

New Estimates of the Nonlinear Fourier Transform for the Defocusing NLS Equation

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The defocusing NLS equation $i u_t = -u_{xx} + 2|u|^2 u$ on the circle admits a global nonlinear Fourier transform, also known as Birkhoff map, linearizing the NLS flow. The regularity properties of u are known to be closely related to the decay properties of the corresponding nonlinear Fourier coefficients. In this paper, we quantify this relationship by providing two-sided polynomial estimates of all integer Sobolev norms $\|u\|_m$, $m \geq 0$, in terms of the weighted norms of the nonlinear Fourier transformed.

1 Introduction

Consider the *defocusing nonlinear Schrödinger* (NLS) equation

$$i \partial_t u = -\partial_x^2 u + 2|u|^2 u, \quad x \in \mathbb{T},$$

on the circle $\mathbb{T} = \mathbb{R}/\mathbb{Z}$ with u taken from $L^2 := L^2(\mathbb{T}, \mathbb{C})$. As is well known, the NLS equation can be written as an *infinite-dimensional Hamiltonian system*. We introduce the phase space $L_c^2 := L^2 \times L^2$ with elements $\varphi = (\varphi_-, \varphi_+)$ and define the inner product

$$\langle \varphi, \psi \rangle := \int_{\mathbb{T}} \varphi_+ \bar{\psi}_+ + \varphi_- \bar{\psi}_- \, dx,$$

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which makes L_c^2 a Hilbert space. The associated norm is denoted by $\|\varphi\|_0$. Further, we endow this space with the Poisson structure given by

$$\{F, G\} := -i \int_{\mathbb{T}} (\partial_{\varphi_-} F \partial_{\varphi_+} G - \partial_{\varphi_+} F \partial_{\varphi_-} G) dx.$$

Here $\partial_{\varphi_-} F$ and $\partial_{\varphi_+} F$ denote the components of the L^2 -gradient $\partial_{\varphi} F$ of a C^1 -functional F . The Hamiltonian system with Hamiltonian

$$\mathcal{H}(\varphi_-, \varphi_+) := \int_{\mathbb{T}} (\partial_x \varphi_- \partial_x \varphi_+ + \varphi_-^2 \varphi_+^2) dx$$

is then given by

$$i\partial_t(\varphi_-, \varphi_+) = (\partial_{\varphi_+} \mathcal{H}, -\partial_{\varphi_-} \mathcal{H}),$$

and the defocusing NLS is obtained by restriction to the invariant subspace of real-type states

$$L_r^2 := \{\varphi \in L_c^2 : \varphi^* = \varphi\}, \quad \varphi^* = (\bar{\varphi}_+, \bar{\varphi}_-).$$

Indeed, with $\varphi = (u, \bar{u})$ the defocusing NLS equation becomes

$$i\partial_t u = i\{u, \mathcal{H}\} = \partial_{\bar{u}} \mathcal{H}(u, \bar{u}).$$

After the KdV equation, the NLS equation was the second evolution equation known to be *completely integrable* by the inverse scattering method [27]. In fact, according to [10] (cf. also [1, 2, 23]), the defocusing NLS equation is integrable in the strongest possible sense meaning that it admits global *Birkhoff coordinates* $(x_n, y_n)_{n \in \mathbb{Z}}$.

To state our main result, let $\sum_{n \in \mathbb{Z}} (\varphi_{2n}^- e^{-i2n\pi x}, \varphi_{2n}^+ e^{i2n\pi x})$ denote the Fourier series of $\varphi = (\varphi_-, \varphi_+) \in L_c^2$ and introduce for any $s \geq 0$ the Sobolev norm $\|\varphi\|_s$ given by

$$\|\varphi\|_s^2 := \sum_{n \in \mathbb{Z}} \langle 2n\pi \rangle^{2s} (|\varphi_{2n}^-|^2 + |\varphi_{2n}^+|^2), \quad \langle x \rangle := 1 + |x|.$$

The space of all $\varphi \in L_c^2$ with $\|\varphi\|_s < \infty$ is denoted by H_c^s , and $H_r^s := L_r^2 \cap H_c^s$. Furthermore, let us introduce the model space

$$h_r^s := \{(x, y) = (x_n, y_n)_{n \in \mathbb{Z}} : \|(x, y)\|_s := (\|x\|_s^2 + \|y\|_s^2)^{1/2} < \infty\},$$

where the norm $\|x\|_s$ is defined by

$$\|x\|_s^2 := \sum_{n \in \mathbb{Z}} \langle 2n\pi \rangle^{2s} |x_n|^2.$$

This space is endowed with the canonical Poisson structure $\{x_n, y_m\} = -\delta_{n,m}$ while all other brackets vanish.

The *Birkhoff map* $\varphi \mapsto (x_n, y_n)_{n \in \mathbb{Z}}$ is a bi-analytic, canonical diffeomorphism $\Omega: H_r^0 \rightarrow h_r^0$, whose restriction $\Omega: H_r^m \rightarrow h_r^m$ is again bi-analytic for any integer $m \geq 1$. On h_r^1 the transformed NLS Hamiltonian $\mathcal{H} \circ \Omega^{-1}$ is a real analytic function of the actions $I_n = (x_n^2 + y_n^2)/2$ alone and the NLS evolution takes the particularly simple form

$$\dot{x}_n = -\omega_n y_n, \quad \dot{y}_n = \omega_n x_n, \quad \omega_n := \partial_{I_n} \mathcal{H}.$$

One may thus think of Ω as a nonlinear Fourier transform for the defocusing NLS equation. Remarkably, the derivative $d_0 \Omega$ of Ω at the origin is the Fourier transform and on L_r^2 , as for the Fourier transform,

$$\|\Omega(\varphi)\|_0 = \|\varphi\|_0,$$

which we also refer to as Parseval's identity—cf., for example, [10, 23]. Our main result says that for higher order Sobolev norms the following version of Parseval's identity holds for the nonlinear map Ω .

Theorem 1.1. For any integer $m \geq 1$, there exist absolute constants $c_m, d_m > 0$ such that the restriction of Ω to H_r^m satisfies the two-sided estimates

$$(i) \quad \|\Omega(\varphi)\|_m \leq c_m (\|\varphi\|_m + (1 + \|\varphi\|_{\tilde{m}-1})^{2m} \|\varphi\|_{\tilde{m}-1}),$$

where $\tilde{m} = \max(m, 2)$, and

$$(ii) \quad \|\varphi\|_m \leq d_m (\|\Omega(\varphi)\|_m + (1 + \|\Omega(\varphi)\|_{m-1})^{4m-3} \|\Omega(\varphi)\|_{m-1}). \quad \square$$

The main feature of Theorem 1.1 is that the constants c_m and b_m can be chosen independently of φ .

Note that the estimate (i) is linear in the highest Sobolev norm $\|\varphi\|_m$ for $m \geq 2$, and that the estimate (ii) is linear in the highest weighted h_r^m -norm $\|\Omega(\varphi)\|_m$ of $\Omega(\varphi)$. Hence, Theorem 1.1 shows that the 1-smoothing property of the Birkhoff map Ω established in [17] holds in a uniform fashion.

The proof of Theorem 1.1 relies on estimates of the action variables $I(\varphi) = (I_n)_{n \in \mathbb{Z}}$ where $I_n = (x_n^2 + y_n^2)/2$, $n \in \mathbb{Z}$. The decay properties of the actions I_n are known to be closely related to the regularity properties of φ —cf. [7, 15, 16]. We need to quantify this relationship by providing two-sided estimates of the Sobolev norms of φ in terms of weighted ℓ^1 -norms of $I(\varphi)$. For that purpose introduce the weighted sequence space

ℓ_s^1 whose norm is defined by

$$\|I(\varphi)\|_{\ell_s^1} := \sum_{n \in \mathbb{Z}} \langle 2n\pi \rangle^s |I_n(\varphi)|.$$

Theorem 1.2. For any integer $m \geq 1$, there exist absolute constants c_m and d_m , such that for all $\varphi \in H_r^m$

$$(i) \quad \|I(\varphi)\|_{\ell_{2m}^1} \leq c_m^2 (\|\varphi\|_m^2 + (1 + \|\varphi\|_{\tilde{m}-1})^{4m} \|\varphi\|_{\tilde{m}-1}^2),$$

where $\tilde{m} = \max(m, 2)$, and

$$(ii) \quad \|\varphi\|_m^2 \leq d_m^2 (\|I(\varphi)\|_{\ell_{2m}^1} + (1 + \|I(\varphi)\|_{\ell_{2m-2}^1})^{4m-3} \|I(\varphi)\|_{\ell_{2m-2}^1}). \quad \square$$

Remark. The same constants c_m, d_m of Theorem 1.1 can be used in Theorem 1.2. \square

It turns out that versions of the estimates (i) of Theorems 1.1 and 1.2 hold for a larger family of spaces referred to *weighted Sobolev spaces*—see [12, 13] for an introduction. A *normalized, submultiplicative, and monotone weight* is a symmetric function $w: \mathbb{Z} \rightarrow \mathbb{R}$ with

$$w_n \geq 1, \quad w_n = w_{-n}, \quad w_{n+m} \leq w_n w_m, \quad w_{|n|} \leq w_{|n|+1},$$

for all $n, m \in \mathbb{Z}$. The class of all such weights is denoted by \mathcal{M} and H_c^w is the space of L_c^2 functions φ with finite w -norm

$$\|\varphi\|_w^2 := \sum_{n \in \mathbb{Z}} w_{2n}^2 (|\varphi_{2n}^-|^2 + |\varphi_{2n}^+|^2).$$

Furthermore, h_r^w denotes the subspace of ℓ_r^2 where $\|x\|_w^2 + \|y\|_w^2 < \infty$,

$$\|x\|_w^2 := \sum_{n \in \mathbb{Z}} w_{2n}^2 |x_n|^2.$$

For any $s \geq 0$, the *Sobolev weight* $\langle n\pi \rangle^s$ gives rise to the usual Sobolev space H_c^s . For $s \geq 0$ and $a > 0$, the *Abel weight* $\langle n\pi \rangle^s e^{a|n|}$ gives rise to the space $H_c^{s,a}$ of L_c^2 -functions, which can be analytically extended to the open strip $\{z : |\Im z| < a/2\pi\}$ of the complex plane with traces in H_c^s on the boundary lines. In between are, among others, the *Gevrey weights*

$$\langle n \rangle^s e^{a|n|^\sigma}, \quad 0 < \sigma < 1, \quad s \geq 0, \quad a > 0,$$

which give rise to the Gevrey spaces $H_c^{s,a,\sigma}$, as well as weights of the form

$$\langle n \rangle^s \exp \left(\frac{a|n|}{1 + \log^\sigma \langle n \rangle} \right), \quad 0 < \sigma < 1, \quad s \geq 0, \quad a > 0,$$

that are lighter than Abel weights but heavier than Gevrey weights.

To avoid certain technicalities in our estimates, we restrict ourselves to weights incorporating a factor which grows at a linear rate. We thus introduce the subclass

$$\mathcal{M}^1 := \{w \in \mathcal{M} : w_n = \langle n \rangle v_n \text{ for all } n \in \mathbb{Z} \text{ with some } v \in \mathcal{M}\}.$$

Finally, we assume that all weights $w \in \mathcal{M}$ are piecewise linearly extended to functions on the real line $w : \mathbb{R} \rightarrow \mathbb{R}_{>0}$, $t \mapsto w[t]$.

Theorem 1.3. For any weight $w \in \mathcal{M}^1$, there exists a complex neighborhood W_w of H_r^w within H_c^w and a constant c_w , such that

$$(i) \quad \sum_{n \in \mathbb{Z}} w_{2n}^2 |I_n| \leq c_w^2 w [16 \|\varphi\|_w^2]^2 \|\varphi\|_w^2.$$

Moreover, the restriction of the Birkhoff map Ω to H_r^w takes values in h_r^w and satisfies

$$(ii) \quad \|\Omega(\varphi)\|_w \leq c_w w [16 \|\varphi\|_w^2] \|\varphi\|_w. \quad \square$$

In this more general set up, the bounds in (i) and (ii) of Theorem 1.3 are of the same type as the weight, and are valid for all submultiplicative weights including those growing exponentially fast. The following version of Theorem 1.3 for Sobolev spaces of real exponent complements the results of Theorems 1.2 and 1.3.

Corollary 1.4. For any real $s \geq 1$, there exist a complex neighborhood W_s of H_r^s and a constant c_s such that

$$\|I(\varphi)\|_{\ell_{2s}^1} \leq c_s^2 (1 + \|\varphi\|_s)^{4s} \|\varphi\|_s^2.$$

Moreover, the restriction of the Birkhoff map Ω to H_r^s takes values in h_r^s and satisfies

$$\|\Omega(\varphi)\|_s \leq c_s (1 + \|\varphi\|_s)^{2s} \|\varphi\|_s. \quad \square$$

Outline. The action variables I_n and more generally the action variables $J_{n,k}$ on levels $k \geq 1$ can be defined entirely in terms of the periodic spectrum of the associated Zakharov–Shabat operator used in the Lax-pair formulation of the NLS equation. More to the point, consider the operator

$$L(\varphi) = \begin{pmatrix} i & \\ & -i \end{pmatrix} \frac{d}{dx} + \begin{pmatrix} & \psi \\ \bar{\psi} & \end{pmatrix},$$

with periodic boundary conditions on the interval $[0, 2]$ of twice the length of the periodicity of $\varphi = (\psi, \bar{\psi}) \in L_r^2$. The spectrum of $L(\varphi)$ is well known to be real, pure point, and

to consist of a double infinite sequence of eigenvalues

$$\cdots \leq \lambda_{n-1}^+ < \lambda_n^- \leq \lambda_n^+ < \lambda_{n+1}^- \leq \cdots$$

with asymptotic behavior $\lambda_n^\pm \sim n\tau$ as $|n| \rightarrow \infty$. The asymptotic behavior of the actions on odd levels $k = 2m + 1$ turns out to be

$$J_{n,2m+1} \sim (\lambda_n^\pm)^{2m} I_n \sim (n\tau)^{2m} I_n, \quad |n| \rightarrow \infty,$$

and they satisfy the trace formula

$$\sum_{n \geq 1} J_{n,2m+1} = \frac{(-1)^{m+1}}{4^m} \mathcal{H}_{2m+1}, \quad m \geq 1,$$

where \mathcal{H}_{2m+1} denotes the $2m + 1$ th Hamiltonian in the NLS hierarchy,

$$\mathcal{H}_1 = \int_{\mathbb{T}} |\psi|^2 \, dx, \quad \mathcal{H}_3 = \int_{\mathbb{T}} (|\psi'|^2 + |\psi|^4) \, dx, \dots$$

Note that \mathcal{H}_3 denotes the Hamiltonian of the NLS equation. On H_r^m , for $m \geq 1$,

$$\mathcal{H}_{2m+1} = (-1)^{(m+1)} \int_{\mathbb{T}} (|\psi^{(m)}|^2 + p_m(\psi, \dots, \psi^{(m-1)})) \, dx,$$

where p_m is a polynomial expression in ψ and its first $m - 1$ derivatives. Viewing \mathcal{H}_{2m+1} as a lower order perturbation of the H_r^m -norm we formally obtain

$$\sum_{n \in \mathbb{Z}} (n\tau)^{2m} I_n \sim \sum_{n \in \mathbb{Z}} J_{n,2m+1} \sim \|\varphi\|_m^2 + O(\|\varphi\|_{m-1}).$$

The essential ingredient to make this formal statement explicit is a sufficiently accurate localization of the periodic eigenvalues λ_n^\pm whose threshold in $|n|$ depends only on $\|\varphi\|_1$. For $|n|$ above this threshold, we can directly compare the weighted action norms and the polynomial expressions in φ as described above, while the remainder for $|n|$ below the threshold can be regarded as an H^1 -error term. Thereof we obtain Theorem 1.2, which directly implies Theorem 1.1. Note that our method of proof completely avoids the use of auxiliary spectral quantities such as *spectral heights* or *conformal mapping theory*.

To prove Theorem 1.3, we take a slightly different approach by quantitatively estimating the action variables in terms of the spacing of the periodic eigenvalues of the associated Zakharov–Shabat operator. For the latter, we obtain estimates in any weighted norm, which allows us to obtain Theorem 1.3.

Related results. Theorem 1.1 for $m = 0$ is referred to as Parseval's identity,

$$\frac{1}{2} \|\Omega(\varphi)\|_0 = \|I(\varphi)\|_{\ell^1} = \frac{1}{2} \|\varphi\|_0^2,$$

and is well known—see [10, 23]. The case $m = 1$ was proved by Korotyaev [19] using conformal mapping theory, see also [21]. However, his method does not seem applicable for the case $m \geq 2$. In fact, it is stated as an open problem in [21].

For the case of the KdV equation

$$\partial_t u = -\partial_x^3 u + 6u\partial_x u, \quad x \in \mathbb{T},$$

Korotyaev [18, 20] obtained polynomial bounds of the Sobolev norms $\|u\|_m$ in terms of the action variables where the order of the polynomials grows factorial in m . Note that the bound in (ii) of Theorem 1.2 is of order 1 in the Sobolev norm $\|\varphi\|_m$ and the order of the remainder grows linearly in m . It turns out that our method can also be applied to the KdV equation. In [25], we improve on the bounds obtained by Korotyaev in [18, 20].

For NLS in weighted Sobolev spaces, the qualitative relationship

$$\varphi \in H_r^w \Leftrightarrow \Omega(\varphi) \in h_r^w$$

is a corollary of the methods presented in [7, 16]—at least for weights growing at subexponential speed. To the best of our knowledge, the estimate of $\|\Omega(\varphi)\|_w$ on H_r^w as well as the estimate of $\|I(\varphi)\|_w$ on a small complex neighborhood of H_r^w as presented in Theorem 1.3 are new.

Viewing the action I_n as a 1-smoothing perturbation of the squared modulus of the n th Fourier coefficient, our method of comparing the weighted action norms with the Hamiltonians of the NLS hierarchy amounts to a separate analysis of Fourier modes of low and high frequencies. This idea has a long history in the analysis of nonlinear partial differential equations. Most recently, it lead Colliander et al. [3–6] to invent the I-Method, which allows to obtain global well posedness of subcritical equations in low regularity regimes where the Hamiltonian (or other integrals) of the equation cease to be well defined. The idea is to damp all sufficiently high Fourier modes of a local solution such that the Hamiltonian can be controlled by weaker norms while still being an almost conserved quantity. The difficulty here is to choose the damping subtle enough such that the nonlinearity of the equation does not create a significant interaction of low and high frequencies. Our situation is so to say opposed to that of the I-Method: As we aim for quantitative global estimates, controlling the modes of low frequencies is the most delicate part. Here the localization of the periodic eigenvalues of the Zakharov–Shabat

operator associated with the NLS equation plays a crucial role. Note that there exists a vast amount of literature on the spectral theory of these operators—see, for example, [7, 10, 11, 13, 21] and the references therein.

Organization of the paper. In Section 2, we recall the definition of the NLS action variables on integer levels $k \geq 1$ as well as the trace formulae relating them to the hierarchy of NLS Hamiltonians. The somewhat lengthy proof of the analyticity properties of the action integrand can be found in Appendix 1. The quadratic localization of the Zakharov–Shabat spectrum is obtained in Section 3, and is subsequently used in Sections 4 and 5 to obtain Theorem 1.2(i) and (ii), respectively. In Section 6, we obtain the estimate of the actions in terms of the spacing of the Zakharov–Shabat eigenvalues which implies Theorem 1.3.

2 Setup

In this section, we briefly recall the definition of NLS action variables from [10], as well as the main properties of the periodic spectrum of Zakharov–Shabat operators used to define them.

For a *potential* $\varphi = (\varphi_-, \varphi_+) \in H_c^0 = L_c^2$, consider the Zakharov–Shabat operator

$$L(\varphi) := \begin{pmatrix} i & \\ & -i \end{pmatrix} \frac{d}{dx} + \begin{pmatrix} & \varphi_- \\ \varphi_+ & \end{pmatrix}$$

on the interval $[0, 2]$ with periodic boundary conditions. The spectrum of $L(\varphi)$ is well known to be pure point, and more precisely, to consist of a sequence of pairs of complex eigenvalues $\lambda_n^+(\varphi)$ and $\lambda_n^-(\varphi)$, listed with algebraic multiplicities, such that

$$\lambda_n^\pm(\varphi) = n\pi + \ell_n^2, \quad n \in \mathbb{Z}.$$

Here ℓ_n^2 denotes a generic ℓ^2 -sequence. We may order the eigenvalues lexicographically—first by their real part, and second by their imaginary part—to represent them as a bi-infinite sequence of complex eigenvalues

$$\cdots \preccurlyeq \lambda_{n-1}^+ \preccurlyeq \lambda_n^- \preccurlyeq \lambda_n^+ \preccurlyeq \lambda_{n+1}^- \preccurlyeq \cdots.$$

By a slight abuse of notation, we call the eigenvalues of $L(\varphi)$ the *periodic spectrum* of φ . Further we introduce the *gap lengths*

$$\gamma_n := \lambda_n^+ - \lambda_n^- = \ell_n^2, \quad n \in \mathbb{Z}.$$

To obtain another characterization of the periodic spectrum, we denote by $M(x, \lambda, \varphi)$ the standard fundamental solution of the ordinary differential equation $L(\varphi)M = \lambda M$, and introduce the discriminant

$$\Delta(\lambda, \varphi) := \operatorname{tr} M(1, \lambda, \varphi).$$

To simplify matters, we may drop some or all of its arguments from the notation, whenever there is no danger of confusion. The periodic spectrum of φ is precisely the zero set of the entire function $\Delta^2(\lambda) - 4$, and we have the product representation

$$\Delta^2(\lambda) - 4 = -4 \prod_{n \in \mathbb{Z}} \frac{(\lambda_n^+ - \lambda)(\lambda_n^- - \lambda)}{\pi_n^2}, \quad \pi_n := \begin{cases} n\pi, & n \neq 0, \\ 1, & n = 0. \end{cases}$$

Hence, this function is a spectral invariant. We also need the λ -derivative $\Delta^\bullet := \partial_\lambda \Delta$ whose zeros are denoted by λ_n^\bullet and satisfy $\lambda_n^\bullet = n\pi + \ell_n^2$. This derivative has the product representation

$$\Delta^\bullet(\lambda) = 2 \prod_{n \in \mathbb{Z}} \frac{\lambda_n^\bullet - \lambda}{\pi_n}.$$

If the potential φ is of real type, as in the case of the defocusing NLS equation, the periodic spectrum is real, the eigenvalues are characterized by Floquet theory, and the lexicographical ordering reduces to the ordering of real numbers

$$\cdots \leq \lambda_{n-1}^+ < \lambda_n^- \leq \lambda_n^\bullet \leq \lambda_n^+ < \lambda_{n+1}^- \leq \cdots.$$

Each $\varphi \in L_r^2$ has an open neighborhood V_φ within L_c^2 for which there exist disjoint closed disks $(U_n)_{n \in \mathbb{Z}}$ centered on the real axis with the properties:

- (i) $\lambda_n^\pm(\psi)$ and $\lambda_n^\bullet(\psi)$ are contained in the interior of U_n for every $\psi \in V_\varphi$,
- (ii) there exists a constant $c \geq 1$ such that for $m \neq n$,

$$c^{-1}|m - n| \leq \operatorname{dist}(U_n, U_m) \leq c|m - n|,$$

- (iii) $U_n = \{|\lambda - n\pi| \leq \pi/4\}$ for $|n|$ sufficiently large.

Such disks are called *isolating neighborhoods*. The union of all V_φ defines an open and connected neighborhood of L_r^2 within L_c^2 and is denoted by W . Throughout this text, V_φ denotes a neighborhood of φ such that a common set of isolating neighborhoods for all $\psi \in V_\varphi$ exists.

Following Flaschka and McLaughlin's [8] approach for the KdV equation, one can define action variables for the defocusing NLS equation by Arnold's formula—see also [23]

$$I_n = \frac{1}{\pi} \int_{a_n} \frac{\lambda \Delta^\bullet(\lambda)}{\sqrt{\Delta^2(\lambda) - 4}} d\lambda.$$

Here a_n denotes a cycle around $(\lambda_n^-, \lambda_n^+)$ on the spectral curve

$$C_\varphi = \{(\lambda, z) : z^2 = \Delta^2(\lambda, \varphi) - 4\} \subset \mathbb{C},$$

on which the square root $\sqrt{\Delta^2(\lambda) - 4}$ is defined. This curve is another spectral invariant associated with φ , and an open Riemann surface of infinite genus if and only if the periodic spectrum of φ is simple. To avoid the technicalities involved with this curve, we follow the approach presented in [10] and fix proper branches of the square root which allows us to reduce the definition of the actions to standard contour integrals in the complex plane.

First, we denote by $\sqrt[+]{}$ the *principal branch* of the square root on the complex plane minus the ray $(-\infty, 0]$. Second, we define the *standard root*

$$\varsigma_n(\lambda) = \sqrt[n]{(\lambda_n^+ - \lambda)(\lambda_n^- - \lambda)}, \quad \lambda \notin [\lambda_n^-, \lambda_n^+],$$

by the condition

$$\varsigma_n(\lambda) := (\tau_n - \lambda) \sqrt[+]{1 - \gamma_n^2/4(\tau_n - \lambda)}, \quad \tau_n = (\lambda_n^- + \lambda_n^+)/2, \quad (1)$$

for all $|\lambda|$ sufficiently large. For any $\varphi \in W$, the standard root is analytic in λ on $\mathbb{C} \setminus [\lambda_n^-, \lambda_n^+]$ and in (λ, ψ) on $(\mathbb{C} \setminus U_n) \times V_\varphi$. Thirdly, we define the *canonical root* $\sqrt[\circ]{\Delta^2(\lambda) - 4}$ by the product representation

$$\sqrt[\circ]{\Delta^2(\lambda) - 4} := 2i \prod_{m \in \mathbb{Z}} \frac{\varsigma_m(\lambda)}{\pi_m}.$$

For any $\varphi \in W$, this root is analytic in λ on $\mathbb{C} \setminus \bigcup_{\gamma_n \neq 0} [\lambda_n^-, \lambda_n^+]$ and in (λ, ψ) on $(\mathbb{C} \setminus \bigcup_{n \in \mathbb{Z}} U_n) \times V_\varphi$ —see [10, Section 12] for all the details.

The n th NLS action variable of $\varphi \in W$ is then given by

$$I_n := \frac{1}{\pi} \int_{\Gamma_n} \frac{\lambda \Delta^\bullet(\lambda)}{\sqrt[\circ]{\Delta^2(\lambda) - 4}} d\lambda,$$

where Γ_n denotes any sufficiently close circuit around $[\lambda_n^-, \lambda_n^+]$. More generally, the n th action on level $k \geq 1$ is given by

$$J_{n,k} := \frac{1}{k\pi} \int_{\Gamma_n} \frac{\lambda^k \Delta^\bullet(\lambda)}{\sqrt[k]{\Delta^2(\lambda) - 4}} d\lambda.$$

It was shown in [10] and, for convenience of the reader, will be reproved in the sequel that each action variable is analytic on W and vanishes if and only if γ_n is zero.

If φ is of real type, then all actions are real and those on odd levels, such as $J_{n,1} = I_n$, are nonnegative. Moreover, the actions on level one are well known to satisfy the trace formula,

$$\sum_{n \in \mathbb{Z}} I_n(\varphi) = \mathcal{H}_1(\varphi) = \frac{1}{2} \|\varphi\|_0^2 = \frac{1}{2} \int_{\mathbb{T}} (|\varphi_-|^2 + |\varphi_+|^2) dx. \quad (2)$$

Similar trace formulae have been derived by McKean and Vaninsky [23] expressing the actions on any level $k \geq 1$ in terms of Hamiltonians of the *NLS hierarchy*. The first three Hamiltonians of this hierarchy are

$$\begin{aligned} \mathcal{H}_1(\varphi) &= \int_{\mathbb{T}} \varphi_- \varphi_+ dx, \\ \mathcal{H}_2(\varphi) &= \frac{1}{2} \int_{\mathbb{T}} (\varphi'_- \varphi_+ - \varphi_- \varphi'_+) dx, \\ \mathcal{H}_3(\varphi) &= \int_{\mathbb{T}} (\varphi'_- \varphi'_+ + \varphi_+^2 \varphi_-^2) dx, \end{aligned}$$

and in general, for any sufficiently regular $\varphi \in L_c^2$,

$$\mathcal{H}_{k+1}(\varphi) = \int_{\mathbb{T}} (-\varphi_- \varphi_+^{(k)} + q_k(\varphi, \dots, \varphi^{(k-1)})) dx, \quad k \geq 1,$$

with q_k being a canonically determined polynomial in φ and its first $k-1$ derivatives—see Appendix 2. The following version of the trace formula is taken from [10].

Theorem 2.1 (Trace formula). For any $k \geq 2$ and any $\varphi \in H_c^{k-1} \cap W$,

$$\sum_{n \in \mathbb{Z}} J_{n,k}(\varphi) = -\frac{1}{(2i)^{k-1}} \mathcal{H}_k(\varphi). \quad (3)$$

□

In particular, for every sufficiently regular real-type potential

$$\sum_{n \in \mathbb{Z}} J_{n,2m+1}(\varphi) = \frac{1}{4^m} \int_{\mathbb{T}} (|\varphi^{(m)}|^2 + \dots) dx, \quad m \geq 0.$$

This identity is used in Sections 4 and 5 to estimate the actions on level $2m + 1$ in terms of $\|\varphi\|_m$. In order to obtain thereof estimates for the action variables on level 1, a detailed analysis of the analytical properties of the action integrand is necessary. To this end, we define for any $\varphi \in W$ on $(\mathbb{C} \setminus \bigcup_{n \in \mathbb{Z}} U_n) \times V_\varphi$ the complex 1-form

$$\omega(\lambda, \psi) := \frac{\Delta^\bullet(\lambda, \psi)}{\sqrt[5]{\Delta^2(\lambda, \psi) - 4}} d\lambda = -i \prod_{m \in \mathbb{Z}} \frac{\lambda_m^\bullet(\psi) - \lambda}{\zeta_m(\lambda, \psi)} d\lambda. \quad (4)$$

We call a path in the complex plane *admissible* for $\varphi \in L_c^2$ if, except possibly at its end-points, it does not intersect any gap $[\lambda_n^-(\varphi), \lambda_n^+(\varphi)]$.

Lemma 2.2. For each $\varphi \in W$, the 1-form ω has the following properties:

- (i) ω is analytic on $(\mathbb{C} \setminus \bigcup_{n \in \mathbb{Z}} U_n) \times V_\varphi$,
- (ii) $\omega(\cdot, \varphi)$ is analytic on $\mathbb{C} \setminus \bigcup_{n \neq 0} [\lambda_n^-, \lambda_n^+]$, and
- (iii) for any admissible path from λ_n^- to λ_n^+ in U_n ,

$$\int_{\lambda_n^-}^{\lambda_n^+} \omega = 0.$$

In particular, for any closed circuit Γ_n in U_n around $[\lambda_n^-, \lambda_n^+]$,

$$\int_{\Gamma_n} \omega = 0. \quad \square$$

Proof. Since $\Delta^2(\lambda, \psi) - 4$ vanishes if and only if λ is a periodic eigenvalue, and numerator and dominator of $\omega(\lambda, \psi)$ are analytic on $(\mathbb{C} \setminus \bigcup_{n \in \mathbb{Z}} U_n) \times V_\varphi$, the first claim follows immediately.

By the same reasoning, $\omega(\cdot, \varphi)$ is analytic on $\mathbb{C} \setminus \bigcup_{n \in \mathbb{Z}} [\lambda_n^-, \lambda_n^+]$. Moreover, if the n th gap is collapsed, that is, $\lambda_n^+ = \lambda_n^- =: \lambda_n$, then $\Delta^2(\lambda) - 4$ has a double root at λ_n , and hence

$$0 = \Delta(\lambda_n) \Delta^\bullet(\lambda_n) = 2(-1)^n \Delta^\bullet(\lambda_n).$$

As $\Delta^\bullet(\lambda)$ has a single root in U_n , namely λ_n^\bullet , we conclude $\lambda_n^\bullet = \lambda_n = \lambda_n^\pm$, and the n th term in the product representations of Δ^\bullet and $\sqrt[5]{\Delta^2 - 4}$ cancels. So $\omega(\cdot, \varphi)$ is analytic on all of U_n , and the second claim follows.

We first consider the case where φ is of real type. Then $(-1)^n \Delta(\lambda, \varphi) \geq 2$ on $[\lambda_n^-, \lambda_n^+]$, as shown in [10], and hence

$$\min_{\lambda_n^- \leq \lambda \leq \lambda_n^+} (-1)^n \Delta(\lambda, \varphi) - \sqrt[5]{\Delta^2(\lambda, \varphi) - 4} > 0.$$

Thus, if we choose a circuit Γ sufficiently close to $[\lambda_n^-, \lambda_n^+]$, and a sufficiently small neighborhood $V \subset V_\varphi$ of φ , then $[\lambda_n^-(\psi), \lambda_n^+(\psi)]$ is enclosed by Γ and the real part of $(-1)^n(\Delta(\lambda, \psi) - \sqrt[n]{\Delta^2(\lambda, \psi) - 4})$ is positive on that circuit for all $\psi \in V$. In consequence, the principal branch of the logarithm

$$l_n(\lambda, \psi) = \log \frac{(-1)^n}{2} (\Delta(\lambda, \psi) + \sqrt[n]{\Delta^2(\lambda, \psi) - 4})$$

is analytic in a neighborhood of Γ and $dl_n = \omega$. It follows that the analytic functional $\psi \mapsto \int_{\partial U_n} \omega$ vanishes on an open neighborhood of φ , and hence on all of V_φ . Since $W = \bigcup_{\varphi \in L_r^2} V_\varphi$, it follows that $\int_{\Gamma_n} \omega = 0$ for every $\varphi \in W$.

Finally, consider $\int_{\lambda_n^-}^{\lambda_n^+} \omega$ with the path of integration chosen to be admissible. As ω is closed around the gap, the integral does not depend on the chosen admissible path. Suppose $\gamma_n \neq 0$, then by the product representation (4) of ω ,

$$\int_{\lambda_n^-}^{\lambda_n^+} \omega = -i \int_{\lambda_n^-}^{\lambda_n^+} \frac{\lambda_n^\bullet - \lambda}{\varsigma_n(\lambda)} \chi_n(\lambda) d\lambda, \quad \chi_n(\lambda) := \prod_{m \neq n} \frac{\lambda_m^\bullet - \lambda}{\varsigma_m(\lambda)}, \quad (5)$$

where $\chi_n(\lambda)$ is analytic on $\mathbb{C} \setminus \bigcup_{m \neq n} [\lambda_m^-, \lambda_m^+]$ —see [10, Section 12]. Further, by the definition of the standard root (1), the function

$$z \mapsto \varsigma_n(\tau_n + z\gamma_n/2) = -z\sqrt[n]{1 - z^{-2}\gamma_n/2}$$

is analytic on $\mathbb{C} \setminus [-1, 1]$. For $s \in (-1, 1)$, consider $z_s = s \pm i\varepsilon$ to conclude

$$\varsigma_n(\tau_n + (s \pm i0)\gamma_n/2) = \mp i\sqrt[n]{1 - s^2\gamma_n/2}.$$

In particular, if α_1 and α_2 are admissible paths from λ_n^- to λ_n^+ running on different sides of $[\lambda_n^-, \lambda_n^+]$, then the integrand takes the opposite sign on these paths and hence

$$\int_{\alpha_1} \omega = - \int_{\alpha_2} \omega.$$

On the other hand, as $\int_{\partial U_n} \omega = 0$, the integral is independent of the chosen path α_i and thus needs to be zero. ■

If we write the action variables as

$$J_{n,k} = \frac{1}{k\pi} \int_{\Gamma_n} \lambda^k \omega, \quad (6)$$

then the analyticity on W is evident. To proceed, we need to find a globally defined primitive of ω . So for $\varphi \in W$, we define on $(\mathbb{C} \setminus \bigcup_{n \in \mathbb{Z}} U_n) \times V_\varphi$ the mapping

$$F(\lambda, \psi) := \frac{1}{2} \left(\int_{\lambda_0^-(\psi)}^{\lambda} \omega(\mu, \psi) + \int_{\lambda_0^+(\psi)}^{\lambda} \omega(\mu, \psi) \right),$$

where the paths of integration are chosen to be admissible. These improper integrals exist, as for $\gamma_0 = 0$ the integrand is analytic on U_0 , while for $\gamma_0 \neq 0$ it is of the form $1/\sqrt{\lambda_0^\pm - \lambda}$ locally around λ_0^\pm . By Lemma 2.2, they are also independent of the chosen admissible path. Hence, F is well defined and one has

$$F(\lambda, \psi) = \int_{\lambda_0^-(\psi)}^{\lambda} \omega(\mu, \psi) = \int_{\lambda_0^+(\psi)}^{\lambda} \omega(\mu, \psi).$$

Even though the eigenvalues λ_0^\pm are, due to their lexicographical ordering, not even continuous on W , the mapping F turns out to be differentiable.

Lemma 2.3. For every $\varphi \in W$, we have that

- (i) F is analytic on $(\mathbb{C} \setminus \bigcup_{n \in \mathbb{Z}} U_n) \times V_\varphi$ and
- (ii) $F(\cdot, \varphi)$ can be continuously extended onto $\mathbb{C} \setminus \bigcup_{\gamma_n \neq 0} (\lambda_n^-, \lambda_n^+)$ with

$$F(\lambda_n^+, \varphi) = F(\lambda_n^-, \varphi), \quad n \in \mathbb{Z}.$$

- (iii) If, in addition, $\varphi \in L_r^2$, then locally around $[\lambda_n^-, \lambda_n^+]$

$$F(\lambda) = l_n(\lambda) - i n \pi, \quad l_n(\lambda) = \log \frac{(-1)^n}{2} (\Delta(\lambda) + \sqrt{\Delta^2(\lambda) - 4}).$$

In particular, for any real $\lambda_n^- < \lambda < \lambda_n^+$,

$$F(\lambda \pm i0) = \pm f_n(\lambda) - i n \pi, \quad f_n(\lambda) = \cosh^{-1}((-1)^n \Delta(\lambda)/2).$$

Clearly, f_n is continuous on $[\lambda_n^-, \lambda_n^+]$, strictly positive on $(\lambda_n^-, \lambda_n^+)$, and vanishes at the boundary points.

- (iv) At the zero potential, one has $F(\lambda, 0) = -i\lambda$. □

Proof. The proof of (i) is standard but a bit technical and can be found in Appendix 1, and claim (ii) follows immediately from the properties of ω .

(iii) Note that locally around $[\lambda_n^-, \lambda_n^+]$ both l_n and F are primitives of ω and hence are identical up to a additive constant which may depend on φ . Clearly, $l_n(\lambda_n^\pm) = 0$. On the

other hand, since $\int_{\lambda_k^-}^{\lambda_k^+} \omega = 0$ for any k ,

$$F(\lambda_n^\pm) = \sum_{k=0}^{n-1} \int_{\lambda_k^\pm}^{\lambda_{k+1}^\pm} \omega, \quad F(\lambda_{-n}^\pm) = \sum_{k=0}^{n-1} \int_{\lambda_{-k}^\pm}^{\lambda_{-k-1}^\pm} \omega, \quad n > 0.$$

As $i(-1)^k \sqrt{\Delta^2(\lambda) - 4} > 0$ for $\lambda_k^+ < \lambda < \lambda_{k+1}^-$ —see [10, Section 12]—we find

$$\int_{\lambda_k^+}^{\lambda_{k+1}^-} \omega = i(-1)^k \int_{\lambda_k^+}^{\lambda_{k+1}^-} \frac{\dot{\Delta}(\lambda)}{\sqrt{4 - \Delta^2(\lambda)}} d\lambda = i(-1)^k \arcsin \frac{\Delta(\lambda)}{2} \Big|_{\lambda_k^+}^{\lambda_{k+1}^-} = -i\pi.$$

Consequently, $F(\lambda_n^\pm) = -in\pi$ and $F - l_n \equiv -in\pi$. Finally, $\pm(-1)^n \sqrt{\Delta^2(\lambda \pm i0) - 4} > 0$ for $\lambda_n^- < \lambda < \lambda_n^+$ hence

$$F(\lambda \pm i0) = \log \frac{1}{2}((-1)^n \Delta(\lambda) \pm \sqrt{\Delta^2(\lambda) - 4}) - in\pi = \pm f_n(\Delta(\lambda)/2) - in\pi.$$

(iv) At the zero potential, $\omega(\lambda, 0) = -i d\lambda$ which is evident from the product representation (4) of ω . ■

Given $\varphi \in W$, we can integrate by parts in (6) to obtain

$$J_{n,k} = \frac{1}{k\pi} \int_{\Gamma_n} \lambda^k \omega = -\frac{1}{\pi} \int_{\Gamma_n} \lambda^{k-1} F(\lambda) d\lambda.$$

Clearly, $J_{n,k}$ vanishes if $\gamma_n = 0$ in view of the analyticity of the integrand. Further, provided φ is of real type, then we may shrink the contour of integration to $[\lambda_n^-, \lambda_n^+]$, and use the properties of F to the effect that

$$J_{n,k} = \frac{2}{\pi} \int_{\lambda_n^-}^{\lambda_n^+} \lambda^{k-1} f_n(\lambda) d\lambda. \quad (7)$$

Thus, on L_r^2 all actions are real and those on odd levels are nonnegative. Moreover, by the mean value theorem,

$$J_{n,2m+1} = \zeta_{n,m}^{2m} I_n, \quad (8)$$

for some $\zeta_{n,m} \in [\lambda_n^-, \lambda_n^+]$. Recall that $\lambda_n^\pm \sim n\pi$ hence for any $m \geq 0$, we have $\zeta_{n,m}^{2m} \sim (n\pi)^{2m}$ asymptotically as $|n| \rightarrow \infty$. A quantitative estimate of the high-level actions $J_{n,2m+1}$ in terms of the actions I_n will be obtained in the next section.

Finally, consider the case of a potential with only finitely many open gaps. Then F is analytic outside some sufficiently large circle and thus admits a Laurent expansion around zero. The coefficients of this expansion turn out to be the Hamiltonians of the NLS hierarchy.

Lemma 2.4. For any finite gap potential $\varphi \in W$, there exists $\Lambda > 0$ such that

$$F(\lambda, \varphi) = -i\lambda - \frac{\mathcal{H}_1(\varphi)}{2i\lambda} + \sum_{n \geq 2} \frac{\mathcal{H}_n(\varphi)}{(2i\lambda)^n}, \quad |\lambda| > \Lambda. \quad \square$$

Proof. We attribute the claim to the asymptotic expansion of $\cosh^{-1}(\Delta/2)$ along the real axis—see [10, 17]. By the basic estimates for the discriminant

$$\Delta(i\tau, \varphi) = 2 \cosh \tau + o(e^\tau), \quad \tau \rightarrow \infty.$$

Hence, on an open neighborhood U of $i[\tau_0, \infty)$ for $\tau_0 > 0$ sufficiently large $\cosh^{-1}(\Delta(\lambda, \varphi)/2)$ is well defined. Here \cosh^{-1} denotes the principal branch of the inverse of \cosh which is defined on $\mathbb{C} \setminus (-\infty, 1)$ and real valued on $[1, \infty)$. On this neighborhood U , the λ -derivatives of $\cosh^{-1}(\Delta/2)$ and F coincide except for possibly the sign of

$$\Re \sqrt{\Delta^2(i\tau) - 4}, \quad \tau \geq \tau_0.$$

This sign is locally constant in φ provided $\tau \geq \tau_0$ and, as the straight line $\{t\varphi : 0 \leq t \leq 1\}$ is compact in L_c^2 , it can be determined by deforming φ to the zero potential. With

$$\sqrt{\Delta^2(i\tau, 0) - 4} = 2i \prod_{m \in \mathbb{Z}} \frac{\zeta_m(i\tau, 0)}{\pi_m} = 2i \prod_{m \in \mathbb{Z}} \frac{m\pi - i\tau}{\pi_m} = 2 \sinh \tau,$$

we conclude the sign is positive on U , and consequently

$$\cosh^{-1} \frac{\Delta(\lambda, \varphi)}{2} = F(\lambda, \varphi) + c(\varphi), \quad \lambda \in U,$$

with an analytic function $c: W \rightarrow \mathbb{C}$. For $\varphi \in W$ fixed, the right-hand side is analytic on $\mathbb{C} \setminus \bigcup_{\gamma_n \neq 0} [\lambda_n^-, \lambda_n^+]$ and continuous on $\mathbb{C} \setminus \bigcup_{\gamma_n \neq 0} (\lambda_n^-, \lambda_n^+)$, hence the left-hand side extends uniquely from U onto the same domain. Moreover, F vanishes at λ_0^\pm , while

$$\cosh^{-1} \frac{\Delta(\lambda_0^\pm, \varphi)}{2} \in \{im\pi : m \in \mathbb{Z}\}.$$

Thus, $c(\varphi) = im\pi$, and since c is continuous, this m is fixed on all of W . Evaluating at the zero potential reveals $\cosh^{-1}(\Delta(\lambda, 0)/2) = -i\lambda = F(\lambda, 0)$ hence $c \equiv 0$.

If φ is a finite gap potential, then $\cosh^{-1}(\Delta/2) = F$ is analytic outside a sufficiently large circle enclosing all open gaps, and has the following asymptotic expansion along the real axis [17]

$$\cosh^{-1} \frac{\Delta(\lambda, \varphi)}{2} = -i\lambda - \frac{\mathcal{H}_1(\varphi)}{2i\lambda} + \sum_{n \geq 2} \frac{\mathcal{H}_n(\varphi)}{(2i\lambda)^n}, \quad \lambda \rightarrow \pm\infty.$$

This proves the claim. ■

3 Localizing the Zakharov–Shabat Spectrum

The goal for this section is to provide a sufficiently accurate localization of the spectrum of the Zakharov–Shabat operator

$$L(\varphi) = \begin{pmatrix} i & \\ & -i \end{pmatrix} \frac{d}{dx} + \begin{pmatrix} & \varphi_- \\ \varphi_+ & \end{pmatrix},$$

allowing to quantify the asymptotic relation $J_{n,2m+1} \sim (n\pi)^{2m} I_n$. Since one can translate the spectrum of φ without changing the L^2 -norm, one cannot obtain a uniform localization on bounded subsets of L_c^2 . Instead, we have to impose some regularity on φ .

Theorem 3.1. Suppose $\varphi \in H_c^1$, then for each $\langle n \rangle \geq 8\|\varphi\|_1^2$,

$$|\lambda_n^\pm - n\pi| \leq \frac{\|\varphi\|_1^2}{\langle n \rangle} + \frac{\sqrt{2}\|\varphi\|_1}{\langle 2n \rangle} \leq \pi/5,$$

while the remaining eigenvalues are contained in the box

$$\{\lambda \in \mathbb{C} : |\Re \lambda| \leq (8\|\varphi\|_1^2 - 1/2)\pi, |\Im \lambda| \leq \|\varphi\|_1\}.$$

□

Remark. In [22], Li and McLaughlin obtained a localization for φ in H_c^1 where the bound on the threshold of $\langle n \rangle$ is exponential in the norm of φ . With a focus on lowering the regularity assumptions on φ rather than improving the threshold for φ smooth, this result was gradually improved by several authors—see, for example, Mityagin [24] and the references therein. The novelty of Theorem 3.1 consists in providing a threshold for $\langle n \rangle$ which is quadratic in the norm of φ . □

The proof is based on a *Lyapunov–Schmidt decomposition* introduced by Kapeler and Mityagin [12]—see also [9, 11]: For the zero potential each $n\pi$, $n \in \mathbb{Z}$, is a double eigenvalue of L with eigenfunctions $e_n^+ := (0, e^{in\pi x})$ and $e_n^- := (e^{-in\pi x}, 0)$. Thus, for a nonzero potential, provided $|n|$ is sufficiently large, we expect exactly two eigenvalues which are close to $n\pi$ and whose eigenfunctions are close to the linear span of e_n^+ and e_n^- . This suggest to separate these modes from the others by a Lyapunov–Schmidt reduction.

To set the stage, we cover the complex plane with the closed strips

$$\mathfrak{U}_n := \{\lambda : |\Re \lambda - n\pi| \leq \pi/2\},$$

and consider the Hilbert space of complex 2-periodic L^2 -functions

$$L_*^2 = \mathcal{P}_n \oplus \mathcal{Q}_n = \text{sp}\{e_n^+, e_n^-\} \oplus \overline{\text{sp}}\{e_k^+, e_k^- : k \neq n\}.$$

The orthogonal projections onto \mathcal{P}_n and \mathcal{Q}_n are denoted by P_n and Q_n , respectively.

To proceed, write the eigenvalue equation $Lf = \lambda f$ in the form

$$A_\lambda f = \Phi f,$$

with the unbounded operators

$$A_\lambda := \lambda - \begin{pmatrix} i & \\ & -i \end{pmatrix} \frac{d}{dx}, \quad \Phi := \begin{pmatrix} & \varphi_- \\ \varphi_+ & \end{pmatrix}.$$

Since A_λ leaves the spaces \mathcal{P}_n and \mathcal{Q}_n invariant, by writing

$$f = u + v = P_n f + Q_n f,$$

we can decompose the equation $A_\lambda f = \Phi f$ into the two equations

$$A_\lambda u = P_n \Phi(u + v), \quad A_\lambda v = Q_n \Phi(u + v),$$

called the P - and the Q -equation, respectively.

We first consider the Q -equation on \mathfrak{U}_n . One checks that for $m \neq n$

$$\min_{\lambda \in \mathfrak{U}_n} |\lambda - m\pi| \geq |n - m| \geq 1,$$

hence it follows with $A_\lambda e_m^\pm = (\lambda - m\pi)e_m^\pm$ that the restriction of A_λ to \mathcal{Q}_n is boundedly invertible for all $\lambda \in \mathfrak{U}_n$. Therefore, multiplying the Q -equation from the left by ΦA_λ^{-1} , gives

$$\Phi v = T_n \Phi(u + v),$$

with $T_n := \Phi A_\lambda^{-1} Q_n$. The latter identity may be written as

$$(I - T_n)\Phi v = T_n \Phi u,$$

hence solving the Q equation amounts to inverting $(I - T_n)$.

We consider operator norms induced by *shifted weighted norms* [13, 26]. Let H_i^w denote the space of complex 2-periodic functions $u = \sum_{m \in \mathbb{Z}} u_m e_m$ equipped with the i -shifted H^w -norm given by

$$\|u\|_{w,i}^2 := \|ue_i\|_w^2 = \sum_{m \in \mathbb{Z}} w_{m+i}^2 |u_m|^2, \quad e_m(x) := e^{im\pi x}.$$

On the space $H_{i,C}^w := H_i^w \times H_i^w$ of 2-periodic functions with values in \mathbb{C}^2 ,

$$f = (f_-, f_+) = \sum_{n \in \mathbb{Z}} (f_n^- e_n^- + f_n^+ e_n^+) = \sum_{n \in \mathbb{Z}} (f_n^- e_{-n}^-, f_n^+ e_n),$$

the i -shifted norm is defined by

$$\|f\|_{w;i}^2 := \|f_-\|_{w;-i}^2 + \|f_+\|_{w;i}^2 = \sum_{m \in \mathbb{Z}} w_{m+i}^2 (|f_m^-|^2 + |f_m^+|^2).$$

Lemma 3.2. If $\varphi \in H_c^w$ with $w \in \mathcal{M}$, then for any $n, i \in \mathbb{Z}$ and any $\lambda \in \mathfrak{U}_n$,

$$T_n = \Phi A_\lambda^{-1} Q_n: H_{i,c}^w \rightarrow H_{-i,c}^w$$

is bounded and satisfies the estimate

$$\|T_n f\|_{w;i} \leq 2 \|\varphi\|_w \|f\|_{w;-i}. \quad \square$$

Proof. Write $T_n f = \Phi g$ with $g = A_\lambda^{-1} Q_n f$. Since the restriction of A_λ to \mathcal{Q}_n is boundedly invertible, the function

$$g = A_\lambda^{-1} Q_n f = \sum_{m \neq n} \left(\frac{f_m^-}{\lambda - m\pi} e_m^- + \frac{f_m^+}{\lambda - m\pi} e_m^+ \right) = (g_-, g_+)$$

is well defined. By Hlder's inequality, we obtain for the weighted ℓ^1 -norm

$$\|g_+ e_{-i}\|_{\ell_w^1} = \sum_{m \neq n} \frac{w_{m-i} |f_m^+|}{|\lambda - m\pi|} \leq \left(\sum_{m \neq n} \frac{1}{|n - m|^2} \right)^{1/2} \|f_+\|_{w;-i} < 2 \|f_+\|_{w;-i},$$

uniformly for $\lambda \in \mathfrak{U}_n$; similarly $\|g_- e_i\|_{\ell_w^1} \leq 2 \|f_-\|_{w;i}$. Since

$$\|T_n f\|_{w;i}^2 = \|\Phi g\|_{w;i}^2 = \|\varphi_- g_+ e_{-i}\|_w^2 + \|\varphi_+ g_- e_i\|_w^2,$$

we can use standard inequalities for the convolution of sequences to obtain

$$\|T_n f\|_{w;i}^2 \leq \|\varphi\|_w^2 (\|g_+ e_{-i}\|_{\ell_w^1}^2 + \|g_- e_i\|_{\ell_w^1}^2) \leq 4 \|\varphi\|_w^2 \|f\|_{w;-i}^2. \quad \blacksquare$$

Note that $T_n f$ is estimated with a shifted H^w -norm where the sign of the shift is opposite to the sign of the shifted H^w -norm of f . This fact will be crucial in the following. In particular, T_n^2 is bounded on $H_{i,c}^w$ and it turns out that $\|T_n^2\|_{w;n} = o(1)$ as $|n| \rightarrow \infty$. Using a Neumann series, we then obtain the bounded invertibility of $(I - T_n)$ for $|n|$ sufficiently large, which solves the Q -equation. To exploit the regularity assumption $\varphi \in H_c^1$ of Theorem 3.1, we restrict ourselves to the subclass \mathcal{M}^1 of weights which have at least a linearly growing factor—see also [7, 9] for weights with factors $\langle n \rangle^\delta$, $0 < \delta < \frac{1}{2}$.

Lemma 3.3. If $\varphi \in H_c^w$ with $w \in \mathcal{M}^1$, then for any $n \in \mathbb{Z}$ and any $\lambda \in \mathfrak{U}_n$,

$$\|T_n^2\|_{w;n} \leq \frac{4}{\langle n \rangle} \|\varphi\|_w^2. \quad \square$$

Proof. As in the preceding lemma, write $T_n^2 f = \Phi g$ with

$$g := A_\lambda^{-1} Q_n \Phi A_\lambda^{-1} Q_n f.$$

A straightforward computation yields

$$g = \sum_{k,l \neq n} \left(\frac{\varphi_{k+l}^-}{\lambda - k\pi} \frac{f_l^+}{\lambda - l\pi} e_k^- + \frac{\varphi_{k+l}^+}{\lambda - k\pi} \frac{f_l^-}{\lambda - l\pi} e_k^+ \right) = (g_-, g_+),$$

and our aim is to estimate the weighted ℓ^1 -norms $\|g_+ e_{-n}\|_{\ell_w^1}$ and $\|g_- e_n\|_{\ell_w^1}$. By assumption $w = \langle n \rangle \cdot v$ with some submultiplicative weight v , hence

$$w_{k-n} \leq \frac{\langle k-n \rangle}{\langle k+l \rangle \langle l+n \rangle} w_{k+l} w_{-l-n}, \quad k, l \in \mathbb{Z}.$$

Consequently, for any $\lambda \in \mathfrak{U}_n$

$$\|g_+ e_{-n}\|_{\ell_w^1} \leq \sum_{k,l \neq n} \frac{\langle k-n \rangle}{\langle k+l \rangle |n-k| \langle l+n \rangle |n-l|} w_{k+l} |\varphi_{k+l}^+| w_{-l-n} |f_l^-|,$$

and with Cauchy–Schwarz and Young’s inequality for convolutions,

$$\leq \left(\sum_{k,l \neq n} \frac{\langle k-n \rangle^2}{\langle k+l \rangle^2 |n-k|^2 \langle l+n \rangle^2 |n-l|^2} \right)^{1/2} \|\varphi\|_w \|f_-\|_{w;-n}.$$

One further checks that

$$\sum_{k \neq n} \frac{\langle k-n \rangle^2}{\langle k+l \rangle^2 |n-k|^2} \leq 32/5, \quad \sum_{l \neq n} \frac{1}{\langle l+n \rangle^2 |n-l|^2} \leq \frac{5/2}{\langle n \rangle^2}.$$

Hence, we obtain for $\|g_+ e_{-n}\|_{\ell_w^1}$ and similarly for $\|g_- e_n\|_{\ell_w^1}$,

$$\|g_+ e_{-n}\|_{\ell_w^1} \leq \frac{4}{\langle n \rangle} \|\varphi\|_w \|f_-\|_{w;-n}, \quad \|g_- e_n\|_{\ell_w^1} \leq \frac{4}{\langle n \rangle} \|\varphi\|_w \|f_+\|_{w;n}.$$

Finally, with $\|T_n^2 f\|_{w;n} = \|\Phi g\|_{w;n}$, this gives

$$\|T_n^2 f\|_{w;n}^2 \leq \|\varphi\|_w^2 \left(\|g_+ e_{-n}\|_{\ell_w^1}^2 + \|g_- e_n\|_{\ell_w^1}^2 \right) \leq \frac{16}{\langle n \rangle^2} \|\varphi\|_w^4 \|f\|_{w;n}^2. \quad \blacksquare$$

Consequently, T_n^2 is a $1/2$ -contraction if $\langle n \rangle \geq 8\|\varphi\|_w^2$. In view of

$$\hat{T}_n := (I - T_n)^{-1} = (I + T_n)(I - T_n^2)^{-1},$$

one then finds a unique solution

$$\Phi v = \hat{T}_n T_n \Phi u$$

of the Q -equation. In turn, as $I + \hat{T}_n T_n = \hat{T}_n$, the P -equation yields

$$A_\lambda u = P_n(I + \hat{T}_n T_n) \Phi u = P_n \hat{T}_n \Phi u.$$

Writing the latter as

$$S_n u = 0, \quad S_n: \mathcal{P}_n \rightarrow \mathcal{P}_n, \quad u \mapsto (A_\lambda - P_n \hat{T}_n \Phi) u,$$

we immediately conclude that there exists the following relationship.

Lemma 3.4. For $\varphi \in H_c^1$ and $\langle n \rangle \geq 8\|\varphi\|_1^2$, a complex number $\lambda \in \mathfrak{U}_n$ is an eigenvalue of L if and only if the determinant of S_n vanishes. \square

Proof. Suppose $Lf = \lambda f$, then by the preceding discussion $S_n u = 0$. Conversely, define for any $u \in \mathcal{P}_n$,

$$v = A_\lambda^{-1} Q_n \hat{T}_n \Phi u \in \mathcal{Q}_n.$$

Then $\Phi v = T_n \hat{T}_n \Phi u$ is well defined, and it follows with $\hat{T}_n = I + T_n \hat{T}_n$ that

$$A_\lambda v = Q_n \hat{T}_n \Phi u = Q_n \Phi u + Q_n T_n \hat{T}_n \Phi u = Q_n \Phi(u + v),$$

so the Q -equation is automatically satisfied. Moreover, if $S_n u = 0$, then

$$A_\lambda u = P_n \hat{T}_n \Phi u = P_n(I + T_n \hat{T}_n) \Phi u = P_n \Phi(u + v).$$

Hence also the P -equation is satisfied, and λ is an eigenvalue of L with eigenfunction $f = u + v$. \blacksquare

Recall that P_n is the orthogonal projection onto the 2D space \mathcal{P}_n . The matrix representation of an operator B on \mathcal{P}_n is given by

$$\left((B e_n^\pm, e_n^\pm) \right)_{\pm, \pm}.$$

Therefore, we find for S_n the representation

$$A_\lambda = \begin{pmatrix} \lambda - n\pi & \\ & \lambda - n\pi \end{pmatrix}, \quad P_n \hat{T}_n \Phi = \begin{pmatrix} a_n^+ & b_n^+ \\ b_n^- & a_n^- \end{pmatrix},$$

with the coefficients of the latter matrix given by

$$a_n^\pm := \langle \hat{T}_n \Phi e_n^\pm, e_n^\pm \rangle, \quad b_n^\pm := \langle \hat{T}_n \Phi e_n^\mp, e_n^\pm \rangle.$$

We point out that these coefficients depend on λ and φ . It has been observed in [7, 16] that these coefficients reflect certain symmetries of the Fourier coefficients of φ . We only need the fact that a_n^+ and a_n^- coincide.

Lemma 3.5. Suppose $\varphi \in H_c^1$ and $\langle n \rangle \geq 8\|\varphi\|_1^2$, then for any $\lambda \in \mathfrak{U}_n$,

$$a_n^+(\lambda) = a_n^-(\lambda) \equiv a_n(\lambda). \quad \square$$

Proof. Recall that $T_n = \Phi A_\lambda^{-1} Q_n$, and denote by $\bar{B}u := \overline{B\bar{u}}$ the complex conjugation of operators. From evaluating the bounded diagonal operators A_λ^{-1} and Q_n at e_m^\pm , and using the identity $\overline{e_m^\pm} = P e_m^\mp$, we conclude

$$(A_\lambda^{-1})^* = P \overline{A_\lambda^{-1}} P = A_\lambda^{-1}, \quad Q_n^* = P \bar{Q}_n P = Q_n, \quad P := \begin{pmatrix} & 1 \\ 1 & \end{pmatrix}.$$

Since A_λ^{-1} leaves Q_n invariant, and $P^2 = I$, we find $(A_\lambda^{-1} Q_n)^* = P \overline{A_\lambda^{-1} Q_n} P$. With $\Phi^* = P \bar{\Phi} P$ this gives

$$(T_n \Phi)^* = \Phi^* (A_\lambda^{-1} Q_n)^* \Phi^* = P \overline{T_n \Phi} P.$$

Inspecting the Neumann expansion of $\hat{T}_n \Phi$ yields $(\hat{T}_n \Phi)^* = P \overline{\hat{T}_n \Phi} P$, thus

$$a_n^+ = \langle \hat{T}_n \Phi e_n^+, e_n^+ \rangle = \langle e_n^+, (\hat{T}_n \Phi)^* e_n^+ \rangle = \langle P e_n^+, \overline{\hat{T}_n \Phi} P e_n^+ \rangle = \langle \hat{T}_n \Phi e_n^-, e_n^- \rangle = a_n^-. \quad \blacksquare$$

It follows that S_n may be written as

$$S_n(\lambda) = \begin{pmatrix} \lambda - n\pi - a_n & -b_n^+ \\ -b_n^- & \lambda - n\pi - a_n \end{pmatrix}.$$

As T_n and Φ are anti-diagonal while I and T_n^2 are diagonal, all even terms $\langle T_n^{2k} \Phi e_n^+, e_n^+ \rangle$ in the expansion of a_n vanish. Using $\hat{T}_n = (I + T_n)(I - T_n^2)^{-1}$, we thus conclude

$$a_n = \langle \hat{T}_n \Phi e_n^+, e_n^+ \rangle = \langle T_n (I - T_n^2)^{-1} \Phi e_n^+, e_n^+ \rangle.$$

On the other hand, all odd terms in the expansion of b_n vanish, such that

$$b_n^\pm - \varphi_{2n}^\pm = \langle (\hat{T}_n - I) \Phi e_n^\mp, e_n^\pm \rangle = \langle T_n^2 (I - T_n^2)^{-1} \Phi e_n^\mp, e_n^\pm \rangle.$$

We introduce the following notion for the sup-norm of a complex-valued function on a domain $U \subset \mathbb{C}$,

$$|f|_U := \sup_{\lambda \in U} |f(\lambda)|.$$

Lemma 3.6. If $\varphi \in H_c^w$ with $w \in \mathcal{M}^1$, then for any $\langle n \rangle \geq 8\|\varphi\|_w^2$ the coefficients a_n and b_n^\pm are analytic functions on \mathfrak{U}_n with bounds

$$|a_n|_{\mathfrak{U}_n} \leq \frac{1}{\langle n \rangle} \|\varphi\|_w^2, \quad w_{2n}|b_n^\pm - \varphi_{2n}^\pm|_{\mathfrak{U}_n} \leq \frac{8}{\langle n \rangle} \|\varphi\|_w^2 \|\varphi_\pm\|_w. \quad \square$$

Proof. Since $\|T_n^2\|_{w;n} \leq 1/2$, the series expansions of a_n and b_n^\pm converge uniformly on \mathfrak{U}_n to analytic functions. Let $u = (I - T_n^2)^{-1} \Phi e_n^+$, then

$$\|u\|_{w;n} \leq \|(I - T_n^2)^{-1}\|_{w;n} \|\Phi e_n^+\|_{w;n} \leq 2\|\varphi_- e_n\|_{w;-n} = 2\|\varphi_-\|_w,$$

and with the series expansion $u = \sum_{m \in \mathbb{Z}} u_m e_m^-$ we may write

$$a_n = \langle T_n u, e_n^+ \rangle = \sum_{m \neq n} \frac{\varphi_{n+m}^-}{\lambda - m\pi} u_m.$$

As $|n - m| \leq |n|$ implies $|n + m| \geq 2|n| - |n - m| \geq |n|$, this gives

$$|a_n|_{\mathfrak{U}_n} \leq \sum_{m \neq n} \frac{1}{\langle n + m \rangle^2 |n - m|} w_{n+m} |\varphi_{n+m}^-| \cdot w_{n+m} |u_m| \leq \frac{1}{\langle n \rangle} \|\varphi_-\|_w \|u\|_{w;n} \leq \frac{2}{\langle n \rangle} \|\varphi_-\|_w^2.$$

Using the representation $a_n = \langle T_n (I - T_n^2) \Phi e_n^-, e_n^- \rangle$, we similarly obtain

$$|a_n|_{\mathfrak{U}_n} \leq \frac{2}{\langle n \rangle} \|\varphi_+\|_w^2.$$

Summing both estimates up gives the first bound. To obtain the second bound, we note that $b_n^- - \varphi_{2n}^- = \langle T_n^2 u, e_n^- \rangle$. Since $\langle f, e_n^- \rangle = \langle f e_{-n}, e_{2n}^- \rangle$ for any function $f \in L_c^2$, we conclude

$$w_{2n} |\langle T_n^2 u, e_n^- \rangle| \leq \|T_n^2 u\|_{w;n} \leq \frac{8}{\langle n \rangle} \|\varphi\|_w^2 \|\varphi_-\|_w.$$

The proof for b_n^+ is the same. ■

In consequence, the determinant of S_n

$$\det S_n = (\lambda - n\pi - a_n)^2 - b_n^+ b_n^-$$

is analytic on \mathfrak{U}_n and close to $(\lambda - n\pi)^2$ provided $\langle n \rangle$ is sufficiently large.

Lemma 3.7. Let $\varphi \in H_c^1$, then for any $\langle n \rangle \geq 8\|\varphi\|_1^2$, the determinant of S_n has exactly two complex roots ξ_+ , ξ_- in \mathfrak{U}_n , which are contained in the disk

$$D_n := \left\{ |\lambda - \pi n| \leq \frac{\|\varphi\|_1^2}{\langle n \rangle} + \frac{\sqrt{2}\|\varphi\|_1}{\langle 2n \rangle} \right\} \subset \left\{ |\lambda - n\pi| \leq \frac{\pi}{5} \right\},$$

and satisfy

$$|\xi^+ - \xi^-|^2 \leq 6|b_n^+ b_n^-|_{\mathfrak{U}_n}. \quad \square$$

Proof. The estimates of the preceding lemma give for $\langle n \rangle \geq 8\|\varphi\|_1^2$

$$|a_n|_{\mathfrak{U}_n} \leq \frac{\|\varphi\|_1^2}{\langle n \rangle}, \quad \langle 2n \rangle^2 |b_n^+ b_n^-|_{\mathfrak{U}_n} \leq \left(1 + \frac{8\|\varphi\|_1}{\langle n \rangle}\right)^2 \|\varphi_+\|_1 \|\varphi_-\|_1,$$

where we used $w_{2n}|b_n^\pm| \leq \|\varphi\|_w + w_{2n}|b_n^\pm - \varphi_{2n}^\pm|$. Therefore,

$$|a_n|_{\mathfrak{U}_n} + |b_n^+ b_n^-|_{\mathfrak{U}_n}^{1/2} \leq \inf_{\lambda \in \mathfrak{U}_n \setminus D_n} |\lambda - n\pi| = \frac{\|\varphi\|_1^2}{\langle n \rangle} + \sqrt{2} \frac{\|\varphi\|_1}{\langle 2n \rangle} \leq \pi/5.$$

It follows from Rouché's Theorem that the function $h = \lambda - n\pi - a_n$ has a single root in D_n , just as $\lambda - n\pi$. Furthermore, h^2 and $\det S_n$ have the same number of roots in D_n , namely two when counted with multiplicity, while $\det S_n$ clearly has no root in $\mathfrak{U}_n \setminus D_n$.

To estimate the distance of the roots, we write $\det S_n = g_+ g_-$ with

$$g_\pm = \lambda - \pi n - a_n \mp \sigma_n, \quad \sigma_n = \sqrt{b_n^+ b_n^-},$$

where the branch of the root is immaterial. Each root ξ of $\det S_n$ is either a root of g_+ or g_- , respectively, and thus satisfies $\xi = \pi n + a_n(\xi) \pm \sigma_n(\xi)$. Therefore,

$$|\xi_+ - \xi_-| \leq |a_n(\xi_+) - a_n(\xi_-)| + |\sigma_n(\xi_+) \pm \sigma_n(\xi_-)| \leq |\partial_\lambda a_n|_{D_n} |\xi_+ - \xi_-| + 2|\sigma_n|_{\mathfrak{U}_n}.$$

Since $\text{dist}(D_n, \partial \mathfrak{U}_n) \geq \pi/2 - \pi/5$, the claim follows with Cauchy's estimate

$$|\partial_\lambda a_n|_{D_n} \leq \frac{|a_n|_{\mathfrak{U}_n}}{\text{dist}(D_n, \partial \mathfrak{U}_n)} \leq \frac{1/8}{\pi/2 - \pi/5} \leq \frac{1}{6}. \quad \blacksquare$$

Proof of Theorem 3.1. For each $\langle n \rangle \geq 8\|\varphi\|_1^2$, Lemma 3.7 applies giving us the two roots ξ_+ and ξ_- of $\det S_n$ in $D_n \subset \mathfrak{U}_n$. As the strips \mathfrak{U}_n cover the complex plane, and since $\lambda_n^\pm \sim n\pi$

asymptotically as $n \rightarrow \pm\infty$ while there are no periodic eigenvalues in $\bigcup_{\langle n \rangle \geq 8\|\varphi\|_1^2} (\mathfrak{U}_n \setminus D_n)$, it follows by a standard counting argument that these roots have to be the periodic eigenvalues λ_n^\pm . In turn, the remaining eigenvalues have to be contained in the strip

$$\{\lambda \in \mathbb{C} : |\Re \lambda| \leq (8\|\varphi\|_1^2 - 1/2)\pi\}.$$

To obtain the estimate for the imaginary part, suppose that f is an L_c^2 normalized eigenfunction for λ , then

$$2\Im \lambda = \lambda - \bar{\lambda} = \langle Lf, f \rangle - \langle f, Lf \rangle = \langle (L - L^*)f, f \rangle.$$

Further, using the L^∞ -estimate $\|g\|_\infty \leq \sqrt{2}\|g\|_1$ for $g \in H^1$, we find

$$\|(L - L^*)f\|_0 \leq \sqrt{2}\|\varphi_+ - \varphi_-\|_1 \|f\|_0 \leq 2\|\varphi\|_1.$$

This completes the proof of the theorem. ■

Incidentally, we obtain the following estimate for the gap lengths, which we will use in Section 6.

Proposition 3.8. Suppose $\varphi \in H_c^w$ with $w \in \mathcal{M}^1$, then for any $\langle N \rangle \geq 8\|\varphi\|_w^2$,

$$\sum_{|n| \geq N} w_{2n}^2 |\gamma_n(\varphi)|^2 \leq 6\|R_N \varphi\|_w^2 + \frac{1152}{\langle N \rangle} \|\varphi\|_w^6,$$

where $R_N \varphi = \sum_{|n| \geq N} (\varphi_{2n}^- e_{-2n}, \varphi_{2n}^+ e_{2n})$. If, in addition, φ is in the complex neighborhood W of L_r^2 , then

$$\sum_{n \in \mathbb{Z}} w_{2n}^2 |\gamma_n(\varphi)|^2 \leq 265\pi^2 w^2 [16\|\varphi\|_w^2 (1 + \|\varphi\|_w^2) \|\varphi\|_w^2]. \quad \square$$

Proof. By Lemma 3.7, we have for any $\langle n \rangle \geq 8\|\varphi\|_w^2$ the estimate

$$|\gamma_n|^2 = |\lambda_n^+ - \lambda_n^-|^2 \leq 6|b_n^+ b_n^-|_{\mathfrak{U}_n} \leq 3(|b_n^+|_{\mathfrak{U}_n}^2 + |b_n^-|_{\mathfrak{U}_n}^2).$$

Using $|b_n^\pm|_{\mathfrak{U}_n} \leq |\varphi_{2n}^\pm| + |b_n^\pm - \varphi_{2n}^\pm|_{\mathfrak{U}_n}$, we thus find for any $\langle N \rangle \geq 8\|\varphi\|_w^2$,

$$\frac{1}{6} \sum_{|n| \geq N} w_{2n}^2 |\gamma_n(\varphi)|^2 \leq \sum_{|n| \geq N} w_{2n}^2 (|\varphi_{2n}^+|^2 + |\varphi_{2n}^-|^2 + |b_n^+ - \varphi_{2n}^+|_{\mathfrak{U}_n}^2 + |b_n^- - \varphi_{2n}^-|_{\mathfrak{U}_n}^2).$$

Further by Lemma 3.6, $w_{2n}^2 |\varphi_{2n}^\pm - b_n^\pm|_{\mathfrak{U}_n}^2 \leq 64\langle n \rangle^{-2} \|\varphi\|_w^4 \|\varphi_\pm\|_w^2$, hence

$$\frac{1}{6} \sum_{|n| \geq N} w_{2n}^2 |\gamma_n(\varphi)|^2 \leq \|R_N \varphi\|_w^2 + 64\|\varphi\|_w^6 \sum_{|n| \geq N} \frac{1}{\langle n \rangle^2},$$

and the first claim follows with $\sum_{|n| \geq N} 1/\langle n \rangle^2 \leq 3/\langle N \rangle$.

If additionally $\varphi \in W$, then each gap is contained in its isolating neighborhood U_n . Those are disjoint complex disks centered on the real line, whose diameters for $|n| < N$ sum up to at most $(2N - 1)\pi$ by Theorem 3.1. Therefore,

$$\sum_{|n| < N} w_{2n}^2 |\gamma_n(\varphi)|^2 \leq w_{2N-2}^2 \left(\sum_{|n| < N} |\gamma_n(\varphi)| \right)^2 \leq w_{2N-2}^2 (2N - 1)^2 \pi^2,$$

and choosing $N + 1 \geq 8\|\varphi\|_w^2 > N$ gives the second claim. \blacksquare

If, for $\langle n \rangle \geq 8\|\varphi\|_w^2$, we use

$$w_{2n}|b_n^\pm| \leq \|\varphi_\pm\|_w + \frac{8}{\langle n \rangle} \|\varphi\|_w^2 \|\varphi_\pm\| \leq 2\|\varphi_\pm\|_w,$$

then we obtain the *individual gap estimate*

$$w_{2n}|\gamma_n(\varphi)| \leq 4\|\varphi\|_w. \quad (9)$$

4 Estimating the Actions

As an immediate corollary to the localization obtained in the previous section, we obtain the following quantitative estimate of the high-level actions.

Proposition 4.1. If $\varphi \in H_r^1$, then for $\langle n \rangle \geq 8\|\varphi\|_1^2$ and $m \geq 0$,

$$J_{n,2m+1} = \zeta_{n,m}^{2m} I_n, \quad |\zeta_{n,m} - n\pi| \leq \frac{\|\varphi\|_1^2}{\langle n \rangle} + \frac{\sqrt{2}\|\varphi\|_1}{\langle 2n \rangle}.$$

In particular, if $n \neq 0$ and $\langle n \rangle \geq 8\|\varphi\|_1^2$, then

$$2^{-m} \langle 2n\pi \rangle^{2m} I_n \leq 4^m J_{n,2m+1} \leq \langle 2n\pi \rangle^{2m} I_n,$$

while the remaining actions for all $\langle n \rangle < 8\|\varphi\|_1^2$ satisfy

$$4^m |J_{n,2m+1}| \leq (16\pi)^{2m} \|\varphi\|_1^{4m} I_n. \quad \square$$

Proof. Recall from (8) that $J_{n,2m+1} = \zeta_{n,m}^{2m} I_n$ with $\zeta_{n,m} \in [\lambda_n^-, \lambda_n^+]$. Provided $\langle n \rangle \geq 8\|\varphi\|_1^2$, the estimate of $|\zeta_{n,m} - n\pi|$ follows from Theorem 3.1. If additionally $n \neq 0$, then $\langle 2n \rangle \geq 3\langle n \rangle/2$ and hence $|\zeta_{n,m} - n\pi| \leq 1/2$. In consequence,

$$\frac{1}{\sqrt{2}} \langle 2n\pi \rangle \leq 2|\zeta_{n,m}| \leq \langle 2n\pi \rangle, \quad n \neq 0.$$

Conversely, if $\langle n \rangle < 8\|\varphi\|_1^2$, then $|\zeta_{n,m}| \leq 8\pi\|\varphi\|_1^2$. \blacksquare

In the sequel, we use Proposition 4.1 to obtain an estimate of

$$\|I(\varphi)\|_{\ell_{2m}^1} = \sum_{n \in \mathbb{Z}} \langle 2n\pi \rangle^{2m} |I_n|$$

in terms of $\sum_{n \in \mathbb{Z}} J_{n,2m+1}$ and a remainder depending solely on $\|\varphi\|_1$. The trace formula and the polynomial structure of the Hamiltonians then allows us to obtain the first part of Theorem 1.2.

Lemma 4.2. For every $m \geq 1$,

$$\|I(\varphi)\|_{\ell_{2m}^1} \leq \langle 16\pi \rangle^{2m} (1 + \|\varphi\|_1)^{4m} \|\varphi\|_0^2 + (-1)^{m+1} 2^m \mathcal{H}_{2m+1}(\varphi),$$

uniformly for all $\varphi \in H_r^m$. □

Proof. Choose $N + 1 \geq 8\|\varphi\|_1^2 > N$, then by Proposition 4.1, the trace formula (3), and the positivity of the actions

$$\sum_{|n| > N} \langle 2n\pi \rangle^{2m} I_n \leq 8^m \sum_{n \in \mathbb{Z}} J_{n,2m+1} = (-1)^{m+1} 2^m \mathcal{H}_{2m+1}.$$

On the other hand, by our choice of N and the trace formula (2)

$$\sum_{|n| \leq N} \langle 2n\pi \rangle^{2m} I_n \leq \langle 2N\pi \rangle^{2m} \sum_{n \in \mathbb{Z}} I_n \leq \langle 16\pi \rangle^{2m} (1 + \|\varphi\|_1)^{4m} \|\varphi\|_0^2. \quad \blacksquare$$

We denote the two components of $\varphi \in L_r^2$ by $\varphi = (\psi, \bar{\psi})$, and to simplify notation write $\psi_{(m)} = \partial_x^m \psi$ such that

$$\int_{\mathbb{T}} |\psi_{(m)}|^2 dx = \frac{1}{2} \|\varphi_{(m)}\|_0^2.$$

Proof of Theorem 1.2(i) for $m = 1$. By the preceding lemma,

$$\sum_{n \in \mathbb{Z}} \langle 2n\pi \rangle^2 I_n \leq \langle 16\pi \rangle^2 (1 + \|\varphi\|_1)^4 \|\varphi\|_0^2 + 2\mathcal{H}_3(\varphi).$$

The Hamiltonian $\mathcal{H}_3(\varphi) = \int_{\mathbb{T}} (|\psi_x|^2 + |\psi|^4) dx$ may be estimated by

$$|\mathcal{H}_3(\varphi)| \leq \frac{1}{2} \|\varphi_x\|_0^2 + \frac{1}{2} \|\psi\|_\infty^2 \|\varphi\|_0^2 \leq \frac{1}{2} (1 + \|\varphi\|_0^2) \|\varphi\|_1^2,$$

where we used the L^∞ -estimate $\|\psi\|_\infty^2 \leq 2\|\psi\|_1^2 = \|\varphi\|_1^2$. So we arrive at

$$\|I(\varphi)\|_{\ell_2^1} \leq (32\pi)^2 (1 + \|\varphi\|_1)^4 \|\varphi\|_1^2. \quad \blacksquare$$

To proceed with the general case $m \geq 2$, we note that on H_r^m ,

$$(-1)^{m+1} \mathcal{H}_{2m+1}(\varphi) = \int_{\mathbb{T}} (|\psi_{(m)}|^2 + p_{2m}(\psi, \bar{\psi}, \dots, \psi_{(m-1)}, \bar{\psi}_{(m-1)})) \, dx,$$

with p_{2m} being a homogenous polynomial of degree $2m + 2$ when ψ , $\bar{\psi}$, and ∂_x each count as one degree. Further, the degree of each monomial of p_{2m} with respect to ψ equals the degree with respect to $\bar{\psi}$ —see Corollary A.4 from the appendix. Consequently, each monomial q of p_{2m} may be estimated by

$$|q| \leq c_q |\psi|^{\mu_0} |\psi_x|^{\mu_1} \cdots |\psi_{(m-1)}|^{\mu_{m-1}},$$

with some positive constant c_q and integers μ_0, \dots, μ_{m-1} . It turns out to be convenient to use exponents which are multiples of 2, that is,

$$|q| \leq c_q |\psi|^{2\sigma_0} |\psi_x|^{2\sigma_1} \cdots |\psi_{(m-1)}|^{2\sigma_{m-1}},$$

with $\sigma_i \in \mathbb{Z}_{\geq 0}/2$. Since q has degree $2m + 2$, we have

$$\sum_{0 \leq i \leq m-1} 2\sigma_i(1+i) = 2m + 2,$$

and as the degree with respect to ψ and $\bar{\psi}$ is the same, $\sum_{0 \leq i \leq m-1} \sigma_i$ is an integer even though $\sigma_i \in \mathbb{Z}_{\geq 0}/2$. Denote by $\mathcal{I}_{2m+2} \subset (\mathbb{Z}_{\geq 0}/2)^m$ the set of all multi-indices $\sigma = (\sigma_i)_{0 \leq i \leq m-1}$ satisfying the constraints

$$\sum_{0 \leq i \leq m-1} 2(1+i)\sigma_i = 2m + 2, \quad |\sigma| := \sum_{0 \leq i \leq m-1} \sigma_i \in \mathbb{Z}. \quad (10)$$

Then, we obtain the estimate

$$|p_{2m}| \leq P_{2m}, \quad P_{2m} := \sum_{\sigma \in \mathcal{I}_{2m+2}} c_\sigma |\psi|^{2\sigma_0} |\psi_x|^{2\sigma_1} \cdots |\psi_{(m-1)}|^{2\sigma_{m-1}}, \quad (11)$$

with positive constants c_σ . The majorant P_{2m} allows us to obtain detailed estimates of p_{2m} in Section 5. For now, we content ourselves with the following preliminary result.

Proof of Theorem 1.2(i) for $m \geq 2$. If $\varphi \in H_r^m$ and $\sigma \in \mathcal{I}_{2m+2}$, then by (10)

$$2m + 2 \geq 2m\sigma_{m-1}.$$

Suppose $\sigma_{m-1} > 1$, then necessarily $m = 2$ and $\sigma_{m-1} = \frac{3}{2}$ while all other indices σ_i vanish. In this case, however, $|\sigma| = \sigma_{m-1}$ is not an integer, giving a contradiction. Therefore, $\sigma_{m-1} \leq 1$ for all $\sigma \in \mathcal{I}_{2m+2}$, that is, P_{2m} is at most quadratic in $|\psi_{(m-1)}|$. With

$\|\psi_{(m-1)}\|_{L^1} \leq \|\psi_{(m-1)}\|_0$, it thus follows

$$\int_{\mathbb{T}} |\psi_{(m-1)}^{2\sigma_{m-1}}| \leq \frac{1}{2} \|\varphi_{(m-1)}\|_0^{2\sigma_{m-1}},$$

and by estimating the remaining factors using the L^∞ -estimate

$$|\psi_{(j)}^{2\sigma_j}| \leq \|\psi_{(j)}\|_\infty^{2\sigma_j} \leq \|\varphi\|_{j+1}^{2\sigma_j},$$

we arrive at

$$\int_{\mathbb{T}} |p_{2m}| \, dx \leq \int_{\mathbb{T}} P_{2m} \, dx \leq \sum_{\mathcal{I}_{2m+2}} c_\sigma \|\varphi\|_{m-1}^{2\sigma_0 + \dots + 2\sigma_{m-1}}.$$

Since $2\sigma_0 + \dots + 2\sigma_{m-1} \leq 2m + 2$, we obtain for the Hamiltonian

$$|\mathcal{H}_{2m+1}(\varphi)| \leq \frac{1}{2} \|\varphi\|_m^2 + b_m^2 (1 + \|\varphi\|_{m-1}^{2m}) \|\varphi\|_{m-1}^2,$$

with some positive constant b_m . Together with Lemma 4.2, this gives

$$\|I(\varphi)\|_{\ell_{2m}^1} \leq 2^{m-1} \|\varphi\|_m^2 + c_m^2 (1 + \|\varphi\|_{m-1}^{4m}) \|\varphi\|_{m-1}^2, \quad m \geq 2. \quad \blacksquare$$

5 Estimating the Sobolev Norms

We now turn to the converse problem of estimating the Sobolev norms of the potential in terms of weighted norms of its actions on level one. Our starting point is the identity

$$\frac{1}{2} \|\varphi_{(m)}\|_0^2 = 4^m \sum_{n \in \mathbb{Z}} J_{n, 2m+1} - \int_{\mathbb{T}} p_{2m} \, dx, \quad m \geq 1, \quad (12)$$

which is obtained by combining Corollary A.4 and the trace formula (3). The key step is estimating the actions $J_{n, 2m+1}$ in terms of I_n . Subsequently, we invoke an inductive argument using that p_{2m} contains only derivatives up to order $m - 1$. The main difficulty is to estimate the actions $J_{n, 2m+1}$ below the threshold of $\langle n \rangle$ provided by Proposition 4.1. Even though there are only finitely many of them, they cannot be controlled by the L^2 -norm $\|\varphi\|_0$ as one can translate the spectrum of φ without changing $\|\varphi\|_0$. Instead, we use the H^1 -norm $\|\varphi\|_1$, and provide estimates of $\|\varphi\|_1$ in terms of I_n by separate arguments which were inspired by work of Korotyaev [19] and Korotyaev and Kargaev [21].

Lemma 5.1. Uniformly for all $\varphi \in H_r^1$,

$$\mathcal{H}_3(\varphi) - 2\mathcal{H}_1^2(\varphi) \leq \sum_{n \in \mathbb{Z}} (2n\pi)^2 I_n(\varphi).$$

In particular, $\mathcal{H}_3(\varphi) \leq \|I(\varphi)\|_{\ell^1_2} + 2\|I(\varphi)\|_{\ell^1}^2$ and

$$\frac{1}{3}\|\varphi\|_1^2 \leq \|I(\varphi)\|_{\ell^1_2} + \|I(\varphi)\|_{\ell^1}^2. \quad \square$$

Proof. As the Hamiltonians and the actions are continuous on H_r^1 , it suffices to consider the case of a finite gap potential. Let C_r denote a sufficiently large circle enclosing all open gaps of φ , then the primitive F of ω defined in Section 2 is analytic outside C_r and its Laurent expansion is given by Lemma 2.4. Thus, by the Residue Theorem,

$$\frac{1}{\pi} \int_{C_r} F^3(\lambda) \, d\lambda = \frac{3}{8i\pi} \int_{C_r} \frac{1}{\lambda} (\mathcal{H}_3 - 2\mathcal{H}_1^2) \, d\lambda = \frac{3}{4} (\mathcal{H}_3 - 2\mathcal{H}_1^2).$$

The right-hand side is real as φ is of real type, and

$$\Re \int_{C_r} F^3(\lambda) \, d\lambda = \sum_{n \in \mathbb{Z}} \int_{\lambda_n^-}^{\lambda_n^+} \Re(F^3(\lambda - i0) - F^3(\lambda + i0)) \, d\lambda.$$

Furthermore, by Lemma 2.3 we have for $\lambda_n^- < \lambda < \lambda_n^+$,

$$\Re(F^3(\lambda - i0) - F^3(\lambda + i0)) = -2f_n^3(\lambda) + 6(n\pi)^2 f_n(\lambda),$$

and since f_n is nonnegative, we conclude with (7)

$$\frac{1}{\pi} \int_{C_r} F^3(\lambda) \, d\lambda \leq \sum_{n \in \mathbb{Z}} \frac{6}{\pi} \int_{\lambda_n^-}^{\lambda_n^+} (n\pi)^2 f_n(\lambda) \, d\lambda = \frac{3}{4} \sum_{n \in \mathbb{Z}} (2n\pi)^2 I_n.$$

This proves the first claim. Note that $\langle 2n\pi \rangle^2 \leq \frac{3}{2}(1 + (2n\pi)^2)$ for $n \in \mathbb{Z}$, so

$$\frac{1}{3}\|\varphi\|_1^2 \leq \frac{1}{2}(\|\varphi_x\|_0^2 + \|\varphi\|_0^2) = \mathcal{H}_3 - \int_{\mathbb{T}} |\psi|^4 \, dx + \mathcal{H}_1.$$

Since $\int_{\mathbb{T}} |\psi|^4 \, dx \geq \mathcal{H}_1^2(\varphi)$, the second claim follows with

$$\frac{1}{3}\|\varphi\|_1^2 \leq \mathcal{H}_3 - 2\mathcal{H}_1^2 + \mathcal{H}_1^2 + \mathcal{H}_1 \leq \|I(\varphi)\|_{\ell^1_2} + \|I(\varphi)\|_{\ell^1}^2. \quad \blacksquare$$

Proof of Theorem 1.2(ii). The case $m = 1$ is an immediate corollary of the lemma above. For the case $m \geq 2$, we find with (12) and Lemma 5.2

$$\frac{1}{2}\|\varphi_{(m)}\|_0^2 = \sum_{n \in \mathbb{Z}} 4^m J_{n,2m+1} - \int_{\mathbb{T}} p_{2m} \, dx \leq \|I\|_{\ell^1_{2m}} + c_m(1 + \|I\|_{\ell^1_{2m-2}})^{4m-3} \|I\|_{\ell^1_{2m-2}} + \int_{\mathbb{T}} |p_{2m}| \, dx.$$

It remains to consider the p_{2m} -term, for which we have the estimate (11)

$$\int_{\mathbb{T}} |p_{2m}| \, dx \leq \int_{\mathbb{T}} P_{2m} \, dx = \sum_{\mathcal{I}_{2m+2}} \int_{\mathbb{T}} c_{\sigma} |\psi|^{2\sigma_0} \cdots |\psi_{(m-1)}|^{2\sigma_{m-1}} \, dx.$$

The right-hand side only contains derivatives of ψ up to order $m - 1$ and the exponents σ_i satisfy the constraints (10). We use this to prove by induction for any $m \geq 2$

$$\frac{1}{2} \|\varphi_{(m)}\|_0^2 \leq \|I\|_{\ell_{2m}^1} + \mathcal{d}_m^2 (1 + \|I\|_{\ell_{2m-2}^1})^{\alpha_m} \|I\|_{\ell_{2m-2}^1},$$

with absolute constants \mathcal{d}_m and exponents $\alpha_m := 4m - 3$. By setting $\alpha_0 = 0$ and $\alpha_1 = 1$, this already holds for $m = 0, 1$.

Now consider the inductive step $m - 1 \mapsto m$. Recall from the proof of Theorem 1.2(i) that P_{2m} is at most quadratic in $|\psi_{(m-1)}|$ for $m \geq 2$. Thus, we have the following expansion:

$$\begin{aligned} P_{2m} &= P_{2m;2}(|\psi|, \dots, |\psi_{(m-2)}|) |\psi_{(m-1)}|^2 + P_{2m;1}(|\psi|, \dots, |\psi_{(m-2)}|) |\psi_{(m-1)}| \\ &\quad + P_{2m;0}(|\psi|, \dots, |\psi_{(m-2)}|), \end{aligned}$$

where for $0 \leq k \leq 2$ the entity $P_{2m;k} |\psi_{(m-1)}|^k$ incorporates those $\sigma \in \mathcal{I}_{2m+2}$ with $2\sigma_{m-1} = k$. In particular, $P_{2m;2} = c |\psi|^2$ with some $c \geq 0$. The three terms in the expansion of P_{2m} will be estimated separately.

Let us begin with $P_{2m;2} |\psi_{(m-1)}|^2 = c |\psi|^2 |\psi_{(m-1)}|^2$. To avoid a flood of constants, we write $a \leq b$ if $a \leq c \cdot b$ with an absolute constant c which is independent of ψ but may depend on m . Using the L^∞ -estimate and the induction hypothesis for $m - 1$, we get

$$\int_{\mathbb{T}} P_{2m;2} |\psi_{(m-1)}|^2 \, dx \leq \|\psi\|_\infty^2 \|\varphi_{(m-1)}\|_0^2 \leq \|\varphi\|_1^2 \|\varphi_{(m-1)}\|_0^2 \leq (1 + \|I\|_{\ell_{2m-2}^1})^{\alpha_{m-1}+2} \|I\|_{\ell_{2m-2}^1},$$

where $\alpha_{m-1} + 2 = 4(m - 1) - 1 \leq 4m - 3 = \alpha_m$.

Next, consider the term $P_{2m;1} |\psi_{(m-1)}|$ which involves all multi-indices $\sigma = (\sigma_i)_{0 \leq i \leq m-1}$ with $2\sigma_{m-1} = 1$. Recall that $|\sigma| = \sum_{0 \leq i \leq m-1} \sigma_i$ is an integer, hence there has to be an index $\sigma_j \neq 0$ with $0 \leq j \leq m - 2$. By Cauchy-Schwarz and the L^∞ -estimate

$$\begin{aligned} \int_{\mathbb{T}} |\psi_{(j)}|^{2\sigma_j} |\psi_{(m-1)}| \, dx &\leq \|\varphi\|_{j+1}^{2\sigma_j-1} \|\varphi_{(j)}\|_0 \|\varphi_{(m-1)}\|_0 \\ &\leq (1 + \|I\|_{\ell_{2m-2}^1})^{\alpha_{j+1}(\sigma_j-1/2) + \alpha_j/2 + \alpha_{m-1}/2} \|I\|_{\ell_{2m-2}^1}^{\sigma_j+1/2} \\ &\leq (1 + \|I\|_{\ell_{2m-2}^1})^{\alpha_{j+1}\sigma_j + \alpha_m\sigma_{m-1} - \delta_j - \delta_{m-1}} \|I\|_{\ell_{2m-2}^1}^{\sigma_j + \sigma_{m-1}}, \end{aligned}$$

where we have set $\delta_k := (\alpha_{k+1} - \alpha_k)/2$, so $\delta_0 = \frac{1}{2}$ and $\delta_k = 2$ for all $k \geq 1$. The remaining factors are estimated by

$$|\psi_{(k)}|^{2\sigma_k} \leq \|\varphi\|_{k+1}^{2\sigma_k} \leq (1 + \|I\|_{\ell_{2m-2}^1})^{\alpha_{k+1}\sigma_k} \|I\|_{\ell_{2m-2}^1}^{\sigma_k}.$$

Both estimates together yield

$$\begin{aligned} \int_{\mathbb{T}} |\psi|^{2\sigma_0} \dots |\psi_{(m-1)}|^{2\sigma_{m-1}} \, dx &\leq (1 + \|I\|_{\ell_{2m-2}^1})^{\alpha_1\sigma_0 + \dots + \alpha_m\sigma_{m-1} - \delta_j - \delta_{m-1}} \|I\|_{\ell_{2m-2}^1}^{\sigma_0 + \dots + \sigma_{m-1}} \\ &\leq (1 + \|I\|_{\ell_{2m-2}^1})^{(\alpha_1+1)\sigma_0 + \dots + (\alpha_m+1)\sigma_{m-1} - \delta_j - \delta_{m-1} - 1} \|I\|_{\ell_{2m-2}^1}. \end{aligned}$$

As $\delta_j + \delta_{m-1} \geq \frac{5}{2}$ and $\sum_{0 \leq i \leq m-1} (1+i)\sigma_i = m+1$, we find for the exponent

$$\sum_{0 \leq i \leq m-1} (\alpha_{i+1} + 1)\sigma_i - \delta_j - \delta_{m-1} - 1 \leq 4m - 2|\sigma| + 1/2.$$

If we had $|\sigma| = 1$, then

$$m+1 = \sum_{0 \leq i \leq m-1} (1+i)\sigma_i \leq \frac{m}{2} + \frac{m-1}{2},$$

which is impossible, thus $|\sigma| \geq 2$. Consequently, the exponent is not greater than $\alpha_m = 4m - 3$ and hence

$$\int_{\mathbb{T}} P_{2m;1} |\psi_{(m-1)}| \, dx \leq (1 + \|I\|_{\ell_{2m-2}^1})^{\alpha_m} \|I\|_{\ell_{2m-2}^1}.$$

It remains to estimate the term $P_{2m;0}$ involving all $\sigma = (\sigma_i)_{0 \leq i \leq m-1}$ with $\sigma_{m-1} = 0$. We first consider those multi-indices σ with $\sigma_i, \sigma_j \neq 0$ for some $i < j \leq m-2$. In that case, by Cauchy–Schwarz,

$$\begin{aligned} \int_{\mathbb{T}} |\psi_{(i)}|^{2\sigma_i} |\psi_{(j)}|^{2\sigma_j} \, dx &\leq \|\varphi\|_{i+1}^{2\sigma_i-1} \|\varphi\|_{j+1}^{2\sigma_j-1} \|\varphi_{(i)}\|_0 \|\varphi_{(j)}\|_0 \\ &\leq (1 + \|I\|_{\ell_{2m-2}^1})^{\alpha_{i+1}\sigma_i + \alpha_{j+1}\sigma_j - \delta_i - \delta_j} \|I\|_{\ell_{2m-2}^1}^{\sigma_i + \sigma_j}, \end{aligned}$$

while the remaining factors are estimated as usual,

$$|\psi_{(k)}|^{2\sigma_k} \leq \|\varphi\|_{k+1}^{2\sigma_k} \leq (1 + \|I\|_{\ell_{2m-2}^1})^{\alpha_{k+1}\sigma_k} \|I\|_{\ell_{2m-2}^1}^{\sigma_k}.$$

Together, both estimates give

$$\begin{aligned} \int_{\mathbb{T}} |\psi|^{2\sigma_0} \dots |\psi_{(m-2)}|^{2\sigma_{m-2}} \, dx &\leq (1 + \|I\|_{\ell_{2m-2}^1})^{\alpha_1\sigma_0 + \alpha_2\sigma_1 + \dots + \alpha_{m-1}\sigma_{m-2} - \delta_i - \delta_j} \|I\|_{\ell_{2m-2}^1}^{\sigma_0 + \dots + \sigma_{m-2}} \\ &\leq (1 + \|I\|_{\ell_{2m-2}^1})^{(\alpha_1+1)\sigma_0 + \dots + (\alpha_{m-1}+1)\sigma_{m-2} - \delta_i - \delta_j - 1} \|I\|_{\ell_{2m-2}^1}. \end{aligned}$$

As before $\delta_i + \delta_j \geq 5/2$ and $|\sigma| \geq 2$, hence the exponent is not greater than $\alpha_m = 4m - 3$.

Finally, consider the case of a monomial $c\psi_{(i)}^{2\sigma_i}$, that is, $(1+i)\sigma_i = m+1$ for some $0 \leq i \leq m-2$. If $i=0$, then $\sigma_0 = m+1$ and

$$\int_{\mathbb{T}} |\psi|^{2m+2} dx \leq \|\varphi\|_1^{2m} \|\varphi\|_0^2 \leq (1 + \|I\|_{\ell_2^1})^{2m} \|I\|_{\ell^1},$$

with $2m \leq \alpha_m$. Conversely, if $i \geq 1$, as necessarily $2\sigma_i \geq 2$, we find

$$\begin{aligned} \int_{\mathbb{T}} |\psi_{(i)}|^{2\sigma_i} dx &\leq \|\varphi\|_{i+1}^{2\sigma_i-2} \|\varphi_{(i)}\|_0^2 \leq (1 + \|I\|_{\ell_{2m-2}^1})^{\alpha_{i+1}(\sigma_i-1)+\alpha_i} \|I\|_{\ell_{2m-2}^1}^{\sigma_i} \\ &\leq (1 + \|I\|_{\ell_{2m-2}^1})^{(\alpha_{i+1}+1)\sigma_i-5} \|I\|_{\ell_{2m-2}^1}^{\sigma_i}, \end{aligned}$$

and further

$$(\alpha_{i+1} + 1)\sigma_i - 5 = 4(m+1) - 5 - 2\sigma_i \leq \alpha_m.$$

Altogether we thus obtain

$$\int_{\mathbb{T}} P_{2m} dx \leq (1 + \|I\|_{\ell_{2m-2}^1})^{\alpha_m} \|I\|_{\ell_{2m-2}^1},$$

which completes the induction step and proves Theorem 1.2(ii). ■

Lemma 5.2. For each $m \geq 1$,

$$\sum_{n \in \mathbb{Z}} 4^m J_{n,2m+1} \leq \|I(\varphi)\|_{\ell_{2m}^1} + (64\pi)^{2m-2} (1 + \|I(\varphi)\|_{\ell_2^1})^{4m-3} \|I(\varphi)\|_{\ell_2^1},$$

uniformly for all $\varphi \in H_r^m$. □

Proof. Let $N+1 \geq 8\|\varphi\|_1^2 > N$, then by Proposition 4.1

$$\sum_{|n| > N} 4^m J_{n,2m+1} \leq \sum_{|n| > N} \langle 2n\pi \rangle^{2m} I_n,$$

while for the remaining actions $J_{n,2m+1} = \tilde{\zeta}_{n,m}^{2m-2} J_{n,3}$ and hence

$$\begin{aligned} \sum_{|n| \leq N} 4^m J_{n,2m+1} &\leq (16\pi)^{2m-2} (1 + \|\varphi\|_1^2)^{2m-2} \sum_{|n| \leq N} 4J_{n,3} \\ &\leq (64\pi)^{2m-2} (1 + \|I(\varphi)\|_{\ell_2^1})^{4m-4} \sum_{|n| \leq N} 4J_{n,3}. \end{aligned}$$

By the trace formula (3) and Lemma 5.1, we finally obtain

$$\sum_{n \in \mathbb{Z}} 4J_{n,3} = \mathcal{H}_3 \leq \|I(\varphi)\|_{\ell_2^1} + 2\|I(\varphi)\|_{\ell^1}^2. \quad \blacksquare$$

6 Actions, Weighted Sobolev Spaces, and Gap Lengths

The case of estimating the action variables of φ in standard Sobolev spaces H_r^m with integers $m \geq 1$ is somewhat special due to the presence of the trace formula (3). When arbitrary weighted Sobolev spaces H_r^w are considered, there is no identity known to exist relating $\|\varphi\|_w$ to Hamiltonians of the NLS hierarchy. Albeit, even in the case of weighted Sobolev spaces, the regularity properties of φ are well known to be closely related to the decay properties of the gap lengths $\gamma_n(\varphi)$ —see, for example, [7, 16] and Section 3. Moreover,

$$\frac{4I_n}{\gamma_n^2} = 1 + \ell_n^2 \quad (13)$$

is known to hold locally uniformly on L_r^2 and hence uniformly on bounded subsets of H_r^1 —see [10]. In this section, we prove a quantitative version of (13) which is quadratic in $\|\varphi\|_1$ on all of H_r^1 . From this and the estimates of the gap lengths given in Section 3, we then obtain Theorem 1.3.

To set the stage, let $\varphi \in W$ and recall from (6),

$$I_n = \frac{1}{\pi} \int_{\Gamma_n} \lambda \omega = -\frac{1}{\pi} \int_{\Gamma_n} (\lambda_n^\bullet - \lambda) \omega.$$

Here the latter identity follows as ω is closed around the gap. In the case $I_n \neq 0$, or equivalently $\gamma_n \neq 0$, we shrink the contour Γ_n to the straight line $[\lambda_n^-, \lambda_n^+]$ and insert the product representation (4) of ω , to obtain

$$I_n = -\frac{2}{\pi} \int_{\lambda_n^-}^{\lambda_n^+} \frac{(\lambda_n^\bullet - \lambda)^2}{i\varsigma_n(\lambda)} \chi_n(\lambda) \, d\lambda, \quad \chi_n(\lambda) = \prod_{m \neq n} \frac{\lambda_m^\bullet - \lambda}{\varsigma_m(\lambda)}.$$

It follows with the substitution $\lambda = \tau_n + t\gamma_n/2$ that

$$\frac{4I_n}{\gamma_n^2} = \frac{2}{\pi} \int_{-1}^1 \frac{(t - t_n)^2}{\sqrt{1 - t^2}} \chi_n(\tau_n + t\gamma_n/2) \, dt, \quad t_n = 2(\lambda_n^\bullet - \tau_n)/\gamma_n.$$

By Lemma A.1, there exists an open connected neighborhood $\hat{W} \subset W$ of L_r^2 such that $|\lambda_m^\bullet - \tau_m| \leq |\gamma_m|$ for all $m \in \mathbb{Z}$, hence $|t_n| \leq 2$, and thus

$$\left| \frac{4I_n}{\gamma_n^2} \right| \leq 9 |\chi_n|_{[\lambda_n^-, \lambda_n^+]}. \quad (14)$$

To get a quantitative version of (13), we thus need a uniform estimate of χ_n .

Lemma 6.1. On $H_c^1 \cap \hat{W}$ for any $|n| \geq \langle N \rangle \geq 8\|\varphi\|_1^2$,

$$|\chi_n|_{[\lambda_n^-, \lambda_n^+]} \leq e^2 \left(\frac{|n| + N + 3/5}{|n| - N - 3/5} \right) \leq 2^9 (1 + \|\varphi\|_1^2). \quad \square$$

Proof. Suppose $\varphi \in H_c^1 \cap \hat{W}$ and choose $\langle N \rangle \geq 8\|\varphi\|_1^2 > N$. For $|n| \geq N$, we split the product χ_n into two parts,

$$\chi_n(\lambda) = \left(\prod_{|m| \leq N} \frac{\lambda_m^\bullet - \lambda}{\varsigma_m(\lambda)} \right) \left(\prod_{\substack{n \neq m \\ |m| > N}} \frac{\lambda_m^\bullet - \lambda}{\varsigma_m(\lambda)} \right).$$

If $|k| > N$, then by Theorem 3.1

$$|\lambda_k^\pm - k\pi| \leq \frac{\|\varphi\|_1^2}{\langle k \rangle} + \frac{\sqrt{2}\|\varphi\|_1}{\langle 2k \rangle} \leq \pi/8,$$

where we used that $\langle 2k \rangle \geq 3\langle k \rangle/2$. Thus, for $|m|, |n| > N$ and $\lambda \in [\lambda_n^-, \lambda_n^+]$,

$$|\tau_m - \lambda| \geq 2|n - m|.$$

Further, $|\gamma_m| \leq 4\|\varphi\|_1/\langle 2m \rangle$ by the individual gap estimate (9), hence

$$\left| \frac{\gamma_m/2}{\tau_m - \lambda} \right| \leq \frac{\|\varphi\|_1}{\langle 2m \rangle |n - m|} \leq 1/4.$$

It follows that $|\varsigma_m(\lambda)| \geq |\tau_m - \lambda| - |\gamma_m|/2$. Moreover, $|\lambda_m^\bullet - \tau_m| \leq |\gamma_m|$ as $\varphi \in \hat{W}$, thus with $(1 + 2r)/(1 - r) \leq 1 + 4r$ for $0 \leq r \leq \frac{1}{4}$, we conclude

$$\left| \frac{\lambda_m^\bullet - \lambda}{\varsigma_m(\lambda)} \right| \leq \frac{|\tau_m - \lambda| + |\gamma_m|}{|\tau_m - \lambda| - |\gamma_m|/2} \leq 1 + \frac{|\gamma_m|}{|n - m|}.$$

It follows with Cauchy–Schwarz that

$$\begin{aligned} \sum_{\substack{m \neq n \\ |m| > N}} \frac{|\gamma_m|}{|n - m|} &\leq \left(\sum_{\substack{m \neq n \\ |m| > N}} \frac{1}{\langle 2m \rangle^2 |n - m|^2} \right)^{1/2} \left(\sum_{|m| > N} \langle 2m \rangle^2 |\gamma_m|^2 \right)^{1/2} \\ &\leq \frac{2}{\langle 2N + 2 \rangle} \left(\sum_{|m| > N} \langle 2m \rangle^2 |\gamma_m|^2 \right)^{1/2}, \end{aligned}$$

and by Proposition 3.8

$$\sum_{|m| > N} \langle 2m \rangle^2 |\gamma_m|^2 \leq 6\|\mathcal{R}_N \varphi\|_1^2 + 144\|\varphi\|_1^4 \leq 4\langle N \rangle^2.$$

Therefore, by the standard estimates for infinite products,

$$\prod_{\substack{|m|>N \\ m \neq n}} \left| \frac{\lambda_m^\bullet - \lambda}{\zeta_m(\lambda)} \right| \leq \exp(4\langle N \rangle / \langle 2N + 2 \rangle) \leq e^2.$$

To estimate the remaining part of the product, we note that λ_m^\bullet and λ_m^\pm are contained in the isolating neighborhood U_m , which is a complex disk centered on the real line. Thus, if $\lambda \in [\lambda_n^-, \lambda_n^+]$ and $n > N$, then

$$|\lambda_{m-1}^\pm - \lambda| > |\lambda_m^\bullet - \lambda|, \quad m \leq N,$$

and consequently

$$\prod_{|m| \leq N} \left| \frac{\lambda_m^\bullet - \lambda}{\zeta_m(\lambda)} \right| = \left| \frac{\lambda_{-N}^\bullet - \lambda}{\zeta_{-N}(\lambda)} \right| \prod_{-N < m \leq N} \left| \frac{\lambda_m^\bullet - \lambda}{\zeta_{m-1}(\lambda)} \right| \leq \frac{|\tau_{-N} - \lambda| + |\gamma_{-N}|}{||\tau_N - \lambda| - |\gamma_N|/2|}.$$

By Theorem 3.1, we have $|\gamma_{\pm N}| \leq \pi/5$, as well as

$$|\pm N - n|\pi - 2\pi/5 \leq |\tau_{\pm N} - \lambda| \leq |\pm N - n|\pi + 2\pi/5.$$

It follows that

$$\frac{|\tau_{-N} - \lambda| + |\gamma_{-N}|}{||\tau_N - \lambda| - |\gamma_N|/2|} \leq \frac{|n| + N + 3/5}{|n| - N - 3/5} \leq 5\langle N \rangle.$$

Similarly, for $\lambda \in [\lambda_n^-, \lambda_n^+]$ with $n < -N$. ■

Proposition 6.2. Suppose $\varphi \in H_c^1 \cap \hat{W}$, then for any $|n| \geq 8\|\varphi\|_1^2$,

$$|I_n| \leq 2^{11}(1 + \|\varphi\|_1^2)|\gamma_n|^2. \quad \square$$

Proof. If $\gamma_n = 0$, then $I_n = 0$ and the estimate clearly holds. If $\gamma_n \neq 0$, then by (14) and the preceding lemma,

$$|4I_n/\gamma_n^2| \leq 9|\chi_n|_{[\lambda_n^-, \lambda_n^+]} \leq 2^{13}(1 + \|\varphi\|_1^2). \quad \blacksquare$$

Proof of Theorem 1.3. Suppose $\varphi \in H_c^w \cap \hat{W}$ and choose $N + 1 \geq 8\|\varphi\|_w^2 > N$. Then by the preceding proposition

$$\sum_{|n| > N} w_{2n}^2 |I_n| \leq 2^{11} \|\varphi\|_1^2 \sum_{|n| > N} w_{2n}^2 |\gamma_n|^2,$$

and the gap lengths may be estimated by Proposition 3.8

$$\sum_{|n|>N} w_{2n}^2 |\gamma_n|^2 \leq 6 \|R_N \varphi\|_w^2 + 144 \|\varphi\|_w^4 \leq 144(1 + \|\varphi\|_w^2) \|\varphi\|_w^2.$$

In particular, the mapping

$$H_c^w \cap \hat{W} \rightarrow [0, \infty), \quad \varphi \mapsto \sum_{n \in \mathbb{Z}} w_{2n}^2 |I_n|,$$

is continuous. Suppose that φ is of real type, then the remaining actions for $|n| \leq N$ may be estimated by

$$\sum_{|n| \leq N} w_{2n}^2 |I_n| \leq w_{2N}^2 \sum_{n \in \mathbb{Z}} I_n \leq w [16 \|\varphi\|_w^2]^2 \|\varphi\|_0^2.$$

Since $w \in \mathcal{M}^1$ is growing with at least linear speed, we thus find on H_r^w

$$\sum_{n \in \mathbb{Z}} w_{2n}^2 |I_n| \leq 2^{20} w [16 \|\varphi\|_w^2]^2 \|\varphi\|_w^2.$$

For any nonzero potential φ , this estimate extends by continuity to a complex neighborhood of φ within H_c^w with just the absolute constant doubled. On the other hand, sufficiently close to the zero potential we have $\|\varphi\|_w^2 < \frac{1}{8}$. In this case, we may choose $N = 0$ such that

$$\sum_{n \neq 0} w_{2n}^2 |I_n| \leq 2^{20} (1 + \|\varphi\|_w)^4 \|\varphi\|_1^2.$$

Consequently, on some sufficiently small open neighborhood of H_r^w in H_c^w ,

$$\|I(\varphi)\|_{\ell_w^1} \leq c_w^2 w [16 \|\varphi\|_w^2]^2 \|\varphi\|_w^2,$$

with a real constant c_w depending only on w_0 . ■

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Appendix 1. Analyticity of the Primitive F

In this appendix, we prove the analyticity of the primitive

$$F(\mu, \psi) = \frac{1}{2} \left(\int_{\lambda_0^-(\psi)}^{\mu} \omega(\lambda, \psi) + \int_{\lambda_0^+(\psi)}^{\mu} \omega(\lambda, \psi) \right),$$

introduced in Section 2 with $\omega = \frac{\Delta^\bullet}{\sqrt[3]{\Delta^2 - 4}} d\lambda$. The proof relies on the following observation.

Lemma A.1. (i) Suppose $\varphi \in W$ and $\gamma_n(\varphi) = 0$ for some $n \in \mathbb{Z}$, then there exists an open neighborhood V_n of φ such that

$$|\lambda_n^\bullet - \tau_n| \leq |\gamma_n|/2, \quad \psi \in V_n.$$

(ii) For each $\varphi \in L_r^2$, there exists an open neighborhood $V \subset L_c^2$ such that

$$|\lambda_m^\bullet - \tau_m| \leq |\gamma_m|, \quad m \in \mathbb{Z}, \quad \psi \in V. \quad \square$$

Proof. (i) For any potential in $V_n = V_\varphi$, we have

$$0 = \frac{1}{4}(\Delta^2 - 4)^\bullet \Big|_{\lambda_n^\bullet} = (2(\tau_n - \lambda)\Delta_n - ((\tau_n - \lambda)^2 - \gamma_n^2/4)\Delta_n^\bullet) \Big|_{\lambda_n^\bullet},$$

where the function Δ_n is analytic on $\mathbb{C} \times W$ and given by

$$\Delta_n(\lambda) := \frac{1}{\pi_n^2} \prod_{m \neq n} \frac{(\lambda_m^+ - \lambda)(\lambda_m^- - \lambda)}{\pi_m^2}.$$

The zeroes of Δ_n are precisely the eigenvalues λ_m^\pm for $m \neq n$, thus Δ_n does not vanish on $U_n \times V_n$ and we have $|\Delta_n| \geq s > 0$ and $|\Delta_n^\bullet| \leq r$ uniformly on $U_n \times V_n$. Since $\gamma_n(\varphi) = 0$, we may shrink V_n , if necessary, to the effect that

$$|\gamma_n(\psi)|r \leq s, \quad \psi \in V_n.$$

To simplify notation, put $f = 2(\tau_n - \lambda)\Delta_n$ and $g = (\Delta^2 - 4)^\bullet/4$. By Lemma 2.2, $\gamma_n(\psi) = 0$ implies $\lambda_n^\bullet(\psi) = \tau_n(\psi)$, hence we may assume $\gamma_n(\psi) \neq 0$. In this case, we find on the boundary of the disk $D_\psi = \{|\lambda - \tau_n(\psi)| \leq |\gamma_n(\psi)|/2\}$,

$$|f - g|_{\partial D_\psi} \leq |\gamma_n(\psi)|^2 r/2 < s|\gamma_n(\psi)| \leq |f|_{\partial D_\psi}.$$

Thus, by Rouch's Theorem, f and g have the same number of roots in D_ψ , and since λ_n^\bullet is the only root of g contained in U_n , the claim (i) follows.

(ii) Recall from [10] that for each $\varphi \in L_c^2$ there exists a neighborhood V and an $M \geq 0$ such that

$$\sum_{|m| \geq M} |\lambda_m^\bullet - \tau_m|^2 / |\gamma_m|^4 \leq 1, \quad \psi \in V.$$

After possibly increasing M , we have $|\lambda_m^\bullet - \tau_m| \leq |\gamma_m|$ on V for all $|m| \geq M$. Suppose that, in addition, $\varphi \in L_r^2$. If $\gamma_m(\varphi) = 0$ for some $|m| < M$, then we may shrink V according to (i) to obtain the desired inequality. On the other hand, if $\gamma_m(\varphi) \neq 0$, then

$$|\lambda_m^\bullet(\varphi) - \tau_m(\varphi)| \leq |\gamma_m(\varphi)|/2,$$

and since λ_m^\bullet , τ_m , and $|\gamma_m|$ are continuous on W , the desired inequality follows after possibly shrinking V . ■

Proposition A.2. The mapping F is analytic on $(\mathbb{C} \setminus \bigcup_{n \in \mathbb{Z}} U_n) \times V_\varphi$. □

Proof. Take any $v \in \partial U_0$, and write the integral as

$$F(\mu) = \frac{1}{2} \left(\int_{\lambda_0^-}^v \omega + \int_{\lambda_0^+}^v \omega \right) + \int_v^\mu \omega,$$

where the admissible path of integration of $\int_{\lambda_0^\pm}^v \omega$ runs in $U_0 \setminus [\lambda_0^-, \lambda_0^+]$ except for its end points, and the one of $\int_v^\mu \omega$ runs in $\mathbb{C} \setminus \bigcup_{n \in \mathbb{Z}} U_n$. Then, by Lemma 2.2, $\int_v^\mu \omega$ is analytic on $(\mathbb{C} \setminus \bigcup_{n \in \mathbb{Z}} U_n) \times V_\varphi$. We now want to prove that $F_v := F(v)$ is analytic on V_φ for an arbitrary $v \in \partial U_0$.

Denote by $Z_0 := \{\psi \in V_\varphi : \gamma_0^2(\psi) = 0\}$ the analytic subvariety of V_φ of potentials with a collapsed zeroth gap. We first prove analyticity on the open set $V_\varphi \setminus Z_0$, secondly continuity on all of V_φ , and finally weak analyticity on Z_0 . The analyticity on V_φ then follows by the general theory of analytic functions—see [14, Appendix A]. Note that by the argument above, we may move v on ∂U_0 after each of the three steps.

Analyticity on $V_\varphi \setminus Z_0$. Fix any ψ^* in $V_\varphi \setminus Z_0$. Since the eigenvalues $\lambda_0^\pm(\psi^*)$ are simple, there exist two functions ρ_1 and ρ_2 , which are analytic on some neighborhood $V \subset V_\varphi \setminus Z_0$ of ψ^* and satisfy the set equality $\{\rho_1, \rho_2\} = \{\lambda_0^-, \lambda_0^+\}$. Provided the straight line $[\rho_1, v]$ is an admissible path, we can use the parameterization

$$\lambda_t = \rho_1 + t\mathbf{z}, \quad \mathbf{z} = v - \rho_1,$$

together with the product representation (4) of ω to get

$$F_\nu = \int_{\rho_1}^\nu \omega = -i \int_0^1 \frac{\lambda_0^\bullet - \lambda_t}{\varsigma_0(\lambda_t)} \chi_0(\lambda_t) z \, dt, \quad \chi_0(\lambda) = \prod_{m \neq 0} \frac{\lambda_m^\bullet - \lambda}{\varsigma_m(\lambda)}. \quad (\text{A.1})$$

Note that χ_0 is analytic on $U_0 \times V_\varphi$ —see [10, Section 12]. Let $\tilde{\gamma}_0 = \rho_2 - \rho_1$ and fix ν on ∂U_0 such that for some real $\sigma(\psi^*) > 0$,

$$\nu = \rho_1(\psi^*) - \sigma(\psi^*)\tilde{\gamma}_0(\psi^*)/2.$$

After possibly shrinking V , we find for the same fixed ν

$$\nu = \rho_1(\psi) - \sigma(\psi)\tilde{\gamma}_0(\psi)/2, \quad \psi \in V,$$

with $\sigma(\psi)$ possibly complex though $\Re \sigma \geq \varepsilon > 0$ on V . As $z = -\sigma\tilde{\gamma}_0/2$, we conclude

$$|\tau_0 - \lambda_t| = |1 + t\sigma| \cdot |\gamma_0/2| > |\gamma_0/2|, \quad 0 < t \leq 1,$$

uniformly on V . Therefore, the path $[\rho_1, \nu]$ is admissible for all ψ in V . In view of (1), the root $\varsigma_0(\lambda_t)$ is continuous in $t \geq 0$, analytic on V for $t > 0$, and satisfies the lower bound

$$|\varsigma_0(\lambda_t)| \geq |\gamma_0/2| \sqrt[4]{(1 + t\varepsilon)^2 - 1} \geq |\gamma_0/2| \sqrt[4]{2t\varepsilon}.$$

It follows that the integrand $(\lambda_0^\bullet - \lambda_t)\chi_0(\lambda_t)z/\varsigma_0(\lambda_t)$ of (A.1) is continuous on $(0, 1] \times V$, analytic on V for fixed $t > 0$, and has an integrable majorant. In consequence, F_ν is analytic on V , and hence on all of $V_\varphi \setminus Z_0$.

Continuity on V_φ . Clearly, F_ν is continuous on $V_\varphi \setminus Z_0$, and its restriction

$$F_\nu \Big|_{Z_0} = -i \int_{\tau_0}^\nu \chi_0(\lambda) \, d\lambda,$$

is continuous, too, as χ_0 and τ_0 are. To establish the continuity of F_ν on V_φ , it thus suffices to show that $F_\nu(\psi^k) \rightarrow F_\nu(\psi^*)$ for every sequence (ψ^k) in $V_\varphi \setminus Z_0$ converging to some ψ^* in Z_0 .

In view of Lemma A.1, we may without loss assume $|\lambda_0^\bullet - \tau_0| \leq |\gamma_0|/2$ for all ψ^k . Substituting $\lambda_t = \tau_0 - t\gamma_0/2$ such that $\varsigma_0(\lambda_t)^2 = -(1 - t^2)\gamma_0^2/4$ gives

$$\left| \int_{\lambda_0^-}^{\tau_0} \frac{\lambda_0^\bullet - \lambda}{i\varsigma_0(\lambda)} \chi_0(\lambda) \, d\lambda \right| \leq \int_0^1 \left| \frac{\lambda_0^\bullet - \tau_0 + t\gamma_0/2}{\sqrt[4]{1 - t^2}} \chi_0(\lambda_t) \right| dt = O(|\gamma_0|),$$

where the implicit constant can be chosen uniformly in k . Therefore,

$$F_\nu(\psi^k) = -i \int_{\tau_0}^\nu \frac{\lambda_0^\bullet - \lambda}{\varsigma_0(\lambda)} \chi_0(\lambda) \, d\lambda \Big|_{\psi^k} + o(1).$$

We can choose v^* on ∂U_0 such that the straight line $[\tau_0, v^*]$ is admissible for any ψ^k . With the parameterization

$$\lambda_t = \tau_0 + tz, \quad z = v^* - \tau_0,$$

we then obtain

$$-i \int_{\tau_0}^{v^*} \frac{\lambda_0^\bullet - \lambda}{\varsigma_0(\lambda)} \chi_0(\lambda) d\lambda = -i \int_0^1 \frac{\lambda_0^\bullet - \lambda_t}{\varsigma_0(\lambda_t)} \chi_0(\lambda_t) z dt.$$

Since $|z|$ is uniformly bounded below on V ,

$$\frac{|\gamma_0/2|}{|\tau_0 - \lambda_t|} = O(|\gamma_0/t|).$$

In view of (1) for any $\delta > 0$, there exists a neighborhood $V_\delta \subset V$ of ψ^* such that the root $\varsigma_0(\lambda_t)$ is continuous and does not vanish on $[\delta, 1] \times V_\delta$. Consequently, for all $\delta > 0$,

$$-i \int_\delta^1 \frac{\lambda_0^\bullet - \lambda_t}{\varsigma_0(\lambda_t)} \chi_0(\lambda_t) z \Big|_{\psi^k} \rightarrow -i \int_\delta^1 \chi_0(\lambda_t) z \Big|_{\psi^*}, \quad k \rightarrow \infty.$$

Let $\varepsilon = |\gamma_0|/2\delta|z|$, then after possibly shrinking V_δ we have $0 < \varepsilon < 1$ and

$$\begin{aligned} \int_0^\delta \left| \frac{\lambda_0^\bullet - \lambda_t}{\varsigma_0(\lambda_t)} \chi_0(\lambda_t) z \right| dt &\leq \int_0^\delta \frac{t|z| + |\gamma_0/2|}{\sqrt{|t^2|z|^2 - |\gamma_0/2|^2}} |\chi_0(\lambda_t) z| dt \\ &= O\left(\delta \int_0^1 \frac{\sqrt[4]{|t+\varepsilon|}}{\sqrt[4]{|t-\varepsilon|}} dt\right), \end{aligned}$$

where the implicit constant is uniform on V_δ . One checks that the latter integral is uniformly bounded for $0 \leq \varepsilon \leq 1$, hence

$$\sup_{\psi \in V_\delta} \int_0^\delta \left| \frac{\lambda_0^\bullet - \lambda_t}{\varsigma_0(\lambda_t)} \chi_0(\lambda_t) z \right| dt = O(\delta).$$

It follows that

$$-i \int_0^1 \frac{\lambda_0^\bullet - \lambda_t}{\varsigma_0(\lambda_t)} \chi_0(\lambda_t) z \Big|_{\psi^k} \rightarrow -i \int_0^1 \chi_0(\lambda_t) z \Big|_{\psi^*}, \quad k \rightarrow \infty,$$

so F_{v^*} is continuous at ψ^* and hence on all of V_φ .

Weak analyticity. The restriction of F_v to Z_0 is given by

$$F_v|_{Z_0} = -i \int_{\tau_0}^v \chi_0(\lambda) d\lambda,$$

and since χ_0 and τ_0 are both analytic, the weak analyticity follows directly. This completes the proof of the analyticity of F_v on V_φ . ■

Appendix 2. Properties of the NLS Hierarchy

In this appendix, we recall some well-known facts about the Hamiltonians of the NLS hierarchy—see [10]. For $\varphi = (\varphi_-, \varphi_+) \in H_c^{k-1}$, the k th NLS Hamiltonian is given by

$$\mathcal{H}_k(\varphi) = \int_{\mathbb{T}} \varphi_- u_k(\varphi_-, \varphi_+, \dots, \varphi_-^{(k-1)}, \varphi_+^{(k-1)}) \, dx, \quad k \geq 1,$$

where $u_1 = -\varphi_+$, and

$$u_{k+1} = u'_k + \varphi_- \sum_{l=1}^{k-1} u_{k-l} u_l, \quad k \geq 1.$$

Lemma A.3. If $\varphi \in H_c^{k-1}$, then

$$u_{k+1} = -\varphi_+^{(k)} + q_k(\varphi_-, \varphi_+, \dots, \varphi_-^{(k-2)}, \varphi_+^{(k-2)}),$$

where q_k is a homogeneous polynomial of degree $k+1$ when φ_- , φ_+ , and ∂_x each count as one degree. Moreover, each term of $\varphi_- q_k$ has at most degree $k-2$ with respect to ∂_x , and the degree with respect to φ_- equals the one with respect to φ_+ . \square

Proof. As is evident from their definition, the polynomials u_k are homogeneous of degree k , and only contain derivatives of φ_- and φ_+ up to order $k-1$. Furthermore, $u_k(\varphi_-, \lambda\varphi_+) = \lambda u_k(\lambda\varphi_-, \varphi_+)$ for all $\lambda \in \mathbb{C}$, which completes the proof. \blacksquare

In case of a smooth real-type potential, that is, $\varphi = (\psi, \bar{\psi})$, the odd Hamiltonians have the form

$$(-1)^{m+1} \mathcal{H}_{2m+1}(\varphi) = \int_{\mathbb{T}} (|\psi^{(m)}|^2 + \psi q_{2m}) \, dx, \quad m \geq 1,$$

where q_{2m} depends on ψ , $\bar{\psi}$, and their derivatives up to order $2m-2$. Suppose that ψq_{2m} contains a monomial $\psi_{(m+n)} \mathfrak{q}(\psi, \bar{\psi}, \dots, \psi_{(2m-2)}, \bar{\psi}_{(2m-2)})$ with $n \geq 0$. Since this monomial has at most degree $2m-2$ with respect to ∂_x , it follows that \mathfrak{q} contains at most $m-2-n$ derivatives. Thus, we can integrate by parts to the effect that each factor of the monomial contains at most $m-1$ derivatives.

Corollary A.4. For any $m \geq 1$, there exists a polynomial p_{2m} such that

$$(-1)^{m+1} \mathcal{H}_{2m+1}(\varphi) = \int_{\mathbb{T}} (|\psi_{(m)}|^2 + p_{2m}(\psi, \bar{\psi}, \dots, \psi_{(m-1)}, \bar{\psi}_{(m-1)})) \, dx,$$

for all $\varphi = (\psi, \bar{\psi})$ in H_r^m . The polynomial p_{2m} is homogenous of degree $2m + 2$ with respect to ψ , $\bar{\psi}$, and ∂_x , and the degree of each term of p_{2m} with respect to ψ equals the one with respect to $\bar{\psi}$. \square

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