

**NONCOMMUTATIVE INSTANTONS AND TWISTOR TRANSFORM**

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*Dedicated to A.N. Tyurin on his 60th birthday***Abstract**

Recently N. Nekrasov and A. Schwarz proposed a modification of the ADHM construction of instantons which produces instantons on a noncommutative deformation of  $\mathbb{R}^4$ . In this paper we study the relation between their construction and algebraic bundles on noncommutative projective spaces. We exhibit one-to-one correspondences between three classes of objects: framed bundles on a noncommutative  $\mathbb{P}^2$ , certain complexes of sheaves on a noncommutative  $\mathbb{P}^3$ , and the modified ADHM data. The modified ADHM construction itself is interpreted in terms of a noncommutative version of the twistor transform. We also prove that the moduli space of framed bundles on the noncommutative  $\mathbb{P}^2$  has a natural hyperkähler metric and is isomorphic as a hyperkähler manifold to the moduli space of framed torsion free sheaves on the commutative  $\mathbb{P}^2$ . The natural complex structures on the two moduli spaces do not coincide but are related by an  $SO(3)$  rotation. Finally, we propose a construction of instantons on a more general noncommutative  $\mathbb{R}^4$  than the one considered by Nekrasov and Schwarz (a  $q$ -deformed  $\mathbb{R}^4$ ).

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## 1. PHYSICAL MOTIVATION

In this section we explain the physical motivation for studying instantons on a noncommutative  $\mathbb{R}^4$ . Readers uninterested in the motivation may skip most of this section and proceed directly to subsection 1.5. Likewise, readers familiar with the way noncommutative instantons arise in string theory may start with subsection 1.5.

**1.1. Instanton equations.** Let  $E$  be a vector bundle with structure group  $G$  on an oriented Riemannian 4-manifold  $X$ , and let  $A$  be a connection on  $E$ . Instanton equation is the equation

$$(1) \quad F_A^+ = 0,$$

where  $F_A$  is the curvature of  $A$ , and  $F_A^+$  denotes the self-dual (SD) part of  $F_A$ . Solutions of this equation are called instantons, or anti-self-dual (ASD) connections. The second Chern class of  $E$  is known in the physics literature as the instanton number. Instantons automatically satisfy the Yang-Mills equation  $d_A(*F) = 0$ , where  $d_A : \Omega^p \otimes \text{End}(E) \rightarrow \Omega^{p+1} \otimes \text{End}(E)$  is the covariant differential, and  $*$  :  $\Omega^p \rightarrow \Omega^{4-p}$  is the Hodge star operator.

There are several physical reasons to be interested in instantons. If one is studying quantum gauge theory on a Riemannian 3-manifold  $M$  (space), then instantons on  $X = M \times \mathbb{R}$  describe quantum-mechanical tunneling between different classical vacua. The possibility of such tunneling has drastic physical effects, some of which can be experimentally observed. If one is studying classical gauge theory on a 5-dimensional space-time  $X \times \mathbb{R}$ , then instantons on  $X$  can be interpreted as solitons, i.e. as static solutions of the Yang-Mills equations of motion. In fact, instantons are the absolute minima of the Yang-Mills energy function of the 5-dimensional theory (with fixed second Chern class).

Both interpretations arise in string theory, but to explain this we need to make a digression and discuss D-branes.

**1.2. D-branes.** It has been discovered in the last few years that string theory describes, besides strings, extended objects (branes) of various dimensions. These extended objects should be regarded as static solutions of (as yet poorly understood) stringy equations of motion. D-branes are a particularly manageable class of branes. Recall that ordinary closed oriented superstrings, known as Type II strings, are described by maps from a Riemann surface  $\Sigma$  (“worldsheet”) to a 10-dimensional manifold  $Z$  (“target”). The physical definition of a D-brane is “a submanifold of  $Z$  on which strings can end.” This means that if a D-brane is present, then one needs to consider maps from a Riemann surface with boundaries to  $Z$  such that the boundaries are mapped to a certain submanifold  $X \subset Z$ . In this case one says

that there is a D-brane wrapped on  $X$ . If  $X$  is connected and has dimension  $p+1$ , then one says that one is dealing with a D $p$ -brane. In general,  $X$  can have several components with different dimensions, and then each component corresponds to a D-brane.

In perturbative string theory, the role of equations of motion is played by the condition that a certain auxiliary quantum field theory on the Riemann surface  $\Sigma$  is conformally invariant. When D-branes are present,  $\Sigma$  has boundaries, and the auxiliary theory must be supplemented with boundary conditions. The requirement that the boundary conditions preserve conformal invariance imposes constraints on the submanifold  $X$ . This constraints should be regarded as equations of motion for D-branes. For example, if we consider a D0-brane wrapped on a 1-dimensional submanifold  $X$ , then conformal invariance requires that  $X$  be a geodesic in  $Z$ . This is the usual equation of motion for a relativistic particle moving in  $Z$ .

An important subtlety is that to specify fully the boundary conditions for the auxiliary theory on  $\Sigma$  it is not sufficient to specify  $X$ ; one should also specify a unitary vector bundle  $E$  on  $X$  and a connection on it. In the simplest case this bundle has rank 1, but one can also have “multiple” D-branes, described by bundles of rank  $r > 1$ . Such bundles describe  $r$  coincident D-branes wrapped on the same submanifold  $X$ . Using the requirement of conformal invariance of the auxiliary two-dimensional quantum field theory, one can derive equations of motion for the Yang-Mills connection on  $E$ . In the low-energy approximation, the equations of motion are the usual Yang-Mills equations  $d_A(*F_A) = 0$ . In particular, instantons are solutions of these equations.

**1.3. Instantons and D-branes.** Let  $Z$  be  $\mathbb{R}^{10}$  with a flat metric, and let  $X \hookrightarrow Z$  be  $\mathbb{R}^5 = \mathbb{R}^4 \times \mathbb{R}$  linearly embedded in  $Z$ . We regard  $\mathbb{R}^4$  as space and  $\mathbb{R}$  as time. Consider  $r$  D4-branes wrapped on  $X$ . This physical system is described by the Yang-Mills action on  $\mathbb{R}^5 = \mathbb{R}^4 \times \mathbb{R}$ . If one is looking for static solutions of the equations of motion, one needs to consider the minima of the Yang-Mills energy function

$$W[A] = \int_{\mathbb{R}^4} \|F_A\|^2,$$

where  $F_A$  is the curvature of a  $U(r)$  connection  $A$ , and  $\|F_A\|^2 = -\text{Tr}(F_A \wedge *F_A)$ . The instanton number of  $A$  is defined by

$$(2) \quad c_2 = \frac{1}{8\pi^2} \int_{\mathbb{R}^4} \text{Tr}(F_A \wedge F_A).$$

If the Yang-Mills energy evaluated on  $A$  is finite, then the bundle  $E$  and the connection  $A$  extend to  $S^4$ , the one-point compactification of  $\mathbb{R}^4$  (see [4] for details). In this case  $c_2$  is the second Chern class of  $E$  and is therefore an integer.

Solutions of instanton equations on  $\mathbb{R}^4$  are precisely the absolute minima of the Yang-Mills energy function. These solutions should be regarded as composed of identical particle-like objects (instantons) on  $X$ , their number being  $c_2$ . Since the energy of the instanton is proportional to  $c_2$ , all “particles” have the same mass. Since the solution is static, the particles neither repel nor attract. This is actually a consequence of supersymmetry: Type II string theory is supersymmetric, and D4-branes with instantons on them leave part of supersymmetry unbroken.

In string theory one may also consider  $k$  D0-branes present simultaneously with  $r$  D4-branes. More specifically, we will consider D0-branes which are at rest, i.e. the corresponding one-dimensional manifolds are straight lines parallel to the time axis. Such a configuration of branes is also supersymmetric, and consequently there are no forces between any of the branes. The positions of D0-branes are not constrained by anything, so their moduli space is  $(\mathbb{R}^9)^k$ . More precisely, since D0-branes are indistinguishable, the moduli space is  $Sym^k(\mathbb{R}^9)$ .

It turns out that an instanton with instanton number  $k$  and  $k$  D0-branes are related: they can be deformed into each other without any cost in energy. A convenient point of view is the following. In the presence of D4-branes wrapped on  $X$  the moduli space of D0-branes has two branches: a branch where their positions are unconstrained and D0-branes are point-like (this branch is isomorphic to  $Sym^k(\mathbb{R}^9)$ ), and the branch where they are constrained to lie on  $X$ . The latter branch is isomorphic to the moduli space  $\mathcal{M}_{r,k}$  of  $U(r)$  instantons on  $X = \mathbb{R}^4$  with  $c_2 = k$ .

The dimension of  $\mathcal{M}_{r,k}$  is known to be  $4rk$  for  $r > 1$  (see for example [4]). For  $r = 1$  instantons do not exist. The translation group of  $\mathbb{R}^4$  acts freely on  $\mathcal{M}_{r,k}$ , and the quotient space describes the relative positions and sizes of instantons. Thus D0-branes are point-like objects when they are away from D4-branes, but when they bind to D4-branes they can acquire finite size.

The “instanton” branch touches the “point-like” branch at submanifolds where some or all of the instantons shrink to zero size. These are the submanifolds where the instanton moduli space is singular. At these submanifolds the point-like instantons can detach from D4-branes and start a new life as D0-branes. This lowers the second Chern class of the bundle on D4-branes. Thus from the string theory perspective it is natural to glue together the moduli spaces of instantons with different Chern classes along singular submanifolds.

**1.4. Noncommutative geometry and D-branes.** Instanton equations (and, more generally, Yang-Mills equations) arise in the low-energy limit of string theory, or equivalently

for large string tension. Recently, another kind of low-energy limit of string theory was discussed in the literature [31]. Consider a trivial  $U(r)$ -bundle on  $X = \mathbb{R}^4$  with a connection  $A$  whose curvature  $F_A$  is of the form  $1 \otimes f$ , where  $1$  is the unit section of  $\text{End}(E)$ , and  $f$  is a constant nondegenerate 2-form. For small  $f$  the D4-branes are described by the ordinary Yang-Mills action, but for large  $F_A$  the stringy equations of motion get complicated. It turns out that the equations of motion simplify again in the limit when both  $F_A$  and the string tension are taken to infinity, with a certain combination of the two kept fixed (one also has to scale the metric appropriately, see [31]). We will call this limit the Seiberg-Witten limit. In this limit the D4-branes are described by Yang-Mills equations on a certain noncommutative deformation of  $\mathbb{R}^4$  (see [31] and references therein).

There is another description of the Seiberg-Witten limit, which is gauge-equivalent to the previous one. Type II string theory reduces at low energies to Type II supergravity in 10 dimensions. The bosonic fields of this low-energy theory include a symmetric rank-two tensor (metric) and a 2-form  $B$ .  $\mathbb{R}^{10}$  with a flat Lorentzian metric and a constant  $B$  is a solution of supergravity equations of motion, as well as full stringy equations of motion. A constant  $B$  can be gauged away, so this is not a very interesting solution. Life gets more interesting if there are D-branes present. For example, consider  $r$  coincident flat D4-branes embedded in  $\mathbb{R}^{10}$  with a constant  $B$ -field. It turns out that one can gauge away a constant  $B$ -field only at the expense of introducing a constant  $F_A$  of the form  $1 \otimes f$  where  $f$  is equal to the pull-back of  $B$  to the worldvolume of the D4-branes. Thus the solution with zero  $F_A$  and nonzero  $B$  is equivalent to the solution with nonzero  $F_A$  and zero  $B$ . Therefore the Seiberg-Witten limit can be described as the limit in which both the  $B$ -field and the string tension become infinite.

The idea that D-branes in a nonzero B-field are described Yang-Mills theory on a noncommutative space was first put forward in [13] for the case of D-branes wrapped on tori.

**1.5. Instanton equations on a noncommutative  $\mathbb{R}^4$ .** The deformed  $\mathbb{R}^4$  that one obtains in the Seiberg-Witten limit is completely characterized by its algebra of functions  $\mathcal{A}$ . It is a noncommutative algebra whose underlying space is a certain subspace of  $C^\infty$  functions on  $\mathbb{R}^4$ . The product is the so-called Wigner-Moyal product formally given by

$$(3) \quad (f \star g)(x) = \lim_{y \rightarrow x} \exp\left(\frac{1}{2} \hbar \theta_{ij} \frac{\partial^2}{\partial x_i \partial y_j}\right) f(x)g(y).$$

Here  $\theta$  is a purely imaginary matrix, and  $\hbar$  is a real parameter (“Planck constant”) which is introduced to emphasize that the Wigner-Moyal product is a deformation of the usual product. In the string theory context  $\theta$  is proportional to  $f^{-1}$ .

Of course, to make sense of this definition we must specify a subspace in the space of  $C^\infty$  functions which is closed under the Wigner-Moyal product. Leaving this question aside for a moment,<sup>1</sup> one can define the exterior differential calculus over  $\mathcal{A}$ . Differential geometry of noncommutative spaces has been developed by A. Connes [12]. In our situation Connes' general theory is greatly simplified. For example, the sheaf of 1-forms  $\Omega^1(\mathcal{A})$  is simply a bimodule  $\mathcal{A}^{\oplus 4}$  (the relation of this definition with the general theory is explained in subsection 8.11). The elements of  $\Omega^1(\mathcal{A})$  will be denoted  $\sum_i f^i(x) dx_i$ , or simply  $f^i(x) dx_i$ . The exterior differential  $d$  is a vector space morphism

$$d : \mathcal{A} \rightarrow \Omega^1(\mathcal{A}), \quad f \mapsto \frac{\partial f}{\partial x_i} dx_i.$$

The exterior differential  $d$  satisfies the Leibniz rule

$$d(f_1 \star f_2) = df_1 \star f_2 + f_1 \star df_2.$$

This makes sense because  $\Omega^1(\mathcal{A})$  is a bimodule.

The sheaf of 2-forms over  $\mathcal{A}$  is a bimodule  $\Omega^2(\mathcal{A}) = \mathcal{A}^{\oplus 6}$  (see subsection 8.11). The definition of the exterior differential can be extended to  $\Omega^1(\mathcal{A})$  in an obvious manner.

Complex conjugation acts as an anti-linear anti-homomorphism of  $\mathcal{A}$ , i.e.  $\overline{(f \star g)} = \bar{g} \star \bar{f}$ . Thus  $\mathcal{A}$  has a natural structure of a  $\star$ -algebra. We will denote the  $\star$ -conjugate of  $f \in \mathcal{A}$  by  $f^\dagger$ .

A trivial bundle over the noncommutative  $\mathbb{R}^4$  is defined as a free  $\mathcal{A}$ -module  $E$ . A trivial unitary bundle over the noncommutative  $\mathbb{R}^4$  is defined as a free module  $V \otimes_{\mathbb{C}} \mathcal{A}$ , where  $V$  is a Hermitean vector space. A connection on a trivial bundle  $E$  is defined as a map

$$\nabla : E \rightarrow E \otimes_{\mathcal{A}} \Omega^1(\mathcal{A}),$$

which is a vector space morphism satisfying the Leibniz rule

$$\nabla(m \star f) = \nabla(m) \star f + m \star df.$$

This formula makes use of the bimodule structure on  $\Omega^1(\mathcal{A})$ .

The curvature  $F_\nabla = [\nabla, \nabla]$  is a morphism of  $\mathcal{A}$ -modules

$$F_\nabla : E \rightarrow E \otimes_{\mathcal{A}} \Omega^2(\mathcal{A}).$$

As in the commutative case, a connection on a trivial bundle  $E$  can be written in terms of a connection 1-form  $A \in \text{End}_{\mathcal{A}}(E) \otimes_{\mathcal{A}} \Omega^1(\mathcal{A})$ :

$$\nabla(m) = dm + A \star m.$$

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<sup>1</sup>String theory considerations do not shed light on this problem.

This formula uses the bimodule structure on  $m$ . If  $E$  is a unitary bundle, and we have  $A^\dagger = -A$ , then we say that  $A$  is a unitary connection.

The curvature is given in terms of  $A$  by the usual formula

$$F_\nabla := F_A = dA + A \wedge A.$$

Here it is understood that

$$f^i dx_i \wedge g^j dx_j = f^i \star g^j dx_i \wedge dx_j.$$

The instanton equation on  $\mathcal{A}$  is again given by (1), and the instanton number is defined by (2).

The most obvious choice of the space of functions closed under the Wigner-Moyal product is the space of polynomial functions. However, this choice is not suitable for our purposes because it precludes the decrease of  $F_A$  at infinity which is necessary for the instanton action to converge. In the commutative case, components of an instanton connection are rational functions [4], so we would like our class of functions to include rational functions on  $\mathbb{R}^4$ . A possible choice for the underlying set of  $\mathcal{A}$  is the set of  $C^\infty$  functions on  $\mathbb{R}^4$  all of whose derivatives are polynomially bounded. Then we face the question of the convergence of the series (3). To avoid dealing with this issue, we modify our definition of the Wigner-Moyal product (see Appendix for details). The modified product makes the space of  $C^\infty$  functions all of whose derivatives are polynomially bounded into an algebra over  $\mathbb{C}$ , and agrees with (3) on polynomial functions.

Polynomial functions form a subalgebra of  $\mathcal{A}$ . This subalgebra is isomorphic to the algebra generated by four variables  $x_i, i = 1, 2, 3, 4$  with relations

$$[x_i, x_j] = \hbar \theta_{ij}.$$

This algebra is usually called the Weyl algebra.

To summarize, there is a limit of string theory in which D4 branes are described by Yang-Mills equations on the noncommutative  $\mathbb{R}^4$  ( $= \mathcal{A}$ ). D0-branes bound to D4-branes are described in this limit by the instanton equations on the noncommutative  $\mathbb{R}^4$ . One can show that, unlike in the commutative case, instantons cannot be deformed to point-like D0-branes without a cost in energy. Thus it is natural to suspect that the moduli space of instantons on the noncommutative  $\mathbb{R}^4$  is metrically complete.

## 2. REVIEW OF THE ADHM CONSTRUCTION AND SUMMARY

All instantons on the commutative  $\mathbb{R}^4$  arise from the so-called ADHM construction. Recently N. Nekrasov and A. Schwarz [28] introduced a modification of this construction which

produces instantons on the noncommutative  $\mathbb{R}^4$ .<sup>2</sup> In the commutative case the completeness of the ADHM construction can be proved using the twistor transform of R. Penrose, so one could hope that the same approach could work in the noncommutative case. In this paper we show that the deformed ADHM data of [28] describe holomorphic bundles on certain noncommutative algebraic varieties and interpret the deformed ADHM construction in terms of noncommutative twistor transform. In this subsection we review both ordinary and deformed ADHM constructions and make a summary of our results.

**2.1. Review of the ADHM construction of instantons.** First let us outline the ADHM construction of  $U(r)$  instantons on the commutative  $\mathbb{R}^4$  following [15]. We assume that the constant metric  $G$  on  $\mathbb{R}^4$  has been brought to the standard form  $G = \text{diag}(1, 1, 1, 1)$  by a linear change of basis. To construct a  $U(r)$  instanton with  $c_2 = k$  one starts with two Hermitean vector spaces  $V \simeq \mathbb{C}^k$  and  $W \simeq \mathbb{C}^r$ . The ADHM data consist of four linear maps  $B_1, B_2 \in \text{Hom}(V, V)$ ,  $I \in \text{Hom}(W, V)$ ,  $J \in \text{Hom}(V, W)$  which satisfy the following two conditions:

$$(i) \quad \mu_c = [B_1, B_2] + IJ = 0, \quad \mu_r = [B_1, B_1^\dagger] + [B_2, B_2^\dagger] + II^\dagger - J^\dagger J = 0.$$

(ii) For any  $\xi = (\xi_1, \xi_2) \in \mathbb{C}^2 \cong \mathbb{R}^4$  the linear map  $\mathcal{D}_\xi \in \text{Hom}(V \oplus V \oplus W, V \oplus V)$  defined by

$$(4) \quad \mathcal{D}_\xi = \begin{pmatrix} B_1 - \xi_1 & -B_2 + \xi_2 & I \\ B_2^\dagger - \bar{\xi}_2 & B_1^\dagger - \bar{\xi}_1 & J^\dagger \end{pmatrix}$$

is surjective.

The equations  $\mu_c = \mu_r = 0$  are called the ADHM equations. They are invariant with respect to the action of the group of unitary transformations of  $V$ . Solutions of these equations are called ADHM data. The space of ADHM data modulo  $U(V)$  transformations has dimension  $4rk$  and carries a natural hyperkähler metric. ADHM construction identifies this moduli space with the moduli space of  $U(r)$  instantons with  $c_2 = k$  and fixed trivialization at infinity. The role of the condition (ii) above is to remove submanifolds in this moduli space where the hyperkähler metric becomes singular (these are point-like instanton singularities mentioned in subsection 1.3). As a result the moduli space of the ADHM data is metrically incomplete.

The instanton connection can be reconstructed from the ADHM data as follows. The condition (ii) implies that the family  $\mathbf{Ker} \mathcal{D}_\xi$  forms a trivial subbundle of  $V \oplus V \oplus W$  of rank  $r$ . Let  $v(\xi)$  be its trivialization, i.e. a linear map  $v(\xi) : \mathbb{C}^r \rightarrow V \oplus V \oplus W$  smoothly

<sup>2</sup>As in the commutative case, one may consider different classes of functions on the noncommutative  $\mathbb{R}^4$ : polynomial,  $C^\infty$  functions rapidly decreasing at infinity,  $C^\infty$  functions all of whose derivatives are polynomially bounded, etc. Our class of functions differs somewhat from that adopted in [28].



depending on  $\xi$  such that  $\mathcal{D}_\xi v(\xi) = 0$  for all  $\xi$ , and  $\rho(\xi) = v(\xi)^\dagger v(\xi)$  is an isomorphism for all  $\xi$ . We set

$$A(\xi) = \rho(\xi)^{-1} v(\xi)^\dagger dv(\xi).$$

The matrix-valued one-form  $A$  is a connection on a trivial unitary bundle of rank  $r$ . One can show that its curvature  $F_A$  is ASD (see [4]). However, it does not satisfy  $A^\dagger = -A$ , because we are not using a unitary gauge. Instead  $A$  satisfies

$$A^\dagger(\xi) = -(\rho(\xi)A(\xi)\rho(\xi)^{-1} + \rho(\xi)d\rho(\xi)^{-1}).$$

To go to a unitary gauge, we must make a gauge transformation

$$A'(\xi) = g(\xi)A(\xi)g(\xi)^{-1} + g(\xi)dg(\xi)^{-1},$$

where  $g(\xi)$  is a function taking values in Hermitian  $r \times r$  matrices and satisfying  $g(\xi)^2 = \rho(\xi)$ .

We now explain, following [28], how to modify the ADHM construction so that it produces rank  $r$  instantons on the noncommutative  $\mathbb{R}^4$  defined in the previous section. It proves convenient to apply an orthogonal transformation which brings the matrix  $\theta$  in (3) to the standard form

$$\theta = \sqrt{-1} \begin{pmatrix} 0 & a & 0 & 0 \\ -a & 0 & 0 & 0 \\ 0 & 0 & 0 & b \\ 0 & 0 & -b & 0 \end{pmatrix}.$$

We will assume that  $a + b \neq 0$ . Since  $\theta$  enters only in the combination  $\hbar\theta$ , we can set  $a + b = 1$  without loss of generality. The relation between the affine coordinates  $\xi_1, \xi_2$  on  $\mathbb{C}^2$  and affine coordinates  $x_1, x_2, x_3, x_4$  on  $\mathbb{R}^4$  is chosen as follows:

$$\xi_1 = x_4 - \sqrt{-1} x_3, \quad \xi_2 = -x_2 + \sqrt{-1} x_1.$$

Then  $\xi_1, \xi_2, \bar{\xi}_1, \bar{\xi}_2$  obey the Weyl algebra relations

$$[\xi_1, \bar{\xi}_1] = 2\hbar b, \quad [\xi_2, \bar{\xi}_2] = 2\hbar a, \quad [\xi_1, \xi_2] = [\xi_1, \bar{\xi}_2] = [\bar{\xi}_1, \xi_2] = [\bar{\xi}_1, \bar{\xi}_2] = 0.$$

The modified ADHM data consist of the same four maps which now satisfy

$$\mu_c = 0, \quad \mu_r = -2\hbar(a + b) \cdot 1_{k \times k}.$$

The instanton connection is given by essentially the same formulas as in the commutative case. The operator  $\mathcal{D}$  is given by the same formula as  $\mathcal{D}_\xi$ , but is now regarded as an element of

$$\text{Hom}_{\mathcal{A}}((V \oplus V \oplus W) \otimes_{\mathbb{C}} \mathcal{A}, (V \oplus V) \otimes_{\mathbb{C}} \mathcal{A}).$$

The module  $\mathbf{Ker} \mathcal{D}$  is a projective module over  $\mathcal{A}$ . Following [10], we assume that it is isomorphic to a free module of rank  $r$ , and  $v$  is the corresponding isomorphism  $v : \mathcal{A}^{\oplus r} \rightarrow \mathbf{Ker} \mathcal{D}$ . We further assume [10] that the morphism

$$\Delta = \mathcal{D}\mathcal{D}^\dagger \in \text{End}_{\mathcal{A}}((V \oplus V) \otimes \mathcal{A})$$

is an isomorphism.<sup>3</sup> Then it is easy to see that  $\rho = v^\dagger v \in \text{End}_{\mathcal{A}}(\mathbb{C}^r \otimes \mathcal{A})$  is an isomorphism too. We set

$$(5) \quad A = \rho^{-1} v^\dagger dv.$$

(The multiplication here and below is understood to be the Wigner-Moyal multiplication.) This formula defines a connection 1-form on a trivial unitary bundle on  $\mathcal{A}$  of rank  $r$ . The curvature of this connection is given by

$$F_A = \rho^{-1} dv^\dagger \wedge (1 - v\rho^{-1}v^\dagger)dv.$$

A short computation (essentially the same as in the commutative case) shows that the curvature can be written in the form

$$F_A = \rho^{-1} v^\dagger d\mathcal{D}^\dagger \Delta^{-1} \wedge d\mathcal{D} v.$$

Furthermore, since  $\mathcal{D}$  and  $\mathcal{D}^\dagger$  are linear in  $\xi_i, \bar{\xi}_i$ , their exterior derivatives have a very simple form:

$$d\mathcal{D} = \begin{pmatrix} -d\xi_1 & d\xi_2 & 0 \\ -d\bar{\xi}_2 & -d\bar{\xi}_1 & 0 \end{pmatrix}, \quad d\mathcal{D}^\dagger = \begin{pmatrix} -d\bar{\xi}_1 & -d\xi_2 \\ d\bar{\xi}_2 & -d\xi_1 \\ 0 & 0 \end{pmatrix}.$$

Note also that by virtue of the deformed ADHM equations  $\Delta$  has a block-diagonal form:

$$\Delta = \begin{pmatrix} \delta & 0 \\ 0 & \delta \end{pmatrix},$$

where  $\delta \in \text{End}_{\mathcal{A}}(V \otimes \mathcal{A})$  is an isomorphism. Using this fact, one can easily see that  $F_A$  is proportional to the 2-forms

$$d\xi_1 \wedge d\bar{\xi}_1 + d\xi_2 \wedge d\bar{\xi}_2, \quad d\xi_1 \wedge d\bar{\xi}_2, \quad d\xi_2 \wedge d\bar{\xi}_1,$$

which are anti-self-dual.

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<sup>3</sup>One can show that the latter assumption is always valid provided  $\hbar \neq 0$ . As for the former one, it is not known what constraints should be imposed on the deformed ADHM data to ensure that  $\mathbf{Ker} \mathcal{D}$  is a free  $\mathcal{A}$ -module of rank  $r$ . For  $r = 1$   $\mathbf{Ker} \mathcal{D}$  is never free [16].

As in the commutative case, the connection  $A$  does not satisfy  $A^\dagger = -A$ . To go to a unitary gauge one has to perform a gauge transformation

$$A' = g \star A \star g^{-1} + g \star dg^{-1}.$$

Here  $g \in \text{Aut}_{\mathcal{A}}(\mathbb{C}^r \otimes \mathcal{A})$  should be found from the conditions  $g^\dagger = g$ ,  $g \star g = \rho$ . The existence of such  $g$  is an additional assumption.

**2.2. Summary of results.** In the commutative case there is a one-to-one correspondence between the following three classes of objects:

- A. Rank  $r$  holomorphic bundles on  $\mathbb{P}^2$  with  $c_2 = k$  and a fixed trivialization on the line at infinity.
- B. The set of ADHM data modulo the natural action of  $U(k)$ .
- C. Rank  $r$  holomorphic bundles on  $\mathbb{P}^3$  with  $c_2 = k$ , a trivialization on a fixed line, vanishing  $H^1(E(-2))$ , and satisfying a certain reality condition.
- D.  $U(r)$  instantons on  $\mathbb{R}^4$  with  $c_2 = k$ .

The correspondence between  $C$  and  $D$  is a particular instance of twistor transform [6]. The correspondence between  $B$  and  $C$  has been proved by Atiyah, Hitchin, Drinfeld, and Manin [5, 4]. Together these two results imply that all instantons on  $\mathbb{R}^4$  arise from the ADHM construction. The correspondence between  $A$  and  $B$  has been proved by Donaldson [15]. One can also prove the correspondence between  $A$  and  $D$  directly [7, 11, 18].

The goal of this paper is to extend some of these results to the noncommutative case. We show that there is a natural one-to-one correspondence between the isomorphism classes of the following objects:

- $A'$ . Algebraic bundles on a noncommutative deformation of  $\mathbb{P}^2$  with  $c_2 = k$  and a fixed trivialization on the line at infinity.
- $B'$ . Deformed ADHM data of Nekrasov and Schwarz modulo the natural  $U(k)$  action.
- $C'$ . Certain complexes of sheaves on a noncommutative deformation of  $\mathbb{P}^3$  satisfying reality conditions.

The moduli space of the deformed ADHM data has a natural hyperkähler metric, and the other two moduli spaces inherit this metric.

Furthermore, we reinterpret the deformed ADHM construction of Nekrasov and Schwarz in terms of a noncommutative deformation of the twistor transform.

It is interesting to note that H. Nakajima [26] studied the same linear algebra data as Nekrasov and Schwarz and showed that their moduli space coincides with the moduli space of torsion free sheaves on a commutative  $\mathbb{P}^2$  with a trivialization on a fixed line. On the other

hand, we show that the same data describe algebraic bundles on a noncommutative  $\mathbb{P}^2$ . As shown below, the interpretation in terms of complexes of sheaves on a noncommutative  $\mathbb{P}^3$  provides a geometric reason for this “coincidence.” We prove that the two moduli spaces are isomorphic as hyperkähler manifolds, but the natural complex structures on them differ by an  $\mathrm{SO}(3)$  rotation.

The rest of the paper is organized as follows. In Section 3 we define noncommutative deformations of certain commutative projective varieties ( $\mathbb{P}^2$ ,  $\mathbb{P}^3$ , and a quadric in  $\mathbb{P}^5$ ). Section 4 is an algebraic preparation for the study of bundles on noncommutative projective spaces. In Section 5 we study the cohomological properties of sheaves on noncommutative  $\mathbb{P}^2$  and  $\mathbb{P}^3$  and define locally free sheaves (i.e. bundles). In Section 6 we show that any bundle on a noncommutative  $\mathbb{P}^2$  trivial on the commutative line at infinity arises as a cohomology of a monad. In Section 7 we exhibit bijections between  $A'$ ,  $B'$ , and  $C'$  and explain the relation with Nakajima’s results. In Section 8 we construct a noncommutative deformation of Grassmannians and flag manifolds and describe a noncommutative version of the twistor transform. We also describe a nice class of noncommutative projective varieties associated with a Yang-Baxter operator and define differential forms on these varieties. In section 9 we consider a more general deformation of  $\mathbb{R}^4$  (a  $q$ -deformed  $\mathbb{R}^4$ ) whose physical significance is obscure at present. We propose an ADHM-like construction of instantons on this space and outline its relation to noncommutative algebraic geometry. In the Appendix we define the Wigner-Moyal product on the space of  $C^\infty$  functions on  $\mathbb{R}^n$  all of whose derivatives are polynomially bounded, and prove that the Wigner-Moyal product provides this space with a structure of an algebra over  $\mathbb{C}$ .

### 3. GEOMETRY OF NONCOMMUTATIVE VARIETIES

**3.1. Algebraic preliminaries.** Let  $\mathbf{k}$  be a base field (we will be dealing only with  $\mathbf{k} = \mathbb{C}$  or  $\mathbf{k} = \mathbb{R}$  in this paper). Let  $A$  be an algebra over  $\mathbf{k}$ . It is called right (left) noetherian if every right (left) ideal is finitely generated, and it is called noetherian if it is both right and left noetherian.

Let  $A = \bigoplus_{i \geq 0} A_i$  be a graded noetherian algebra. We denote by  $\mathrm{mod}(A)$  the category of finitely generated right  $A$ -modules, by  $\mathrm{gr}(A)$  the category of finitely generated graded right  $A$ -modules, and by  $\mathrm{tors}(A)$  the full subcategory of  $\mathrm{gr}(A)$  which consists of finite dimensional graded  $A$ -modules.

An important role will be played by the quotient category  $\mathrm{qgr}(A) = \mathrm{gr}(A)/\mathrm{tors}(A)$ . It has the following explicit description. The objects of  $\mathrm{qgr}(A)$  are the objects of  $\mathrm{gr}(A)$  (we denote by  $\widetilde{M}$  the object in  $\mathrm{qgr}(A)$  which corresponds to a module  $M$ ). The morphisms

in  $qgr(A)$  are given by

$$\mathrm{Hom}_{qgr}(\widetilde{M}, \widetilde{N}) = \varinjlim_{M'} \mathrm{Hom}_{gr}(M', N)$$

where  $M'$  runs over submodules of  $M$  such that  $M/M'$  is finite dimensional.

On the category  $gr(A)$  there is a shift functor: for a given graded module  $M = \bigoplus_{i \geq 0} M_i$  the shifted module  $M(r)$  is defined by  $M(r)_i = M_{r+i}$ . The induced shift functor on the quotient category  $qgr(A)$  sends  $\widetilde{M}$  to  $\widetilde{M}(r) = \widetilde{M(r)}$ .

Similarly, we can consider the category  $Gr(A)$  of all graded right  $A$ -modules. It contains the subcategory  $Tors(A)$  of torsion modules. Recall that a module  $M$  is called torsion if for any element  $x \in M$  one has  $xA_{\geq s} = 0$  for some  $s$ , where  $A_{\geq s} = \bigoplus_{i \geq s} A_i$ . We denote by  $QGr(A)$  the quotient category  $Gr(A)/Tors(A)$ . The category  $QGr(A)$  contains  $qgr(A)$  as a full subcategory. Sometimes it is convenient to work in  $QGr(A)$  instead of  $qgr(A)$ .

Henceforth, all graded algebras will be noetherian algebras generated by the first component  $A_1$  with  $A_0 = \mathbf{k}$ .

Sometimes we use subscripts  $R$  or  $L$  for categories  $gr(A)$ ,  $qgr(A)$ , etc., to specify whether right or left modules are considered. If the subscript is omitted, the modules are taken to be right modules. For the same reason for an  $A$ -bimodule  $M$  we sometimes write  $M_A$  or  ${}_A M$  to specify whether the right or left module structure is considered.

**3.2. Noncommutative varieties.** A variety in commutative geometry is a topological space with a sheaf of functions (continuous, smooth, analytic, algebraic, etc.) which is, obviously, a sheaf of algebras. One of the main objects in geometry (algebraic or differential) is a bundle or, more generally, a sheaf. To any variety  $X$  we can associate an abelian category of sheaves of modules (maybe with some additional properties) over the sheaf of algebras of functions. Given a sheaf of modules on  $X$ , the space of its global sections is a module over the algebra of global functions on  $X$ . Thus the functor of global sections associates to every  $X$  an algebra and a certain category of modules over it. Under favorable circumstances, much of the information about the geometry of  $X$  is contained in this purely algebraic datum. Let us give a few examples.

If  $X$  is a compact Hausdorff topological space, then the category of vector bundles over  $X$  is equivalent to the category of finitely generated projective modules over the algebra of continuous functions on  $X$  [33, 35]. The equivalence is given by the functor which maps a vector bundle to the module of its global sections.

It is well known that if  $A$  is a commutative noetherian algebra, the category of coherent sheaves on the noetherian affine scheme  $Spec(A)$  is equivalent to the category of finitely

generated modules over  $A$ . The equivalence is again given by the functor which attaches to a coherent sheaf the module of its global sections.

In the case of projective varieties the only global functions are constants, so one has to act somewhat differently. Since a projective variety  $X$  is by definition a subvariety of a projective space, it inherits from it the line bundle  $\mathcal{O}_X(1)$  and its tensor powers  $\mathcal{O}_X(i)$ . We can consider a graded algebra

$$\Gamma(X) = \bigoplus_{i \geq 0} H^0(X, \mathcal{O}_X(i)).$$

This algebra is called the homogeneous coordinate algebra of  $X$ . Furthermore, for any sheaf  $\mathcal{F}$  we can define a graded  $A$ -module

$$\Gamma(\mathcal{F}) = \bigoplus_{i \geq 0} H^0(X, \mathcal{F}(i)).$$

It can be checked that  $\Gamma$  is a functor from the category of coherent sheaves on  $X$   $coh(X)$  to  $gr(\Gamma(X))$ . In a brilliant paper [32], J-P. Serre described the category of coherent sheaves on a projective scheme  $X$  in terms of graded modules over the graded algebra  $\Gamma(X)$ . He proved that the category  $coh(X)$  is equivalent to the quotient category  $qgr(\Gamma(X)) = gr(\Gamma(X))/tors(\Gamma(X))$ . The equivalence is given by the composition of the functor  $\Gamma$  with the projection  $\pi : gr(A) \rightarrow qgr(A)$ . On other hand, let  $A = \bigoplus_{i \geq 0} A_i$  be a graded commutative algebra generated over  $\mathbb{k}$  by the first component (which is assumed to be finite dimensional). We can associate to  $A$  a projective scheme  $X = Proj(A)$ . Serre proved that the category  $coh(X)$  is equivalent to the category  $qgr(A)$ .

The equivalence also holds for the category of quasicoherent sheaves on  $X$  and the category  $QGr(A) = Gr(A)/Tors(A)$ .

In all of the above examples it turned out that the natural category of sheaves or bundles on a variety is equivalent to a certain category defined in terms of (graded) modules over some (graded) algebra. On the other hand, “as A. Grothendieck taught us, to do geometry you really don’t need a space, all you need is a category of sheaves on this would-be space” ([24], p.83).

For this reason, in the realm of algebraic geometry it is natural to regard a noncommutative noetherian algebra as a coordinate algebra of a noncommutative affine variety; then the category of finitely generated right modules over this algebra is identified with the category of coherent sheaves on the corresponding variety. Similarly, a noncommutative graded noetherian algebra is regarded as a homogeneous coordinate algebra of a noncommutative projective variety. The category of finitely generated graded right modules over this algebra

modulo torsion modules is identified with the category of coherent sheaves on this variety (see [3], [24], [34]).

A different approach to noncommutative geometry has been pursued by A. Connes [12].

**3.3. Noncommutative deformations of commutative varieties.** Many important noncommutative varieties arise as deformations of commutative ones.

Let  $X$  be a commutative variety (affine or projective). Let  $A$  be the corresponding commutative (graded) algebra. A noncommutative deformation of  $X$  is a deformation of the algebra structure on  $A$ , that is, a deformation of the multiplication law. Usually it is not easy to write down an explicit formula for the deformed product.

There is a more algebraic way to describe noncommutative deformations of commutative varieties. Assume that the algebra  $A$  is given in terms of generators and relations. This means that  $A$  is given as a quotient  $A = T(V)/\langle R \rangle$ , where  $V$  is the vector space spanned by the generators,  $T(V)$  is the tensor algebra of  $V$ , and  $\langle R \rangle$  is a two-sided ideal in  $T(V)$  generated by a subspace of relations  $R \subset T(V)$ . Assume that  $R_{\hbar} \subset T(V)$  is a one-parameter deformation of the subspace  $R$ . Then  $A_{\hbar} = T(V)/\langle R_{\hbar} \rangle$  is a one-parameter deformation of  $A$ . (If  $A$  is graded, then we assume that  $R$  is a graded subspace, and the deformation preserves the grading).

We denote by  $X_{\hbar}$  the noncommutative variety corresponding to the algebra  $A_{\hbar}$ . Thus  $X_{\hbar}$  is a noncommutative one-parameter deformation of  $X$ .

If  $X$  is projective and  $A$  is a graded algebra, then we denote by  $\text{coh}(X_{\hbar})$  the category  $\text{qgr}(A_{\hbar})$ . Furthermore, as in the commutative case, we will write  $\mathcal{O}(r)$  for the object  $\widetilde{A}_{\hbar}(r)$ .

Now we define noncommutative varieties which are going to be used in this paper.

**3.4. Noncommutative  $\mathbb{C}^4$ .** Denote by  $A(\mathbb{C}^4)$  the algebra of polynomial functions on  $\mathbb{C}^4$ . Let  $\theta$  be a skew-symmetric  $4 \times 4$  matrix.

Let us define the algebra  $A(\mathbb{C}_{\hbar}^4)$  as an algebra over  $\mathbb{C}$  generated by  $x_i$  ( $i = 1, 2, 3, 4$ ) with relations  $[x_i, x_j] = \hbar\theta_{ij}$ :

$$(6) \quad A(\mathbb{C}_{\hbar}^4) = T(x_1, x_2, x_3, x_4)/\langle [x_i, x_j] = \hbar\theta_{ij} \rangle_{1 \leq i, j \leq 4}.$$

We will regard  $A(\mathbb{C}_{\hbar}^4)$  as the algebra of polynomial functions on a noncommutative affine variety  $\mathbb{C}_{\hbar}^4$ .

**3.5. Noncommutative 4-dimensional quadric.** Let  $G$  be a  $4 \times 4$  symmetric nondegenerate matrix. Consider a graded algebra  $Q_{\hbar} = \bigoplus_{i \geq 0} Q_i$  over  $\mathbb{C}$  generated by the elements

$X_1, X_2, X_3, X_4, D, T$  of degree 1 with the following quadratic relations:

$$(7) \quad \begin{aligned} [T, D] &= [T, X_i] = 0, \\ [X_i, X_j] &= \hbar \theta_{ij} T^2, \\ [D, X_i] &= 2\hbar \sum_{lk} \theta_{il} G^{lk} X_k T, \\ \sum_{ij} G^{ij} X_i X_j &= DT. \end{aligned}$$

We denote by  $\mathbb{Q}_{\hbar}^4$  the noncommutative projective variety corresponding to the algebra  $Q_{\hbar}$ . It is evident that  $\mathbb{Q}_{\hbar}^4$  is a deformation of a 4-dimensional commutative quadric  $\mathbb{Q}^4 = \{\sum_{ij} G^{ij} X_i X_j = DT\} \subset \mathbb{C}\mathbb{P}^5$ .

**3.6. Embedding**  $\mathbb{C}_{\hbar}^4 \hookrightarrow \mathbb{Q}_{\hbar}^4$ . Let  $Q_{\hbar}[T^{-1}]$  be the localization of the algebra  $Q_{\hbar}$  with respect to  $T$ . Elements of degree 0 in  $Q_{\hbar}[T^{-1}]$  form a subalgebra which will be denoted by  $Q_{\hbar}[T^{-1}]_0$ .

**Lemma 3.1.** *The map  $x_i \mapsto T^{-1}X_i$  ( $i = 1, 2, 3, 4$ ) induces an isomorphism of the algebra  $A(\mathbb{C}_{\hbar}^4)$  with the algebra  $Q_{\hbar}[T^{-1}]_0$ .*

*Proof.* Obvious. □

This means that  $\mathbb{C}_{\hbar}^4$  can be identified with the open subset  $\{T \neq 0\}$  in  $\mathbb{Q}_{\hbar}^4$ . For this reason,  $\mathbb{Q}_{\hbar}^4$  may be regarded as a compactification of  $\mathbb{C}_{\hbar}^4$  which is compatible with the bilinear form  $G$ . Note also that the complement of  $\mathbb{C}_{\hbar}^4$  in  $\mathbb{Q}_{\hbar}^4$  corresponds to the algebra

$$Q_{\hbar}/\langle T \rangle = \mathbb{T}(X_1, X_2, X_3, X_4, D) / \langle [X_i, X_j] = [D, X_i] = 0, \sum_{ij} G^{ij} X_i X_j = 0 \rangle.$$

Since this algebra is commutative, the complement is the usual 3-dimensional commutative quadratic cone. Thus one may say that  $\mathbb{Q}_{\hbar}^4$  is obtained from  $\mathbb{C}_{\hbar}^4$  by adding a cone ‘‘at infinity’’. This is in complete analogy with the commutative case.

**3.7. Noncommutative  $\mathbb{P}_{\hbar}^2$  and  $\mathbb{P}_{\hbar}^3$ .** Noncommutative deformations of the projective plane have been classified in [1], [2], [9]. We will need one of them, namely the one whose homogeneous coordinate algebra is a graded algebra  $PP_{\hbar} = \bigoplus_{i \geq 0} PP_{\hbar i}$  over  $\mathbb{C}$  generated by the elements  $w_1, w_2, w_3$  of degree 1 with the relations:

$$(8) \quad \begin{aligned} [w_3, w_i] &= 0 \text{ for any } i = 1, 2, 3, \\ [w_1, w_2] &= 2\hbar w_3^2. \end{aligned}$$

We will also need a noncommutative deformation of the 3-dimensional projective space, whose homogeneous coordinate algebra will be denoted  $PS_{\hbar} = \bigoplus_{i \geq 0} PS_{\hbar i}$ . It is a graded



algebra over  $\mathbb{C}$  generated by  $PS_{\hbar 1} = U$ , where the vector space  $U$  is spanned by elements  $z_1, z_2, z_3, z_4$  obeying the relations

$$(9) \quad \begin{aligned} [z_3, z_i] &= [z_4, z_i] = 0 \quad \text{for any } i = 1, 2, 3, 4, \\ [z_1, z_2] &= 2\hbar z_3 z_4. \end{aligned}$$

The noncommutative projective varieties corresponding to  $PP_{\hbar}$  and  $PS_{\hbar}$  will be denoted  $\mathbb{P}_{\hbar}^2$  and  $\mathbb{P}_{\hbar}^3$ , respectively.

Note that for  $\hbar \neq 0$  all algebras  $PS_{\hbar}$  are isomorphic, and therefore the varieties  $\mathbb{P}_{\hbar}^3$  are the same for all  $\hbar \neq 0$ . The same is true for  $\mathbb{P}_{\hbar}^2$ .

**3.8. Subvarieties in  $\mathbb{P}_{\hbar}^3$  and  $\mathbb{P}_{\hbar}^2$ .** If  $I \subset PS_{\hbar}$  is a graded two-sided ideal, then the quotient algebra  $PS_{\hbar}/I$  corresponds to a closed subvariety  $X(I) \subset \mathbb{P}_{\hbar}^3$ . Let us describe some of them.

Let  $J$  be the graded two-sided ideal generated by  $z_3$  and  $z_4$ . Then  $PS_{\hbar}/J = \mathbb{T}(z_1, z_2)/\langle [z_1, z_2] = 0 \rangle$ , hence  $X(J)$  is the commutative projective line.

For each point  $p = (\lambda : \mu) \in \mathbb{P}^1$  let  $J_p$  denote the graded two-sided ideal generated by  $\lambda z_3 + \mu z_4$ . If  $p = (0 : 1)$  or  $p = (1 : 0)$ , then it is easy to see that  $X(J_p)$  is the commutative projective plane. For all other  $p \in \mathbb{P}^1$  we have

$$PS_{\hbar}/J_p = \mathbb{T}(z_1, z_2, z_3)/\left\langle [z_1, z_3] = [z_2, z_3] = 0, [z_1, z_2] = -2\hbar \frac{\lambda}{\mu} z_3^2 \right\rangle,$$

hence  $X(J_p)$  is a noncommutative projective plane isomorphic to  $\mathbb{P}_{\hbar}^2$ .

We have  $J_p \subset J$  for all  $p \in \mathbb{P}^1$ , hence all planes  $X(J_p)$  pass through the line  $X(J)$ . Thus we see that  $\mathbb{P}_{\hbar}^3$  is a pencil of noncommutative projective planes passing through a fixed commutative projective line.

Similarly, the two-sided ideal generated by  $w_3$  in  $PP_{\hbar}$  corresponds to a commutative projective line  $l = \{w_3 = 0\} \subset \mathbb{P}_{\hbar}^2$ .

#### 4. PROPERTIES OF ALGEBRAS $PS_{\hbar}$ AND $PP_{\hbar}$ AND THE RESOLUTION OF THE DIAGONAL

This section is a preparation for the study of sheaves on  $\mathbb{P}_{\hbar}^3$  and  $\mathbb{P}_{\hbar}^2$ . We show that the algebras  $PS_{\hbar}$  and  $PP_{\hbar}$  are regular and Koszul and construct the resolution of the diagonal, which will enable us to associate monads to certain bundles on  $\mathbb{P}_{\hbar}^2$ .

**4.1. Quadratic algebras.** A graded algebra  $A = \bigoplus_{i \geq 0} A_i$  over a field  $\mathbb{k}$  is called quadratic if it is connected (i.e.  $A_0 = \mathbb{k}$ ), is generated by the first component  $A_1$ , and the ideal of relations is generated by the subspace of quadratic relations  $R(A) \subset A_1 \otimes A_1$ .

Therefore the algebra  $A$  can be represented as  $T(A_1)/\langle R(A) \rangle$ , where  $T(A_1)$  is a free tensor algebra generated by the space  $A_1$ .

The algebras  $PS_{\hbar}$  and  $PP_{\hbar}$  are quadratic algebras. For example,  $PS_{\hbar}$  can be represented as  $T(U)/\langle W \rangle$ , where  $U = PS_{\hbar 1}$  is a 4-dimensional vector space and  $W$  is the 6-dimensional subspace of  $U \otimes U$  spanned by the relations (9).

**4.2. The dual algebra.** For any quadratic algebra  $A = T(A_1)/\langle R(A) \rangle$  we can define its dual algebra which is also quadratic.

Let us identify  $A_1^* \otimes A_1^*$  with  $(A_1 \otimes A_1)^*$  by  $(l \otimes m)(a \otimes b) = m(a)l(b)$ . Denote by  $R(A)^\perp$  the annihilator of  $R(A)$  in  $A_1^* \otimes A_1^*$ , i.e. the subspace which consists of such  $q \in (A_1^*)^{\otimes 2}$  that  $q(r) = 0$  for any  $r \in R(A)$ .

**Definition 4.1.** ([24]) *The algebra  $A^! = T(A_1^*)/\langle R(A)^\perp \rangle$  is called the dual algebra of  $A$ .*

**Example 4.2.** *Let  $\{\check{z}_i\}, i = 1, 2, 3, 4$ , be the basis of  $PS_{\hbar 1}^! = U^*$  which is dual to  $\{z_i\}$ . By definition,  $PS_{\hbar}^!$  is generated by  $\{\check{z}_i\}$  with defining relations*

$$\begin{aligned} \check{z}_i^2 &= 0 && \text{for all } i = 1, \dots, 4; \\ \check{z}_i \check{z}_j + \check{z}_j \check{z}_i &= 0 && \text{for all } i < j, (i, j) \neq (3, 4); \\ \check{z}_3 \check{z}_4 + \check{z}_4 \check{z}_3 &= \hbar[\check{z}_1, \check{z}_2] = 2\hbar \check{z}_1 \check{z}_2. \end{aligned}$$

In the commutative case the dual algebra of the symmetric algebra  $S(U)$  is isomorphic to the exterior algebra  $\Lambda(U^*)$ . Obviously, the algebras  $PS_{\hbar}^!$  and  $PP_{\hbar}^!$  are deformations of exterior algebras. For example, the vector space  $PS_{\hbar k}^!$  is spanned by the elements  $\check{z}_{i_1} \cdots \check{z}_{i_k}$  with  $i_1 < \cdots < i_k$ . In particular, the dimension of the vector space  $PS_{\hbar k}^!$  is equal to  $\binom{4}{k}$ . Similarly, the dimension of  $PP_{\hbar k}^!$  is equal to  $\binom{3}{k}$ .

**Proposition 4.3.** *Let  $A$  be  $PS_{\hbar}$  or  $PP_{\hbar}$ , and let  $n$  be 4 or 3, respectively. The multiplication map  $A_k^! \otimes A_{n-k}^! \longrightarrow A_n^!$  is a non-degenerate pairing. Hence the dual algebra  $A^!$  is a Frobenius algebra, i.e.  $(A^!)_{A^!} \cong (A^! A^!)^*$  as right  $A^!$ -modules.*

*Proof.* The proposition holds for the exterior algebra, and therefore also for the algebra  $A^!$ , since the latter is a “small” deformation of the exterior algebra.  $\square$

**4.3. The Koszul complex.** Consider right  $A$ -modules  $(A_k^!)^* \otimes A$ . The following complex  $K(A)$  is called the (right) Koszul complex of a quadratic algebra:

$$\cdots \xrightarrow{d} (A_3^!)^* \otimes A(-3) \xrightarrow{d} (A_2^!)^* \otimes A(-2) \xrightarrow{d} (A_1^!)^* \otimes A(-1) \xrightarrow{d} (A_0^!)^* \otimes A \longrightarrow 0,$$

where the map  $d : (A_k^!)^* \otimes A \rightarrow (A_{k-1}^!)^* \otimes A$  is a composition of two natural maps:

$$(A_k^!)^* \otimes A \longrightarrow (A_k^!)^* \otimes A_1^! \otimes A_1 \otimes A \longrightarrow (A_k^!)^* \otimes A.$$

Here the first arrow sends  $\alpha \otimes a$  to  $\alpha \otimes e \otimes a$  with  $e$  defined as

$$e = \sum_i y_i \otimes x_i \in A_1^! \otimes A_1,$$

and  $\{x_i\}$  and  $\{y_i\}$  being the dual bases of  $A_1$  and  $A_1^!$ , respectively. The second map is determined by the algebra structures on  $A^!$  and  $A$ .

It is a well-known fact that  $d^2 = 0$  (see, for example, [24]).

Let  $\mathbf{k}_A$  be the trivial right  $A$ -module. The Koszul complex  $K.(A)$  possesses a natural augmentation  $K. \xrightarrow{\varepsilon} \mathbf{k}_A \rightarrow 0$ .

**Definition 4.4.** (see [30]) *A quadratic algebra  $A = \bigoplus_{i \geq 0} A_i$  is called a Koszul algebra if the augmented Koszul complex  $K.(A) \xrightarrow{\varepsilon} \mathbf{k}_A \rightarrow 0$  is exact.*

In the same manner one can define the left Koszul complex of a quadratic algebra. It is well known that the exactness of the right Koszul complex is equivalent to the exactness of the left Koszul complex (see, for example, [21]).

**Proposition 4.5.** *The algebras  $PS_{\hbar}$  and  $PP_{\hbar}$  are Koszul algebras.*

*Proof.* For  $\hbar = 0$  this is a well-known fact about the symmetric algebra  $S(U)$ . Since the augmented Koszul complex is exact for  $\hbar = 0$ , it is also exact for small  $\hbar$ , and consequently for all  $\hbar$ .  $\square$

Since the dual algebras  $PS_{\hbar}^!$  and  $PP_{\hbar}^!$  are finite, the Koszul resolutions for the algebras  $PS_{\hbar}$  and  $PP_{\hbar}$  are finite too and have the same form as the resolutions for ordinary symmetric algebras. For example, the Koszul resolution for  $A = PP_{\hbar}$  is:

$$\{0 \rightarrow (A_3^!)^* \otimes A(-3) \rightarrow (A_2^!)^* \otimes A(-2) \rightarrow (A_1^!)^* \otimes A(-1) \rightarrow (A_0^!)^* \otimes A\} \rightarrow \mathbb{C}.$$

**4.4. Resolution of the diagonal.** Consider a bigraded vector space

$$K_{\cdot, \cdot}^2(A) = \bigoplus_{k, l \geq 0} K_{k, l}^2(A) \quad \text{with} \quad K_{k, l}^2(A) = A(k) \otimes (A_{l-k}^!)^* \otimes A(-l).$$

Consider morphisms  $d_R : K_{k, l}^2 \rightarrow K_{k, l-1}^2$  and  $d_L : K_{k, l}^2 \rightarrow K_{k+1, l}^2$  given by the following compositions

$$\begin{aligned} d_R &: A \otimes (A_k^!)^* \otimes A \rightarrow A \otimes (A_k^!)^* \otimes A_1^! \otimes A_1 \otimes A \rightarrow A \otimes (A_{k-1}^!)^* \otimes A, \\ d_L &: A \otimes (A_k^!)^* \otimes A \rightarrow A \otimes A_1 \otimes A_1^! \otimes (A_k^!)^* \otimes A \rightarrow A \otimes (A_{k-1}^!)^* \otimes A. \end{aligned}$$

Here the leftmost maps are given by

$$e_R = \sum_i y_i \otimes x_i \in A_1^! \otimes A_1 \quad \text{and} \quad e_L = \sum_i x_i \otimes y_i \in A_1 \otimes A_1^!,$$

where  $\{x_i\}$  and  $\{y_i\}$  are the dual bases of  $A_1$  and  $A_1^!$ , respectively, while the rightmost maps are induced by the algebra structures of  $A^!$  and  $A$ . It is easy to show that

$$d_R^2 = d_L^2 = 0 \quad \text{and} \quad d_R d_L = d_L d_R,$$

hence  $K^2(A)$  is a bicomplex. It is called the *double Koszul bicomplex* of the quadratic algebra  $A$ .

The topmost part of the bicomplex looks as follows:

$$\begin{array}{ccccc} \dots & \xrightarrow{d_R} & A \otimes (A_{l+1}^!)^* \otimes A(-1-l) & \xrightarrow{d_R} & A \otimes (A_l^!)^* \otimes A(-l) & \xrightarrow{d_R} & \dots \\ & & \downarrow d_L & & \downarrow d_L & & \\ \dots & \xrightarrow{d_R} & A(1) \otimes (A_l^!)^* \otimes A(-1-l) & \xrightarrow{d_R} & A(1) \otimes (A_{l-1}^!)^* \otimes A(-l) & \xrightarrow{d_R} & \dots \end{array}$$

Each term of the bicomplex  $K^2(A)$  has an obvious structure of a bigraded  $A$ -bimodule, and it is clear that the differentials are morphisms of bigraded  $A$ -bimodules.

Let

$$\mathcal{K}_l(A) = \text{Ker } d_L : K_{0,l}^2(A) \rightarrow K_{1,l}^2(A).$$

Then  $\mathcal{K}(A)$  is a complex of bigraded  $A$ -bimodules (with respect to the differential  $d_R$ ).

Consider a bigraded algebra  $\Delta = \bigoplus_{i,j} \Delta_{ij}$  with  $\Delta_{ij} = A_{i+j}$  and with the multiplication induced from  $A$ . The algebra  $\Delta$  is called the diagonal bigraded algebra of  $A$ . Note that the multiplication map induces a surjective morphism of  $A$ -bimodules  $\delta : A \otimes A \rightarrow \Delta$ .

**Lemma 4.6.** *The map*

$$\delta : \mathcal{K}_0(A) = A \otimes A \rightarrow \Delta$$

*gives an augmentation of the complex  $\mathcal{K}(A)$ .*

*Proof.* We have to check that  $\delta \cdot d_R : \mathcal{K}_1(A) \rightarrow A$  vanishes. Note that  $K_{0,1}^2(A) = A \otimes A_1 \otimes A(-1)$ , and that the differentials  $d_R$  and  $d_L$  restricted to  $K_{0,1}^2(A)$  coincide with the multiplication maps  $m_{1,2}$  and  $m_{2,3}$ , respectively. Thus we have the following commutative diagram:

$$\begin{array}{ccccc} \mathcal{K}_1(A) & \xrightarrow{d_R} & \mathcal{K}_0(A) & \xrightarrow{\delta} & \Delta \\ \downarrow & & \parallel & & \parallel \\ A \otimes A_1 \otimes A(-1) & \xrightarrow{m_{1,2}} & A \otimes A & \xrightarrow{\delta} & \Delta \\ m_{2,3} \downarrow & & & & \\ A(1) \otimes A(-1) & & & & \end{array}$$

Now the Lemma follows because  $\delta \cdot m_{1,2} = \delta \cdot m_{2,3}$  (associativity) obviously annihilates  $\text{Ker } m_{2,3} = \mathcal{K}_1(A)$ .  $\square$

**Proposition 4.7.** *If  $A$  is Koszul, then  $\mathcal{K}.(A) \xrightarrow{\delta} \Delta$  is exact.*

*Proof.* The  $(p, q)$ -bigraded component of  $K_{k,l}^2(A)$  is equal to  $A_{p+k} \otimes (A_{l-k}^!)^* \otimes A_{q-l}$ , hence the  $(p, q)$ -bigraded component of the bicomplex  $K_{\cdot,\cdot}^2(A)$  vanishes for  $l < k$  or  $l > q$ . Thus the  $(p, q)$ -bigraded component of the bicomplex  $K_{\cdot,\cdot}^2(A)$  is bounded. Therefore both spectral sequences of the bicomplex  $K_{\cdot,\cdot}^2(A)$  converge to the cohomology of the total complex  $\text{Tot}(K_{\cdot,\cdot}^2(A))$ . The first term of the first spectral sequence reads

$$E_{k,l}^1 = \begin{cases} A(l) \otimes \mathbf{k}(-l), & \text{if } k = l \\ 0, & \text{otherwise} \end{cases}$$

Hence the spectral sequence degenerates in the first term, and we have

$$H^0(\text{Tot}(K_{\cdot,\cdot}^2(A))) = \bigoplus_{l=0}^{\infty} A(l) \otimes \mathbf{k}(-l), \quad H^{\neq 0}(\text{Tot}(K_{\cdot,\cdot}^2(A))) = 0.$$

On the other hand, the first term of the second spectral sequence reads

$$E_{k,l}^1 = \begin{cases} \mathbf{k}(l) \otimes A(-l), & \text{if } k = l > 0 \\ \mathcal{K}_l(A), & \text{if } k = 0 \\ 0, & \text{otherwise} \end{cases}$$

Hence the spectral sequence degenerates in the second term, and we have

$$H^0(\text{Tot}(K_{\cdot,\cdot}^2(A))) = H^0(\mathcal{K}.(A)) \oplus \left( \bigoplus_{l=1}^{\infty} \mathbf{k}(l) \otimes A(-l) \right), \quad H^l(\text{Tot}(K_{\cdot,\cdot}^2(A))) = H^l(\mathcal{K}.(A)).$$

Therefore  $H^{\neq 0}(\mathcal{K}.(A)) = 0$ , and we have an exact sequence

$$0 \rightarrow H^0(\mathcal{K}.(A)) \rightarrow \bigoplus_{l=0}^{\infty} A(l) \otimes \mathbf{k}(-l) \rightarrow \bigoplus_{l=1}^{\infty} \mathbf{k}(l) \otimes A(-l) \rightarrow 0.$$

Looking at  $(p, q)$ -bigraded component of this sequence we see that

$$(H^0(\mathcal{K}.(A)))_{p,q} = \begin{cases} A_{p+q}, & \text{if } p, q \geq 0 \\ 0, & \text{otherwise} \end{cases}$$

Thus  $H^0(\mathcal{K}.(A)) = \Delta$ . □

**Definition 4.8.** *Define the left  $A$ -module  $\Omega^k$  as the cohomology of the left Koszul complex, truncated in the term  $K_k$ . In particular,  $\Omega^1$  is defined by the so-called Euler sequence*

$$(10) \quad 0 \rightarrow \Omega^1 \rightarrow A(-1) \otimes A_1 \xrightarrow{m} A \xrightarrow{\varepsilon} \mathbf{k} \rightarrow 0.$$

In section 8.11 we will show that for noncommutative projective spaces the sheaves corresponding to the modules  $\Omega^k$  can be regarded as sheaves of differential forms.

**Proposition 4.9.** *We have  $\mathcal{K}_k(A) = \Omega^k(k) \otimes A(-k)$  .*

*Proof.* This follows immediately from the definition of  $\Omega^k$  and  $\mathcal{K}_k(A)$  . □

Combining Propositions 4.7 and 4.9, we obtain the following resolution of the diagonal:

$$(11) \quad \dots \longrightarrow \Omega^2(2) \otimes A(-2) \longrightarrow \Omega^1(1) \otimes A(-1) \longrightarrow A \otimes A \longrightarrow \Delta \longrightarrow 0.$$

**4.5. Cohomological properties of the algebras  $PS_{\hbar}$  and  $PP_{\hbar}$  .** First we note that both algebras  $PS_{\hbar}$  and  $PP_{\hbar}$  are noetherian. This follows from the fact that they are Ore extensions of commutative polynomial algebras (see for example, [25]). For the same reason the algebras  $PS_{\hbar}$  and  $PP_{\hbar}$  have finite right (and left) global dimension, which is equal to 4 and 3, respectively (see [25], p. 273).

We remind that the global dimension of a ring  $A$  is the minimal number  $n$  (if it exists) such that for any two modules  $M$  and  $N$  we have  $\text{Ext}_A^{n+1}(M, N) = 0$  .

In the paper [1] the notion of a regular algebra has been introduced. Regular algebras have many good properties (see [3], [2], [39], etc.).

**Definition 4.10.** *A graded algebra  $A$  is called regular of dimension  $d$  if it satisfies the following conditions:*

- (1)  $A$  has global dimension  $d$ ,
- (2)  $A$  has polynomial growth, i.e.  $\dim A_n \leq cn^\delta$  for some  $c, \delta \in \mathbb{R}$ ,
- (3)  $A$  is Gorenstein, meaning that  $\text{Ext}_A^i(\mathbf{k}, A) = 0$  if  $i \neq d$ ,  
and  $\text{Ext}_A^d(\mathbf{k}, A) = \mathbf{k}(l)$  for some  $l$ .

Here  $\text{Ext}_A$  stands for the Ext functor in the category  $\text{mod}(A)$  .

It is easy to see that these properties are verified for  $PS_{\hbar}$  and  $PP_{\hbar}$  . Property (2) holds because our algebras grow as ordinary polynomial algebras. Property (3) follows from the fact that  $PS_{\hbar}$  and  $PP_{\hbar}$  are Koszul algebras and the dual algebras are Frobenius resolutions. In this case the Gorenstein parameter  $l$  in (3) is equal to the global dimension  $d$  . Thus we have

**Proposition 4.11.** *The algebras  $PS_{\hbar}$  and  $PP_{\hbar}$  are noetherian regular algebras of global dimension 4 and 3, respectively. For these algebras the Gorenstein parameter  $l$  coincides with the global dimension  $d$  .*

## 5. COHOMOLOGICAL PROPERTIES OF SHEAVES ON $\mathbb{P}_{\hbar}^2$ AND $\mathbb{P}_{\hbar}^3$

**5.1. Ampleness and cohomology of  $\mathcal{O}(i)$  .** Let  $A$  be a graded algebra and  $X$  be the corresponding noncommutative projective variety. Consider the sequence of sheaves  $\{\mathcal{O}(i)\}_{i \in \mathbb{Z}}$  in the category  $\text{coh}(X) \cong \text{qgr}(A)$ , where  $\mathcal{O}(i) = \widetilde{A(i)}$  .

This sequence is called ample if the following conditions hold:

- (a) For every coherent sheaf  $\mathcal{F}$  there are integers  $k_1, \dots, k_s$  and an epimorphism
 
$$\bigoplus_{i=1}^s \mathcal{O}(-k_i) \longrightarrow \mathcal{F}.$$
- (b) For every epimorphism  $\mathcal{F} \longrightarrow \mathcal{G}$  the induced map
 
$$\mathrm{Hom}(\mathcal{O}(-n), \mathcal{F}) \longrightarrow \mathrm{Hom}(\mathcal{O}(-n), \mathcal{G})$$
 is surjective for  $n \gg 0$ .

It is proved in [3] that the sequence  $\{\mathcal{O}(i)\}$  is ample in  $qgr(A)$  for a graded right noetherian  $\mathbf{k}$ -algebra  $A$  if it satisfies the extra condition:

$$(\chi_1): \quad \dim_{\mathbf{k}} \mathrm{Ext}_A^1(\mathbf{k}, M) < \infty$$

for any finitely generated graded  $A$ -module  $M$ .

This condition can be verified for all noetherian regular algebras (see [3], Theorem 8.1). In particular, the categories  $\mathrm{coh}(\mathbb{P}_{\hbar}^3)$ ,  $\mathrm{coh}(\mathbb{P}_{\hbar}^2)$  have ample sequences.

For any sheaf  $\mathcal{F} \in qgr(A)$  we can define a graded module  $\Gamma(\mathcal{F})$  by the rule:

$$\Gamma(\mathcal{F}) := \bigoplus_{i \geq 0} \mathrm{Hom}(\mathcal{O}(i), \mathcal{F})$$

It is proved in [3] that for any noetherian algebra  $A$  that satisfies the condition  $\chi_1$  the correspondence  $\Gamma$  is a functor from  $qgr(A)$  to  $gr(A)$  and the composition of  $\Gamma$  with the natural projection  $\pi: gr(A) \longrightarrow qgr(A)$  is isomorphic to the identity functor (see [3], ch. 3,4).

Now we formulate a result about the cohomology of sheaves on noncommutative projective spaces. This result is proved in [3] for a general regular algebra and parallels the commutative case.

**Proposition 5.1.** (Theorem 8.1. [3]) *Let  $A$  be  $PS_{\hbar}$  or  $PP_{\hbar}$ , and  $X$  be  $\mathbb{P}_{\hbar}^3$  or  $\mathbb{P}_{\hbar}^2$ , respectively. Denote by  $n$  the dimension of  $X$  (in our case  $n = 3$  or  $n = 2$ , respectively). Then*

- 1) *The cohomological dimension of  $\mathrm{coh}(X)$  is equal to  $\dim(X)$ , i.e. for any two coherent sheaves  $\mathcal{F}$  and  $\mathcal{G}$   $\mathrm{Ext}^i(\mathcal{F}, \mathcal{G})$  vanishes if  $i > n$ .*
- 2) *There are isomorphisms*

$$(12) \quad H^p(X, \mathcal{O}(i)) = \begin{cases} A_k & \text{for } p = 0, i \geq 0 \\ A_{-i-1-n}^* & \text{for } p = n, i \leq -n - 1 \\ 0 & \text{otherwise} \end{cases}$$

This proposition and the ampleness of the sequence  $\{\mathcal{O}(i)\}$  implies the following corollary:

**Corollary 5.2.** *Let  $X$  be either  $\mathbb{P}_h^3$  or  $\mathbb{P}_h^2$ . Then for any sheaf  $\mathcal{F} \in \text{coh}(X)$  and for all sufficiently large  $i \geq 0$  we have*

$$\text{Hom}(\mathcal{F}, \mathcal{O}(i)) = 0.$$

*Proof.* By ampleness a sheaf  $\mathcal{F}$  can be covered by a finite sum of sheaves  $\mathcal{O}(k_j)$ . Now the statement follows from the Proposition, because  $\text{Hom}(\mathcal{O}(k_j), \mathcal{O}(i)) = 0$  for all  $i < k_j$ .  $\square$

**Corollary 5.3.** *Let  $X$  be either  $\mathbb{P}_h^3$  or  $\mathbb{P}_h^2$ . Then for any sheaf  $\mathcal{F} \in \text{coh}(X)$  and for all sufficiently large  $i \geq 0$  we have*

$$H^k(X, \mathcal{F}(i)) = 0$$

for all  $k \geq 1$ .

*Proof.* The group  $H^k(X, \mathcal{F}(i))$  coincides with  $\text{Ext}^k(\mathcal{O}(-i), \mathcal{F})$ . Let  $k$  be the maximal integer (it exists because the global dimension is finite) such that for some  $\mathcal{F}$  there exists arbitrarily large  $i$  such that  $\text{Ext}^k(\mathcal{O}(-i), \mathcal{F}) \neq 0$ . Assume that  $k \geq 1$ . Choose an epimorphism  $\bigoplus_{j=1}^s \mathcal{O}(-k_j) \rightarrow \mathcal{F}$ . Let  $\mathcal{F}_1$  denote its kernel. Then for  $i > \max\{k_j\}$  we have  $\text{Ext}^{>0}(\mathcal{O}(-i), \bigoplus_{j=1}^s \mathcal{O}(-k_j)) = 0$ , hence  $\text{Ext}^k(\mathcal{O}(-i), \mathcal{F}) \neq 0$  implies  $\text{Ext}^{k+1}(\mathcal{O}(-i), \mathcal{F}) \neq 0$ . This contradicts the assumption of the maximality of  $k$ .  $\square$

**5.2. Serre duality and the dualizing sheaf.** A very useful property of commutative smooth projective varieties is the existence of the dualizing sheaf. Recall that a sheaf  $\omega$  is called dualizing if for any  $\mathcal{F} \in \text{coh}(X)$  there are natural isomorphisms of  $\mathbf{k}$ -vector spaces

$$H^i(X, \mathcal{F}) \cong \text{Ext}^{n-i}(\mathcal{F}, \omega)^*,$$

where  $*$  denotes the  $\mathbf{k}$ -dual. The Serre duality theorem asserts the existence of the dualizing sheaf for smooth projective varieties. In this case the dualizing sheaf is a line bundle and coincides with the sheaf of differential forms  $\Omega_X^n$  of top degree.

Since the definition of  $\omega$  is given in abstract categorical terms, it can be extended to the noncommutative case. More precisely, we will say that  $qgr(A)$  satisfies classical Serre duality if there is an object  $\omega \in qgr(A)$  together with natural isomorphisms

$$\text{Ext}^i(\mathcal{O}, -) \cong \text{Ext}^{n-i}(-, \omega)^*.$$

Our noncommutative varieties  $\mathbb{P}_h^3$  and  $\mathbb{P}_h^2$  satisfy classical Serre duality, with dualizing sheaves being  $\mathcal{O}_{\mathbb{P}_h^3}(-4)$  and  $\mathcal{O}_{\mathbb{P}_h^2}(-3)$ , respectively. This follows from the paper [39], where the existence of a dualizing sheaf in  $qgr(A)$  has been proved for a general class of algebras which includes all noetherian regular algebras. In addition, the authors of [39] showed that



the dualizing sheaf coincides with  $\tilde{A}(-l)$ , where  $l$  is the Gorenstein parameter for  $A$  (see condition (3) of Definition 4.10).

**5.3. Bundles on noncommutative projective spaces.** To any graded right  $A$ -module  $M$  one can attach a left  $A$ -module  $M^\vee = \text{Hom}_A(M, A)$  which is also graded. Note that under this correspondence the right module  $A_A(r)$  goes to the left module  ${}_A A(-r)$ .

It is known that if  $A$  is a noetherian regular algebra, then  $\text{Hom}_A(-, A)$  is a functor from the category  $gr(A)_R$  to the category  $gr(A)_L$ . Moreover, its derived functor  $\mathbf{R}\text{Hom}_A(-, A)$  gives an anti-equivalence between the derived categories of  $gr(A)_R$  and  $gr(A)_L$  (see [38], [39], [37]).

If we assume that the composition of the functor  $\text{Hom}_A(-, A)$  with the projection  $gr(A)_L \rightarrow qgr(A)_L$  factors through the projection  $gr(A)_R \rightarrow qgr(A)_R$ , then we obtain a functor from  $qgr(A)_R$  to  $qgr(A)_L$  which is denoted by  $\underline{\mathcal{H}om}(-, \mathcal{O})$ . This functor is not right exact and has right derived functors  $\underline{\mathcal{E}xt}^i(-, \mathcal{O})$ ,  $i > 0$ , from  $qgr(A)_R$  to  $qgr(A)_L$ .

For a noetherian regular algebra the functor  $\underline{\mathcal{H}om}(-, \mathcal{O})$  and its right derived functors exist. This follows from the fact that the functors  $\text{Ext}_A^i(-, A)$  send a finite dimensional module to a finite dimensional module (see condition (3) of Definition 4.10).

Moreover, in this case the functor  $\underline{\mathcal{H}om}(-, \mathcal{O})$  can be represented as the composition of the functor  $\Gamma : qgr(A)_R \rightarrow gr(A)_R$ , the functor  $\text{Hom}_A(-, A) : gr(A)_R \rightarrow gr(A)_L$ , and the projection  $\pi : gr(A)_L \rightarrow qgr(A)_L$ . This can be illustrated by the following commutative diagram:

$$(13) \quad \begin{array}{ccc} gr(A)_R & \xrightarrow{\text{Hom}_A(-, A)} & gr(A)_L \\ \pi \downarrow \uparrow \Gamma & & \downarrow \pi \\ qgr(A)_R & \xrightarrow{\underline{\mathcal{H}om}(-, \mathcal{O})} & qgr(A)_L \end{array}$$

For a noetherian regular algebra the functor  $\mathbf{R}\text{Hom}_A(-, A)$  is an anti-equivalence between the derived categories of  $gr(A)_R$  and  $gr(A)_L$  and takes complexes of finite dimensional modules over  $gr(A)_R$  to complexes of finite dimensional modules over  $gr(A)_L$ . This implies that the functor  $\mathbf{R}\underline{\mathcal{H}om}(-, \mathcal{O})$  gives an anti-equivalence between the derived categories of  $qgr(A)_R$  and  $qgr(A)_L$ . (Note that for derived functors  $\mathbf{R}\text{Hom}_A(-, A)$  and  $\mathbf{R}\underline{\mathcal{H}om}(-, \mathcal{O})$  there is also a commutative diagram like (13)).

The functors  $\underline{\mathcal{E}xt}^j(-, \mathcal{O})$  can be described more explicitly. Let  $M$  be an  $A$ -bimodule. Regarding it as a right module, we see that for any  $\mathcal{F} \in QGr(A)_R$  the groups  $\text{Ext}^i(\mathcal{F}, \tilde{M})$  have the structure of left  $A$ -modules. We can project them to  $QGr(A)_L$ . Thus each

bimodule  $M$  defines functors from  $QGr(A)_R$  to  $QGr(A)_L$ , which will be denoted by  $\pi\text{Ext}^i(-, \widetilde{M})$ .

Now, using  $\pi\Gamma = id$  and the commutativity of the diagram (13) for the derived functors  $\text{Ext}_A^j(-, A)$  and  $\underline{\mathcal{E}xt}^j(-, \mathcal{O})$ , we obtain isomorphisms

$$(14) \quad \underline{\mathcal{E}xt}^j(\mathcal{F}, \mathcal{O}) \cong \pi\text{Ext}_A^j(\Gamma(\mathcal{F}), A) \cong \pi\text{Ext}_{gr(A)}^j(\Gamma(\mathcal{F}), \bigoplus_{i \geq 0} A(i)) \cong \pi\text{Ext}^j(\mathcal{F}, \bigoplus_{i \geq 0} \mathcal{O}(i))$$

for any sheaf  $\mathcal{F} \in qgr(A)_R$ .

**Definition 5.4.** *We call a coherent sheaf  $\mathcal{F} \in qgr(A)_R$  locally free (or a bundle) if  $\underline{\mathcal{E}xt}^j(\mathcal{F}, \mathcal{O}) = 0$  for any  $j \neq 0$ .*

REMARK. In the commutative case this definition is equivalent to the usual definition of a locally free sheaf.

**Definition 5.5.** *The dual sheaf  $\underline{\mathcal{H}om}(\mathcal{F}, \mathcal{O}) \in qgr(A)_L$  will be denoted by  $\mathcal{F}^\vee$ .*

If  $\mathcal{F} \in qgr(A)_L$  is a bundle, then the dual sheaf  $\mathcal{F}^\vee$  is a bundle in  $qgr(A)_L$ , because  $\mathbf{R}\underline{\mathcal{H}om}(\mathcal{F}^\vee, \mathcal{O}) = \mathcal{F}$  in the derived category, and  $\underline{\mathcal{E}xt}^j(\mathcal{F}^\vee, \mathcal{O}) = 0$  for  $j \neq 0$ .

Thus we have a good definition of locally free sheaves on  $\mathbb{P}_\hbar^3$  and  $\mathbb{P}_\hbar^2$ . Since the derived functor  $\mathbf{R}\underline{\mathcal{H}om}(-, \mathcal{O})$  gives an anti-equivalence between the derived categories of  $qgr(A)_R$  and  $qgr(A)_L$ , there is an isomorphism:

$$(15) \quad \text{Hom}(\mathcal{F}, \mathcal{G}) \cong \text{Hom}(\mathcal{G}^\vee, \mathcal{F}^\vee)$$

for any two bundles  $\mathcal{F}$  and  $\mathcal{G}$  on  $\mathbb{P}_\hbar^3$  or  $\mathbb{P}_\hbar^2$ .

## 6. BUNDLES ON $\mathbb{P}_\hbar^2$

**6.1. Bundles on  $\mathbb{P}_\hbar^2$  with a trivialization on the commutative line.** In this section we study bundles on  $\mathbb{P}_\hbar^2$ . By definition, a bundle is an object  $\mathcal{E} \in \text{coh}(\mathbb{P}_\hbar^2)$  satisfying the additional condition  $\underline{\mathcal{E}xt}^i(\mathcal{E}, \mathcal{O}) = 0$  for all  $i > 0$  (see (5.4)).

The noncommutative plane  $\mathbb{P}_\hbar^2$  contains the commutative projective line  $l \cong \mathbb{P}^1$  given by the equation  $w_3 = 0$ . If  $M$  is a  $PP_\hbar$ -module, then the quotient module  $M/Mw_3$  is a  $PP_\hbar/\langle w_3 \rangle$ -module. This gives a functor  $\text{coh}(\mathbb{P}_\hbar^2) \rightarrow \text{coh}(\mathbb{P}^1)$ ,  $\mathcal{F} \mapsto \mathcal{F}|_l$ . The sheaf  $\mathcal{F}|_l$  is referred to as the restriction of  $\mathcal{F}$  to the line  $l$ .

**Lemma 6.1.** *If  $\mathcal{F}$  is a bundle, there is an exact sequence:*

$$(16) \quad 0 \longrightarrow \mathcal{F}(-1) \xrightarrow{\cdot w_3} \mathcal{F} \longrightarrow \mathcal{F}|_l \longrightarrow 0.$$

*Proof.* To prove this we only need to show that multiplication by  $w_3$  is a monomorphism. If  $\mathcal{F}$  is a bundle, it can be embedded into a direct sum  $\bigoplus_{i=1}^s \mathcal{O}(k_i)$ , because by ampleness the dual bundle  $\mathcal{F}^\vee$  is covered by a direct sum of line bundles. Now, since the morphism  $\mathcal{O}(k_i - 1) \xrightarrow{w_3} \mathcal{O}(k_i)$  is mono for any  $i$ , the same is true for the morphism  $\mathcal{F}(-1) \xrightarrow{w_3} \mathcal{F}$ .  $\square$

**Lemma 6.2.** *Let  $\mathcal{E}$  be a bundle on  $\mathbb{P}_h^2$  such that its restriction  $\mathcal{E}|_l$  to the commutative line  $l$  is isomorphic to a trivial bundle  $\mathcal{O}_l^{\oplus r}$ . Then*

$$H^0(\mathbb{P}_h^2, \mathcal{E}(-1)) = H^0(\mathbb{P}_h^2, \mathcal{E}(-2)) = H^2(\mathbb{P}_h^2, \mathcal{E}(-1)) = H^2(\mathbb{P}_h^2, \mathcal{E}(-2)) = 0.$$

*Proof.* We have the following exact sequence in the category  $\text{coh}(\mathbb{P}_h^2)$ :

$$(17) \quad 0 \longrightarrow \mathcal{E}(-2) \longrightarrow \mathcal{E}(-1) \longrightarrow \mathcal{E}(-1)|_l \longrightarrow 0.$$

Since  $\mathcal{E}(-1)|_l \cong \mathcal{O}_l(-1)^{\oplus r}$ , we have  $H^0(\mathcal{E}(-1)|_l) = 0$ .

Assume that  $\mathcal{E}(-1)$  has a nontrivial section. Then  $\mathcal{E}(-2)$  has a nontrivial section too. For the same reason  $\mathcal{E}(-3)$  has a nontrivial section, and so on. Thus for any  $n < 0$  the bundle  $\mathcal{E}(-n)$  has a nontrivial section. By (15) we have isomorphisms:

$$H^0(\mathcal{E}(-n)) \cong \text{Hom}(\mathcal{O}(n), \mathcal{E}) \cong \text{Hom}(\mathcal{E}^\vee, \mathcal{O}(-n)).$$

On the other hand, by Corollary 5.2 the last group is trivial for  $n \gg 0$ . Hence  $H^0(\mathcal{E}(-n)) = 0$  for all  $n \gg 0$ , and consequently  $H^0(\mathcal{E}(-2)) = H^0(\mathcal{E}(-1)) = 0$ .

Further, assume that  $H^2(\mathcal{E}(-2))$  is nontrivial. Since  $H^1(\mathcal{E}(i)|_l) = 0$  for all  $i \geq -1$  we have from the exact sequence (16) with  $\mathcal{F} = \mathcal{E}(i)$  that  $H^2(\mathcal{E}(i))$  is nontrivial too for all  $i \geq -1$ . But this contradicts Corollary 5.3. Therefore  $H^2(\mathcal{E}(-2)) = H^2(\mathcal{E}(-1)) = 0$ . This completes the proof.  $\square$

**6.2. Monads on  $\mathbb{P}_h^2$  and  $\mathbb{P}_h^3$ .** As in the commutative case, a non-degenerate monad on  $\mathbb{P}_h^2$  or  $\mathbb{P}_h^3$  is a complex over  $\text{coh}(\mathbb{P}_h^2)$

$$0 \longrightarrow H \otimes \mathcal{O}(-1) \xrightarrow{m} K \otimes \mathcal{O} \xrightarrow{n} L \otimes \mathcal{O}(1) \longrightarrow 0$$

for which the map  $n$  is an epimorphism and  $m$  is a monomorphism. (Note that there is another more restrictive definition of a monad, according to which the dual map  $(m)^*$  has to be an epimorphism, see [29]). The coherent sheaf

$$E = \mathbf{Ker}(n)/\mathbf{Im}(m)$$

is called the cohomology of a monad. A morphism between two monads is a morphism of complexes. The following lemma is proved in [29](Lemma 4.1.3) in the commutative case, but the proof is categorical and applies to the noncommutative case as well.

**Lemma 6.3.** *Let  $X$  be either  $\mathbb{P}_{\hbar}^2$  or on  $\mathbb{P}_{\hbar}^3$ , and let  $E$  and  $E'$  be the cohomology bundles of two monads*

$$\begin{aligned} M : 0 &\longrightarrow H \otimes \mathcal{O}(-1) \xrightarrow{m} K \otimes \mathcal{O} \xrightarrow{n} L \otimes \mathcal{O}(1) \longrightarrow 0, \\ M' : 0 &\longrightarrow H' \otimes \mathcal{O}(-1) \xrightarrow{m'} K' \otimes \mathcal{O} \xrightarrow{n'} L' \otimes \mathcal{O}(1) \longrightarrow 0 \end{aligned}$$

on  $X$ . Then the natural mapping

$$\mathrm{Hom}(M, M') \longrightarrow \mathrm{Hom}(E, E')$$

is bijective.

The proof is based on the fact that

$$\mathrm{Ext}^j(\mathcal{O}, \mathcal{O}(-1)) = \mathrm{Ext}^j(\mathcal{O}(1), \mathcal{O}(-1)) = \mathrm{Ext}^j(\mathcal{O}(1), \mathcal{O}) = 0$$

for all  $j$  (see [29], Lemma 4.1.3).

**6.3. Non-degeneracy conditions.** In the definition of a monad we require that the map  $n$  be an epimorphism. In the commutative case this condition must be verified pointwise. In the noncommutative case the situation is simpler in some sense, because the complement of the commutative line  $l$  does not have points.

**Lemma 6.4.** *If the restriction of a sheaf  $\mathcal{F} \in \mathrm{coh}(\mathbb{P}_{\hbar}^2)$  to the projective line  $l$  is the zero object, then  $\mathcal{F}$  is also the zero object.*

*Proof.* Let  $M$  be a finitely generated graded  $PP_{\hbar}$ -module such that  $\mathcal{F} \cong \widetilde{M}$ . Consider an exact sequence:

$$M \xrightarrow{w_3} M(1) \longrightarrow N \longrightarrow 0.$$

Since  $\widetilde{N} = \mathcal{F}(1)|_l = 0$ , the module  $N$  is finite dimensional. This implies that for  $i \gg 0$  the map  $M_i \xrightarrow{w_3} M_{i+1}$  is surjective. Moreover, these maps are isomorphisms for  $i \gg 0$ , because all  $M_i$  are finite dimensional vector spaces. Let us identify all  $M_i$  for  $i \gg 0$  with respect to these isomorphisms. Using the  $A$ -module structure on  $M$ , we obtain a representation of the Weyl algebra  $\mathrm{T}(X, Y)/\langle [X, Y] = 2\hbar \rangle$  on the vector space  $M_i$ . But it is well known that the Weyl algebra does not have finite dimensional representations. Thus  $M_i = 0$  for all  $i \gg 0$ , and  $M$  is finite dimensional. Therefore  $\mathcal{F} = 0$ .  $\square$

The following corollary is an immediate consequence of the Lemma.

**Corollary 6.5.** *Let  $f : \mathcal{F} \longrightarrow \mathcal{G}$  be a morphism in  $\mathrm{coh}(\mathbb{P}_{\hbar}^2)$ . Suppose its restriction  $\bar{f} : \mathcal{F}|_l \longrightarrow \mathcal{G}|_l$  is an epimorphism. Then  $f$  is an epimorphism too.*

**6.4. From the resolution of the diagonal to a monad.** Let  $M$  be an  $A$ -bimodule. Regarding it as a left module, we see that for any  $\mathcal{F} \in QGr(A)_L$  the groups  $\text{Ext}^i(\mathcal{F}, \widetilde{M})$  have the structure of right  $A$ -modules. We can project them to  $QGr(A)_R$ . Thus each bimodule  $M$  defines functors  $\pi\text{Ext}^i(-, \widetilde{M})$  from  $QGr(A)_L$  to  $QGr(A)_R$ .

Let  $\mathcal{E}$  be a bundle on  $\mathbb{P}_\hbar^2$  such that its restriction to the line  $l$  is a trivial bundle. Let us consider the bundle  $\mathcal{E}^\vee(1) \in qgr(PP_\hbar)_L$  and the resolution of the diagonal  $\mathcal{K}(PP_\hbar)$ , which has only three terms:

$$\{0 \longrightarrow PP_\hbar(-1) \otimes PP_\hbar(-2) \longrightarrow \Omega^1(1) \otimes PP_\hbar(-1) \longrightarrow PP_\hbar \otimes PP_\hbar\} \longrightarrow \Delta.$$

The resolution of the diagonal is a complex of bimodules. It induces a complex  $\widetilde{\mathcal{K}}$  over  $QGr(PP_\hbar)_L$ :

$$(18) \quad \{0 \longrightarrow \mathcal{O}(-1) \otimes PP_\hbar(-2) \longrightarrow \Omega^1(1) \otimes PP_\hbar(-1) \longrightarrow \mathcal{O} \otimes PP_\hbar\} \longrightarrow \widetilde{\Delta}$$

where  $\Omega^1$  is a sheaf on  $\mathbb{P}_\hbar^2$  corresponding to the  $PP_\hbar$ -module  $\Omega^1$ .

As described above, each  $A$ -bimodule  $M$  gives the functors  $\pi\text{Ext}^i(-, \widetilde{M})$  from  $QGr(A)_L$  to  $QGr(A)_R$ . In particular, each object of the resolution of the diagonal induces such functors.

First we calculate these functors for the object  $\widetilde{\Delta}$ . Note that the object  $\widetilde{\Delta}$  coincides with  $\bigoplus_{i \geq 0} \mathcal{O}(i)$ . Hence by (14) we have

$$\pi\text{Ext}^j(\mathcal{E}^\vee(1), \widetilde{\Delta}) = 0$$

if  $j > 0$ , while  $\pi\text{Ext}^0(\mathcal{E}^\vee(1), \widetilde{\Delta}) \cong \mathcal{E}(-1)$ .

The resolution of the diagonal (18) gives us a spectral sequence with the  $E_1$  term

$$E_1^{pq} = \pi\text{Ext}^q(\mathcal{E}^\vee(1), \widetilde{\mathcal{K}}_{-p}) \implies \pi\text{Ext}^{p+q}(\mathcal{E}^\vee(1), \widetilde{\Delta}),$$

which converges to

$$E_\infty^i = \begin{cases} \mathcal{E}(-1) & \text{if } i = 0 \\ 0 & \text{otherwise} \end{cases}$$

Now we describe all terms  $E_1^{pq}$  of this spectral sequence. First we have

$$\pi\text{Ext}^j(\mathcal{E}^\vee(1), \mathcal{O} \otimes PP_\hbar) \cong \text{Ext}^j(\mathcal{E}^\vee(1), \mathcal{O}) \otimes \widetilde{PP_\hbar} \cong \text{Ext}^j(\mathcal{E}^\vee(1), \mathcal{O}) \otimes \mathcal{O} \cong H^j(\mathbb{P}_\hbar^2, \mathcal{E}(-1)) \otimes \mathcal{O}.$$

By Lemma 6.2, these groups are trivial for  $j \neq 1$ . For the same reason we have

$$\pi\text{Ext}^j(\mathcal{E}^\vee(1), \mathcal{O}(-1) \otimes PP_\hbar(-2)) = H^j(\mathbb{P}_\hbar^2, \mathcal{E}(-2)) \otimes \mathcal{O}(-2) = 0$$

for  $j \neq 1$  and

$$\pi\text{Ext}^1(\mathcal{E}^\vee(1), \mathcal{O}(-1) \otimes PP_\hbar(-2)) \cong H^1(\mathbb{P}_\hbar^2, \mathcal{E}(-2)) \otimes \mathcal{O}(-2).$$

Now let us consider the functors which are associated with the object  $\Omega^1(1) \otimes PP_{\hbar}(-1)$ . We have

$$\pi \text{Ext}^j(\mathcal{E}^\vee(1), \Omega^1(1) \otimes PP_{\hbar}(-1)) \cong \text{Ext}^j(\mathcal{E}^\vee, \Omega^1) \otimes \mathcal{O}(-1).$$

It follows from the Koszul complex that the sheaf  $\Omega^1$  can be included in two exact sequences:

$$\begin{aligned} 0 &\longrightarrow \Omega^1 \longrightarrow \mathcal{O}(-1) \otimes PP_{\hbar 1} \longrightarrow \mathcal{O} \longrightarrow 0, \\ 0 &\longrightarrow \mathcal{O}(-3) \longrightarrow \mathcal{O}(-2) \otimes (PP_{\hbar 1})^* \longrightarrow \Omega^1 \longrightarrow 0. \end{aligned}$$

Applying the functor  $\text{Hom}(\mathcal{E}^\vee, -)$  to the first sequence and taking into account that  $\text{Hom}(\mathcal{E}^\vee, \mathcal{O}(-1)) = 0$ , we obtain  $\text{Hom}(\mathcal{E}^\vee, \Omega^1) = 0$ . Similarly, we deduce from the second sequence that  $\text{Ext}^2(\mathcal{E}^\vee, \Omega^1) = 0$ , because  $\text{Ext}^2(\mathcal{E}^\vee, \mathcal{O}(-2)) = 0$ . This implies that the object  $\pi \text{Ext}^j(\mathcal{E}^\vee(1), \Omega^1(1) \otimes PP_{\hbar}(-1))$  is trivial for all  $j \neq 1$ .

Thus our spectral sequence is nothing more than the complex

$$\pi \text{Ext}^1(\mathcal{E}^\vee(1), \tilde{\mathcal{K}}_2) \longrightarrow \pi \text{Ext}^1(\mathcal{E}^\vee(1), \tilde{\mathcal{K}}_1) \longrightarrow \pi \text{Ext}^1(\mathcal{E}^\vee(1), \tilde{\mathcal{K}}_0),$$

which is isomorphic to the complex

$$H^1(\mathbb{P}_{\hbar}^2, \mathcal{E}(-2)) \otimes \mathcal{O}(-2) \longrightarrow \text{Ext}^1(\mathcal{E}^\vee, \Omega^1) \otimes \mathcal{O}(-1) \longrightarrow H^1(\mathbb{P}_{\hbar}^2, \mathcal{E}(-1)) \otimes \mathcal{O}.$$

It has only one cohomology which coincides with  $\mathcal{E}(-1)$ .

**Theorem 6.6.** *Let  $\mathcal{E}$  be a bundle on  $\mathbb{P}_{\hbar}^2$  such that its restriction to the commutative line  $l$  is isomorphic to the trivial bundle  $\mathcal{O}_l^{\oplus r}$ . Then  $\mathcal{E}$  is the cohomology of a monad*

$$0 \longrightarrow H \otimes \mathcal{O}(-1) \xrightarrow{m} K \otimes \mathcal{O} \xrightarrow{n} L \otimes \mathcal{O}(1) \longrightarrow 0$$

with  $H = H^1(\mathbb{P}_{\hbar}^2, \mathcal{E}(-2))$ ,  $L = H^1(\mathbb{P}_{\hbar}^2, \mathcal{E}(-1))$ , and such monad is unique up to an isomorphism. Moreover, in this case the vector spaces  $H$  and  $L$  have the same dimension.

*Proof.* The existence of such a monad was proved above. The uniqueness follows from Lemma 6.3. The equality of dimensions of  $H$  and  $L$  follows immediately from the exact sequence (17).  $\square$

**6.5. Barth description of monads.** Now following Barth [8], we give a description of the moduli space of vector bundles on  $\mathbb{P}_{\hbar}^2$  trivial on the line  $l$  in terms of linear algebra (see also [15]).

Denote by  $\mathfrak{M}_{\hbar}(r, 0, k)$  the moduli space of bundles on the noncommutative  $\mathbb{P}_{\hbar}^2$  trivial on the line  $l$  and with a fixed trivialization there (i.e. with a fixed isomorphism  $\mathcal{E}|_l \cong \mathcal{O}_l^{\oplus r}$ ). Let  $\dim H^1(\mathbb{P}_{\hbar}^2, \mathcal{E}(-1)) = k$ . As in the commutative case, the numbers  $r, 0, k$  can be regarded as the rank, first Chern class, and second Chern class of  $\mathcal{E}$ , respectively.

The following theorem gives a description of this moduli space which is similar to the description given by Barth in the commutative case.

**Theorem 6.7.** *Let  $\{(b_1, b_2; j, i)\}$  be the set of quadruples of matrices  $b_1, b_2 \in M_{k \times k}(\mathbb{C})$ ,  $j \in M_{r \times k}(\mathbb{C})$ ,  $i \in M_{k \times r}(\mathbb{C})$ , which satisfy the condition*

$$[b_1, b_2] + ij + 2\hbar \cdot 1_{k \times k} = 0.$$

*Then the space  $\mathfrak{M}_\hbar(r, 0, k)$  is the quotient of this set with respect to the following free action of  $\mathrm{GL}(k, \mathbb{C})$  :*

$$b_i \mapsto gb_i g^{-1}, \quad j \mapsto jg^{-1}, \quad i \mapsto gi, \quad \text{where } g \in \mathrm{GL}(k, \mathbb{C}).$$

*Proof.* Let  $\mathcal{E}$  be a bundle on  $\mathbb{P}_\hbar^2$  trivial on the line  $l$ . We showed above that any such bundle comes from a monad unique up to an isomorphism. Conversely, suppose we have a monad

$$(19) \quad 0 \longrightarrow H \otimes \mathcal{O}(-1) \xrightarrow{m} K \otimes \mathcal{O} \xrightarrow{n} L \otimes \mathcal{O}(1) \longrightarrow 0$$

with  $\dim H = \dim L = k$  such that its restriction to the line  $l$  is a monad with the cohomology  $\mathcal{O}_l^{\oplus r}$ . Then the cohomology of this monad is a bundle on  $\mathbb{P}_\hbar^2$  which belongs to  $\mathfrak{M}_\hbar(r, 0, k)$ . Indeed, the cohomologies of the dual complex

$$0 \longrightarrow \mathcal{O}(-1) \otimes L^* \xrightarrow{n^*} \mathcal{O} \otimes K^* \xrightarrow{m^*} \mathcal{O}(1) \otimes H^* \longrightarrow 0$$

coincide with  $\underline{\mathrm{Hom}}(\mathcal{E}, \mathcal{O})$  and  $\underline{\mathrm{Ext}}^1(\mathcal{E}, \mathcal{O})$ . Hence, to prove that  $\mathcal{E}$  is a bundle, it is sufficient to show that the dual complex is a monad too, i.e. that the map  $m^*$  is an epimorphism. The restriction of the dual complex to  $l$  is a monad which is dual to the restriction of the monad (19) to  $l$ . Hence the restriction of  $m^*$  on  $l$  is an epimorphism. Then, by Lemma 6.5,  $m^*$  is an epimorphism as well. Thus to describe the moduli space  $\mathfrak{M}_\hbar(r, 0, k)$  we have to describe the space of all monads (19) modulo isomorphisms preserving trivialization on  $l$ .

Consider a monad

$$0 \longrightarrow H \otimes \mathcal{O}(-1) \xrightarrow{m} K \otimes \mathcal{O} \xrightarrow{n} L \otimes \mathcal{O}(1) \longrightarrow 0$$

with  $\dim H = \dim L = k$  and  $\dim K = 2k + r$ . Denote by  $\mathcal{E}$  its cohomology bundle.

The maps  $m$  and  $n$  can be regarded as elements of  $H^* \otimes K \otimes W$  and  $K^* \otimes L \otimes W$ , respectively, where  $W = H^0(\mathbb{P}_\hbar^2, \mathcal{O}(1))$  is the vector space spanned by  $w_1, w_2, w_3$ . The maps  $m$  and  $n$  can be written as

$$m_1 w_1 + m_2 w_2 + m_3 w_3, \quad n_1 w_1 + n_2 w_2 + n_3 w_3,$$

where  $m_i : H \rightarrow K$  and  $n_i : K \rightarrow L$  are constant linear maps.

Let us restrict the monad to the line  $l$ . The monadic condition  $nm = 0$  implies now:

$$n_1 m_2 + n_2 m_1 = 0, \quad n_1 m_1 = 0, \quad n_2 m_2 = 0.$$

Moreover, since the restriction of  $\mathcal{E}$  to  $l$  is trivial, the composition  $n_1 m_2$  is an isomorphism (see [29], Lemma 4.2.3). We can choose bases for  $H, K, L$  so that  $n_1 m_2 = 1_{k \times k}$  (the identity matrix) and

$$m_1 = \begin{pmatrix} 1_{k \times k} \\ 0_{k \times k} \\ 0_{r \times k} \end{pmatrix}, \quad m_2 = \begin{pmatrix} 0_{k \times k} \\ 1_{k \times k} \\ 0_{r \times k} \end{pmatrix}, \quad n_1 = \begin{pmatrix} 0_{k \times k} & 1_{k \times k} & 0_{k \times r} \end{pmatrix}, \quad n_2 = \begin{pmatrix} -1_{k \times k} & 0_{k \times k} & 0_{k \times r} \end{pmatrix}.$$

Using the equations  $n_3 m_1 + n_1 m_3 = 0$  and  $n_3 m_2 + n_2 m_3 = 0$  we can write:

$$m_3 = \begin{pmatrix} b_1 \\ b_2 \\ j \end{pmatrix}, \quad n_3 = \begin{pmatrix} -b_2 & b_1 & i \end{pmatrix}.$$

Now the monadic condition  $nm = 0$  can be written as:

$$(n_3 m_3) \cdot w_3^2 + 1_{k \times k} \cdot [w_1, w_2] = 0.$$

Therefore we obtain the following matrix equation:

$$[b_1, b_2] + ij + 2\hbar \cdot 1_{k \times k} = 0.$$

Note that the last  $r$  basis vectors of  $K$  give us a trivialization of the restriction of  $\mathcal{E}$  to the line  $l$ . It is easy to check that any isomorphism of a monad which preserves trivialization on  $l$  and the choice of the bases of  $H, K, L$  made above has the form

$$b_i \mapsto g b_i g^{-1}, \quad j \mapsto j g^{-1}, \quad i \mapsto g i, \quad \text{where } g \in \text{GL}(k, \mathbb{C}).$$

This proves the theorem.  $\square$

## 7. THE NONCOMMUTATIVE VARIETY $\mathbb{P}_\hbar^3$ AS A TWISTOR SPACE

**7.1. Real structures.** A  $*$ -algebra is, by definition, an algebra over  $\mathbb{C}$  with an anti-linear anti-homomorphism  $*$  satisfying  $*^2 = id$ . A  $*$ -structure on a (graded) algebra is regarded as a real structure on the corresponding (projective) noncommutative variety.

Let us introduce real structures on the complex varieties  $\mathbb{C}_\hbar^4$  and  $\mathbb{Q}_\hbar^4$  defined in section 3. Assume that in (6), (7) the skew-symmetric matrix  $\theta$  is purely imaginary and  $\hbar$  is real. Then there is a unique  $*$ -structure on the algebra  $A(\mathbb{C}_\hbar^4)$  such that  $x_i^* = x_i$ . We denote the corresponding noncommutative variety by  $\mathbb{R}_\hbar^4$ .



Assume in addition that the symmetric matrix  $G$  in (7) is real and positive definite. There is a unique  $*$ -structure on the algebra  $Q_{\hbar}$  such that  $X_i^* = X_i$ ,  $D^* = D$ , and  $T^* = T$ . The corresponding noncommutative real variety will be called the noncommutative sphere and denoted by  $S_{\hbar}^4$ . The embedding of  $C_{\hbar}^4$  into  $Q_{\hbar}^4$  induces an embedding  $\mathbb{R}_{\hbar}^4 \hookrightarrow S_{\hbar}^4$ . Recall that the complement of  $C_{\hbar}^4$  in  $Q_{\hbar}^4$  is a commutative quadratic cone  $\sum_{kl} G^{kl} X_k X_l = 0$  which has only one real point. Thus  $S_{\hbar}^4$  can be regarded as a one-point compactification of  $\mathbb{R}_{\hbar}^4$ .

By a linear change of basis one can bring the pair  $(G, \theta)$  to the standard form

$$(20) \quad G = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}, \quad \theta = \sqrt{-1} \begin{pmatrix} 0 & a & 0 & 0 \\ -a & 0 & 0 & 0 \\ 0 & 0 & 0 & b \\ 0 & 0 & -b & 0 \end{pmatrix}.$$

Furthermore, since  $\hbar$  and  $\theta$  enter only in the combination  $\hbar \cdot \theta$ , and we assume that  $a + b \neq 0$ , we can set  $a + b = 1$  without loss of generality.

**7.2. Realification of  $\mathbb{P}_{\hbar}^3$ .** Recall that the noncommutative projective space  $\mathbb{P}_{\hbar}^3$  corresponds to the algebra  $PS_{\hbar}$  with generators  $z_i, i = 1, 2, 3, 4$ , and relations (9). Consider an algebra  $\widehat{PS}_{\hbar}$  with generators  $z_i, \bar{z}_i, i = 1, 2, 3, 4$ , and relations

$$(21) \quad \begin{aligned} [z_1, z_2] &= 2\hbar(a+b)z_3z_4, & [z_1, \bar{z}_1] &= 2\hbar bz_3\bar{z}_3 - 2\hbar az_4\bar{z}_4, & [z_1, \bar{z}_2] &= 0, \\ [\bar{z}_1, \bar{z}_2] &= -2\hbar(a+b)\bar{z}_3\bar{z}_4, & [z_2, \bar{z}_2] &= 2\hbar az_3\bar{z}_3 - 2\hbar bz_4\bar{z}_4, & [z_2, \bar{z}_1] &= 0, \\ [z_i, z_j] &= [z_i, \bar{z}_j] = [\bar{z}_i, z_j] = [\bar{z}_i, \bar{z}_j] = 0 & \text{for all } i = 3, 4; j = 1, 2, 3, 4 \end{aligned}$$

There is a unique  $*$ -structure on this algebra such that  $z_i^* = \bar{z}_i$ ,  $\bar{z}_i^* = z_i$ . We denote the corresponding real variety  $\mathbb{P}_{\hbar}^3(\mathbb{R})$ . This variety can be considered a realification of  $\mathbb{P}_{\hbar}^3$ . **REMARK.** In contrast to the commutative situation, a noncommutative complex variety in general has many different realifications. We have an ambiguity in the choice of relations involving both  $z_i$  and  $\bar{z}_j$ . The realification (21) is distinguished by the fact that it is the twistor space of the noncommutative sphere  $S_{\hbar}^4$ , as explained below.

In the commutative case there is a map from  $\mathbb{P}^3(\mathbb{R})$  to the sphere  $S^4$  which is a  $\mathbb{P}^1$  fibration. The corresponding  $\mathbb{P}^1$  bundle is the projectivization of a spinor bundle on  $S^4$ . This map is known as the Penrose map. In the noncommutative case we have a similar picture. The analogue of the Penrose map is a map  $\Pi : \mathbb{P}_{\hbar}^3(\mathbb{R}) \longrightarrow S_{\hbar}^4$  which is associated

with the homomorphism of  $*$ -algebras  $Q_{\hbar} \longrightarrow \widehat{PS}_{\hbar}$  :

$$\begin{aligned} X_1 &\mapsto -\frac{\sqrt{-1}}{2} (z_1 \bar{z}_4 - \bar{z}_1 z_4 - \bar{z}_2 z_3 + z_2 \bar{z}_3), & D &\mapsto -\frac{1}{2} (z_1 \bar{z}_1 + \bar{z}_1 z_1 + z_2 \bar{z}_2 + \bar{z}_2 z_2), \\ X_2 &\mapsto \frac{1}{2} (z_1 \bar{z}_4 + \bar{z}_1 z_4 - \bar{z}_2 z_3 - z_2 \bar{z}_3), & T &\mapsto - (z_3 \bar{z}_3 + z_4 \bar{z}_4), \\ X_3 &\mapsto -\frac{\sqrt{-1}}{2} (\bar{z}_1 z_3 - z_1 \bar{z}_3 + z_2 \bar{z}_4 - \bar{z}_2 z_4), \\ X_4 &\mapsto \frac{1}{2} (z_1 \bar{z}_3 + \bar{z}_1 z_3 + \bar{z}_2 z_4 + z_2 \bar{z}_4). \end{aligned}$$

Note that for  $\hbar = 0$  we obtain the homomorphism of commutative algebras which corresponds to the usual Penrose map. This means that  $\mathbb{P}_{\hbar}^3(\mathbb{R})$  is the twistor space of  $\mathbb{S}_{\hbar}^4$ .

The variety  $\mathbb{P}_{\hbar}^3(\mathbb{R})$  is a twistor space in yet another sense. For the commutative  $\mathbb{R}^4$  the complex structures compatible with the symmetric bilinear form  $G$  and orientation are parametrized by points of a  $\mathbb{P}^1$ . This remains true in the noncommutative case. A complex structure (resp. orientation) on  $\mathbb{R}_{\hbar}^4$  is defined as a complex structure (resp. orientation) on the real vector space  $U$  spanned by  $x_1, \dots, x_4$ . We will choose an orientation on  $U$  and require that the complex structure be compatible with it. All such complex structures are parametrized by points of a  $\mathbb{P}^1$ .

Recall now that  $\mathbb{P}_{\hbar}^3$  is a pencil of noncommutative projective planes passing through the commutative line. Let us pick any one of them. The realification of  $\mathbb{P}_{\hbar}^3$  defined above induces a realification of the noncommutative projective plane. It is easy to see that the complement of the commutative line  $w_3 = \bar{w}_3 = 0$  in the realified projective plane is isomorphic to  $\mathbb{R}_{\hbar}^4$ . Furthermore, the complement carries a natural complex structure defined by

$$w_3^{-1} w_i \mapsto \sqrt{-1} w_3^{-1} w_i, \quad \bar{w}_3^{-1} \bar{w}_i \mapsto -\sqrt{-1} \bar{w}_3^{-1} \bar{w}_i, \quad i = 1, 2.$$

The Penrose map induces an identification between the complement and  $\mathbb{R}_{\hbar}^4 \subset \mathbb{S}_{\hbar}^4$ , and therefore induces a complex structure on the latter. Varying the noncommutative projective plane, one obtains all possible complex structures on  $\mathbb{R}_{\hbar}^4$  compatible with a particular orientation. This is completely analogous to the commutative case.

### 7.3. Connection between sheaves on commutative and noncommutative planes.

In this subsection we are going to connect the moduli space  $\mathfrak{M}_{\hbar}(r, 0, k)$  of bundles on  $\mathbb{P}_{\hbar}^2$  with a trivialization on the line  $l$  with the moduli space  $\mathfrak{M}(r, 0, k)$  of torsion free sheaves on the commutative  $\mathbb{P}^2$  with a trivialization on a fixed line. The bridge between bundles on  $\mathbb{P}_{\hbar}^2$  and torsion free sheaves on  $\mathbb{P}^2$  is provided by the twistor variety  $\mathbb{P}_{\hbar}^3$ . This gives a geometrical interpretation of Nakajima's results (the description of the moduli space  $\mathfrak{M}(r, 0, k)$  by the deformed ADHM data [27, 26]). We will construct a hyperkähler manifold  $\mathcal{M}$  parametrizing certain complexes on  $\mathbb{P}_{\hbar}^3$  which is isomorphic to  $\mathfrak{M}(r, 0, k)$

(which is also a hyperkähler manifold [27]). The isomorphism is given by the restriction of complexes to one of the commutative  $\mathbb{P}^2$ 's. On the other hand, the restriction of complexes to a noncommutative plane  $\mathbb{P}_\hbar^2$  yields an isomorphism between  $\mathcal{M}$  with a particular choice of complex structure and the moduli space  $\mathfrak{M}_\hbar(r, 0, k)$ . Thus  $\mathfrak{M}_\hbar(r, 0, k)$  can be obtained from  $\mathfrak{M}(r, 0, k)$  by a rotation of complex structure.

Consider complexes  $\mathcal{C}$  on  $\mathbb{P}_\hbar^3$  of the form

$$(22) \quad 0 \longrightarrow H \otimes \mathcal{O}(-1) \xrightarrow{M} K \otimes \mathcal{O} \xrightarrow{N} L \otimes \mathcal{O}(1) \longrightarrow 0$$

with  $\dim H = \dim L = k$ ,  $\dim K = 2k + r$ , which satisfies the condition that its restriction to the line  $l$  has only one cohomology which is a trivial bundle (with a fixed trivialization). This condition implies that  $M$  is a monomorphism. Note that  $N$  is not an epimorphism in general, so (22) is not a monad. But the restriction of the complex (22) to any noncommutative plane is a monad by Corollary 6.5. Thus  $N$  can fail to be surjective only on the commutative planes  $z_3 = 0$  and  $z_4 = 0$ .

Now we introduce a real structure on  $\mathbb{P}_\hbar^3$  (this is different from the real structure on the realification of  $\mathbb{P}_\hbar^3$  defined above). Assume that  $\hbar$  is a real number. Consider an anti-linear anti-homomorphism  $\bar{\mathcal{J}}$  of  $PS_\hbar$  defined by

$$\bar{\mathcal{J}}(z_1) = z_2, \quad \bar{\mathcal{J}}(z_2) = -z_1, \quad \bar{\mathcal{J}}(z_3) = z_4, \quad \bar{\mathcal{J}}(z_4) = -z_3, \quad \bar{\mathcal{J}}(\lambda) = \bar{\lambda}, \quad \lambda \in \mathbb{C}.$$

Thus  $\bar{\mathcal{J}}$  is a homomorphism of  $\mathbb{R}$ -algebras from  $PS_\hbar$  to the opposite algebra  $PS_\hbar^{op}$ . (The notation  $\bar{\mathcal{J}}$  is used by analogy with the commutative case, where this anti-homomorphism is a composition of a complex structure  $J$  with complex conjugation [15].)

The anti-homomorphism  $\bar{\mathcal{J}}$  induces a functor  $\bar{\mathcal{J}}^*$  from  $qgr(PS_\hbar)_R$  to  $qgr(PS_\hbar^{op})_R$ . The latter category is naturally identified with the category  $qgr(PS_\hbar)_L$ . Using this identification we can consider the composition of  $\bar{\mathcal{J}}^*$  with the dualization functor  $\underline{\text{Hom}}(-, \mathcal{O})$  as a functor from  $qgr(PS_\hbar)_R$  to itself. For any bundle  $\mathcal{E}$  we denote by  $\bar{\mathcal{J}}^*(\mathcal{E})^\vee$  its image under this functor. The functor can be extended to complexes of bundles. It takes the complex  $\mathcal{C}$  (22) to the complex  $\bar{\mathcal{J}}^*(\mathcal{C})^\vee$

$$0 \longrightarrow \bar{L}^* \otimes \mathcal{O}(-1) \xrightarrow{\bar{\mathcal{J}}^*(N)^\vee} \bar{K}^* \otimes \mathcal{O} \xrightarrow{\bar{\mathcal{J}}^*(M)^\vee} \bar{H}^* \otimes \mathcal{O}(1) \longrightarrow 0.$$

Let us consider complexes  $\mathcal{C}$  on  $\mathbb{P}_\hbar^3$  with an isomorphism

$$(23) \quad \bar{\mathcal{J}}^*(\mathcal{C})^\vee \cong \mathcal{C}$$

and trivialization on the line  $l$ . Then the space  $K$  acquires a hermitian metric and  $L$  becomes isomorphic to  $\bar{H}^*$ . The reasoning of section 6 shows that we can represent the

maps  $M$  and  $N$  as

$$M_1 z_1 + M_2 z_2 + M_3 z_3 + M_4 z_4, \quad N_1 z_1 + N_2 z_2 + N_3 z_3 + N_4 z_4,$$

where  $M_i$  and  $N_i$  are constant maps. By a suitable choice of bases we can put these maps into the form

$$(24) \quad M_1 = \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix}, \quad M_2 = \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix}, \quad M_3 = \begin{pmatrix} B_1 \\ B_2 \\ J \end{pmatrix}, \quad M_4 = \begin{pmatrix} B_1' \\ B_2' \\ J' \end{pmatrix},$$

$$N_1 = \begin{pmatrix} 0 & 1 & 0 \end{pmatrix}, \quad N_2 = \begin{pmatrix} -1 & 0 & 0 \end{pmatrix}, \quad N_3 = \begin{pmatrix} -B_2 & B_1 & I \end{pmatrix}, \quad N_4 = \begin{pmatrix} -B_2' & B_1' & I' \end{pmatrix}.$$

Using the reality conditions  $\bar{J}^*(N)^\vee = M$  and  $\bar{J}^*(M)^\vee = -N$  we find that

$$(25) \quad B_1' = -B_2^\dagger, \quad B_2' = B_1^\dagger, \quad J' = I^\dagger, \quad I' = -J^\dagger.$$

Finally the condition  $NM = 0$  gives

- a)  $\mu_c = [B_1, B_2] + IJ = 0,$
- b)  $\mu_r = [B_1, B_1^\dagger] + [B_2, B_2^\dagger] + II^\dagger - J^\dagger J = -2\hbar \cdot 1_{k \times k}.$

These matrix equations are invariant under the following action of  $U(k)$  :

$$(26) \quad B_i \mapsto gB_i g^{-1}, \quad I \mapsto gI, \quad J \mapsto Jg^{-1}, \quad \text{where } g \in U(k).$$

Denote by  $\mathbf{M}$  the vector space of complex matrices  $(B_1, B_2, I, J)$ . It has a structure of a quaternionic vector space defined by

$$(B_1, B_2, I, J) \mapsto (-B_2^\dagger, B_1^\dagger, -J^\dagger, I^\dagger),$$

and, moreover, it is a flat hyperkähler manifold (see [27]). The map  $\mu = (\mu_r, \mu_c)$  is a hyperkähler moment map for the action of  $U(k)$  defined in (26) (see [19]). Since the action of  $U(k)$  on  $\mu_c^{-1}(0) \cap \mu_r^{-1}(-2\hbar \cdot 1)$  is free, the quotient  $\mathcal{M} = \mu_c^{-1}(0) \cap \mu_r^{-1}(-2\hbar \cdot 1)/U(k)$  is a smooth hyperkähler manifold. This manifold parametrizes complexes (22) with a real structure (23) and a trivialization on the line  $l$ .

On the other hand, it was proved in [27, 26] that the moduli space  $\mathfrak{M}(r, 0, k)$  of torsion free sheaves on the commutative  $\mathbb{P}^2$  with a trivialization on a fixed line can be identified with  $\mathcal{M}$ .

This identification can be described geometrically as follows. Let us assume that  $\hbar$  is positive. It can be checked that in this case the map  $N$  can fail to be surjective only on the plane  $z_4 = 0$ . We can restrict the complex (22) to the commutative plane  $z_3 = 0$ . The restriction is a monad and its cohomology sheaf is a torsion free sheaf. It is easy to see that this yields a complex isomorphism from  $\mathcal{M}$  to  $\mathfrak{M}(r, 0, p)$ .

The restriction of the complex (22) to a noncommutative plane is a monad as well. This yields a map from  $\mathcal{M}$  to the moduli space  $\mathfrak{M}_{\hbar}(r, 0, k)$  of bundles on the noncommutative plane. Let us show that this map is an isomorphism. To this end we note that on the level of the linear algebra data this map sends a quadruple  $(B_1, B_2, I, J)$  to the quadruple  $(b_1, b_2, i, j)$  with

$$b_1 = B_1 - B_2^\dagger, \quad b_2 = B_2 + B_1^\dagger, \quad i = I - J^\dagger, \quad j = J + I^\dagger.$$

Further, note that the equations  $\mu_c = 0$ ,  $\mu_r = -2\hbar \cdot 1$  are equivalent to the equation  $[b_1, b_2] + i \cdot j + 2\hbar \cdot 1 = 0$  and the vanishing of the moment map for the action of the group  $U(k)$  on the space of quadruples  $(b_1, b_2, i, j)$ . Now it follows from the theorem of Kempf and Ness ([27], [20]) that the map  $\mathcal{M} \rightarrow \mathfrak{M}_{\hbar}(r, 0, k)$  is a diffeomorphism. It becomes a complex isomorphism if we replace the natural complex structure of the space  $\mathcal{M}$  with another one within the  $\mathbb{P}^1$  of complex structures on  $\mathcal{M}$ .

Thus we have

**Theorem 7.1.** *The moduli space  $\mathfrak{M}_{\hbar}(r, 0, k)$  is a smooth hyperkähler manifold of real dimension  $4rk$ , and as a hyperkähler manifold it is isomorphic to the moduli space  $\mathfrak{M}(r, 0, k)$  of torsion free sheaves on the commutative  $\mathbb{P}^2$  with a trivialization on a fixed line. As a complex manifold  $\mathfrak{M}_{\hbar}(r, 0, k)$  is obtained from  $\mathfrak{M}(r, 0, k)$  by a rotation of the complex structure.*

The above discussion shows that there are natural bijections between

$A'$ . Bundles on  $\mathbb{P}_{\hbar}^2$  with a trivialization on the commutative line  $l$  and  $c_2 = k$ .

$B'$ . Solutions of the equations  $\mu_c = 0$ ,  $\mu_r = -2\hbar \cdot 1$  modulo the action of  $U(k)$ .

$C'$ . Complexes of sheaves on  $\mathbb{P}_{\hbar}^3$  of the form (22) with a trivialization on the commutative line  $l$  satisfying the reality condition (23).

One can show that for  $r > 1$  a generic complex (22) is a monad and its cohomology is a bundle  $\mathcal{E}$  on  $\mathbb{P}_{\hbar}^3$  such that

$$(27) \quad H^1(\mathbb{P}_{\hbar}^3, \mathcal{E}(-2)) = 0, \quad \bar{\mathcal{J}}^*(\mathcal{E})^\vee \cong \mathcal{E}.$$

Moreover, it can be shown that any bundle  $\mathcal{E}$  satisfying the conditions (27) can be represented as a cohomology of a monad of the form (22).

## 8. NONCOMMUTATIVE TWISTOR TRANSFORM

**8.1. Review of the twistor transform.** In the commutative case the ADHM construction of instantons has the following geometric interpretation. Consider the double fibration

$$(28) \quad \mathbf{G}(2; 4) \xleftarrow{p} \mathbf{F}\mathbf{I}(1, 2; 4) \xrightarrow{q} \mathbb{P}^3,$$

where  $\mathbf{G}(2;4)$  is the Grassmannian and  $\mathbf{Fl}(1,2;4)$  is the partial flag variety. The Grassmannian  $\mathbf{G}(2;4)$  has a real structure with  $\mathbb{S}^4$  as the set of real points. For any bundle  $\mathcal{E}$  on  $\mathbb{P}^3$  its twistor transform is defined as a sheaf  $p_*q^*\mathcal{E}$  on  $\mathbf{G}(2;4)$ . Given ADHM data we have a monad on  $\mathbb{P}^3$  whose cohomology is a bundle. It can be shown that the restriction of its twistor transform to the sphere  $\mathbb{S}^4$  coincides with the instanton bundle corresponding to these ADHM data. The instanton connection can also be reconstructed from the bundle on  $\mathbb{P}^3$  (see [4, 23] for details).

In this section we show that one can consider the noncommutative quadric introduced in section 3 as a noncommutative Grassmannian  $\mathbf{G}(2;4)$ . We also construct a noncommutative flag variety  $\mathbf{Fl}(1,2;4)$  and projections  $p, q$  giving a noncommutative analogue of the twistor diagram (28). The twistor transform can be defined in the same way as above. It produces a bundle on the noncommutative sphere from the deformed ADHM data. We show that this bundle is precisely the kernel of the map  $\mathcal{D}$  defined in section 2.

It should also be possible to construct the instanton connection on the noncommutative  $\mathbb{R}^4$  from the complex of sheaves on  $\mathbb{P}_h^3$ . To do this, one needs to develop differential geometry of noncommutative affine and projective varieties. We go some way in this direction by defining differential forms and spinors.

Since the goal of this section is mainly illustrative, we limit ourselves to stating the results. An interested reader should be able to fill in the proofs.

**8.2. Tensor categories.** A good way to construct noncommutative varieties with properties similar to those of commutative varieties is to start with a tensor category (see [24, 22]). Let  $\mathcal{T}$  be an abelian tensor category. Consider a tensor functor  $\Phi : \mathcal{T} \rightarrow \mathcal{Vect}$  to the abelian tensor category of vector spaces compatible with the associativity constraint but not compatible with the commutativity constraint. If  $A$  is a commutative algebra in the tensor category  $\mathcal{T}$ , then  $\Phi(A)$  is a noncommutative algebra in the tensor category  $\mathcal{Vect}$ . If  $M \in \mathcal{T}$  is a right  $A$ -module, then  $\Phi(M)$  is a right  $\Phi(A)$ -module. Any right  $A$ -module (in the category  $\mathcal{T}$ ) has a natural structure of a left  $A$ -module (and hence an  $A$ -bimodule). Thus any right  $\Phi(A)$ -module of the form  $\Phi(M)$  has a natural structure of a  $\Phi(A)$ -bimodule.

Consider the category  $Comm_{\mathcal{T}}$  of all finitely generated (graded) commutative algebras in the tensor category  $\mathcal{T}$ . Then under  $\Phi$  the category  $Comm_{\mathcal{T}}$  is mapped to a subcategory of the category of finitely generated (graded) algebras. This subcategory enjoys many properties of the category of commutative (graded) algebras. For example, for all  $A, B \in Comm_{\mathcal{T}}$  there is a natural algebra structure on  $\Phi(A) \otimes \Phi(B)$  coming from the algebra structure on  $A \otimes B$ . The corresponding subcategory in the category of noncommutative affine (resp.

projective) varieties shares a lot of properties with the category of commutative varieties. For example, if  $X$  and  $Y$  are varieties in this category, then using the tensor product of the corresponding algebras one can define the ‘‘Cartesian’’ product  $X \times Y$ . More generally, given a pair of morphisms  $X \rightarrow Z$  and  $Y \rightarrow Z$  one can define the fiber product  $X \times_Z Y$ . Further, starting from the module of differential forms of  $A$  one can construct the sheaf of differential forms on the corresponding noncommutative variety.

The category  $qgr(\Phi(A))$  has a nice subcategory which consists of modules of the form  $\Phi(M)$ , where  $M \in \mathcal{T}$  is an  $A$ -module. To any object  $\Phi(M)$  of this subcategory one can associate its symmetric and exterior powers. The symmetric powers of  $\Phi(M)$  form a noncommutative graded algebra. This enables one to define the projectivization of the sheaf corresponding to the module  $\Phi(M)$ .

**8.3. Yang-Baxter operators.** One way to construct an abelian tensor category  $\mathcal{T}$  with a functor  $\Phi : \mathcal{T} \rightarrow \mathcal{Vect}$  is to consider a Yang-Baxter operator (see [24], [22]).

A Yang-Baxter operator on a vector space  $V$  is an operator  $R : V \otimes V \rightarrow V \otimes V$ , such that

$$(29) \quad R^2 = \mathbf{id}_{V \otimes V}, \quad (R \otimes \mathbf{id}_V)(\mathbf{id}_V \otimes R)(R \otimes \mathbf{id}_V) = (\mathbf{id}_V \otimes R)(R \otimes \mathbf{id}_V)(\mathbf{id}_V \otimes R).$$

A Yang-Baxter operator induces an action of the permutation group  $\mathfrak{S}_n$  on the tensor power  $V^{\otimes n}$ , where the transposition  $(i, i+1) \in \mathfrak{S}_n$  acts as the operator

$$R_{i,i+1} = \mathbf{id}_{V^{\otimes(i-1)}} \otimes R \otimes \mathbf{id}_{V^{\otimes(n-i-1)}} : V^{\otimes n} \rightarrow V^{\otimes n}.$$

Equations (29) ensure that operators  $R_{i,i+1}$  satisfy the relations between the transpositions  $(i, i+1)$  in the group  $\mathfrak{S}_n$ .

If  $R$  is a Yang-Baxter operator on a vector space  $V$ , then the dual operator  $R^\vee : V^* \otimes V^* \rightarrow V^* \otimes V^*$  is also a Yang-Baxter operator.

Given a Yang-Baxter operator  $R : V \otimes V \rightarrow V \otimes V$ , one can construct an abelian tensor category  $\mathcal{T}_R$  and a functor  $\Phi_R : \mathcal{T}_R \rightarrow \mathcal{Vect}$  such that  $V$  is a  $\Phi_R$ -image of some object of  $\mathcal{T}_R$ , and the commutativity morphism in the category  $\mathcal{T}_R$  is mapped by  $\Phi_R$  to  $R$  [22]. As mentioned above, given any two objects  $A, B$  of the category  $\mathcal{Comm}_{\mathcal{T}_R}$ , one has a natural algebra structure on the vector space  $\Phi(A) \otimes \Phi(B)$ . This algebra will be denoted  $\Phi(A) \otimes_R \Phi(B)$  and called the  $R$ -tensor product of  $\Phi(A)$  and  $\Phi(B)$ .

It is well known that there is a one-to-one correspondence between irreducible representations of the group  $\mathfrak{S}_n$  and partitions of  $n$  (Young diagrams). Under this correspondence the trivial partition  $(n)$  corresponds to the sign representation, while the maximal partition  $\underbrace{(1, 1, \dots, 1)}_{n \text{ times}}$  corresponds to the identity representation. Given a partition  $(k_1, \dots, k_r)$

of  $n$  ( $k_1 \geq k_2 \geq \dots \geq k_r$ ) we denote by  $(k_1, \dots, k_r)$  the corresponding irreducible representation and by  $\Sigma_R^{(k_1, \dots, k_r)} V$  (resp.  $\Sigma_R^{(k_1, \dots, k_r)} V^*$ ) the  $(k_1, \dots, k_r)$ -isotypical component of  $V^{\otimes n}$  (resp.  $(V^*)^{\otimes n}$ ), i.e. the sum of all subrepresentations of  $V^{\otimes n}$  (resp.  $(V^*)^{\otimes n}$ ) isomorphic to  $(k_1, \dots, k_r)$ . We also put  $\Lambda_R^n V = \Sigma_R^{(n)} V$ ,  $\Lambda_R^n V^* = \Sigma_R^{(n)} V^*$  for brevity.

REMARK. The subspaces  $\Sigma_R^\lambda V \subset V^{\otimes n}$  are the  $\Phi_R$ -images of some objects of the category  $\mathcal{T}_R$ .

Let  $\lambda, \mu$  be partitions of  $n$  and  $m$  respectively. It is clear that the action of the permutation  $\sigma_{n,m} \in \mathfrak{S}_{n+m}$

$$\sigma_{n,m}(i) = \begin{cases} i + m, & \text{if } 1 \leq i \leq n \\ i - n, & \text{if } n + 1 \leq i \leq n + m \end{cases}$$

gives an isomorphism

$$R_{n,m} : \Sigma_R^\lambda V \otimes \Sigma_R^\mu V \rightarrow \Sigma_R^\mu V \otimes \Sigma_R^\lambda V.$$

REMARK. This isomorphism is the image of an isomorphism in the category  $\mathcal{T}_R$ .

The trivial example of a Yang-Baxter operator is the usual transposition

$$R_0(v_1 \otimes v_2) = v_2 \otimes v_1.$$

We will say that  $R$  is a deformation-trivial Yang-Baxter operator if  $R$  is an algebraic deformation of  $R_0$  in the class of Yang-Baxter operators. For a deformation-trivial Yang-Baxter operator  $R$  we have

$$\dim \Sigma_R^\lambda V = \dim \Sigma_{R_0}^\lambda V$$

for any partition  $\lambda$ .

**8.4. The noncommutative projective space.** Let  $R$  be a deformation-trivial Yang-Baxter operator on the vector space  $V^*$ . Then the graded algebra

$$S_R V^* = T(V^*) / \langle \Lambda_R^2 V^* \rangle$$

is a noncommutative deformation of the coordinate algebra of the projective space  $\mathbb{P}(V)$ . We denote by  $\mathbb{P}_R(V)$  the corresponding noncommutative variety. Thus  $\mathbb{P}_R(V)$  is a noncommutative deformation of the projective space  $\mathbb{P}(V)$ .

**Example 8.1.** *The operator*

$$(30) \quad \begin{aligned} R(z_i \otimes z_j) &= z_j \otimes z_i, & \text{if } (i, j) \neq (1, 2), (2, 1) \\ R(z_1 \otimes z_2) &= z_2 \otimes z_1 + 2\hbar(a z_3 \otimes z_4 + b z_4 \otimes z_3) \\ R(z_2 \otimes z_1) &= z_1 \otimes z_2 - 2\hbar(b z_3 \otimes z_4 + a z_4 \otimes z_3). \end{aligned}$$



is a deformation trivial Yang-Baxter operator on the 4 -dimensional vector space  $Z^*$  with the basis  $\{z_1, z_2, z_3, z_4\}$ . By definition the homogeneous coordinate algebra of  $\mathbb{P}_R(Z)$  is generated by  $z_1, z_2, z_3, z_4$  with relations (9) (we set  $a + b = 1$  as before). Hence  $\mathbb{P}_R(Z)$  is isomorphic to the noncommutative projective space  $\mathbb{P}_\hbar^3$  defined in section 3. The space  $Z^*$  was denoted  $U$  in that section.

The above example shows that part of the data encoded in the Yang-Baxter operator  $R$  is lost in the structure of the corresponding noncommutative projective space. We will see below that this data appears in the structure of other noncommutative varieties associated with  $R$ .

**8.5. Noncommutative Grassmannians.** It is well known that the homogeneous coordinate algebra of the Grassmann variety  $\mathbf{G}(k; V)$  is a graded quadratic algebra with  $\Lambda^k V^*$  as the space of generators and

$$\mathbf{Ker} \left( \Lambda^k V^* \otimes \Lambda^k V^* \rightarrow (V^*)^{\otimes 2k} \rightarrow \Sigma^{(k,k)} V^* \right)$$

as the space of relations. This description justifies the following definition.

**Definition 8.2.** Let  $R$  be a Yang-Baxter operator on the space  $V^*$ . The noncommutative Grassmann variety  $\mathbf{G}_R(k; V)$  is the noncommutative projective variety corresponding to the quadratic algebra

$$\mathcal{G}_R(k; V) = T(\Lambda_R^k V^*) / \left\langle \mathbf{Ker}(\Lambda_R^k V^* \otimes \Lambda_R^k V^* \rightarrow \Sigma_R^{(k,k)} V^*) \right\rangle$$

The algebra  $\mathcal{G}_R(k; V)$  is the  $\Phi_R$ -image of a commutative algebra in the category  $\mathcal{T}_R$ .

If  $R$  is deformation-trivial, then  $\mathbf{G}_R(k; V)$  is a noncommutative deformation of  $\mathbf{G}(k; V)$ . Note that  $\mathbf{G}_R(1; V) = \mathbb{P}_R(V)$  by definition.

**Example 8.3.** Consider the noncommutative Grassmannian  $\mathbf{G}_R(2; Z)$  corresponding to the Yang-Baxter operator (30). Let

$$z_{ij} = \frac{1}{2}((z_i \otimes z_j - z_j \otimes z_i) - R(z_i \otimes z_j - z_j \otimes z_i)) \in \Lambda_R^2 Z^*.$$

Then it is easy to check that  $\mathcal{G}_R(2; Z)$  is generated by the elements

$$Y_1 = z_{13}, \quad Y_2 = -z_{24}, \quad Y_3 = z_{23}, \quad Y_4 = z_{14}, \quad D = -z_{12}, \quad T = z_{34},$$

with relations

$$(31) \quad \begin{aligned} [Y_1, Y_2] &= 2\hbar a T^2, & [Y_3, Y_4] &= 2\hbar b T^2, \\ [D, Y_1] &= -2\hbar a Y_1 T, & [D, Y_2] &= 2\hbar a Y_2 T, \\ [D, Y_3] &= -2\hbar b Y_3 T, & [D, Y_4] &= 2\hbar b Y_4 T, \\ DT &= \frac{1}{2}(Y_1 Y_2 + Y_2 Y_1 + Y_3 Y_4 + Y_4 Y_3), \end{aligned}$$

$[Y_i, Y_j] = [T, Y_j] = [T, D] = 0$  for all  $i = 3, 4, j = 1, 2, 3, 4$ . Comparing with (7) one can see that the algebra  $\mathcal{G}_R(2; Z)$  is isomorphic to  $\mathbb{Q}_\hbar$  with  $G$  and  $\theta$  given by

$$G = \frac{1}{2} \begin{pmatrix} 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{pmatrix}, \quad \theta = 2\hbar \begin{pmatrix} 0 & a & 0 & 0 \\ -a & 0 & 0 & 0 \\ 0 & 0 & 0 & b \\ 0 & 0 & -b & 0 \end{pmatrix}.$$

Note that the variables  $X_i, i = 1, 2, 3, 4$ , used in section 7 to describe the quadric are related to  $Y_i, i = 1, 2, 3, 4$ , by the following formulas:

$$(32) \quad \begin{aligned} Y_1 &= X_2 + \sqrt{-1} X_1, & Y_2 &= -X_2 + \sqrt{-1} X_1, \\ Y_3 &= X_4 + \sqrt{-1} X_3, & Y_4 &= -X_4 + \sqrt{-1} X_3. \end{aligned}$$

**8.6. Products of Grassmannians and flag varieties.** Let  $R$  be a Yang-Baxter operator on the vector space  $V^*$ . Consider a sequence  $k_1, \dots, k_r$  of integers. Let  $\mathbb{Z}^r$  be the free abelian group with  $r$  generators  $e_1, \dots, e_r$ . The  $R$ -tensor product

$$\mathcal{G}_R(k_1; V) \otimes_R \dots \otimes_R \mathcal{G}_R(k_r; V)$$

is a  $\mathbb{Z}^r$ -graded algebra generated by the vector spaces  $\Lambda_R^{k_i} V^*$  in degree  $e_i$ , with relations

$$\mathbf{Ker} \left( \Lambda_R^{k_i} V^* \otimes \Lambda_R^{k_i} V^* \rightarrow \Sigma_R^{(k_i, k_i)} V^* \right)$$

in degree  $2e_i$  for all  $i$  and

$$\mathbf{Ker} \left( (\Lambda_R^{k_i} V^* \otimes \Lambda_R^{k_j} V^*) \oplus (\Lambda_R^{k_j} V^* \otimes \Lambda_R^{k_i} V^*) \xrightarrow{(\mathbf{id}, -R_{k_j, k_i})} \Lambda_R^{k_i} V^* \otimes \Lambda_R^{k_j} V^* \right)$$

in degree  $e_i + e_j$  for all  $i > j$ . For any increasing sequence  $k_1, \dots, k_r$  we define also a  $\mathbb{Z}^r$ -graded algebra  $\mathcal{FL}_R(k_1, \dots, k_r; V)$ . It has the same generators as the algebra  $\mathcal{G}_R(k_1; V) \otimes_R \dots \otimes_R \mathcal{G}_R(k_r; V)$ , subject to the same relations in degrees  $2e_i$  and to relations

$$\mathbf{Ker} \left( (\Lambda_R^{k_i} V^* \otimes \Lambda_R^{k_j} V^*) \oplus (\Lambda_R^{k_j} V^* \otimes \Lambda_R^{k_i} V^*) \xrightarrow{(\mathbf{id}, -R_{k_j, k_i})} \Lambda_R^{k_i} V^* \otimes \Lambda_R^{k_j} V^* \longrightarrow \Sigma_R^{(k_i, k_j)} V^* \right)$$

in degree  $e_i + e_j$  for all  $i > j$ . This definition is suggested by the Borel-Weil-Bott theorem (see [14]). In particular, for  $R = R_0$  we get the algebra corresponding to the commutative flag variety.

We define the  $R$ -Cartesian product  $\mathbf{G}_R(k_1; V) \times_R \dots \times_R \mathbf{G}_R(k_r; V)$  and the noncommutative flag variety  $\mathbf{Fl}_R(k_1, \dots, k_r; V)$  as the noncommutative varieties corresponding to the algebras  $\mathcal{G}_R(k_1; V) \otimes_R \dots \otimes_R \mathcal{G}_R(k_r; V)$  and  $\mathcal{FL}_R(k_1, \dots, k_r; V)$  respectively.

To make this compatible with our definition of a noncommutative variety, we consider instead of a  $\mathbb{Z}^r$ -graded algebra its diagonal subalgebra. The diagonal subalgebra is a graded

algebra whose  $n$ -th graded component is the  $n(e_1 + \dots + e_r)$ -graded component of the  $\mathbb{Z}^r$ -graded algebra. Thus according to section 3 the category of coherent sheaves on the  $R$ -Cartesian product of Grassmannians (or the flag variety) is the category  $qgr$  of the corresponding diagonal subalgebra.

The algebra  $\mathcal{F}\mathcal{L}_R(k_1, \dots, k_r; V)$  is the  $\Phi_R$ -image of a commutative algebra in the category  $\mathcal{T}_R$ . Hence one can define the  $R$ -Cartesian product of several flag varieties.

If  $R$  is deformation-trivial, then  $\mathbf{G}_R(k_1; V) \times_R \dots \times_R \mathbf{G}_R(k_r; V)$  and  $\mathbf{F}\mathbf{L}_R(k_1, \dots, k_r; V)$  are noncommutative deformations of the corresponding commutative varieties.

Note that we have a canonical embedding of the graded algebra  $\mathcal{G}_R(k_i; V)$  into the graded algebra  $\mathcal{F}\mathcal{L}_R(k_1, \dots, k_i, \dots, k_r; V)$  inducing the canonical projections

$$p_i : \mathbf{F}\mathbf{L}_R(k_1, \dots, k_i, \dots, k_r; V) \rightarrow \mathbf{G}_R(k_i; V).$$

On the other hand, by definition  $\mathcal{F}\mathcal{L}_R(k_1, \dots, k_r; V)$  is a quotient algebra of the algebra  $\mathcal{G}_R(k_1; V) \otimes_R \dots \otimes_R \mathcal{G}_R(k_r; V)$ . Hence  $\mathbf{F}\mathbf{L}_R(k_1, \dots, k_r; V)$  can be regarded as a closed subvariety in  $\mathbf{G}_R(k_1; V) \times_R \dots \times_R \mathbf{G}_R(k_r; V)$ .

**Example 8.4.** *The algebra  $\mathcal{G}_R(1; Z) \otimes_R \mathcal{G}_R(2; Z)$  corresponding to the Yang-Baxter operator (30) is generated by the elements  $z_1, z_2, z_3, z_4, Y_1, Y_2, Y_3, Y_4, D, T$  with relations (9), (31), and*

$$\begin{aligned} [z_1, Y_2] &= -2\hbar a z_4 T, & [z_2, Y_1] &= 2\hbar a z_3 T, \\ [z_1, Y_3] &= -2\hbar b z_3 T, & [z_2, Y_4] &= -2\hbar b z_4 T, \\ [z_1, D] &= -2\hbar b z_3 Y_4 - 2\hbar a z_4 Y_1, & [z_2, D] &= 2\hbar a z_3 Y_2 - 2\hbar b z_4 Y_3, \end{aligned}$$

$[z_1, Y_1] = [z_2, Y_2] = 0$ ,  $[z_3, Y_i] = [z_3, D] = 0$ ,  $[z_4, Y_i] = [z_4, D] = 0$ ,  $[z_i, T] = 0$  for all  $i = 1, 2, 3, 4$ . The algebra  $\mathcal{F}\mathcal{L}_R(1, 2; Z)$  is given by the same generators subject to the same relations, as well as the additional relations

$$(33) \quad \begin{pmatrix} 0 & T & Y_2 & Y_3 \\ T & 0 & -Y_4 & Y_1 \\ Y_2 & Y_4 & 0 & D - \hbar(a+b)T \\ Y_3 & -Y_1 & -D - \hbar(a+b)T & 0 \end{pmatrix} \begin{pmatrix} z_1 \\ z_2 \\ z_3 \\ z_4 \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \\ 0 \\ 0 \end{pmatrix}.$$

As explained above, we have projections

$$\mathbb{Q}_{\hbar} \longleftarrow \mathbf{G}_R(2; Z) \xleftarrow{p} \mathbf{F}\mathbf{L}_R(1, 2; Z) \xrightarrow{q} \mathbb{P}_R(Z) \longleftarrow \mathbb{P}_{\hbar}^3$$

and a closed embedding

$$\mathbf{F}\mathbf{L}_R(1, 2; Z) \subset \mathbf{G}_R(2; Z) \times_R \mathbb{P}_R(Z) = \mathbb{Q}_{\hbar} \times_R \mathbb{P}_{\hbar}^3.$$

**8.7. Tautological bundles.** Let  $\mathcal{V}$  (resp.  $\mathcal{V}^*$ ,  $\Sigma_R^\lambda \mathcal{V}$ ,  $\Sigma_R^\lambda \mathcal{V}^*$ ) denote the coherent sheaf on  $\mathbf{G}_R(k; V)$  corresponding to the free right  $\mathcal{G}_R(k; V)$ -module  $V \otimes \mathcal{G}_R(k; V)$  (resp.  $V^* \otimes \mathcal{G}_R(k; V)$ ,  $\Sigma_R^\lambda V \otimes \mathcal{G}_R(k; V)$ ,  $\Sigma_R^\lambda V^* \otimes \mathcal{G}_R(k; V)$ ). Since the space of global sections of the sheaf  $\mathcal{O}(1)$  on the Grassmannian  $\mathbf{G}_R(k; V)$  is  $\Lambda_R^k V^*$ , the maps  $\Lambda_R^{k-1} V^* \rightarrow V \otimes \Lambda_R^k V^*$  and  $\Lambda_R^{k+1} V^* \rightarrow V^* \otimes \Lambda_R^k V^*$  induce morphisms of sheaves

$$\Lambda_R^{k-1} \mathcal{V}^*(-1) \xrightarrow{\phi} \mathcal{V} \quad \text{and} \quad \Lambda_R^{k+1} \mathcal{V}^*(-1) \xrightarrow{\psi} \mathcal{V}^*.$$

We put  $S = \mathbf{Im} \phi$ ,  $\mathcal{V}/S = \mathbf{Coker} \phi$ ,  $S' = \mathbf{Im} \psi$ ,  $\mathcal{V}^*/S' = \mathbf{Coker} \psi$ .

REMARK. For  $k = 1$  we have  $S = \mathcal{O}(-1)$ ,  $\mathcal{V}^*/S' = \mathcal{O}(1)$ .

One can show that these sheaves are locally free. We refer to them as tautological bundles.

The free  $\mathcal{G}_R(k; V)$ -modules, corresponding to the sheaves  $\Sigma_R^\lambda \mathcal{V}$ ,  $\Sigma_R^\lambda \mathcal{V}^*$  are the  $\Phi_R$ -images of free modules over the corresponding algebra in the category  $\mathcal{T}_R$ . Furthermore, the morphisms  $\phi$  and  $\psi$  are  $\Phi_R$ -images. This implies that the  $\mathcal{G}_R(k; V)$ -modules corresponding to the tautological bundles are  $\Phi_R$ -images as well. Therefore they all have a natural structure of  $\mathcal{G}_R(k; V)$ -bimodules. This allows to define  $R$ -symmetric powers  $S_R^k(-)$  (resp.  $R$ -exterior powers  $\Lambda_R^k(-)$ ) of the tautological bundles as the corresponding  $\Phi_R$ -images.

One can check that we have canonical isomorphisms of bimodules

$$\mathcal{V}^*/S' \cong S^\vee, \quad S' \cong (\mathcal{V}/S)^\vee.$$

**Example 8.5.** Let  $R$  be the Yang-Baxter operator (30) and  $k = 2$ . Let  $\check{z}_1, \check{z}_2, \check{z}_3, \check{z}_4$  be the dual basis of  $Z$ . Then the twisted maps

$$\begin{aligned} \phi(1) &: Z^* \otimes \mathcal{O}_{\mathbf{G}_R} \rightarrow Z \otimes \mathcal{O}_{\mathbf{G}_R}(1), \\ \psi(1) &: Z \otimes \mathcal{O}_{\mathbf{G}_R} \cong \Lambda_R^3 Z^* \otimes \mathcal{O}_{\mathbf{G}_R} \rightarrow Z^* \otimes \mathcal{O}_{\mathbf{G}_R}(1) \end{aligned}$$

are given by

$$\begin{aligned} \phi(1) : \begin{pmatrix} z_1 \\ z_2 \\ z_3 \\ z_4 \end{pmatrix} &\mapsto \begin{pmatrix} 0 & D + \hbar(a-b)T & -Y_1 & -Y_4 \\ D - \hbar(a-b)T & 0 & -Y_3 & Y_2 \\ -Y_1 & Y_3 & 0 & -T \\ -Y_4 & -Y_2 & T & 0 \end{pmatrix} \begin{pmatrix} \check{z}_1 \\ \check{z}_2 \\ \check{z}_3 \\ \check{z}_4 \end{pmatrix} \\ \psi(1) : \begin{pmatrix} \check{z}_1 \\ \check{z}_2 \\ \check{z}_3 \\ \check{z}_4 \end{pmatrix} &\mapsto \begin{pmatrix} 0 & T & Y_2 & Y_3 \\ T & 0 & -Y_4 & Y_1 \\ Y_2 & Y_4 & 0 & D - \hbar(a+b)T \\ Y_3 & -Y_1 & -D - \hbar(a+b)T & 0 \end{pmatrix} \begin{pmatrix} z_1 \\ z_2 \\ z_3 \\ z_4 \end{pmatrix} \end{aligned}$$

Note that  $\psi(1)\phi = 0$  and  $\phi(1)\psi = 0$ . Hence we have isomorphisms

$$S'(1) \cong \mathcal{V}/S, \quad S(1) \cong S^\vee.$$

Note also that on the open subset  $T \neq 0$  elements  $(z_3, z_4)$  give a trivialization of the tautological bundle  $S^\vee$ . More precisely, the restriction of the sections  $z_1, z_2$  of  $S^\vee$  can be expressed as

$$(34) \quad z_1 = y_4 z_3 - y_1 z_4, \quad z_2 = -y_2 z_3 - y_3 z_4$$

where  $y_i = T^{-1}Y_i$ . Similarly, the elements  $(\check{z}_1, \check{z}_2)$  give a trivialization of  $\mathcal{V}/S$  on  $T \neq 0$ . Thus the restrictions of all tautological bundles to the open subset  $T \neq 0$  correspond to the free rank two bimodule over the Weyl algebra  $A(\mathbb{C}_\hbar^4)$ .

**8.8. Pull-back and push-forward.** Recall that we have canonical projections  $p_i : \mathbf{Fl}_R(k_1, k_2; V) \rightarrow \mathbf{G}_R(k_i; V)$  ( $i = 1, 2$ ). Given a right graded  $\mathcal{G}_R(k_i; V)$ -module  $E$  we consider the right bigraded  $\mathcal{FL}_R(k_1, k_2; V)$ -module  $E \otimes_{\mathcal{G}_R(k_i; V)} \mathcal{FL}_R(k_1, k_2; V)$ . The diagonal subspace of this module is a graded module over the diagonal subalgebra of  $\mathcal{FL}_R(k_1, k_2; V)$ . This gives the pull-back functor

$$p_i^* : \text{coh}(\mathbf{G}_R(k_i; V)) \rightarrow \text{coh}(\mathbf{Fl}_R(k_1, k_2; V)).$$

The pull-back functor is exact and takes a  $\Phi_R$ -image to a  $\Phi_R$ -image. In particular, the pull-backs of the tautological bundles have a canonical bimodule structure.

The pull-back functor  $p_i^*$  admits a right adjoint functor  $p_{i*} : \text{coh}(\mathbf{Fl}_R(k_1, k_2; V)) \rightarrow \text{coh}(\mathbf{G}_R(k_i; V))$ , called the push-forward functor. It also takes a  $\Phi_R$ -image to a  $\Phi_R$ -image.

The line bundles  $p_1^*\mathcal{O}(i)$  and  $p_2^*\mathcal{O}(j)$  on the flag variety  $\mathbf{Fl}_R(k_1, k_2; V)$  are  $\Phi_R$ -images, hence they have a canonical bimodule structure. Therefore, we have a well-defined tensor

product

$$\mathcal{O}(i, j) = p_1^* \mathcal{O}(i) \otimes p_2^* \mathcal{O}(j).$$

The line bundle  $\mathcal{O}(i, j)$  is also a  $\Phi_R$ -image and has a canonical bimodule structure. The  $n$ -th graded component of the corresponding module over the diagonal subalgebra of  $\mathcal{FL}_R(k_1, k_2; V)$  is the  $((n+i)e_1 + (n+j)e_2)$ -graded component of the algebra  $\mathcal{FL}_R(k_1, k_2; V)$ .

One can check that the push-forward of the line bundle  $\mathcal{O}(j_1, j_2)$  with respect to  $p_2$  is given by the formula

$$p_{2*} \mathcal{O}(j_1, j_2) = S_R^{j_1}(S^\vee)(j_2).$$

**8.9.  $\mathbf{Fl}_R(1, 2; Z)$  as the projectivization of the tautological bundle.** The  $R$ -symmetric powers of the tautological bundle form a sheaf of graded algebras on the Grassmannian  $\mathbf{G}_R(k; V)$

$$S_R(S^\vee) = T(S^\vee) / \langle \Lambda_R^2 S^\vee \rangle.$$

The corresponding  $\mathcal{G}_R(k; V)$ -module

$$\bigoplus_{i,j=0}^{\infty} \Gamma(\mathbf{G}_R(k; V), S_R^j(S^\vee)(i))$$

is a bigraded module with a structure of a bigraded algebra. One can check that this bigraded algebra is isomorphic to the bigraded algebra  $\mathcal{FL}_R(1, k; V)$ . Thus we can regard the flag variety  $\mathbf{Fl}_R(1, k; V)$  as the projectivization of the tautological bundle  $S$  on the Grassmannian  $\mathbf{G}_R(k; V)$ . In particular,  $\mathbf{Fl}_R(1, 2; Z)$  is the projectivization of the tautological bundle  $S$  on the Grassmannian  $\mathbf{G}_R(2; Z)$ .

**8.10. Noncommutative twistor transform.** If  $E$  is a coherent sheaf on the noncommutative projective space  $\mathbb{P}_R(Z) = \mathbb{P}_\hbar^3$ , we define its twistor transform as the sheaf  $p_* q^* E$  on  $\mathbf{G}_R(2; Z) = \mathbb{Q}_\hbar$ , where  $q$  is the projection  $\mathbf{Fl}_R(1, 2; Z) \rightarrow \mathbb{P}_R(Z) = \mathbb{P}_\hbar^3$  and  $p$  is the projection  $\mathbf{Fl}_R(1, 2; Z) \rightarrow \mathbf{G}_R(2; Z) = \mathbb{Q}_\hbar$ . Similarly, we can define the twistor transform of a complex of sheaves on  $\mathbb{P}_\hbar^3$ . Actually, it is more natural to consider the derived twistor transform, i.e. the derived functor of the ordinary twistor transform.

Consider a complex  $\mathcal{C}$  of the form

$$0 \longrightarrow H \otimes \mathcal{O}(-1) \xrightarrow{M} K \otimes \mathcal{O} \xrightarrow{N} L \otimes \mathcal{O}(1) \longrightarrow 0$$

on the projective space  $\mathbb{P}_\hbar^3$ . One can check that under the twistor transform one has

$$\mathcal{O}_{\mathbb{P}_\hbar^3}(-1) \mapsto 0, \quad \mathcal{O}_{\mathbb{P}_\hbar^3} \mapsto \mathcal{O}_{\mathbf{G}_R}, \quad \mathcal{O}_{\mathbb{P}_\hbar^3}(1) \mapsto S^\vee.$$

In fact, for these sheaves the derived twistor transform coincides with the ordinary one. Thus the (derived) twistor transform takes the complex  $\mathcal{C}$  to the complex

$$0 \longrightarrow K \otimes \mathcal{O} \xrightarrow{\mathcal{N}} L \otimes S^\vee \longrightarrow 0.$$

Let  $\mathcal{E}$  denote the middle cohomology of the complex  $\mathcal{C}$ . It follows that the twistor transform takes  $\mathcal{E}$  to the kernel of the map  $\mathcal{N} : K \otimes \mathcal{O} \longrightarrow L \otimes S^\vee$ .

One can describe  $\mathcal{N}$  without reference to the twistor transform. The morphism  $\mathcal{N}$  is the same thing as a vector space morphism

$$(35) \quad N_1 z_1 + N_2 z_2 + N_3 z_3 + N_4 z_4 : K \longrightarrow Z^* \otimes L.$$

Here the maps  $N_i$  are given in terms of the deformed ADHM data according to (24) and (25). The map  $\mathcal{N}$  is a composition of two maps

$$K \otimes \mathcal{O}_{\mathbf{G}_R} \longrightarrow L \otimes Z^* \otimes \mathcal{O}_{\mathbf{G}_R} \longrightarrow L \otimes S^\vee,$$

where the first map is given by (35), while the second map comes from the canonical projection  $Z^* \otimes \mathcal{O}_{\mathbf{G}_R} \rightarrow S^\vee$ . (We remind that  $S^\vee$  is the cokernel of the map  $\psi : Z \otimes \mathcal{O}_{\mathbf{G}_R}(-1) \longrightarrow Z^* \otimes \mathcal{O}_{\mathbf{G}_R}$ .)

Recall that on the open subset  $\{T \neq 0\}$  the bundle  $S^\vee$  is trivial, and the elements  $(z_3, z_4)$  give its trivialization (see (34)). Hence the restriction of the twistor transform of the complex (22) to this open subset is isomorphic to the complex

$$(36) \quad 0 \longrightarrow K \otimes \mathcal{O} \xrightarrow{\begin{pmatrix} N_3 + y_4 N_1 - y_2 N_2 \\ N_4 - y_1 N_1 - y_3 N_2 \end{pmatrix}} (L \oplus L) \otimes \mathcal{O} \longrightarrow 0.$$

Assume now that the complex (22) is given by the deformed ADHM data  $(B_1, B_2, I, J)$  (see section (7)). Applying the formulas (24) and (25), we see that with respect to the chosen bases of  $L$  and  $K$  the map  $\mathcal{N}$  is given by the matrix

$$\begin{pmatrix} -B_2 + y_2 & B_1 + y_4 & I \\ -B_1^\dagger + y_3 & -B_2^\dagger - y_1 & -J^\dagger \end{pmatrix}.$$

It is evident that this operator is related to the operator  $\mathcal{D}$  in (4) by a change of basis. In particular, the Nekrasov-Schwarz coordinates  $\xi_1, \xi_2, \bar{\xi}_1, \bar{\xi}_2$  (see section 2) can be expressed through  $x_i = T^{-1} X_i$  as follows:

$$\begin{aligned} \xi_1 &= -y_4 = x_4 - \sqrt{-1} x_3, & \xi_2 &= y_2 = -x_2 + \sqrt{-1} x_1, \\ \bar{\xi}_1 &= y_3 = x_4 + \sqrt{-1} x_3, & \bar{\xi}_2 &= -y_1 = -x_2 - \sqrt{-1} x_1. \end{aligned}$$

Thus the twistor transform of the complex corresponding to the deformed ADHM data coincides with the instanton bundle corresponding to these data (see section 2). This gives

a geometric interpretation of the deformed ADHM construction of the noncommutative instanton bundle.

**8.11. Differential forms.** Let an algebra  $A$  be the  $\Phi_R$ -image of a commutative algebra in the category  $\mathcal{T}_R$ . This means that there exists an operator  $R : A^{\otimes 2} \rightarrow A^{\otimes 2}$  compatible with the multiplication law of  $A$ . Above we have defined the  $R$ -tensor product  $A \otimes_R A$  which is also an algebra with a Yang-Baxter operator. Explicitly, the multiplication law of  $A \otimes_R A$  is defined as follows. Let  $m$  be the multiplication map from  $A \otimes A$  to  $A$ . Then the multiplication map from  $(A \otimes A) \otimes (A \otimes A)$  to  $A \otimes A$  is given by  $m_{12}m_{34}R_{23}$  in the obvious notation. It is easy to see that the multiplication map  $m$  is a homomorphism of algebras.

Let  $I$  denote the kernel of the map  $m : A \otimes_R A \rightarrow A$ . Then  $I$  is a two-sided ideal of the algebra  $A \otimes_R A$ .

**Definition 8.6.** We define the bimodule of  $R$ -differential forms of the algebra  $A$  by

$$\Omega_A^1 = I/I^2.$$

For a motivation of this definition, see [12]. Furthermore, suppose  $A$  is a graded algebra. Consider the total grading of the bigraded algebra  $A \otimes_R A$ . The two-sided ideal  $I$  inherits the grading. Therefore the bimodule  $\Omega_A^1$  is graded too.

In the graded case, besides  $\Omega_A^1$ , we can define the module of projective differential forms of  $A$  in the following way. Let  $\chi : A \otimes_R A \rightarrow A \otimes_R A$  be the linear operator which acts on the  $(p, q)$ -th graded component of the algebra  $A \otimes_R A$  as a scalar multiplication by  $q$ . Since  $\chi$  is a derivation, we have  $\chi(I^2) \subset I$ . Therefore  $m(\chi(I^2)) = 0$ . Furthermore, the induced map  $\Omega_A^1 = I/I^2 \xrightarrow{m \cdot \chi} A$  is a morphism of graded  $A$ -bimodules.

**Definition 8.7.** We define the  $A$ -bimodule of projective differential forms of the algebra  $A$  by

$$\widehat{\Omega}_A^1 = \mathbf{Ker}(\Omega_A^1 \xrightarrow{m \cdot \chi} A).$$

First, let us apply this construction of differential forms to the noncommutative affine variety  $\mathbb{C}_\hbar^4$  (subsection 3.4). The algebra  $A(\mathbb{C}_\hbar^4)$  of polynomial functions on  $\mathbb{C}_\hbar^4$  is the Weyl algebra:

$$A(\mathbb{C}_\hbar^4) = \mathbb{T}(x_1, x_2, x_3, x_4) / \langle [x_i, x_j] = \hbar \theta_{ij} \rangle_{1 \leq i, j \leq 4}.$$

Let us define the Yang-Baxter operator on the tensor square of the subspace of  $A(\mathbb{C}_\hbar^4)$  spanned by  $1, x_1, x_2, x_3, x_4$  by the formula

$$1 \otimes x_i \mapsto x_i \otimes 1, \quad x_i \otimes 1 \mapsto 1 \otimes x_i, \quad x_i \otimes x_j \mapsto x_j \otimes x_i + \hbar \theta_{ij} \cdot 1 \otimes 1 \quad \text{for all } 1 \leq i, j \leq 4.$$



This Yang-Baxter operator has a unique extension to the whole  $A(\mathbb{C}_\hbar^4)$  compatible with the multiplication law.

There is another way to look at this Yang-Baxter operator. Recall that  $\mathbb{C}_\hbar^4$  is an open subset  $T \neq 0$  in the noncommutative Grassmannian  $\mathbf{G}_R(2; Z)$  where  $R$  is defined by (30). The Yang-Baxter operator on the quadratic algebra  $\mathcal{G}_R(2; Z)$  has the property that  $R(T \otimes a) = a \otimes T$  for any  $a \in \mathcal{G}_R(2; Z)$ . Hence it descends to a Yang-Baxter operator on  $A(\mathbb{C}_\hbar^4)$ . It is easy to see that it acts on the tensor square of the subspace spanned by  $1, x_1, x_2, x_3, x_4$  in the above manner.

We define the sheaf of differential forms  $\Omega_{\mathbb{C}_\hbar^4}^1$  as the bimodule of  $R$ -differential forms of the algebra  $A(\mathbb{C}_\hbar^4)$ . It is easy to check that  $\Omega_{\mathbb{C}_\hbar^4}^1$  is isomorphic to the bimodule  $A(\mathbb{C}_\hbar^4)^{\oplus 4}$ . Furthermore, we can take any  $R$ -exterior power of  $\Omega_{\mathbb{C}_\hbar^4}^1$  and thereby define  $\Omega_{\mathbb{C}_\hbar^4}^p$ . This enables us to define a connection and its curvature on any bundle on the noncommutative affine space. The relevant formulas were written above (see subsection 1.5).

Second, we define the sheaf of differential forms  $\Omega_{\mathbf{G}_R}^1$  on the noncommutative Grassmannian  $\mathbf{G}_R(k; V)$  as the sheaf corresponding to the module of projective differential forms  $\widehat{\Omega}_{\mathcal{G}_R}^1$ .

It can be shown that as in the commutative case we have an isomorphism of coherent sheaves on the noncommutative Grassmannian  $\mathbf{G}_R(k; V)$ :

$$\Omega_{\mathbf{G}_R}^1 \cong S \otimes S'.$$

It follows that for  $k = 1$  that we have an exact sequence

$$0 \longrightarrow \Omega_{\mathbb{P}_R(V)}^1 \longrightarrow \mathcal{V}^*(-1) \longrightarrow \mathcal{O} \longrightarrow 0.$$

Thus this definition of the sheaf of differential forms  $\Omega_{\mathbb{P}_R(V)}^1$  is consistent with Definition 4.8.

Similarly, one can define the sheaf of differential forms  $\Omega_{\mathbf{F}\mathbf{1}_R}^1$  on the noncommutative flag variety  $\mathbf{F}\mathbf{1}_R(k_1, \dots, k_r; V)$ . One can check that the projection

$$p_i : \mathbf{F}\mathbf{1}_R(k_1, \dots, k_i, \dots, k_r; V) \rightarrow \mathbf{G}_R(k_i; V)$$

induces a morphism of bundles  $p_i^* : \Omega_{\mathbf{G}_R}^1 \rightarrow \Omega_{\mathbf{F}\mathbf{1}_R}^1$ .

In the commutative case the ADHM construction of the instanton connection can be interpreted in terms of twistor transform (see [4, 23] for details). We believe that this can be done in the noncommutative case as well. It appears that the most convenient definition of connection on a bundle on a noncommutative projective variety is in terms of jet bundles (see, for example, [23]).

9. INSTANTONS ON A  $q$ -DEFORMED  $\mathbb{R}^4$ 

In this paper we have focused on a particular noncommutative deformation of  $\mathbb{R}^4$  related to the Wigner-Moyal product (3). This is the only deformation of  $\mathbb{R}^4$  which is known to arise in string theory. But most of our constructions work for more general deformations which do not have a clear physical interpretation. For example, let us replace  $\mathbb{C}_\hbar^4$  with a noncommutative affine variety whose coordinate ring is generated by  $z_1, z_2, z_3, z_4$  subject to the following quadratic relations:

$$qz_1z_2 - q^{-1}z_2z_1 = \hbar, \quad qz_3z_4 - q^{-1}z_4z_3 = \hbar, \quad [z_1, z_3] = [z_1, z_4] = [z_2, z_3] = [z_2, z_4] = 0.$$

We will denote this noncommutative affine variety by  $\mathbb{C}_{q,\hbar}^4$ , and its coordinate algebra by  $\mathcal{A}_{q,\hbar}$ . If  $\hbar$  and  $q$  are real, we can define a  $*$ -operation on  $\mathcal{A}_{q,\hbar}$  by  $z_1^* = z_2, z_3^* = z_4$ . The corresponding real noncommutative affine variety will be denoted by  $\mathbb{R}_{q,\hbar}^4$ .

Consider now the following deformation of the ADHM equations:

$$(37) \quad [B_1, B_2]_{q^{-1}} + IJ = 0, \quad [B_1, B_1^\dagger]_{q^{-1}} + [B_2, B_2^\dagger]_q + II^\dagger - J^\dagger J = -2\hbar \cdot 1_{k \times k}.$$

Here  $B_1, B_2 \in \text{Hom}(V, V)$ ,  $I \in \text{Hom}(W, V)$ ,  $J \in \text{Hom}(V, W)$ , as usual, and by  $[A, B]_q$  we mean a  $q$ -commutator:

$$[A, B]_q = qAB - q^{-1}BA.$$

We claim that solutions of these “ $q$ -deformed” ADHM equations can be used as an input for the construction of instantons on  $\mathbb{R}_{q,\hbar}^4$  of rank  $r = \dim W$  and instanton charge  $k = \dim V$ . Let us sketch this construction. Define an operator

$$\mathcal{D} \in \text{Hom}_{\mathcal{A}_{q,\hbar}}((V \oplus V \oplus W) \otimes_{\mathbb{C}} \mathcal{A}_{q,\hbar}, (V \oplus V) \otimes_{\mathbb{C}} \mathcal{A}_{q,\hbar})$$

by the formula

$$\mathcal{D} = \begin{pmatrix} B_1 - qz_1 & -qB_2 + qz_2 & I \\ B_2^\dagger - \bar{z}_2 & qB_1^\dagger - \bar{z}_1 & J^\dagger \end{pmatrix}.$$

Now we can go through the same manipulations as in section 2: assume that  $\mathcal{D}$  is surjective, and its kernel is a free module, and define a connection 1-form by the expression (5). The same formal computation as in section 2 shows that the curvature of this connection is anti-self-dual.

In order to ensure that  $\mathcal{D}$  is surjective, it is probably necessary to replace the algebra  $\mathcal{A}_{q,\hbar}$  with some bigger algebra containing  $\mathcal{A}_{q,\hbar}$  as a subalgebra. This bigger algebra should play the role of the algebra of smooth functions on our noncommutative  $\mathbb{R}^4$ . For  $\hbar = 0$ ,  $q \neq 1$  there is even a natural candidate for this bigger algebra: it should consist of  $C^\infty$

functions on  $\mathbb{C}^2$  with some suitable growth conditions at infinity and the product defined by

$$(38) \quad (f \star g)(z_1, z_2, \bar{z}_1, \bar{z}_2) = \exp \left( -\ln(q) \left( z_1 \bar{z}'_1 \frac{\partial^2}{\partial z_1 \partial \bar{z}'_1} + z_2 \bar{z}'_2 \frac{\partial^2}{\partial z_2 \partial \bar{z}'_2} - z'_1 \bar{z}_1 \frac{\partial^2}{\partial z'_1 \partial \bar{z}_1} - z'_2 \bar{z}_2 \frac{\partial^2}{\partial z'_2 \partial \bar{z}_2} \right) \right) f(z_1, z_2, \bar{z}_1, \bar{z}_2) g(z'_1, z'_2, \bar{z}'_1, \bar{z}'_2) \Big|_{z'_1=z_1, z'_2=z_2}.$$

Assuming that this formal expression exists, it is easy to check that the product is associative, that polynomial functions form a subalgebra with respect to it, and that this subalgebra is isomorphic to  $\mathcal{A}_{q, \hbar}$ .

It is natural to conjecture that all instantons on  $\mathbb{R}_{q, \hbar}^4$  arise from this deformed ADHM construction. Note that in this case the deformed ADHM equations are not hyperkähler moment map equations, and one cannot use the hyperkähler quotient construction to infer the existence of hyperkähler metric on the quotient space.

The algebro-geometric part of the story can also be generalized. We did not go through this carefully, but nevertheless would like to indicate one result. It appears that the  $q$ -deformed ADHM data can be interpreted in terms of sheaves on a more general noncommutative  $\mathbb{P}^2$  than the one defined in section 3. The graded algebra corresponding to this noncommutative  $\mathbb{P}^2$  is generated by degree one elements  $z_1, z_2, z_3$  with the quadratic relations

$$qz_1z_2 - q^{-1}z_2z_1 = 2\hbar z_3^2, \quad [z_i, z_3] = 0, \quad i = 1, 2.$$

This algebra is one of the Artin-Schelter regular algebras of dimension three [1, 2]. It is characterized by the fact that the corresponding noncommutative variety  $\mathbb{P}_{q, \hbar}^2$  contains as subvarieties a commutative quadric and a noncommutative line. The latter is given by the equation  $z_3 = 0$ . In the limit  $q \rightarrow 1$  the plane  $\mathbb{P}_{q, \hbar}^2$  reduces to  $\mathbb{P}_{\hbar}^2$ , and the union of the quadric and the line turns into the triple commutative line  $l$  which played such a prominent role in this paper. If  $q \neq 1$ , then in the limit  $\hbar \rightarrow 0$  the quadric turns into a union of two intersecting commutative lines  $z_1 = 0$  and  $z_2 = 0$ .

For any  $q$  the line  $z_3 = 0$  should be regarded as “the line at infinity” (which is noncommutative for  $q \neq 1$ ). It is plausible that the  $q$ -deformed ADHM data are in one-to-one correspondence with bundles, or maybe torsion-free sheaves, on  $\mathbb{P}_{q, \hbar}^2$  with a trivialization on this line.

## 10. APPENDIX

In this section we define a  $\star$ -product on the space of complex-valued  $C^\infty$  functions on  $\mathbb{R}^n$  whose derivatives of arbitrary order are polynomially bounded. The  $\star$ -product endows

this space with a structure of a  $\mathbb{C}$ -algebra and reduces to the Wigner-Moyal product (3) on polynomial functions.

**Definition 10.1.** Let  $\Phi$  be a topological vector space which is a subspace of the space of  $C^\infty$  functions on  $\mathbb{R}^n$ , and let  $\Phi'$  be the space of distributions on  $\Phi$ . Let  $f$  be a  $\mathbb{C}$ -valued function on  $\mathbb{R}^n$  which simultaneously is a distribution in  $\Phi'$ .  $f$  is called a multiplier if for any  $\phi \in \Phi$   $f\phi \in \Phi$ .

The set of multipliers of  $\Phi'$  is obviously a subspace of  $\Phi'$ .

**Definition 10.2.** Let  $f \in \Phi'$ .  $f$  is called a convolute if for any  $\phi \in \Phi$  we have

$$(f * \phi)(x) \equiv (f(\xi), \phi(x + \xi)) \in \Phi,$$

and this expression depends continuously on  $\phi$ . The above expression is called the convolution of  $f$  with  $\phi$ .

The set of convolutes is obviously a subspace of  $\Phi'$ .

We will denote the Fourier duals of  $\Phi$  and  $\Phi'$  by  $\tilde{\Phi}$  and  $\tilde{\Phi}'$ , respectively. If  $f \in \Phi$ , then  $\tilde{f} \in \tilde{\Phi}$  will be the Fourier transform of  $f$ , etc.

**Definition 10.3.** The Schwartz space  $\mathcal{S}(\mathbb{R}^n)$  is the space of  $\mathbb{C}$ -valued  $C^\infty$  functions on  $\mathbb{R}^n$  such that  $\phi \in \mathcal{S}$  if and only if all the norms

$$(39) \quad \sup_x x^k D^m \phi(x), \quad k = 0, 1, 2, \dots,$$

are finite. Here  $m = (m_1, \dots, m_n)$  is an arbitrary polyindex.

Convergence on  $\mathcal{S}$  is defined using the family of norms (39). Then  $\mathcal{S}$  becomes a complete countably normed space [17].

**Proposition 10.4.** A function  $f \in \mathcal{S}'$  is a multiplier if and only if it is a  $C^\infty$  function on  $\mathbb{R}^n$  all of whose derivatives are polynomially bounded.

*Proof.* Obvious. □

The following theorem proved in [36] describes the subspace of convolutes of  $\mathcal{S}'$ :

**Theorem 10.5.** A distribution  $f \in \mathcal{S}'$  is a convolute if and only if it has the form

$$f = \sum_{|\alpha| < r} D^\alpha f_\alpha(x),$$

where  $r$  is a positive integer, and  $f_\alpha$  are  $C^0$  functions on  $\mathbb{R}^n$  which decrease at infinity faster than any negative power of  $x$ .

The functions which decrease at infinity faster than any negative power will be called rapidly decreasing.

The following theorem is proved in [17], vol. 2, ch. III:

**Theorem 10.6.** *Fourier transform and its inverse act as automorphisms on both  $\mathcal{S}$  and  $\mathcal{S}'$ .*

From now on we identify  $\mathcal{S} \cong \tilde{\mathcal{S}}$ ,  $\mathcal{S}' \cong \tilde{\mathcal{S}'}$ .

**Theorem 10.7.** *Fourier transform and its inverse establish an isomorphism between the space of multipliers and the space of convolutes of  $\mathcal{S}'$ .*

*Proof.* By the preceding theorem, it is sufficient to show that the Fourier transform of every multiplier is a convolute, and vice versa. The former fact is proved in [17], vol. 2, ch. III. Let us prove the converse.

By theorem 10.5, every convolute has the form

$$f(x) = \sum_{|\alpha| < r} D^\alpha f_\alpha(x)$$

for some  $r$  and rapidly decreasing continuous functions  $f_\alpha$ . Let

$$\tilde{f}_\alpha(p) = \int f_\alpha(x) e^{\sqrt{-1}px} d^n x.$$

be the Fourier transform of  $f_\alpha(x)$ . Since the integrals

$$\int x^\beta f_\alpha(x) e^{\sqrt{-1}px} d^n x$$

are absolutely convergent, the functions  $\tilde{f}_\alpha$  are  $C^\infty$  functions. Furthermore, the Fourier transform of  $f$  is equal to

$$\tilde{f}(p) = \sum_{|\alpha| < r} (-\sqrt{-1} p)^\alpha \tilde{f}_\alpha(p)$$

(see [17], vol. 2, ch. III), hence  $\tilde{f}$  is also a  $C^\infty$  function. Finally, since by the preceding theorem the Fourier transform of any element of  $\mathcal{S}'$  is again an element of  $\mathcal{S}'$ ,  $\tilde{f}$  and all its derivatives are polynomially bounded. Hence  $\tilde{f}$  is a multiplier.  $\square$

**Definition 10.8.** *Let  $\omega$  be a skew-symmetric real-valued bilinear form on  $\mathbb{R}^n$ . The  $\diamond$ -product on the space of convolutes of  $\mathcal{S}'$  is defined by*

$$(\tilde{f} \diamond \tilde{g})(p) = \int \tilde{f}(q) \tilde{g}(p - q) e^{\sqrt{-1}\omega(p,q)} \frac{d^n q}{(2\pi)^n}.$$

**Theorem 10.9.** *The  $\diamond$ -product is well-defined and makes the space of convolutes of  $\mathcal{S}'$  into an algebra over  $\mathbb{C}$ .*

*Proof.* We will prove that the  $\diamond$ -product of two convolutes of  $\mathcal{S}'$  is well-defined, and is again a convolute of  $\mathcal{S}'$ . The rest is obvious.

It is sufficient to consider the case when

$$\tilde{f}(p) = D^\alpha \tilde{f}_0(p), \quad \tilde{g}(p) = D^\beta \tilde{g}_0(p).$$

Then, integrating by parts, we may rewrite the  $\diamond$ -product in the following form:

$$(-1)^{|\alpha|} \int \tilde{f}_0(q) \frac{\partial^\alpha}{\partial q^\alpha} \left[ \frac{\partial^\beta}{\partial p^\beta} \tilde{g}_0(p-q) e^{\sqrt{-1}\omega(p,q)} \right] \frac{d^n q}{(2\pi)^n}.$$

Derivatives acting on the exponential bring down powers of  $p$ , so the integral can be rewritten as

$$P\left(p, \frac{\partial}{\partial p}\right) \int \tilde{f}_0(q) \frac{\partial^\beta}{\partial p^\beta} \tilde{g}_0(p-q) e^{\sqrt{-1}\omega(p,q)} \frac{d^n q}{(2\pi)^n},$$

where  $P(u, v)$  is a homogeneous polynomial of degree  $|\alpha|$ . We now use the Leibniz rule repeatedly to rewrite the expression above as

$$P\left(p, \frac{\partial}{\partial p}\right) \int Q\left(q, \frac{\partial}{\partial p}\right) \left[ \tilde{f}_0(q) \tilde{g}_0(p-q) e^{\sqrt{-1}\omega(p,q)} \right] \frac{d^n q}{(2\pi)^n},$$

where  $Q(u, v)$  is a homogeneous polynomial of degree  $|\beta|$ . Because both  $\tilde{f}_0$  and  $\tilde{g}_0$  are rapidly decreasing, the integral converges absolutely and defines a  $C^0$  function of  $q$  which is rapidly decreasing. Hence the  $\diamond$ -product of  $\tilde{f}_0$  and  $\tilde{g}_0$  has the form

$$\sum_{|m| \leq |\alpha| + |\beta|} D^m \tilde{h}_m(p),$$

where the functions  $\tilde{h}_m(p)$  are continuous and rapidly decreasing. It follows that the space of convolutes is closed under the  $\diamond$ -product.  $\square$

**Corollary 10.10.** *The space of multipliers of  $\mathcal{S}'$  inherits a product from the  $\diamond$ -product on the space of convolutes of  $\mathcal{S}'$ , and this product makes the space of multipliers into an algebra over  $\mathbb{C}$ . Polynomials form a subalgebra of this algebra isomorphic to the Weyl algebra with generators  $x_i$ ,  $i = 1, \dots, n$ , and relations*

$$[x_i, x_j] = 2\sqrt{-1}\omega_{ij}.$$

*Proof.* The first statement is an immediate consequence of theorems 10.7 and 10.9. The second statement follows from a simple computation.  $\square$

It is this product on the space of multipliers that we call the Wigner-Moyal product and denote with  $\star$ .

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