



Neutron Optical Beam Splitter from Holographically Structured Nanoparticle-Polymer Composites

M. Fally,^{1,*} J. Klepp,¹ Y. Tomita,² T. Nakamura,² C. Pruner,³ M. A. Ellabban,⁴ R. A. Rupp,^{1,5} M. Bichler,¹ I. Drevenšek Olenik,^{6,7} J. Kohlbrecher,⁸ H. Eckerlebe,⁹ H. Lemmel,¹⁰ and H. Rauch^{10,11}

¹Faculty of Physics, University of Vienna, Boltzmanngasse 5, A-1090 Wien, Austria

²Department of Electronics Engineering, University of Electro-Communications, 1-5-1 Chofugaoka, Chofu, Tokyo 182, Japan

³Department of Materials Science and Physics, University of Salzburg, A-5020 Salzburg, Austria

⁴Physics Department, Faculty of Science, Tanta University, Tanta 31527, Egypt

⁵Nankai University, TEDA Applied Physics Faculty, Tianjin 300457, People's Republic of China

⁶Faculty of Mathematics and Physics, University of Ljubljana, Jadranska 19, SI 1001 Ljubljana, Slovenia

⁷J. Stefan Institute, Jamova 39, SI 1001 Ljubljana, Slovenia

⁸Laboratory for Neutron Scattering, ETH Zurich & Paul Scherrer Institut, 5232 Villigen PSI, Switzerland

⁹GKSS Forschungszentrum, D-21502 Geesthacht, Germany

¹⁰Institut Laue Langevin, Boîte Postale 156, F-38042 Grenoble Cedex 9, France

¹¹Atominstytut, Technische Universität Wien, Stadionallee 2, 1020 Wien, Austria

(Received 16 February 2010; published 17 September 2010)

We report a breakthrough in the search for versatile diffractive elements for cold neutrons. Nanoparticles are spatially arranged by holographical means in a photopolymer. These grating structures show remarkably efficient diffraction of cold neutrons up to about 50% for effective thicknesses of only 200 μm . They open up a profound perspective for next generation neutron-optical devices with the capability to tune or modulate the neutron diffraction efficiency.

DOI: 10.1103/PhysRevLett.105.123904

PACS numbers: 82.35.Np, 42.40.Eq, 61.05.F–

Neutron-optical phenomena arise from coherent elastic scattering. They are described by the wave equation for neutrons. The essential material property is the neutron-optical potential or, equivalently, the neutron refractive index at wavelength λ [1]. The basic diffractive element is a (one-dimensional) sinusoidal grating characterized by a refractive index $n(x) = n_0 + n_1 \cos(Kx)$ which is periodically modulated with the spatial frequency K and amplitude n_1 . Depending on its diffraction efficiency, it can be used as beam splitter, mirror, or monochromator, and it can be altered and arranged to form more complex devices for imaging, spectroscopy, and in particular, cold neutron interferometry [2–4]. Efficient neutron-optical devices are vital for any neutron experiment, either as a part of the instrumentation (e.g., monochromators, guides, calibration standards) or as a tool to investigate fundamental physical questions (e.g., interferometers). Neutron optics—and, in particular, interferometry with thermal neutrons using perfect crystals as diffractive elements—has reached a satisfying level giving rise to an appealing insight on quantum mechanical problems [5]. As the difference of the refractive index for neutrons from unity depends quadratically on the wavelength the latter should be further increased to maximize the neutron-optical potential. On the other hand the neutron flux decreases dramatically at low energies. On top of that, interferometry experiments with very cold neutrons, so far, employed ruled or blazed gratings as thin diffraction elements. Consequently, various diffraction orders are excited at the same time resulting in even lower measuring signals

[6]. Therefore, cold neutrons offer a fair tradeoff between flux and magnitude of the neutron-optical potential. Moreover, small angle neutron scattering (SANS) beam lines are workhorses at neutron facilities due to the interdisciplinary interests of biologists or medical scientists when exploring complex large-scale soft matter; see, e.g., [7]. A few attempts have been conducted to manufacture diffractive elements for cold neutrons. However, the results on the efficiency were rather discouraging [2,8–10].

In this Letter we introduce holographically patterned nanoparticle-polymer composite systems as efficient and extremely flexible diffractive elements for cold neutrons.

Holographic production of neutron-optical elements can be performed in materials where the neutron refractive index can be changed by photon irradiation. This is the so-called photo-neutron refractive effect (PNRE) [8,11]. Upon illumination with a spatial light pattern a corresponding refractive-index modulation $\Delta n = (\lambda^2/2\pi)\Delta\mathcal{B}$ occurs that is proportional to the local change $\Delta\mathcal{B}$ of the coherent scattering length density \mathcal{B} . In the simplest case $\mathcal{B} = b_c\rho$ is the product of the coherent scattering length b_c and the number density of atoms ρ . The PNRE may result from a photoinduced modulation of either of them.

Up to now only single-component systems were investigated and no suitable light sensitive materials have been found so far. Either fundamental problems on the neutron-optical or on the light-optical side turned out to be too difficult to overcome. This has recently led us to turn the attention to multicomponent systems. Recently, a success-

ful attempt was made by employing holographic polymer dispersed liquid crystals [10]. The figure of merit—the coherent scattering length density modulation—is strong ($\Delta\mathcal{B} \approx 10 \mu\text{m}^{-2}$) in these materials. However, the disadvantage is its restricted neutron diffraction efficiency due to a limited grating thickness d of about only $30 \mu\text{m}$. The reason for this is a strong diffuse scattering and/or coherent light-induced scattering (holographic scattering [12]) originating from the anisotropy of the liquid crystals component which destroys the incident light pattern when propagating through the sample. For low efficiencies of less than a few percent the relation to the coherent scattering length density modulation is independent of the diffraction regime and reads $\eta_1 = \nu^2$ with the normalized coupling length $\nu = \lambda\Delta\mathcal{B}d/2$.

Preparation and handling of various nanoparticle-polymer composites nowadays have reached a pretty mature state [13]. In particular, the groups of Tomita *et al.* [14–17] and Stumpe *et al.* [18,19] have intensively investigated inorganic nanoparticles (SiO_2 and ZrO_2) embedded in a photopolymer matrix. They serve as proof of principle for the essential advantages offered by this class of materials for neutron optics over previously used pure photopolymer systems [8,11,20]: the refractive-index modulation can be tuned by the type of nanoparticles, their size, and volume ratio. An additional benefit of including nanoparticles constitutes the increased mechanical stability because shrinkage, which is typical for polymerization processes, is strongly reduced [14]. Stability is a decisive advantage when it comes to the construction of cold neutron interferometers. Moreover, the composites permit us to produce thicker gratings with mirrorlike reflection. For light-optical purposes it has indeed already been shown that grating thicknesses up to 0.2 mm can be achieved with only moderate disturbance of the interference pattern by holographic scattering [21,22].

SiO_2 (ZrO_2) nanoparticles having an average core diameter of approximately 13 (3) nm were prepared by liquid-phase synthesis and dissolved in a methyl isobutyl ketone (toluene) solution. The SiO_2 sol was dispersed to methacrylate monomer 2-methyl-acrylic acid 2-4-[2-(2-methyl-acryloyloxy)-ethylsulfanyl-methyl]-benzylsulfanyl-ethyl ester, whose refractive index is 1.59 in the solid phase at 589 nm . The ZrO_2 sol was dispersed to acrylate monomer 2-propenoic acid, (octahydro-4,7-methano-1*H*-indene-2,5-diyl)bis(methylene) ester, with an index of 1.53 in the same phase. At this wavelength the refractive index of the surface-treated SiO_2 (ZrO_2) nanoparticles is 1.46 (1.71)

We used $1 \text{ wt } \%$ titanocene (Irgacure784, Ciba) as radical photoinitiator enabling the monomer to photopolymerize at wavelengths shorter than 550 nm . The mixed syrup was cast on a spacer loaded glass plate and was dried before covering it with another glass plate to obtain film samples for light-optical characterization as well as

neutron-optical measurements. Both sample types, the one incorporating SiO_2 nanoparticles (referred to as samples S below) and the one with ZrO_2 nanoparticles (referred to as sample Z below), contained the same nanoparticle concentration of $20 \text{ vol } \%$. Further details on our sample preparation technique can be found in Refs. [16,22]. A conventional two-beam interference setup was used for recording an unslanted transmission grating employing two mutually coherent and *s*-polarized laser beams of equal intensities at $\lambda_L = 532 \text{ nm}$. A series of such S samples with different thicknesses d_0 and grating spacings Λ and a Z sample were prepared. In what follows we focus our investigations on only three of them: Z0 and S0 ($d_0 = 50 \mu\text{m}$, $\Lambda = 1 \mu\text{m}$) as well as S7 ($d_0 = 100 \mu\text{m}$, $\Lambda = 0.5 \mu\text{m}$).

In order to relate the properties of the actual gratings to those previously investigated [16,21] we first characterized them by light-optical diffraction experiments. They reveal two major common features for the gratings employing different nanoparticle species: (i) diffraction occurs clearly in the two-wave coupling regime (Bragg regime, i.e., only two diffracted beams propagate simultaneously) and (ii) the gratings are obviously nonsinusoidal for both nanoparticle species which is indicated by the existence of higher diffraction orders. The latter is due to a nonlinear phase separation mechanism between nanoparticles and polymerizing monomer during holographic recording [15]. Further, we noticed that (i) the refractive-index modulation for light is by about a factor of 2 larger for the sample Z0 as compared to the corresponding sample S0, and (ii) a smaller grating spacing leads to a reduction of the refractive-index modulation [14,16,19,21].

The SANS experiments were conducted using cold neutrons at the spallation source SINQ, SANS-1 (Paul-Scherrer Institute) and the Geesthacht Neutron Facility, SANS-2 instrument. Measurements were performed at various mean neutron wavelengths between $\lambda = 1.16 \text{ nm}$ and 2 nm having a distribution of about $\Delta\lambda/\lambda = 10\%$. For the benefit of sufficient lateral coherence, the angular spread was restricted to less than 1 mrad by using two narrow slits and full collimation distances. The diffracted intensities were measured at a sample-detector distance of about 20 m by a two-dimensional detector.

The experiments at $\lambda = 1.16 \text{ nm}$ yield first order diffraction efficiencies of only $\eta_1 = 3\%$ and $\eta_1 = 1\%$ for the samples S0 and Z0, respectively. Evaluating the coherent scattering length density modulations, we find $\Delta\mathcal{B} = 6.0 \mu\text{m}^{-2}$ and $\Delta\mathcal{B} = 3.1 \mu\text{m}^{-2}$. The obtained values are in the range of those found in holographic polymer dispersed liquid crystals with hitherto strongest neutron diffraction [10]. However, nanoparticle-polymer composites do not share the drawbacks of those materials such as the holographic light scattering obstacle [12]. Moreover, neutron diffraction is more efficient for the S than for the Z sample which is contrary to the light optic result.

Therefore, we focused our further investigations only on the SiO_2 nanoparticle composites having a different grating spacing ($\Lambda = 0.5 \mu\text{m}$) and thickness. From the light-optical experiments we expect that a smaller grating spacing also reduces the coherent scattering length density modulation for neutrons, an assumption that is completely confirmed by measuring the diffraction signal from S7: a diffraction efficiency of $\eta_1 = 6\%$ at a wavelength of 1.7 nm yields a value of $\Delta\mathcal{B} = 3.2 \mu\text{m}^{-2}$.

To facilitate the search for a route to create beam splitters or mirrors the parameter space ($\log(\nu), \log(Q)$) is divided into three diffraction regimes, determined by the equalities $Q\nu = 1 \wedge Q = 20\nu$ as shown in Fig. 1 [23]. Here, the coupling length ν and $Q = 2\pi\lambda d/\Lambda^2$ are the decisive parameters [23,24]. Each of these regimes is governed by a particular diffraction theory: the two-wave coupling regime (Bragg) described by Kogelnik's theory [25], the multiwave coupling regime neglecting dephasing (Raman-Nath theory [26]), and the regime to be evaluated according to the rigorous coupled wave analysis (RCWA) [27]. Starting from the (ν, Q) values for the samples Z0, S0 we see that even at the larger wavelengths we reside in the Raman-Nath diffraction regime. Nevertheless, the good news is that from these first experiments we could deduce a road map to produce beam splitters operating in the two-wave coupling regime. To move from the Raman-Nath to the Bragg diffraction regime just the thickness or wavelength (diagonal motion) must be increased and—most importantly—the grating spacing Λ decreased (horizontal motion). In Fig. 1 the experimentally obtained (ν, Q) values are shown together with arrows indicating the

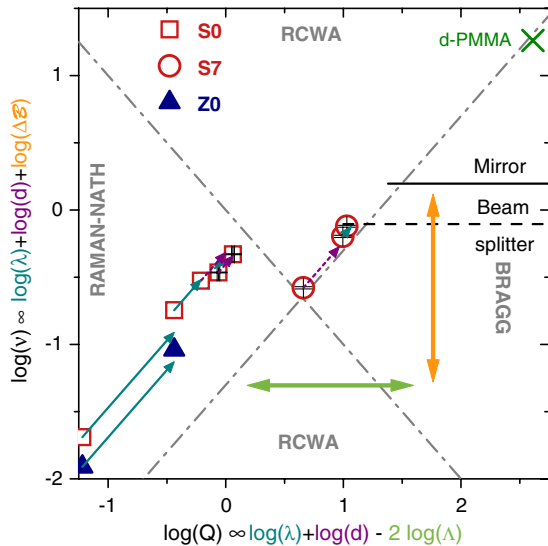


FIG. 1 (color online). Diffraction regimes separated by gray dash-dotted lines and experimental results for S0 (red squares), Z0 (blue triangles), and S7 (red circles) measured with thermal and cold neutrons. For comparison d -PMMA ($d = 3500 \mu\text{m}$, $\Lambda = 381 \text{ nm}$, $\lambda = 2.7 \text{ nm}$) is shown. Dashed or solid diagonal arrows indicate a change in d or λ .

type of parameter that was changed. A further possible path (vertical motion), i.e., in- or decreasing $\Delta\mathcal{B}$, can be easily realized by simply altering the volume ratio or the species of the nanoparticles. We would like to mention that measurements at thermal neutron wavelengths for S0 and Z0 ($\lambda = 0.19 \text{ nm}$) unexpectedly yield 25% lower values of $\Delta\mathcal{B}$ which is far out of the experimental error and not yet understood.

To demonstrate the high potential of these types of gratings we carefully chose the tunable parameters d and λ using sample S7 with $\Lambda = 0.5 \mu\text{m}$. The results are shown in Fig. 2. Here, the measured angular dependence of the efficiencies for 0th and ± 1 st diffraction orders of sample S7 is depicted. It is obvious that the grating serves nearly as a perfect 50:50 beam splitter. Neutron diffraction in this case occurs essentially in the two-wave coupling (Bragg) regime. Then the corresponding coupling length ν is given by $\nu = \arcsin(\sqrt{\eta_1})$ for the first-order diffraction [25]. For an accurate evaluation taking into account also the small contributions from multiwave coupling we employed a fit based on the RCWA [27] included in Fig. 2.

Switching back to Fig. 1, the three experimental points for S7 correspond to the parameter pairs ($\lambda = 1.7 \text{ nm}$, $d = d_0$), ($\lambda = 1.7 \text{ nm}$, $d \sim 2d_0$), and ($\lambda = 2 \text{ nm}$, $d \sim 2d_0$). These latter values allowed us to realize a beam splitter by employing (1) a reduced grating spacing as compared to S0, Z0, (2) an increased wavelength of $\lambda = 2 \text{ nm}$ as well as (3) an increased effective thickness $d = d_0/\cos(\zeta)$ by an inclination of the sample for an angle ζ around the grating vector [28].

Let us make a final note on why it is suboptimal to try the apparently simplest solution for producing a mirror, namely, to use, e.g., deuterated poly-methylmethacrylate (d -PMMA) and to increase the grating thickness to a few millimeters like done in a previous study [4]. The experi-

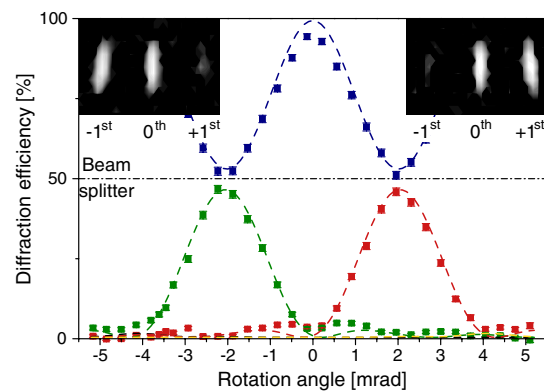


FIG. 2 (color online). Angular dependence of the 0th and ± 1 st order diffraction efficiencies for cold neutrons at $\lambda = 2 \text{ nm}$ for gratings assembled from a SiO_2 nanoparticle-polymer composite with a spacing of $\Lambda = 500 \text{ nm}$ and an effective thickness of $d = 203 \mu\text{m}$ (S7). Lines are a simultaneous RCWA fit to the data. Inset: Detector matrix (gray scale of logarithmic box counts at Bragg positions).

mentally obtained value is also shown in Fig. 1. The disadvantage of such geometrically thick gratings is their extremely high angular selectivity. In this case the rocking curve becomes a very rapidly oscillating function of the Off-Bragg parameter. Since the angular width of the rocking curve is usually much narrower than the angular spread of the beam, this results in averaging over the “Pendellösung” fringes [1,29] with a mean reflectivity that cannot exceed 50%—an issue that is well known for standard perfect crystal silicon lattices used in thermal neutron interferometry [5].

In the future we could even take advantage of the spin dependent interaction potential between nanoparticles carrying a nuclear spin and polarized neutrons, e.g., superparamagnetic or multiferroic nanoparticles. This possibility opens up new perspectives of switchable gratings by external fields.

In summary, we have realized a beam splitter for cold neutrons on the basis of nanoparticle-polymer composites, i.e., a first order diffraction efficiency of about 50% for a sample with 200 μm effective thickness. By reducing the grating spacing, increasing the grating thickness as well as the volume ratio of the nanoparticles at the same time even mirrors could be produced. The versatility (species of nanoparticles with distinct properties, volume ratio) as well as the excellent optical quality and reduced shrinkage during preparation [17] comprise the major advantages over any other system investigated up to now. Thus the experimental demonstration that nanoparticle-photopolymer composites exhibit a strong PNRE can be regarded as a breakthrough for the future development of versatile neutron diffractive elements and a new generation of cold neutron interferometers.

We are grateful to K. Momose and S. Yasui for contributing to the sample preparation and Xu Jingjun for discussion. Financial support by the Austrian Science Fund (P-18988, P-20265) and from the Ministry of Education, Culture, Sports, Science and Technology of Japan (Grant No. 20360028) is acknowledged. This research has been supported by the European Commission under the 6th Framework Programme through the Key Action: Strengthening the European Research Area, Research Infrastructures. Contract No. RII3-CT-2003-505925.

*martin.fally@univie.ac.at

<http://fun.univie.ac.at>

[1] V.F. Sears, *Neutron Optics*, Neutron Scattering in Condensed Matter Vol. 3 (Oxford University Press, New York, 1989).

- [2] U. Schellhorn, R.A. Rupp, S. Breer, and R.P. May, *Physica (Amsterdam)* **234B–236B**, 1068 (1997).
- [3] F. Pfeiffer, C. Grünzweig, O. Bunk, G. Frei, E. Lehmann, and C. David, *Phys. Rev. Lett.* **96**, 215505 (2006).
- [4] C. Pruner, M. Fally, R.A. Rupp, R.P. May, and J. Vollbrandt, *Nucl. Instrum. Methods Phys. Res., Sect. A* **560**, 598 (2006).
- [5] H. Rauch and S.A. Werner, *Neutron Interferometry*, Neutron Scattering in Condensed Matter Vol. 12 (Oxford University Press, New York, 2000).
- [6] G. van der Zouw, M. Weber, J. Felber, R. Gähler, P. Geltenbort, and A. Zeilinger, *Nucl. Instrum. Methods Phys. Res., Sect. A* **440**, 568 (2000).
- [7] S. Teixeira *et al.*, *Chem. Phys.* **345**, 133 (2008).
- [8] R.A. Rupp, J. Hehmann, R. Matull, and K. Ibel, *Phys. Rev. Lett.* **64**, 301 (1990).
- [9] H. Funahashi, T. Ebisawa, T. Haseyama, M. Hino, A. Masaike, Y. Otake, T. Tabaru, and S. Tasaki, *Phys. Rev. A* **54**, 649 (1996).
- [10] M. Fally, I. Drevenšek-Olenik, M.A. Ellabban, K.P. Pranzas, and J. Vollbrandt, *Phys. Rev. Lett.* **97**, 167803 (2006).
- [11] M. Fally, *Appl. Phys. B* **75**, 405 (2002).
- [12] M.A. Ellabban, M. Fally, H. Uršič, and I. Drevenšek-Olenik, *Appl. Phys. Lett.* **87**, 151101 (2005).
- [13] Y. Zhao *et al.*, *Nature Mater.* **8**, 979 (2009).
- [14] N. Suzuki, Y. Tomita, and T. Kojima, *Appl. Phys. Lett.* **81**, 4121 (2002).
- [15] Y. Tomita, N. Suzuki, and K. Chikama, *Opt. Lett.* **30**, 839 (2005).
- [16] N. Suzuki, Y. Tomita, K. Ohmori, M. Hidaka, and K. Chikama, *Opt. Express* **14**, 12712 (2006).
- [17] T. Nakamura, J. Nozaki, Y. Tomita, K. Ohmori, and M. Hidaka, *J. Opt. A* **11**, 024010 (2009).
- [18] O.V. Sakhno, L.M. Goldenberg, J. Stumpe, and T.N. Smirnova, *Nanotechnology* **18**, 105704 (2007).
- [19] O.V. Sakhno, L.M. Goldenberg, J. Stumpe, and T.N. Smirnova, *J. Opt. A* **11**, 024013 (2009).
- [20] R.A. Rupp, U. Hellwig, U. Schellhorn, J. Kohlbrecher, A. Wiedenmann, and J. Woods, *Polymer* **38**, 2299 (1997).
- [21] N. Suzuki and Y. Tomita, *Appl. Opt.* **43**, 2125 (2004).
- [22] N. Suzuki and Y. Tomita, *Appl. Opt.* **46**, 6809 (2007).
- [23] T.K. Gaylord and M.G. Moharam, *Appl. Opt.* **20**, 3271 (1981).
- [24] R. Extermann and G. Wannier, *Helv. Phys. Acta* **IX**, 520 (1936).
- [25] H. Kogelnik, *Bell Syst. Tech. J.* **48**, 2909 (1969).
- [26] C.V. Raman and N.S.N. Nath, *Proc. Indian Acad. Sci. A* **2**, 406 (1935).
- [27] M.G. Moharam and T.K. Gaylord, *J. Opt. Soc. Am.* **71**, 811 (1981).
- [28] V. Somenkov, S. Shilstein, N. Belova, and K. Utemisov, *Solid State Commun.* **25**, 593 (1978).
- [29] C.G. Shull, *Phys. Rev. Lett.* **21**, 1585 (1968).