

Journal Club by Oleg Chalaev

Nonequilibrium mesoscopic conductance fluctuations

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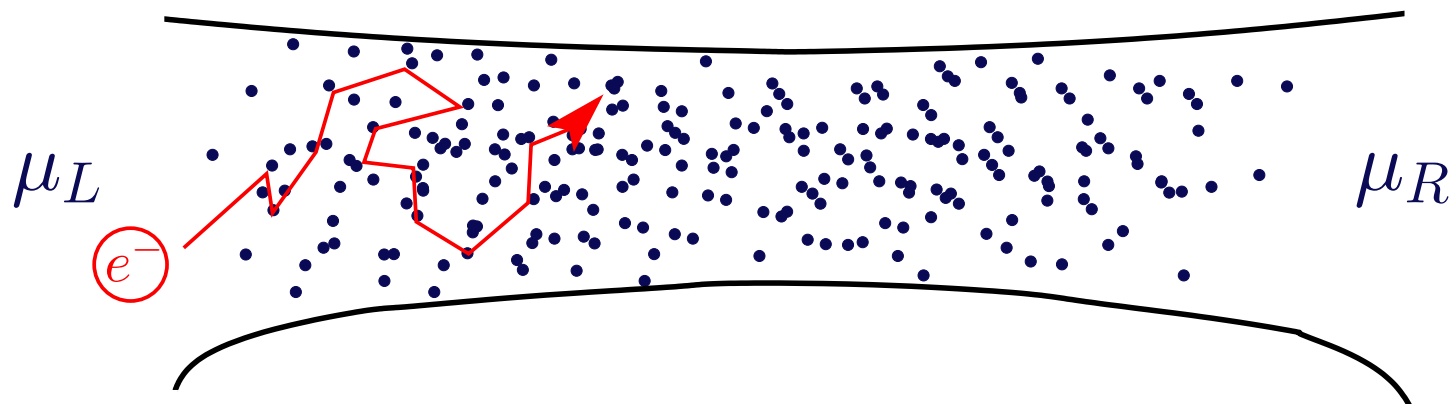
Universal conductance fluctuations in equilibrium

Altshuler'85; Lee& Stone'85:

The conductance of a metallic sample with a fixed concentration of impurities exhibits fluctuations of the order of e^2/h , when $L_\phi \gtrsim L$.

$$\langle \delta\sigma^2 \rangle = \frac{8 e^4}{15 h^2}$$

What happens out of equilibrium?



A few words about G_K

- Keldysh technique is very similar to the $T = 0$ equilibrium technique described in Abrikosov, Gor'kov, Dzyaloshinskii
- The major difference from $T = 0$ technique: now instead of scalar Green's function we operate with $G = \begin{pmatrix} G_R & G_K \\ 0 & G_A \end{pmatrix}$
- Without interaction $G_K = (1 - 2f_E)(G_R - G_A)$,
 $f_E =$ energy distribution function.

Connection with the density of states ν :

$$\frac{1}{V} \sum_p [G_R^E(\vec{p}) - G_A^E(\vec{p})] = -2\pi i \nu_E$$

The average current $\vec{j} \sim \vec{p} \langle G_K(\vec{p}) \rangle$

The two-step energy distribution function

Without interaction: kinetic equation or charge conservation

$$\implies \nabla^2 \langle G^K \rangle = 0$$

Boundary conditions:

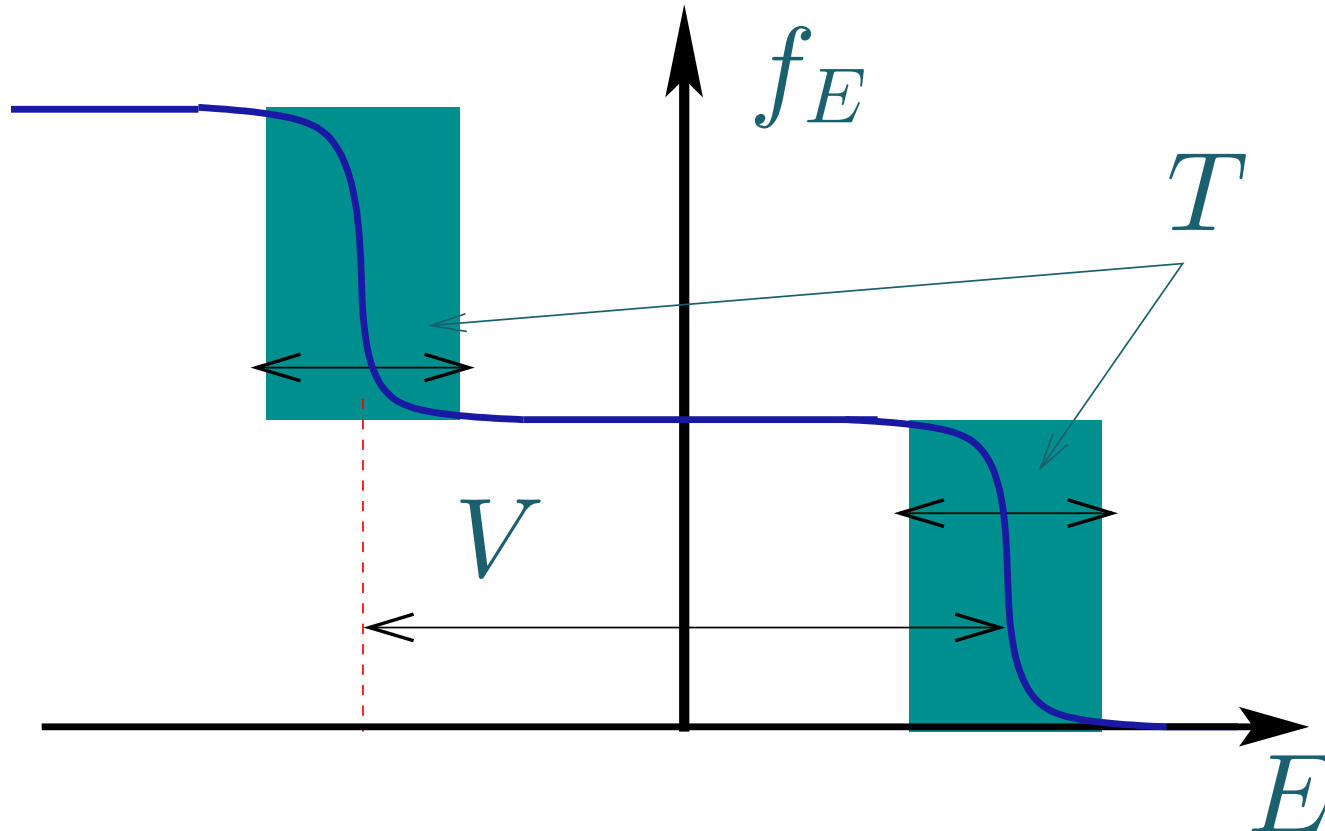
$$\langle G_K^E(x) \rangle = -2\pi i\nu \times \begin{cases} 1 - 2f(E), & x = 0 \\ 1 - 2f(E - eV), & x = L \end{cases}$$

\implies we get double-step distribution function:

$$\langle G_K^E(x) \rangle = -2\pi i\nu \left\{ 1 - 2f(E) + 2\frac{x}{L} [f(E) - f(E - eV)] \right\}$$

This derivation of the double-step f_E I don't understand; My understanding is more primitive, see pp. 65-66 in [1].

The two-step energy distribution function



⇒ we get double-step distribution function:

$$\langle G_{\text{K}}^E(x) \rangle = -2\pi i\nu \left\{ 1 - 2f(E) + 2\frac{x}{L} [f(E) - f(E - eV)] \right\}$$

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What are the consequences of non-equilibrium

- $G_{R/A}$ are not altered since they are short range objects:

$$G_{R/A}(\vec{r}, \vec{r}') \sim \exp\left[-\frac{|\vec{r}-\vec{r}'|}{l}\right], \quad l \ll L.$$
- G_K is changed since it depends on the energy distribution
- Cooperon/diffuson are changed since they are long-range objects, $\sim L_\omega = \sqrt{iD/\omega}$

An equation for the **diffusion propagator**:

$$\left\{ \partial_x^2 + \frac{i\omega}{D} + \frac{ie}{D} [\phi_1(x) - \phi_2(x)] \right\} \Pi_\omega(x, x') = -\delta(x - x')$$

Boundary conditions: $\Pi = 0$ at $x = 0$ and $x = L$

The effect of interaction between electrons

Ideologically the authors have done (in reality – more rigorously) the following:

- got expressions with cooperons from nasty diagrams
- inserted $\tau_\phi \equiv \tau_\phi(T^*(x))$ into the denominator of the cooperon;

$$T^*(x) = eVx(1-x), \quad \tau_\phi[T^*] \sim (D\nu^2/T^{*2})^{1/3}$$

- claiming that in this way they've taken interactions into account

Is this really all what interaction does out of equilibrium????

The result

Out of equilibrium, additional contributions appear:

$$\langle \delta g \delta g \rangle = \langle \delta g \delta g \rangle_0 + \langle \delta g \delta g \rangle_1 + \langle \delta g \delta g \rangle_2 ,$$

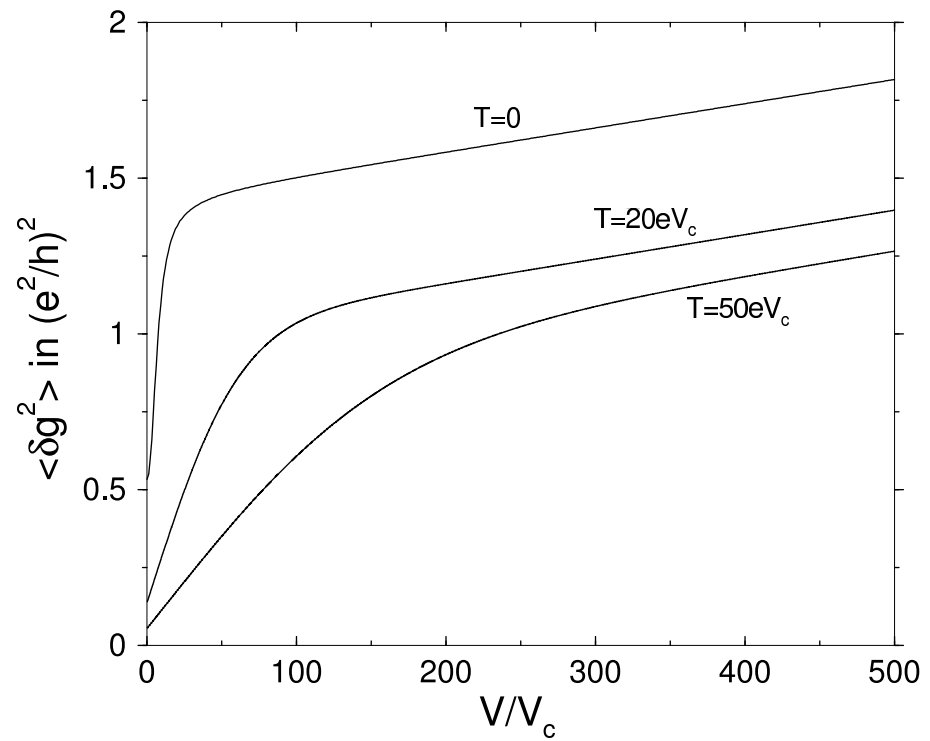
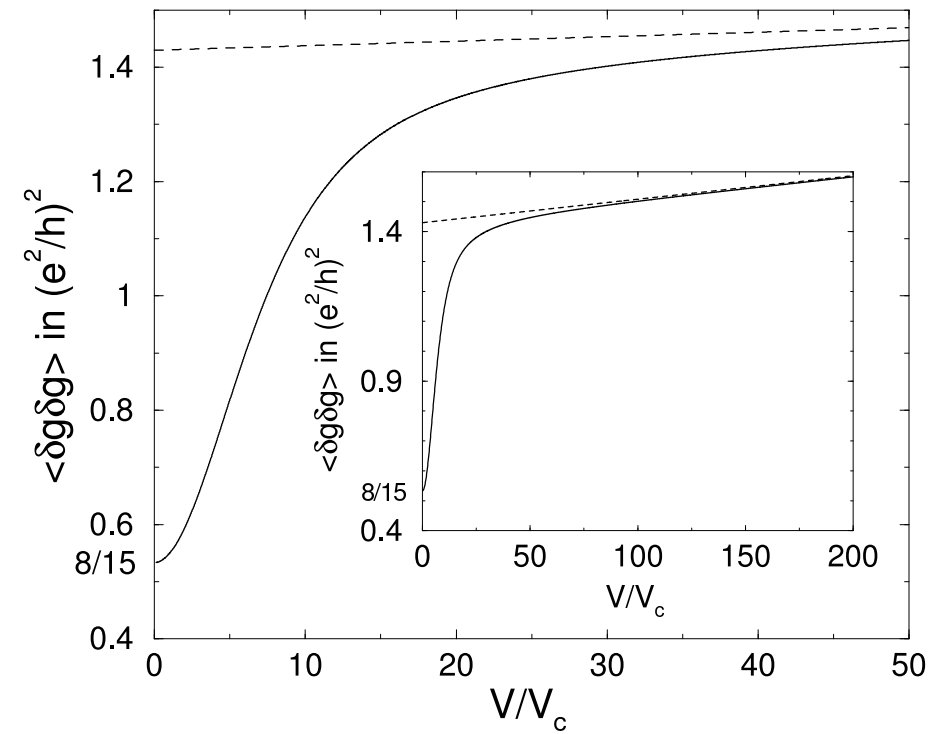
$$\langle \delta g(V) \delta g(V) \rangle_0 = 16 \Xi_0 \Big|_{\alpha=0} = \frac{8}{15} ,$$

$$\langle \delta g(V) \delta g(V) \rangle_1 = 32 \int_0^{V/V_c} dz \frac{\partial}{\partial \alpha} \Xi_{z - \frac{V}{V_c}} \Big|_{\alpha=0} , \quad eV_c = D/L^2$$

$$\langle \delta g(V) \delta g(V) \rangle_2 = -16 \int_0^{V/V_c} dz_1 dz_2 \frac{\partial^2}{\partial \alpha^2} \Xi_{z_1 - z_2} \Big|_{\alpha=0} ,$$

$$\Xi_z = \int_0^1 dy_1 dy_2 \left[2 |\Pi_z(y_1, y_2)|^2 + \text{Re} \Pi_z^2(y_1, y_2) \right]$$

The voltage and temperature dependence



Conclusions

- “nonequilibrium” terms are calculated.
 - interaction is taken into account via dephasing
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References

- [1] Supriyo Datta. *Electronic Transport in Mesoscopic Systems*. Cambridge uni. press, 1997.