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New developments in the production of theranostic pairs of radionuclides

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13 Abstract

10

14 A brief historical background of the development of the theranostic approach in nuclear medicine is given and seven theranostic pairs of radionuclides, namely ^{44g}Sc/⁴⁷Sc, 15 ⁶⁴Cu/⁶⁷Cu, ⁸³Sr/⁹⁰Sr, ⁸⁶Y/⁹⁰Y, ¹²⁴I/¹³¹I, ¹⁵²Tb/¹⁶¹Tb and ¹⁵²Tb/¹⁴⁹Tb, are considered. The first 16 six pairs consist of a positron and a β^- -emitter whereas the seventh pair consists of a 17 18 positron and an α -particle emitter. The decay properties of all those radionuclides are briefly mentioned and their production methodologies are discussed. The positron emitters 19 64 Cu, 86 Y and 124 I are commonly produced in sufficient quantities via the (p,n) reaction on 20 the respective highly enriched target isotope. A clinical scale production of the positron 21 22 emitter ^{44g}Sc has been achieved via the generator route as well as via the (p,n) reaction, but further development work is necessary. The positron emitters ⁸³Sr and ¹⁵²Tb are under 23 development. Among the therapeutic radionuclides, ⁸⁹Sr, ⁹⁰Y and ¹³¹I are commercially 24 25 available and ¹⁶¹Tb can also be produced in sufficient quantity at a nuclear reactor. Great efforts are presently underway to produce ⁴⁷Sc and ⁶⁷Cu via neutron, photon and charged 26 particle induced reactions. The radionuclide ¹⁴⁹Tb is unique because it is an α -particle 27 emitter. The present method of production of ¹⁵²Tb and ¹⁴⁹Tb involves the use of the 28 spallation process in combination with an on-line mass separator. The role of some 29 30 emerging irradiation facilities in the production of special radionuclides is discussed.

31 Keywords

32 Theranostic pair of radionuclides. Decay data. Cross section and excitation function.

33 Production methodology. Yield and purity. Specific activity.

34 **1. Introduction**

35 Radioactivity is unique in the sense that it can be routinely used in nuclear medicine both 36 for diagnosis and therapy [1]. Each application, however, demands a special type of 37 radionuclide, the choice being dependent on its decay properties. Thus, γ -ray emitters like 99m Tc (T_{1/2} = 6.0 h), 123 I (T_{1/2} = 13.2 h) and 201 Tl (T_{1/2} = 3.06 d), and positron emitters, like 38 ¹¹C ($T_{\frac{1}{2}} = 20.4 \text{ min}$), ¹⁸F ($T_{\frac{1}{2}} = 109.6 \text{ min}$) and ⁶⁸Ga ($T_{\frac{1}{2}} = 1.13 \text{ h}$) are commonly used in 39 40 diagnostic studies utilizing Single Photon Emission Computed Tomography (SPECT) or 41 Positron Emission Tomography (PET), respectively. As regards internal radionuclide 42 therapy (endoradiotherapy), in general, radionuclides emitting low-range highly ionizing 43 radiation, i.e., α - or β -particles, conversion and/or Auger electrons, are of great interest. 44 The major problem in internal radiotherapy, however, has been the quantification of 45 radiation dose caused to various organs, mainly due to uncertainties in the measurement of 46 radioactivity from outside the body of the patient. Although in the case of a few therapeutic radionuclides, e. g., 131 I (T_{1/2} = 8.02 d) and 188 Re (T_{1/2} = 17.0 h), γ -scanning or SPECT has 47 48 been used to determine the radioactivity distribution in the body, the methodology lacks 49 precision. The uncertainty in radioactivity distribution is still higher for radionuclides decaying by pure β^- -emission, e.g., ${}^{32}P(T_{\frac{1}{2}} = 14.3 \text{ d}), {}^{89}Sr(T_{\frac{1}{2}} = 50.5 \text{ d})$ and ${}^{90}Y(T_{\frac{1}{2}} = 2.7 \text{ d})$ 50 51 d), because imaging is usually done through the use of bremsstrahlung.

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53 In the early 1990s, thoughts started developing in several laboratories to use a SPECT radionuclide as a surrogate of a therapeutic radionuclide [2], e.g., ¹¹¹In ($T_{\frac{1}{2}} = 2.8$ 54 d), a trivalent metal, as a surrogate of ⁹⁰Y, another trivalent metal. There has also been 55 56 discussion about the use of several other metallic radionuclides [3]. However, none of those 57 approaches provided patient-individual quantitative data on radiation doses. In 1992, a few 58 researchers at the Forschungszentrum Jülich, Germany, came to the idea of combining PET 59 and endoradiotherapy by using a pair of radionuclides of the same element, one emitting positrons and the other β^- -particles. The choice fell on the pair ${}^{86}Y/{}^{90}Y$. To this end, the 60

61 β^+ -emitting radionuclide ⁸⁶Y (T_{1/2} = 14.7 h) was developed and produced in sufficient 62 quantity [4, 5] and it was applied together with the β^- -emitting radionuclide ⁹⁰Y (T_{1/2} = 2.7 63 d) in a tumour patient study [6]. That investigation is regarded today as the beginning of 64 the theranostic concept. The development of this concept has been recently described in 65 detail [7].

66

67 By administering to a specific patient a positron-emitting radioisotope of an 68 element together with a therapeutic radioisotope of the same element (which emits β^{-} or 69 α -particles, or low-energy Auger/conversion electrons), it is possible to measure the uptake 70 kinetics in an organ of the patient via PET imaging, thereby allowing an accurate 71 dosimetric calculation, which leads to quantification of therapy. This concept is now called 72 "theranostic approach" and it is finding increasing application. The methodology of using 73 "matched-pair" of radionuclides in patient care studies is known as "personalized 74 medicine".

75

76 There are several suitable or potentially suitable theranostic pairs of radionuclides, e. g. ^{44g}Sc/⁴⁷Sc; ⁶⁴Cu/⁶⁷Cu; ⁶⁸Ga/⁶⁷Ga, ⁷²As/⁷⁷As; ⁸³Sr/⁸⁹Sr; ⁸⁶Y/⁹⁰Y; ^{110g}In/¹¹¹In; ¹²⁴I/¹³¹I; 77 ¹⁵²Tb/¹⁶¹Tb and ¹⁵²Tb/¹⁴⁹Tb. Some of them have already found application in clinical 78 79 research while the others are being developed. In recent years there is also an increasing 80 tendency to handle only one radionuclide as a theranostic agent, especially if it is readily available. One example is 177g Lu. The dosimetry is based on γ -ray spectrometry or SPECT 81 82 and the therapy effect is well known. However, in comparison to the PET technique, 83 SPECT is not quantitative, though in recent years high-quality SPECT systems have been 84 developed.

85

In this review we discuss seven rather established pairs of radionuclides where a combination of PET and internal radiotherapy is involved. Their production methods are described and the prospects of their availability on a clinical scale are considered.

89

2. Choice of radionuclides: decay data

The decay properties of the seven pairs of radionuclides under consideration in this review are given in **Table 1**. The major decay data were taken from refs. [8-10] and they represent the commonly accepted values. Only in a few individual cases, e.g., ⁶⁴Cu and ¹²⁴I, own recently measured data [11] are given. The positron emission intensities for ⁸³Sr, ⁸⁶Y and ¹⁵²Tb are rather uncertain.

95 The positron endpoint energy and the associated γ -rays play important roles in PET 96 measurements. Whereas a high positron endpoint energy affects the resolution of a scan, 97 the γ -rays present in the vicinity of the annihilation radiation may altogether distort the image. From this point of view the positron emitter ⁸⁶Y is far from ideal, but it could be 98 99 used after many scattering corrections [12, 13]. There is some problem with ¹²⁴I as well, 100 but the corrections needed are much smaller [12-14]. Somewhat similar result was obtained with ^{44g}Sc [15]. The positron emitter ⁶⁴Cu is almost ideal for PET imaging because of its 101 102 low positron endpoint energy and almost no emitted γ -ray, the abundance of the 1346 keV 103 γ -ray being negligibly low. It has been therefore extensively used in PET studies related to 104 radioimmunotherapy. As far as the other two β^+ -emitters are concerned (i.e. ⁸³Sr and ¹⁵²Tb), very few PET measurements have been reported. The radionuclide ⁸³Sr appears to 105 be promising because its positron endpoint energy is comparable to that of ^{44g}Sc. The 106 radionuclide ¹⁵²Tb has somewhat higher positron endpoint energy but since the associated 107 108 γ -rays are not too many, it has been used in PET measurements after applying scattering corrections similar to those in the case of ¹²⁴I. As regards therapeutic radionuclides, ⁸⁹Sr 109 and 90 Y are pure β^- -emitters. The radionuclide 149 Tb is an exotic α -emitter. The 110 radionuclides 47 Sc, 67 Cu, 131 I and 161 Tb emit β^- -particles with relatively low endpoint 111 energies and a few associated γ -rays. 112

3. Production methodologies

The development of production methodology of a novel radionuclide involves work in several directions, e.g., nuclear data, irradiation technology, chemical separation and quality control of the product. We consider several of those aspects below for each individual radionuclide. For a few radionuclides, some production details were recently

	l	8 ⁺ -emitting 1	adionucli	de		Therapeutic radionuclide					
Radio- nuclide	T ½	Mode of decay	E ^{β+} (max) (keV)	Main γ-rays		Radio- nuclide	T ¹ / ₂	Mode of decay	Corpuscular radiation	Main	γ -rays
		(%)		Energy (keV)	Intensity (%)	-		(%)	E _{max} (keV)	Energy (keV)	Intensity (%)
⁴⁴ Sc	3.9 h	EC (5.7) β ⁺ (94.3)	1470	1157.0	99.9	⁴⁷ Sc	3.35 d	β (100)	610	159.4	68
⁶⁴ Cu ^{b)}	12.7 h	EC (43.8) β^+ (17.8) β^- (38.4)	653 571	1346.0	0.53	⁶⁷ Cu	2.58 d	β (100)	577	184.6	48.6
⁸³ Sr	32.4 h	EC (74) β^+ (26)	1274	762.7 381.6	30.0 19.6	⁸⁹ Sr	50.5 d	β (100)	1470		
⁸⁶ Y	14.7 h	EC (67) β ⁺ (33)	2335	627.8 1076.7 1153.2	32.6 82.5 30.5	⁹⁰ Y ^{c)}	2.7 d	β (100)	2290		
¹²⁴ I ^{b)}	4.18 d	EC (78) β ⁺ (22)	2137	602.7 722.8	61 10	¹³¹ I	8.02 d	β (100)	607	364.5 637.0	82 7.3
¹⁵² Tb	17.5 h	EC (82) β ⁺ (18)	2500	344.3	57	¹⁶¹ Tb ¹⁴⁹ Tb	6.9 d 4.1 h	β (100) α (16.7) β^+ (4.3) EC (79)	590 α: 5830 600	74.6 165.0 352.2	9.8 27.8 33.0

Table 1. Major decay data^{a)} of the theranostic pairs of radionuclides

^{a)} Data taken from Refs. [8-10], unless otherwise stated.
^{b)} Decay data based partly on own measurement [11].
^{c)} Obtained generally from a generator system.

reported [16, 17]. For those radionuclides, therefore, the present review gives only some updated information.

124 **3.1 Theranostic pair** ^{44g}Sc/⁴⁷Sc

125 The trivalent element scandium forms very useful metal complexes with many oxygen-126 containing bifunctional chelators. This pair of radionuclides is therefore of great potential 127 value in theranostic investigations. Although the positron emitter ⁴³Sc ($T_{\frac{1}{2}} = 3.9$ h) is also 128 very interesting and is presently attracting considerable attention, we limit our discussion 129 to ^{44g}Sc because it has been more thoroughly investigated.

130

131 **Production of** ⁴⁴g**Sc**

For the production of the positron emitter ⁴⁴gSc in no-carrier-added form, two routes have
been investigated:

134

a)
$${}^{45}Sc(p,2n){}^{44}Ti (60.4 a) \xrightarrow{EC} {}^{44g}Sc$$
 generator system

- 135 b) Direct production of 44g Sc.
- 136

The first route involves the production of the long-lived parent ⁴⁴Ti at an intermediate 137 138 energy accelerator. The cross sections of the ${}^{45}Sc(p,2n){}^{44}Ti$ nuclear reaction have been well investigated [18, 19] and the energy range $E_p = 35 \rightarrow 15$ MeV appears to be very suitable 139 for production purposes. The calculated thick target yield of ⁴⁴Ti over this energy range 140 amounts to ~ 4 kBg μ A⁻¹ h⁻¹ (for 1 h irradiation). Due to the long half-life of ⁴⁴Ti, its 141 production is a rather difficult proposition. Although it was proposed a long time ago [20], 142 143 hitherto only a 185 MBq generator has been reported [21] and some post-elution purification of ^{44g}Sc has been described [22]. In recent years, more effort has been devoted 144 to the separation of the parent ⁴⁴Ti via anion-exchange chromatography [23] and the 145 daughter ^{44g}Sc through cation-exchange chromatography [24]. The generator activity, 146 however, has still been limited to about 175 MBq. The separated ^{44g}Sc is free of ^{44m}Sc (T_{1/2} 147 148 = 2.44 d).

150 The second route of production of ^{44g}Sc entails the utilization of either the 44 Ca(p,n) 44g Sc or the 44 Ca(d,2n) 44g Sc reaction. The excitation functions of those reactions 151 have been measured [25-30]. A third reaction, namely ${}^{41}K(\alpha,n){}^{44g}Sc$, is also possible. Its 152 cross sections have also been measured [26, 31, 32]. The thick target yields of ^{44g}Sc 153 154 calculated from the excitation functions are given in **Fig. 1**. The data for the (p,n) reaction were taken from refs. [25, 26, 28] whereby the Levkovskii data [26] were reduced by a 155 156 factor of 0.82 [33]. The cross section data adopted for the (d,2n) reaction were from [30] 157 and those for the (α, n) reaction from refs. [26, 31, 32]. Evidently, the yield from the (p, n)158 reaction is higher than that from the (d,2n) reaction up to about 30 MeV; thereafter the 159 (d,2n) reaction appears to give a higher yield. The yield from the (α,n) process is much 160 lower. In each case a highly enriched target is necessary to achieve clinically relevant yields of ^{44g}Sc. 161

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163 Several groups measured cross sections of a large number of charged particle induced 164 reactions in which ${}^{44g}Sc$ was formed as a subsidiary product. Furthermore, a few groups 165 investigated the production of ${}^{44g}Sc$ (together with other Sc isotopes) using ^{nat}Ca as the 166 target material [cf. 34, 35]. The formation of ${}^{44g}Sc$ as a side product was also investigated 167 in studies primarily done on the formation of ${}^{43}Sc$ in α -particle induced reactions on ${}^{nat}K$ 168 and ${}^{nat,44}Ca$ [36-38]. All those studies are helpful in optimizing the production of ${}^{44g}Sc$.

169

For clinical scale production of ^{44g}Sc, targets consisting of ⁴⁴CaO (enrichment 95%) 170 171 and ${}^{44}CaCO_3$ (enrichment > 99%) have been used [27, 30, 39, 40]. Irradiations were done with protons ($E_p = 11 \rightarrow 5 \text{ MeV}$) [27, 40] or deuterons ($E_d = 16 \rightarrow 10 \text{ MeV}$) [30, 41] at beam 172 173 currents of up to 50 μ A and 2 μ A, respectively. The separation of ^{44g}Sc and the recovery of 174 the target material were achieved through ion-exchange chromatography. By using the 175 (d,2n) reaction, a batch yield of about 50 MBq of ^{44g}Sc was achieved [41] but it could be 176 increased by increasing the beam current. In the case of the (p,n) reaction, on the other 177 hand, a batch yield of up to 2 GBq ^{44g}Sc has been reported [40]. The product is of high 178 radiochemical purity and can be used immediately for preparing radiometal complexes. 179 The only drawback of the direct method of production of ^{44g}Sc is the associated longer lived metastable state 44m Sc (T_{1/2} = 2.44 d), amounting to < 1% and ~ 2.5% in the (p,n) and 180

181 (d,2n) reactions, respectively [30]. On the other hand, this drawback is positively used in some laboratories to prepare a so-called "in-vivo generator" [41]. The longer lived ^{44m}Sc 182 decays 100% by isomeric transition to ^{44g}Sc which can be measured via PET. Since the 183 spin of the 44m Sc isomer is relatively high (6⁺) as compared to that of 44g Sc (2⁺), it was 184 predicted [42] that an α -particle induced reaction would lead to a higher yield of ^{44m}Sc. 185 This has been experimentally observed in the ${}^{42}Ca(\alpha,d){}^{44m,g}Sc$ process [38]. The ratio of 186 187 ^{44m}Sc to ^{44g}Sc increased to about 11% at E_{α} = 29 MeV. On the other hand, the thick target yields of both ^{44m}Sc and ^{44g}Sc in the α -particle induced reaction [38] are much lower than 188 those in the (p,n) and (d,2n) reactions discussed above. 189

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In summary, both the direct and indirect methods of production of ^{44g}Sc are interesting, but further development work is needed. A new aspect with regard to the direct production is the development of a solution target for use at a medical cyclotron. By irradiating a solution of ^{nat}Ca(NO₃₎₂ with 13 MeV protons, ^{44g}Sc was produced in quantities up to 28 MBq, sufficient for local radiochemical and possibly animal studies [43].

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197 **Production of** ⁴⁷**Sc**

198 The production methods for the β^- -emitting therapeutic radionuclide ⁴⁷Sc in no-carrier-199 added form have been under investigation for more than 40 years but in recent years, with 200 the developing concept of theranostic application, the efforts have been intensified. Since 201 in most cases Ti is used as a target material, a large number of radiochemical separation 202 methods for no-carrier-added ⁴⁷Sc from products formed in the interaction of Ti with 203 neutrons, photons and charged particles have been developed [cf. 44-54]. Good summaries of those methods have been given [49, 50]. Similarly, separation methods of ⁴⁷Sc from an 204 205 irradiated Ca target have also been described [55-58].

A summary of the routes used to date for the production of 47 Sc is given in **Table 2**. An old but very successful method has been the 47 Ti(n,p) 47 Sc reaction [45-52, 59-61]. The cross section averaged for the fission neutron spectrum (σ_{FS}) amounts to 20 ± 2 mb [62]. By irradiating 200 mg of 94.5% enriched 47 TiO₂ target in a high flux nuclear reactor for

about 3.6 days it was possible to obtain a batch yield of 1.6 GBq of ⁴⁷Sc of high

211	Table 2. Routes for production of ⁴⁷ Sc
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Nuclear process	Target (enrichment)	Cross section or projectile energy	Production related work	Separation yield (%)	Purity (%)	Batch yield GBq [Ref.]	Other references
⁴⁷ Ti(n,p) ⁴⁷ Sc	^{nat} TiO ₂ ; ⁴⁷ TiO ₂ (94.5 %)	$\sigma_{FS}: 20 \pm 2 \text{ mb*}$	Irradiation in a high-flux reactor; chemical processing	> 97	> 99.5	1.6 [49]	[44-48, 50-52] [59]
⁴⁸ Ti(γ,p) ⁴⁷ Sc	⁴⁸ TiO ₂ (99.1 %) ^{nat} TiO ₂	Photons: 60 MeV Photons: 40 MeV	Irradiation in photon field; chemical processing	> 90	> 95	11×10^{-3} [54] (for 100 mg target) 186×10^{-3} [54] (for 3 g target)	
	⁴⁸ TiO ₂ (96.2 %)	Photons: 40 MeV	Simulation; benchmarking				[63]
$\stackrel{^{46}\text{Ca}(n,\gamma)^{47}\text{Ca}}{\xrightarrow{\beta^-} {}^{47}\text{Sc}}$	⁴⁶ Ca(NO ₃) ₂ (31.7 %)	$\begin{split} \sigma_{th} &: 0.7 \pm 0.2 \; b^{\dagger} \\ I_o &: 0.32 \pm 0.12 \; b^{\dagger} \end{split}$	Irradiation in a high-flux reactor; chemical processing	> 80	> 99	0.6 [58] (for 1 mg target)	[55, 57, 60]
$\stackrel{^{48}\text{Ca}(\gamma,n)^{47}\text{Ca}}{\xrightarrow{\beta^-}} \stackrel{^{47}\text{Sc}}{\xrightarrow{47}\text{Sc}}$	^{nat} Ca	Photons: 40 MeV	Simulation; benchmarking; yield measurement				[64, 65]
⁴⁸ Ti(p,2p) ⁴⁷ Sc	⁴⁸ TiO ₂ (98.5 %)	48 < E _p < 150 MeV	High-current proton irradiation; chemical processing	> 90	Not accept- able	< 1 [48]	[49, 60, 61]

* Value from A. Calamand, IAEA Technical Report-156 (1974) 273; (σ_{FS} is fission neutron spectrum averaged cross section). [†] Value from S.F. Mughabghab and D.I. Garber, BNL-325 (1973) 20-6; (σ_{th} is thermal cross section; I₀ is resonance integral). 212

radionuclidic and chemical purity [49]. Higher yields are possible, if thicker targets would be used. Other groups used $^{nat}TiO_2$ as target material and the neutron flux was not very high, so the resulting yield of ^{47}Sc was lower.

217

Another old method is the ${}^{48}\text{Ti}(\gamma,p){}^{47}\text{Sc}$ reaction using high-energy photons [53]. In recent years investigations on the formation of a few therapeutic radionuclides using highly powerful accelerators (which deliver high-intensity, high-energy photons) have been intensified. In a most recent work at the Argonne National Laboratory [54] a batch yield of 187 MBq of ${}^{47}\text{Sc}$ has been achieved by using photons generated by an electron beam of 40 MeV (incident on a convertor) at a maximum power of about 3 kW. Further studies to increase the yields are in progress in several laboratories [cf. 63].

225

A third method of ⁴⁷Sc production utilizes the decay of ⁴⁷Ca ($T_{\frac{1}{2}} = 4.54$ d). The nuclear 226 process generally used is ${}^{46}Ca(n,\gamma){}^{47}Ca \xrightarrow{\beta^-} {}^{47}Sc$ [55, 57, 58, 60]. The method has two 227 limitations: a) the abundance of 46 Ca in nat Ca is only 0.004%, so that an enriched target is 228 229 absolutely necessary, which is very expensive, b) the cross section of the (n, γ) reaction is 230 not high (see **Table 2**). Nonetheless, the methodology has been recently well developed by using a 31.7% enriched 46 Ca(NO₃)₂ target and irradiating it at the neutron high flux reactor 231 in Grenoble. The ⁴⁷Sc activity was separated from calcium by column chromatography, 232 similar to the method developed for the separation of ⁴⁴Sc from a ⁴⁴Ca target (see above). 233 From a 1 mg ⁴⁶Ca target, a batch yield of 600 MBq of ⁴⁷Sc was obtained. A higher yield 234 could be achieved by increasing the amount of the target material. Besides the neutron 235 activation of ⁴⁶Ca, the production of ⁴⁷Ca is also being investigated via the ⁴⁸Ca(γ ,n) –route 236 237 [64, 65], especially in view of the increasing potential of high power electron linear 238 accelerators. Irradiations were done with photons obtained from a 40 MeV, 1 kW beam of electrons on a convertor, and the radioactivity of the product ⁴⁷Ca was assayed. Further 239 240 simulation, benchmarking and separation studies are continuing.

241

The production of ⁴⁷Sc has been attempted using charged particles as well, particularly via intermediate energy protons on ^{nat}Ti using the accelerator BLIP at Brookhaven National Laboratory [48, 49, 61]. The ⁴⁷Sc yields determined over the energy region $48 \le E_p \le 150$

MeV were on the order of a few GBq. The level of other Sc isotopes, especially ⁴⁶Sc, 245 246 however, was rather high. More recent studies in a few other laboratories are concentrating 247 on optimization of the energy range for production of this radionuclide. Two other methods investigated for the production of ⁴⁷Sc at the research level consist of the reactions 248 $^{44}Ca(\alpha,p)^{47}Sc$ and $^{48}Ca(p,2n)^{47}Sc$. In the former case, using a 97.0% enriched $^{44}CaCO_3$ 249 target [37] high-purity ⁴⁷Sc was obtained in low yield which was, however, sufficient for a 250 251 preclinical study. In the latter case [66], only the (p,2n) reaction cross section was 252 measured.

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Thus, in summary, considerable effort is presently being devoted to obtain high-quality 4⁷Sc in quantities sufficient for medical applications. In particular the photon induced reactions are receiving great attention.

257 **3.2 Theranostic pair ⁶⁴Cu/⁶⁷Cu**

The element copper has a versatile co-ordination chemistry. In the no-carrier-added form copper radioisotopes are able to bind with biologically relevant small molecules as well as with some antibodies and proteins. It is thus very suitable for preparing metal-chelates for medical use [67, 68]. Two positron emitters of copper, namely ⁶¹Cu ($T_{1/2} = 3.4$ h) and ⁶⁴Cu ($T_{1/2} = 12.7$ h), have been used in PET studies. For theranostic applications, however, the radionuclide ⁶⁴Cu appears to be more suitable because of its longer half-life. We therefore concentrated on this radionuclide.

265

266 **Production of ⁶⁴Cu**

Several routes have been investigated for the production of no-carrier-added ⁶⁴Cu. The oldest among them is the ⁶⁴Zn(n,p)⁶⁴Cu reaction in a nuclear reactor (for a brief summary see [69–71]). The fission neutron spectrum averaged cross section (σ_{FS}) amounts to 31 ± 2.3 mb [62] and sufficient quantities of ⁶⁴Cu could be produced in a medium to high-flux reactor. The purity of the product achieved, however, did not meet the stringent demands for medical applications. In recent years some further efforts have been made to produce better quality ⁶⁴Cu via the above reaction in a nuclear reactor [70, 71], in particular by using

99.4% enriched ⁶⁴ZnO as target material in a thermal neutron shielded sample holder and
efficient separation methods for radiocopper [71]. Furthermore, accelerator produced
neutrons have also been used, e. g. d(Be) break up neutrons [72] or 14 MeV neutrons [73].
In the latter two cases the (n,p) reaction cross section is higher. However, due to low
neutron fluxes the yield of ⁶⁴Cu was low.

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The emphasis regarding the production of ⁶⁴Cu got shifted over the last several years 280 281 from a reactor to a cyclotron. Proton and deuteron induced reactions on several target isotopes, especially the reactions ${}^{64}Ni(p,n){}^{64}Cu$, ${}^{64}Ni(d,2n){}^{64}Cu$, ${}^{68}Zn(p,\alpha n){}^{64}Cu$, 282 66 Zn(p,2pn) 64 Cu, 64 Zn(d,2p) 64 Cu and 66 Zn(d, α) 64 Cu were investigated till 2009 over a wide 283 284 energy range of up to 80 MeV using highly enriched target isotopes, with the aim of obtaining data for the production of ⁶⁴Cu. Based on a critical analysis of the published 285 286 nuclear reaction cross section data, Aslam et al. [74] presented a comparison of the various production reactions of ⁶⁴Cu and came to the conclusion that the ⁶⁴Ni(p,n)⁶⁴Cu reaction 287 288 over the energy range of $E_p = 12 \rightarrow 8$ MeV would be the best choice. The calculated thick target yield amounts to 304 MBq μ A⁻¹ h⁻¹ (for 1h irradiation) and no radionuclidic impurity 289 occurs. In recent years some further measurements near the threshold of the ${}^{64}Ni(p,n){}^{64}Cu$ 290 reaction have been carried out [75] and the reaction ${}^{67}Zn(p,\alpha){}^{64}Cu$ has also been studied 291 [76]. Furthermore, in connection with the specific activity of ⁶⁴Cu, the formation of non-292 radioactive copper during the production of ⁶⁴Cu via proton and deuteron-induced reactions 293 on enriched ⁶⁴Ni has also been considered [77]. The nuclear process ⁶⁴Ni(p,n)⁶⁴Cu, 294 295 developed at the Forschungszentrum Jülich [78], has now become the standard procedure for the production of ⁶⁴Cu. The major features were the preparation of a target via 296 electrodeposition of ⁶⁴Ni on a Au backing, a clean separation of ⁶⁴Cu via ion-exchange 297 298 chromatography, and an efficient recovery of the enriched target material. The technology 299 was further developed in some laboratories [79-81] and batch yields of up to 40 GBq of ⁶⁴Cu were achieved. Several other optimization studies have also been performed [82-87]. 300 301 Many small hospital-based laboratories are now producing this radionuclide in amounts 302 sufficient for local use. A few newer developments are related to more efficient chemical 303 separation and purification of ⁶⁴Cu [88-91]. There has been some emphasis on automation 304 of the production procedure as well [92-96]. Thus, considerable interest has been aroused

in recent years in the production of 64 Cu via this route. Due to the increasing demand for this radionuclide, on one hand solution targets similar to those for 44g Sc mentioned above are being developed [97] and, on the other, a commercialization of the process is being pursued. However, it should be mentioned that small amounts of 64 Cu have also been produced via the nuclear processes 64 Zn(d,2p) 64 Cu [98, 99] and 68 Zn(p, α n) 64 Cu [100-103], the latter partly as a by-product in the production of 67 Ga via the 68 Zn(p,2n) 67 Ga reaction.

311

312 **Production of** ⁶⁷Cu

The production of the therapeutic radionuclide 67 Cu (T_{1/2} = 2.58 d) in no-carrier-added form has also been under consideration for more than 40 years and the knowledge available till 2011 was critically reviewed [104]. A few other later reviews dealt with the newer information [17, 105-107]. In this work therefore only some salient features are mentioned.

Similar to ⁶⁴Cu, the production of ⁶⁷Cu in neutron induced reactions, especially in a 318 nuclear reactor via the 67 Zn(n,p) 67 Cu reaction ($\sigma_{FS} = 1.07 \pm 0.04$ mb) has received some 319 new attention [69, 71], in particular by using 93% enriched ⁶⁷ZnO as target material [71]. 320 321 The same threshold reaction has also been investigated with 14 MeV neutrons; however, by using a ^{nat}ZnO target [73]. A yet another method making use of the 68 Zn(n,np) 67 Cu 322 323 reaction induced by fast neutrons, generated by breakup of 40 MeV deuterons on a graphite 324 target, has also been utilized [108]. In those two works [73, 108] the fundamental separation and purification procedures were established. The ⁶⁷Cu obtained via the latter 325 326 process using a 99.29% enriched ⁶⁸ZnO target was shown to be suitable for preclinical 327 studies [109]. For large scale production, however, further development work using high 328 neutron fluxes is needed.

329

Another reaction which has been under investigation for a long time is the $^{68}Zn(\gamma,p)^{67}Cu$ process. In one early study ^{nat}Zn was used as target material [110] and in another 98.97% enriched ⁶⁸ZnO was employed [111]. In both cases chemical separation of the product ⁶⁷Cu was carried out. The batch yield achieved was up to 185 MBq but the chemical purity would not meet the standard required today. With the increasing significance of ⁶⁷Cu combined with the development of powerful electron accelerators, in

recent years the efforts to utilize the 68 Zn(γ ,p) 67 Cu reaction for 67 Cu production have been 336 intensified [64, 112-115]. Production yields of ⁶⁷Cu have been measured experimentally 337 338 and compared with theoretically calculated values [112, 113], extensive purification 339 methodology was developed [114], simulation studies were performed and predicted activities were verified with experimental data [64, 115]. The yield of ⁶⁷Cu achieved 340 amounts to about 1 MBq g⁻¹ kW⁻¹ h⁻¹. Thus, tens of MBq of ⁶⁷Cu can easily be produced. 341 342 It is expected that with further intensification of technological efforts to develop highintensity accelerators (possibly up to 100 kW power), it should be possible to produce ⁶⁷Cu 343 344 in GBg quantities.

345

346 In addition to the neutron and photon induced reactions described above for the production of ⁶⁷Cu, considerable effort has been invested over the years to make use of 347 348 charged-particle induced reactions as well. The four nuclear processes investigated are 349 listed in **Table 3**. The suitable energy ranges and the calculated thick target yields are based 350 on evaluated excitation functions [116] and a few other measurements. However, it should be mentioned that a new measurement on the ${}^{68}Zn(p,2p){}^{67}Cu$ reaction [117] gives cross 351 352 section values which are lower than the evaluated data up to 60 MeV by about 10%. If those values are accepted, the calculated yield of ⁶⁷Cu would decrease slightly. The yield 353 values for the 70 Zn(d,an) 67 Cu and 64 Ni(a,p) 67 Cu reactions given in **Table 3** were derived 354 355 from individual experimental cross section curves, for the former reaction from ref. [118] 356 and for the latter from refs. [119,120].

357

As far as the practical production of 67 Cu is concerned, in the case of the 70 Zn(p, α) 67 Cu reaction two studies were performed, one using a 99.7% enriched 70 ZnO target [121] and the other using a 70% enriched 70 Zn electroplated target [122]. The separation yields were comparable but, as understandable, the radionuclidic purity of 67 Cu achieved was higher in the first study due to the higher enrichment of the target. The batch yield of 67 Cu obtained via this production route was, however, quite low. With

Nuclear reaction	Energy range (MeV)	Calculated thick target yield (MBq/µAh)	Target (enrichment)	Production related work	Separation yield (%)	Radionuclidic purity (%)	Batch yield MBq [Ref.]
⁷⁰ Zn(p,α) ⁶⁷ Cu	$18 \rightarrow 12$	2.2	⁷⁰ ZnO (99.7 %)	Irradiation at 4 µA; anion- exchange separation	> 80	> 99	0.8 [121] for 10 mg target
			⁷⁰ Zn electroplated (70 %)	Irradiation at 20 µA; solvent extraction and anion-exchange separation	> 80	> 85	14 [122]
⁷⁰ Zn(d,nα) ⁶⁷ Cu	$20 \rightarrow 10$	4.2	⁷⁰ Zn metal (95.35 %)	Low current irradiation of thin target; consective cation- and anion-exchange separation	> 90	> 90	0.95 [118]
⁶⁸ Zn(p,2p) ⁶⁷ Cu	$70 \rightarrow 30$	30	⁶⁸ ZnO (99.0 %)	Irradiation at 3 µA; ion- exchange chromatography	83	> 97	117 [127]
			⁶⁸ ZnO (99.7 %)	Irradiation at 100 µA; extensive chemical processing	> 92	mixture of ⁶⁴ Cu and ⁶⁷ Cu ^{a)}	1.6×10^{3} [128]
⁶⁴ Ni(α,p) ⁶⁷ Cu	$35 \rightarrow 10$	0.8	⁶⁴ Ni electroplated (99.07 %)	Irradiation at 15 μA; cation- exchange separation	> 90	> 75	55 [123]

Table 3. Charged-particle induced nuclear reactions used for the production of ⁶⁷Cu.

^{a)} Using an incident proton beam of 92 MeV.

- 366 regard to the 70 Zn(d, α n) 67 Cu reaction, the production test involved only low current
- 367 irradiation of a very thin target and so the batch yield achieved was very low [118]. There
- 368 is the possibility to produce larger quantities of 67 Cu if thicker targets are used. The reaction
- 369 ${}^{64}\text{Ni}(\alpha,p){}^{67}\text{Cu}$ also leads to a

370 relatively low yield of 67 Cu because of the low cross section and the low range of α -371 particles. Nonetheless, a suitable target was prepared and, after a 7 hour irradiation with 36 372 MeV α -particles at 15 μ A, followed by chemical separation, a total of 55 MBq of 67 Cu was 373 achieved [123]. The product was chemically very pure and was used in preclinical studies 374 [123]. The level of 64 Cu impurity was, however, somewhat high.

375

376 In contrast to the above mentioned three low yield processes, the reaction 68 Zn(p,2p) 67 Cu at intermediate energies leads to a much higher yield. It has therefore been 377 378 receiving more attention. It was originally utilized for production of ⁶⁷Cu by irradiation 379 with protons of energies about 180 MeV followed by chemical separation [48, 61, 124]. 380 The yield was very high but the specific activity was low. Later investigations concentrated more over the energy region up to 70 MeV, utilizing highly enriched ⁶⁸Zn as target material 381 382 and extensive chemical processing [125-127]. Further extensive work has recently been 383 reported using about 100 MeV protons [128]. The suggested production energy range is, 384 however, $E_p = 70 \rightarrow 30$ MeV [105]; at higher energies a considerable amount of inactive 65 Cu is formed via the 68 Zn(p,2p2n) 65 Cu reaction which decreases the specific activity of 385 ⁶⁷Cu. Using an incident proton energy of about 92 MeV, batch yields of a few GBq of ⁶⁷Cu 386 have been achieved at BNL. However, the product contains about 5 times more ⁶⁴Cu than 387 ⁶⁷Cu. Thus further optimization work utilizing lower proton energies is needed. A further 388 newer approach is to harvest ⁶⁷Cu from the cooling loop of the Facility for Rare Isotopes 389 390 (FRIB) presently under construction; some preliminary results have been obtained by 391 analysis of a few samples from the aqueous beam stop at the National Superconducting 392 Cyclotron Laboratory (NSCL) [129].

393

From the above discussion it is obvious that the development of production methods of ⁶⁷Cu is of great timely interest because it is one of the most important theranostic radionuclides. Diversified efforts are underway to obtain it in sufficient quantity and good quality for medical applications.

398 **3.3 Theranostic pair ⁸³Sr/⁸⁹Sr**

Strontium is an important bone seeking element. The radionuclides of strontium could therefore be used in diagnostic and therapeutic studies related to bone. The β^- -emitting 89 Sr (T_{1/2} = 50.5 d) is one of the earliest known radionuclides to cure metastases in bone. It also finds application in palliation studies. The β^+ -emitting analogue 83 Sr (T_{1/2} = 32.4 h) should be suitable for theranostic application. As far as we know, to date no PET measurement has been reported using 83 Sr; yet its decay properties suggest that it is potentially suitable.

406

407 **Production of ⁸³Sr**

408 Regarding the production of no-carrier-added ⁸³Sr, excitation functions were measured for 409 the ⁸⁵Rb(p,xn)⁸¹⁻⁸⁵Sr processes up to 100 MeV [130, 131] and ⁸²Kr(³He,xn)^{82,83}Sr reactions up to 36 MeV [132]. Therefrom the suitable energy ranges for the production of ⁸³Sr via 410 411 those two processes were deduced. The calculated thick target yields of the radionuclides 412 formed in the interactions of protons with ⁸⁵Rb are [131] shown in **Fig. 2**. The optimum energy range for the production of ⁸³Sr is $E_p = 37 \rightarrow 30$ MeV, whereby the yield of ⁸³Sr 413 amounts to 160 MBq μ A⁻¹ h⁻¹ (for 1 h irradiation) and the levels of the two long-lived 414 impurities 85 Sr (T¹/₂ = 64.9 d) and 82 Sr (T¹/₂ = 25.3 d) are 0.24% and 0.04%, respectively. 415 A similar analysis for the ³He-particle induced reactions on ⁸²Kr showed that the optimum 416 417 energy range for the production of ⁸³Sr is $E_{3He} = 18 \rightarrow 10$ MeV, whereby the yield of ⁸³Sr amounts to 5.1 MBq μ A⁻¹ h⁻¹ (for 1 h irradiation) and the level of the only impurity ⁸²Sr is 418 0.20%. The method of choice for the production of ⁸³Sr is thus the ⁸⁵Rb(p,3n)-reaction, 419 420 although the availability of 40 MeV protons is often a problem.

421

422 Irradiations of several targets with low beam currents of 40 MeV protons and 18 423 MeV ³He-particles were carried out to measure experimental thick target yields. In the former case, pressed ⁸⁵RbCl pellets absorbing about 5 MeV of the proton beam were used 424 425 and, in the latter, ⁸²Kr gas absorbing about 8 MeV of the ³He-particle energy was irradiated 426 in a special target system [133]. Highly efficient separation methods, using high 427 performance liquid chromatography, were developed to obtain radiostrontium of high quality [131]. The results were compared with the theoretical data. The radionuclide ⁸³Sr 428 429 was obtained in quantities of up to 20 MBq via the (p,3n) process and up to 5 MBq via the

(³He,2n) reaction [131]. A clinical scale production was, however, not demonstrated.
Nevertheless, it should be possible to obtain ⁸³Sr in quantities sufficient for medical
application by using the technology developed for the production of ⁸²Sr (parent of ⁸²Sr/
⁸²Rb generator system), except that the proton energy incident on the ⁸⁵RbCl target should
be 40 MeV instead of 70 MeV used in the ⁸²Sr production.

435

436 **Production of ⁸⁹Sr**

437 As far as the production of the therapeutic radionuclide 89 Sr is concerned, some use has been made of the ⁸⁸Sr(n,γ)⁸⁹Sr reaction. However, due to the very low specific activity, the 438 product ⁸⁹SrCl₂ has been used only in palliative therapy of malignant metastases to the 439 440 skeleton. For preparation of radiopharmaceuticals with high specific activity, a production route involving the neutron threshold reaction 89 Y(n,p) 89 Sr has been developed. The cross 441 442 section averaged for the fission neutron spectrum is low ($\sigma_{FS} = 0.31 \pm 0.06$ mb [62]); therefore long irradiations are needed. The target material consisting of Y2O3 powder, 443 pressed to a pellet, is placed in an Al capsule. The irradiation is done for several weeks at 444 a high fast neutron flux of 1-2 x10¹⁵ n cm⁻² s⁻¹. Thereafter the chemical processing starts 445 by dissolving the irradiated target in HNO₃ and extracting the bulk of yttrium in 446 tributylphosphate. The purification of ⁸⁹Sr is done by incorporating several cation-447 448 exchange chromatographic steps. The finally purified product is then obtained as ⁸⁹SrCl₂ 449 in dilute HCl in a batch yield of about 20 GBq. Large quantities of this radionuclide are 450 produced mainly at the reactor RIAR in Dimitovgrad, Russia [134, 135]. It is then shipped 451 to various parts of the world.

452 **3.4 Theranostic pair ⁸⁶Y/⁹⁰Y**

As mentioned in the introduction, this was the first pair of radionuclides used for
theranostic studies. Its development has been described in detail in a recent publication [7].
In this article therefore only a very brief account is given.

456

457 For the production of the positron emitter ⁸⁶Y ($T_{\frac{1}{2}} = 14.7$ h), the nuclear reactions 458 ⁸⁶Sr(p,n)⁸⁶Y, ⁸⁸Sr(p,3n)⁸⁶Y, ^{nat}Zr(p,x)⁸⁸Y and ^{nat}Rb(³He,xn)⁸⁶Y were investigated (for

references see [136]). Very recently the nuclear process ${}^{89}Y(p,4n){}^{86}Zr \xrightarrow{EC,\beta^+} {}^{86}Y$ has also 459 been reported [137]. The method of choice for production of ⁸⁶Y, however, is the 460 ⁸⁶Sr(p,n)⁸⁶Y reaction on a highly enriched target, originally reported by the Jülich group 461 [5, 6]. Over the optimum energy range of $E_p = 14 \rightarrow 7$ MeV the expected thick target yield 462 463 of ⁸⁶Y amounts to 371 MBq μ A⁻¹ h⁻¹ (for 1 h irradiation). Although an evaluation revealed discrepancy in nuclear data [136], the production technology has been well developed. For 464 irradiation mostly solid 97% enriched ⁸⁶SrCO₃ target is used at a proton beam current of 465 466 about 10 μ A. For the chemical separation of radioyttrium, two methods have been 467 advantageously used:

468 a) Co-precipitation with La(OH)₃, followed by cation-exchange chromatography,

- b) Electrolytic removal of radioyttrium.
- 470

471 A detailed discussion of the separation procedures is given in ref. [7]. Batch yields of a few 472 GBq of 86 Y have been reported. At a few medical cyclotrons, solution targets have been 473 developed to produce small quantities of 86 Y for local use. The radionuclidic purity of 86 Y 474 amounts to > 97%; the major impurity 87m Y originates from the small amount of the isotope 475 87 Sr present in the enriched 86 Sr target. Due to great demand for this radionuclide, efforts 476 are underway to commercialize its production.

477

As regards the production of the β^- -emitter ⁹⁰Y (T_{1/2} = 2.7 d), it could be done via the ⁸⁹Y(n, γ)⁹⁰Y process, but the specific activity is very low. No-carrier-added ⁹⁰Y is therefore generally obtained via the ⁹⁰Sr/⁹⁰Y generator system. The parent activity ⁹⁰Sr (T_{1/2} = 28.6 a) is separated from the fission products and fixed on a generator column. The daughter ⁹⁰Y is eluted about once a week using 2N HCl as eluent. About 3-5 GBq quantities of ⁹⁰Y are collected in 0.5 mL of the eluent. Such generator systems are commercially available.

485 **3.5 Theranostic pair** ¹²⁴I/¹³¹I

486 This is a unique pair of radionuclides. In contrast to the four metallic pairs discussed above, 487 namely ${}^{44g}Sc/{}^{47}Sc$, ${}^{64}Cu/{}^{67}Cu$, ${}^{83}Sr/{}^{89}Sr$ and ${}^{86}Y/{}^{90}Y$, this pair belongs to the group of halogens which form a rather strong covalent bond and have therefore been frequently
applied following the "analogue" approach. A large number of radiopharmaceuticals have
been developed using halogens. Thus, both ¹²⁴I and ¹³¹I find applications both individually
and collectively as a theranostic pair.

492

The therapeutic use of ¹³¹I has been successfully practised for more than 70 years, especially in treatment of thyroid diseases. The use of ¹²⁴I is relatively new. It was first proposed in 1988 by Lambrecht et al. [138]. Since then extensive studies on its production and preparation of radiopharmaceuticals have been performed. Today it is widely used in tumour targeting as well as in thyroid dosimetry.

498

The various methods investigated for the production of 124 I (T_{1/2} = 4.18 d) have been 499 500 extensively reviewed [139]. A critical analysis of the cross section data was performed 501 [140, 141]. A summary of the results was given [106]. It was concluded that the 124 Te(p,n) 124 I reaction, originally suggested by Scholten et al. [142] is the method of choice 502 for the production of ¹²⁴I. For a 99.8% enriched ¹²⁴Te target over the energy range $E_p = 12$ 503 \rightarrow 8 MeV the expected ¹²⁴I yield is 16 MBq μ A⁻¹ h⁻¹ (for 1h irradiation). This yield is not 504 505 very high, but the product obtained is of the highest radionuclidic purity, the level of the associated long-lived ¹²⁵I ($T_{\frac{1}{2}}$ = 60.0 d) impurity being < 0.1%. On the other hand, it is felt 506 that the ¹²⁵Te(p,2n)¹²⁴I reaction [143] over the energy range $E_p = 21 \rightarrow 15$ MeV may also 507 be quite useful; the yield of 124 I is 5 times higher than that via the (p,n) reaction and the 508 level of the ¹²⁵I Impurity is < 1%. Today, for clinical scale production of ¹²⁴I, the 509 ¹²⁴Te(p.n)¹²⁴I reaction is almost universally applied and batch yields of a few GBq are 510 511 obtained. The procedure commonly involves irradiation of a ¹²⁴TeO₂ target and removal of 512 radioiodine by a distillation process at about 750 °C [144-150]. A detailed review of the 513 distillation parameters used by various groups was presented [139]. Radioiodine is 514 generally collected almost quantitatively in 0.3 mL of 0.02 M NaOH solution. Its radiochemical form is checked by high performance liquid chromatography (HPLC); it is 515 516 > 98% iodide which is very suitable for subsequent synthesis steps. The enriched target 517 material is regenerated (without any substantial loss) for reuse.

519 In recent years the separation of radioiodine from α -particle irradiated antimony 520 was also investigated using solvent extraction and ion-chromatographic techniques [151-521 153]. The radionuclidic purity of the product achieved was quite high. However, due to the 522 low batch yield of ¹²⁴I, those methods have not found much practical application.

523

As far as the production of ¹³¹I ($T_{\frac{1}{2}} = 8.02 \text{ d}$) is concerned, the methodology is well established [cf. 154]. It is a reactor radionuclide and is produced either via the fission process (as a subsidiary of ⁹⁹Mo production) or via the route ¹³⁰Te(n, γ)^{131m,g}Te $\stackrel{\beta^-}{\rightarrow}$ ¹³¹I. In the latter case, both dry and wet distillation methods have been used for the separation of radioiodine. Large quantities of ¹³¹I are commercially available.

529 **3.6 Theranostic pairs** ¹⁵²Tb/¹⁶¹Tb and ¹⁵²Tb/¹⁴⁹Tb

530 These two pairs of radionuclides are rather exotic but very promising. In recent years there 531 has been an increasing interest in the application of radiolanthanides in imaging and 532 therapy, especially because a trivalent lanthanide forms stable complexes with many 533 oxygen-containing bifunctional chelators. The imaging is generally done by SPECT which, 534 however, is not quantitative. The radionuclide ¹⁵²Tb ($T_{\frac{1}{2}} = 17.5$ h) is the only suitable β^+ -535 emitter in the region of lanthanides which has been successfully developed for PET 536 measurements. It can thus serve as an exact diagnostic match to the β^{-} -emitting therapeutic 537 radionuclide ¹⁶¹Tb ($T_{\frac{1}{2}} = 6.9 \text{ d}$) as well as to the α -particle emitting therapeutic radionuclide ¹⁴⁹Tb ($T_{\frac{1}{2}}$ = 4.1 h), whose potential in therapy was first suggested by Allen and Blagojevic 538 539 [155]. In fact these three radionuclides together with the Auger electron emitter ¹⁵⁵Tb ($T_{\frac{1}{2}}$ 540 = 5.3 d) make the element terbium very versatile for medical applications, somewhat 541 similar to copper and iodine.

542

543 **Development of ¹⁵²Tb and ¹⁴⁹Tb**

544 Work on the development of the β^+ -emitter ¹⁵²**Tb** and the α -particle emitter ¹⁴⁹**Tb** has been 545 going on for quite some time and two rather uncommon reactions have been investigated 546 for their production.

548 a) *Heavy-ion induced reactions*, first studied in Sydney [156,157]. Using a natural Nd 549 target, ¹⁵²Dy was produced over the energy range of 80 to 110 MeV. The contributing reactions were ${}^{142}Nd({}^{12}C.2n){}^{152}Dy$, ${}^{143}Nd({}^{12}C.3n){}^{152}Dy$, ${}^{144}Nd({}^{12}C.4n){}^{152}Dy$ and 550 ¹⁴⁵Nd(¹²C, 5n)¹⁵²Dy. The product ¹⁵²Dy decays with a half-life of 2.4 h to ¹⁵²Tb. After 551 552 irradiation the thick Nd metal target was therefore allowed to decay for about 12 hours, 553 thereafter it was dissolved in 6 M HNO₃, evaporated to dryness and the residue 554 redissolved in α -hydroxyisobutyric acid (α -HIBA). The separation of no-carrier added ¹⁵²Tb was then achieved through cation-exchange chromatography. The batch yield of 555 ¹⁵²Tb amounted to a few MBq. It was sufficient for tracer studies but not for a PET 556 557 phantom measurement. In the same Nd target irradiated with ${}^{12}C$ ions, the α -particle emitting ¹⁴⁹Tb was formed via the ¹⁴²Nd(¹²C, 5n)¹⁴⁹Dy \rightarrow ¹⁴⁹Tb process. Its batch yield 558 559 amounted to a few MBq [157].

560

561 b) Spallation reaction, first studied at CERN [156]. A tantalum foil was irradiated with 562 1000 MeV protons. The spallation products were released from the target at 2400 °C. 563 The ionized products were separated electromagnetically at the ISOLDE facility. The 564 spallation products of mass number 152 were collected and subjected to a two-step separation procedure, similar to the one used in the separation of 86 Y [5], viz. at first 565 566 coprecipitation of radioterbium with La(OH)₃, then removal of radioterbium from lanthanum by cation-exchange chromatography. The batch yield of ¹⁵²Tb amounted to 567 568 770 MBq [156]. A PET phantom measurement demonstrated the feasibility of using ¹⁵²Tb for monitoring the behavior of therapeutic terbium radionuclides [156]. 569

570

571 Following the successful production of ¹⁵²Tb via the spallation process, several 572 optimization studies and further development work were carried out, in particular with 573 regard to on-line mass separation [158, 159]. To demonstrate the utility of ¹⁵²Tb, a proof 574 of concept study was performed with ¹⁵²Tb-labelled folate in a mouse bearing folate 575 receptor (FR)-positive tumours [158]. A more detailed in vivo imaging study using several 576 other ¹⁵²Tb-labelled compounds showed the potential of this radionuclide for PET studies 577 [159]. Very recently the first application of this positron emitter in human PET/CT has been convincingly demonstrated [160]. The significance of this radionuclide is thusincreasing.

580

Besides the application of the spallation process to the production of ¹⁵²Tb, many 581 582 investigations on other possible production reactions have also been carried out. They deal 583 either with cross section measurements of proton and deuteron induced reactions on 584 gadolinium and dysprosium [161-166] or with chemical separation of radioterbium from 585 gadolinium irradiated with protons [167], europium irradiated with α -particles [168] or lanthanum and cerium irradiated with ¹⁶O-ions [169, 170]. The (p,xn) reactions on 586 587 gadolinium isotopes in the intermediate energy range appear to be promising. An example 588 is given in Fig. 3, which has been adapted from the data of Steyn et al. [162]. The cross section of the ¹⁵⁵Gd(p,4n)¹⁵²Tb reaction is fairly high and over the energy range of E_p = 589 50 \rightarrow 30 MeV, the calculated yield of ¹⁵²Tb amounts to about 1.45 GBq μ A⁻¹ h⁻¹ (for 1 h 590 irradiation). Thus using an enriched ¹⁵⁵Gd target, in principle, it should be possible to 591 592 produce ¹⁵²Tb in quantities sufficient for medical applications.

593

594 With regard to the production of the therapeutic radionuclides of terbium, the case of the α -particle emitter ¹⁴⁹Tb has been mentioned above. Its production in tracer quantities 595 596 via the heavy-ion induced reaction was reported [157]. Subsequently, Beyer et al. [171, 597 172] produced this radionuclide on a clinical scale via spallation of tantalum with 1400 598 MeV protons in conjunction with on-line isotope separation at CERN, and demonstrated direct evidence for single cancer cell killing using ¹⁴⁹Tb-rituximab. In general, however, 599 600 the availability of this radionuclide is rare. On the other hand the cross sections of a few 601 (p,xn) reactions on a few gadolinium isotopes, leading to the formation of ¹⁴⁹Tb, have been 602 described [162]. They appear to be interesting for production purposes but specific 603 production methodology needs to be developed.

604

605 **Production of ¹⁶¹Tb**

606 The production of the β⁻-emitting therapeutic radionuclide ¹⁶¹Tb is usually done in a 607 nuclear reactor via the sequence ¹⁶⁰Gd(n,γ)¹⁶¹Gd $\stackrel{\beta^-}{\rightarrow}$ ¹⁶¹Tb. In general, an enriched ¹⁶⁰Gd 608 target is irradiated with a high neutron flux and separation of ¹⁶¹Tb from the gadolinium target is done by cation-exchange chromatography with α-HIBA, followed by concentration of ¹⁶¹Tb solution [158, 173, 174], There is, however, some difficulty in the production process. The intermediate nuclide ¹⁶¹Gd ($T_{\frac{1}{2}} = 3.7 \text{ min}$) has a very high neutron capture cross section ($\sigma_{\text{th}} \approx 20000 \text{ b}$) so that the formation of ¹⁶¹Tb through the β⁻-decay of ¹⁶¹Gd is in strong competition with the formation of ¹⁶²Gd through the (n,γ) reaction. A short irradiation with a high neutron flux is advantageous. In general, the radionuclide ¹⁶¹Tb could be made available in sufficient quantities.

616 **4. Concluding remarks**

617 The theranostic approach in nuclear medicine, i.e. administering to a specific person two 618 radionuclides of the same element in the same chemical form, one emitting positrons and 619 the other highly-ionizing low-range radiation to cause therapeutic effect, is gaining 620 increasing significance because it constitutes "personalized medicine". In this review seven 621 such pairs have been dealt with and their production methods have been discussed. The positron emitters ⁶⁴Cu, ⁸⁶Y and ¹²⁴I are well characterized and the respective production 622 623 technology using the (p,n) reaction on the respective highly enriched target isotope is well developed. The positron emitter ^{44g}Sc is presently attracting great attention. Though its 624 clinical scale production has been achieved via two routes, namely the ⁴⁴Ti/^{44g}Sc generator 625 626 system and the direct production via the (p,n) reaction, further development work is 627 necessary to ensure its large scale production. The basic methodology for production of the positron emitter ⁸³Sr has also been demonstrated but due to the need of an intermediate 628 629 energy cyclotron, not much progress has been made with regard to its production on a clinical scale. The positron emitter ¹⁵²Tb is potentially very interesting. The production 630 631 methodology developed so far, however, is rather exotic because it makes use of the 632 spallation process in combination with on-line mass separation. Attempts are presently 633 underway to produce it at an intermediate energy cyclotron/accelerator. All those positron 634 emitters have either been shown to be, or are expected to be, suitable for PET 635 measurements; only in the case of ⁸⁶Y the large number of associated γ -rays cause some 636 difficulty, but after proper corrections, the images can be satisfactorily interpreted.

Regarding the therapeutic radionuclides, ⁸⁹Sr and ⁹⁰Y decay by emission of β^{-} -638 particles of intermediate energy. Both are produced in a nuclear reactor, the former via the 639 (n,p) reaction and the latter via the 90 Sr/ 90 Y generator system. The generator parent 90 Sr is 640 separated from fission products. Both ⁸⁹Sr and ⁹⁰Y are commercially available. The β^- -641 particle endpoint energies of the remaining four radionuclides, namely ⁴⁷Sc, ⁶⁷Cu, ¹³¹I and 642 ¹⁶¹Tb are relatively low (< 610 keV). The radionuclide ¹³¹I is produced in a nuclear reactor 643 644 either via fission or more commonly via the sequence ${}^{130}\text{Te}(n,\gamma){}^{131\text{m},g}\text{Te} \rightarrow {}^{131}\text{I}$. It has been known for a very long time and is extensively used in internal radiotherapy. It is 645 commercially available. The radionuclide ¹⁶¹Tb is also produced in a nuclear reactor 646 through the sequence ${}^{160}\text{Gd}(n,\gamma){}^{161}\text{Gd} \rightarrow {}^{161}\text{Tb}$ and it is available in sufficient quantities. In 647 648 recent years interest has also been growing in the comparison of the therapeutic effect of 649 the four very similar β^- -particle emitters, namely ⁴⁷Sc, ⁶⁷Cu, ¹⁶¹Tb, and ¹⁷⁷Lu [173-175]. The radionuclides ⁴⁷Sc and ⁶⁷Cu are very interesting but difficult to produce. Therefore 650 651 presently strong efforts are underway to produce them through neutron, photon and charged 652 particle induced reactions.

653

In contrast to the above mentioned theranostic pairs of radionuclides consisting of a β^+ -emitter and a β^- -emitter, the pair ¹⁵²Tb/¹⁴⁹Tb is unique in that the radionuclide ¹⁵²Tb is a β^+ -emitter and ¹⁴⁹Tb is an α -emitter. The efficacy of ¹⁴⁹Tb for targeted α -therapy has been demonstrated but the exotic production route, involving spallation and on-line mass separation, makes its availability very rare. Further development work is called for.

659

Besides the 7 rather established theranostic pairs of radionuclides discussed in this review, the pair 72 As/ 77 As is in development [cf. 176-178]. Furthermore, there are 3 other pairs where the combination consists of a positron emitter and an Auger electron emitter as a therapeutic partner. They are 68 Ga/ 67 Ga, 110g In/ 111 In and 152 Tb/ 155 Tb. However, since Auger therapy using the radionuclides 67 Ga, 111 In and 155 Tb is still developing, those pairs have not been considered in this review.

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667 In conclusion, it may be stated that the field of theranostics is attracting tremendous 668 attention today, but the availability of the respective radionuclides plays a very important role. Concerted efforts are needed to produce several of the above mentioned radionuclides in quantities sufficient for clinical studies. Enhanced utilization of intermediate energy cyclotrons/accelerators would be very advantageous. Furthermore, for production of a few special radionuclides, use of powerful electron linear accelerators may be beneficial. Similarly, the use of some rather unconventional methods, like heavy-ion induced reactions and on-line mass separation of radioactive products, may also be worthwhile, especially for small scale production of some exotic radionuclides for tracer studies.

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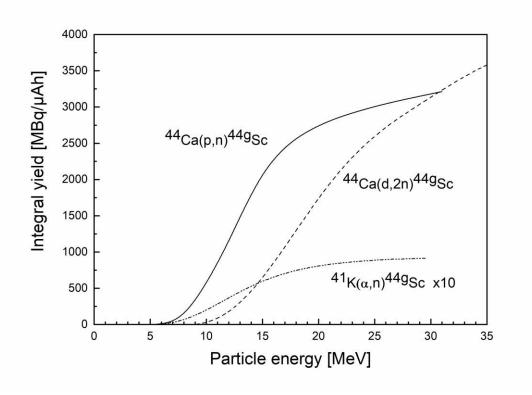
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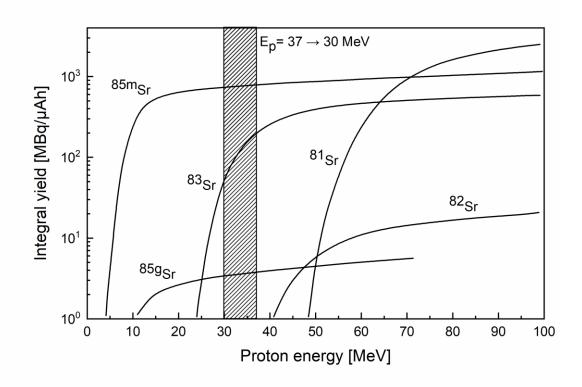
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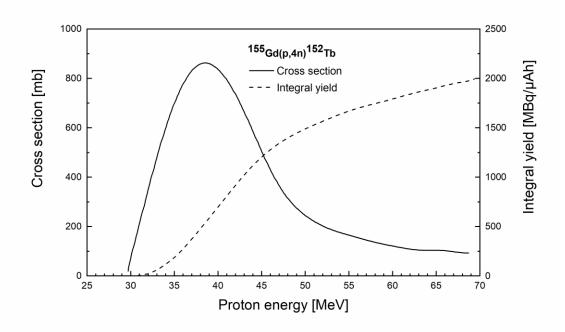
Figures



1216Fig. 1Thick target yields of ${}^{44g}Sc$ calculated from the excitation functions of1217 ${}^{44}Ca(p,n){}^{44g}Sc, {}^{44}Ca(d,2n){}^{44g}Sc$ and ${}^{41}K(\alpha,n){}^{44g}Sc$ reactions reported in refs. [25,121826, 28, 30-32]. The values are shown as curves as a function of the particle1219energy.



1222Fig 2. Calculated integral yields of radionuclides of Sr formed in the interaction of 85 Rb1223with protons of increasing energies. The optimum energy range for the production1224of 83 Sr is $E_p = 37 \rightarrow 30$ MeV (after Kastleiner et al. [131]).



1227Fig 3. Excitation function of the 155Gd(p,4n)152Tb reaction and the calculated integral1228yield of 152Tb assuming a 100 % enrichment of the target (adapted from Steyn et1229al. [162]).