Non-local SFT tachyon and cosmology

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Alexey Koshelev
University of Crete

Mainly based on JHEP 02 (2007) 041, hep-th/0605085 by I.Ya. Aref'eva, A. K. and JHEP 04 (2007) 029, hep-th/0701103 by A. K.

Plan

- Overview of the problem
 - Cosmological motivations
 - Problems and challenge
 - Why String Field Theory?
- Tachyon spectroscopy
- Infinitely many scalars vs. the non-locality
- Emergence of a phantom
- Real Cosmology
- Summary and Outlook

Cosmological motivation

- Data on Ia supernovae
- Galaxy clusters measurements
- WMAP

Universe exhibits an accelerated expansion

Equation of state: $p = w\rho$, w < 0 — Dark Energy

Perlmutter et. al., 1999 Riess et. al., 2004 Spergel et. al., 2006

Theoretical issues

- w > -1 Quintessence models
- w = -1 Cosmological constant
- w < -1 Phantom models

Our universe is known to be homogeneous, isotropic and with high accuracy spatially flat.

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(1+3) dimensional spatially flat FRW universe

$$ds^2 = -dt^2 + a(t)^2 d\vec{x}^2$$

Theoretical problems

- Just a cosmological constant has no theoretical explanation so far
- It is difficult to cook a Phantom divide (w = -1) crossing.
- $w = \mathbf{const} < -1 \Rightarrow \mathbf{Big} \ \mathbf{Rip} \ \mathbf{singularity}$.
- Phantoms (ghosts) being physical particles look harmful for the theory.
- It is possible that w changes with time.

Challenge

We need a dynamical model of Dark Energy which might be able to cross the Phantom divide.

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Proposal

Derive a scalar field model of Dark Energy starting from initially reliable theory with (probably) <u>non-local</u> interaction for this scalar field.

SFT (p-adic) Tachyon

Tachyon effective action $(\alpha' = 1)$

$$S = \frac{1}{g_4^2} \int dx \sqrt{-\eta} \left(\frac{1}{2} \Phi \mathcal{F}(\square) \Phi - \frac{1}{p+1} \Phi^{p+1}(x) \right)$$

Cubic Fermionic SFT: $\mathcal{F}(z) = (\xi^2 z + 1) e^{-\frac{1}{4}z}, \ \xi^2 \approx 0.9556, \ p = 3$

Aref'eva, Belov, A.K.

Medvedev, NPB638 (2002) 3

Tachyon EOM looks very simple: $\mathcal{F}\Phi = \Phi^p$

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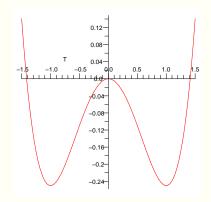
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Tachyon potential (odd p)

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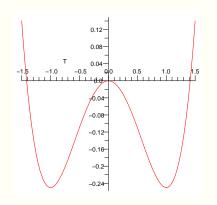
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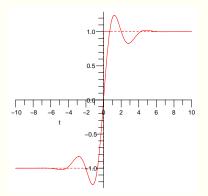
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Tachyon potential (odd p)



Rolling solution

Aref'eva, Joukovskaya, A.K., JHEP 09 (2003) 012

Good points of SFT

- The theory is UV complete
- Quantum computations in p-adic action ($\xi = 0$) can be carried out analytically up to all orders and the resulting finite effective action can be constructed

Brekke, Freund, Olson, Witten, 1988

• The interaction is non-local thus giving a chance for the Phantom divide crossing

Minimal coupling to gravity

$$S = \int dx \sqrt{-\mathbf{g}} \left(\frac{R}{2\kappa^2} + \frac{1}{g_4^2} \left(\frac{1}{2} \Phi \mathcal{F}(\square_{\mathbf{g}}) \Phi - \frac{1}{p+1} \Phi^{p+1}(x) - \frac{p-1}{2(p+1)} - \tau \right) \right)$$

- $\bullet \ \kappa^2 = 8\pi G = \frac{1}{M_P^2}$
- \bullet τ is a correction to the brane tension dictated by an existence of the rolling solution

Aref'eva, astro-ph/0410443; Aref'eva, A.K., Vernov, astro-ph/0412619

- \bullet τ is expected to be generated through coupling to closed string excitations
- We introduce $\Lambda = \frac{\tau}{g_4^2}$. Its value can be adjusted to a realistic one, e.g. giving the Hubble parameter $\sim 10^{-60} M_P$.

Late time tachyon spectroscopy

We consider a generalization:

- $\mathcal{F}(z)$ is analytic in \mathbb{C} , i.e. $\mathcal{F}(z) = c_n z^n$, $\mathcal{F}(0) = 1$, $c_n \in \mathbb{R}$
- Any p > 1

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Our expectation:

Tachyon rolls down to the minimum and is expected to stop at the bottom in infinite time

$$\Phi = 1 - \psi \Rightarrow S_{\psi} = \frac{1}{g_4^2} \int dx \sqrt{-g} \left(\frac{1}{2} \psi \mathcal{F}(\Box_g) \psi - \frac{p}{2} \psi^2 - \tau \right)$$

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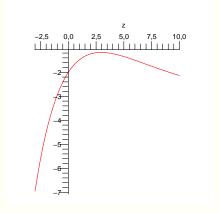
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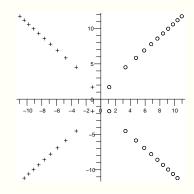


$$\mathcal{F}(z) - p$$
 in CSSFT

EOM:
$$(\mathcal{F} - p)\psi = 0$$

Characteristic equation:

$$\mathcal{F}(\omega^2) = p$$



Roots ω in CSSFT

Infinitely many scalars vs. the non-locality

New action

$$S = \frac{1}{g_4^2} \int dx \sqrt{-g} \frac{1}{2} \sum_k \left(\mathcal{F}'(\omega_k^2) \psi_k(\square_g - \omega_k^2) \psi_k + \mathcal{F}'(\omega_k^{2*}) \bar{\psi}_k(\square_g - \omega_k^{2*}) \bar{\psi}_k \right)$$

- EOMs are manifestly local and linear.
- Sum over k is indefinite until \mathcal{F} is not specified explicitly
- On the solution $\psi_k = \psi_{k+} + \psi_{k-}$ because \square_g is the second order differential operator

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- EOMs are manifestly local and linear.
- Sum over k is indefinite until \mathcal{F} is not specified explicitly
- On the solution $\psi_k = \psi_{k+} + \psi_{k-}$ because \square_g is the second order differential operator
- Spectrum and Energy momentum tensor are reproduced in this way
- The construction does not depend on a particular metric
- It is consistent to keep only one mode, say ψ_{k+} afterwards

Phantom emergence

Simplest consistent possibility: only single $\psi_{k+} \not\equiv 0$

We put
$$\psi = \alpha + i\beta$$
, $\omega^2 = M + iN$, $\mathcal{F}'(\omega^2) = x + iy$

Action for fields α and β becomes

$$S = \frac{1}{g_4^2} \int dx \sqrt{-g} \left(\alpha(x\mathcal{D} - xM + yN)\alpha - \beta(x\mathcal{D} - xM + yN)\beta - 2\alpha(y\mathcal{D} - yM - xN)\beta \right).$$

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- For any signs of parameters one normal and one phantom field present in the system

 Ostrogradski, 1850
- Only field α is physical one since $\alpha = \frac{\psi + \psi^*}{2}$
- $N \neq 0$ because there are no real roots
- \bullet M, x, y are not restricted but at least one of x or y is non-zero

This action may serve as a toy model for the tachyon around its vacuum.

Tachyon at large times must have phantom properties

Cosmological scenarios

$$S = \int dx \sqrt{-g} \left(\frac{R}{2\kappa^2} + \frac{1}{g_4^2} \left(\frac{1}{2} \psi \mathcal{F}(\Box_g) \psi - \frac{p}{2} \psi^2(x) \right) - \Lambda \right)$$

Using the developed machinery we pass to a local theory with many scalars and under an assumption that only one specific mode $\psi_{k+} \not\equiv 0$ one has as first approximation

$$\psi = \alpha e^{-rt} \cos(\nu t + \varphi)$$

$$a = a_0 e^{H_0 t} + \frac{e^{(H_0 - 2r)t}}{g_4^2 M_P^2} (s \sin(2\nu t) + c \cos(2\nu t))$$

where
$$r + i\nu = \frac{3}{2}H_0 \pm \sqrt{\frac{9}{4}H_0^2 - \omega_k^2}$$
 and $H_0 = \sqrt{\frac{\Lambda}{3M_P^2}}$

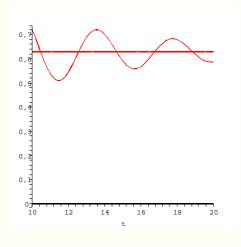
For $r = H_0/2$ oscillations in a(t) will not die despite the fact that oscillations in Φ vanish.

Cosmological properties

Generic parameters, i.e. not necessarily $r = H_0/2$.

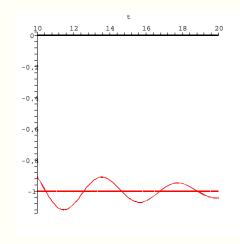
Hubble parameter

$$H = \frac{\dot{a}}{a}$$



Total effective state parameter

$$w = -1 - \frac{2}{3} \frac{\dot{H}}{H^2}$$



Quintessence and Phantom phases change one each other.

No Big Rip singularity
Crossing of the phantom divide

Summary

• Rolling SFT tachyon is a reliable candidate for the Dark Energy

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- Rolling SFT tachyon is a reliable candidate for the Dark Energy
- Non-local action with a general operator \mathcal{F} is analyzed and a local formulation for a linearization near a non-perturbative vacuum is given.
- The energy and pressure can be easily computed for a general function \mathcal{F} without specifying its explicit form as well as an arbitrary metric.
- It is shown that tachyon scalar field generates a crossing of the phantom divide in the cosmological constant background. This crossing is periodic one and a condition of non-vanishing oscillations is formulated.
- There is no Big Rip singularity in the model.

Further directions

• Coupling to dilaton.

- A.K., in progress
- Coupling to vector field etc.
- Proof of stability of found solutions
- Cosmological perturbations of theories with infinitely many derivatives
- Numeric and may be analytic solution to full equations
- As well as many other questions

Thank you!