

Electron. J. Probab. **23** (2018), no. 110, 1–61. ISSN: 1083-6489 https://doi.org/10.1214/18-EJP230

Non-Hermitian random matrices with a variance profile (I): deterministic equivalents and limiting ESDs

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Abstract

For each n, let $A_n = (\sigma_{ij})$ be an $n \times n$ deterministic matrix and let $X_n = (X_{ij})$ be an $n \times n$ random matrix with i.i.d. centered entries of unit variance. We study the asymptotic behavior of the empirical spectral distribution μ_n^Y of the rescaled entry-wise product

$$Y_n = \left(\frac{1}{\sqrt{n}}\sigma_{ij}X_{ij}\right).$$

For our main result we provide a deterministic sequence of probability measures μ_n , each described by a family of *Master Equations*, such that the difference $\mu_n^Y - \mu_n$ converges weakly in probability to the zero measure. A key feature of our results is to allow some of the entries σ_{ij} to vanish, provided that the standard deviation profiles A_n satisfy a certain quantitative irreducibility property. An important step is to obtain quantitative bounds on the solutions to an associate system of Schwinger-Dyson equations, which we accomplish in the general sparse setting using a novel graphical bootstrap argument.

Keywords: large random matrix; non-hermitian matrix; variance profile; empirical spectral distribution.

AMS MSC 2010: Primary 15B52, Secondary 15A18; 60B20.

Submitted to EJP on March 2, 2018, final version accepted on October 3, 2018.

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1 Introduction

For an $n \times n$ matrix M with complex entries and eigenvalues $\lambda_1, \ldots, \lambda_n \in \mathbb{C}$ (counted with multiplicity and labeled in some arbitrary fashion), the *empirical spectral distribution (ESD)* is given by

$$\mu_M = \frac{1}{n} \sum_{i=1}^n \delta_{\lambda_i} . \tag{1.1}$$

A seminal result in non-Hermitian random matrix theory is the *circular law*, which describes the asymptotic global distribution of the spectrum for matrices with i.i.d. entries of finite variance. The following strong form of the circular law was established by Tao and Vu [60], and is the culmination of the work of many authors [30, 48, 31, 15, 17, 33, 50, 61] – see the survey [19] for a detailed historical account.

Theorem 1.1 (Circular law). Let ξ be a complex random variable of zero mean and unit variance, and for each n let $X_n = (X_{ij}^{(n)})$ be an $n \times n$ matrix whose entries are i.i.d. copies of ξ . Then almost surely, the ESDs $\mu_{\frac{1}{2}X_n}$ converge weakly to the circular measure

$$\mu_{\rm circ}(dx\,dy):=\frac{1}{\pi} \mathbf{1}_{\{|x|^2+|y|^2\leq 1\}}\,dx\,dy.$$

One of the remarkable features of the circular law is that the asymptotic behavior of ESDs is insensitive to specific details of the entry distributions, apart from the first two moments. This is an instance of the *universality phenomenon* in random matrix theory.

The circular law has been an important tool for understanding the stability of dynamical systems on complex networks, going back to work of May in ecology [47], and later work of Sompolinski et al. in neuroscience [57]. May used an i.i.d. matrix X_n to model the *community matrix* for a food network of *n* species, where the entry X_{ij} determines the rate of growth (or decay) of the population of species *i* due to species *j*. The stability of the system is determined by the spectrum of X_n – specifically by whether it has eigenvalues with sufficiently large real part – and May used the circular law¹ to derive a criterion for stability.

¹At the time the circular law was only known to hold in the complex Gaussian case thanks to work of Ginibre and Mehta [30, 48]. Strictly speaking, May's argument assumes that there are asymptotically no eigenvalues outside the limiting support, which is now known to hold under some moment hypotheses [16, 21].

Recently there has been increasing interest in extending the arguments of [47, 57] to matrix models with more structured distributions. In neural networks, where random matrices are used to model the synaptic matrix, the work [53] considered perturbed i.i.d. matrices of the form $X_n + M_n$, where M_n is a fixed matrix with all entries within a fixed proportion of columns taking a fixed positive value μ^+ , and all remaining entries taking a fixed negative value μ^- . Their motivation was to conform to Dale's Law, stating that neurons are either inhibitory or excitatory. In this case M_n is a rank-one perturbation; as was later shown rigorously in work of Tao [58], low rank perturbations do not affect the limiting spectral distribution, but may lead to the creation of outlier eigenvalues.

Several recent works have studied the limiting spectral distribution for random matrices of the form

$$A_n \odot X_n = (\sigma_{ij} X_{ij}) \tag{1.2}$$

(suitably rescaled) where $A_n = (\sigma_{ij})$ is a fixed (deterministic) standard deviation profile. From a modeling perspective the σ_{ij} can reflect the varying degrees of interaction between species/neurons. In theoretical ecology the works [7, 8] considered asymmetric standard deviation profiles, i.e. taking $\sigma_{ij} \neq \sigma_{ji}$, in order to create more realistic predator-prey cascading relationships. In neuroscience the works [6, 5] considered matrices A_n partitioned into a bounded number of block submatrices having constant entries within each block, in order to model networks with a bounded number of cell types. We also note that predating these works, Girko [32, Chap. 25, 26] (see also the references therein) studied non-Hermitian matrices with standard deviation and mean profiles and provided *canonical equations* to describe the limiting spectral densities.

Some works have also gone beyond matrices with independent entries of specified mean and variance, for instance considering products and sums of deterministic matrices with a random matrix having i.i.d. entries [1], or allowing correlations between entries X_{ij}, X_{ji} [9]. We also mention that parallel to the study of non-Hermitian matrices there have been many works devoted to the study of Hermitian random matrices with a variance profile, both Wigner and Gram-type – see for instance Girko [32, Chapter 7, 8], Shlyakhtenko [56], Guionnet [35], Anderson and Zeitouni [12], Hachem et al. [36], Ajanki et al. [2].

As has been pointed out in the ecology literature [8], a key feature that is missing from the literature on models of the form (1.2) is to allow A_n to have zero entries. Indeed, the nodes in large real-world ecological or neural networks do not interact with all other nodes. One fix has been to take A_n to have i.i.d. Bernoulli(p) indicator entries, independent of X_n , i.e. to model the support of the network by a sparse Erdős–Rényi digraph. As was shown by Wood [64] the circular law still holds for $A_n \odot X_n$ (after rescaling by $(pn)^{-1/2}$) if $p \ge n^{\alpha-1}$ for any fixed $\alpha \in (0, 1]$. However, the valence of the nodes in the resulting network is highly concentrated around pn, while the valence distribution for real-world networks is highly non-uniform [8]. With the ability to set A_n deterministically one can reflect some known underlying geometry of the network.

In the present work, our main focus is to understand the asymptotic global spectral distribution under minimal assumptions on the variance profile. Key features of our results, which also present significant challenges in the proofs, are to allow a large number of the entries σ_{ij} to be zero, as well as to allow asymmetry (i.e. $\sigma_{ij} \neq \sigma_{ji}$). As an example, we are able to handle variance profiles satisfying a *robust irreducibility* condition (see Definition 2.7), which includes matrices that are close in the cut norm to an irreducible graphon (Lemma 2.5).

As compared with the proof of the circular law, the identification and description of a limiting measure is significantly more involved. In this article we prove the existence of a sequence of deterministic measures – called *deterministic equivalents* – which asymptotically approximate the random ESDs. In particular we obtain non-trivial information

even when the ESDs themselves do not converge to a limit. The identification of the deterministic equivalents involves analysis of a (cubic) polynomial system of *Master Equations* determined by the variance profile. A relative of the Master Equations known as the Quadratic Vector Equation was studied in recent work of Ajanki, Erdős and Krüger and Alt, Erdős and Krüger on the spectrum of Hermitian matrices with a variance profile [3, 10].

Since the initial release of this paper, a local law version of our main statement (Theorem 2.3) was proved in [11] under the restriction that the standard deviation profile σ_{ij} is uniformly strictly positive and that the distribution of the matrix entries possesses a bounded density and has all its moments finite. In this case, it is also proved that the density of the deterministic equivalents is positive and bounded on its support. The latter properties may no longer hold if the standard deviation profile has zero entries or is not uniformly lower bounded. In these cases, the limiting distribution may offer a wider variety of behavior, such as a blowup or vanishing density at zero, or a point mass at zero; see [22].

It is by now well known that the study of ESDs for non-Hermitian random matrices is intimately connected with proving the quantitative invertibility of such matrices – that is, establishing lower tail estimates for small singular values. The possible sparsity of the matrices considered here gives rise to significant challenges for this task. Bounds on the smallest singular value sufficient for our purposes were established by the first author in [23]. In the present work we obtain control on the remaining small singular values from Wegner-type bounds, which are established by a quantitative analysis of the Master Equations. Specifically, the key is to show solutions to the *Regularized Master Equations* (see (2.7)) are uniformly bounded in the spectral scale parameter t and the dimension n. For this task we use an iterative graphical bootstrap argument (reminiscent of bootstrap arguments from the theory of differential equations) that exploits expansion properties of a directed graph naturally associated to the variance profile. (Graph expansion properties were also key for the analysis of the smallest singular value in [23].) We discuss these aspects of the proof in more detail after presenting the results in Section 2.

1.1 The model

In this article we study the following general class of random matrices with nonidentically distributed entries.

Definition 1.2 (Random matrix with a variance profile). For each $n \ge 1$, let A_n be a (deterministic) $n \times n$ matrix with entries $\sigma_{ij}^{(n)} \ge 0$, let X_n be a random matrix with i.i.d. entries $X_{ij}^{(n)} \in \mathbb{C}$ satisfying

$$\mathbb{E}X_{11}^{(n)} = 0, \quad \mathbb{E}|X_{11}^{(n)}|^2 = 1$$
 (1.3)

and set

$$Y_n = \frac{1}{\sqrt{n}} A_n \odot X_n \tag{1.4}$$

where \odot is the matrix Hadamard product, i.e. Y_n has entries $Y_{ij}^{(n)} = \frac{1}{\sqrt{n}} \sigma_{ij}^{(n)} X_{ij}^{(n)}$. The empirical spectral distribution of Y_n is denoted by μ_n^Y . We refer to A_n as the standard deviation profile and to $A_n \odot A_n = ((\sigma_{ij}^{(n)})^2)$ as the variance profile. We additionally define the normalized variance profile as $V_n = n^{-1}A_n \odot A_n$. When no ambiguity occurs, we will drop the index n and simply write σ_{ij}, X_{ij}, V , etc.

Remark 1.1. Note we do not assume the $(X_{ij}^{(n)})$'s are independent or identically distributed for different n's.

Our goal is to describe the asymptotic behavior of the ESDs μ_n^Y given a sequence of standard deviation profiles A_n which can be sparse, and may not converge in any sense to a limiting standard deviation profile.

1.2 Master equations and deterministic equivalents

The main result of this article states that under certain assumptions on the sequence of standard deviation profiles A_n and the distribution of the entries of X_n , there exists a sequence of deterministic probability measures μ_n that are *deterministic equivalents* of the spectral measures μ_n^Y such that

$$\mu_n^Y \sim \mu_n$$
 in probability $(n \to \infty)$,

see Section 2.1.2 for notation.

The measures μ_n are described by a polynomial system of *Master Equations*. Denote by V_n^{T} the transpose matrix of V_n and by $\rho(V_n)$ its spectral radius. For a parameter $s \ge 0$, the Master Equations are the following system of 2n + 1 equations in 2n unknowns $q_1, \ldots, q_n, \tilde{q}_1, \ldots, \tilde{q}_n$:

$$\begin{cases} q_i &= \frac{(V_n^{\mathsf{T}} \boldsymbol{q})_i}{s^2 + (V_n \tilde{\boldsymbol{q}})_i (V_n^{\mathsf{T}} \boldsymbol{q})_i} \\ \tilde{q}_i &= \frac{(V_n \tilde{\boldsymbol{q}})_i}{s^2 + (V_n \tilde{\boldsymbol{q}})_i (V_n^{\mathsf{T}} \boldsymbol{q})_i} , \quad q_i, \tilde{q}_i \ge 0, \ i \in [n], \\ \sum_{i \in [n]} q_i = \sum_{i \in [n]} \tilde{q}_i \end{cases}$$
(1.5)

where q, \tilde{q} are the $n \times 1$ column vectors with components q_i, \tilde{q}_i , respectively.

If $s \ge \sqrt{\rho(V_n)}$, it can be proved that the Master Equations admit the unique trivial solution $(q, \tilde{q}) = 0$. Provided that $0 < s < \sqrt{\rho(V_n)}$ and that the matrix V_n is irreducible, it can be shown that the Master Equations admit a unique positive solution (q, \tilde{q}) which depends only on s. This solution $s \mapsto (q(s), \tilde{q}(s))$ is continuous on $(0, \infty)$. With this definition of q and \tilde{q} , the deterministic equivalent μ_n is defined as the radially symmetric probability distribution over \mathbb{C} satisfying

$$\mu_n\{z \in \mathbb{C}, |z| \le s\} = 1 - \frac{1}{n} \boldsymbol{q}^\mathsf{T}(s) V_n \widetilde{\boldsymbol{q}}(s) , \quad s > 0 .$$

It readily follows that the support of μ_n is contained in the disk of radius $\sqrt{\rho(V_n)}$. **Acknowledgments.** The work of NC was partially supported by NSF grants DMS-1266164 and DMS-1606310. The work of WH and JN was partially supported by the program "modèles numériques" of the French Agence Nationale de la Recherche under the grant ANR-12-MONU-0003 (project DIONISOS). Support of the Labex BEZOUT from the Université Paris Est is also acknowledged. The work of DR was partially supported by Austrian Science Fund (FWF): M2080-N35. The authors thank the referee for their careful reading and insights which enabled to substantially shorten the paper.

2 Presentation of the results

2.1 Notational preliminaries

Denote by [n] the set $\{1, \dots, n\}$ and let $\mathbb{C}_+ = \{z \in \mathbb{C}, \operatorname{Im}(z) > 0\}$. For $\mathcal{X} = \mathbb{C}$ or \mathbb{R} , let $C_c(\mathcal{X})$ (resp. $C_c^{\infty}(\mathcal{X})$) the set of $\mathcal{X} \to \mathbb{R}$ continuous (resp. smooth) and compactly supported functions. Let $\mathcal{B}(z, r)$ be the open ball of \mathbb{C} with center z and radius r. If $z \in \mathbb{C}$, then \bar{z} is its complex conjugate; let $i^2 = -1$. The Lebesgue measure on \mathbb{C} will be either

denoted by $\ell(dz)$ or dxdy. For $x, y \in \mathbb{R}$ we write $\max(x, y) = x \lor y$ and $\min(x, y) = x \land y$. The cardinality of a finite set S is denoted by |S|. For $S \subset [n]$ and when clear from the context we will abbreviate $S^c = [n] \setminus S$.

2.1.1 Matrices

We denote by $\mathbf{1}_n$ the $n \times 1$ vector of 1's. Given two $n \times 1$ vectors $\boldsymbol{u}, \boldsymbol{v}$, we denote their scalar product $\langle \boldsymbol{u}, \boldsymbol{v} \rangle = \sum_{i \in [n]} \bar{u}_i v_i$.

For a given matrix A, denote by A^{T} its transpose and by A^* its conjugate transpose. Denote by I_n the $n \times n$ identity matrix. If clear from the context, we omit the dimension. For $a \in \mathbb{C}$ and when clear from the context, we sometimes write a instead of a I and similarly write a^* instead of $(aI)^* = \bar{a}I$. For matrices A, B of the same dimensions we denote by $A \odot B$ their Hadamard, or entrywise, product (i.e. $(A \odot B)_{ij} = A_{ij}B_{ij}$).

Given two Hermitian matrices A and B, the notations $A \ge B$ and A > B refer to the usual positive semidefinite ordering. Notations \succ and \succcurlyeq refer to the elementwise inequalities for real matrices or vectors. Namely, if A and B are real matrices,

$$A \succ B \iff A_{ij} > B_{ij} \quad \forall i, j \quad \text{and} \quad A \succcurlyeq B \iff A_{ij} \ge B_{ij} \quad \forall i, j.$$

The notation $A \succcurlyeq_{\neq} 0$ stands for $A \succcurlyeq 0$ and $A \neq 0$. Given a matrix A, ||A|| refers to its spectral norm, and $||A|||_{\infty}$ to its max-row norm, defined as: $|||A|||_{\infty} := \max_{i \in [n]} \sum_{j=1}^{n} |A_{ij}| = \max_{\|u\|_{\infty} \leq 1} ||Au||_{\infty}$. We denote the spectral radius of an $n \times n$ matrix A by $\rho(A) := \max \{ |\lambda| : \lambda \text{ is an eigenvalue of } A \}$. If A is a square matrix, we write $\operatorname{Im}(A) = (A - A^*)/(2i)$. For an n-dimension vector a, $\operatorname{diag}(a_i)$ denotes the $n \times n$ diagonal matrix with a as its diagonal elements.

2.1.2 Convergence of measures

Given probability distributions ν_n, ν over some set \mathcal{X} (= \mathbb{R} or \mathbb{C}), we will denote the weak convergence of ν_n to ν by $\nu_n \xrightarrow[n \to \infty]{\mathbf{w}} \nu$. If ν_n is random, $\nu_n \xrightarrow[n \to \infty]{\mathbf{w}} \nu$ almost surely (resp. in probability) stands for the fact that for all $f \in C_c(\mathcal{X})$,

$$\int f \, d\nu_n \xrightarrow[n \to \infty]{} \int f \, d\nu \quad \text{almost surely (resp. in probability)}.$$

Let (μ_n) and (ν_n) be deterministic sequences of probability distributions over \mathcal{X} , and let (ν_n) be tight, i.e. for all $\varepsilon > 0$, one can find a compact set $\mathcal{K}_{\varepsilon}$ such that $\sup_n \nu_n(\mathcal{X} \setminus \mathcal{K}_{\varepsilon}) \leq \varepsilon$. We will denote by

$$\mu_n \sim \nu_n$$
 as $n \to \infty$

the fact that the signed measure $\mu_n - \nu_n$ weakly converges to zero, i.e. $\int f d\mu_n - \int f d\nu_n \rightarrow 0$ for all $f \in C_c(\mathcal{X})$. If the sequence (μ_n) is random while (ν_n) is deterministic and tight, then

 $\mu_n \sim \nu_n$ almost surely (resp. in probability)

stands for

$$\int f \, d\mu_n - \int f \, d\nu_n \xrightarrow[n \to \infty]{} 0 \quad \text{almost surely (resp. in probability)},$$

for all $f \in C_c(\mathcal{X})$. Notice that $\mu_n \sim \nu_\infty$ is equivalent to $\mu_n \xrightarrow[n \to \infty]{\mathbf{w}} \nu_\infty$.

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2.1.3 Stieltjes transforms

Let μ be a nonnegative finite measure on $\mathbb R$ and

$$g_{\mu}(\eta) = \int \frac{\mu(d\lambda)}{\lambda - \eta} , \quad \eta \in \mathbb{C}_+$$
 (2.1)

its Stieltjes transform. Then the following properties are standard

(i)
$$g_{\mu}(\eta) \in \mathbb{C}_+$$
, (ii) $|g_{\mu}(\eta)| \leq \frac{\mu(\mathbb{R})}{\operatorname{Im}(\eta)}$, (iii) $\lim_{y \to +\infty} -iyg_{\mu}(iy) = \mu(\mathbb{R})$.

Moreover, $-(z + g_{\mu}(z))^{-1}$ is the Stieltjes transform of a probability measure, see for instance [63, Theorem B.3]. In particular

$$|z + g_{\mu}(z)|^{-1} \le \text{Im}^{-1}(z), \quad z \in \mathbb{C}_{+}.$$
 (2.2)

2.1.4 Graph-theoretic notation

Given an $n \times n$ non-negative matrix $A = (\sigma_{ij})$ we form a directed graph $\Gamma = \Gamma(A)$ on the vertex set [n] that puts an edge $i \to j$ whenever $\sigma_{ij} > 0$. We denote the out-neighborhood of a vertex $i \in [n]$ in the graph Γ by

$$\mathcal{N}_A(i) := \{ j \in [n] : \sigma_{ij} > 0 \}.$$
(2.3)

Consequently, the in-neighborhood is denoted $\mathcal{N}_{A^{\mathsf{T}}}(i)$. For a set $S \subset [n]$ we write

$$\mathcal{N}_A(S) := \bigcup_{i \in S} \mathcal{N}_A(i) = \{ j \in [n] : \mathcal{N}_{A^{\mathsf{T}}}(j) \cap S \neq \emptyset \}.$$
(2.4)

For $\delta \in (0,1)$ we denote the associated *densely-connected out-neighbors* of a set $S \subset [n]$ by

$$\mathcal{N}_A^{(\delta)}(S) = \{ j \in [n] : |\mathcal{N}_{A^{\mathsf{T}}}(j) \cap S| \ge \delta |S| \}.$$

$$(2.5)$$

To obtain quantitative results we will generally work with the graph associated to the matrix

$$A(\sigma_0) = (\sigma_{ij} \mathbf{1}_{\sigma_{ij} \ge \sigma_0}) \tag{2.6}$$

which only keeps the entries exceeding a fixed cutoff parameter $\sigma_0 > 0$, setting the remaining entries to zero.

2.2 Model assumptions

We will establish results concerning sequences of matrices Y_n as in Definition 1.2 under various additional assumptions on A_n and X_n , which we now summarize. We note that many of our results only require a subset of these assumptions.

For our main result we need the following additional assumption on the distribution of the entries of X_n .

A0 (Moments). We have $\mathbb{E}|X_{11}^{(n)}|^{4+\varepsilon} \leq M_0$ for all $n \geq 1$ and some fixed $\varepsilon > 0$, $M_0 < \infty$.

Remark 2.1. Assumption A0 is needed to apply the results from [23] to bound the smallest singular value as in (2.10). It is also used in Section 3 to quantitatively bound the difference between our random measures and their deterministic equivalents, which is crucial for obtaining logarithmic integrability of singular value distributions. This latter step can likely be accomplished with fewer moments, but we do not pursue this.

A1 (Bounded variances). There exists $\sigma_{\max} \in (0,\infty)$ such that

$$\sup_{n} \max_{1 \le i,j \le n} \sigma_{ij}^{(n)} \le \sigma_{\max}$$

Remark 2.2 (Convention). While we will keep the value σ_{max} generic in the statements, we will always set it to 1 in the proofs, with no loss of generality.

In order to express the next key assumption, we need to introduce the following *Regularized Master Equations* which are a specialization of the Schwinger–Dyson equations of Girko's Hermitized model associated to Y_n (see Section 2.5 for further discussion). The following is proved in Section 4.1.

Proposition 2.1 (Regularized Master Equations). Let $n \ge 1$ be fixed, let A_n be an $n \times n$ nonnegative matrix and write $V_n = \frac{1}{n}A_n \odot A_n$. Let s, t > 0 be fixed, and consider the following system of equations

$$\begin{cases} r_i = \frac{(V_n^{\mathsf{T}} \boldsymbol{r})_i + t}{s^2 + ((V_n \widetilde{\boldsymbol{r}})_i + t)((V_n^{\mathsf{T}} \boldsymbol{r})_i + t)} \\ \widetilde{r}_i = \frac{(V_n \widetilde{\boldsymbol{r}})_i + t}{s^2 + ((V_n \widetilde{\boldsymbol{r}})_i + t)((V_n^{\mathsf{T}} \boldsymbol{r})_i + t)} \end{cases},$$
(2.7)

where $\mathbf{r} = (r_i)$ and $\tilde{\mathbf{r}} = (\tilde{r}_i)$ are $n \times 1$ vectors. Denote by $\mathbf{\vec{r}} = \begin{pmatrix} \mathbf{r} \\ \tilde{\mathbf{r}} \end{pmatrix}$. Then this system admits a unique solution $\mathbf{\vec{r}} = \mathbf{\vec{r}}(s,t) \succ 0$. This solution satisfies the identity

$$\sum_{i \in [n]} r_i = \sum_{i \in [n]} \widetilde{r}_i.$$
(2.8)

A2 (Admissible variance profile). Let $\vec{r}(s,t) = \vec{r}_n(s,t) \succ 0$ be the solution of the Regularized Master Equations for given $n \ge 1$. For all s > 0, there exists a constant C = C(s) > 0 such that

$$\sup_{n \ge 1} \sup_{t \in (0,1]} \frac{1}{n} \sum_{i \in [n]} r_i(s,t) \le C .$$

A variance profile V_n for which the previous estimate holds is called *admissible*.

Remark 2.3. Assumption A2 may seem obscure at first sight as it necessitates to solve the regularized master equations to check if a variance profile is admissible. In particular, it is not clear if this assumption is compatible with some sparsity. In Section 2.4, we provide sufficient conditions on the variance profile V_n which imply A2, namely A3 (lower bound on V_n), A4 (symmetric V_n) and A5 (robust irreducibility for V_n).

2.3 Statement of the results

Recall the Master Equations (1.5), and notice that these equations are obtained from the Regularized Master Equations (2.7) by letting the parameter t go to zero. Notice however that condition $\sum q_i = \sum \tilde{q}_i$ is now required for uniqueness and not a consequence as in (2.7).

In order to prove existence of solutions q, \tilde{q} to the Master Equations we need to assume the standard deviation profile A_n is irreducible. This is equivalent to assuming the associated digraph $\Gamma(A_n)$ is strongly connected (recall the notations from Section 2.1.4). This will cause no loss of generality, as we can conjugate Y_n by an appropriate permutation matrix to put A_n in block-upper-triangular form with irreducible blocks on the diagonal – the spectrum of Y_n is then the union of the spectra of the block diagonal submatrices. **Theorem 2.2** (Master equations). Let $n \ge 1$ be fixed, let A_n be an $n \times n$ nonnegative matrix and write $V_n = \frac{1}{n}A_n \odot A_n$. Assume that A_n is irreducible. Then the following hold:

- 1. For $s \ge \sqrt{\rho(V_n)}$ the system (1.5) has the unique solution $\vec{q}(s) = 0$.
- 2. For $s \in (0, \sqrt{\rho(V)_n})$ the system (1.5) has a unique non-trivial solution $\vec{q}(s) \geq 0$. Moreover, this solution satisfies $\vec{q}(s) > 0$.
- 3. The function $s \mapsto \vec{q}(s)$ defined in parts (1) and (2) is continuous on $(0, \infty)$ and is continuously differentiable on $(0, \sqrt{\rho(V)}) \cup (\sqrt{\rho(V)}, \infty)$.

Remark 2.4 (Convention). Above and in the sequel we abuse notation and write $\vec{q} = \vec{q}(s)$ to mean a solution of the equation (1.5), understood to be the nontrivial solution for $s \in (0, \sqrt{\rho(V_n)})$.

The main result of this paper is the following.

Theorem 2.3 (Main result). Let $(Y_n)_{n\geq 1}$ be a sequence of random matrices as in Definition 1.2, and assume A0, A1 and A2 hold. Assume moreover that A_n is irreducible for all $n \geq 1$.

1. There exists a sequence of deterministic measures $(\mu_n)_{n\geq 1}$ on \mathbb{C} such that

$$\mu_n^Y \sim \mu_n$$
 in probability.

2. Let $\vec{q}(s)^{\mathsf{T}} = (q(s)^{\mathsf{T}} \ \widetilde{q}(s)^{\mathsf{T}})$ be as in Theorem 2.2, and for $s \in (0, \infty)$ put

$$F_n(s) = 1 - \frac{1}{n} \langle \boldsymbol{q}(s), V_n \widetilde{\boldsymbol{q}}(s) \rangle.$$
(2.9)

Then F_n extends to an absolutely continuous function on $[0, \infty)$ which is the CDF of a probability measure with support contained in $[0, \sqrt{\rho(V_n)}]$ and continuous density on $(0, \sqrt{\rho(V_n)})$.

3. For each $n \ge 1$ the measure μ_n from part (1) is the unique radially symmetric probability measure on \mathbb{C} with $\mu_n(\{z : |z| \le s\}) = F_n(s)$ for all $s \in (0, \infty)$.

Remark 2.5 (Almost sure convergence under different hypotheses). As with Theorem 1.1, a key component of the proof is a lower tail estimate for the smallest singular value of scalar shifts $Y_n - zI_n$ of the form

$$\mathbb{P}(s_n(Y_n - zI_n) \le n^{-\beta}) = O(n^{-\alpha})$$
(2.10)

holding for a.e. fixed $z \in \mathbb{C}$. (Crucially, we do not need such an estimate for every $z \in \mathbb{C}$, as **A2** only requires s = |z| > 0 and allows variance profiles for which (2.10) is false when z = 0.) Such a bound is available for arbitrary fixed $z \neq 0$ and $\alpha > 0$ a small constant by recent work of the first author [23]; see Proposition 6.1. Obtaining (2.10) with $\alpha > 1$ would immediately improve conclusion (1) to almost sure convergence by an application of the Borel–Cantelli lemma.

Such improvements are already available under stronger assumptions on A_n and X_n . For instance, under A3 and replacing A0 with a bounded density assumption, an easy argument gives (2.10) for any fixed $\alpha > 0$ and some $\beta = \beta(\alpha) > 0$; see [19, Section 4.4]. Remark 2.6 (Density of μ_n versus density of F_n). In the previous theorem, by the classical polar change of coordinates, the density φ_n of μ_n for $0 < |z| < \sqrt{\rho_V}$ is given by the formula:

$$\varphi_n(|z|) = \frac{1}{2\pi|z|} \frac{d}{ds} F_n(s) \Big|_{s=|z|} = -\frac{1}{2\pi n|z|} \frac{d}{ds} \langle \boldsymbol{q}(s), V \widetilde{\boldsymbol{q}}(s) \rangle \Big|_{s=|z|}$$

As an illustration we show how our results recover the circular law for matrices with i.i.d. entries.

Example 2.1 (The circular law). Consider a standard deviation profile A_n with all elements equal to 1 and assume that **A0** holds. It is well known in this case that $\mu_n^Y \xrightarrow{w} \mu_{\text{circ}}$ in probability (and even almost surely), where μ_{circ} stands for the circular law with density $\pi^{-1}1_{\{|z|\leq 1\}}$. We can recover this result with Theorem 2.3. In this case, both systems (2.7) and (1.5) simplify into a single equation:

$$r_i \equiv r = \frac{r+t}{s^2 + (r+t)^2}$$
, $r > 0$ and $q_i \equiv q = \frac{q}{s^2 + q^2}$, $q \ge 0$. (2.11)

From the first equation, one can prove that $r(s,t) \leq 1$ for $t \in (0,1]$. In fact,

$$r = \frac{r+t}{s^2 + (r+t)^2} \le \frac{1}{r+t} \quad \Rightarrow \quad r^2 + rt \le 1 \quad \Rightarrow \quad r^2 \le 1 \; .$$

Hence A2 is fulfilled. The second equation has the unique nontrivial solution

$$q(s) = \begin{cases} \sqrt{1 - s^2} & 0 \le s \le 1\\ 0 & s \ge 1 \end{cases}.$$
 (2.12)

Consequently, $F_n(s) = s^2$ for $s \le 1$. From Remark 2.6, we conclude the desired convergence.

In the next example, we prove that a doubly stochastic normalized variance profile is admissible and that the associated deterministic equivalent μ_n is the circular law.

Example 2.2 (Doubly stochastic variance profile). Assume that matrix V_n is doubly stochastic, i.e. $\frac{1}{n} \sum_i \sigma_{ij}^2 = \frac{1}{n} \sum_j \sigma_{ij}^2 = 1$ for all $1 \le i, j \le n$. Then, one quickly verifies that the vectors $\vec{r} = r\mathbf{1}$ and $\vec{q} = q\mathbf{1}$ with r, q as in (2.11) respectively satisfy the Regularized Master Equations and the Master Equations. As a consequence **A2** can be established as in Example 2.1. Let now **A1** hold and assume that the variance profile V_n is irreducible for all $n \ge 1$ then one can apply Theorem 2.3 with μ_n equal to the circular law.

Remark 2.7. Note that under A1 the doubly stochastic condition implies that the number of non-zero entries in each row and column is linear in n.

In the following theorem, we relax the irreducibility assumption, which requires some additional argument.

Theorem 2.4 (The circular law for doubly stochastic variance profiles). Let $(Y_n)_{n\geq 1}$ be a sequence of random matrices as in Definition 1.2, and assume **A0** and **A1** hold. Suppose also that the normalized variance profiles V_n are doubly stochastic, i.e. $\frac{1}{n}\sum_i \sigma_{ij}^2 = \frac{1}{n}\sum_j \sigma_{ij}^2 = 1$ for all $1 \leq i, j \leq n$. Then $\mu_n^Y \xrightarrow{w} \mu_{\text{circ}}$ in probability.

This parallels results of Girko [32, §7.11, §8.2] and Anderson and Zeitouni [12] (see also [56], [35] for the Gaussian case) for random Hermitian matrices with a doubly-stochastic variance profile which obey the Marchenko–Pastur or Wigner semi-circle laws.

2.4 Sufficient conditions for admissibility

Hereafter we introduce a series of assumptions directly checkable over matrices (V_n) without solving a priori the regularized master equations. These assumptions enforce **A2**.

The simplest such assumption is to enforce uniform positivity of the variances, which allows one to bypass some of the most technical portions of our argument. This assumption was also made in the recent work [11].

A3 (Lower bound on variances). There exists $\sigma_{\min} > 0$ such that

$$\inf_{n} \min_{1 \le i,j \le n} \sigma_{ij}^{(n)} \ge \sigma_{\min}$$

We generalize A3 below with the expansion-type condition A5.

Proposition 2.5. Let $A = (\sigma_{ij})$ be an $n \times n$ matrix with entries $\sigma_{ij} \ge \sigma_{\min} > 0$ for some $\sigma > 0$. Let $\vec{r} \succ 0$ be the unique solution of the Regularized Master Equations (2.7). Then

$$\frac{1}{n}\sum_{i=1}^{n}r_i \le \frac{1}{\sigma_{\min}}$$

In particular, if $A_n = (\sigma_{ij}^{(n)})$ is a sequence of standard deviation profiles as in Definition 1.2 for which **A3** holds, then **A2** is satisfied, i.e. V_n is admissible.

A4 (Symmetric variance profile). For all $n \ge 1$, the normalized variance profile (or equivalently the standard deviation profile) is symmetric:

$$V_n = V_n^\mathsf{T}$$

Proposition 2.6. Let $A = (\sigma_{ij})$ be a symmetric matrix with nonnegative entries, and let $\vec{r} \succ 0$ be the unique solution of the Regularized Master Equations (2.7). Then

$$\frac{1}{n}\sum_{i=1}^{n}r_i \le \frac{1}{2s}.$$

In particular, if $A_n = (\sigma_{ij}^{(n)})$ is a sequence of standard deviation profiles as in Definition 1.2 for which **A4** holds, then **A2** is satisfied.

We now introduce the following strengthening of the irreducibility assumption, which can be understood as a kind of expansion condition on an associated directed graph.

Definition 2.7 (Robust irreducibility). For $\delta, \kappa \in (0, 1)$ we say that a nonnegative $n \times n$ matrix A is (δ, κ) -robustly irreducible if the following hold:

1. For all $i \in [n]$,

$$\left|\mathcal{N}_{A}(i)\right|, \left|\mathcal{N}_{A^{\mathsf{T}}}(i)\right| \geq \delta n.$$
(2.13)

2. For all $S \subset [n]$ with $1 \leq |S| \leq n - 1$,

$$\left|\mathcal{N}_{A^{\mathsf{T}}}^{(\delta)}(S) \cap S^{c}\right| \geq \min(\kappa|S|, |S^{c}|). \tag{2.14}$$

For comparison, a nonnegative $n \times n$ matrix A is irreducible if and only if $\mathcal{N}_{A^{\mathsf{T}}}(S) \cap S^c \neq \emptyset$ for all $S \subset [n]$ with $1 \leq |S| \leq n-1$. Thus, a matrix A satisfying the conditions of Definition 2.7 is "robustly irreducible" in the sense that A remains irreducible even after setting a small linear proportion of entries equal to zero.

Remark 2.8 (Relation to broad connectivity). In their work on permanent estimators, Rudelson and Zeitouni assume a stronger expansion-type condition on $A(\sigma_0)$ which they call broad connectivity [54]. The conditions for a nonnegative square matrix A to be (δ, κ) -broadly connected are the same as in Definition 2.7, except (2.14) is replaced by the stronger condition

$$\left|\mathcal{N}_{A^{\mathsf{T}}}^{(\delta)}(S)\right| \ge \min(n, (1+\kappa)|S|) \tag{2.15}$$

for all nonempty $S \subset [n]$.

Recall the definition (2.6) of $A(\sigma_0)$.

A5 (Robust irreducibility). There exist constants $\sigma_0, \delta, \kappa \in (0, 1)$ such that for all $n \ge 1$, $A_n(\sigma_0)$ is (δ, κ) -robustly irreducible.

Note that **A5** enables variance profiles with a large proportion of vanishing entries and is implied by **A3**.

Theorem 2.8. Consider a sequence of standard deviation profiles $A_n = (\sigma_{ij}^{(n)})$ as in Definition 1.2, and assume that **A1** and **A5** hold. Then **A2** holds, i.e. V_n is admissible.

It turns out that the mere irreducibility of V_n provides a weaker form of A2.

Proposition 2.9. Let V_n be an irreducible variance profile and let $\vec{r} = \vec{r}(s,t)$ be the solution of the associated Regularized Master Equations (2.7). Then there exists C = C(s,n) such that

$$\sup_{t \in (0,1]} \frac{1}{n} \sum_{i \in [n]} r_i(s,t) \leq C \, .$$

The main difference here is that constant C depends on n and may blow up with n. Depending on the variance profile, this proposition is sometimes sufficient to verify A2. **Example 2.3** (Variance profile with a block structure). Let $k \ge 1$ be a fixed integer, and $M = (m_{ij})_{i,j \in [k]}$ be a $k \times k$ irreducible matrix with nonnegative elements. Let $J_m = \mathbf{1}_m \mathbf{1}_m^{\mathsf{T}}$. Assume that n = km ($m \ge 1$) and consider the $n \times n$ matrix

$$V_{n} = \frac{1}{n} \begin{pmatrix} m_{11}J_{m} & \cdots & m_{1k}J_{m} \\ \vdots & & \\ m_{k1}J_{m} & \cdots & m_{kk}J_{m} \end{pmatrix}.$$
 (2.16)

Then V_n is admissible, i.e. **A2** is fulfilled. In fact, V_n is irreducible and its block structure implies that

$$\boldsymbol{r}^{\mathsf{T}} = (\underbrace{\rho_1, \cdots, \rho_1}_{m \text{ times}}, \cdots, \underbrace{\rho_k, \cdots, \rho_k}_{m \text{ times}}), \quad \tilde{\boldsymbol{r}}^{\mathsf{T}} = (\underbrace{\tilde{\rho}_1, \cdots, \tilde{\rho}_1}_{m \text{ times}}, \cdots, \underbrace{\tilde{\rho}_k, \cdots, \tilde{\rho}_k}_{m \text{ times}})$$

where $\rho = (\rho_i)$ and $\tilde{\rho} = (\tilde{\rho}_i)$ satisfy the 2k equations

$$\rho_i = \frac{(M_k^{\mathsf{T}} \boldsymbol{\rho})_i + t}{s^2 + ((M_k^{\mathsf{T}} \boldsymbol{\rho})_i + t)(M_k \tilde{\boldsymbol{\rho}})_i + t)}, \quad \tilde{\rho}_i = \frac{(M_k \tilde{\boldsymbol{\rho}})_i + t}{s^2 + ((M_k^{\mathsf{T}} \boldsymbol{\rho})_i + t)(M_k \tilde{\boldsymbol{\rho}})_i + t)}, \quad i \in [k]$$

with $M_k = \frac{1}{k}M$. In particular,

$$\sup_{t \in (0,1]} \frac{1}{n} \sum_{i \in [n]} r_i(s,t) = \sup_{t \in (0,1]} \frac{1}{k} \sum_{i \in [k]} \rho_i(s,t) \,,$$

where the latter is finite by Proposition 2.9 and does not depend on n, hence A2.

We now state two general classes of standard deviation profiles satisfying A5 that include many examples not covered by A3, A4 or Example 2.3. The classes are in a similar spirit, essentially saying that $A = (\sigma_{ij})$ is approximately controlled from below by a block matrix whose nonzero blocks form an irreducible "pattern", and which are "regular" in the sense that the nonzero entries in the blocks are uniformly distributed in a certain sense. In what follows we write

$$e_A(S,T) = \sum_{i \in S, j \in T} 1_{\sigma_{ij} > 0}, \qquad \rho_A(S,T) = \frac{e_A(S,T)}{|S||T|}.$$

We recall a common assumption from the extremal combinatorics literature (see for instance [45, Definition 1.6]). For $\varepsilon, \rho_0 \in (0, 1)$, we say that an $m \times n$ matrix A is (ε, ρ_0) -super-regular if for every $S \subset [m], T \subset [n]$ with $|S| \ge \varepsilon m, |T| \ge \varepsilon n$ we have $\rho_A(S,T) \ge \rho_0$. Informally, this says that all sub-matrices of A of relative size ε have density (proportion of non-zero entries) at least ρ_0 .

Definition 2.10 (Block-pseudorandom irreducibility). For $\delta, \varepsilon, \rho_0 \in (0, 1)$ and $K \in \mathbb{N}$, we say that a non-negative $n \times n$ matrix A is $(\delta, K, \varepsilon, \rho_0)$ -block-pseudorandomly irreducible, or $(\delta, K, \varepsilon, \rho_0)$ -BPI for short, if (2.13) holds, and if there is a partition of [n] into sets T_1, \ldots, T_K with $T_q \ge n/(2K)$ for each $1 \le q \le K$, and a $K \times K$ irreducible 0/1 matrix $M = (m_{pq})$, such that for each $(p,q) \in [K]^2$ with $m_{pq} = 1$, the sub-matrix A_{T_p,T_q} is (ε, ρ_0) -super-regular.

Thus, the BPI condition is met if we can partition [n] into a bounded number of parts of roughly equal size, such that the partitioned matrix contains an irreducible "pattern" of super-regular blocks. We note the ways in which the assumption that $A(\sigma_0)$ is BPI generalizes Example 2.3:

- 1. The blocks do not need to be the same size.
- 2. The matrix A does not have to be constant within blocks (and can in fact have a large proportion of zero entries in every block, as well as an arbitrary proportion of entries larger that σ_0).
- 3. *A* is arbitrary outside the sub-matrices A_{T_p,T_q} with $m_{pq} = 1$ (whereas in Example 2.3 the entries of *A* were zero there).

We also note that this assumption generalizes the *block fully indecomposable* assumption from [3] (see Definition 2.9 there), which makes a stronger assumption on the connectivity pattern M than irreducibility (that it is *fully indecomposable*) and also requires entries to be uniformly bounded below in blocks for which $m_{pg} = 1$.

Lemma 2.4. Let $A = (\sigma_{ij})$ be an $n \times n$ non-negative matrix that is $(\delta, K, \varepsilon, \rho)$ -BPI for some $\delta, \rho \in (0, 1), \varepsilon \in (0, \delta/8)$, and $K \in \mathbb{N}$. Then A is (δ_0, κ) -robustly irreducible with $\delta_0 = \min(\delta/2, \rho/(4K))$ and $\kappa = \min(\delta/4, 1/(8K))$.

Proof. The condition (2.13) is immediate. For (2.14), we first consider the case that $|S| \le \delta n/4$. From (2.13) we have

$$\delta n|S| \le \sum_{i=1}^{n} \sum_{j \in S} 1_{\sigma_{ij} > 0} = \sum_{i=1}^{n} |\mathcal{N}_A(i) \cap S| \le |\mathcal{N}_{A^{\mathsf{T}}}^{(\delta/2)}(S)||S| + n(\delta/2)|S|$$

which rearranges to give

$$|\mathcal{N}_{A^{\mathsf{T}}}^{(\delta/2)}(S) \cap S^c| \geq |\mathcal{N}_{A^{\mathsf{T}}}^{(\delta/2)}(S)| - |S| \geq \delta n/2 - |S| \geq \delta n/4$$

whenever $|S| \leq \delta n/4$. Since $n \geq |S|$ we conclude (2.14) holds in this case with $\kappa = \delta/4$.

Assume now that $|S| > \delta n/4$. Let $G = \{q \in [K] : |S \cap T_q| \ge \frac{1}{2}|S||T_q|/n\}$. From the pigeonhole principle we know that G is nonempty. Suppose first that $G \neq [K]$. Then since M is irreducible, there exists $(p,q) \in G^c \times G$ such that $m_{pq} = 1$. Now since $|S \cap T_q| \ge |S||T_q|/(2n) \ge (\delta/8)|T_q|$, and since $\varepsilon < \delta/8$, we deduce from the super-regularity of A_{T_p,T_q} that $|\mathcal{N}_A(i) \cap (S \cap T_q)| \ge \rho |S \cap T_q|$ for at least $(1-\varepsilon)|T_p|$ values of $i \in T_p$. Indeed, supposing otherwise, if T'_p is the set of $i \in T_p$ for which this lower bound fails then we have $e_A(T'_p, S \cap T_q) = \sum_{i \in T_q} |\mathcal{N}_A(i) \cap (S \cap T_q)| < \rho |T'_p||S \cap T_q|$, which is only possible if $|T'_p| < \varepsilon |T_p|$. Thus, we have

$$|\mathcal{N}_A(i) \cap S| \ge |\mathcal{N}_A(i) \cap (S \cap T_q)| \ge \rho |S \cap T_q| \ge \rho |S| |T_q| / (2n) \ge \frac{\rho}{4K} |S|$$

for at least

$$(1-\varepsilon)|T_p| - |S \cap T_p| \ge \left(1-\varepsilon - \frac{|S|}{2n}\right)|T_p| \ge \left(\frac{1}{2} - \varepsilon\right)|T_p| \ge \frac{1}{2K}\left(\frac{1}{2} - \varepsilon\right)n \ge \frac{1}{8K}n$$

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values of $i \in S^c$, where in the penultimate inequality of the former display we used that $q \in G$, and in the first inequality of the latter display we used that $p \in G^c$. The claim follows in this case.

It remains to handle the case that $|S| > \delta n/4$ and G = [K]. In this case, repeating the lines above shows that $|N_A(i) \cap S| \ge \frac{\rho}{4K} |S|$ for at least $(1 - \varepsilon)|T_p|$ values of $i \in T_p$ for every $1 \le p \le K$ (because for every $p \in [K]$ we must have $m_{pq} = 1$ for at least one value of q). We hence have $|\mathcal{N}_{A^{\mathsf{T}}}^{(\rho/(4K))}(S)| \ge (1 - \varepsilon)n$. But since $\varepsilon < \delta/8$ and $|S| > \delta n/4$ we conclude that $|\mathcal{N}_{A^{\mathsf{T}}}^{(\rho/(4K))}(S) \cap S^c| \ge \delta n/4 \ge (\delta/4)|S|$ as desired. \Box

Our next sufficient condition for robust irreducibility involves notions from the theory of graph limits, in particular the concept of a graphon. For further background we refer to [46]. For our purposes, let us say that a *graphon* is a (Lebesgue) measurable function $W : [0, 1]^2 \rightarrow [0, 1]$. We equip the set of graphons with the cut norm:

$$\|W\|_{\square} = \sup_{S,T \subseteq [0,1]} \left| \int_{S \times T} W(x,y) dx dy \right|,$$

where the integral is taken with respect to product Lebesgue measure, and the supremum is taken over all measurable subsets of [0, 1]. To a non-negative $n \times n$ matrix $A = (\sigma_{ij})$ we associate the graphon W_A which is equal to σ_{ij} on the set $[\frac{i-1}{n}, \frac{i}{n}) \times [\frac{j-1}{n}, \frac{j}{n})$ for each $1 \leq i, j \leq n$ (and set $W_A(x, 1) = W_A(1, x) \equiv 0$).

We say that a graphon U is a *stepfunction* if there is a measurable partition $\mathcal{P} = \{S_1, \ldots, S_K\}$ of [0, 1] and a $K \times K$ matrix $M = (m_{pq})$ such that $U(x, y) = m_{pq}$ for all $p, q \in [K]$ and all $(x, y) \in S_p \times S_q$. We call \mathcal{P}, M the partition and pattern, respectively, associated with U.

Definition 2.11. Let $K \in \mathbb{N}$, and $\alpha, \sigma_* \in (0, 1)$. Say that a graphon W is (K, α, σ_*) -*irreducible* if there exists a stepfunction U with $W \ge U$ pointwise, and such that the following hold for the partition $\mathcal{P} = \{S_1, \ldots, S_K\}$ and pattern M associated with U:

- 1. Leb $(S_p) \ge \alpha/K$ for all $p \in [K]$.
- 2. $m_{pq} \in \{0\} \cup [\sigma_*, 1]$ for all $p, q \in [K]$.
- 3. M is irreducible.

Lemma 2.5. Let $A = (\sigma_{ij})$ be a non-negative $n \times n$ matrix and let $K \in N$ and $\delta, \alpha, \sigma_* \in (0,1)$. Assume (2.13) holds, and that there is a (K, α, σ_*) -irreducible graphon W such that $||W_A - W||_{\Box} \leq \frac{\delta^2 \alpha^2 \sigma_*}{32K^2} =: \delta_0$. Then A is $(\delta_0/2, \delta_0 \delta/4)$ -robustly irreducible.

Proof. Arguing as in the proof of Lemma 2.4 it suffices to verify the condition (2.14) for $|S| \ge \delta n/4$. If $|S| \ge (1 - \delta/2)n$ then from (2.13) it follows from the triangle inequality that $|\mathcal{N}_A(i) \cap S| \ge \delta n/2$ for all $i \in [n]$, and (2.14) follows in this case. We henceforth assume $|S^c| \ge \delta n/2$. We have

$$\frac{1}{n^2}e_A(S^c,S) = \int_{\hat{S}\times\hat{S}^c} W_A(x,y)dxdy \ge \int_{\hat{S}\times\hat{S}^c} W(x,y)dxdy - \delta_0 \ge \int_{\hat{S}\times\hat{S}^c} U(x,y)dxdy - \delta_0,$$

where for $S \subset [n]$ we denote the dilation $\hat{S} := \frac{1}{n} \cdot S$. Let $L \subset [K]$ such that $1 \leq |L| \leq K-1$, and such that $\operatorname{Leb}(\hat{S} \cap S_p) \geq (\delta/4) \operatorname{Leb}(S_p)$ for all $p \in L$ and $\operatorname{Leb}(\hat{S}^c \cap S_p) \geq (\delta/4) \operatorname{Leb}(S_p)$ for all $p \in L^c$. It is possible to choose L such that both L and L^c are nonempty since $|\hat{S}|, |\hat{S}^c| \geq \delta/4$. Since M is irreducible we must have $m_{p'q'} \geq \sigma_*$ for some $(p,q) \in L \times L^c$.

Now we have

$$\begin{aligned} \int_{\hat{S}\times\hat{S}^c} U(x,y)dxdy &\geq \sum_{p,q=1}^K m_{p,q}\operatorname{Leb}(\hat{S}\cap S_p)\operatorname{Leb}(\hat{S}^c\cap S_q) \\ &\geq \sigma_*(\delta/4)^2\operatorname{Leb}(S_{p'})\operatorname{Leb}(S_{q'}) \geq \frac{\delta^2\alpha^2\sigma_*}{16K^2} = 2\delta_0 \end{aligned}$$

Combined with the previous display, we conclude $e_A(S^c, S) \ge \delta_0 n^2$. On the other hand,

$$e_A(S^c, S) = \sum_{i \in S^c} |\mathcal{N}_A(i) \cap S| \le |S| |\mathcal{N}_{A^{\mathsf{T}}}^{(\delta_0/2)}(S)| + \frac{\delta_0}{2} |S| |S^c|$$

which rearranges to $|\mathcal{N}_{A^{\mathsf{T}}}^{(\delta_0/2)}(S)| \geq (\delta_0/2)|S^c| \geq (\delta\delta_0/4)|S|$, as desired.

2.5 Outline of the proof

As is well known in the literature devoted to large non-Hermitian random matrices, the spectral behavior of such matrices can be studied with the help of the so-called Girko's Hermitization procedure, which is intimately related with the *logarithmic potential* of their spectral measure [31]. By definition, the logarithmic potential U_{μ} of a probability measure μ on \mathbb{C} which integrates $\log |\cdot|$ near infinity is the function from \mathbb{C} to $(-\infty, \infty]$ defined by

$$U_{\mu}(z) = -\int_{\mathbb{C}} \log |\lambda - z| \, \mu(d\lambda) \, .$$

Writing z = x + iy the Laplace operator is $\Delta = \partial_{xx}^2 + \partial_{yy}^2 = 4 \partial_z \partial_{\bar{z}}$ where $\partial_z = \frac{1}{2} (\partial_x - i\partial_y)$ and $\partial_{\bar{z}} = \frac{1}{2} (\partial_x + i\partial_y)$. The probability measure μ can be recovered by the formula $\mu = -\frac{1}{2\pi}\Delta U_{\mu}$, valid in the set $\mathcal{D}'(\mathbb{C})$ of the Schwartz distributions, which means that

$$\int \psi(z)\,\mu(dz) = -\frac{1}{2\pi} \int_{\mathbb{C}} \Delta \psi(z)\,U_{\mu}(z)\,\ell(dz) \qquad \text{for all} \quad \psi \in C_{c}^{\infty}(\mathbb{C}) \;.$$

Now, setting $\mu = \mu_n^Y$, the logarithmic potential can be written as

$$U_{\mu_n^Y}(z) = -\frac{1}{n} \sum_{i=1}^n \log|\lambda_i - z| = -\frac{1}{n} \log|\det(Y_n - z)| = -\frac{1}{n} \log\sqrt{\det(Y_n - z)(Y_n - z)^*}$$
$$= -\int_0^\infty \log(x) L_{n,z}(dx),$$

where $L_{n,z} := \frac{1}{n} \sum_{i=1}^{n} \delta_{s_{i,z}}$ is the empirical distribution of the singular values $s_{1,z} \ge \cdots \ge s_{n,z} \ge 0$ of $Y_n - z$. For technical reasons, it is easier to consider the symmetrized empirical distribution of the singular values

$$\check{L}_{n,z} := \frac{1}{2n} \sum_{i=1}^{n} \delta_{-s_{i,z}} + \frac{1}{2n} \sum_{i=1}^{n} \delta_{s_{i,z}}$$
(2.17)

for which a similar identity holds:

$$U_{\mu_n^Y}(z) = -\int_{\mathbb{R}} \log |x| \,\check{L}_{n,z}(dx) \,.$$

This identity is at the heart of Girko's strategy. In a word, it shows that in order to evaluate the asymptotic behavior of the spectral distribution μ_n^Y , we can focus on the asymptotic behavior of $\check{L}_{n,z}$ for almost all $z \in \mathbb{C}$. By considering $\check{L}_{n,z}$, we are in the more

familiar world of Hermitian matrices. Informally, for all $z \in \mathbb{C}$, we will find a sequence $(\check{\nu}_{n,z})_{n \in \mathbb{N}}$ of deterministic probability measures such that $\check{L}_{n,z} \sim \check{\nu}_{n,z}$, and

$$\int_{\mathbb{R}} \log |x| \, \check{L}_{n,z}(dx) \approx \int_{\mathbb{R}} \log |x| \, \check{\nu}_{n,z}(dx) \quad \text{for large n .}$$

Setting

$$h_n(z) := -\int_{\mathbb{R}} \log |x| \,\check{\nu}_{n,z}(dx),$$

gives that for all $\psi \in C_c^{\infty}(\mathbb{C})$,

$$\int \psi(z)\,\mu_n^Y(dz)\approx -\frac{1}{2\pi}\int_{\mathbb{C}}\Delta\psi(z)\,h_n(z)\,\ell(dz)\quad\text{for large n}\,,$$

showing that $h_n(z)$ is the logarithmic potential of a probability measure μ_n , and in particular that $\mu_n^Y \sim \mu_n$. Further smoothness properties of $h_n(z)$ will finally yield the properties of μ_n stated in Theorem 2.3.

We provide more details hereafter with precise pointers to the article's results.

1. Study of the associated Hermitian model

This topic is covered in Section 3.

Given $z \in \mathbb{C}$, we establish the existence of a sequence $(\check{\nu}_{n,z})_{n\in\mathbb{N}}$ of deterministic probability measures such that $\check{L}_{n,z} \sim \check{\nu}_{n,z}$ almost surely. To this end, we introduce the $2n \times 2n$ Hermitian matrix

$$\boldsymbol{Y}_{n}^{z} := \begin{pmatrix} 0 & Y-z \\ Y^{*}-z^{*} & 0 \end{pmatrix}, \quad z \in \mathbb{C},$$
(2.18)

with spectral measure $L_{n,z}$ [42, Theorem 7.3.7]. The asymptotic analysis of $L_{n,z}$ relies on the *resolvent*:

$$\boldsymbol{R}_{n}(z,\eta) := \begin{pmatrix} -\eta & Y-z \\ Y^{*}-z^{*} & -\eta \end{pmatrix}^{-1}, \quad \eta \in \mathbb{C}_{+},$$
(2.19)

of \boldsymbol{Y}_{n}^{z} . The mere definition (2.1) of a Stieltjes transform yields

$$g_{\check{L}_{n,z}}(\eta) = \frac{1}{2n} \operatorname{tr} \boldsymbol{R}_n(z,\eta).$$

The rigorous use of the Stieltjes transform for the study of ESDs of Hermitian random matrices goes back to Pastur [52], and was further developed by Bai to obtain quantitative results [13, 14]. Beginning with the seminal works [26, 27] of Erdő, Schlein and Yau this approach has been used to show that the semicircle law governs the spectral distribution for Wigner matrices down to near-optimal scales. In these works, the basic strategy is to use resolvent identities to show that the Stieltjes transform approximately satisfies a fixed-point equation, sometimes called the *Schwinger-Dyson* (or *master-loop*) equation. This approach was extended to Hermitian matrices with doubly stochastic variance profile [28]. However, for Hermitian matrices with more general variance profiles and non-zero mean it becomes necessary to consider a *system* of equations that are approximately satisfied by individual diagonal entries of the resolvent.

In Section 3 we derive the Schwinger–Dyson equations for our setting, cf. Proposition 3.1:

$$\begin{cases} p_{i} = \frac{(V_{n}^{\mathsf{T}}\boldsymbol{p})_{i} + \eta}{|z|^{2} - ((V_{n}\tilde{\boldsymbol{p}})_{i} + \eta)((V_{n}^{\mathsf{T}}\boldsymbol{p})_{i} + \eta)} \\ \tilde{p}_{i} = \frac{(V_{n}\tilde{\boldsymbol{p}})_{i} + \eta}{|z|^{2} - ((V_{n}\tilde{\boldsymbol{p}})_{i} + \eta)((V_{n}^{\mathsf{T}}\boldsymbol{p})_{i} + \eta)} \end{cases}, \qquad \boldsymbol{p} = (p_{i}), \ \tilde{\boldsymbol{p}} = (\tilde{p}_{i}), \tag{2.20}$$

for $\eta \in \mathbb{C}_+$, with unique solution the $2n \times 1$ vector $\vec{p} = (p \ \tilde{p})$ satisfying $\operatorname{Im} \vec{p} \succ 0$. We prove that there exists a probability distribution $\check{\nu}_{n,z}$ whose Stieltjes transform $g_{\check{\nu}_{n,z}}$ is defined as $g_{\check{\nu}_{n,z}} = \frac{1}{n} \sum_{i \in [n]} p_i$.

In Theorem 3.2, it is established that for all $\eta \in \mathbb{C}_+$, $g_{\check{L}_{n,z}}(\eta) - g_{\check{\nu}_{n,z}}(\eta) \to 0$ almost surely, which in particular implies that $\check{L}_{n,z} \sim \check{\nu}_{n,z}$ a.s. In Proposition 3.6 we obtain a quantitative estimate along the imaginary axis of the form

$$\mathbb{E}g_{\check{L}_{n,z}}(\mathrm{i}t) - g_{\check{\nu}_{n,z}}(\mathrm{i}t) = \mathcal{O}\left(\frac{1}{t^c\sqrt{n}}\right), \quad t > 0,$$
(2.21)

for some integer *c*.

We note that recently there has been much work analyzing the Schwinger–Dyson equations corresponding to Hermitian random matrices with mean and variance profiles satisfying a range of assumptions. For the centered case one is led to the so-called *Quadratic Vector Equation*, which has been thoroughly analyzed in works of Ajanki, Erdős and Krüger and Alt, Erdős and Krüger [3, 10]; the application to universality for local spectral statistics was carried out in [2]. Very recently they have made the extension to matrices with correlated entries, which involves the study of the so-called *Matrix Dyson Equation* [4]. In another recent work, He, Knowles and Rosenthal prove an approximate (matrix-valued) self-consistent equation for resolvents of Hermitian random matrices with arbitrary mean and variance profile, which covers the structure of the model (2.18) [39]. However, they assume the entries have all moments finite, and their aim is to obtain a local law at the optimal scale. In the present work, a sub-optimal quantitative analysis of the system (2.20) under few moments will suffice for our purposes of understanding the spectrum of the associated non-Hermitian model Y_n at global scale.

2. From the spectral measures $\check{L}_{n,z} \sim \check{\nu}_{n,z}$ to the spectral measures $\mu_n^Y \sim \mu_n$ via the associated logarithmic potentials

This topic is covered in Sections 3 (partly) and 6.

The fact that $\tilde{L}_{n,z} \sim \check{\nu}_{n,z}$ a.s. does not ensure that the random logarithmic potential $U_{\mu_n^Y}(z)$ becomes close to the deterministic logarithmic potential $h_n(z)$ (assuming the latter is well defined). Essentially, this is due to the fact that $x \mapsto \log |x|$ is unbounded near zero and infinity. While the singularity at infinity is easily handled using the almost sure tightness of the measures $\check{L}_{n,z}$, the singularity at zero presents a major technical challenge (indeed, this hurdle was the reason it took decades to establish the circular law under the optimal hypotheses). We show that under the admissibility assumption **A2**, $x \mapsto \log |x|$ is $\check{\nu}_{n,z}$ -integrable, and that for all $\tau, \tau' > 0$, there is $\varepsilon > 0$ small enough such that

$$\mathbb{P}\left\{\left|\int_{-\varepsilon}^{\varepsilon} \log|x|\,\check{L}_{n,z}(dx)\right| > \tau\right\} < \tau' \quad \text{for all large } n.$$
(2.22)

Together with the almost sure tightness and weak convergence $L_{n,z} \sim \check{\nu}_{n,z}$, we can show that $U_{\mu_n^Y}(z) - h_n(z) \rightarrow_n 0$ in probability. The almost sure convergence is an open problem not covered in this article.

The proof of (2.22) is based on two ingredients. The first one is a result from [23] by the first author giving a lower tail estimate for the smallest singular value of $Y_n - z$ for arbitrary fixed $z \in \mathbb{C} \setminus \{0\}$ under the sole assumption **A1** on the standard deviation profile A_n . For the second result, established in Section 6, we show that there exist two constants $C, \gamma_0 > 0$ such that for all x > 0,

$$\mathbb{E}\mathring{L}_{n,z}((-x,x)) \leq C(x \vee n^{-\gamma_0}) .$$

Bounds of this form on the expected density of states for random Hermitian operators are sometimes referred to as *local Wegner estimates* (after the work [62]). Their application to the convergence of the empirical spectral distribution for non-Hermitian random matrices goes back to Bai's proof of the circular law [15]; our presentation of the argument is closer to the one in [34].

Such an estimate follows from control on $\text{Im }\mathbb{E}g_{\tilde{L}_{n,z}}(\text{i}t)$ for small t > 0, is obtained in two steps. First, sufficient control of $\mathbb{E}g_{\tilde{L}_{n,z}}(\text{i}t) - g_{\tilde{\nu}_{n,z}}(\text{i}t)$ is already provided by the estimate (2.21).

Second, we rely on A2 to state that $g_{\vec{\nu}_{n,z}}(it)$ is bounded independent of n and t. For this task, a variation of the Schwinger–Dyson equations, namely the *Regularized Master Equations*, introduced in Proposition 2.1, is obtained simply by setting $\eta = it$ in the Schwinger–Dyson equations (2.20) – in this case, $\vec{p} \in i\mathbb{R}^{2n}$ and $\vec{r} = (r^{\mathsf{T}} \tilde{r}^{\mathsf{T}})^{\mathsf{T}}$ is defined as $\vec{r}(t) = \operatorname{Im} \vec{p}(it)$. Hence by A2,

$$\operatorname{Im} g_{\check{\nu}_{n,z}}(\mathrm{i}t) = \frac{1}{n} \sum_{i \in [n]} r_i(t) \le C$$
(2.23)

for some $C < \infty$ independent of n and t (depending only on |z| and the parameters in our assumptions).

3. Description of the deterministic probability measure μ_n

This is covered in Sections 4 and 7.

We have proved so far that $\mu_n^Y \sim \mu_n$ in probability, where μ_n is the probability measure whose logarithmic potential $U_{\mu_n}(z)$ coincides with $h_n(z)$. It remains to establish the properties of μ_n that are stated by Theorem 2.3.

In Section 4 we prove Theorem 2.2. Our approach to obtaining the solution $\vec{q}(s)$ of (1.5) is through the Regularized Master Equations (2.7). Since these equations are obtained by a simple transformation of the Schwinger–Dyson equations (2.20), by our work in Section 3 we know that (2.7) has a unique solution $\vec{r}(s,t)$ satisfying $\vec{r}(s,t) \succ 0$, where we write s = |z|. We then show that the pointwise limit $\vec{r}_*(s) := \lim_{t\downarrow 0} \vec{r}(s,t)$ exists, and moreover that $\vec{q} = \vec{r}_*$ and is the unique solution to (1.5). Having properly defined $\vec{q}(s)$, our main task is to show that the distribution $-(2\pi)^{-1}\Delta h_n(z)$ in fact defines a density on the set $\mathcal{D} := \{z \in \mathbb{C} : |z| \neq 0, |z| \neq \sqrt{\rho(V_n)}\}$ and to provide an expression for this density. The general approach towards solving this problem can be found in the physics literature (see [29]). Define on $\mathbb{C} \times (0, \infty)$ the functions

$$\mathcal{U}_{n}^{Y}(z,t) := -\frac{1}{2n} \log \det((Y-z)^{*}(Y-z) + t^{2}) \quad \text{and} \quad \mathcal{U}_{n}(z,t) := -\frac{1}{2} \int_{\mathbb{R}} \log(x^{2} + t^{2}) \check{\nu}_{n,z}(dx) \,.$$
(2.24)

For fixed t > 0, these functions can be seen as regularized versions of the logarithmic potentials $U_{\mu_n^Y}(z)$ and $U_{\mu_n}(z)$ respectively, which converge back as $t \downarrow 0$, in $\mathcal{D}'(\mathbb{C})$.

On the other hand, let us consider again the resolvent $R_n(z,\eta)$ introduced in (2.19). Using the well-known formula for the inverse of a partitioned matrix [42, §0.7.3] and writing

$$\boldsymbol{R}_{n}(z,\eta) =: \begin{pmatrix} G_{n}(z,\eta) & F_{n}(z,\eta) \\ F_{n}'(z,\eta) & \widetilde{G}_{n}(z,\eta) \end{pmatrix},$$
(2.25)

we get that by setting $\eta = it$, that $\partial_{\bar{z}} \mathcal{U}_n^Y(z,t)$ coincides with $n^{-1} \text{tr} F_n(z,it)$. Relying on the asymptotic analysis made in Section 3 on the resolvent \mathbf{R}_n , we can easily obtain an expression for $\partial_{\bar{z}} \mathcal{U}_n(z,t)$ by considering the asymptotic behavior of $n^{-1} \text{tr} F_n(z,it)$. We

then conclude by studying the equation

$$\Delta U_{\mu_n} = 4 \lim_{t \downarrow 0} \partial_z \partial_{\bar{z}} \mathcal{U}_n(z, t).$$

Section 7.1 is devoted to these questions.

In Section 7.2 we conclude the proof of Theorem 2.4. As we noted above, the key is that (2.23) is easily obtained under the double stochasticity assumption by examining the explicit solution \vec{r} to the Regularized Master Equations.

4. Sufficient conditions for A2 to hold

This topic is covered in Section 5.

While (2.23) can be proved in a few lines under **A3** (see Proposition 2.5) or **A4** (see Proposition 2.6), establishing such a bound under the more general robust irreducibility assumption **A5** is significantly more technical. Here it is helpful to view the standard deviation profile in terms of the associated directed graph $\Gamma(A(\sigma_0))$ (which was defined in Section 2.1). The basic idea is that the equations (2.7) encode relationships between the size of components $r_i(t)$, $\tilde{r}_i(t)$ at a vertex *i* to the sizes of the components at neighboring vertices. Assuming toward a contradiction that $r_{i_0}(t)$ is large at some vertex i_0 , we can use the robust irreducibility assumption to propagate the property of having large $r_i(t)$ to most of the other vertices *i*. We can also use the equations (2.7) to show that $\tilde{r}_i(t)$ will consequently be small for most *i*. However, this contradicts the crucial trace identity $\sum_{i=1}^{n} r_i(t) = \sum_{i=1}^{n} \tilde{r}_i(t)$, which derives from the fact that the matrix **R** in (2.19) satisfies

$$\sum_{i=1}^{n} \mathbf{R}_{ii} = \sum_{i=1}^{n} \mathbf{R}_{n+i,n+i}.$$

See Section 5 for further details.

We remark that under the stronger broad connectivity assumption on the standard deviation profile (see Remark 2.8), Wegner-type estimates that are sufficient for the purposes of this paper were obtained by the first author by a completely different argument, following a geometric approach introduced by Tao and Vu in [60] – see [24, Theorem 4.5.1].

2.6 **Open questions**

Relaxing the robust irreducibility assumption While control on the smallest singular value is proved under very general conditions (see Proposition 6.1), we have made the additional robust irreducibility assumption **A5** in order to handle the other small singular values via Wegner estimates. Would it be possible to lighten this assumption?

Almost sure convergence One may want to upgrade the convergence $\mu_n^Y \sim \mu_n$ in probability in Theorem 2.3 to almost sure convergence, as discussed in Remark 2.5.

Extension to sparse models While our assumptions allow any fixed proportion of the entries σ_{ij} to be zero, Assumption A1 requires the number of non-zero entries to be a constant proportion of the total number of entries. Indeed, otherwise by the Weyl comparison inequality (cf. e.g. [41, Theorem 3.3.13]), the empirical spectral distributions μ_n^Y converge weakly in probability to δ_0 , the point mass at the origin. To obtain a nontrivial limit would require a rescaling of the matrices Y_n , which amounts to rescaling A_n to have entries of growing size.

A1 is required both to bound the smallest singular value of the shifted random matrices and to prove effective bounds on the Stieltjes transform. We expect that

our results should extend to certain matrices with density $\sim n^{\varepsilon - 1}$ for arbitrary fixed $\varepsilon \in (0, 1)$, suitably rescaled. An interesting first case to consider is random band matrices with shrinking bandwidth. The limit of the empirical distribution of the singular values was recently computed in [43], but bounds on the smallest singular value were not considered.

Extension to allow heavy-tailed entries In a similar direction, it would be interesting to prove an analogue of Theorem 2.3 for the case that the entries X_{ij} lie in the basin of attraction of an α -stable law for some $\alpha \in (0, 2)$. In this case we expect the deterministic equivalents μ_n will not have compact support. The limiting empirical distribution of singular values for such matrices (allowed to be rectangular with bounded eccentricity) was studied by Belinschi, Dembo and Guionnet in [18]. For the case that the entries are i.i.d. the limiting empirical spectral distribution was established by Bordenave, Caputo and Chafaï in [20].

3 Asymptotics of singular values distributions

Recall that in Section 2.5 we introduced the Hermitian matrix \boldsymbol{Y}_{n}^{z} in (2.18) whose spectral measure is $\check{L}_{n,z}$ in (2.17). We also introduced the resolvent of \boldsymbol{Y}_{n}^{z} , $\boldsymbol{R}(z,\eta)$ in (2.19) and labeled its blocks in (2.25). By the well-known formula for the inverse of a partitioned matrix [42, §0.7.3],

$$G(z,\eta) = \eta \left((Y-z)(Y-z)^* - \eta^2 \right)^{-1},$$

$$\widetilde{G}(z,\eta) = \eta \left((Y-z)^*(Y-z) - \eta^2 \right)^{-1},$$

$$F(z,\eta) = (Y-z) \left((Y-z)^*(Y-z) - \eta^2 \right)^{-1},$$

$$F'(z,\eta) = \left((Y-z)^*(Y-z) - \eta^2 \right)^{-1} (Y-z)^*.$$
(3.1)

The main objective of this section is to provide deterministic counterparts of the normalized traces of these matrix functions. Given the matrices G and \tilde{G} , it will be convenient to introduce the vectors

$$\boldsymbol{g} = (G_{11} \cdots G_{nn})^{\mathsf{T}}, \quad \tilde{\boldsymbol{g}} = (\widetilde{G}_{11} \cdots \widetilde{G}_{nn})^{\mathsf{T}} \quad \text{and} \quad \boldsymbol{\vec{g}}^{\mathsf{T}} = (\boldsymbol{g}^{\mathsf{T}} \, \boldsymbol{\tilde{g}}^{\mathsf{T}}).$$
(3.2)

We begin by deriving the Schwinger-Dyson equations, a system of equations approximately satisfied by the diagonal entries of the matrices in (3.1). We then show the Schwinger-Dyson equations have a unique solution corresponding to Stieltjes transforms of probability measures and analyze the properties of the solution. Finally we estimate the difference between (3.1) and the true solution of the Schwinger-Dyson equations, which in turn is used to estimate the difference between the empirical spectral measure of \boldsymbol{Y}_n^z and its deterministic counterpart.

Notation 3.1. Let $\alpha_n = \alpha_n(z, \eta)$ and $\beta_n = \beta_n(z, \eta)$ be complex sequences such that there exist some constant C > 0 and some integers c_0, c_1 all independent from η and n but which may depend on z such that

$$|\alpha_n| \le \frac{C|\eta|^{c_1}}{\operatorname{Im}^{c_0}(\eta) \wedge 1} |\beta_n| .$$

We denote this by $\alpha_n = \mathcal{O}_\eta(\beta_n)$. If $\alpha_n = \alpha_n^i$ and $\beta_n = \beta_n^i$ depend on some extra parameter $i \in \mathcal{I}$, then the notation $\mathcal{O}_\eta()$ in $\alpha_n^i = \mathcal{O}_\eta(\beta_n^i)$ must be understood uniform in *i*. If α_n and β_n are vectors or matrices, the notation $\alpha_n = \vec{\mathcal{O}}_\eta(\beta_n)$ corresponds to a uniform entrywise relation.

3.1 Derivation of the Schwinger-Dyson equations

In this subsection we specialize to the case that the entries of X are i.i.d. standard complex Gaussian variables. Later we will compare a general matrix with a Gaussian matrix, at which point we will label the Gaussian matrix and associated quantities with a superscript \mathcal{N} ; however, we omit the superscript in the present subsection.

For a resolvent R as defined in (2.19) with complex entries, the following differentiation formulas hold true and will be needed in the sequel:

$$\frac{\partial \boldsymbol{R}_{ij}}{\partial Y_{k\ell}} = -\boldsymbol{R}_{ik}\boldsymbol{R}_{n+\ell,j} , \quad \frac{\partial \boldsymbol{R}_{ij}}{\partial \overline{Y}_{\ell k}} = -\boldsymbol{R}_{i,n+k}\boldsymbol{R}_{\ell j} , \quad 1 \le k, \ell \le n , \quad 1 \le i,j \le 2n .$$
(3.3)

We will heavily rely on the variance estimates provided in Proposition A.1 and Corollary A.2. Denote by $\boldsymbol{Y} = \begin{bmatrix} 0 & Y \\ Y^* & 0 \end{bmatrix}$. The equation $\boldsymbol{R}^{-1}\boldsymbol{R} = I_{2n}$ yields

$$-\eta \mathbf{R} + \mathbf{Y}\mathbf{R} + \begin{bmatrix} -zF' & -z\widetilde{G} \\ -z^*G & -z^*F \end{bmatrix} = I_{2n} .$$
(3.4)

Taking $i \in [n]$ yields

$$-\eta \mathbb{E}G_{ii} + \mathbb{E}(\boldsymbol{Y}\boldsymbol{R})_{ii} - z\mathbb{E}F'_{ii} = 1.$$
(3.5)

Applying the integration by part formula for complex Gaussian random variables (see for instance [51, (2.1.40)]) together with (3.3) yields

$$\mathbb{E}(\boldsymbol{Y}\boldsymbol{R})_{ii} = \sum_{\ell=1}^{n} \mathbb{E}Y_{i\ell}\boldsymbol{R}_{n+\ell,i} = \sum_{\ell=1}^{n} \frac{\sigma_{i\ell}^{2}}{n} \mathbb{E}\left[\frac{\partial \boldsymbol{R}_{n+\ell,i}}{\partial \overline{Y_{i\ell}}}\right] = -\sum_{\ell=1}^{n} \frac{\sigma_{i\ell}^{2}}{n} \mathbb{E}(\boldsymbol{R}_{n+\ell,n+\ell}\boldsymbol{R}_{ii}) .$$

Plugging this into (3.5) yields

$$-\mathbb{E}\left(\eta + [V_n \tilde{\boldsymbol{g}}]_i\right)G_{ii} - z\mathbb{E}F'_{ii} = 1.$$
(3.6)

Specializing again Equation (3.4) for $i \in [n]$ yields, with similar arguments,

$$-\eta \mathbb{E}F_{ii} - \mathbb{E}[V_n \tilde{\boldsymbol{g}}]_i F_{ii} - z \mathbb{E}\widetilde{G}_{ii} = 0.$$
(3.7)

Arguing similarly with the help of the following integration by parts formula, valid for i > n,

$$\mathbb{E}(\boldsymbol{Y}\boldsymbol{R})_{ij} = \sum_{\ell=1}^{n} \mathbb{E}(\overline{Y}_{\ell,i-n}\boldsymbol{R}_{\ell j}) = \sum_{\ell=1}^{n} \frac{\sigma_{\ell,i-n}^{2}}{n} \mathbb{E}\left[\frac{\partial \boldsymbol{R}_{\ell j}}{\partial Y_{\ell,i-n}}\right] = -\sum_{\ell=1}^{n} \frac{\sigma_{\ell,i-n}^{2}}{n} \mathbb{E}(\boldsymbol{R}_{\ell \ell}\boldsymbol{R}_{ij})$$

yields the following equations

$$-\mathbb{E}(\eta + [V_n^{\mathsf{T}}\boldsymbol{g}]_i)\widetilde{G}_{ii} - z^*\mathbb{E}F_{ii} = 1, \qquad (3.8)$$

$$-\mathbb{E}(\eta + [V_n^{\dagger} g]_i)F'_{ii} - z^* \mathbb{E}G_{ii} = 0.$$
(3.9)

Notice that equations (3.6)–(3.9) can be compactly written

$$\mathbb{E}\begin{bmatrix} -\operatorname{diag}(V_n \tilde{\boldsymbol{g}}) - \eta & -z \\ -z^* & -\operatorname{diag}(V_n^{\mathsf{T}} \boldsymbol{g}) - \eta \end{bmatrix} \boldsymbol{R} = I_{2n} .$$
(3.10)

Using Cauchy-Schwarz inequality and the estimates in Proposition A.1, we get

$$\mathbb{E}[V_{n}\tilde{\boldsymbol{g}}]_{i}F_{ii} - \mathbb{E}[V_{n}\tilde{\boldsymbol{g}}]_{i}\mathbb{E}F_{ii} = \mathcal{O}_{\eta}\left(\frac{1}{n^{3/2}}\right) \quad \text{thus} \quad -(\eta + \mathbb{E}[V_{n}\tilde{\boldsymbol{g}}]_{i})\mathbb{E}F_{ii} = z\mathbb{E}\widetilde{G}_{ii} + \mathcal{O}_{\eta}\left(\frac{1}{n^{3/2}}\right)$$

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by (3.7). In particular,

$$-\mathbb{E}F_{ii} = z \frac{\mathbb{E}G_{ii}}{\eta + \mathbb{E}[V_n \tilde{g}]_i} + \mathcal{O}_\eta \left(\frac{1}{n^{3/2}}\right)$$

since $|\eta + \mathbb{E}[V_n \tilde{g}]_i|^{-1} \leq \text{Im}^{-1}(\eta)$. On the other hand, using the same decorrelation argument in equation (3.8), we obtain

$$-\mathbb{E}(\eta + [V_n^{\mathsf{T}}\boldsymbol{g}]_i)\mathbb{E}\widetilde{G}_{ii} - z^*\mathbb{E}F_{ii} = 1 + \mathcal{O}_\eta\left(\frac{1}{n^{3/2}}\right) \ .$$

Combining these two equations, we finally get

$$\mathbb{E}\widetilde{G}_{ii}\left\{-(\eta + \mathbb{E}[V_n^{\mathsf{T}}\boldsymbol{g}]_i) + \frac{|z|^2}{\eta + \mathbb{E}[V_n\tilde{\boldsymbol{g}}]_i}\right\} = 1 + \mathcal{O}_\eta\left(\frac{1}{n^{3/2}}\right) \ .$$

Using the property (2.2) twice, one has

$$\left| -(\eta + \mathbb{E}[V_n^{\mathsf{T}}\boldsymbol{g}]_i) + \frac{|z|^2}{\eta + \mathbb{E}[V_n \tilde{\boldsymbol{g}}]_i} \right|^{-1} \leq \frac{1}{\mathrm{Im}(\eta)}$$

Hence

$$\mathbb{E}\widetilde{G}_{ii} = \frac{1}{-(\eta + \mathbb{E}[V_n^{\mathsf{T}}\boldsymbol{g}]_i) + \frac{|\boldsymbol{z}|^2}{\eta + \mathbb{E}[V_n \boldsymbol{\tilde{g}}]_i}} + \mathcal{O}_{\eta}\left(\frac{1}{n^{3/2}}\right) .$$
(3.11)

Combining similarly equations (3.6) and (3.9) and decorrelating when needed with the help of Proposition A.1, we obtain the companion equation:

$$\mathbb{E}G_{ii} = \frac{1}{-(\eta + \mathbb{E}[V_n \tilde{\boldsymbol{g}}]_i) + \frac{|z|^2}{\eta + \mathbb{E}[V_n^{\mathsf{T}} \boldsymbol{g}]_i}} + \mathcal{O}_\eta\left(\frac{1}{n^{3/2}}\right) .$$
(3.12)

We now introduce an unperturbed version of equations (3.11) and (3.12).

3.2 Schwinger-Dyson equations

In this section we introduce the Schwinger–Dyson equations (2.20). Notice that these equations already appear in [4], from which we will deduce properties for their solutions. To write this more compactly, we introduce the notation $\vec{b} = \begin{pmatrix} b \\ \tilde{b} \end{pmatrix}$ for any two $n \times 1$ vectors b and \tilde{b} with complex components and the following definitions:

$$\Upsilon(\vec{\boldsymbol{b}},\eta) := \operatorname{diag}\left(\frac{1}{|z|^2 - ([V_n \tilde{\boldsymbol{b}}]_i + \eta)([V_n^{\mathsf{T}} \boldsymbol{b}]_i + \eta)}; i \in [n]\right) := \operatorname{diag}(\Upsilon_i(\vec{\boldsymbol{b}},\eta); i \in [n]),$$
(3.13)

and

$$\mathcal{J}(\vec{\boldsymbol{b}},\eta) := \begin{pmatrix} \boldsymbol{\Upsilon}(\vec{\boldsymbol{b}},\eta)V_n^{\mathsf{T}} & 0\\ 0 & \boldsymbol{\Upsilon}(\vec{\boldsymbol{b}},\eta)V_n \end{pmatrix} \vec{\boldsymbol{b}} + \eta \begin{pmatrix} \boldsymbol{\Upsilon}(\vec{\boldsymbol{b}},\eta)\mathbf{1}\\ \boldsymbol{\Upsilon}(\vec{\boldsymbol{b}},\eta)\mathbf{1} \end{pmatrix} .$$
(3.14)

Then (2.20) can be compactly written as

$$\vec{\boldsymbol{p}} = \mathcal{J}(\vec{\boldsymbol{p}}, \eta). \tag{3.15}$$

Both Υ and \mathcal{J} depend on z as well (and to be even more precise, on |z|). We will not indicate this dependence in the sequel.

We now collect properties of solutions to (2.20).

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Proposition 3.1 (Schwinger–Dyson equations). For all fixed $\eta \in \mathbb{C}_+$ and $z \in \mathbb{C}$, let $p = (p_i)$ and $\tilde{p} = (\tilde{p}_i)$ be two $n \times 1$ vectors which solve (2.20).

- 1. The system (2.20) admits a unique solution \vec{p} satisfying $\text{Im}\,\vec{p} \succ 0$.
- 2. For any initial vector \vec{p}_0 with $\operatorname{Im} \vec{p}_0 \geq 0$, the iterations $\vec{p}_{k+1} = \mathcal{J}(\vec{p}_k)$ converge to this solution \vec{p} as $k \to \infty$.
- 3. For all $z \in \mathbb{C}$ and $i \in [n]$, the functions $\eta \mapsto p_i(\eta)$ and $\eta \mapsto \tilde{p}_i(\eta)$ are Stieltjes transforms of symmetric probability measures on \mathbb{R} respectively denoted by μ_i and $\tilde{\mu}_i$. In particular, let $r_i(t) := \operatorname{Im} p_i(it)$ and $\tilde{r}_i(t) := \operatorname{Im} \tilde{p}_i(it)$ for t > 0 then

$$p_i(\mathrm{i}t) = \mathrm{i}r_i(t)$$
 and $\tilde{p}_i(\mathrm{i}t) = \mathrm{i}\tilde{r}_i(\mathrm{i}t)$

Otherwise stated, p_i and \tilde{p}_i are purely imaginary complex numbers along the imaginary axis.

4. Moreover, $\sum_{i=1}^{n} p_i = \sum_{i=1}^{n} \tilde{p}_i$ and the common value

$$\eta \longmapsto \frac{1}{n} \sum_{i=1}^{n} p_i(\eta) = \frac{1}{n} \sum_{i=1}^{n} \tilde{p}_i(\eta)$$

is the Stieltjes transform if a symmetric probability measure $\check{\nu}_{n,z}$. We denote this Stieltjes transform by $\eta \mapsto g_{\check{\nu}_{n,z}}(\eta)$.

5. The sequences of probability measures $(\mu_i; i \leq n; n \geq 1)$, $(\tilde{\mu}_i; i \leq n; n \geq 1)$ and $(\check{\nu}_{n,z}; n \geq 1)$ are tight. In particular,

$$\sup_{n\geq 1}\int |x|^2\check{\nu}_{n,z}(\,dx)\,<\,\infty\,.$$

We henceforth refer to the solution $\vec{p} = \vec{p}(\eta)$, $\operatorname{Im} \vec{p} \succ 0$ as the solution to the Schwinger-Dyson equations.

Proof. Proofs of parts (1) and (2) are a direct application of Earle-Hamilton's theorem[38], which was first used (to the authors' knowledge) in random matrix theory by [40].We now address part (3) and (5).

The fact that the p_i 's and \tilde{p}_i 's are Stieltjes transforms of probability measures immediately follows from [4, Proposition 2.1]. This result also yields the desired tightness properties.

In order to complete the proof of part (3), we now prove that the probability measures associated to the p_i 's are symmetric. To this end, simply observe that, given $\eta \in \mathbb{C}_+$, if $\vec{p} = (p, \tilde{p})$ is the solution with $\operatorname{Im} \vec{p} \succ 0$ of the Schwinger–Dyson equations (2.20), then $-\vec{p}$ is the unique solution with $\operatorname{Im}(-\vec{p}) \prec 0$ of the analogous system obtained by replacing η with $-\eta$. The result follows from the application of Lemma A.1.

It remains to prove that along the imaginary axis (it, t > 0), p_i and \tilde{p}_i are purely imaginary complex numbers.

$$p_{i}(it) = \int_{(-\infty,0)} \frac{\mu_{i}(d\lambda)}{\lambda - it} + \int_{(0,\infty)} \frac{\mu_{i}(d\lambda)}{\lambda - it} - \frac{\mu_{i}(\{0\})}{it},$$

$$= \int_{(0,\infty)} \left(\frac{1}{\lambda - it} - \frac{1}{\lambda + it}\right) \mu_{i}(d\lambda) + \frac{i\mu_{i}(\{0\})}{t},$$

$$= 2i \int_{(0,\infty)} \frac{t}{\lambda^{2} + t^{2}} \mu_{i}(d\lambda) + \frac{i\mu_{i}(\{0\})}{t} =: ir_{i}(t).$$
(3.16)

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Similarly, $\tilde{p}(it) = i\tilde{r}(t)$ with $\tilde{r}(t) = (\tilde{r}_i(t))$ and $\tilde{r}_i(t) > 0$. Notice for future use that

$$\lim_{t \to \infty} tr_i(t) = 1 \quad \text{and} \quad \lim_{t \to \infty} t\tilde{r}_i(t) = 1 .$$
(3.17)

Parts (3) and (5) of the theorem are established.

We now prove part (4), that is $\sum_i p_i = \sum_i \tilde{p}_i$. Getting back to the system (2.20), we have

$$\sum_{i=1}^{n} p_i([V\tilde{\boldsymbol{p}}]_i + \eta) = \sum_{i=1}^{n} \frac{([V\tilde{\boldsymbol{p}}]_i + \eta)([V_n^{\mathsf{T}}\boldsymbol{p}]_i + \eta)}{-([V\tilde{\boldsymbol{p}}]_i + \eta)([V_n^{\mathsf{T}}\boldsymbol{p}]_i + \eta) + |z|^2} = \sum_{i=1}^{n} \tilde{p}_i([V_n^{\mathsf{T}}\boldsymbol{p}]_i + \eta).$$

But

$$\sum_{i=1}^{n} p_i [V \tilde{\boldsymbol{p}}]_i = \sum_{i,\ell=1}^{n} p_i \sigma_{i,\ell}^2 \tilde{p}_i = \sum_{i=1}^{n} \tilde{p}_i [V_n^{\mathsf{T}} \boldsymbol{p}]_i.$$

Since $\eta \neq 0$, we get the desired result.

3.3 Asymptotics of the spectral measure $\check{L}_{n,z}$ and the Hermitian resolvent

Theorem 3.2. Assume **A0** and **A1** hold, and let $\check{\nu}_{n,z}$ be defined as in Proposition 3.1-(4). Then for all $z \in \mathbb{C}$, $(\check{\nu}_{n,z})_n$ is tight, and

$$\check{L}_{n,z} \sim \check{\nu}_{n,z}$$

almost surely. Moreover, for any $\varepsilon > 0$, $x \mapsto \log |x|$ is $\check{\nu}_{n,z}$ -integrable on the set $\{|x| \ge \varepsilon\}$ and

$$\int_{\{|x|\geq\varepsilon\}} \log|x|\,\check{L}_{n,z}(dx) - \int_{\{|x|\geq\varepsilon\}} \log|x|\,\check{\nu}_{n,z}(dx) \xrightarrow[n\to\infty]{a.s.} 0$$

We will sometimes refer to $\check{\nu}_{n,z}$ as the deterministic equivalent of $\check{L}_{n,z}$.

The proof of Theorem 3.2 is postponed to Section 3.5. Notice that the first part $(\check{L}_{n,z} \sim \check{\nu}_{n,z})$ is a variation of classical results, see for example [36]. It will be a direct consequence of the forthcoming theorem on the asymptotics of Hermitian resolvent.

In order to get some insight on the asymptotics of the spectral measure μ_n^Y , we need more than the asymptotics of $\check{L}_{n,z}$. We rewrite hereafter the Schwinger–Dyson equations of Proposition 3.1 in a more suitable way for the forthcoming analysis. In what follows, the dependence in |z| is implicit and will be recalled if necessary.

We now introduce the deterministic equivalents to F and G, defined in (3.1). Let $\vec{p} = (p, \tilde{p})$ be the solution of the Schwinger–Dyson equations (2.20). Define the $n \times n$ diagonal matrices P, \tilde{P}, Θ and $\tilde{\Theta}$ by

$$\begin{aligned} P &:= \operatorname{diag}(\boldsymbol{p}) ,\\ \widetilde{P} &:= \operatorname{diag}(\widetilde{\boldsymbol{p}}) ,\\ \Theta &:= \operatorname{diag}\left((V_n \, \widetilde{\boldsymbol{p}})_i, \ i \in [n]\right) ,\\ \widetilde{\Theta} &:= \operatorname{diag}\left((V_n^\mathsf{T} \boldsymbol{p})_i, \ i \in [n]\right) . \end{aligned}$$

After easy massaging, the Schwinger–Dyson equations $\vec{p} = \mathcal{J}(\vec{p}, \eta)$ are equivalent to:

$$P = \left(-(\Theta + \eta) + |z|^2(\widetilde{\Theta} + \eta)^{-1}\right)^{-1} \quad \text{and} \qquad \widetilde{P} = \left(-(\widetilde{\Theta} + \eta) + |z|^2(\Theta + \eta)^{-1}\right)^{-1}.$$

Consider $2n \times 2n$ matrix \boldsymbol{S} defined as

$$\boldsymbol{S} := - \begin{pmatrix} \Theta(|z|,\eta) + \eta & z \\ z^* & \widetilde{\Theta}(|z|,\eta) + \eta \end{pmatrix}^{-1}.$$
(3.18)

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This definition is similar to equation (3.10) satisfied by the entries of the resolvent R. By the formula for the inverse of a partitioned matrix [42, §0.7.3], it holds that

$$oldsymbol{S} = egin{pmatrix} P(|z|,\eta) & B(z,\eta) \ B'(z,\eta) & \widetilde{P}(|z|,\eta) \end{pmatrix},$$

where

$$B(z,\eta) = -z \Big(\Theta(|z|,\eta) + \eta\Big)^{-1} \widetilde{P}(|z|,\eta) = -zP(|z|,\eta) \left(\widetilde{\Theta}(|z|,\eta) + \eta\Big)^{-1} , \qquad (3.19)$$

and $B'(z,\eta)$ can be made explicit in a similar fashion, but will not be used.

Theorem 3.3. Assume **A0** and **A1** hold. Then almost surely, for every $z \in \mathbb{C}$ and $\eta \in \mathbb{C}_+$,

$$\frac{1}{n} \begin{pmatrix} \operatorname{tr} G(z,\eta) & \operatorname{tr} F(z,\eta) \\ \operatorname{tr} F'(z,\eta) & \operatorname{tr} \widetilde{G}(z,\eta) \end{pmatrix} - \frac{1}{n} \begin{pmatrix} \operatorname{tr} P(|z|,\eta) & \operatorname{tr} B(z,\eta) \\ \operatorname{tr} B'(z,\eta) & \operatorname{tr} \widetilde{P}(|z|,\eta) \end{pmatrix} \xrightarrow[n \to \infty]{} 0 \ .$$

Moreover, there exist $\alpha, \beta > 0$ such that for $t \in (n^{-\alpha}, n^{\beta})$ and n large enough, with $\eta = it$ we have

$$\frac{1}{n} \begin{pmatrix} \operatorname{tr} \mathbb{E} G(z,\eta) & \operatorname{tr} \mathbb{E} F(z,\eta) \\ \operatorname{tr} \mathbb{E} F'(z,\eta) & \operatorname{tr} \mathbb{E} \widetilde{G}(z,\eta) \end{pmatrix} - \frac{1}{n} \begin{pmatrix} \operatorname{tr} P(|z|,\eta) & \operatorname{tr} B(z,\eta) \\ \operatorname{tr} B'(z,\eta) & \operatorname{tr} \widetilde{P}(|z|,\eta) \end{pmatrix} = \vec{\mathcal{O}}_{\eta} \left(n^{-1/2} \right) \,.$$

The rate provided along the imaginary axis is not optimal, but sufficient for our purposes.

The proof of Theorem 3.3 immediately follows from Propositions 3.4, 3.5 and 3.6 stated hereafter.

Proposition 3.4. Assume **A0** holds and let $z \in \mathbb{C}$ and $\eta \in \mathbb{C}_+$. Then almost surely,

$$\frac{1}{n} \begin{pmatrix} \operatorname{tr} G(z,\eta) & \operatorname{tr} F(z,\eta) \\ \operatorname{tr} F'(z,\eta) & \operatorname{tr} \widetilde{G}(z,\eta) \end{pmatrix} - \frac{1}{n} \begin{pmatrix} \operatorname{tr} \mathbb{E}G(z,\eta) & \operatorname{tr} \mathbb{E}F(z,\eta) \\ \operatorname{tr} \mathbb{E}F'(z,\eta) & \operatorname{tr} \mathbb{E}\widetilde{G}(z,\eta) \end{pmatrix} \xrightarrow[n \to \infty]{} 0 \ .$$

Proof. This is a direct application of [19, Lemma 4.21].

To manage the expectation terms n^{-1} tr $\mathbb{E}(\cdot)$, we introduce the Gaussian counterparts of the quantities of interest. Consider a family of i.i.d. standard complex random variables $(X_{ij}^{\mathcal{N}}; 1 \leq i, j \leq n)$, where $X_{ij}^{\mathcal{N}} = (U + iU')/\sqrt{2}$, with U, U' being independent real $\mathcal{N}(0,1)$ random variables. Notice in particular that

$$\mathbb{E}X_{ij}^{\mathcal{N}} = 0$$
, $\mathbb{E}\left(X_{ij}^{\mathcal{N}}\right)^2 = 0$ and $\mathbb{E}|X_{ij}^{\mathcal{N}}|^2 = 1$.

Similarly, let $Y_{ij}^{\mathcal{N}} = \frac{\sigma_{ij}}{\sqrt{n}} X_{ij}^{\mathcal{N}}$, and let $\mathbf{R}^{\mathcal{N}}$, $G^{\mathcal{N}}$, $\tilde{G}^{\mathcal{N}}$, $F^{\mathcal{N}}$, and $F'^{\mathcal{N}}$ the matrix functions associated with the matrix $Y^{\mathcal{N}} = (Y_{ij}^{\mathcal{N}})$ as in (2.19),(3.1). Then we have the following proposition.

Proposition 3.5. Assume **A0** and **A1** hold. Let $z \in \mathbb{C}$ and $\eta \in \mathbb{C}_+$. Then

$$\frac{1}{n} \begin{pmatrix} \operatorname{tr} \mathbb{E}G(z,\eta) & \operatorname{tr} \mathbb{E}F(z,\eta) \\ \operatorname{tr} \mathbb{E}F'(z,\eta) & \operatorname{tr} \mathbb{E}\widetilde{G}(z,\eta) \end{pmatrix} - \frac{1}{n} \begin{pmatrix} \operatorname{tr} \mathbb{E}G^{\mathcal{N}}(z,\eta) & \operatorname{tr} \mathbb{E}F^{\mathcal{N}}(z,\eta) \\ \operatorname{tr} \mathbb{E}F'^{\mathcal{N}}(z,\eta) & \operatorname{tr} \mathbb{E}\widetilde{G}^{\mathcal{N}}(z,\eta) \end{pmatrix} = \vec{\mathcal{O}}_{\eta} \begin{pmatrix} n^{-1/2} \end{pmatrix} \,.$$

The proof of Proposition 3.5 relies on fairly standard arguments and is thus postponed to Appendix A.3.

Proposition 3.6. Assume A1 holds, and let $z \in \mathbb{C}$ and $\eta \in \mathbb{C}_+$. Then

$$\frac{1}{n} \begin{pmatrix} \operatorname{tr} \mathbb{E} G^{\mathcal{N}}(z,\eta) & \operatorname{tr} \mathbb{E} F^{\mathcal{N}}(z,\eta) \\ \operatorname{tr} \mathbb{E} F^{'\mathcal{N}}(z,\eta) & \operatorname{tr} \mathbb{E} \widetilde{G}^{\mathcal{N}}(z,\eta) \end{pmatrix} - \frac{1}{n} \begin{pmatrix} \operatorname{tr} P(|z|,\eta) & \operatorname{tr} B(z,\eta) \\ \operatorname{tr} B'(z,\eta) & \operatorname{tr} \widetilde{P}(|z|,\eta) \end{pmatrix} \xrightarrow[n \to \infty]{} 0 \,.$$

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Moreover, there exist $\alpha, \beta > 0$ such that for $t \in (n^{-\alpha}, n^{\beta})$ and n large enough, with $\eta = \mathrm{i}t$ we have

$$\frac{1}{n} \begin{pmatrix} \operatorname{tr} \mathbb{E}G^{\mathcal{N}}(z,\eta) & \operatorname{tr} \mathbb{E}F^{\mathcal{N}}(z,\eta) \\ \operatorname{tr} \mathbb{E}F^{'\mathcal{N}}(z,\eta) & \operatorname{tr} \mathbb{E}\widetilde{G}^{\mathcal{N}}(z,\eta) \end{pmatrix} - \frac{1}{n} \begin{pmatrix} \operatorname{tr} P(|z|,\eta) & \operatorname{tr} B(z,\eta) \\ \operatorname{tr} B'(z,\eta) & \operatorname{tr} \widetilde{P}(|z|,\eta) \end{pmatrix} = \vec{\mathcal{O}}_{\eta} \begin{pmatrix} n^{-3/2} \end{pmatrix}$$

The proof of Proposition 3.6 follows hereafter in Section 3.4.

From Theorem 3.3, we will deduce the asymptotic behavior of the empirical distribution $L_{n,z}$ of the singular values of $Y_n - z$ by analyzing the convergence of $n^{-1} \operatorname{tr} G(z, \eta)$. Moreover, for any $\varepsilon > 0$, the asymptotic behavior of $\int_{\{|x| \ge \varepsilon\}} \log |x| \check{L}_{n,z}(dx)$ will be identified.

3.4 Proof of Proposition 3.6

We first prove the convergence to zero. Recall that $V = \frac{A}{\sqrt{n}} \odot \frac{A}{\sqrt{n}}$, that \vec{g} is introduced in (3.2), that \vec{p} is the solution of the Schwinger–Dyson equations and define

$$ec{arepsilon} = \mathbb{E}ec{ec{g}} - ec{p}$$

Recall the definition of Υ in (3.13). For given $n \times 1$ vectors \boldsymbol{b} and $\tilde{\boldsymbol{b}}$, define

$$\begin{cases} \Delta(\boldsymbol{b}) = \operatorname{diag}(\eta + (V_n^{\mathsf{T}}\boldsymbol{b})_i) := \operatorname{diag}(\Delta_i(\boldsymbol{b})) \\ \widetilde{\Delta}(\widetilde{\boldsymbol{b}}) = \operatorname{diag}(\eta + (V_n\widetilde{\boldsymbol{b}})_i) := \operatorname{diag}(\widetilde{\Delta}_i(\widetilde{\boldsymbol{b}})) \end{cases} .$$
(3.20)

Then

$$\mathbb{E}G_{ii} - p_i = \frac{1}{-(\eta + [V\mathbb{E}\tilde{\boldsymbol{g}}]_i) + \frac{|\boldsymbol{z}|^2}{\eta + [V^{\mathsf{T}}\mathbb{E}\boldsymbol{g}]_i}} - \frac{1}{-(\eta + [V\tilde{\boldsymbol{p}}]_i) + \frac{|\boldsymbol{z}|^2}{\eta + [V^{\mathsf{T}}\boldsymbol{p}]_i}} + \mathcal{O}_{\eta}\left(n^{-3/2}\right)$$
$$= \Upsilon_i(\mathbb{E}\boldsymbol{\vec{g}})\Upsilon_i(\boldsymbol{\vec{p}})\left\{|\boldsymbol{z}|^2[V^{\mathsf{T}}(\mathbb{E}\boldsymbol{g} - \boldsymbol{p})]_i + \Delta_i(\mathbb{E}\boldsymbol{g})\Delta_i(\boldsymbol{p})[V(\mathbb{E}\tilde{\boldsymbol{g}} - \tilde{\boldsymbol{p}})]_i\right\} + \mathcal{O}_{\eta}\left(n^{-3/2}\right),$$
(3.21)

and a similar expression holds for $\mathbb{E}\widetilde{G}_{ii} - \widetilde{p}_i$. Taking into account the straightforward estimates

$$|\Delta_i(\mathbb{E}oldsymbol{g})|, \; |\Delta_i(oldsymbol{p})| \leq |\eta| + rac{\sigma_{\max}^2}{\mathrm{Im}(\eta)} \; \; \; ext{and} \; \; \; |\Upsilon_i(\mathbb{E}oldsymbol{g})|, \; |\Upsilon_i(oldsymbol{p})| \leq rac{1}{\mathrm{Im}^2(\eta)} \, ,$$

we end up with

$$\|\vec{\boldsymbol{\varepsilon}}\|_{\infty} \leq K \left(\frac{|\eta|^2 + |z|^2}{\operatorname{Im}^4(\eta)} + \frac{1}{\operatorname{Im}^6(\eta)} \right) \|\vec{\boldsymbol{\varepsilon}}\|_{\infty} + \vec{\mathcal{O}}_{\eta} \left(n^{-3/2} \right) ,$$

where K is an absolute constant (depending on σ_{\max}). Letting $\eta \in \mathbb{C}_+$ be chosen in such a way that $K\left(\frac{|\eta|^2+|z|^2}{\operatorname{Im}^4(\eta)}+\frac{1}{\operatorname{Im}^6(\eta)}\right) < 1$, one has $\|\vec{\varepsilon}'\|_{\infty}(\eta) \xrightarrow[n \to \infty]{} 0$. On the other hand, $\mathbb{E}G_{ii} - p_i$ is analytic and uniformly bounded over the compact subsets of \mathbb{C}_+ . Therefore, every converging subsequence converges toward an analytic function which coincides with the zero function on a sufficiently large open subset of \mathbb{C}_+ , and hence is equal to the zero function over \mathbb{C}_+ . The same applies to $\mathbb{E}\tilde{G}_{ii} - \tilde{p}_i$. This proves that $\|\vec{\varepsilon}'\|_{\infty} \xrightarrow[n \to \infty]{} 0$ for every $\eta \in \mathbb{C}_+$.

We now focus on $\eta = it$. For a given $2n \times 1$ vector $\vec{b} = \begin{pmatrix} b \\ \tilde{b} \end{pmatrix}$, introduce the $2n \times 2n$ matrix

$$\mathcal{A}(\vec{\boldsymbol{b}}) = \frac{1}{\sqrt{n}} \begin{pmatrix} |z| \boldsymbol{\Upsilon}(\vec{\boldsymbol{b}}) A^{\mathsf{T}} & \boldsymbol{\Upsilon}(\vec{\boldsymbol{b}}) \Delta(\boldsymbol{b}) A \\ \boldsymbol{\Upsilon}(\vec{\boldsymbol{b}}) \widetilde{\Delta}(\tilde{\boldsymbol{b}}) A^{\mathsf{T}} & |z| \boldsymbol{\Upsilon}(\vec{\boldsymbol{b}}) A \end{pmatrix} .$$
(3.22)

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We can compactly express this equation as $\vec{\epsilon} = \mathcal{A}(\mathbb{E}\vec{g}) \odot \mathcal{A}(\vec{p}) \vec{\epsilon} + \vec{\mathcal{O}}_{\eta} (n^{-3/2})$ and easily prove that

$$\vec{\boldsymbol{\varepsilon}} = \mathcal{A}(\vec{\boldsymbol{p}}) \odot \mathcal{A}(\vec{\boldsymbol{p}}) \,\vec{\boldsymbol{\varepsilon}} + \vec{\mathcal{O}}_{\eta} \left(\|\vec{\boldsymbol{\varepsilon}}\|_{\infty}^2 + n^{-3/2} \right) \,. \tag{3.23}$$

In Appendix A.4 we prove that $I - \mathcal{A}(\vec{p}) \odot \mathcal{A}(\vec{p})$ is invertible and that

$$\left\| \left(I - \mathcal{A}(\vec{p}) \odot \mathcal{A}(\vec{p}) \right)^{-1} \right\|_{\infty} = \mathcal{O}_{\eta} (1) , \quad \eta = \mathrm{i}t .$$
(3.24)

Combining (3.23) and (3.24), we finally end up with $\|\vec{e}\|_{\infty} = O_{\eta} \left(\|\vec{e}\|_{\infty}^2 + n^{-3/2} \right)$, for $\eta = \mathrm{i}t$. Following Notation 3.1, we rewrite this estimate as

$$\|\vec{\varepsilon}\|_{\infty} \le \kappa(\mathrm{i}t) \left(\|\vec{\varepsilon}\|_{\infty}^2 + n^{-3/2}\right) \quad \text{for} \quad \kappa(\mathrm{i}t) = \frac{Ct^{c_1}}{t^{c_0} \wedge 1} \,. \tag{3.25}$$

The quadratic polynomial $P(X) = \kappa(it)X^2 - X + \kappa(it)n^{-3/2}$ admits two distinct real roots as long as $1 - 4\kappa^2(it)n^{-3/2} > 0$. Denote by $\Gamma_n(\alpha, \beta)$ the set

$$\Gamma_n(\alpha,\beta) = \left\{\eta = \mathrm{i}t,\, t > 0,\, t \in (n^{-\alpha},n^\beta)\right\}$$

If $\eta = it \in \Gamma_n(\alpha, \beta)$, then $4\kappa^2(it)n^{-3/2} \leq C^2 n^{2c_0\alpha + 2c_1\beta - 2/3}$. The r.h.s. goes to zero as long as $c_0\alpha + c_1\beta < 1/3$. This can be fulfilled for $\alpha, \beta > 0$ small enough, which is supposed to hold henceforth.

Then for $n \ge n_0(\alpha, \beta)$ large enough, $1 - 4\kappa^2(\mathrm{i}t)n^{-3/2} > 0$, thus P(X) admits two distinct real roots and (3.25) yields

$$\|\vec{\varepsilon}\|_{\infty} \leq X_{1}(\mathrm{i}t) := \frac{1 - \sqrt{1 - 4\kappa^{2}(\mathrm{i}t)n^{-3/2}}}{2\kappa(\mathrm{i}t)}$$

or $\|\vec{\varepsilon}\|_{\infty} \geq X_{2}(\mathrm{i}t) := \frac{1 + \sqrt{1 - 4\kappa^{2}(\mathrm{i}t)n^{-3/2}}}{2\kappa(\mathrm{i}t)}$. (3.26)

Since for $n \ge n_0(\alpha, \beta)$ and $\eta \in \Gamma_n(\alpha, \beta)$ the functions $\eta \mapsto \|\vec{\varepsilon}\|_{\infty}(\eta), X_1(\eta), X_2(\eta)$ are continuous, only one of the two conditions (3.26) can hold uniformly in $\Gamma_n(\alpha, \beta)$. For $\eta = i$ we have that $\kappa(i) = C$ and $X_2(i) = \frac{1+\sqrt{1-4C^2n^{2/3}}}{2C} = \mathcal{O}(1)$, but $\|\vec{\varepsilon}\|_{\infty} \xrightarrow[n \to \infty]{n \to \infty} 0$ by the first part of the proposition. Hence the condition $\|\vec{\varepsilon}\|_{\infty}(i) \ge X_2(i)$ cannot hold and necessarily

$$\|\vec{\boldsymbol{\varepsilon}}\|_{\infty} \leq X_1(\eta) = \frac{2\kappa(\eta)n^{-3/2}}{1 + \sqrt{1 - 4\kappa^2(\eta)n^{-3/2}}} \leq 2\kappa(\eta)n^{-3/2} \quad \forall \eta \in \Gamma_n(\alpha,\beta) + \frac{1}{2\kappa(\eta)n^{-3/2}} \leq 2\kappa(\eta)n^{-3/2} \quad \forall \eta \in \Gamma_n(\alpha,\beta) + \frac{1}{2\kappa(\eta)n^{-3/2}} \leq 2\kappa(\eta)n^{-3/2} < 2\kappa(\eta)n$$

We have proved so far that for $n \ge n_0(\alpha, \beta)$ large enough and $\eta \in \Gamma_n(\alpha, \beta)$, $\|\vec{e}\|_{\infty} = O_\eta(n^{-3/2})$.

We are now in position to conclude:

$$\left|\frac{1}{n}\sum_{i\in[n]}\mathbb{E}G_{ii} - \frac{1}{n}\sum_{i\in[n]}p_i\right| \leq \|\vec{\varepsilon}\|_{\infty} \leq \mathcal{O}_{\eta}\left(n^{-3/2}\right)$$

The same arguments apply verbatim for the term $\frac{1}{n}\sum_i(\mathbb{E}\,\widetilde{G}_{ii}-\widetilde{p}_i).$ Consider now the term

$$\frac{1}{n} \operatorname{tr} \mathbb{E}F - \frac{1}{n} \operatorname{tr} B \stackrel{(a)}{=} \frac{1}{n} \sum_{i=1}^{n} \left\{ \frac{z}{-([V \mathbb{E}\tilde{\boldsymbol{g}}]_{i} + \eta)} \mathbb{E}\tilde{G}_{ii} - \frac{z}{-([V \tilde{\boldsymbol{p}}]_{i} + \eta)} \tilde{p}_{i} \right\} + \mathcal{O}_{\eta} \left(n^{-3/2} \right) ,$$

where (a) follows from (3.7) and (3.19). One can now apply the same arguments as previously and handle similarly the term $\frac{1}{n} \operatorname{tr} \mathbb{E}F'(z,\eta) - \frac{1}{n} \operatorname{tr} B'(z,\eta)$. This completes the proof of Proposition 3.6.

3.5 Proof of Theorem 3.2

The convergence $\check{L}_{n,z} \sim \check{\nu}_{n,z}$ is a direct consequence of Theorem 3.3. Now, it is easy to prove with the help of the law of large numbers that a.s.

$$\limsup_n \int |x|^2 \check{L}_{n,z}(dx) < \infty .$$

This, together with Proposition 3.1-(5), yields

$$\int_{\{|x|\geq\varepsilon\}} \log |x|\,\check{L}_{n,z}(dx) - \int_{\{|x|\geq\varepsilon\}} \log |x|\,\check{\nu}_{n,z}(dx) \xrightarrow[n\to\infty]{a.s.} 0.$$

The proof is complete.

4 Proof of Theorem 2.2: analysis of the Master Equations

4.1 Regularized Master Equations

Our first step is to introduce the so-called *Regularized Master Equations*, which are obtained from the Schwinger–Dyson equations (2.20) by taking $\eta = it$ and substituting p(it) = ir(t).

Given two $n \times 1$ vectors \boldsymbol{a} and $\tilde{\boldsymbol{a}}$ with nonnegative components and fixed numbers s > 0 and $t \ge 0$, let $\boldsymbol{\vec{a}} = \begin{pmatrix} \boldsymbol{a} \\ \tilde{\boldsymbol{a}} \end{pmatrix}$ and define the following quantities

$$\Psi(\vec{a},t) := \operatorname{diag}\left(\frac{1}{s^2 + ((V_n \tilde{a})_i + t)((V_n^{\mathsf{T}} a)_i + t)}; i \in [n]\right) := \operatorname{diag}(\psi_i(\vec{a},t); i \in [n]),$$
(4.1)

and

$$\mathcal{I}(\vec{a},t) := \begin{pmatrix} \Psi(\vec{a},t)V_n^{\mathsf{T}} & 0\\ 0 & \Psi(\vec{a},t)V_n \end{pmatrix} \vec{a} + t \begin{pmatrix} \Psi(\vec{a},t)\mathbf{1}_n\\ \Psi(\vec{a},t)\mathbf{1}_n \end{pmatrix}$$

We also define $\Psi(\vec{a}) := \Psi(\vec{a}, 0)$.

The proof of Proposition 2.1 amounts to showing that the vector equation $\vec{r} = \mathcal{I}(\vec{r}, t)$ admits a unique solution $\vec{r} = \vec{r}(s, t) \succ 0$.

Remark 4.1. We shall also prove that for any initial vector $\vec{r}_0 \geq 0$, the iterations $\vec{r}_{k+1} = \mathcal{I}(\vec{r}_k, t)$ converge to \vec{r} as $k \to \infty$.

Proof of Proposition 2.1. This follows immediately from Proposition 3.1 by setting p(it) = ir(t) and $\tilde{p}(it) = i\tilde{r}(t)$.

4.2 **Proof of Theorem 2.2**

In the following propositions, we recall some known properties of nonnegative and irreducible matrices.

Proposition 4.1 ([55, Theorems 1.1 and 5.5]). Let A and B be two square matrices such that $0 \preccurlyeq A \preccurlyeq B$. Then $\rho(A) \le \rho(B)$. Moreover, if B is irreducible, then $\rho(A) = \rho(B)$ implies that A = B.

Proposition 4.2 ([55, Theorem 1.6]). Let $A \succeq 0$ be a square and irreducible matrix, and let $x \succeq 0$ be a vector satisfying $Ax \preccurlyeq x$. Then $x \succ 0$ and $\rho(A) \le 1$. Moreover, $\rho(A) = 1$ if and only if Ax = x.

The proof of the following lemma is deferred to Section 5 – see Proposition 5.2.

Lemma 4.1. Let *V* be a nonnegative and irreducible $n \times n$ matrix, and let $\vec{r}(s,t)$ be the solution of the regularized master equations (2.7). Let $[a,b] \subset (0,\infty)$ and $\varepsilon > 0$, then

$$\sup_{(s,t)\in[a,b]\times[0,\varepsilon]}\|\vec{r}(s,t)\|<\infty$$

In particular, $\vec{r}(s,t)$ admits an accumulation point for fixed s > 0 as $t \downarrow 0$.

Next we show that any accumulation point provided by the above lemma constitutes a solution to the Master Equations (1.5):

Lemma 4.2 (Existence of solutions to the Master Equations). Let V and $\vec{r}(s,t)$ be as in Lemma 4.1.

1. Let s > 0. If $\vec{r}_* = (r_*, \tilde{r}_*) \succeq 0$ is an accumulation point for $\vec{r}(s, t)$ as $t \downarrow 0$, then

$$\begin{pmatrix} ec{r}_* \\ 0 \end{pmatrix} = \begin{pmatrix} \Psi(ec{r}_*)V^{\intercal} & 0 \\ 0 & \Psi(ec{r}_*)V \\ \mathbf{1}^{\intercal} & -\mathbf{1}^{\intercal} \end{pmatrix} ec{r}_* \ ,$$

where we recall that $\Psi(\vec{r}_*) = \operatorname{diag}(\psi_i(\vec{r}_*))_{i=1}^n$ with $\psi_i(\vec{r}_*) = (s^2 + (V\widetilde{r}_*)_i(V^{\mathsf{T}}r_*)_i)^{-1}$.

2. If moreover $s^2 \in (0, \rho(V))$, then $\vec{r}_* \succcurlyeq_{\neq} 0$.

Proof. The proof of part (1) is straightforward.

We now prove part (2) of the lemma. Let (t_k) be a positive sequence converging to zero in such a way that $\lim_{k\to\infty} \vec{r}(s,t_k) = \vec{r}_*$. Since $\vec{r}(s,t_k)$ satisfies (2.7), we have in particular that $\Psi(\vec{r})(s,t_k)V\tilde{r} \prec \tilde{r}$ and $\Psi(\vec{r})(s,t_k)V^{\mathsf{T}}r \prec r$. From Proposition 4.2 it follows that that $\rho(\Psi(\vec{r})(s,t_k)V) = \rho(\Psi(\vec{r})(s,t_k)V^{\mathsf{T}}) < 1$, and by the continuity of the spectral radius that $\rho(\Psi(\vec{r}_*)V) \leq 1$. If $\vec{r}_* = 0$, then $\Psi(\vec{r}_*) = s^{-2}I$ and $\rho(\Psi(\vec{r}_*)V) = s^{-2}\rho(V) > 1$ since $s^2 \in (0, \rho(V))$, which yields a contradiction. Necessarily, $\vec{r}_* \succcurlyeq \neq 0$.

Theorem 2.2–(1) and (2) are now consequences of the following lemma.

Lemma 4.3 (Uniqueness of solutions to the Master Equations). Let V be a nonnegative and irreducible $n \times n$ matrix, and let $\vec{q} \succeq 0$ be a solution of the system (1.5), which exists by the previous lemma.

- 1. If $s^2 \ge \rho(V)$, then $\vec{q} = 0$.
- 2. If $s^2 \in (0, \rho(V))$, then \vec{q} is unique as a solution of (1.5) satisfying $\vec{q} \succcurlyeq_{\neq} 0$. This solution satisfies $\vec{q} \succ 0$.

Proof. We first prove Item (1). Observe that $\Psi(\vec{q})V^{\mathsf{T}}$ is a nonnegative irreducible matrix for all $\vec{q} \succeq 0$. Assume that $q \succeq 0$. Then $\rho(\Psi(\vec{q})V^{\mathsf{T}}) = 1$ and $q \succ 0$ by Proposition 4.2. From the equation $\mathbf{1}^{\mathsf{T}} q = \mathbf{1}^{\mathsf{T}} \widetilde{q}$ obtained from (1.5), we have $\widetilde{q} \succeq 0$. Therefore, $\widetilde{q} \succ 0$ by an argument similar to the one used for q. Consequently, $(V^{\mathsf{T}} q)_i (V \widetilde{q})_i > 0$ for all $i \in [n]$, leading to the contradiction

$$1 = \rho(\Psi(\vec{q})V^{\mathsf{T}}) < \rho(s^{-2}V^{\mathsf{T}}) = s^{-2}\rho(V) \le 1$$

where the strict inequality is due to Proposition 4.1. Hence $\vec{q} = 0$.

We now turn to Item (2). The argument $\vec{q} \succcurlyeq_{\neq} 0 \Rightarrow \vec{q} \succ 0$ is identical to Item (1).

The first step towards establishing uniqueness of the solution is showing that if $\vec{q} = (q^{\mathsf{T}}, \tilde{q}^{\mathsf{T}})^{\mathsf{T}}$ and $\vec{q'} = ((q')^{\mathsf{T}}, (\tilde{q'})^{\mathsf{T}})^{\mathsf{T}}$ are two positive solutions such that $\vec{q} \neq \vec{q'}$, then $\Psi(\vec{q}) \neq \Psi(\vec{q'})$. Assume the contrary. The equation $q = \Psi(\vec{q})V^{\mathsf{T}}q$ shows that 1 is the Perron-Frobenius eigenvalue of the irreducible matrix $\Psi(\vec{q})V^{\mathsf{T}}$ (Proposition 4.2). Since its eigenspace has the dimension one, we get that $q = \alpha q'$ for some $\alpha > 0$. A similar

argument shows that $\tilde{q} = \tilde{\alpha} \tilde{q}'$ for some $\tilde{\alpha} > 0$. Using the assumption $\Psi(\vec{q}) = \Psi(\vec{q'})$ again and inspecting the expressions of these terms, we get that $\alpha = \tilde{\alpha}^{-1}$. Moreover, the equations $\mathbf{1}^{\mathsf{T}} q = \mathbf{1}^{\mathsf{T}} \tilde{q}$ and $\mathbf{1}^{\mathsf{T}} q' = \mathbf{1}^{\mathsf{T}} \tilde{q}'$ show that $\alpha = \tilde{\alpha}$. This implies that $\vec{q} = \vec{q'}$, a contradiction.

To establish the uniqueness, let us still consider the two positive solutions $\vec{q} \neq \vec{q'}$. Write $\Psi = \text{diag}(\psi_i) = \Psi(\vec{q})$ and $\Psi' = \text{diag}(\psi'_i) = \Psi(\vec{q'})$, and define the vectors

$$\boldsymbol{\varphi} = \begin{pmatrix} \varphi_1 \\ \vdots \\ \varphi_n \end{pmatrix} = V \tilde{\boldsymbol{q}}, \quad \tilde{\boldsymbol{\varphi}} = \begin{pmatrix} \tilde{\varphi}_1 \\ \vdots \\ \tilde{\varphi}_n \end{pmatrix} = V^{\mathsf{T}} \boldsymbol{q}, \quad \boldsymbol{\varphi} = \begin{pmatrix} \boldsymbol{\varphi} \\ \tilde{\boldsymbol{\varphi}} \end{pmatrix},$$

and their similarly defined analogues φ' , $\tilde{\varphi}'$, and $\tilde{\varphi}'$. It holds by the irreducibility of V that $\tilde{\varphi}, \tilde{\varphi}' \succ 0$. We now write

$$\varphi_i = \frac{1}{n} \sum_{\ell=1}^n \sigma_{i,\ell}^2 \psi_\ell \varphi_\ell = \frac{1}{n} \sum_{\ell=1}^n \sigma_{i,\ell}^2 \psi_\ell^2 (s^2 \varphi_\ell + \varphi_\ell^2 \tilde{\varphi}_\ell)$$

and a similar equation for $\tilde{\varphi}_i$, giving rise to the identity

$$\vec{\varphi} = \begin{pmatrix} s^2 V \Psi^2 & V \Psi^2 \Phi^2 \\ V^{\mathsf{T}} \Psi^2 \widetilde{\Phi}^2 & s^2 V^{\mathsf{T}} \Psi^2 \end{pmatrix} \vec{\varphi}$$

where $\mathbf{\Phi} = \operatorname{diag}(\varphi_i)$ and $\widetilde{\mathbf{\Phi}} = \operatorname{diag}(\widetilde{\varphi}_i)$. Equivalently, the nonnegative matrix

$$K_{\vec{q}} := \begin{pmatrix} s^2 \Phi^{-1} V \Psi^2 \Phi & \Phi^{-1} V \Psi^2 \Phi^2 \widetilde{\Phi} \\ \widetilde{\Phi}^{-1} V^{\mathsf{T}} \Psi^2 \widetilde{\Phi}^2 \Phi & s^2 \widetilde{\Phi}^{-1} V^{\mathsf{T}} \Psi^2 \widetilde{\Phi} \end{pmatrix}$$

satisfies $K_{\vec{q}}\mathbf{1} = \mathbf{1}$. Considering now the two solutions $\vec{\varphi}$ and $\vec{\varphi}'$, we can write

$$\begin{split} \varepsilon_{i} &:= \Big| \frac{\varphi_{i} - \varphi_{i}'}{\sqrt{\varphi_{i}\varphi_{i}'}} \Big| \\ &= \frac{1}{\sqrt{\varphi_{i}\varphi_{i}'}} \frac{1}{n} \Big| \sum_{\ell=1}^{n} \sigma_{i,\ell}^{2} (\psi_{\ell}\varphi_{\ell} - \psi_{\ell}'\varphi_{\ell}') \Big| = \frac{1}{\sqrt{\varphi_{i}\varphi_{i}'}} \frac{1}{n} \Big| \sum_{\ell=1}^{n} \sigma_{i,\ell}^{2} \psi_{\ell} \psi_{\ell}' ((\psi_{\ell}')^{-1}\varphi_{\ell} - \psi_{\ell}^{-1}\varphi_{\ell}') \Big| \\ &= \frac{1}{n} \Big| \sum_{\ell=1}^{n} \Big(\frac{\sigma_{i,\ell}^{2} s^{2} \psi_{\ell} \psi_{\ell}' \sqrt{\varphi_{\ell}\varphi_{\ell}'}}{\sqrt{\varphi_{i}\varphi_{i}'}} \frac{\varphi_{\ell} - \varphi_{\ell}'}{\sqrt{\varphi_{\ell}\varphi_{\ell}'}} + \frac{\sigma_{i,\ell}^{2} \psi_{\ell} \psi_{\ell}' \varphi_{\ell} \varphi_{\ell}' \sqrt{\tilde{\varphi_{\ell}}\tilde{\varphi_{\ell}'}}}{\sqrt{\varphi_{i}\varphi_{\ell}'}} \Big) \Big| \\ &\leq \frac{1}{n} \sum_{\ell=1}^{n} \Big(\frac{\sigma_{i,\ell}^{2} s^{2} \psi_{\ell} \psi_{\ell}' \sqrt{\varphi_{\ell}\varphi_{\ell}'}}{\sqrt{\varphi_{i}\varphi_{i}'}} \Big| \frac{\varphi_{\ell} - \varphi_{\ell}'}{\sqrt{\varphi_{\ell}\varphi_{\ell}'}} \Big| + \frac{\sigma_{i,\ell}^{2} \psi_{\ell} \psi_{\ell}' \varphi_{\ell} \varphi_{\ell}' \sqrt{\tilde{\varphi_{\ell}}\tilde{\varphi_{\ell}'}}}{\sqrt{\varphi_{i}\varphi_{\ell}'}} \Big| \frac{\tilde{\varphi_{\ell}} - \tilde{\varphi_{\ell}'}}{\sqrt{\tilde{\varphi_{\ell}}\tilde{\varphi_{\ell}'}}} \Big| \Big) \end{split}$$

for every $i \in [n]$, and we also have a similar inequality for $\tilde{\varepsilon}_i := |(\tilde{\varphi}_i - \tilde{\varphi}'_i)/\sqrt{\tilde{\varphi}_i \tilde{\varphi}'_i}|$. It results that the vector $\boldsymbol{\varepsilon} = (\varepsilon_1, \ldots, \varepsilon_n, \tilde{\varepsilon}_1, \ldots, \tilde{\varepsilon}_n)^{\mathsf{T}}$ satisfies the inequality $\boldsymbol{\varepsilon} \preccurlyeq K_{\vec{q}, \vec{q}'} \boldsymbol{\varepsilon}$, where

$$K_{\vec{q},\vec{q}'} := \begin{pmatrix} s^2 (\Phi \Phi')^{-1/2} V \Psi \Psi' (\Phi \Phi')^{1/2} & (\Phi \Phi')^{-1/2} V \Psi \Psi' \Phi \Phi' (\widetilde{\Phi} \widetilde{\Phi}')^{1/2} \\ (\widetilde{\Phi} \widetilde{\Phi}')^{-1/2} V^{\mathsf{T}} \Psi \Psi' \widetilde{\Phi} \widetilde{\Phi}' (\Phi \Phi')^{1/2} & s^2 (\widetilde{\Phi} \widetilde{\Phi}')^{-1/2} V^{\mathsf{T}} \Psi \Psi' (\widetilde{\Phi} \widetilde{\Phi}')^{1/2} \end{pmatrix},$$

and $\Phi' = \operatorname{diag}(\varphi'_i)$, $\widetilde{\Phi}' = \operatorname{diag}(\widetilde{\varphi}'_i)$. By applying the Cauchy-Schwarz inequality to the scalar products $\boldsymbol{x}_m \mathbf{1}, m = 1, \ldots, n$, where \boldsymbol{x}_m is the row m of $K_{\vec{q},\vec{q}'}$, we get that $K_{\vec{q},\vec{q}'} \mathbf{1} \preccurlyeq (K_{\vec{q}}\mathbf{1}) \odot (K_{\vec{q}'}\mathbf{1}) = \mathbf{1}$.

Now, for any $k \in \mathbb{N}$, we have

$$\begin{pmatrix} V & V \\ V^{\mathsf{T}} & V^{\mathsf{T}} \end{pmatrix}^{k} \succcurlyeq \begin{pmatrix} V^{k} & V^{k} \\ (V^{\mathsf{T}})^{k} & (V^{\mathsf{T}})^{k} \end{pmatrix}.$$

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Since $\vec{\varphi}, \vec{\varphi}' \succ 0$ and V is irreducible, it holds that for any $(i, j) \in [2n]^2$, there exists $k \in [n]$ such that $[K_{\vec{q},\vec{q}'}^k]_{ij} > 0$, implying that $K_{\vec{q},\vec{q}'}$ is irreducible. Relying on these results, we will show that there exists $i \in [n]$ such that $x_i \mathbf{1} < 1$, i.e. Cauchy-Schwarz inequality is strict for this row vector. Proposition 4.2 will then show that $\rho(K_{\vec{q},\vec{q}'}) < 1$. By consequence, the only solution to the inequality $\varepsilon \preccurlyeq K_{\vec{q},\vec{q}'}\varepsilon$ will be $\varepsilon = 0$, contradicting the assertion $\vec{q} \neq \vec{q'}$, and the uniqueness of the solution of (1.5) for $\vec{q} \succcurlyeq_{\neq} 0$ will follow.

Recalling that $\Psi(\vec{q}) \neq \Psi(\vec{q'})$, there exists $\ell \in [n]$ such that $\varphi_{\ell} \tilde{\varphi}_{\ell} \neq \varphi'_{\ell} \tilde{\varphi}'_{\ell}$. Since V is irreducible, no column of this matrix is zero. Therefore, we can choose $i \in [n]$ such that $\sigma_{i,\ell}^2 > 0$. Consider the vector

$$\boldsymbol{v}_i := \left(\left(\frac{s\sigma_{i,m}\psi_m(\vec{\boldsymbol{q}})\sqrt{\varphi_m}}{\sqrt{n}\sqrt{\varphi_i}} \right)_{m=1}^n, \left(\frac{\sigma_{i,m}\psi_m(\vec{\boldsymbol{q}})\varphi_m\sqrt{\tilde{\varphi}_m}}{\sqrt{n}\sqrt{\varphi_i}} \right)_{m=1}^n \right) = \left((v_{1,m})_{m=1}^n, (v_{2,m})_{m=1}^n \right)$$

and his analogue v'_i (with the obvious notations). Consider also the 2×2 matrix

$$M = \begin{pmatrix} v_{1,\ell} & v_{2,\ell} \\ v_{1,\ell}' & v_{2,\ell}' \end{pmatrix} = \begin{pmatrix} \frac{\sigma_{i,\ell}\psi_\ell(\vec{q})}{\sqrt{n}\sqrt{\varphi_i}} & \\ & \frac{\sigma_{i,\ell}\psi_\ell(\vec{q}')}{\sqrt{n}\sqrt{\varphi_i'}} \end{pmatrix} \begin{pmatrix} \sqrt{\varphi_\ell} & \sqrt{\tilde{\varphi}_\ell}\varphi_\ell \\ \sqrt{\varphi'_\ell} & \sqrt{\tilde{\varphi}'_\ell}\varphi'_\ell \end{pmatrix} \begin{pmatrix} s \\ & 1 \end{pmatrix} := M_1 M_2 M_3.$$

Since $\det(M_2) = \sqrt{\varphi_\ell \varphi'_\ell} (\sqrt{\varphi'_\ell \tilde{\varphi}'_\ell} - \sqrt{\varphi_\ell \tilde{\varphi}_\ell}) \neq 0$, the vectors v_i and v'_i are not collinear. Therefore,

$$x_i \mathbf{1} = \langle v_i, v'_i \rangle < \|v_i\|^2 \|v'_i\|^2 = (K_{\vec{q}} \mathbf{1})_i (K_{\vec{q}'} \mathbf{1})_i = 1$$

Lemma 4.3 is proved.

It remains to establish Theorem 2.2–(3). This is a consequence of following lemma, which also provides an expression for $\nabla \vec{q}(s)$ on $(0, \rho(V)^{1/2})$.

Lemma 4.4. Assume that the nonnegative $n \times n$ matrix V is irreducible. Then the function $s \mapsto \vec{q}(s)$ is continuous on $(0, \infty)$, and is continuously differentiable on $(0, \rho(V)^{1/2}) \cup (\rho(V)^{1/2}, \infty)$. Setting

$$\begin{split} M(s) &= \begin{pmatrix} s^2 \Psi(\vec{q}(s))^2 V^{\mathsf{T}} & -\Psi(\vec{q}(s))^2 \tilde{\Phi}(s)^2 V \\ -\Psi(\vec{q}(s))^2 \Phi(s)^2 V^{\mathsf{T}} & s^2 \Psi(\vec{q}(s))^2 V \end{pmatrix}, \\ \Phi(s) &= \operatorname{diag}(\varphi_i(s))_{i=1}^n, \ \tilde{\Phi}(s) = \operatorname{diag}(\tilde{\varphi}_i(s))_{i=1}^n, \ \varphi_i(s) = (V\tilde{q}(s))_i, \ \tilde{\varphi}_i(s) = (V^{\mathsf{T}}q(s))_i, \\ A(s) &= \begin{pmatrix} I_{2n} - M(s) \\ (\mathbf{1}_n^{\mathsf{T}} & -\mathbf{1}_n^{\mathsf{T}}) \end{pmatrix} \in \mathbb{R}^{(2n+1)\times 2n}, \quad \text{and} \quad b(s) = - \begin{pmatrix} \Psi(\vec{q}(s))^2 V^{\mathsf{T}}q(s) \\ \Psi(\vec{q}(s))^2 V \tilde{q}(s) \\ 0 \end{pmatrix} \in \mathbb{R}^{2n+1}, \end{split}$$

the matrix A(s) has a full column rank, and

$$\nabla \vec{q}(s) = 2sA(s)^{-L}b(s),$$

where $A(s)^{-L}$ is the left inverse of A(s). On $(\rho(V)^{1/2}, \infty)$, it holds that $\vec{q}(s) = \nabla \vec{q}(s) = 0$.

Proof. We already know that $\vec{q}(s) = 0$ on $[\rho(V)^{1/2}, \infty)$, and we can easily check that $\nabla \vec{q}(s) = 2sA(s)^{-L}b(s)$ on $(\rho(V)^{1/2}, \infty)$.

Let us show that q(s) is continuous on $(0, \rho(V)^{1/2})$. Fix $s \in (0, \rho(V)^{1/2})$. When u belongs to a small neighborhood of s in $(0, \rho(V)^{1/2})$, the function $\vec{q}(u)$ is bounded by Lemma 4.1, since $\vec{q}(u)$ is the limit as $t \downarrow 0$ of the bounded function $\vec{q}(u,t)$. Let $u_k \to_k s$ be such that $\vec{q}(u_k) \to_k \vec{q}_*$. The vector \vec{q}_* is clearly a solution to (1.5). Observing that $\rho(\Psi(\vec{q}(u_k))V) = 1$, we get by the continuity of the spectral radius that $\rho(\Psi(\vec{q}_*)V) = 1$. If \vec{q}_* were equal to zero, then we would have $\rho(\Psi(\vec{q}_*)V) = \rho(s^{-2}V) > 1$, a contradiction. Therefore $\vec{q} \succcurlyeq_{\neq} 0$, and by Lemma 4.3–(2), $\vec{q}_* = \vec{q}(s)$ since $\vec{q}(s)$ is the only nonnegative and non zero solution to (1.5).

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 \square

To obtain the continuity of $\vec{q}(s)$ on $(0, \infty)$, all what remains to prove is that $\vec{q}(u) \to 0$ as $u \uparrow \rho(V)^{1/2}$. Relying on Lemma 4.1, take a sequence $u_k \uparrow_k \rho(V)^{1/2}$ such that $\vec{q}(u_k) \to_k \vec{q}_*$. Then, the same argument as in the proof of Lemma 4.3–(1) shows that $\vec{q}_* = 0$.

To establish the differentiability of $ec{q}(s)$ on $(0,
ho(V)^{1/2})$, we start by writing

$$q_i(s) = \psi_i \tilde{\varphi}_i = \psi_i^2 (s^2 \tilde{\varphi}_i + \tilde{\varphi}_i^2 \varphi_i) = \psi_i^2 (s^2 (V^{\mathsf{T}} \boldsymbol{q})_i + \tilde{\varphi}_i^2 (V \boldsymbol{\tilde{q}})_i).$$

Doing a similar derivation for $\tilde{q}_i(s)$, we get the equation $\vec{q}(s) = N(s)\vec{q}(s)$, where

$$N(s) = \begin{pmatrix} s^2 \Psi(\vec{q}(s))^2 V^{\mathsf{T}} & \Psi(\vec{q}(s))^2 \tilde{\Phi}(s)^2 V \\ \Psi(\vec{q}(s))^2 \Phi(s)^2 V^{\mathsf{T}} & s^2 \Psi(\vec{q}(s))^2 V \end{pmatrix},$$

As in the proof of Lemma 4.3–(2), we can show that N(s) is irreducible. Thus, the Perron–Frobenius eigenvalue of N(s) is equal to one, it is algebraically simple, and its associated eigenspace is generated by $\vec{q}(s)$.

Now, given two real numbers $s,s' \in (0,\rho(V)^{1/2})$ with $s \neq s'$, we have

$$q_i - q'_i = \psi_i \tilde{\varphi}_i - \psi'_i \tilde{\varphi}'_i = \psi_i \psi'_i ((\psi'_i)^{-1} \tilde{\varphi}_i - \psi_i^{-1} \tilde{\varphi}'_i) = \psi_i \psi'_i (-(s^2 - s'^2) \tilde{\varphi}_i + s^2 (\tilde{\varphi}_i - \tilde{\varphi}'_i) - \tilde{\varphi}_i \tilde{\varphi}'_i (\varphi_i - \varphi'_i))$$

where we set $q_i = q_i(s)$ and $q'_i = q(s')$, and we used the same notational shortcut for all the other quantities. We thus have

$$\frac{q_i - q_i'}{s^2 - s'^2} = -\left(\Psi \Psi' V^\mathsf{T} \boldsymbol{q}\right)_i + \left(s^2 \Psi \Psi' V^\mathsf{T} \frac{\boldsymbol{q} - \boldsymbol{q}'}{s^2 - s'^2}\right)_i - \left(\Psi \Psi' \widetilde{\boldsymbol{\Phi}} \widetilde{\boldsymbol{\Phi}}' V \frac{\widetilde{\boldsymbol{q}} - \widetilde{\boldsymbol{q}}'}{s^2 - s'^2}\right)_i.$$

Doing a similar derivation for $\tilde{q}_i - \tilde{q}'_i$, we obtain the system

$$(I - \boldsymbol{M}(s, s'))\frac{\boldsymbol{\vec{q}} - \boldsymbol{\vec{q}}'}{s^2 - s'^2} = \boldsymbol{a}(s, s')$$

where

$$\boldsymbol{M}(\boldsymbol{s},\boldsymbol{s}') = \begin{pmatrix} \boldsymbol{s}^2 \boldsymbol{\Psi} \boldsymbol{\Psi}' \boldsymbol{V}^\mathsf{T} & -\boldsymbol{\Psi} \boldsymbol{\Psi}' \boldsymbol{\widetilde{\Phi}} \boldsymbol{\widetilde{\Phi}}' \boldsymbol{V} \\ -\boldsymbol{\Psi} \boldsymbol{\Psi}' \boldsymbol{\Phi} \boldsymbol{\Phi}' \boldsymbol{V}^\mathsf{T} & \boldsymbol{s}^2 \boldsymbol{\Psi} \boldsymbol{\Psi}' \boldsymbol{V} \end{pmatrix} \quad \text{and} \quad \boldsymbol{a}(\boldsymbol{s},\boldsymbol{s}') = -\begin{pmatrix} \boldsymbol{\Psi} \boldsymbol{\Psi}' \boldsymbol{V}^\mathsf{T} \boldsymbol{q} \\ \boldsymbol{\Psi} \boldsymbol{\Psi}' \boldsymbol{V} \boldsymbol{\widetilde{q}} \end{pmatrix}$$

Using in addition the identity $\sum q_i = \sum \tilde{q}_i$, we get the system $A(s, s')(\vec{q} - \vec{q}')/(s^2 - s'^2) = b(s, s')$, where

$$\boldsymbol{A}(s,s') = \begin{pmatrix} I - \boldsymbol{M}(s,s') \\ \begin{pmatrix} \mathbf{1}_n^\mathsf{T} & -\mathbf{1}_n^\mathsf{T} \end{pmatrix} \end{pmatrix} \in \mathbb{R}^{(2n+1) \times 2n} \quad \text{and} \quad \boldsymbol{b}(s,s') = \begin{pmatrix} \boldsymbol{a}(s,s') \\ 0 \end{pmatrix}.$$

By the continuity of $\vec{q}(s)$, $A(s,s') \to A(s)$ and $b(s,s') \to b(s)$ as $s' \to s$. It is easy to see that $(\boldsymbol{x}^{\mathsf{T}}, \tilde{\boldsymbol{x}}^{\mathsf{T}})^{\mathsf{T}}$ is an eigenvector of M(s) if and only if $(\boldsymbol{x}^{\mathsf{T}}, -\tilde{\boldsymbol{x}}^{\mathsf{T}})^{\mathsf{T}}$ is an eigenvector of N(s). Thus, the right null space I - M(s) is spanned by $\boldsymbol{v}(s) := (\boldsymbol{q}(s)^{\mathsf{T}}, -\tilde{\boldsymbol{q}}(s)^{\mathsf{T}})^{\mathsf{T}}$. But we clearly have $A(s)\boldsymbol{v}(s) \neq 0$, hence the matrix A(s) has a full column rank. Thus, for s' close enough to s,

$$\frac{\vec{q} - \vec{q}'}{s^2 - s'^2} = \boldsymbol{A}(s, s')^{-L} \begin{pmatrix} \boldsymbol{a}(s, s') \\ 0 \end{pmatrix} \xrightarrow[s' \to s]{} A(s)^{-L} b(s)$$

which shows that $\vec{q}(s)$ is differentiable for any $s \in (0, \rho(V)^{1/2})$, with the gradient $2sA(s)^{-L}b(s)$. The continuity of this gradient follows from the continuity of A(s) and b(s) and the fact that A(s) has a full column rank for any $s \in (0, \rho(V)^{1/2})$.

5 Bounding solutions to the Regularized Master Equations via graphical bootstrapping

In this section we are concerned with establishing bounds on the solution $\vec{r}(s,t) \succ 0$ to the Regularized Master Equations (2.7) that are uniform in the regularization parameter t > 0. Here we will view the standard deviation profile A and all parameters as fixed. Hence, we fix $n \ge 1$ and consider an arbitrary nonnegative $n \times n$ matrix $A = (\sigma_{ij})$. Putting $V = \frac{1}{n}A \odot A$ and fixing s, t > 0, we let $\vec{r} = \vec{r}(s,t)$ denote the unique solution to the Regularized Master Equations satisfying $\vec{r} \succ 0$, as is provided by Proposition 2.1.

5.1 Some preparation and proofs of Propositions 2.5 and 2.6

We begin by recording some key estimates and identities that will be used repeatedly in the sequel. We may write the Regularized Master Equations (2.7) in component form as

$$\frac{1}{r_i} = \varphi_i + \frac{s^2}{\widetilde{\varphi}_i}, \qquad \frac{1}{\widetilde{r}_i} = \widetilde{\varphi}_i + \frac{s^2}{\varphi_i}$$
(5.1)

where

$$\widetilde{\varphi}_i := t + (V^{\mathsf{T}} \boldsymbol{r})_i, \qquad \varphi_i := t + (V \widetilde{\boldsymbol{r}})_i$$
(5.2)

(from Proposition 2.1 we have $r_i, \tilde{r}_i > 0$ for all $i \in [n]$, so we are free to take reciprocals). Additionally from Proposition 2.1 we have trace identity:

$$\frac{1}{n}\sum_{j=1}^{n}r_{j} = \frac{1}{n}\sum_{j=1}^{n}\widetilde{r}_{j}.$$
(5.3)

From (5.1) it is immediate that

$$\frac{1}{2} \le \frac{r_i}{\min\left(\tilde{\varphi}_i/s^2, 1/\varphi_i\right)}, \ \frac{\tilde{r}_i}{\min\left(\varphi_i/s^2, 1/\tilde{\varphi}_i\right)} \le 1$$
(5.4)

for all $i \in [n]$. We can similarly bound the product

$$r_i \widetilde{r}_i = \frac{\varphi_i \widetilde{\varphi}_i}{(s^2 + \varphi_i \widetilde{\varphi}_i)^2} \le \frac{1}{s^2} \min\left(\frac{\varphi_i \widetilde{\varphi}_i}{s^2}, \frac{s^2}{\varphi_i \widetilde{\varphi}_i}\right) \le 1/s^2.$$
(5.5)

Hence, if for some $i \in [n]$ one of r_i, \tilde{r}_i is large, the other is small. Finally, we note the trivial upper bounds

$$r_i, \, \widetilde{r}_i \le 1/t \quad \forall i \in [n]. \tag{5.6}$$

We now prove Propositions 2.5 and 2.6.

Proof of Proposition 2.5. Assume towards a contradiction that $\frac{1}{n} \sum_{j=1}^{n} r_j > 1/\sigma_{\min}$. Then there exists $i \in [n]$ such that $r_i > 1/\sigma_{\min}$. From (5.4) it follows that

$$1/\sigma_{\min} < 1/\varphi_i = 1/(t + (V\widetilde{\boldsymbol{r}})_i)$$

and so

$$\sigma_{\min} > (V\widetilde{\boldsymbol{r}})_i = \frac{1}{n} \sum_{j=1}^n \sigma_{ij}^2 \widetilde{r}_j \ge \frac{\sigma_{\min}^2}{n} \sum_{j=1}^n \widetilde{r}_j.$$

Rearranging, we find $\frac{1}{n} \sum_{j=1}^{n} \widetilde{r}_j < 1/\sigma_{\min}$. From (5.3) it follows that $\frac{1}{n} \sum_{j=1}^{n} r_j < 1/\sigma_{\min}$, a contradiction.

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Proof of Proposition 2.6. If $V = \frac{1}{n}A \odot A$ is symmetric, then the 2n regularized master equations merge into n equations since $\mathbf{r} = \tilde{\mathbf{r}}$. In fact since $V^{\mathsf{T}} = V$ then if $\vec{\mathbf{r}} = (\mathbf{r}^{\mathsf{T}} \, \tilde{\mathbf{r}}^{\mathsf{T}})^{\mathsf{T}}$ is a solution of the Regularized Master Equations, so is $\check{\mathbf{r}} = (\tilde{\mathbf{r}}^{\mathsf{T}} \, \mathbf{r}^{\mathsf{T}})^{\mathsf{T}}$. By uniqueness, $\mathbf{r} = \tilde{\mathbf{r}}$. Hence the Regularized Master Equations write

$$r_i = \frac{(Vr)_i + t}{s^2 + (Vr)_i^2}, \quad i \in [n].$$

An elementary analysis of the function $f(x) = \frac{x}{s^2 + x^2}$ yields $\sup_{x \in [0,\infty)} f(x) \le (2s)^{-1}$. Hence

$$\frac{1}{n}\sum_{i\in[n]}r_i=\frac{1}{n}\sum_{i\in[n]}f((V\boldsymbol{r})_i+t) \ \leq \ \frac{1}{2s} \qquad \text{for all }t>0. \qquad \qquad \square$$

Our main objective now is to establish the following, which immediately yields Theorem 2.8.

Proposition 5.1. Assume $\sigma_{ij} \leq \sigma_{\max}$ for all $i, j \in [n]$ and some $\sigma_{\max} < \infty$, and that $A(\sigma_0)$ is (δ, κ) -robustly irreducible for some $\sigma_0, \delta, \kappa \in (0, 1)$ (see Definition 2.7). For every fixed $z \in \mathbb{C} \setminus \{0\}$ there exists a constant $K = K(z, \sigma_0, \sigma_{\max}, \delta, \kappa) < \infty$ such that

$$\frac{1}{n}\sum_{i=1}^{n}r_{i}\leq K.$$

In the proof of Proposition 2.5 we were able to pass from a lower bound on a single component r_i to an upper bound on the average $\frac{1}{n} \sum_{j=1}^{n} \tilde{r}_j$ in one line. This will not be possible when we allow some of the variances σ_{ij}^2 to be zero. We will employ a bootstrap-type argument which we roughly outline as follows:

- 1. Assume towards a contradiction that $\frac{1}{n}\sum_{i=1}^{n}r_i$ is large. By the pigeonhole principle there exists $i_0 \in [n]$ such that r_{i_0} is large.
- 2. Use the estimates (5.4)–(5.6) together with assumptions on the connectivity properties of the associated directed graph to iteratively "grow" the set of indices i for which we know r_i is large.
- 3. Once we have shown r_i is large for (almost) all $i \in [n]$, by (5.5) it follows that \tilde{r}_i is small for (almost) all $i \in [n]$. We then apply the trace constraint (5.3) to derive a contradiction.

We emphasize that the key idea for our proofs to bound $\frac{1}{n}\sum_{i}r_{i}$ will be to play (5.5) against the trace constraint (5.3).

Recall the graph-theoretic notation from Section 2.1. In this section we abbreviate

$$\mathcal{N}_{+}(i) := \mathcal{N}_{A(\sigma_0)}(i), \quad \mathcal{N}_{-}(i) := \mathcal{N}_{A(\sigma_0)^{\mathsf{T}}}(i)$$
(5.7)

for the in- and out-neighborhoods of a vertex i in the graph $\Gamma = \Gamma(A(\sigma_0))$, and similarly define $\mathcal{N}_{-}^{(\delta)}(i)$. For parameters $\alpha, \beta > 0$ we define the sets

$$S_{\alpha} = \{ i \in [n] : r_i \ge \alpha \| \widetilde{\varphi} \|_{\infty} \}, \qquad T_{\beta} = \{ i \in [n] : \widetilde{r}_i < \beta / \| \widetilde{\varphi} \|_{\infty} \}$$
(5.8)

where here and in the sequel we write $\|\widetilde{\varphi}\|_{\infty} = \max_{i \in [n]} \widetilde{\varphi}_i$ and similarly $\|\varphi\|_{\infty} = \max_{i \in [n]} \varphi_i$. (Note that by (5.2) we have $\|\widetilde{\varphi}\|_{\infty} \ge t > 0$.)

5.2 Qualitative boundedness

In this subsection we establish Proposition 2.9 and Lemma 4.1, which give an n-dependent bound on the components of r, \tilde{r} assuming only that the standard deviation profile is irreducible. This was used in Section 4 with a compactness argument to establish existence of solutions to the Master Equations. While Lemma 4.1 follows from Proposition 5.1 under the robust irreducibility assumption A5 (which we assume for our main result), we prove this lemma separately for two reasons:

- to show that A5 is not needed for the conclusion of Theorem 2.2, and
- to provide a cartoon for the more technical proof of Proposition 5.1.

We will also establish some auxiliary lemmas that will be reused in the proof of Proposition 5.1 (in particular Lemmas 5.3 and 5.4).

Proposition 2.9 and Lemma 4.1 are an immediate consequence of the following:

Proposition 5.2. For any fixed $0 < \sigma_0 \leq \sigma_{\max}$ and $s_0 > 0$ there is a constant $K_0(n, s_0, \sigma_0, \sigma_{\max})$ such that the following holds. Let $s \geq s_0$, and suppose A is irreducible with $\sigma_{ij} \in \{0\} \cup [\sigma_0, \sigma_{\max}]$ for all $1 \leq i, j \leq n$. Then $\frac{1}{n} \sum_{i=1}^n r_i \leq K_0$.

We now begin the proof of Proposition 5.2. First we dispose of the case that $s > 2\sigma_{\max}$: Claim 5.1. Suppose $s > 2\sigma_{\max}$. Then

$$r_i, \tilde{r}_i \le \min\left(\frac{t}{s^2 - \sigma_{\max}^2}, \frac{1}{t}\right) \le (s^2 - \sigma_{\max}^2)^{-1/2}.$$
 (5.9)

Proof. Suppose $r_{i^*} = \max_{i \in [n]} r_i$. From (5.4) we have

$$r_{i^*} \leq \frac{1}{s^2} \widetilde{\varphi}_i = \frac{1}{s^2} (t + (V^{\mathsf{T}} \boldsymbol{r})_i) \leq \frac{1}{s^2} (t + \sigma_{\max}^2 r_{i^*}).$$

Rearranging we obtain $r_{i^*} \leq t/(s^2 - \sigma_{\max}^2)$, which combines with (5.6) to give the desired uniform bound for r_i , $i \in [n]$. The same bound is obtained for \tilde{r}_i by similar lines.

Without loss of generality we take $\sigma_{\max} = 1$. By Claim 5.1 we may assume $0 < s_0 \le s \le 2$. We may also assume $t \le 1$. Indeed, otherwise it follows from (5.6) that $\frac{1}{n} \sum_{i=1}^{n} r_i < 1$ and we are done. Let K > 0 to be chosen later depending on n, σ_0 and s_0 , but independent of t, and assume

$$\frac{1}{n}\sum_{i=1}^{n}r_{i} \ge K.$$
(5.10)

We will derive a contradiction for *K* sufficiently large.

In the following lemma we use the irreducibility of A_n to show that if T_β is non-empty for some β sufficiently small, then $T_{\beta'} = [n]$ for a somewhat larger value of β' . This will allow us to assume a uniform lower bound on the components \tilde{r}_i .

Lemma 5.2. There are positive constants $C_0(\sigma_0, n)$, $\beta_0(\sigma_0, n)$ such that for all $\beta \leq \beta_0$, if T_β is non-empty then $T_{C_0\beta} = [n]$.

Proof. Let $\beta > 0$ to be taken sufficiently small depending on σ_0, n , and suppose T_β is non-empty. Then there exists $i \in [n]$ such that $\tilde{r}_i < \beta/\|\tilde{\varphi}\|_{\infty}$. From (5.4) it follows that

$$\frac{1}{2}\min(\varphi_i/s^2, 1/\widetilde{\varphi}_i) < \beta/\|\widetilde{\varphi}\|_{\infty}.$$

Assuming $\beta \leq 1/2$ it follows that

$$2s^{2}\beta/\|\widetilde{\varphi}\|_{\infty} > \varphi_{i} \ge (V\widetilde{r})_{i} \ge \sigma_{0}^{2}\frac{1}{n}\sum_{j\in\mathcal{N}_{+}(i)}\widetilde{r}_{j}$$

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and hence

$$\widetilde{r}_j < \frac{2s^2n}{\sigma_0^2} \frac{\beta}{\|\widetilde{\varphi}\|_{\infty}} \qquad \forall j \in \mathcal{N}_+(i).$$
(5.11)

Again from (5.4), if we further assume $\beta \leq \sigma_0^2/4s^2n$ then it follows that

$$\frac{4s^4n}{\sigma_0^2} \frac{\beta}{\|\widetilde{\varphi}\|_{\infty}} > \varphi_j \qquad \forall j \in \mathcal{N}_+(i).$$
(5.12)

Now let $k \in [n]$ be arbitrary. By the irreducibility of A_n there exists a directed path in the associated digraph Γ_n from vertex k to vertex i of length at most n. Applying the above lines iteratively along each edge of the path we find

$$\widetilde{r}_k \le \left(\frac{2s^2n}{\sigma_0^2}\right)^n \frac{\beta}{\|\widetilde{\varphi}\|_{\infty}}$$

if we take $\beta \leq \frac{1}{2} (\frac{\sigma_0^2}{2s^2n})^{n-1}$. Since k was arbitrary, the result follows by setting $C_0 = (8n/\sigma_0^2)^n$ and $\beta_0 = \frac{1}{2} (\sigma_0^2/(8n))^{n-1}$ (here we have used our assumption $s \leq 2$).

If $T_{C_0\beta} = [n]$ for some $\beta \leq \beta_0$ then by the trace identity (5.3),

$$\frac{1}{n}\sum_{i=1}^{n}r_{i} = \frac{1}{n}\sum_{i=1}^{n}\widetilde{r}_{i} \le C_{0}\beta/\|\widetilde{\varphi}\|_{\infty}.$$
(5.13)

On the other hand, from (5.4) we have

$$r_j \le \widetilde{\varphi}_j / s^2 \le \|\widetilde{\varphi}\|_{\infty} / s^2 \tag{5.14}$$

for all $j \in [n]$. In particular,

$$\frac{1}{n}\sum_{j=1}^{n}r_{j} \le \|\widetilde{\varphi}\|_{\infty}/s^{2}$$
(5.15)

Hence,

$$\frac{1}{n}\sum_{i=1}^{n}r_{i} \le \min\left(\frac{C_{0}\beta}{\|\tilde{\varphi}\|_{\infty}}, \frac{\|\tilde{\varphi}\|_{\infty}}{s^{2}}\right) \le (C_{0}\beta/s^{2})^{1/2} \le (C_{0}\beta_{0}/s_{0}^{2})^{1/2},$$
(5.16)

which contradicts (5.10) if K is sufficiently large. Hence we may assume T_{β_0} is empty for $\beta_0(\sigma_0, n)$ as in Lemma 5.2. Thus,

$$\widetilde{r}_i \ge \beta_0 / \|\widetilde{\varphi}\|_{\infty} \quad \forall i \in [n].$$
(5.17)

Now we find a value of α for which S_{α} is already of linear size:

Lemma 5.3. Assume $K \ge 2/s^2$. Then $|S_{1/4}| \ge (s^2/4)n$.

Proof. From our assumption and (5.15),

$$2/s^2 \le K \le \frac{1}{n} \sum_{j=1}^n r_j \le \|\widetilde{\varphi}\|_{\infty}/s^2,$$
(5.18)

so $\|\widetilde{\varphi}\|_{\infty} \geq 2$. Let $i \in [n]$ such that $\widetilde{\varphi}_i = \|\widetilde{\varphi}\|_{\infty}$. We have

$$\|\widetilde{\varphi}\|_{\infty} = \widetilde{\varphi}_i = t + \frac{1}{n} \sum_{j=1}^n \sigma_{ji}^2 r_j \le t + \frac{1}{n} \sum_{j=1}^n r_j.$$

Since $\|\widetilde{\varphi}\|_{\infty} \ge 2$ and $t \le 1$, $\frac{1}{n} \sum_{j=1}^n r_j \ge \frac{1}{2} \|\widetilde{\varphi}\|_{\infty}$. Now again by (5.14),

$$\|\widetilde{\boldsymbol{\varphi}}\|_{\infty}n/2 \leq \sum_{j \in S_{1/4}} r_j + \sum_{j \in S_{1/4}^c} r_j \leq (\|\widetilde{\boldsymbol{\varphi}}\|_{\infty}/s^2)|S_{1/4}| + \|\widetilde{\boldsymbol{\varphi}}\|_{\infty}n/4,$$

and the result follows by rearranging.

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Next we seek to show that we can enlarge S_{α} by lowering α . By irreducibility we can find a vertex $i^* \in S_{\alpha}$ that is connected to S_{α}^c . We can use this to show that the average of the components r_k over S_{α}^c is bounded below by cr_{i^*} for some small c > 0 depending on α, n, s, σ_0 . From the pigeonhole principle we obtain $k \in S_{\alpha}^c$ with $r_k \ge cr_{i^*} \ge c\alpha \|\tilde{\varphi}\|_{\infty}$. Taking $\alpha' = c\alpha$, we will then have shown $|S_{\alpha'}| \ge |S_{\alpha}| + 1$.

We begin by relating the values of r, \tilde{r} on a fixed set of vertices $S \subset [n]$ to the values taken on S^c . For an $n \times n$ matrix M and $S, T \subset [n]$ nonempty we write $M_{S \times T}$ for the $|S| \times |T|$ submatrix of M with entries indexed by $S \times T$. The following lemma will also be used in the proof of Proposition 5.1.

Lemma 5.4. Fix a nonempty set $S \subset [n]$, and recall the diagonal matrix Ψ from (4.1). The $|S| \times |S|$ matrix $\Psi_{S \times S}^{-1} - V_{S \times S}^{\mathsf{T}}$ is invertible, and we denote its inverse

$$W^{S} = (\Psi_{S \times S}^{-1} - V_{S \times S}^{\mathsf{T}})^{-1}.$$
(5.19)

In terms of W^S the restrictions of r, \tilde{r} to S and S^c satisfy

$$\boldsymbol{r}_{S} = W^{S}(t + V_{S^{c} \times S}^{\mathsf{T}} \boldsymbol{r}_{S^{c}}), \quad \widetilde{\boldsymbol{r}}_{S} = (W^{S})^{\mathsf{T}}(t + V_{S \times S^{c}} \widetilde{\boldsymbol{r}}_{S^{c}}).$$
(5.20)

Furthermore, the entries of W^S satisfy the following bounds. For all $i, j \in S$,

$$W_{ij}^S \ge 0. \tag{5.21}$$

For all $j \in S$,

$$\sum_{i \in S} W_{ij}^S \le \tilde{r}_j / t \tag{5.22}$$

and if (5.17) holds for some $\beta_0 > 0$.

$$\sum_{i\in S} W_{ij}^S |\mathcal{N}_+(i) \cap S^c| \le \left(\frac{n\|\widetilde{\varphi}\|_{\infty}}{\beta_0 \sigma_0^2}\right) \widetilde{r}_j.$$
(5.23)

Proof. Arguing as in the proof of Lemma 4.2(2) (using Proposition 4.2) we find the spectral radius of $\Psi_{S\times S}V_{S\times S}^{\mathsf{T}}$ is strictly less than 1. Hence, $(I - \Psi_{S\times S}V_{S\times S}^{\mathsf{T}})^{-1}$ has a convergent Neumann series, and it follows that

$$W^S = \Psi_{S \times S} (I - \Psi_{S \times S} V_{S \times S}^{\mathsf{T}})^{-1} = \Psi_{S \times S} \sum_{k=0}^{\infty} (\Psi_{S \times S} V_{S \times S}^{\mathsf{T}})^k$$

is well defined. Furthermore, since all of the matrices in the above series have nonnegative entries, (5.21) follows.

(5.20) is quickly obtained by rearranging the equations (2.7).

Now for (5.22) and (5.23), let $j \in S$ be arbitrary. From the second equation in (5.20),

$$\widetilde{r}_j = \sum_{i \in S} W_{ij}^S \Big(t + \frac{1}{n} \sum_{k \in S^c} \sigma_{ik}^2 \widetilde{r}_k \Big).$$

In particular, $\widetilde{r}_j \geq t \sum_{i \in S} W^S_{ij}$, giving (5.22), and

$$\widetilde{r}_j \ge \frac{1}{n} \sum_{i \in S} W_{ij}^S \sum_{k \in S^c} \sigma_{ik}^2 \widetilde{r}_k \ge \frac{\beta_0 \sigma_0^2}{n \|\widetilde{\varphi}\|_{\infty}} \sum_{i \in S} W_{ij}^S |\mathcal{N}_+(i) \cap S^c|$$

which rearranges to give (5.23).

We can use Lemma 5.4 and the irreducibility of A to establish the following:

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Lemma 5.5 (Incrementing α). Let $\alpha > 0$ such that $1 \le |S_{\alpha}| \le n - 1$. If $K > \sqrt{2n}/(s_0^3 \alpha)$ then there exists $\alpha' = \alpha'(\alpha, s_0, \sigma_0, n) \in (0, \alpha)$ such that $|S_{\alpha'}| \ge |S_{\alpha}| + 1$.

Let us conclude the proof of Proposition 5.2 on the above lemma. Putting $\alpha_0 = 1/4$, by Lemma 5.3 we have $S_{\alpha_0} \neq \emptyset$. Taking K sufficiently large depending on s_0, σ_0 and n we can iterate Lemma 5.5 at most n times to find $\alpha = \alpha(s_0, \sigma_0, n) > 0$ such that $S_\alpha = [n]$. Then by (5.3) and (5.4),

$$\alpha \|\widetilde{\boldsymbol{\varphi}}\|_{\infty} \leq \frac{1}{n} \sum_{j=1}^{n} r_j = \frac{1}{n} \sum_{j=1}^{n} \widetilde{r}_j \leq \frac{1}{s^2 \alpha \|\widetilde{\boldsymbol{\varphi}}\|_{\infty}},$$
(5.24)

so we have $\|\widetilde{\varphi}\|_{\infty} \leq 1/(s\alpha)$. Then again by (5.4) we have

$$\frac{1}{n}\sum_{j=1}^{n}r_{j} \leq \frac{\|\widetilde{\varphi}\|_{\infty}}{s^{2}} \leq \frac{1}{s^{3}\alpha} \leq \frac{1}{s_{0}^{3}\alpha}$$
(5.25)

and we are done.

Proof of Lemma 5.5. We write $W = W^{S_{\alpha}}$. From the first equation in (5.20), for any $i \in S_{\alpha}$ we have

$$\alpha \|\widetilde{\varphi}\|_{\infty} \le r_i = \sum_{j \in S_{\alpha}} W_{ij} \left(t + \sum_{k \in S_{\alpha}^c} \sigma_{kj}^2 r_k \right).$$
(5.26)

Suppose first that

$$\sum_{\in S_{\alpha}} W_{ij} > \frac{\alpha \|\tilde{\varphi}\|_{\infty}}{2t}$$
(5.27)

for some $i \in S_{\alpha}$. Then from (5.22),

$$\frac{\alpha \|\widetilde{\varphi}\|_{\infty}}{2t} < \sum_{j \in S_{\alpha}} W_{ij} \le \frac{1}{t} \sum_{j \in S_{\alpha}} \widetilde{r}_j \le \frac{1}{s^2 t} \sum_{j \in S_{\alpha}} \frac{1}{r_j} \le \frac{|S_{\alpha}|}{t s^2 \alpha \|\widetilde{\varphi}\|_{\infty}},$$

where in the third inequality we applied (5.4). Rearranging we have $\|\tilde{\varphi}\|_{\infty} \leq \sqrt{2|S|}/(s\alpha) \leq \sqrt{2n}/(s_0\alpha)$ in this case. On the other hand, from (5.15) we have $K \leq \|\tilde{\varphi}\|_{\infty}/s_0^2$, and we obtain a contradiction if $K > \sqrt{2n}/(s_0^3\alpha)$.

Suppose now that (5.27) does not hold for any $i \in S_{\alpha}$. Then rearranging (5.26) we have

$$\frac{\alpha \|\tilde{\varphi}\|_{\infty}}{2} \le \sum_{j \in S_{\alpha}} \sum_{k \in S_{\alpha}^c} \sigma_{kj}^2 r_k W_{ij}$$
(5.28)

for any $i \in S_{\alpha}$. From the assumption that A_n is irreducible there exists $(i^*, j^*) \in S_{\alpha} \times S_{\alpha}^c$ such that $\sigma_{i^*j^*} \ge \sigma_0$, i.e. $|\mathcal{N}_+(i^*) \cap S_{\alpha}^c| \ge 1$. From (5.23) it follows that

$$W_{i^*j} \le \left(\frac{n\|\tilde{\varphi}\|_{\infty}}{\beta_0 \sigma_0^2}\right) \tilde{r}_j \tag{5.29}$$

for all $j \in S_{\alpha}$. Inserting this bound in (5.28) we have

$$\frac{\alpha \|\widetilde{\varphi}\|_{\infty}}{2} \leq \frac{n \|\widetilde{\varphi}\|_{\infty}}{\beta_0 \sigma_0^2} \sum_{j \in S_{\alpha}} \sum_{k \in S_{\alpha}^c} \sigma_{kj}^2 r_k \widetilde{r}_j \leq \frac{n |S_{\alpha}|}{\beta_0 \sigma_0^2 s^2 \alpha} \sum_{k \in S_{\alpha}^c} r_k,$$

where in the second inequality we applied the bounds $\sigma_{ij} \leq 1$ for all $i, j \in [n]$ and $\tilde{r}_j \leq (s^2 \alpha \|\tilde{\varphi}\|_{\infty})^{-1}$ for all $j \in S_{\alpha}$ (by (5.5)). Rearranging we have

$$\sum_{k \in S_{\alpha}^{c}} r_{k} \ge \frac{\alpha^{2} s^{2} \beta_{0} \sigma_{0}^{2}}{2n |S_{\alpha}|} \| \widetilde{\varphi} \|_{\infty}.$$
(5.30)

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By the pigeonhole principle there exists $k \in S^c_{\alpha}$ such that

$$r_k \ge \frac{\alpha^2 s^2 \beta_0 \sigma_0^2}{2n |S_\alpha| |S_\alpha^c|} \|\widetilde{\varphi}\|_{\infty} \ge \frac{\alpha^2 s_0^2 \beta_0 \sigma_0^2}{2n^3} \|\widetilde{\varphi}\|_{\infty}.$$
(5.31)

Setting $\alpha' = \alpha^2 s_0^2 \beta_0 \sigma_0^2 / (2n^3)$ we have $k \in S_{\alpha'}$. Also, since $\alpha' < \alpha$ we have $S_{\alpha} \subset S_{\alpha'}$. Hence, $|S_{\alpha'}| \ge |S_{\alpha}| + 1$ as desired.

5.3 Quantitative boundedness

Now we prove Proposition 5.1. By rescaling the variance profile V we may take $\sigma_{\max} = 1$. By Claim 5.1 we may assume $s \in (0, 2]$. As in the proof of Proposition 5.2 we may assume $t \leq 1$. We may also assume n is sufficiently large depending on s, σ_0, δ and κ . In the remainder of the section we make use of asymptotic notation $O(), \leq, \geq$, allowing implied constants to depend on the parameters s, σ_0, δ and κ , but not on n and t.

As in the proof of Proposition 5.2 we assume (5.10) holds for some K > 0 and aim to derive a contradiction for K sufficiently large depending on s, σ_0, δ and κ . The argument follows the same general outline as the proof in the previous subsection. We will reuse Lemmas 5.3 and 5.4 as stated, but we will need versions of Lemmas 5.2 and 5.5 with constants independent of n.

5.3.1 Lower bounding \tilde{r}_i

The following is an analogue of Lemma 5.2.

Lemma 5.6. There are positive constants $C_0(s, \sigma_0, \delta, \kappa)$, $\beta_0(s, \sigma_0, \delta, \kappa)$ such that for all $\beta \leq \beta_0$, if T_β is non-empty then $T_{C_0\beta} = [n]$.

Proof. Let $\beta > 0$ to be taken sufficiently small and assume T_{β} is non-empty. Fix an element $i_0 \in T_{\beta}$. We will grow the set T_{β} in stages by enlarging β by appropriate constant factors. We do this by iterative application of the following:

Claim 5.7. Let $\beta, \varepsilon_0 \in (0, 1/2]$, and assume $0 < |T_\beta| \le (1 - \varepsilon_0)n$. There exists $C = C(\sigma_0, \delta, \varepsilon_0) > 0$ such that if n is sufficiently large depending on κ and ε_0 then $|T_{C\beta} \setminus T_\beta| \ge (\delta \varepsilon_0 / 2)n$.

Proof. By the assumption that $A(\sigma_0)$ is (δ, κ) -robustly irreducible we have

$$|\mathcal{N}^{(\delta)}_{-}(T^{c}_{\beta}) \cap T_{\beta}| \geq \min(\kappa |T^{c}_{\beta}|, |T_{\beta}|) \geq \min(\kappa \varepsilon_{0} n, 1) \geq 1$$

if n is sufficiently large. Fix an element $i \in \mathcal{N}_{-}^{(\delta)}(T_{\beta}^{c}) \cap T_{\beta}$. By definition we have

$$\mathcal{N}_{+}(i) \cap T^{c}_{\beta} \geq \delta |T^{c}_{\beta}| \geq \delta \varepsilon_{0} n.$$

Next, we claim that for any C > 0 we have

$$|\mathcal{N}_{+}(i) \cap T_{C\beta}^{c}| \le \frac{2s^2}{C\sigma_0^2}n.$$
(5.32)

Indeed, since $i \in T_{\beta}$, by (5.8) and (5.4) we have

$$\frac{\beta}{|\widetilde{\varphi}\|_{\infty}} > \widetilde{r}_i \ge \frac{1}{2} \min\left(\frac{\varphi_i}{s^2}, \frac{1}{\widetilde{\varphi}_i}\right)$$

Since $\beta \leq 1/2$ it follows that the minimum is attained by the first argument. Thus

$$\frac{\beta}{\|\widetilde{\varphi}\|_{\infty}} > \frac{\varphi_i}{2s^2} > \frac{1}{2s^2} \frac{1}{n} \sum_{j=1}^n \sigma_{ij}^2 \widetilde{r}_j \ge \frac{\sigma_0^2}{2s^2} \frac{1}{n} \sum_{j \in \mathcal{N}_+(i)} \widetilde{r}_j.$$

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From Markov's inequality it follows that for any C > 0, $\tilde{r}_j < C\beta/\|\tilde{\varphi}\|_{\infty}$ for all but at most $(2s^2/C\sigma_0^2)n$ values of $j \in \mathcal{N}_+(i)$, which gives (5.32). Combining these estimates and taking $C = 8s^2/(\sigma_0^2\delta \varepsilon_0)$ we have

$$|T_{C\beta} \setminus T_{\beta}| \ge |\mathcal{N}_{+}(i) \cap T_{\beta}^{c}| - |\mathcal{N}_{+}(i) \cap T_{C\beta}^{c}| \ge \left(\delta \varepsilon_{0} - \frac{2s^{2}}{C\sigma_{0}^{2}}\right)n \ge (\delta \varepsilon_{0}/2)n$$

as desired.

Applying the above claim iteratively with $\varepsilon_0 = s^2/8$ we obtain $C'(\sigma_0, \delta, s) < \infty$ such that if β is sufficiently small depending on σ_0, δ, s and n is sufficiently large depending on κ, s , then

$$|T_{C'\beta}| \ge (1 - s^2/8)n. \tag{5.33}$$

 \square

Now let $C_0 > 0$ to be chosen later, and towards a contradiction suppose $T_{C_0\beta} \neq [n]$. Then there exists $i \in [n]$ such that $\tilde{r}_i \geq C_0\beta/\|\tilde{\varphi}\|_{\infty}$. From (5.4) we have the upper bound $\tilde{r}_i \leq \|\varphi\|_{\infty}/s^2$, so we conclude $\|\varphi\|_{\infty} \geq C_0s^2\beta/\|\tilde{\varphi}\|_{\infty}$. Now from our assumption $K \leq \frac{1}{n}\sum_{j=1}^n r_j = \sum_{j=1}^n \tilde{r}_j$, if K is sufficiently large depending on s then the same argument as in the proof of Lemma 5.3 shows that $\tilde{r}_j \geq \|\varphi\|_{\infty}/4$ for at least $(s^2/4)n$ values of $j \in [n]$. Thus, $\tilde{r}_j \geq C_0s^2\beta/(4\|\tilde{\varphi}\|_{\infty})$ for at least $(s^2/4)n$ values of $j \in [n]$, i.e. $|T_{C_0s^2\beta/4}| < (1 - \frac{s^2}{4})n$. Taking $C_0 = 4C'/s^2$ we contradict (5.33), and we conclude $T_{C_0\beta} = [n]$.

Now by the same lines as in (5.13)–(5.17) we conclude

$$\widetilde{r}_i \ge \beta_0 / \|\widetilde{\varphi}\|_{\infty} \quad \forall i \in [n]$$
(5.34)

for some $\beta_0(s, \sigma_0, \delta, \kappa) > 0$. Note we are now free to use the estimates in Lemma 5.4 with this value of β_0 .

5.3.2 Upper bounding r_i

Here our task is essentially to modify the proof of Lemma 5.5 to show we can take α' sufficiently small and *independent of* n such that $|S_{\alpha'} \setminus S_{\alpha}| \gtrsim n$, rather than merely nonempty. We can then conclude the proof by iterating this fact a bounded number of times.

Let us summarize the key new ideas. In the proof of Lemma 5.5 we used the irreducibility assumption to find an element $i^* \in S_{\alpha}$ such that the average of the components r_k over S_{α}^c was of order $\geq_{\alpha,n} r_{i^*} \geq \alpha \|\widetilde{\varphi}\|_{\infty}$ (see (5.30)). In a similar spirit, Lemma 5.8 below controls the average of r_k over $k \in S_{\alpha}^c$ from below by the average of r_i over $i \in U_0$, for a set $U_0 \subset S_{\alpha}$ that is *densely connected* to S_{α}^c . By averaging over a large set U_0 we are able to use the full strength of the bounds in Lemma 5.4 and avoid any dependence of the constants on n.

Proceeding naïvely, one can then use Lemma 5.8 to deduce $|S_{c_0\alpha^2} \setminus S_{\alpha}| \ge c_0\alpha |S_{\alpha}^c|$ for a sufficiently small constant $c_0 = c_0(s, \sigma_0, \delta, \kappa) > 0$. However, when iterating this bound over a sequence of values $\alpha_{k+1} = c_0\alpha_k^2$, the sets S_{α_k} grow by an exponentially decreasing proportion of n, so this is not enough to find a value of α for which $|S_{\alpha}|$ is close to n.

Instead, in Lemma 5.9 we are able to grow S_{α} by a constant factor using a nested iteration argument, which we now describe. We would like to find some value of $\alpha' \in (0, \alpha)$ for which

$$|S_{\alpha'} \setminus S_{\alpha}| \ge c|S_{\alpha'}| \tag{5.35}$$

where c > 0 is small constant. Suppose that (5.35) fails. By the expansion assumption, we know that $S_{\alpha'}$ contains a fairly large set $U = \mathcal{N}_{-}^{(\delta)}(S_{\alpha'}^c) \cap S_{\alpha'}$ (of size at least

 $\min(|S_{\alpha'}|, \kappa | S_{\alpha'}^c|))$ that is densely connected to $S_{\alpha'}^c$. In particular, if c is sufficiently small depending on κ , then U must have large overlap with S_{α} . Denoting the overlap by U_0 , Lemma 5.8 can now be applied to deduce

$$|S_{c_0\alpha\alpha'}\setminus S_{\alpha'}|\geq c_0\alpha|S_{\alpha'}^c|$$

for some $c_0 = c_0(s, \sigma_0, \delta, \kappa) > 0$ sufficiently small. The key is that the constant of proportionality on the right hand side is independent of α' . Hence, for fixed α , as long as (5.35) fails we can iteratively lower α' to increase $|S_{\alpha'} \setminus S_{\alpha}|$ by an amount $\gtrsim \alpha |S_{\alpha'}^c|$, until eventually (5.35) holds. This whole procedure can then be iterated a bounded number of times to obtain α'' such $|S_{\alpha''}|$ is close to n.

Having motivated the key ideas, we turn now to the proofs.

Lemma 5.8. Let $\alpha \in (0,1)$ and suppose that $0 < |S_{\alpha}| \le (1 - \delta/2)n$. If K is sufficiently large depending on $\alpha, s, \sigma_0, \delta, \kappa$, then for any $U_0 \subset \mathcal{N}^{(\delta)}_-(S^c_{\alpha}) \cap S_{\alpha}$ with $|U_0| \ge \frac{1}{10} |\mathcal{N}^{(\delta)}_-(S^c_{\alpha}) \cap S_{\alpha}|$ we have

$$\frac{1}{|S_{\alpha}^{c}|} \sum_{k \in S_{\alpha}^{c}} r_{k} \gtrsim \frac{\alpha}{|S_{\alpha}|} \sum_{i \in U_{0}} r_{i}.$$
(5.36)

Proof. First we prove the comparison

$$\sum_{i \in U_0} r_i \ge 2 \sum_{j \in S_\alpha} \widetilde{r}_j \tag{5.37}$$

assuming K is sufficiently large depending on $\alpha, s, \sigma_0, \delta, \kappa$. Indeed, if (5.37) does not hold, then by the fact that $U_0 \subset S_{\alpha}$ and (5.5),

$$\alpha \|\widetilde{\boldsymbol{\varphi}}\|_{\infty} |U_0| \leq \sum_{i \in U_0} r_i < 2 \sum_{j \in S_{\alpha}} \widetilde{r}_j \leq \frac{2|S_{\alpha}|}{s^2 \alpha \|\widetilde{\boldsymbol{\varphi}}\|_{\infty}}.$$

Rearranging we have

$$\|\widetilde{\boldsymbol{\varphi}}\|_{\infty} \leq \frac{1}{\alpha} \left(\frac{2|S_{\alpha}|}{s^{2}|U_{0}|}\right)^{1/2} \lesssim \frac{1}{\alpha} \left(\frac{|S_{\alpha}|}{\min(|S_{\alpha}|, |S_{\alpha}^{c}|)}\right)^{1/2} \lesssim 1/\alpha$$

where in the second bound we applied the robust irreducibility assumption and our assumed bounds on U_0 and S_{α} , and in the bound we used that both S_{α} and its complement are of linear size in *n*. From our assumption (5.10), (5.4) and the above it follows that

$$K \le \frac{1}{n} \sum_{i=1}^{n} r_i \le \|\widetilde{\varphi}\|_{\infty} / s^2 \lesssim 1/\alpha.$$
(5.38)

Taking K sufficiently large depending on $\alpha, s, \sigma_0, \delta$ and κ , we may assume (5.37) holds. From (5.20) and Lemma 5.4 we have

$$\sum_{i \in U_0} r_i = t \left(\sum_{j \in S_\alpha} \sum_{i \in U_0} W_{ij}^{S_\alpha} \right) + \frac{1}{n} \sum_{j \in S_\alpha} \sum_{k \in S_\alpha^c} \left(\sum_{i \in U_0} W_{ij}^{S_\alpha} \right) \sigma_{kj}^2 r_k$$
$$\leq \sum_{j \in S_\alpha} \widetilde{r}_j + \frac{1}{n} \sum_{j \in S_\alpha} \sum_{k \in S_\alpha^c} \left(\sum_{i \in U_0} W_{ij}^{S_\alpha} \right) \sigma_{kj}^2 r_k,$$

where in the bound we have applied (5.21) and (5.22). Applying (5.37) and rearranging yields

$$\sum_{i \in U_0} r_i \le \frac{2}{n} \sum_{j \in S_\alpha} \sum_{k \in S_\alpha^c} \left(\sum_{i \in U_0} W_{ij}^{S_\alpha} \right) \sigma_{kj}^2 r_k.$$
(5.39)

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Now since $U_0 \subset \mathcal{N}_{-}^{(\delta)}(S_{\alpha}^c)$, for any $i \in U_0$ we have $|\mathcal{N}_{+}(i) \cap S_{\alpha}^c| \geq \delta |S_{\alpha}^c|$. Together with (5.21) and (5.23) this implies

$$\delta|S_{\alpha}^{c}|\sum_{i\in U_{0}}W_{ij}^{S_{\alpha}}\leq \sum_{i\in S_{\alpha}}W_{ij}^{S_{\alpha}}|\mathcal{N}_{+}(i)\cap S_{\alpha}^{c}|\leq \left(\frac{n\|\widetilde{\varphi}\|_{\infty}}{\beta_{0}\sigma_{0}^{2}}\right)\widetilde{r}_{j}.$$

Rearranging we obtain a bound on $\sum_{i \in U_0} W_{ij}^{S_{\alpha}}$, which we substitute in (5.39) to obtain

$$\sum_{i \in U_0} r_i \leq \frac{2 \|\widetilde{\varphi}\|_{\infty}}{\beta_0 \sigma_0^2 \delta |S_{\alpha}^c|} \sum_{j \in S_{\alpha}} \sum_{k \in S_{\alpha}^c} \sigma_{kj}^2 \widetilde{r}_j r_k \leq \frac{2 |S_{\alpha}|}{s^2 \alpha \beta_0 \sigma_0^2 \delta |S_{\alpha}^c|} \sum_{k \in S_{\alpha}^c} r_k,$$

where in the second inequality we applied (5.5) to bound $\tilde{r}_j \leq 1/s^2 \alpha \|\tilde{\varphi}\|_{\infty}$ for all $j \in S_{\alpha}$. The result now follows by rearranging.

Lemma 5.9. For any $\alpha \in (0,1)$ there exists $\alpha' = \alpha'(\alpha, s, \sigma_0, \delta, \kappa) > 0$ such that either

$$|S_{\alpha'}| \ge (1 - \delta/2)n \tag{5.40}$$

or

$$|S_{\alpha'} \setminus S_{\alpha}| \ge \frac{1}{2} \min(|S_{\alpha'}|, \kappa | S_{\alpha'}^c|)$$
(5.41)

(or both).

Proof. For $\alpha' \in (0, \alpha)$ denote by $P(\alpha')$ the statement that at least one of (5.40) and (5.41) holds. We will show that while $P(\alpha')$ fails, we can lower α' by a controlled amount to increase the size of $S_{\alpha'} \setminus S_{\alpha}$ by a little bit. We can then iterate this until $P(\alpha')$ holds.

Let $\alpha' \in (0, \alpha)$ be arbitrary and assume $P(\alpha')$ fails. We claim there exists $c_0(s, \sigma_0, \delta, \kappa) > 0$ such that

$$\frac{1}{|S_{\alpha'}^c|} \sum_{k \in S_{\alpha'}^c} r_k \ge c_0 \alpha \alpha' \|\widetilde{\varphi}\|_{\infty}.$$
(5.42)

Put $U_0 = \mathcal{N}_{-}^{(\delta)}(S_{\alpha'}^c) \cap S_{\alpha}$. By the robust irreducibility assumption and the fact that (5.41) fails,

$$|U_0| \ge |\mathcal{N}_{-}^{(\delta)}(S_{\alpha'}^c) \cap S_{\alpha'}| - |S_{\alpha'} \setminus S_{\alpha}| \ge \frac{1}{2}|\mathcal{N}_{-}^{(\delta)}(S_{\alpha'}^c) \cap S_{\alpha'}|$$

$$(5.43)$$

$$\geq \frac{1}{2}\min(|S_{\alpha'}|,\kappa|S_{\alpha'}^{c}|).$$
(5.44)

By (5.43) and Lemma 5.8,

$$\frac{1}{|S_{\alpha'}^c|} \sum_{k \in S_{\alpha'}^c} r_k \gtrsim \frac{\alpha'}{|S_{\alpha'}|} \sum_{i \in U_0} r_i \ge \alpha \alpha' \|\widetilde{\varphi}\|_{\infty} \frac{|U_0|}{|S_{\alpha'}|} \gtrsim \alpha \alpha' \|\widetilde{\varphi}\|_{\infty},$$

where in the last inequality we applied (5.44) and the fact that (5.40) fails. This gives (5.42) as desired.

Now denoting

$$U' = \left\{ k \in S_{\alpha'}^c : r_k \ge \frac{1}{2} c_0 \alpha \alpha' \| \widetilde{\varphi} \|_{\infty} \right\}$$

we have

$$\sum_{k \in S_{\alpha'}^c} r_k \leq \sum_{k \in U'} r_k + \sum_{k \in S_{\alpha'}^c \setminus U'} r_k \leq \alpha' \|\widetilde{\varphi}\|_{\infty} |U'| + \frac{1}{2} c_0 \alpha \alpha' \|\widetilde{\varphi}\|_{\infty} |S_{\alpha'}^c|,$$

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where we used that by definition, $r_k \leq \alpha' \|\varphi\|_{\infty}$ for all $k \in S_{\alpha'}^c$. Combining with (5.42) and rearranging gives

$$|S_{c_0\alpha\alpha'/2} \setminus S_{\alpha'}| \ge |U'| \ge \frac{1}{2}c_0\alpha|S_{\alpha'}^c|.$$
(5.45)

Since (5.45) holds as long as $P(\alpha')$ fails, we can repeatedly lower α' by a factor $c_0\alpha/2$ to obtain $\alpha' = \alpha'(\alpha, s, \sigma_0, \delta, \varepsilon)$ such that $P(\alpha')$ holds. More explicitly, for each $k \ge 0$ put $\alpha_k = (c_0\alpha/2)^k \alpha$ and abbreviate $S_k := S_{\alpha_k}$. Then for all $k \ge 1$ such that $P(\alpha_k)$ fails we have $|S_{k+1} \setminus S_k| \ge \frac{1}{2}c_0\alpha|S_k^c|$, so

$$|S_{k+1} \setminus U_0| = |S_{k+1} \setminus S_k| + \dots + |S_1 \setminus U_0| \ge \frac{1}{2}c_0\alpha(|S_k^c| + \dots + |S_0^c|) \ge (k+1)\frac{1}{2}c_0\alpha|S_{k+1}^c|.$$

Thus, we must have that $P(\alpha_k)$ holds for some $k \leq 2\kappa/c_0\alpha$. (This gives α' of size $O(1/\alpha)^{-O(1/\alpha)}$.)

Now we conclude the proof of Proposition 5.1. From Lemma 5.3 we have $|S_{1/4}| \ge (s^2/4)n$. Applying Lemma 5.9 O(1) times we obtain $\alpha'' \gtrsim 1$ such that

$$|S_{\alpha''}| \ge (1 - \delta/2)n.$$
 (5.46)

Now from (5.5) we have

$$\widetilde{r}_j \le \frac{1}{s^2 \alpha'' \|\widetilde{\varphi}\|_{\infty}} \tag{5.47}$$

for all $j \in S_{\alpha''}$. On the other hand, for any $j \in S_{\alpha''}^c$,

$$1/\widetilde{r}_j \ge \widetilde{\varphi}_j \ge (V^{\mathsf{T}}r)_j \ge \frac{1}{n} \sum_{i \in S_{\alpha''}} \sigma_{ij}^2 r_i \ge \frac{1}{n} \sigma_0^2 \alpha'' \|\widetilde{\varphi}\|_{\infty} |\mathcal{N}_{-}(j) \cap S_{\alpha''}|.$$

From (5.46) and the robust irreducibility assumption (specifically the condition (2.13)),

$$|\mathcal{N}_{-}(j) \cap S_{\alpha''}| \ge \delta n - |S_{\alpha''}^c| \ge \delta n/2.$$

Combining the previous two displays we obtain

$$\widetilde{r}_j \le \frac{2}{\delta \sigma_0^2 \alpha'' \|\widetilde{\boldsymbol{\varphi}}\|_{\infty}}$$

for all $j \in S^c_{\alpha''}$. Together with (5.47) we have

$$\widetilde{r}_j \lesssim \frac{1}{\alpha'' \|\widetilde{\varphi}\|_{\infty}}$$

for all $j \in [n]$. Applying (5.3),

$$\alpha'' \|\widetilde{\varphi}\|_{\infty} n/2 \le \alpha'' \|\widetilde{\varphi}\|_{\infty} |S_{\alpha''}| \le \sum_{j=1}^n r_j = \sum_{j=1}^n \widetilde{r}_j \lesssim \frac{n}{\alpha'' \|\widetilde{\varphi}\|_{\infty}}$$

and rearranging gives $\|\widetilde{\varphi}\|_{\infty} \lesssim 1/\alpha'' \lesssim 1$. Finally, since

$$K \le \frac{1}{n} \sum_{j=1}^{n} r_j \le \|\widetilde{\varphi}\|_{\infty} / s^2$$

by (5.4), we obtain a contradiction if K is sufficiently large depending on s, σ_0, δ and κ . It follows that (5.10) fails for sufficiently large K, which concludes the proof of Proposition 5.1.

Remark 5.1. We note that in the above proof we only applied the expansion bound (2.14) to sets of size at least $\delta n/10$.

6 Proof of Theorem 2.3-(1): tail estimates and asymptotics of the logarithmic potential

The main purpose of this section is to show that the logarithmic potential $U_{\mu_n^Y}$ is close to $h_n(z) = -\int_{\mathbb{R}} \log |x| \check{\nu}_{n,z}(dx)$ for large *n*, and moreover that h_n is the logarithmic potential of a probability measure μ_n . Recall that in Theorem 3.2 we have already established the almost sure convergence of the truncated potentials:

$$\int_{\{|x|\geq \varepsilon\}} \log |x|\,\check{L}_{n,z}(dx) - \int_{\{|x|\geq \varepsilon\}} \log |x|\,\check{\nu}_{n,z}(dx) \xrightarrow[n\to\infty]{a.s.} 0 \ .$$

Thus, we need to show that these measures uniformly integrate the singularity of $x \mapsto \log |x|$ at 0. The proof has two main ingredients. The first is a result from [23] by the first author (stated in Proposition 6.1 below) that provides control on the smallest singular value of $Y_n - z$.

The second is the control of the remaining small singular values of $Y_n - z$ via the quantity $\mathbb{E} \operatorname{Im} g_{\tilde{L}_{n,z}}(it)$ when t is close to zero. We refer to step 2 in Section 2.5 for an outline of this argument.

Finally, to obtain the deterministic equivalents μ_n for the ESDs μ_n^Y we rely on a meta-model argument, which has been used before in [25, 49]. The idea is that for fixed *n* we can define a sequence $\{Y_n^{(m)}\}_{m\geq 1}$ of $nm \times nm$ random matrices as in Definition 1.2, where the standard deviation profiles $A_n^{(m)}$ are obtained by replacing each entry σ_{ij} of A_n by an $m \times m$ block with entries all equal to σ_{ij} . We can then show that the logarithmic potentials of the associated ESDs converge to h_n as $m \to \infty$, which will allow us to deduce that h_n is itself the logarithmic potential of a probability measure. This argument is described in more detail in Section 6.2 below.

6.1 Control on small singular values

The following result, obtained by one of the authors in [23], gives an estimate on the lower tail of the smallest singular value $s_{n,z}$ of $Y_n - z$.

Proposition 6.1 ([23], Theorem 1.19 and Corollary 1.22). Assume A0 and A1 hold, and fix $z \in \mathbb{C} \setminus \{0\}$. There exist constants $C(|z|, M_0, \sigma_{\max}), \alpha(\varepsilon), \beta(|z|, \varepsilon, M_0, \sigma_{\max}) > 0$ such that for all $n \geq 1$,

$$\mathbb{P}\left(s_{n,z} \le n^{-\beta}\right) \le C n^{-\alpha}.$$
(6.1)

Remark 6.1. Similar bounds have been obtained under stronger assumptions on the standard deviation profile. For instance, (6.1) follows from [19, Lemma A.1] if we additionally assume A3 (and in fact this result does not require A0). Further assuming that A_n is composed of a bounded number of blocks of equal size with constant entries, [5, Corollary 5.2] gives (6.1) with $\alpha > 0$ as large as we please (and $\beta = \beta(\alpha)$). An easy argument also gives (6.1) for arbitrary fixed $\alpha > 0$ and $\beta(\alpha)$ under **A3** and replacing **A0** with a bounded density assumption - see [19, Section 4.4]. For the case that the entries X_{ij} are real Gaussian variables and $A_n(\sigma_0)$ is (δ, κ) -broadly connected for some fixed $\sigma_0, \delta, \kappa \in (0,1)$ (see Definition 2.15), (6.1) holds with arbitrary lpha > 0 and eta = lpha + 1 by [54, Theorem 2.3].

We now consider the other small singular values of $Y_n - z$. The key is the uniform control on solutions to the Regularized Master Equations (2.7) provided by Assumption A2 combined with Theorem 3.3.

Corollary 6.2 (Wegner estimates). Let A0, A1 and A2 hold. Then, for all $z \in \mathbb{C} \setminus \{0\}$ there exist constants $C, \gamma_0 > 0$ such that for all x > 0,

$$\check{\nu}_{n,z}((-x,x)) \le Cx \tag{6.2}$$

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and

$$\mathbb{E}\check{L}_{n,z}((-x,x)) \le C(x \lor n^{-\gamma_0}).$$
(6.3)

Proof. We rely on the following elementary estimate for the Stieltjes transform of a probability measure μ (see for instance [34, Lemma 15]):

Im
$$g_{\mu}(it) = t \int \frac{\mu(d\lambda)}{\lambda^2 + t^2} \ge t \int_{-t}^{t} \frac{\mu(d\lambda)}{\lambda^2 + t^2} \ge \frac{1}{2t} \mu((-t,t)).$$
 (6.4)

Recall that $\text{Im } g_{\check{\nu}_{n,z}}(it) = n^{-1} \sum_{i \in [n]} r_i(|z|, t)$. The first Wegner estimate (6.2) is a straightforward consequence of Assumption **A2** and the estimate (6.4).

We now establish the second Wegner estimate (6.3) and first prove that there exists $\gamma_0>0$ such that

$$\sup_{t \ge n^{-\gamma_0}} \mathbb{E} \operatorname{Im} g_{\check{L}_{n,z}}(\mathrm{i}t) \le C$$
(6.5)

for all $n \ge 1$. For $t \ge 1$, $\mathbb{E} \operatorname{Im} g_{\tilde{L}_{n,z}}(\mathrm{i}t) \le 1$ by the mere definition of a Stieltjes transform. Assume t < 1 and recall that $\mathbb{E} \operatorname{Im} g_{\tilde{L}_{n,z}}(\mathrm{i}t) = n^{-1} \operatorname{tr} \operatorname{Im} \mathbb{E}G(z, \mathrm{i}t)$. By Theorem 3.3, there exist constants $c_0, C > 0$ such that

$$\left|\frac{1}{n} \mathrm{tr}\, \operatorname{Im} \mathbb{E} G(z, \mathrm{i} t) - \frac{1}{n} \sum_{i=1}^{n} r_i(|z|, t)\right| \leq \frac{C}{\sqrt{n} t^{c_0}} \,.$$

By A2, we therefore get that

$$\mathbb{E} \operatorname{Im} g_{\check{L}_{n}}(it) \leq C(t^{-c_0}n^{-1/2}+1).$$

By letting now $t \ge n^{-\gamma_0}$ with $\gamma_0 = 1/(2c_0)$, we obtain (6.5). Combining this result with (6.4), we get

$$\mathbb{E}\check{L}_{n,z}((-x,x)) \leq \mathbb{E}\check{L}_{n,z}((-(x \vee n^{-\gamma_0}), x \vee n^{-\gamma_0})) \leq 2C(x \vee n^{-\gamma_0})$$

which is the desired result.

6.2 Comparison of logarithmic potentials via a meta-model

We now turn to the task of finding the measures μ_n from Theorem 2.3 which serve as a sequence of deterministic equivalents for the ESDs μ_n^Y . A first idea is to try to show that for every $\psi \in \mathbb{C}^{\infty}_c(\mathbb{C})$,

$$\int \psi(z)\,\mu_n^Y(dz) = -\frac{1}{2\pi} \int \Delta \psi(z)\,U_{\mu_n^Y}(z)\,\ell(dz) = \frac{1}{2\pi} \int \Delta \psi(z) \Big(\int_{\mathbb{R}} \log |x|\,\check{L}_{n,z}(dx)\Big)\ell(dz)$$

is "close" to $-(2\pi)^{-1}\int \Delta\psi(z)h_n(z)dz$. However, there is a difficulty in directly applying this approach, related to the fact that $\check{L}_{n,z}$ does not converge in general with no further assumption on the variance profile matrices V_n .

To circumvent this difficulty, we rely on a *meta-model argument*, which has been used in [25, 49], and which we now describe. Let n be fixed, consider the standard deviation profile $A_n = (\sigma_{ij})$ and the normalized variance profile $V_n = (\frac{1}{n}\sigma_{ij}^2)$. Recall the associated

Schwinger–Dyson equations as provided in Proposition 3.1 and the solution $\vec{p} = \begin{pmatrix} p \\ \tilde{p} \end{pmatrix}$, of dimension $2n \times 1$. Define the meta-model in the following way: for an integer $m \ge 1$, consider the $nm \times nm$ standard deviation profile matrix defined as

$$\boldsymbol{A}_{n}^{(m)} = \begin{pmatrix} A_{n} & \cdots & A_{n} \\ \vdots & \vdots \\ A_{n} & \cdots & A_{n} \end{pmatrix} = (\boldsymbol{1}_{m} \boldsymbol{1}_{m}^{\mathsf{T}}) \otimes A_{n} ,$$

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associated to the normalized variance profile $V_n^{(m)} = (\mathbf{1}_m \mathbf{1}_m^{\mathsf{T}}) \otimes m^{-1} V_n$, and the random matrix

$$\boldsymbol{Y}_{n}^{(m)} = \left(\frac{[\boldsymbol{A}_{n}^{(m)}]_{ij}}{\sqrt{mn}} X_{ij}^{(nm)}\right)_{i,j \in [nm]}.$$
(6.6)

Denote by $\check{\mathbf{L}}_{n,z}^{(m)}$ the symmetrized empirical distribution of the singular values of $\mathbf{Y}_n^{(m)} - zI_{mn}$. Due to the specific form of $\mathbf{V}_n^{(m)}$, it is straightforward to check that the solutions of the Schwinger–Dyson equations associated to this model are provided by

$$ec{p}_m = \left(egin{array}{c} p_m \ ilde{p}_m \end{array}
ight) \qquad ext{where} \qquad p_m^\mathsf{T} = \left(p^\mathsf{T}, \cdots, p^\mathsf{T}
ight) \quad ext{and} \quad ilde{p}_m^\mathsf{T} = \left(ilde{p}^\mathsf{T}, \cdots, ilde{p}^\mathsf{T}
ight) \;,$$

where ${m p}_m$ and ${m ilde p}_m$ are nm imes 1 vectors. As an important consequence, we have:

$$g_{\check{\nu}_{n,z}^{(m)}}(\eta) = \frac{1}{mn} \sum_{i=1}^{mn} [\vec{p}_m]_i = \frac{1}{n} \sum_{i=1}^n [\vec{p}]_i = g_{\check{\nu}_{n,z}}(\eta).$$

Hence the Stieltjes transform $g_{\check{\nu}_{n,z}^{(m)}}$ of $\check{\nu}_{n,z}^{(m)}$ does not depend on m and is equal to $\check{\nu}_{n,z}$. Finally, if **A0** and **A1** are satisfied for Y_n , they are also satisfied for $Y_n^{(m)}$. In particular, $\check{L}_{n,z}^{(m)} \sim \check{\nu}_{n,z}^{(m)}$ admits a genuine limit as $m \to \infty$:

$$\check{\boldsymbol{L}}_{n,z}^{(m)} \xrightarrow[m \to \infty]{} \check{\boldsymbol{\nu}}_{n,z}$$
 a.s. (6.7)

since $\check{\nu}_{n,z}^{(m)} = \check{\nu}_{n,z}$.

We now state our proposition giving the existence of the measures μ_n , the proof of which will occupy the main part of the remainder of this section.

Proposition 6.3. Let A0, A1 and A2 hold. Then the following hold:

1. For all $n \ge 1$ and $z \in C \setminus \{0\}$, the function

$$h_n(z) = -\int_{\mathbb{R}} \log |x| \,\check{\nu}_{n,z}(dx)$$

is well defined and for every compact set $\mathcal{K}\subset\mathbb{C}$,

$$\sup_{n} \int_{\mathcal{K}} |h_n(z)|^2 \,\ell(dz) < \infty \,.$$

Moreover, $h_n(z)$ coincides with the logarithmic potential $U_{\mu_n}(z)$ of a probability measure μ_n on \mathbb{C} .

2. For μ_n as defined in part (1), there exists a constant C > 0, independent of n, such that for all M > 0,

$$\mu_n(\{z \in \mathbb{C}; |z| > M\}) \le \frac{C}{M^2}.$$

We will rely on the following lemmas, whose proofs are respectively deferred to Appendices A.5 and A.6.

Lemma 6.1. Let $(\Omega, \mathcal{F}, \mathbb{P})$ be a given probability space, ζ a finite positive measure on \mathbb{C} and $f_n : \Omega \times \mathbb{C} \to \mathbb{R}$ measurable functions satisfying

$$\sup_{n} \int_{\mathbb{C}} |f_{n}(\omega, z)|^{1+\alpha} \mathbb{P} \otimes \zeta(d\omega \times dz) \leq C$$
(6.8)

for some constants $\alpha, C > 0$.

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Let $g : \mathbb{C} \to \mathbb{R}$ be a measurable function such that for ζ -almost all $z \in \mathbb{C}$,

$$f_n(\omega, z) \xrightarrow[n \to \infty]{\mathcal{P}} g(z)$$

Then $\int_{\mathbb{C}} |g(z)|^{1+\alpha} \zeta(dz) \leq C$, and

$$\int_{\mathbb{C}} f_n(\omega, z) \,\zeta(dz) \xrightarrow[n \to \infty]{\mathcal{P}} \int g(z) \,\zeta(dz).$$
(6.9)

Lemma 6.2. Let (ζ_n) be a sequence of random probability measures on \mathbb{C} . Assume that a.s. (ζ_n) is tight and that there exists a locally integrable function $h : \mathbb{C} \to \mathbb{R}$ such that for all $\psi \in C_c^{\infty}(\mathbb{C})$,

$$\int \psi(z)\,\zeta_n(dz) \xrightarrow[n \to \infty]{\mathcal{P}} -\frac{1}{2\pi} \int \Delta\psi(z)h(z)\,\ell(dz) \;. \tag{6.10}$$

Then there exists a non-random probability measure ζ on \mathbb{C} with logarithmic potential h, i.e.

$$h(z) = -\int_{\mathbb{C}} \log|z - u| \,\zeta(\,du)$$

for almost all $z \in \mathbb{C}$ such that $\zeta_n \xrightarrow[n \to \infty]{\mathbf{w}} \zeta$ in probability.

Proof of Proposition 6.3. We prove the first point of the proposition.

With $\boldsymbol{Y}_n^{(m)}$ as in (6.6), denote by $\mu_{n,m}^{\boldsymbol{Y}}$ the spectral measure of $\boldsymbol{Y}_n^{(m)}$, and recall that

$$U_{\mu_{n,m}^{\mathbf{Y}}}(z) = -\int_{\mathbb{R}} \log |x| \,\check{\boldsymbol{L}}_{n,z}^{(m)}(dx).$$

The proof consists of the following three steps:

1. To show that for every $z \in \mathbb{C} \setminus \{0\}$, $x \mapsto \log |x|$ is $\check{\nu}_{n,z}$ -integrable and

$$U_{\mu_{n,m}}(z) \xrightarrow[m \to \infty]{\mathcal{P}} h_n(z).$$
(6.11)

- 2. To show that the function $h_n(z)$ is measurable.
- 3. To show that for any compact set $\mathcal{K} \subset \mathbb{C}$,

$$\sup_{m} \mathbb{E} \int_{\mathcal{K}} |U_{\mu_{n,m}^{\mathbf{Y}}}(z)|^{2} \ell(dz) \le C$$
(6.12)

for some constant C > 0 independent of n.

The three previous steps being proved, the assumptions of Lemma 6.1 are fulfilled and the lemma yields

$$\int_{\mathcal{K}} |h_n(z)|^2 \,\ell(dz) \leq C$$

where C does not depend upon n. Moreover,

$$\int \psi(z) \,\mu_{n,m}^{\boldsymbol{Y}}(dz) \;=\; -\frac{1}{2\pi} \int \Delta \psi(z) \,U_{\mu_{n,m}^{\boldsymbol{Y}}}(z) \,\ell(dz) \xrightarrow{\mathcal{P}} -\frac{1}{2\pi} \int \Delta \psi(z) \,h_n(z) \,\ell(dz)$$

for every $\psi \in \mathbb{C}^{\infty}_{c}(\mathbb{C})$.

It remains to apply Lemma 6.2 to conclude that h_n is the logarithmic potential of a probability distribution μ_n on \mathbb{C} and point (1) of Proposition 6.3 will be proved.

Let us address Step 1, to prove the convergence in (6.11), we separately consider the integrals defining the logarithmic potentials in the regions |x| greater than and less than ε .

Taking into account the convergence (6.7) and applying Theorem 3.2 to the sequence $(\check{L}_{n,z}^{(m)})_m$, we deduce that $x \mapsto \log |x|$ is $\check{\nu}_{n,z}$ -integrable near infinity, and that

$$\int_{\{|x|\geq\varepsilon\}} \log|x|\,\check{\boldsymbol{L}}_{n,z}^{(m)}(dx) \xrightarrow[m\to\infty]{} \int_{\{|x|\geq\varepsilon\}} \log|x|\,\check{\nu}_{n,z}(dx) \,. \tag{6.13}$$

We now handle the remaining regions.

$$\begin{aligned} \left| \int_{\{|x|<\varepsilon\}} \log |x| \,\check{\nu}_{n,z}(dx) \right| &\leq \int_{\{|x|<\varepsilon\}} |\log |x|| \,\check{\nu}_{n,z}(dx) \\ &= \int_{\mathbb{R}} \check{\nu}_{n,z} \left\{ |x|<\varepsilon, \log |x|\leq -y \right\} dy \stackrel{(a)}{\leq} C \int_{0}^{\infty} (\exp(-y)\wedge\varepsilon) \,dy = C\varepsilon(1-\log\varepsilon) \,, \end{aligned}$$

$$(6.14)$$

where (a) follows from Wegner's estimate (6.2) in Corollary 6.2.

.

Let $(s_{i,z})_{i\in[mn]}$ be the singular values of $Y_n^{(m)} - z$ ordered as $s_{1,z} \ge \cdots \ge s_{mn,z}$. For $z \in \mathbb{C} \setminus \{0\}$ and $\beta > 0$ the exponent as in Proposition 6.1, we introduce the event

$$\mathcal{G}_m := \left\{ \boldsymbol{s}_{i,z} \ge (mn)^{-\beta} \,, \; i \in [mn] \right\} \;.$$

For all $\tau > 0$,

$$\mathbb{P}\left\{ \left| \int_{\{|x|<\varepsilon\}} \log|x|\,\check{\boldsymbol{L}}_{n,z}^{(m)}(dx) \right| > \tau \right\}$$

$$\leq \mathbb{P}\left\{ \left| \mathbbm{1}_{\mathcal{G}_m} \int_{\{|x|<\varepsilon\}} \log|x|\,\check{\boldsymbol{L}}_{n,z}^{(m)}(dx) \right| > \frac{\tau}{2} \right\} + \mathbb{P}\left\{ \left| \mathbbm{1}_{\mathcal{G}_m^c} \int_{\{|x|<\varepsilon\}} \log|x|\,\check{\boldsymbol{L}}_{n,z}^{(m)}(dx) \right| > \frac{\tau}{2} \right\} .$$

Noticing that

$$\left\{ \left| \mathbb{1}_{\mathcal{G}_m^c} \int_{\{|x| < \varepsilon\}} \log |x| \,\check{\boldsymbol{L}}_{n,z}^{(m)}(dx) \right| > \frac{\tau}{2} \right\} \subset \left\{ \boldsymbol{s}_{mn,z} \le (mn)^{-\beta} \right\}$$

for m large enough, Proposition 6.1 yields that

$$\mathbb{P}\left\{\left|\mathbbm{1}_{\mathcal{G}_m^c}\int_{\{|x|<\varepsilon\}}\log|x|\,\check{\boldsymbol{L}}_{n,z}^{(m)}(dx)\right|>\frac{\tau}{2}\right\} \leq \frac{C}{(mn)^{\alpha}}\;.$$

Recall the constant γ_0 in Corollary 6.2. Choose m large enough and $\gamma \leq \gamma_0$ small enough so that $(nm)^{-\beta} \leq (mn)^{-\gamma} \leq \varepsilon \leq 1$. We now estimate

$$\begin{split} \mathbb{E} \mathbb{1}_{\mathcal{G}_{m}} \int_{\{|x|<\varepsilon\}} |\log|x| \, |\, \check{\boldsymbol{L}}_{n,z}^{(m)}(dx) \\ &= \int_{\{(mn)^{-\beta} \le |x| \le (mn)^{-\gamma}\}} |\log|x| \, |\, \mathbb{E} \check{\boldsymbol{L}}_{n,z}^{(m)}(dx) + \int_{\{(mn)^{-\gamma} < |x|<\varepsilon\}} |\log|x| \, |\, \mathbb{E} \check{\boldsymbol{L}}_{n,z}^{(m)}(dx) \\ &:= I_{1} + I_{2} \; . \end{split}$$

By the Wegner estimate (6.3), we obtain

$$I_1 \leq \beta \log(mn) \mathbb{E} \check{\boldsymbol{L}}_{n,z}^{(m)}([-(mn)^{-\gamma}, (mn)^{-\gamma}]) \leq C\beta(mn)^{-\gamma} \log(mn).$$

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On the other hand, another application of the same Wegner estimate yields

$$I_{2} = \int_{0}^{\infty} \mathbb{E}\check{\boldsymbol{L}}_{n,z}^{(m)} \left(\left\{ x : |\log|x| |\mathbb{1}_{[(mn)^{-\gamma},\varepsilon]}(|x|) \ge y \right\} \right) dy$$

$$= \int_{0}^{\infty} \mathbb{E}\check{\boldsymbol{L}}_{n,z}^{(m)} \left([-e^{-y} \wedge \varepsilon, -(mn)^{-\gamma}] \cup [(mn)^{-\gamma}, e^{-y} \wedge \varepsilon] \right) dy$$

$$\leq C \int_{0}^{\infty} (\exp(-y) \wedge \varepsilon) dy = C\varepsilon (1 - \log \varepsilon) .$$

Therefore, by Markov's inequality, we finally obtain

$$\mathbb{P}\left\{ \left| \int_{\{|x|<\varepsilon\}} \log |x| \,\check{\boldsymbol{L}}_{n,z}^{(m)}(dx) \right| > \tau \right\} \leq \frac{1}{\tau} 2C \big[\beta(mn)^{-\gamma} \log(mn) + \varepsilon (1 - \log \varepsilon) \big] + C(mn)^{-\alpha}.$$

Thus, for all $\tau, \tau' > 0$, we can choose $\varepsilon > 0$ small enough so that

$$\mathbb{P}\bigg\{ \bigg| \int_{\{|x|<\varepsilon\}} \log |x| \,\check{\boldsymbol{L}}_{n,z}^{(m)}(dx) \bigg| > \tau \bigg\} < \tau'$$

for m large enough. Gathering this result with (6.13) and (6.14) yields (6.11), and Step 1 is proved.

We now address Step 2 and study the measurability of $h_n(z)$. Recall the regularized logarithmic potentials defined in (2.24) and consider also the following function:

$$\mathcal{U}_{n,m}^{\boldsymbol{Y}}(z,t) := -\frac{1}{2nm} \log \det((\boldsymbol{Y}_n^{(m)} - z)^* (\boldsymbol{Y}_n^{(m)} - z) + t^2) = -\frac{1}{2} \int_{\mathbb{R}} \log(x^2 + t^2) \, \check{\boldsymbol{L}}_{n,z}^{(m)}(dx) \, .$$

Given z and $z' \in \mathbb{C}$, Hoffman-Wielandt's theorem applied to $m{Y}_n^{(m)} - z$ and $m{Y}_n^{(m)} - z'$ yields

$$\max_{i \in [mn]} |s_{i,z} - s_{i,z'}| \le |z - z'| .$$

Thus

$$\begin{aligned} \left| \mathcal{U}_{n,m}^{\boldsymbol{Y}}(z,t) - \mathcal{U}_{n,m}^{\boldsymbol{Y}}(z',t) \right| \\ &= \left| \frac{1}{2nm} \left| \sum_{i \in [mn]} \log \left(1 + \frac{\boldsymbol{s}_{i,z}^2}{t^2} \right) - \log \left(1 + \frac{\boldsymbol{s}_{i,z'}^2}{t^2} \right) \right| \le \left| \frac{1}{2t^2} \max_{i \in [mn]} |\boldsymbol{s}_{i,z} - \boldsymbol{s}_{i,z'}| \right| \le \left| \frac{|z - z'|}{2t^2} \right| \end{aligned}$$

and it follows that for any fixed t > 0 the family $\{z \mapsto \mathcal{U}_{n,m}^{\mathbf{Y}}(z,t)\}_{m \ge 1}$ is uniformly equicontinuous. Since from Theorem 3.2 we have $\mathcal{U}_{n,m}^{\mathbf{Y}}(z,t) \xrightarrow[m \to \infty]{} \mathcal{U}_n(z,t)$ almost surely, it follows that $z \mapsto \mathcal{U}_n(z,t)$ is continuous for any fixed t > 0. Finally, since $x \mapsto \log |x|$ is $\check{\nu}_{n,z}$ -integrable near zero for any $z \neq 0$ by (6.14),

$$\mathcal{U}_n(z,t) \xrightarrow[t \to 0]{} h_n(z).$$

The measurability of h_n follows and Step 2 is proved.

We now address Step 3 and prove (6.12). Observe that on any compact set $\mathcal{K} \in \mathbb{C}$, there exists a constant $C_{\mathcal{K}}$ such that

$$\int_{\mathcal{K}} (\log |\lambda - z|)^2 \, \ell(dz) \; \leq \; C_{\mathcal{K}} (1 + |\lambda|^2)$$

for all $\lambda \in \mathbb{C}$. Denote by $(\boldsymbol{\lambda}_i; i \in [mn])$ the eigenvalues of $\boldsymbol{Y}_n^{(m)}$. We have

$$\mathbb{E} \int_{\mathcal{K}} |U_{\mu_{n,m}^{\mathbf{Y}}}(z)|^{2} \ell(dz)$$

$$\leq \mathbb{E} \left(\frac{1}{mn} \sum_{i \in [mn]} \int_{\mathcal{K}} (\log |\boldsymbol{\lambda}_{i}^{(m)} - z|)^{2} \ell(dz) \right) \leq C_{\mathcal{K}} \left(1 + \mathbb{E} \int |\lambda|^{2} \mu_{m,n}^{\mathbf{Y}}(d\lambda) \right).$$

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By the Weyl comparison inequality for eigenvalues and singular values (cf. e.g. [41, Theorem 3.3.13]),

$$\begin{split} \int |\lambda|^2 \mu_{n,m}^{\boldsymbol{Y}}(d\lambda) &= \frac{1}{nm} \sum_{i=1}^{nm} |\boldsymbol{\lambda}_i|^2 \\ &\leq \frac{1}{nm} \sum_{i=1}^{nm} \boldsymbol{s}_{i,0}^2 = \frac{1}{nm} \operatorname{tr} \left(\boldsymbol{Y}_n^{(m)} (\boldsymbol{Y}_n^{(m)})^* \right) \leq \frac{\sigma_{\max}^2}{(nm)^2} \sum_{i,j=1}^{nm} |X_{ij}^{(nm)}|^2. \end{split}$$

Taking the expectation of the previous inequality finally yields

$$\mathbb{E}\int_{\mathcal{K}}|U_{\mu_{n,m}^{\mathbf{Y}}}(z)|^{2}\,\ell(dz)\leq C.$$

Step 3 is proved.

We now prove point (2) of Proposition 6.3. Given M > 0, we get from Lemma 6.1 that

$$\limsup_m \mu_{n,m}^{\boldsymbol{Y}}(\{z\in\mathbb{C};\;|z|>M\})\leq\limsup_m \frac{1}{M^2}\int_{\mathbb{C}}|\lambda|^2\,\mu_{n,m}^{\boldsymbol{Y}}(d\lambda)\leq \frac{C}{M^2}\quad\text{a.s.}$$

where C > 0 is independent of n. Let ψ be a nonnegative $C_c^{\infty}(\mathbb{C})$ function equal to one for |z| < M and to zero if |z| > M + 1. As a byproduct of Lemma 6.2,

$$\mu_{n,m}^{\boldsymbol{Y}} \xrightarrow[m \to \infty]{w} \mu_n$$

almost surely. Consequently, on a set of probability one,

$$\mu_n(\{z \in \mathbb{C}; |z| \le M+1\}) \ge \int \psi(z) \,\mu_n(dz) = \lim_m \int \psi(z) \,\mu_{n,m}^{\mathbf{Y}}(dz) \ge 1 - \frac{C}{M^2}.$$
position 6.3 is proved.

Prop

6.3 Conclusion of the proof of Theorem 2.3-(i)

We can now complete the proof of Theorem 2.3–(i) and prove that $\mu_n^Y \sim \mu_n$ in probability, with μ_n defined in Proposition 6.3.

By Proposition 6.3, the sequence (μ_n) is tight. It remains to prove that for all $\varphi \in C_c(\mathbb{C})$, $\int \varphi d\mu_n^Y - \int \varphi d\mu_n \to 0$ in probability. By the density of $C_c^{\infty}(\mathbb{C})$ in $C_c(\mathbb{C})$, it is enough to show that

$$\int \psi(z)\mu_n^Y(dz) - \int \psi(z)\mu_n(dz) = -\frac{1}{2\pi} \int \Delta\psi(z)(U_{\mu_n^Y}(z) - U_{\mu_n}(z))\,\ell(dz) \xrightarrow[n \to \infty]{\mathcal{P}} 0$$

for all $\psi \in C^\infty_c(\mathbb{C})$. By mimicking the proof of Proposition 6.3, where $m{Y}^{(m)}_n$ and m are replaced with Y_n and n respectively, we straightforwardly obtain that $U_{\mu_n^Y}(z) - U_{\mu_n}(z) \to 0$ in probability for every $z \in \mathbb{C} \setminus \{0\}$. This proof also shows that $\sup_n \mathbb{E} \int_{\mathcal{K}} |U_{\mu_n^Y}(z)|^2 \ell(dz) < \infty$ for all compact sets $\mathcal{K} \subset \mathbb{C}$. We also know by Proposition 6.3 that $\sup_n \int_{\mathcal{K}} |U_{\mu_n}(z)|^2 \ell(dz) < 1$ ∞ . The result now follows from Lemma 6.1.

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7.1 **Proof of Theorem 2.3: Identification of** μ_n

We established in the previous section that $\mu_n^Y \sim \mu_n$ in probability. To conclude the proof of Theorem 2.3, it remains to show that μ_n is rotationally invariant, and that its radial cumulative distribution function

$$\mu_n\{z \in \mathbb{C} \; ; \; |z| \le r\}$$

coincides with the function F_n specified in the statement of the theorem. These facts, along with the properties of F_n , are established in Lemma 7.2 below.

For the remainder of this section, we set

$$b_n(z,t) := -\frac{z}{2n} \operatorname{tr} \Psi(\vec{r}(|z|,t),t)$$
 and $b_n(z) = -\frac{z}{2n} \operatorname{tr} \Psi(\vec{q}(|z|))$, (7.1)

where $\Psi(\cdot, t)$ and $\Psi(\cdot)$ are defined in (4.1), $\vec{r}(\cdot, t)$ is defined in Proposition 2.1, and $\vec{q}(\cdot)$ is defined in 2.2.

Lemma 7.1. Under the same assumptions as in Theorem 2.3, the function $z \mapsto b_n(z)$ is locally integrable on \mathbb{C} , and

$$\partial_{\bar{z}} U_{\mu_n}(z) = b_n(z)$$

in $\mathcal{D}'(\mathbb{C})$.

Proof. Recall the definition of $U_n(z,t)$ in (2.24). Recall also that by Proposition 6.3-(i), the probability measure μ_n is such that:

$$U_{\mu_n}(z) = -\int_{\mathbb{R}} \log |x| \check{\nu}_{n,z}(dx) \; .$$

We first prove that

$$\mathcal{U}_n(z,t) \xrightarrow[t\downarrow 0]{\mathcal{D}'(\mathbb{C})} U_{\mu_n}(z) .$$
(7.2)

Recall from the proof of Proposition 6.3 that $z \mapsto U_n(z,t)$ is continuous for any fixed t > 0. It is moreover clear from the expressions of $U_n(z,t)$ and $U_{\mu_n}(z)$ that $U_n(z,t) \uparrow U_{\mu_n}(z)$ as $t \downarrow 0$. Recall that $U_{\mu_n}(z)$, being a logarithmic potential, is locally integrable (as can be seen by Fubini's theorem). Thus, given a fixed $t_0 > 0$,

$$0 \leq \mathcal{U}_n(t,z) - \mathcal{U}_n(t_0,z) \leq U_{\mu_n}(z) - \mathcal{U}_n(t_0,z),$$

for $0 < t \le t_0$, and (7.2) immediately follows from the monotone convergence theorem. By a property of the convergence in $\mathcal{D}'(\mathbb{C})$, this implies the convergence of the distributional derivative

$$\partial_{\bar{z}} \mathcal{U}_n(z,t) \xrightarrow[t\downarrow 0]{\mathcal{D}'(\mathbb{C})} \partial_{\bar{z}} U_{\mu_n}(z).$$
 (7.3)

We now prove that for all t > 0,

$$\partial_{\bar{z}}\mathcal{U}_n(z,t) = b_n(z,t) \tag{7.4}$$

in $\mathcal{D}'(\mathbb{C})$. We shall rely on a meta-model argument. Recall the meta-model $Y_n^{(m)}$ introduced in (6.6), its limiting property (6.7), and the definition of $\mathcal{U}_{n,m}^{Y}(z,t)$ in (2.24).

Fix t > 0. By Theorem 3.2, $\mathcal{U}_{n,m}^{\boldsymbol{Y}}(z,t) \to \mathcal{U}_n(z,t)$ almost surely as $m \to \infty$ for all $z \in \mathbb{C}$. Furthermore, recalling the notation $(s_{i,z})_{i \in [mn]}$ for the singular values of $\boldsymbol{Y}_n^{(m)} - z$, we have

$$\left|\mathcal{U}_{n,m}^{\boldsymbol{Y}}(z,t)\right|^{2} = \left|\log(t) + \frac{1}{2mn}\sum \log\left(1 + \boldsymbol{s}_{i,z}^{2}/t^{2}\right)\right|^{2} \stackrel{(a)}{\leq} 2\left|\log(t)\right|^{2} + \frac{1}{t^{2}mn}\sum_{i\in[mn]}\boldsymbol{s}_{i,z}^{2}$$

where (a) follows from the elementary inequality $2^{-1} \log^2(1+x) \le x$, valid for $x \ge 0$. In particular, this implies that

$$\mathbb{E}\int_{\mathcal{K}} |\mathcal{U}_{n,m}^{\boldsymbol{Y}}(z,t)|^2 \,\ell(dz) \leq (\log t)^2 + \frac{1}{t^2} \mathbb{E}\int_{\mathcal{K}} \frac{\operatorname{tr}\,(\boldsymbol{Y}_n^{(m)} - z)^*(\boldsymbol{Y}_n^{(m)} - z)}{mn} \,\ell(dz) \leq C$$

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on every compact set $\mathcal{K} \subset \mathbb{C}$. By Lemma 6.1, we get that $\mathcal{U}_n(\cdot, t)$ is locally integrable on \mathbb{C} , and that

$$\int \partial_{\bar{z}} \psi(z) \mathcal{U}_{n,m}^{\boldsymbol{Y}}(z,t) \,\ell(dz) \xrightarrow[m \to \infty]{\mathcal{P}} \int \partial_{\bar{z}} \psi(z) \,\mathcal{U}_n(z,t) \,\ell(dz) \tag{7.5}$$

for all $\psi \in C_c^{\infty}(\mathbb{C})$. An integration by parts along with Jacobi's formula shows that for all $\omega \in \Omega$, the distributional derivative $\partial_{\bar{z}} \mathcal{U}_{n,m}^{Y}(t,z)$ coincides with the pointwise derivative, which is given by

$$\partial_{\bar{z}} \mathcal{U}_{n,m}^{\boldsymbol{Y}}(z,t) = \frac{1}{2nm} \operatorname{tr} (\boldsymbol{Y}_{n}^{(m)} - z) ((\boldsymbol{Y}_{n}^{(m)} - z)^{*} (\boldsymbol{Y}_{n}^{(m)} - z) + t^{2})^{-1}$$

On the other hand, we know from Theorem 3.3 that $\partial_{\bar{z}} \mathcal{U}_{n,m}^{\boldsymbol{Y}}(z,t) \to b_n(z,t)$ almost surely as $m \to \infty$, for all $z \in \mathbb{C}$. Moreover, from a singular value decomposition of $\boldsymbol{Y}_n^{(m)} - z$ we easily see that $|\partial_{z^*} \mathcal{U}_{n,m}^{\boldsymbol{Y}}(z,t)| \leq (4t)^{-1}$. Consequently, we get by Lemma 6.1 again that

$$\int \partial_{\bar{z}} \psi(z) \, \mathcal{U}_{n,m}^{\boldsymbol{Y}}(t,z) \, \ell(dz) = -\int \psi(z) \, \partial_{\bar{z}} \mathcal{U}_{n,m}^{\boldsymbol{Y}}(t,z) \, \ell(dz) \xrightarrow{\mathcal{P}} -\int \psi(z) \, b_n(z,t) \, \ell(dz).$$

Comparing with (7.5), we obtain that $\partial_{\bar{z}}\mathcal{U}_n(z,t) = b_n(z,t)$ in $\mathcal{D}'(\mathbb{C})$.

We now consider the limit in $t \downarrow 0$ in (7.4). Since $|b_n(z,t)| \le |2z|^{-1}$, the dominated convergence theorem yields

$$b_n(z,t) \xrightarrow{\mathcal{D}'(\mathbb{C})} b_n(z)$$

Combining this convergence together with (7.3) and (7.4), we obtain the desired result. $\hfill\square$

In order to characterize the probability measure μ_n , we use the equation $\mu_n = -(2\pi)^{-1}\Delta U_{\mu_n}$ and rely on the smoothness properties of ΔU_{μ_n} that can be deduced from Lemma 4.4. We recall that $\vec{q}(s)$ is defined in the statement of Theorem 2.2.

Lemma 7.2. The probability measure μ_n is rotationally invariant. On $(0, \infty)$, the distribution function $F_n(s) := \mu_n(\{z : |z| \le s\})$ satisfies

$$F_n(s) = 1 - \frac{1}{n} \langle \boldsymbol{q}(s), V \widetilde{\boldsymbol{q}}(s) \rangle.$$

The support of μ_n is contained in $\{z : |z| \le \sqrt{\rho(V)}\}$. Finally, F_n is absolutely continuous on $(0, \infty)$, and has a continuous density on $(0, \sqrt{\rho(V)})$.

Before entering the proof, we note that the rotational invariance of μ_n can be "guessed" from the form of the Schwinger-Dyson equations of Proposition 3.1. Indeed, from this one sees that the Stieltjes transform $g_{\nu_{n,z}}(\eta) = n^{-1} \operatorname{tr} P(|z|, \eta)$ of $\nu_{n,z}$ depends on z only through its absolute value, and this is therefore also the case for $U_{\mu_n}(z)$. It is easy to check that this yields the rotational invariance of μ_n .

Proof. First, we show that $\mu_n(C) = 0$, where C is the circle with center zero and radius $\sqrt{\rho(V)}$.

Consider a smooth function $\phi : \mathbb{R} \to [0,1]$ with support in [-1,1] and value $\phi(0) = 1$, and the function

$$g_{\varepsilon}(z) = \phi\left(\varepsilon^{-1}\left(|z| - \sqrt{\rho(V)}\right)\right), \quad z \in \mathbb{C}$$

with support in the annulus $\{z : \sqrt{\rho(V)} - \varepsilon \le |z| \le \sqrt{\rho(V)} + \varepsilon\}$. We have

$$\int g_{\varepsilon}(z)\,\mu_n(dz) = -\frac{1}{2\pi}\int g_{\varepsilon}(z)\,\Delta U_{\mu_n}(z)\,\ell(dz) = \frac{2}{\pi}\int \partial_z g_{\varepsilon}(z)\,b_n(z)\,\ell(dz)$$

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where b_n is defined in (7.1). Notice that

$$\lim_{\varepsilon \downarrow 0} \int g_{\varepsilon}(z) \,\mu_n(dz) = \mu_n(C). \tag{7.6}$$

By replacing $|z|=\sqrt{x^2+y^2}$ and computing $\partial_z=rac{1}{2}(\partial_x-\mathrm{i}\partial_y)$, we get

$$\partial_z g_{\varepsilon}(z) = \frac{\overline{z}}{2\varepsilon |z|} \phi' \left(\varepsilon^{-1} \left(|z| - \sqrt{\rho(V)} \right) \right).$$

Hence, replacing b_n by its expression in (7.1), we obtain

$$\frac{2}{\pi} \int \partial_z g_{\varepsilon}(z) \, b_n(z) \, \ell(dz) = -\frac{1}{2\varepsilon n\pi} \int |z| \phi' \left(\varepsilon^{-1} \left(|z| - \sqrt{\rho(V)} \right) \right) \operatorname{tr} \Psi(\vec{q}(|z|) \, \ell(dz))$$

$$\stackrel{(a)}{=} -\frac{1}{2\varepsilon n\pi} \int_{\theta=0}^{2\pi} \int_{\rho=\sqrt{\rho(V)}-\varepsilon}^{\sqrt{\rho(V)}+\varepsilon} \phi' \left(\varepsilon^{-1} \left(|z| - \sqrt{\rho(V)} \right) \right) \operatorname{tr} \Psi(\vec{q}(\rho)) \, \rho^2 \, d\rho \, d\theta$$

$$\stackrel{(b)}{=} -\frac{1}{\varepsilon n} \int_{-1}^{1} \left(\sqrt{\rho(V)} + \varepsilon u \right)^2 \phi'(u) \operatorname{tr} \Psi(\vec{q}(\sqrt{\rho(V)} + \varepsilon u)) \varepsilon \, du$$

where (a) follows from a change of variables in polar coordinates and (b), from the change of variable $u = \frac{\rho - \sqrt{\rho(V)}}{\varepsilon}$. Since $n^{-1} \text{tr} \Psi(\vec{q}(|z|)) \leq |z|^{-2}$, the dominated convergence theorem yields

$$\begin{split} &-\frac{1}{n}\int_{-1}^{1}\left(\sqrt{\rho(V)}+\varepsilon u\right)^{2}\phi'(u)\mathrm{tr}\,\Psi(\vec{q}(\sqrt{\rho(V)}+\varepsilon u))\,du\\ &\xrightarrow[\varepsilon\downarrow 0]{} -\frac{\rho(V)}{n}\mathrm{tr}\,\Psi(\vec{q}(\sqrt{\rho(V)}))\int_{-1}^{1}\phi'(u)\,du\\ &=-\frac{\rho(V)}{n}\mathrm{tr}\,\Psi(\vec{q}(\sqrt{\rho(V)}))\left[\phi(1)-\phi(-1)\right]\ =\ 0\ .\end{split}$$

Equating with (7.6), we finally conclude $\mu_n(C) = 0$.

By Theorem 2.2–(3), the mapping $z \mapsto \vec{q}(|z|)$ is continuously differentiable on the open set $\mathcal{D} := \{z \in \mathbb{C} : |z| \neq 0, |z| \neq \rho(V)^{1/2}\}$. Therefore, $b_n(z)$ is continuously differentiable on this set, and for any $g \in C_c^{\infty}(\mathcal{D})$, we get

$$\begin{split} \int_{\mathbb{C}} g(z) \,\mu_n(dz) &= -\frac{1}{2\pi} \int_{\mathbb{C}} g(z) \,\Delta U_{\mu_n}(z) \,\ell(dz) \,= \, -\frac{2}{\pi} \int_{\mathbb{C}} g(z) \,\partial_z b_n(z) \,\ell(dz) \\ &= \, \int_{\mathbb{C}} g(z) \,f_n(z) \,\ell(dz) \end{split}$$

where the density $f_n(z)$ is given by

$$\begin{split} f_n(z) &:= -\frac{2}{\pi} \partial_z b_n(z) = \frac{1}{n\pi} \partial_z \left(z \operatorname{tr} \Psi(\vec{q}(|z|)) \right) = \frac{1}{n\pi} \left\{ \operatorname{tr} \Psi(\vec{q}(|z|)) + |z|^2 \partial_{|z|^2} \operatorname{tr} \Psi(\vec{q}(|z|)) \right\} \\ &= \frac{1}{n\pi} \partial_{|z|^2} \left\{ |z|^2 \operatorname{tr} \Psi(\vec{q}(|z|)) \right\} = \frac{-1}{\pi n} \sum_{i=1}^n \partial_{|z|^2} \left(\frac{\varphi_i \tilde{\varphi}_i}{|z|^2 + \varphi_i \tilde{\varphi}_i} \right) \\ &= \frac{-1}{\pi n} \partial_{|z|^2} \langle q(|z|), V \tilde{q}(|z|) \rangle. \end{split}$$

Since $f_n(z)$ depends only on |z|, this density is rotationally invariant. From Theorem 2.2-(1), $f_n(z) = 0$ for $|z| > \sqrt{\rho(V)}$. Thus, the support of μ_n is contained in $\overline{\mathcal{B}(0,\sqrt{\rho(V)})}$. Moreover, $F_n(\sqrt{\rho(V)}) = 1 = \lim_{v \uparrow \sqrt{\rho(V)}} F_n(v)$, since $\mu_n(C) = 0$. Given

 $0 < s \leq v < \sqrt{\rho(V)}$, we have

$$F_n(v) - F_n(s) = \int_{\mathcal{B}(0,v) \setminus \mathcal{B}(0,s)} f_n(z) \,\ell(dz) = \int_0^{2\pi} d\theta \int_s^v \frac{-1}{2\pi r n} \partial_r \langle \boldsymbol{q}(r), V \widetilde{\boldsymbol{q}}(r) \rangle \, r \, dr$$
$$= \frac{1}{n} \langle \boldsymbol{q}(s), V \widetilde{\boldsymbol{q}}(s) \rangle - \frac{1}{n} \langle \boldsymbol{q}(v), V \widetilde{\boldsymbol{q}}(v) \rangle.$$

By taking $v \uparrow \sqrt{\rho(V)}$, $\langle q(v), V \tilde{q}(v) \rangle \to 0$, and we get the expression (2.9). Finally, the continuity of the density of F_n on $(0, \sqrt{\rho(V)})$ follows from Theorem 2.2–(3).

7.2 Proof of Theorem 2.4

In Example 2.2, it has been proved that Theorem 2.4 holds under the additional assumption that the matrices A_n are irreducible. Now for the general case, by conjugating Y_n by a permutation matrix we may assume A_n takes the form

$$A_{n} = \begin{pmatrix} A_{n}^{(1)} & \cdots & 0\\ \vdots & \ddots & \vdots\\ 0 & \cdots & A_{n}^{(m)} \end{pmatrix}$$
(7.7)

where $A_n^{(1)}, \ldots, A_n^{(m)}$ are square irreducible matrices of respective dimension $n_1 \ge \cdots \ge n_m$. Indeed, for A_n a general nonnegative matrix we can achieve this with the upper triangular blocks not necessarily zero, but these are forced to be zero by the stochasticity condition. Also, by **A1** and the row-sum constraint applied to the last row of $A_n = (\sigma_{ij})$, $1 = \frac{1}{n} \sum_{j=1}^n \sigma_{ij}^2 \le \frac{n_m}{n} \sigma_{\max}^2$ so in fact we have

$$n_1, \dots, n_m \ge n/\sigma_{\max}^2. \tag{7.8}$$

Denote the corresponding submatrices of X_n by $X_n^{(k)}$ and set

$$Y_n^{(k)} = \frac{1}{\sqrt{n}} A_n^{(k)} \odot X_n^{(k)} = \frac{1}{\sqrt{n_k}} B_n^{(k)} \odot X_n^{(k)}$$
(7.9)

where we set $B_n^{(k)} = (n_k/n)^{1/2} A_n^{(k)}$. For each k we have:

- 1. $A_n^{(k)}$ is irreducible,
- 2. $\frac{1}{n}B_n^{(k)} \odot B_n^{(k)}$ is doubly stochastic, and
- 3. $n_k \to \infty$ as $n \to \infty$ (by (7.8)).

Thus, for each k the ESD $\mu_n^{(k)}$ of $Y_n^{(k)}$ converges weakly in probability to $\mu_{\rm circ}$. Since the μ_n^Y is the weighted sum:

$$\mu_n^Y = (n_1/n)\mu_n^{(1)} + \dots + (n_m/n)\mu_n^{(m)}$$

we get that μ_n^Y converges weakly in probability to $\mu_{\rm circ}$. This concludes the proof of Theorem 2.4.

A Remaining proofs

A.1 Stieltjes transform of a symmetric probability measure

We note that a symmetric probability distribution $\check{\nu}$ on \mathbb{R} satisfies $\check{\nu}(A) = \check{\nu}(-A)$ for each Borel set $A \subset \mathbb{R}$.

Lemma A.1. A probability measure $\check{\nu}$ is symmetric if and only if its Stieltjes transform $g_{\check{\nu}}$, seen as an analytic function on $\mathbb{C} \setminus \mathbb{R}$, satisfies $g_{\check{\nu}}(-\eta) = -g_{\check{\nu}}(\eta)$.

Proof. The necessity is obvious from the definition of the Stieltjes transform and from the fact that $\check{\nu}(d\lambda) = \check{\nu}(-d\lambda)$. To prove the sufficiency, we use the Perron inversion formula, that says that for any function $\varphi \in C_c(\mathbb{R})$,

$$\int_{\mathbb{R}} \varphi(x) \,\check{\nu}(dx) = \lim_{\varepsilon \downarrow 0} \frac{1}{\pi} \int_{\mathbb{R}} \varphi(x) \,\operatorname{Im} g_{\check{\nu}}(x + \mathrm{i}\varepsilon) \, dx.$$

By a simple variable change at the right hand side, and by using the equalities $g_{\check{\nu}}(-\eta) = -g_{\check{\nu}}(\eta)$ and $g_{\check{\nu}}(\bar{\eta}) = \bar{g}_{\check{\nu}}(\eta)$, we obtain that $\int \varphi(x) \check{\nu}(dx) = \int \varphi(-x) \check{\nu}(dx)$, as desired. \Box

A.2 Variance estimates

In this section we collect without proofs a number of standard variance estimates. Let (Y_n) be a sequence of matrices as in Definition 1.2. In the sequel we drop the subscript n. Denote by $(\vec{e_i})$ the standard vector basis. We introduce the following notations:

$$Y = (\vec{y}_1, \cdots, \vec{y}_n) \text{ and } Q(\eta^2) = \left[\sum_{i=2}^n (\vec{y}_i - z\vec{e}_i)(\vec{y}_i - z\vec{e}_i)^* - \eta^2\right]^{-1}$$

Recall the definition of matrices \boldsymbol{R} and \boldsymbol{G} in (2.19).

Proposition A.1. Let A0 and A1 hold. Let Δ be a $n \times n$ deterministic diagonal matrix, then the following estimates hold:

$$\operatorname{var}(\boldsymbol{R}_{ij}) = \mathcal{O}_{\eta}\left(n^{-1}\right) \text{ for } 1 \leq i, j \leq 2n , \qquad (A.1)$$

$$\operatorname{var}\left(\frac{1}{n}\operatorname{tr}\Delta G\right) = \mathcal{O}_{\eta}\left(\|\Delta\|^{2} n^{-2}\right) , \qquad (A.2)$$

$$\operatorname{var}\left[(\vec{y}_{1} - z\vec{e}_{1})^{*}[\eta Q]^{\alpha}(\vec{y}_{1} - z\vec{e}_{1})\right] = \mathcal{O}_{\eta}\left(n^{-1}\right) \text{ for } \alpha = 1, 2.$$
 (A.3)

Similar estimates hold true if G is replaced by \tilde{G} , if one considers the columns of Y^* instead of those of Y, etc.

These estimates can be obtained as in the proof of [49, Proposition 6.3], see also the references therein.

As a direct corollary of the previous proposition, we have:

Corollary A.2. Let A0 and A1 hold.

$$\operatorname{var}\left[\left(\eta + \frac{1}{n}\operatorname{tr}\Delta G\right)^{-1}\right] = \mathcal{O}_{\eta}\left(\|\Delta\|^{2} n^{-2}\right) , \qquad (A.4)$$

$$\operatorname{var}\left[\left(\eta + (\vec{y}_1 - z\vec{e}_1)^*[\eta Q](\vec{y}_1 - z\vec{e}_1)\right)^{-1}\right] = \mathcal{O}_\eta\left(n^{-1}\right) \,. \tag{A.5}$$

Proof. Let us establish (A.4). Notice first that $\left|\eta + \frac{1}{n} \operatorname{tr} \Delta G\right|^{-1} \leq \operatorname{Im}^{-1}(\eta)$ by (2.2).

$$\operatorname{var} \left(\frac{1}{\eta + \frac{1}{n} \operatorname{tr} \Delta G} \right) \stackrel{(a)}{\leq} \mathbb{E} \left| \frac{1}{\eta + \frac{1}{n} \operatorname{tr} \Delta G} - \frac{1}{\mathbb{E} \left(\eta + \frac{1}{n} \operatorname{tr} \Delta G \right)} \right|^{2}, \\ \leq \frac{1}{\operatorname{Im}^{4}(\eta)} \operatorname{var} \left(\frac{1}{n} \operatorname{tr} \Delta G \right) = \mathcal{O}_{\eta} \left(\frac{\|\Delta\|^{2}}{n^{2}} \right),$$

where (a) follows from the fact that $var(X) = inf_a \mathbb{E}|X - a|^2$. Estimate (A.5) can be established similarly.

A.3 Proof of Proposition 3.5

Notice first that it is sufficient to prove

$$\frac{1}{n} \operatorname{tr} \mathbb{E} \boldsymbol{C} \left(\boldsymbol{R}(z, \eta) - \boldsymbol{R}^{\mathcal{N}}(z, \eta) \right) = \vec{\mathcal{O}}_{\eta} \left(\frac{1}{\sqrt{n}} \right) \,,$$

where C is any of the following $2n \times 2n$ matrices:

$$\begin{pmatrix} I_n & 0 \\ 0 & 0 \end{pmatrix}, \begin{pmatrix} 0 & 0 \\ I_n & 0 \end{pmatrix}, \begin{pmatrix} 0 & I_n \\ 0 & 0 \end{pmatrix}, \begin{pmatrix} 0 & 0 \\ 0 & I_n \end{pmatrix}.$$

We follow the approach used in [59] and perform an entry-by-entry interpolation between Y and $Y^{\mathcal{N}}$. Let $\alpha : [n^2] \to [n]^2$ be the bijection defined as $\alpha(m) = (i, j)$, where m - 1 = (i-1)n + j - 1 is the Euclidean division of m - 1 by n. Given $m \in [n^2]$, define the $n \times n$ matrices $Z^m = (Z^m_{i,j})$ and $W^m = (W^m_{i,j})$ by

$$Z_{i,j}^m := \left\{ \begin{array}{ll} Y_{i,j}^{\mathcal{N}} & \text{if } \alpha^{-1}(i,j) < m, \\ Y_{i,j} & \text{if } \alpha^{-1}(i,j) \geq m, \end{array} \right. \quad \text{and} \quad W_{i,j}^m := \left\{ \begin{array}{ll} Y_{i,j}^{\mathcal{N}} & \text{if } \alpha^{-1}(i,j) < m, \\ 0 & \text{if } \alpha^{-1}(i,j) = m, \\ Y_{i,j} & \text{if } \alpha^{-1}(i,j) > m. \end{array} \right.$$

Notice that $Z^1 = Y$. By convention, we denote $Z^{n^2+1} = Y^N$. Redenoting the resolvent $\mathbf{R}(z,\eta)$ defined in (2.19) as \mathbf{R}_Y to express the dependence on Y (thus, $\mathbf{R}^N = \mathbf{R}_{Y^N}$), we have

$$\frac{1}{n}\operatorname{tr} \mathbb{E} \boldsymbol{C} \left(\boldsymbol{R}_{Y} - \boldsymbol{R}_{Y^{\mathcal{N}}}\right) = \frac{1}{n} \sum_{m=1}^{n^{2}} \operatorname{tr} \mathbb{E} \boldsymbol{C} \left(\boldsymbol{R}_{Z^{m}} - \boldsymbol{R}_{Z^{m+1}}\right)$$
$$= \frac{1}{n} \sum_{m=1}^{n^{2}} \operatorname{tr} \mathbb{E} \boldsymbol{C} \left(\boldsymbol{R}_{Z^{m}} - \boldsymbol{R}_{W^{m}}\right) - \operatorname{tr} \mathbb{E} \boldsymbol{C} \left(\boldsymbol{R}_{Z^{m+1}} - \boldsymbol{R}_{W^{m}}\right). \quad (A.6)$$

For $(i, j) = \alpha(m)$, the matrices $\Delta_m := Z^m - W^m$ and $\Delta_m^N := Z^{m+1} - W^m$ are given by the equations

$$\mathbf{\Delta}_m = \begin{bmatrix} 0 & Y_{ij} \boldsymbol{e}_i \boldsymbol{e}_j^* \\ \bar{Y}_{ij} \boldsymbol{e}_j \boldsymbol{e}_i^* & 0 \end{bmatrix} \quad \text{and} \quad \mathbf{\Delta}_m^{\mathcal{N}} = \begin{bmatrix} 0 & Y_{ij}^{\mathcal{N}} \boldsymbol{e}_i \boldsymbol{e}_j^* \\ \bar{Y}_{ij}^{\mathcal{N}} \boldsymbol{e}_j \boldsymbol{e}_i^* & 0 \end{bmatrix},$$

where e_i is the i^{th} canonical vector of \mathbb{C}^n . These matrices are (at most) rank-two matrices that are both independent of W^m . We now use the identity $\mathbf{R}_{Z^m} = \mathbf{R}_{W^m} - \mathbf{R}_{Z^m} \Delta_m \mathbf{R}_{W^m}$ three times to obtain

$$\mathbb{E}C(\boldsymbol{R}_{Z^m}-\boldsymbol{R}_{W^m})=-\mathbb{E}C\boldsymbol{R}_{W^m}\boldsymbol{\Delta}_m\boldsymbol{R}_{W^m}+\mathbb{E}\boldsymbol{R}_{W^m}(\boldsymbol{\Delta}_m\boldsymbol{R}_{W^m})^2-\mathbb{E}\boldsymbol{R}_{Z^m}(\boldsymbol{\Delta}_m\boldsymbol{R}_{W^m})^3.$$

Similarly,

$$\mathbb{E}C(\boldsymbol{R}_{Z^{m+1}}-\boldsymbol{R}_{W^m})=-\mathbb{E}C\boldsymbol{R}_{W^m}\boldsymbol{\Delta}_m^{\mathcal{N}}\boldsymbol{R}_{W^m}+\mathbb{E}\boldsymbol{R}_{W^m}(\boldsymbol{\Delta}_m^{\mathcal{N}}\boldsymbol{R}_{W^m})^2-\mathbb{E}\boldsymbol{R}_{Z^m}(\boldsymbol{\Delta}_m^{\mathcal{N}}\boldsymbol{R}_{W^m})^3.$$

The first two terms at the right hand sides of these equations are identical since the first two moments of Y_{ij} and $Y_{ij}^{\mathcal{N}}$ are equal, thus, they cancel out in (A.6). Moreover, it is easy to see from the general properties of the resolvents and from the expressions of Δ_m and $\Delta_m^{\mathcal{N}}$ that the traces of the terms with the cubes are bounded by $Cn^{-3/2}/\eta^4$. The result follows.

A.4 Proof of estimate (3.24)

This estimate relies on the following standard proposition, whose proof can be found in [37, Lemma 5.2].

Proposition A.3. Let $C \geq 0$ be a $n \times n$ matrix and $\boldsymbol{u} = (u_{\ell}) \succ 0$ and $\boldsymbol{v} = (v_{\ell}) \succ 0$ two $n \times 1$ vector. Assume that the following equality holds true: $\boldsymbol{u} = C\boldsymbol{u} + \boldsymbol{v}$. Then $\rho(C) < 1$, matrix I - C is invertible, $(I - C)^{-1} \geq 0$ and

$$\left\| \| (I-C)^{-1} \| \right\|_{\infty} \leq \frac{\max(u_{\ell}; \ell \in [n])}{\min(v_{\ell}; \ell \in [n])}.$$

Recall the definitions (3.13) and (3.20) of $\Upsilon(\vec{b})$ and $\Delta(b)$, $\widetilde{\Delta}(\tilde{b})$. In order to study the properties of matrix $\mathcal{A}(p) \odot \mathcal{A}(p)$, we introduce an auxiliary system. Since the p_i 's satisfy (2.20), we immediately obtain

$$\operatorname{Im}(p_{i}) = \frac{\left[V^{\mathsf{T}}\operatorname{Im}(\boldsymbol{p})\right]_{i} |z|^{2}}{||z|^{2} - (\eta + [V\tilde{\boldsymbol{p}}]_{i})(\eta + [V^{\mathsf{T}}\boldsymbol{p}]_{i})|^{2}} + \frac{\left[V\operatorname{Im}(\tilde{\boldsymbol{p}})\right]_{i}}{|-(\eta + [V\tilde{\boldsymbol{p}}]_{i}) + \frac{|z|^{2}}{\eta + [V^{\mathsf{T}}\boldsymbol{p}]_{i}}|^{2}} + \frac{\operatorname{Im}(\eta)}{|-(\eta + [V\tilde{\boldsymbol{p}}]_{i}) + \frac{|z|^{2}}{\eta + [V^{\mathsf{T}}\boldsymbol{p}]_{i}}|^{2}} \left(\frac{|z|^{2}}{|\eta + [V^{\mathsf{T}}\boldsymbol{p}]_{i})|^{2}} + 1\right)$$

and its counterpart for $\text{Im}(\tilde{p}_i)$. Denote by $\boldsymbol{v}(\boldsymbol{\vec{p}})$ the $2n \times 1$ vector defined by

$$\left[\boldsymbol{v}(\vec{\boldsymbol{p}})\right]_{i} = \frac{\mathrm{Im}(\eta)}{\left|-(\eta + [V\tilde{\boldsymbol{p}}]_{i}) + \frac{|\boldsymbol{z}|^{2}}{\eta + [V^{\mathsf{T}}\boldsymbol{p}]_{i}}\right|^{2}} \left(\frac{|\boldsymbol{z}|^{2}}{\left|\eta + [V^{\mathsf{T}}\boldsymbol{p}]_{i}\right|^{2}} + 1\right)$$

for $i \in [n]$ and

$$\left[\boldsymbol{v}(\boldsymbol{\vec{p}})\right]_{i} = \frac{\mathrm{Im}(\eta)}{\left|-(\eta + [V^{\mathsf{T}}\boldsymbol{p}]_{i}) + \frac{|\boldsymbol{z}|^{2}}{\eta + [V\boldsymbol{\tilde{p}}]_{i}}\right|^{2}} \left(\frac{|\boldsymbol{z}|^{2}}{|\eta + [V\boldsymbol{\tilde{p}}]_{i})|^{2}} + 1\right)$$

for $i \in \{n + 1, \dots, 2n\}$. Then the system satisfied by $\text{Im}(\vec{p})$ writes

$$\operatorname{Im}(\vec{p}) = \mathcal{A}(\vec{p}) \odot \overline{\mathcal{A}(\vec{p})} \operatorname{Im}(\vec{p}) + \boldsymbol{v}(\vec{p}) , \qquad (A.7)$$

where matrix $\mathcal{A}(\vec{p})$ has been defined in (3.22). Since matrix $\mathcal{A}(\vec{p}) \odot \overline{\mathcal{A}(\vec{p})}$ has nonnegative entries, we will rely on Proposition A.3 to evaluate $\left\| \left(I - \mathcal{A}(\vec{p}) \odot \overline{\mathcal{A}(\vec{p})} \right)^{-1} \right\|_{\infty}$. We need to check that $\operatorname{Im}(\vec{p}), v(\vec{p}) \succ 0$, to upper bound $\operatorname{Im}(p_{\ell}), \operatorname{Im}(\tilde{p}_{\ell})$ and to lower bound $\left[v(\vec{\theta}) \right]_{i}$. Since p_{ℓ} and \tilde{p}_{ℓ} are Stieltjes transform, we have $|\operatorname{Im}(p_{\ell})| \lor |\operatorname{Im}(\tilde{p}_{\ell})| \le \operatorname{Im}(\eta)^{-1}$ which we write in short $\operatorname{Im}(\vec{p}) \prec \operatorname{Im}(\eta)^{-1}$. Now, if $i \le n$

$$\left[\boldsymbol{v}(\vec{\boldsymbol{p}})\right]_i \geq \frac{\mathrm{Im}(\eta)}{\left|-(\eta + [V\tilde{\boldsymbol{p}}]_i) + \frac{|\boldsymbol{z}|^2}{\eta + [V^{\mathsf{T}}\boldsymbol{p}]_i}\right|^2}$$

with

$$\frac{1}{\operatorname{Im}(\eta)} \left| -(\eta + [V\tilde{\boldsymbol{p}}]_i) + \frac{|z|^2}{\eta + [V^{\mathsf{T}}\boldsymbol{p}]_i} \right|^2 \leq 2 \left(\frac{|\eta|^2}{\operatorname{Im}(\eta)} + \frac{\sigma_{\max}^4}{\operatorname{Im}^3(\eta)} + \frac{|z|^4}{\operatorname{Im}^3(\eta)} \right) = \mathcal{O}_{\eta}(1) .$$

The case where $n+1 \leq i \leq 2n$ being handled similarly, we finally get

$$\min_{i} \left[\boldsymbol{v}(\vec{\boldsymbol{\theta}}) \right]_{i} \geq \frac{1}{\mathcal{O}_{\eta}\left(1\right)} . \tag{A.8}$$

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Hence, $\boldsymbol{v}(\vec{\boldsymbol{\theta}}) \succ 0$ and $\left[\boldsymbol{v}(\vec{\boldsymbol{\theta}})\right]_i$ is lower bounded away from zero.

In order to prove $\vec{p} \succ 0$, we argue as follows: the p_i 's are Stieltjes transforms of probability measures μ_i . These probability measures are tight, see for instance Proposition 3.1-(v). In particular, there exists a real number K such that $\mu_i([-K, K]) \ge \frac{1}{2}$. Hence,

$$\operatorname{Im}(p_i) = \operatorname{Im}(\eta) \int_{\mathbb{R}} \frac{\mu_i(d\lambda)}{|\lambda - \eta|^2} \ge \operatorname{Im}(\eta) \int_{-K}^{K} \frac{\mu_i(d\lambda)}{|\lambda - \eta|^2} \ge \frac{\mu_i([-K, K])}{2(K^2 + |\eta|^2)} \ge \frac{1}{4(K^2 + |\eta|^2)} .$$

We are now in position to apply Proposition A.3. This proposition yields in particular that $\rho(\mathcal{A}(\vec{p}) \odot \overline{\mathcal{A}(\vec{p})}) < 1$ and gathering the estimates $\operatorname{Im}(\vec{p}) \prec \operatorname{Im}(\eta)^{-1}$ and (A.8), we obtain

$$\left\| \left(I - \mathcal{A}(\vec{p}) \odot \overline{\mathcal{A}(\vec{p})} \right)^{-1} \right\|_{\infty} = \mathcal{O}_{\eta} (1) .$$
(A.9)

We can now conclude. Notice that along the imaginary axis, p(it) = ir(t) and $\tilde{p}(it) = i\tilde{r}(t)$ by Proposition 3.1–(3). Hence, straightforward computations yield that

$$I - \mathcal{A}(\vec{p}) \odot \mathcal{A}(\vec{p}) = \begin{pmatrix} A & B \\ C & D \end{pmatrix}$$
 and $I - \mathcal{A}(\vec{p}) \odot \overline{\mathcal{A}(\vec{p})} = \begin{pmatrix} A & -B \\ -C & D \end{pmatrix}$.

By standard block inversion formulas, the same structure occurs for the inverses

$$(I - \mathcal{A}(\vec{p}) \odot \mathcal{A}(\vec{p}))^{-1} = \begin{pmatrix} \widetilde{A} & \widetilde{B} \\ \widetilde{C} & \widetilde{D} \end{pmatrix}$$
 and $(I - \mathcal{A}(\vec{p}) \odot \overline{\mathcal{A}(\vec{p})})^{-1} = \begin{pmatrix} \widetilde{A} & -\widetilde{B} \\ -\widetilde{C} & \widetilde{D} \end{pmatrix}$,

from which we immediately deduce that

$$\left\| \left(I - \mathcal{A}(\vec{p}) \odot \overline{\mathcal{A}(\vec{p})} \right)^{-1} \right\|_{\infty} = \left\| \left(I - \mathcal{A}(\vec{p}) \odot \mathcal{A}(\vec{p}) \right)^{-1} \right\|_{\infty}$$

by either optimizing over $\vec{u} = (u, \tilde{u})$ with $\|\vec{u}\|_{\infty} = 1$ or $(u, -\tilde{u})$. Estimate (3.24) is established.

A.5 Proof of Lemma 6.1

Assume without loss of generality that ζ is a probability measure. Let $\varphi \in C_c(\mathbb{C})$. Then $\mathbb{E} \varphi(f_n(\cdot, z) - g(z)) \xrightarrow[n \to \infty]{} \varphi(0)$ for ζ -almost all $z \in \mathbb{C}$ since the convergence in probability induces the convergence in distribution. Thus, by the dominated convergence and Fubini's theorems,

$$\int_{\Omega \times \mathbb{C}} \varphi(f_n(\omega, z) - g(z)) \left(\mathbb{P} \otimes \zeta \right) (d\omega \times dz) \xrightarrow[n \to \infty]{} \varphi(0) \,.$$

In other words, $f_n - g$ converges to 0 in distribution, hence in probability, for the probability measure $\mathbb{P} \otimes \zeta$. As a consequence (see for instance [44, Lemma 3.11]),

$$\int_{\mathbb{C}} |g(z)|^{1+\alpha} \zeta(dz) = \int_{\Omega \times \mathbb{C}} |g(z)|^{1+\alpha} (\mathbb{P} \otimes \zeta) (d\omega \times dz) \leq C .$$

By (6.8), the sequence (f_n) is $\mathbb{P} \otimes \zeta$ -uniformly integrable, hence $\int_{\Omega \times \mathbb{C}} |f_n(\omega, z) - g(z)| (\mathbb{P} \otimes \zeta) (d\omega \times dz) \xrightarrow[n \to \infty]{} 0$, see for instance [44, Proposition 3.12]. Convergence (6.9) follows from Markov's inequality.

A.6 Proof of Lemma 6.2

Let $(\psi_k)_{k\geq 1}$ be a sequence of smooth compactly supported functions, dense in $C_c(\mathbb{C})$ for the supremum norm $\|\psi\|_{\infty} = \sup_{z\in\mathbb{C}} |\psi(z)|$. By the diagonal extraction procedure, one can find a subsequence $(\zeta_{n'})$ such that with probability one $(\zeta_{n'})$ is tight and $\int \psi_k d\zeta_{n'} \xrightarrow[n'\to\infty]{} -\frac{1}{2\pi} \int \Delta \psi_k(z)h(z)\ell(dz)$ for all $k\geq 1$. Thus, on this set of probability one, the tight sequence $(\zeta_{n'})$ has a unique non-random limit point ζ , and this limit point satisfies $\zeta = -(2\pi)^{-1}\Delta h$ in $\mathcal{D}'(\mathbb{C})$, the set of Schwartz distributions. With this at hand, we get from the assumption that

$$\int \psi_k(z)\,\zeta_n(dz) \xrightarrow[n\to\infty]{\mathcal{P}} \int \psi_k(z)\,\zeta(dz)$$

for all $k \ge 1$. By a density argument, we thus get that

$$\int \varphi(z)\,\zeta_n(dz) \xrightarrow[n\to\infty]{\mathcal{P}} \int \varphi(z)\,\zeta(dz)$$

for every $\varphi \in C_c(\mathbb{C})$.

References

- [1] Y. Ahmadian, F. Fumarola, and K. D. Miller, *Properties of networks with partially structured and partially random connectivity*, Physical Review E **91** (2015), no. 1, 012820. MR-3416669
- [2] O. Ajanki, L. Erdős, and T. Krüger, Universality for general wigner-type matrices, Probab. Theory Relat. Fields, 169 (2017), no. 3–4, 667–727. MR-3719056
- [3] O. Ajanki, L. Erdős, and T. Krüger, Quadratic vector equations on complex upper half-plane, To be published in Memoirs of the A.M.S., arXiv:1506.05095.
- [4] O. Ajanki, L. Erdős, and T. Krüger, Stability of the matrix dyson equation and random matrices with correlations, Probability Theory and Related Fields (2018), 1–81.
- [5] J. Aljadeff, D. Renfrew, and M. Stern, Eigenvalues of block structured asymmetric random matrices, J. Math. Phys. 56 (2015), no. 10, 103502, 14. MR-3403052
- [6] J. Aljadeff, M. Stern, and T. Sharpee, *Transition to chaos in random networks with cell-type-specific connectivity*, Phys. Rev. Lett. **114** (2015), 088101.
- [7] S. Allesina, J. Grilli, G. Barabás, S. Tang, J. Aljadeff, and A. Maritan, Predicting the stability of large structured food webs, Nature communications 6 (2015).
- [8] S. Allesina and S. Tang, *The stability-complexity relationship at age 40: a random matrix perspective*, Population Ecology **57** (2015), no. 1, 63–75.
- [9] S. Allesina and S. Tang, Stability criteria for complex ecosystems, Nature 483 (2012), no. 7388, 205–208.
- [10] J. Alt, L. Erdős, and T. Krüger, Local law for random Gram matrices, Electron. J. Probab. 22 (2017), Paper no. 21, 41pp.. MR-3622895
- [11] J. Alt, L. Erdős, and T. Krüger, Local inhomogeneous circular law, Ann. Appl. Probab. 28 (2018), no. 1, 148–203. MR-3770875
- [12] G. W. Anderson and O. Zeitouni, A CLT for a band matrix model, Probab. Theory Related Fields 134 (2006), no. 2, 283–338. MR-2222385
- [13] Z. D. Bai, Convergence rate of expected spectral distributions of large random matrices. I. Wigner matrices, Ann. Probab. 21 (1993), no. 2, 625–648. MR-1217559
- [14] Z. D. Bai, Convergence rate of expected spectral distributions of large random matrices.
 II. Sample covariance matrices, Ann. Probab. 21 (1993), no. 2, 649–672. MR-1217560
- [15] Z. D. Bai, Circular law, Ann. Probab. 25 (1997), no. 1, 494–529. MR-1428519
- [16] Z. D. Bai and J. W. Silverstein, No eigenvalues outside the support of the limiting spectral distribution of information-plus-noise type matrices, Random Matrices: Theory and Applications 1 (2012), no. 1. MR-2930382
- [17] Z. Bai and J. W. Silverstein, Spectral analysis of large dimensional random matrices, second ed., Springer Series in Statistics, Springer, New York, 2010. MR-2567175

- [18] S. Belinschi, A. Dembo, and A. Guionnet, Spectral measure of heavy tailed band and covariance random matrices, Comm. Math. Phys. 289 (2009), no. 3, 1023–1055. MR-2511659
- [19] C. Bordenave and D. Chafaï, Around the circular law, Probab. Surv. 9 (2012), 1–89. MR-2908617
- [20] C. Bordenave, P. Caputo, and D. Chafaï, Spectrum of non-Hermitian heavy tailed random matrices, Comm. Math. Phys. 307 (2011), no. 2, 513–560. MR-2837123
- [21] C. Bordenave, P. Caputo, and D. Chafaï, and K. Tikhomirov, On the spectral radius of a random matrix: An upper bound without fourth moment, Ann. Probab. 46 (2018), no. 4, 2268–2286. MR-3813992
- [22] N. A. Cook, W. Hachem, J. Najim, and D. Renfrew, *Non-Hermitian Random Matrices With a Variance Profile (II): Properties and Examples*, In preparation (2018).
- [23] N. A. Cook, Lower bounds for the smallest singular value of structured random matrices, Ann. Probab. 46 (2018), no. 6, 3442–3500. MR-3857860
- [24] N. A. Cook, Spectral properties of non-hermitian random matrices, Ph.D. thesis, University of California, Los Angeles, 2016. MR-3306005
- [25] J. Dumont, W. Hachem, S. Lasaulce, P. Loubaton, and J. Najim, On the capacity achieving covariance matrix for Rician MIMO channels: an asymptotic approach, IEEE Trans. Inform. Theory 56 (2010), no. 3, 1048–1069. MR-2723661
- [26] L. Erdős, B. Schlein, and H-T. Yau, Local semicircle law and complete delocalization for Wigner random matrices, Comm. Math. Phys. 287 (2009), no. 2, 641–655. MR-2481753
- [27] L. Erdős, B. Schlein, and H-T. Yau, Semicircle law on short scales and delocalization of eigenvectors for Wigner random matrices, Ann. Probab. 37 (2009), no. 3, 815–852. MR-2537522
- [28] L. Erdős, H-T. Yau, and J. Yin, Bulk universality for generalized Wigner matrices, Probab. Theory Related Fields 154 (2012), no. 1–2, 341–407. MR-2981427
- [29] J. Feinberg and A. Zee, Non-Hermitian random matrix theory: method of Hermitian reduction, Nuclear Phys. B 504 (1997), no. 3, 579–608. MR-1488584
- [30] J. Ginibre, Statistical ensembles of complex, quaternion, and real matrices, J. Mathematical Phys. 6 (1965), 440–449. MR-0173726
- [31] V. L. Girko, Circular law, Theory of Probability and Its Applications 29 (1985), no. 4, 694–706. MR-0773436
- [32] V. L. Girko, Theory of stochastic canonical equations. Vol. I, Mathematics and its Applications, vol. 535, Kluwer Academic Publishers, Dordrecht, 2001. MR-1887675
- [33] F. Götze and A. Tikhomirov, *The circular law for random matrices*, The Annals of Probability 38 (2010), no. 4, 1444–1491. MR-2663633
- [34] A. Guionnet, M. Krishnapur, and O. Zeitouni, *The single ring theorem*, Ann. of Math. (2) **174** (2011), no. 2, 1189–1217. MR-2831116
- [35] A. Guionnet, Large deviations upper bounds and central limit theorems for non-commutative functionals of Gaussian large random matrices, Ann. Inst. H. Poincaré Probab. Statist. 38 (2002), no. 3, 341–384. MR-1899457
- [36] W. Hachem, P. Loubaton, and J. Najim, Deterministic equivalents for certain functionals of large random matrices, Ann. Appl. Probab. 17 (2007), no. 3, 875–930. MR-2326235
- [37] W. Hachem, P. Loubaton, and J. Najim, A CLT for information-theoretic statistics of gram random matrices with a given variance profile, Ann. Appl. Probab. 18 (2008), no. 6, 2071– 2130. MR-2473651
- [38] L. A. Harris, Fixed points of holomorphic mappings for domains in Banach spaces [mr1981265], Proceedings of the International Conference on Fixed-Point Theory and its Applications, Hindawi Publ. Corp., Cairo, 2003, pp. 261–274. MR-2093334
- [39] Y. He, A. Knowles, and R. Rosenthal, Isotropic self-consistent equations for mean-field random matrices, Probability Theory and Related Fields, 171 (2018), no. 1–2, 203–249. MR-3800833
- [40] J. W. Helton, R. Rashidi Far, and R. Speicher, Operator-valued semicircular elements: solving a quadratic matrix equation with positivity constraints, Int. Math. Res. Not. 15 (2007), no. 22. MR-2376207

- [41] R. A. Horn and C. R. Johnson, *Topics in matrix analysis*, Cambridge University Press, Cambridge, 1994, Corrected reprint of the 1991 original. MR-1288752
- [42] R. A. Horn and C. R. Johnson, *Matrix analysis*, second ed., Cambridge University Press, Cambridge, 2013. MR-2978290
- [43] I. Jana and A. Soshnikov, ESD of singular values of random band matrices; Marchenko-Pastur law and more, arXiv preprint arXiv:1610.02153 (2016). MR-3682960
- [44] O. Kallenberg, Foundations of modern probability, second ed., Probability and its Applications (New York), Springer-Verlag, New York, 2002. MR-1876169
- [45] J. Komlós and M. Simonovits, Szemerédi's regularity lemma and its applications in graph theory, Combinatorics, Paul Erdős is eighty, Vol. 2 (Keszthely, 1993), Bolyai Soc. Math. Stud., vol. 2, János Bolyai Math. Soc., Budapest, 1996, pp. 295–352. MR-1395865
- [46] L. Lovász, Large networks and graph limits, American Mathematical Society Colloquium Publications, vol. 60, American Mathematical Society, Providence, RI, 2012. MR-3012035
- [47] R. M. May, Will a large complex system be stable?, Nature 238 (1972), 413–414.
- [48] M. L. Mehta, Randon matrices and the statistical theory of energy level, Academic Press, 1967. MR-0220494
- [49] J. Najim and J. Yao, Gaussian fluctuations for linear spectral statistics of large random covariance matrices, Ann. Appl. Probab. 26 (2016), no. 3, 1837–1887. MR-3513608
- [50] G. Pan and W. Zhou, Circular law, extreme singular values and potential theory, J. Multivariate Anal. 101 (2010), no. 3, 645–656. MR-2575411
- [51] L. Pastur and M. Shcherbina, Eigenvalue distribution of large random matrices, Mathematical Surveys and Monographs, vol. 171, American Mathematical Society, Providence, RI, 2011. MR-2808038
- [52] L. A. Pastur, Spectra of random self adjoint operators, Uspehi Mat. Nauk 28 (1973), no. 1(169), 3–64. MR-0406251
- [53] K. Rajan and L. F. Abbott, Eigenvalue spectra of random matrices for neural networks, Phys. Rev. Lett. 97 (2006), 188104.
- [54] M. Rudelson and O. Zeitouni, Singular values of Gaussian matrices and permanent estimators, Random Structures Algorithms 48 (2016), no. 1, 183–212. MR-3432577
- [55] E. Seneta, Non-negative matrices and Markov chains, Springer Series in Statistics, Springer, New York, 2006, Revised reprint of the second (1981) edition [Springer-Verlag, New York; MR-0719544]. MR-2209438
- [56] D. Shlyakhtenko, Random Gaussian band matrices and freeness with amalgamation, Internat. Math. Res. Notices (1996), no. 20, 1013–1025. MR-1422374
- [57] H. Sompolinsky, A. Crisanti, and H.-J. Sommers, *Chaos in random neural networks*, Phys. Rev. Lett. **61** (1988), no. 3, 259–262. MR-0949871
- [58] T. Tao, Outliers in the spectrum of iid matrices with bounded rank perturbations, Probab. Theory Related Fields 155 (2013), no. 1–2, 231–263. MR-3010398
- [59] T. Tao and V. Vu, Random matrices: Universality of local eigenvalue statistics, Acta Mathematica 206 (2011), no. 1, 127. MR-2784665
- [60] T. Tao, V. Vu, and M. Krishnapur, Random matrices: universality of esds and the circular law, The Annals of Probability 38 (2010), no. 5, 2023–2065. MR-2722794
- [61] T. Tao and V. H. Vu, Random matrices: the circular law, Commun. Contemp. Math. 10 (2008), no. 2, 261–307. MR-2409368
- [62] F. Wegner, Bounds on the density of states in disordered systems, Z. Phys. B 44 (1981), no. 1–2, 9–15. MR-0639135
- [63] J. Weidmann, Linear operators in Hilbert spaces, Graduate Texts in Mathematics, vol. 68, Springer-Verlag, New York-Berlin, 1980, Translated from the German by Joseph Szücs. MR-0566954
- [64] P. M. Wood, Universality and the circular law for sparse random matrices, Ann. Appl. Probab.
 22 (2012), no. 3, 1266–1300. MR-2977992

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