

NEW INTEGRABILITY CONDITIONS OF DERIVATIONAL EQUATIONS OF A SUBMANIFOLD IN A GENERALIZED RIEMANNIAN SPACE

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Abstract

The present work is a continuation of [5] and [6]. In [5] we have obtained derivational equations of a submanifold X_M of a generalized Riemannian space GR_N . Since the basic tensor in GR_N is asymmetric and in this way the connection is also asymmetric, in a submanifold the connection is generally asymmetric too. By reason of this, we define 4 kinds of covariant derivative and obtain 4 kinds of derivational equations. In [6] we have obtained integrability conditions and Gauss-Codazzi equations using the 1st and the 2st kind of covariant derivative.

The present work deals in the cited matter, using the 3rd and the 4th kind of covariant derivative. One obtains three new integrability conditions for derivational equations of tangents and three such conditions for normals of the submanifold, as the corresponding Gauss-Codazzi equations too.

1 Introduction

1.1. A generalized Riemannian space GR_N is a differentiable manifold equipped with an asymmetric basic tensor $G_{ij}(x^1, \dots, x^N)$ (the components) where x^i are the local coordinates. The symmetric, respectively antisymmetric part of G_{ij} are H_{ij} and K_{ij} .

For the lowering and raising of indices in GR_N one uses H_{ij} , respectively H^{ij} , where

$$(1.1) \quad (H^{ij}) = (H_{ij})^{-1}, \quad (\det(H_{ij}) \neq 0).$$

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Cristoffel symbols at GR_N are

$$(1.2) \quad \Gamma_{i,jk} = \frac{1}{2}(G_{ji,k} - G_{jk,i} + G_{ik,j}), \quad \Gamma_{jk}^i = H^{ip}\Gamma_{p,jk},$$

where, for example, $G_{ji,k} = \partial G_{ji}/\partial x^k$. Based on the asymmetry of G_{ij} , it follows that the Cristoffel symbols are also asymmetric with respect to j, k in (1.2).

By equations

$$(1.3) \quad x^i = x^i(u^1, \dots, u^M) \equiv x^i(u^\alpha), \quad i = 1, \dots, N,$$

a submanifold X_M is defined in local coordinates. If $\text{rank}(B_\alpha^i) = M$ ($B_\alpha^i = \partial x^i/\partial u^\alpha$) and

$$(1.4) \quad g_{\alpha\beta} = B_\alpha^i B_\beta^j G_{ij},$$

X_M becomes $GR_M \subset GR_N$, with **induced basic tensor** (1.4), which is generally also asymmetric. Note that in the present work Latin indices i, j, \dots take values $1, \dots, N$ and refer to the GR_N , while the Greek ones take values $1, \dots, M$ and refer to the GR_M .

In the GR_M are valid the relations similar to (1.1) and (1.2). The symmetric part of $g_{\alpha\beta}$ is denoted with $h_{\alpha\beta}$, and antisymmetric one with $k_{\alpha\beta}$, where e.g.

$$(1.5) \quad h_{\alpha\beta} = B_\alpha^i B_\beta^j H_{ij}, \quad (h^{\alpha\beta}) = (h_{\alpha\beta})^{-1}.$$

Cristoffel symbols $\tilde{\Gamma}_{\alpha,\beta\gamma}$, $\tilde{\Gamma}_{\beta\gamma}^\alpha = h^{\alpha\pi}\tilde{\Gamma}_{\pi,\beta\gamma}$ are expressed by $g_{\alpha\beta}$ analogously to (1.2).

For the unit, mutually orthogonal vectors N_A^i , which are orthogonal to the GR_M too, we have [1]

$$(1.6) \quad H_{ij}N_A^i N_B^j = e_A \delta_B^A = h_{AB}, \quad e_A \in \{-1, 1\}, \quad H_{ij}N_A^i B_\alpha^j = 0,$$

where $A, B, \dots \in \{M+1, \dots, N\}$.

As it is known, the following relations between Cristoffel symbols of a generalized Riemannian space and its subspace are valid:

$$(1.7) \quad \tilde{\Gamma}_{\alpha,\beta\gamma} = \Gamma_{i,jk} B_\alpha^i B_\beta^j B_\gamma^k + H_{ij} B_\alpha^i B_{\beta,\gamma}^j,$$

$$(1.8) \quad \tilde{\Gamma}_{\beta\gamma}^\alpha = h^{\pi\alpha}\tilde{\Gamma}_{\pi,\beta\gamma} = h^{\pi\alpha}(\Gamma_{i,jk} B_\pi^i B_\beta^j B_\gamma^k + H_{ij} B_\pi^i B_{\beta,\gamma}^j),$$

i.e.

$$(1.8') \quad \tilde{\Gamma}_{\beta\gamma}^\alpha = h^{\pi\alpha} H_{pi} B_\pi^p (\Gamma_{jk}^i B_\beta^j B_\gamma^k + B_{\beta,\gamma}^i).$$

1.2. The set of normals of the submanifold $X_M \subset GR_N$ make a **normal bundle** for X_M , and we note it X_{N-M}^N . One can introduce a metric tensor on X_{N-M}^N

$$(1.9) \quad g_{AB} = G_{ij} N_A^i N_B^j,$$

which is asymmetric in a general case.

The symmetric part is

$$(1.10) \quad h_{AB} = H_{ij} N_A^i N_B^j \stackrel{(1.5)}{=} e_A \delta_B^A = h_{BA} = \begin{cases} e_A, & A=B, \\ 0, & \text{otherwise.} \end{cases}, \quad e_A \in \{-1, 1\}.$$

If

$$(h^{AB}) = (h_{AB})^{-1},$$

we have

$$h^{AB} = e_A \delta_B^A = h_{AB} = h^{BA}.$$

On X_{N-M}^N one can define in two manners connection coefficients

$$(1.11) \quad \bar{\Gamma}_{1 \frac{1}{2} \frac{3}{4}}^A = H_{ij} h^{AQ} N_Q^j (N_{B,\mu}^i + \Gamma_{pq}^i N_B^p B_\mu^q).$$

Being the coefficients $\Gamma, \tilde{\Gamma}, \bar{\Gamma}$ non-symmetric in general, for a tensor, defined at points of GR_M , is possible define four kinds of covariant derivative. For example

$$(1.12) \quad \begin{array}{c} \nabla_{\frac{1}{2} \frac{3}{4}} t_{j\beta B}^{i\alpha A} \equiv t_{j\beta B|\mu}^{i\alpha A} = t_{j\beta B,\mu}^{i\alpha A} + \Gamma_{\substack{pm \\ pm \\ mp}}^i t_{j\beta B}^{p\alpha A} B_\mu^m - \Gamma_{\substack{jm \\ mj \\ jm}}^p t_{p\beta B}^{i\alpha A} B_\mu^m \\ + \tilde{\Gamma}_{\substack{\mu\pi \\ \pi\mu \\ \mu\pi}}^\alpha t_{j\beta B}^{i\pi A} - \tilde{\Gamma}_{\substack{\mu\beta \\ \mu\beta \\ \beta\mu}}^\pi t_{j\pi B}^{i\alpha A} + \bar{\Gamma}_{\substack{1 \\ 1}}^A t_{j\beta B}^{i\alpha P} - \bar{\Gamma}_{\substack{1 \\ 1}}^P t_{j\beta P}^{i\alpha A} \end{array}$$

In this way four connection $\nabla_\theta, \theta \in \{1, \dots, 4\}$, on $X_M \subset GR_N$ are defined. We shall note the obtained structures $(X_M \subset GR_N, \nabla_\theta, \theta \in \{1, \dots, 4\})$.

2 New first and second kind integrability conditions of derivational equations

2.0. In [5] are obtained derivational equations of a submanifold in a GR_N , and in [6] integrability conditions of these equations in the structure $(X_M \subset GR_N, \nabla_\theta, \theta \in \{1, 2\})$. In the present work we engage in this problem for the structure $(X_M \subset GR_N, \nabla_\theta, \theta \in \{3, 4\})$.

As it is proved in [5] (Th. 1.2.), *derivational equations* in the considered case for a tangent are

$$(2.1) \quad B_{\alpha|\mu}^i = \sum_P \Omega_{P\alpha\mu}^\theta N_P^i, \quad \theta \in \{3, 4\},$$

and then for induced torsion in X_M is valid

$$(2.2) \quad \tilde{T}_{\beta\gamma}^\alpha = 0 \quad (\tilde{\Gamma}_{\beta\gamma}^\alpha = \tilde{\Gamma}_{\gamma\beta}^\alpha).$$

By virtue of the Th. 2.3. in [5], for unit normal is

$$(2.3) \quad N_{A|\mu}^i = -e_A \Omega_{A\rho\mu} h^{\pi\rho} B_\pi^i, \quad \theta \in \{3, 4\},$$

and

$$(2.4) \quad \bar{\Gamma}_1^A B_\mu = \bar{\Gamma}_2^A B_\mu = \bar{\Gamma}^A B_\mu,$$

in (1.12), and based on (1.8) in [5]

$$(2.5) \quad \frac{\Omega_{P\alpha\mu}}{\frac{1}{2}} = e_P H_{ij} N_P^i (B_{\alpha,\mu}^j + \Gamma_{\frac{pm}{mp}}^j B_\alpha^p B_\mu^m) = \frac{\Omega_{P\alpha\mu}}{\frac{3}{4}}.$$

In relation with (2.2,4), the addends in (1.12), related to X_M and to X_{N-M}^N are not different for separate kinds of derivatives, and (1.12) now becomes

$$(2.6) \quad \begin{aligned} t_{j\beta B|\mu}^{i\alpha A} &= t_{j\beta B,\mu}^{i\alpha A} + \Gamma_{\frac{mp}{mp}}^i t_{j\beta B}^{p\alpha A} B_\mu^m - \Gamma_{\frac{mj}{jm}}^p t_{p\beta B}^{i\alpha A} B_\mu^m \\ &+ \tilde{\Gamma}_{\pi\mu}^\alpha t_{j\beta B}^{i\pi A} - \tilde{\Gamma}_{\beta\mu}^\pi t_{j\pi B}^{i\alpha A} + \bar{\Gamma}_{P\mu}^A t_{j\beta B}^{i\alpha P} - \bar{\Gamma}_{B\mu}^P t_{j\beta P}^{i\alpha A}, \end{aligned}$$

where the coefficients $\tilde{\Gamma}$ are symmetric, and $\bar{\Gamma}$ are unique ($\bar{\Gamma}_1 = \bar{\Gamma}_2 = \bar{\Gamma}$). If in a differentiated tensor no exists indices as i, j, \dots , we write $|\mu$ instead of $|\mu_\theta$.

Using (2.1,3), we get (see (2.4) in [6])

$$(2.7) \quad \begin{aligned} B_{\alpha|\mu|\nu}^i - B_{\alpha|\nu|\mu}^i &= \sum_P [e_P h^{\pi\rho} (-\Omega_{P\alpha\mu} \Omega_{P\rho\nu} + \Omega_{P\alpha\nu} \Omega_{P\rho\mu}) B_\pi^i \\ &+ (\Omega_{P\alpha\mu|\nu} - \Omega_{P\alpha\nu|\mu}) N_P^i], \quad \theta, \omega \in \{3, 4\}. \end{aligned}$$

2.1. With respect of Ricci-type identities (12) and (13) from [2], and taking into consideration (2.2), we have

$$(2.8) \quad B_{\alpha|\mu|\nu}^i - B_{\alpha|\nu|\mu}^i = R_{\theta-2}^i{}^{pnm} B_\alpha^p B_\mu^m B_\nu^n - \tilde{R}_{\alpha\mu\nu}^\pi B_\pi^i, \quad \theta \in \{3, 4\},$$

where

$$(2.9a) \quad R_{jmn}^i = \Gamma_{jm,n}^i - \Gamma_{jn,m}^i + \Gamma_{jm}^p \Gamma_{pn}^i - \Gamma_{jn}^p \Gamma_{pm}^i,$$

$$(2.9b) \quad R_{2jmn}^i = \Gamma_{mj,n}^i - \Gamma_{nj,m}^i + \Gamma_{mj}^p \Gamma_{np}^i - \Gamma_{nj}^p \Gamma_{mp}^i$$

are **curvature tensors of the 1st, respectively 2nd kind** of GR_N and $\tilde{R}_{\beta\mu\nu}^\alpha$ is, with respect of (2.2), curvature tensor of $R_M \subset GR_N$.

We obtained in [6] three kinds integrability conditions for derivational equation of a tangent B_α^i , i.e. for $B_{\alpha|\mu}^i$, $\theta \in \{1, 2\}$. We shall consider here such conditions for $\theta \in \{3, 4\}$.

If one substitutes $\theta = \omega \in \{3, 4\}$ into (2.7) and compares with (2.8), taking into consideration (2.5) and (2.6), we get

$$(2.10) \quad \begin{aligned} R_{\theta-2jpmn}^i B_\alpha^p B_\mu^m B_\nu^n &= [\tilde{R}_{\alpha\mu\nu}^\pi - \sum_P e_P h^{\pi\rho} (\Omega_{P\alpha\mu}^\theta \Omega_{P\rho\nu}^\theta - \Omega_{P\alpha\nu}^\theta \Omega_{P\rho\mu}^\theta)] B_\pi^i \\ &+ \sum_P [\Omega_{P\alpha\mu|\nu}^\theta - \Omega_{P\alpha\nu|\mu}^\theta] N_P^i, \quad \theta \in \{3, 4\}, \end{aligned}$$

which are **the 1st and the 2nd integrability conditions** of derivational equation (2.1) in the structure $(X_M \subset GR_N, \nabla_\theta, \theta \in \{3, 4\})$.

a) Composing the previous equation with $H^{ij} B_\beta^j$, one gets

$$(2.11) \quad R_{\theta-2jpmn} B^j \beta B_\alpha^p B_\mu^m B_\nu^n = \tilde{R}_{\beta\alpha\mu\nu} - \sum_P e_P (\Omega_{P\alpha\mu}^\theta \Omega_{P\beta\nu}^\theta - \Omega_{P\alpha\nu}^\theta \Omega_{P\beta\mu}^\theta), \quad \theta \in \{3, 4\},$$

where

$$(2.12 a, b) \quad R_{\theta-2jpmn} = H_{ij} R_{\theta-2pmn}^i, \quad \tilde{R}_{\beta\alpha\mu\nu} = h_{\pi\beta} \tilde{R}_{\alpha\mu\nu}^\pi, \quad \theta \in \{3, 4\}.$$

Taking into count the antisymmetry of the tensors (2.12) with respect of the first two indices and substituting i in place of p , the equation (2.11) becomes

$$(2.13) \quad \tilde{R}_{\alpha\beta\mu\nu} = R_{\theta-2ijmn} B_\alpha^i B_\beta^j B_\mu^m B_\nu^n - \sum_P e_P (\Omega_{P\alpha\mu}^\theta \Omega_{P\beta\nu}^\theta - \Omega_{P\alpha\nu}^\theta \Omega_{P\beta\mu}^\theta), \quad \theta \in \{3, 4\},$$

which are **Gauss equations of the 1st and the 2nd kind** in the structure $(X_M \subset GR_N, \nabla_\theta, \theta \in \{3, 4\})$.

b) Composing the equation (2.10) with $H_{ij} N_Q^j$ we obtain finally

$$(2.14) \quad R_{\theta-2ijmn} B_\alpha^i N_Q^j B_\mu^m B_\nu^n = e_Q (\Omega_{Q\alpha\nu|\mu}^\theta - \Omega_{Q\alpha\mu|\nu}^\theta), \quad \theta \in \{3, 4\},$$

and that are **the 1st Codazzi equations of the 1st and the 2nd kind** at the cited structure.

2.2. Consider the same matter for the unit normal N_A^i . Using (2.3,1), we obtain (see (2.13) in [6]):

$$(2.15) \quad \begin{aligned} N_{A|\mu|\nu}^i - N_{A|\nu|\mu}^i &= -e_A h^{\pi\rho} [(\Omega_{A\rho\mu|\nu} - \Omega_{A\rho\nu|\mu})B_\pi^i \\ &+ \sum_P (\Omega_{A\rho\mu}\Omega_{P\pi\nu} - \Omega_{A\rho\nu}\Omega_{P\pi\mu})N_P^i]. \end{aligned}$$

In order to find corresponding Ricci-type identity for the left side of this equation for $\theta = \omega \in \{3, 4\}$, we use (2.6). Firstly, we have

$$(2.16) \quad N_{A|_3}^i = N_{A,\mu}^i + \Gamma_{pm}^i N_A^p B_\mu^m - \bar{\Gamma}_{A\mu}^P N_P^i,$$

and further

$$\begin{aligned} N_{A|_3|_3}^i &= (N_{A|_3}^i)_{,\nu} + \Gamma_{sn}^i N_{A|_3}^s B_\nu^n - \tilde{\Gamma}_{\mu\nu}^\sigma N_{A|_3}^s - \bar{\Gamma}_{A\nu}^S N_{S|_3}^i \\ &= N_{A,\mu\nu}^i + \Gamma_{pm,n}^i N_A^p B_\mu^m B_\nu^n + \Gamma_{pm}^i N_{A,\nu}^p B_\mu^m + \Gamma_{pm}^i N_A^p B_{\mu,\nu}^m \\ &\quad - \bar{\Gamma}_{A\mu,\nu}^P N_P^i - \bar{\Gamma}_{A\mu}^P N_{P,\nu}^i + \Gamma_{sn}^i N_{A,\mu}^s B_\nu^n + \Gamma_{sn}^i \Gamma_{pm}^s B_\nu^p N_A^m B_\mu^m \\ &\quad - \Gamma_{sn}^i N_P^s \bar{\Gamma}_{A\mu}^P B_\nu^n - \tilde{\Gamma}_{\mu\nu}^\sigma N_{A,\sigma}^i - \tilde{\Gamma}_{\mu\nu}^\sigma \Gamma_{pm}^i N_A^p B_\sigma^m + \tilde{\Gamma}_{\mu\nu}^\sigma \bar{\Gamma}_{A\sigma}^P N_P^i \\ &\quad - \bar{\Gamma}_{A\nu}^S N_{S,\mu}^i - \bar{\Gamma}_{A\nu}^S \Gamma_{pm}^i N_S^p B_\mu^m + \bar{\Gamma}_{A\nu}^S \bar{\Gamma}_{S\mu}^P N_P^i, \end{aligned}$$

wherefrom

$$(2.17) \quad N_{A|_3|\nu}^i - N_{A|\nu|_3}^i = \bar{R}_{1pmn}^i N_A^p B_\mu^m B_\nu^n - \bar{R}_{A\mu\nu}^P N_P^i,$$

where

$$(2.18) \quad \bar{R}_{B\mu\nu}^A = \bar{\Gamma}_{B\mu,\nu}^A - \bar{\Gamma}_{B\nu,\mu}^A + \bar{\Gamma}_{B\mu}^P \bar{\Gamma}_{P\nu}^A - \bar{\Gamma}_{B\nu}^P \bar{\Gamma}_{P\mu}^A,$$

is **curvature tensor of the space GR_N with respect to the normal submanifold** in the structure $(X_M \subset GR_N, \nabla, \theta \in \{3, 4\})$.

By means of the 4th kind of covariant derivative we obtain an equation corresponding to (2.17), and we conclude

$$(2.19) \quad N_{A|\mu|\nu}^i - N_{A|\nu|\mu}^i = R_{\theta-2pmn}^i N_A^p B_\mu^m B_\nu^n - \bar{R}_{A\mu\nu}^P N_P^i, \quad \theta \in \{3, 4\}.$$

If one substitutes into (2.15) $\theta = \omega \in \{3, 4\}$ and equalizes the right sides of obtained equation and (2.19), we get **the 1st and the 2nd kind integrability conditions of derivational equation (2.3)** in the structure $(X_M \subset GR_N, \nabla, \theta \in \{3, 4\})$:

$$(2.20) \quad \begin{aligned} R_{\theta-2pmn}^i N_A^p B_\mu^m B_\nu^n &= e_A h^{\pi\rho} (\Omega_{A\rho\mu|\nu} - \Omega_{A\rho\nu|\mu}) B_\pi^i \\ &+ [\bar{R}_{A\mu\nu}^P - e_A h^{\pi\rho} \sum_P (\Omega_{A\rho\mu}\Omega_{P\pi\nu} - \Omega_{A\rho\nu}\Omega_{P\pi\mu})] N_P^i, \quad \theta \in \{3, 4\}. \end{aligned}$$

a) If we compose this equation with $H_{ij}B_\beta^j$ one obtains an equation equivalent with (2.14), that is the 1st Codazzi equation of the 1st and the 2nd kind for the structure $(X_M \subset GR_N, \nabla, \theta \in \{3, 4\})$.

b) By composing the equation (2.20) with $H_{ij}N_B^j$, one obtains endly

$$(2.21) \quad R_{\theta-2}{}_{ijmn} N_A^i N_B^j B_\mu^m B_\nu^n = \bar{R}_{AB\mu\nu} + e_A e_B h^{\pi\rho} (\Omega_{A\pi\mu} \Omega_{B\rho\nu} - \Omega_{A\pi\nu} \Omega_{B\rho\mu}),$$

where

$$(2.22) \quad \bar{R}_{AB\mu\nu} = h_{AP} \bar{R}_{B\mu\nu}^P.$$

The equation (2.21) is the **2nd Codazzi equation of the 1st and the 2nd kind** for the structure $(X_M \subset GR_N, \nabla, \theta \in \{3, 4\})$.

Based on expressed above, the next theorems are valid:

Theorem 2.1. *The 1st and the 2nd kind integrability conditions for derivational equations (2.1), (2.3) in the in the structure $(X_M \subset GR_N, \nabla, \theta \in \{3, 4\})$ are given by equations (2.10), (2.20) respectively, where Ω_θ is given in (2.5), R_1, R_2 in (2.9), \tilde{R} is curvature tensor of the symmetric connection $\tilde{\Gamma}$, while \bar{R} is given in (2.18), (2.22).*

Theorem 2.2. *The Gauss equations of the 1st and the 2nd kind in the structure $(X_M \subset GR_N, \nabla, \theta \in \{3, 4\})$ are given in (2.13), the 1st Codazzi equations of the 1st and the 2nd kind in (2.14), and the 2nd Codazzi equations of the 1st and the 2nd kind in (2.21) in the same structure.*

3 Third kind integrability condition of derivational equations

3.1. Using simultaneously the 3rd and the 4th kind of covariant derivative by virtue of (2.6), we obtain Ricci-type identity (eq. (46) in [2]):

$$(3.1) \quad B_{\alpha|_3\mu|_4\nu}^i - B_{\alpha|_4\nu|_3\mu}^i = R_{4p\mu\nu}^i B_\alpha^p - \tilde{R}_{\alpha\mu\nu}^\pi B_\pi^i,$$

where

$$(3.2) \quad R_{4j\mu\nu}^i = (\Gamma_{jm,n}^i - \Gamma_{nj,m}^i + \Gamma_{jm}^p \Gamma_{np}^i - \Gamma_{nj}^p \Gamma_{pm}^i) B_\mu^m B_\nu^n + T_{jm}^i (B_{\mu,\nu}^m - \tilde{\Gamma}_{\nu\mu}^\pi B_\pi^m)$$

is curvature tensor of the 4th kind of GR_N with respect to $X_M \subset GR_N$.

On the other hand, if we put into (2.7) $\theta = 3$, $\omega = 4$ and compare the obtained equation with (3.1), we obtain **the 3rd kind integrability condition** of derivational equation (2.1) in the structure $(X_M \subset GR_N, \nabla, \theta \in \{3, 4\})$:

$$(3.3) \quad \begin{aligned} R_{4p\mu\nu}^i B_\alpha^p &= [\tilde{R}_{\alpha\mu\nu}^\pi - \sum_P e_P h^{\pi\rho} (\Omega_{1P\alpha\mu} \Omega_{2P\rho\nu} - \Omega_{2P\alpha\nu} \Omega_{1P\rho\mu})] B_\pi^i \\ &+ \sum_P (\Omega_{1P\alpha\mu|\nu} - \Omega_{2P\alpha\nu|\mu}) N_P^i. \end{aligned}$$

a) Composing previous equation with $H_{ij} B_\beta^j$, we get

$$R_{4jp\mu\nu} B_\beta^j B_\alpha^p = \tilde{R}_{\beta\alpha\mu\nu} - \sum_P e_P (\Omega_{1P\alpha\mu} \Omega_{2P\beta\nu} - \Omega_{2P\alpha\nu} \Omega_{1P\beta\mu}),$$

i.e., exchanging $j \rightarrow i$, $p \rightarrow j$, $\alpha \leftrightarrow \beta$, it follows that

$$(3.4) \quad \tilde{R}_{\alpha\beta\mu\nu} = R_{4ij\mu\nu} B_\alpha^i B_\beta^j - \sum_P e_P (\Omega_{1P\alpha\mu} \Omega_{2P\beta\nu} - \Omega_{2P\alpha\nu} \Omega_{1P\beta\mu}),$$

where

$$(3.5) \quad R_{4ij\mu\nu} = H_{ip} R_{4j\mu\nu}^p.$$

The equation (3.4) is **Gauss equation of the 3rd kind in the structure** $(X_M \subset GR_N, \nabla, \theta \in \{3, 4\})$.

b) Composing (3.4) with $H_{ij} N_Q^j$, we obtain

$$R_{4ij\mu\nu} N_Q^i B_\alpha^j = e_Q (\Omega_{1Q\alpha\mu|\nu} - \Omega_{2Q\alpha\nu|\mu}).$$

This is **the 1st Codazzi equation of the 3rd kind** in the cited structure.

3.2. On the base of (2.6) and (2.16) we have

$$\begin{aligned} N_{34}^i{}_{A|\mu|\nu} &= (N_{34}^i{}_{A|\mu}),_\nu + \Gamma_{ns}^i N_{A|\mu}^s B_\nu^n - \tilde{\Gamma}_{\mu\nu}^\sigma N_{A|\sigma}^i - \bar{\Gamma}_{A\nu}^S N_{S|\mu}^i \\ &= N_{A,\mu\nu}^i + \Gamma_{pm,n}^i N_A^p B_\mu^m B_\nu^n + \Gamma_{pm}^i N_{A,\nu}^p B_\mu^m + \Gamma_{pm}^i N_A^p B_{\mu,n}^m \\ &- \bar{\Gamma}_{A\mu,\nu}^P N_P^i - \bar{\Gamma}_{A\mu}^P N_{P,\nu}^i + \Gamma_{ns}^i N_{A,\mu}^s B_\nu^n + \Gamma_{ns}^i \Gamma_{pm}^s B_\nu^n N_A^p B_\mu^m \\ &- \Gamma_{ns}^i N_P^s \bar{\Gamma}_{A\mu}^P B_\nu^n - \tilde{\Gamma}_{\mu\nu}^\sigma N_{A,\sigma}^i - \tilde{\Gamma}_{\mu\nu}^\sigma \Gamma_{pm}^i N_A^p B_\sigma^m + \tilde{\Gamma}_{\mu\nu}^\sigma \bar{\Gamma}_{A\sigma}^P N_P^i \\ &- \bar{\Gamma}_{A\nu}^S N_{S,\mu}^i - \bar{\Gamma}_{A\nu}^S \Gamma_{pm}^i N_S^p B_\mu^m + \bar{\Gamma}_{A\nu}^S \bar{\Gamma}_{S\mu}^P N_P^i, \end{aligned}$$

and

$$(3.6) \quad N_{34}^i{}_{A|\mu|\nu} - N_{43}^i{}_{A|\nu|\mu} = R_{4p\mu\nu}^i N_A^p - \bar{R}_{A\mu\nu}^P N_P^i,$$

where R_4 is given in (3.2), and \bar{R} in (2.18).

By substituting into (2.15) $\theta = 3$, $\omega = 4$ and comparing the obtained equation with (3.6), we obtain **the 3rd kind integrability condition of derivational equation (2.3)** in the structure $(X_M \subset GR_N, \nabla_\theta, \theta \in \{3, 4\})$:

$$(3.11) \quad R_{4p\mu\nu}^i N_A^p = -e_A h^{\pi\rho} (\Omega_{1A\rho\mu|\nu} - \Omega_{2A\rho\nu|\mu}) B_\pi^i + [\bar{R}_{A\mu\nu}^P - e_A h^{\pi\rho} \sum_P (\Omega_{1A\rho\mu} \Omega_{2P\pi\nu} - \Omega_{2A\rho\nu} \Omega_{1P\pi\mu})] N_P^i.$$

a) Composing this equation with $H_{ij} B_\beta^j$ one obtains the equation of the form (3.5), that is the 1st Codazzi of the 3rd kind.

b) Composing (3.7) with $H_{ij} N_B^j$, we obtain **the 2nd Codazzi equation of the 3rd kind** in the above cited structure:

$$(3.8) \quad R_{4ij\mu\nu} N_A^i N_B^j = \bar{R}_{AB\mu\nu} + e_A e_B h^{\pi\rho} (\Omega_{1A\rho\mu} \Omega_{2B\pi\nu} - \Omega_{2A\rho\nu} \Omega_{1B\pi\mu}).$$

From exposed, the following theorems are valid.

Theorem 3.1. *The 3rd kind integrability conditions of derivational equations (2.1, 3) for $(X_M \subset GR_N, \text{with the structure } (X_M \subset GR_N, \nabla_\theta, \theta \in \{3, 4\}), \text{ where the connection } \nabla_\theta \text{ is defined in (2.6), are given:$*

- for tangents B_α^i by equation (3.3),
- for normals N_A^i by equation (3.7).

Theorem 3.2. *In the same structure (from the previous theorem) the Gauss equation of the 3rd kind for $X_M \subset GR_N$ is given in (3.4), the 1st Codazzi equation of the 3rd kind by (3.5), and the 2nd Codazzi equation of the 3rd kind by (3.8).*

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