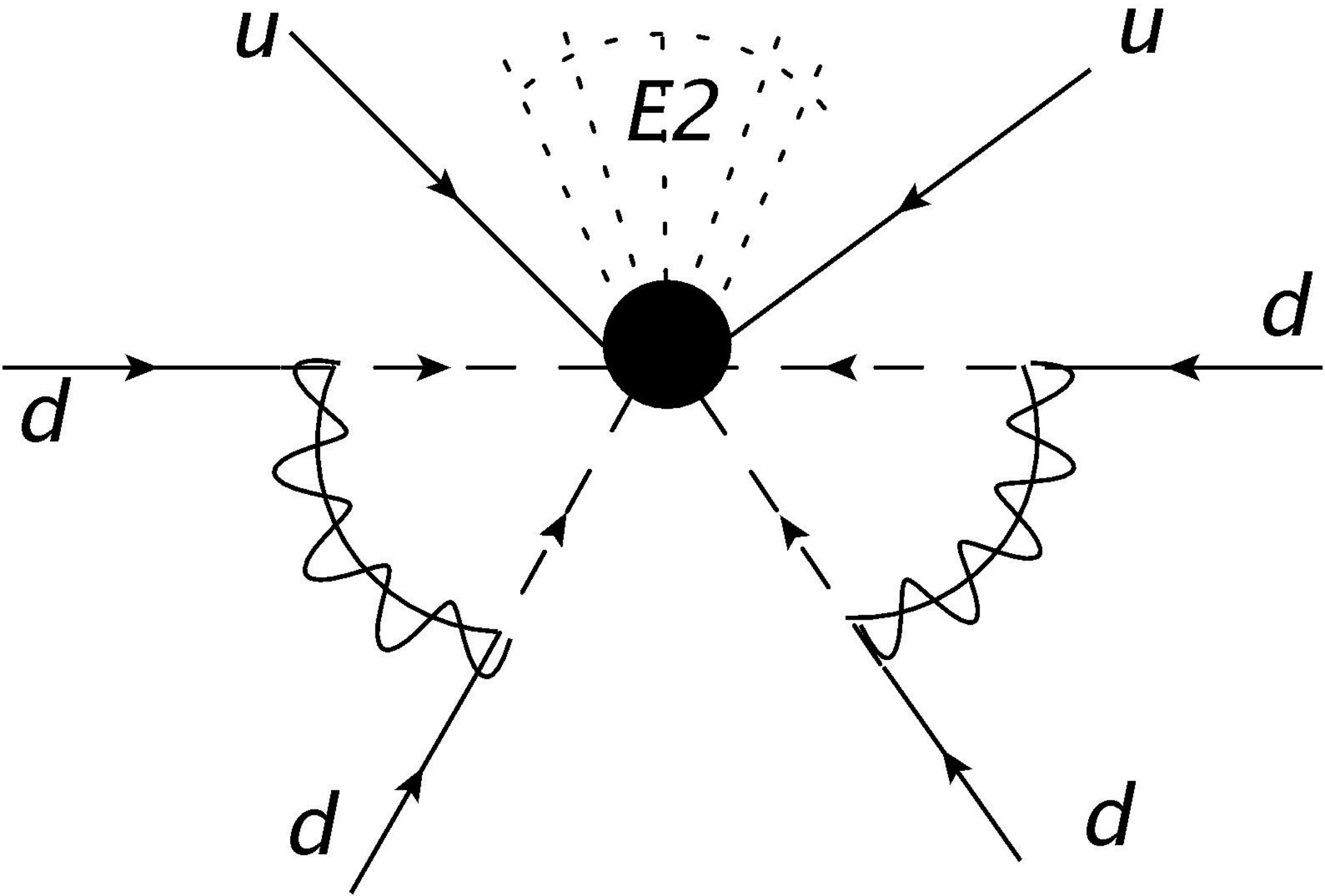
*D. Cremades, L. E. Ibanez and F. Marchesano,  
JHEP 0307, 038 (2003) [hep-th/0302105]. (And other many papers)*

**Non perturbative Mixing generated by exotic instantons**



# Direct Generation





# Calculations from mixed disk amplitudes

$$\mathcal{L}_{eff} \sim k_f^{(1)} U_f^i \tau_i \alpha + k_f^{(2)} D_f^i \tau_i \beta$$

$$W_{E2-D6-\hat{D}6} = \int d^6 \tau d^4 \beta d^2 \alpha e^{\mathcal{L}_{eff}} = \mathcal{Y}^{(1)} \frac{e^{-S_{E2}}}{M_S^3} \epsilon_{ijk} \epsilon_{i'j'k'} U^i D^j D^k U^{i'} D^{j'} D^{k'}$$

$$\mathcal{Y}_{f_1 f_2 f_3 f_4 f_5 f_6}^{(1)} = k_{f_1}^{(1)} k_{f_2}^{(1)} k_{f_3}^{(2)} k_{f_4}^{(2)} k_{f_5}^{(2)} k_{f_6}^{(2)}$$

$$\mathcal{O}_{n\bar{n}} + \mathcal{O}_{\Lambda\bar{\Lambda}} = \frac{y_1}{\mathcal{M}_{E2}^3 M_{SUSY}^2} (u^c d^c d^d)(u^c d^c d^c) + \frac{y_2}{\mathcal{M}_{E2}^3 M_{SUSY}^2} (u^c d^c s^c)(u^c d^c s^c)$$

$$\mathcal{M}_{E2}^3 = e^{+S_{E2}} M_S^3, \quad y_1 = \mathcal{Y}_{111111}^{(1)} \quad \text{and} \quad y_2 = \mathcal{Y}_{112112}^{(1)}$$

# Phenomenology

*Next future on  $NN\bar{b}$ : 1000 TeV*

*Compatible with:*

- ***TeV-susy,  $M_S=10^5$  TeV***  
*with large (3-cycles)*
- ***$M_S=M_{SUSY}=10$  TeV, factor***  
*100 (3-cycles).*
- ***$M_S=M_{SUSY}=100-1000$  TeV factor***  
*10-1 (3-cycles).*



- **Stringy resonances and anomalous  $Z'$**  for LHC or future collider
- Exotic instantons are classical configuration in B-violating scattering amplitudes. **Cutoff of the cross section expected**
- **No proton destabilization, no FCNCs** related to  $NN\bar{b}$  diagram

# Other considerations

- This is a **non-Wilsonian UV completion** of a six quark effective operator.
- R-parity is **dynamically broken**.  
Subtly compatible with gauge invariance
- Other operator like **Weinberg's  $W=HLHL/M$**  can be similarly generated and tested in colliders.

# Other alternative models

- Models with a B-violating mixing of a vector-like pair  $3(B=-2/3)-3\bar{3}(B=-1/3)$  (A.A, M.Bianchi, JHEP 2014; A.A JHEP 2015; A.A, M.Bianchi, JHEP 2015)
- Pati-Salam like with 10-plets  
(A.A, M.Bianchi JHEP 2015)
- uddX/M, with X a singlet fermion (A.A, 2015)

# Conclusions

String theory could have a peculiar phenomenology in neutron physics.

In particular Exotic Instantons could generate a Majorana mass term for the neutron, without other dangerous consequences like fast proton decay or strong FCNCs.

In some subregions of parameters we have other channels interesting for LHC or future colliders.

Exotic Instantons could be the portal to Quantum Gravity, inducing no-perturbative coupling at scale like 1000TeV.

**THANK YOU!!!**