

Nonlinear thermal radiation in flow induced by a slendering surface accounting thermophoresis and Brownian diffusion

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Abstract. Our attention here in this research is on scrutinizing the nonlinear convection characteristics in a flow induced by a slendering surface. Flow expression is developed through electrically conducting Williamson nanomaterial. Nonlinear forms of stretching and free stream velocities are imposed. Consideration of nonlinear thermal radiation, non-uniform heat generation/absorption, Joule and convective heating aspects describe the phenomenon of heat transfer. The zero-mass condition for concentration is also considered. The compatible transformations produce strong nonlinear differential systems. The problems are computed analytically utilizing the *bvp4c* procedure. Heat transfer rate and drag force are also explained for various physical variables. Our analysis reveals that the heat transfer rate augments via larger radiation parameter and Biot number. Moreover, larger Brownian motion and thermophoresis parameters have opposite characteristics on concentration field. For the verification of the present findings, the results of the presented analysis have been compared with the available works in particular situations and reasonable agreement is noted.

1 Introduction

The flow investigation in the region of boundary layer via a stretchable sheet is a fascinating research field because of its significance in several utilizations in engineering systems and industries, like wire drawing, glass fiber production, hot rolling, metal spinning, contraction process of metallic plates and aerodynamic extrusion of plastic sheets. In view of such applications, Sakiadis [1] initially modeled the problem of boundary layer approximations by stretchable surfaces. Afterwards, several attempts regarding stretchable flow have been reported (see [2–6]). Besides this, when the temperature difference between the surface and ambient fluid is significantly large, the nonlinear density temperature (NDT) variations influence the flow and heat transfer distributions. Therefore in the buoyancy force expression, nonlinear density temperature variations cannot be neglected. Information related to such phenomenon can be seen in refs. [7–12].

Undoubtedly, energy generation is the most important issue in industrial requirements. Thus scientists are focused to improve heat transport rate of systems which comprise chemical plants, power stations, air conditioning and petrochemical industry. In this regard, several techniques were proposed but are not appropriate because of less thermal conductivity. Consequently thermal scientists established energy materials known as nanomaterials comprising nanoscale particles [13]. The tiny-scale particles in a base fluid are termed as nanofluid. A nanofluid is an improved type of fluid comprising nanometer-sized particles (diameter < 100 nm) or fibers suspended in the ordinary fluid. Such feature of high thermal conductivity has been noticed by Masuda *et al.* [14]. Choi *et al.* [15] disclosed that the thermal conductivity of the fluid can be improved up to twice that of the base fluid by the addition of a small quantity (< 1% by volume) of nanoparticles to common heat transport liquids. Nanoliquids have a wide range of applications in microelectronics, microfluidics, solid-state lighting, transportation, medical applications, biomedical industry, detergency,

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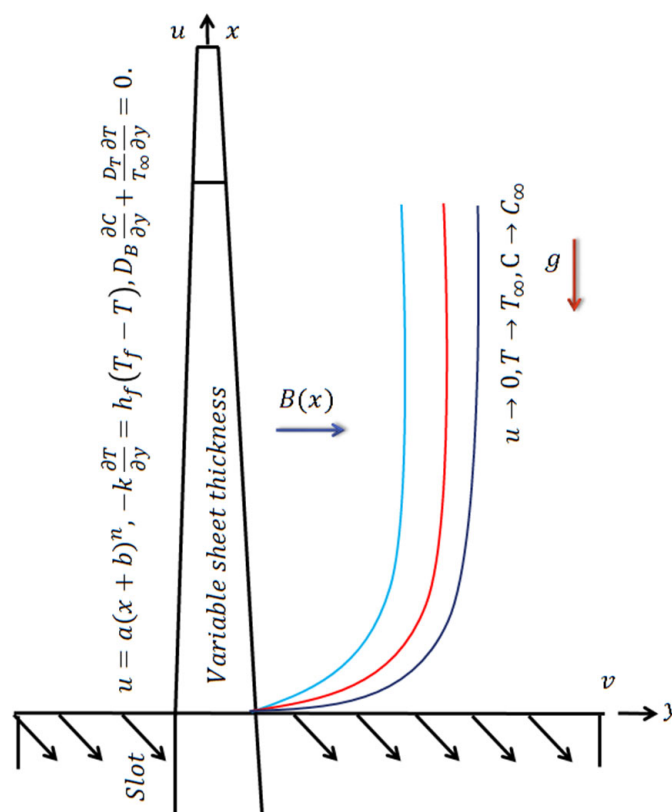


Fig. 1. Geometry of flow problem.

power generation, etc. Besides this, the interaction of magnetic field with nanoliquids has numerous potential applications and may be utilized to deal with problems like chilling of nuclear reactors by liquid sodium and induction flow meter, which relies on the potential difference of the fluid in the direction normal to the motion and to the magnetic field. A number of articles on magneto nanofluids are available in the literature and some are given in refs. [16–25].

The consideration of the nonlinear thermal radiative fluid flow by a stretchable surface is an area of potential interest for researchers due to its demands in several engineering and physical processes. Applicable areas of such processes are solar-power technology, propulsion devices, nuclear plants, combustion chambers, aircraft and chemical processes at high operating temperatures. Especially the thermal radiation aspect plays an important role in regulating the heat transport process in the polymer processing industry. Keeping this in mind, the effect of nonlinear thermal radiation on the Sakiadis flow were firstly reported by Pantokratoras and Fang [26]. Cortell [27] addressed the importance of thermal radiation effect in the stretchable flow of a viscous liquid. Khan *et al.* [28] developed the series solutions to analyze the nonlinear thermal radiation effect on the Burgers nanofluid flow with convective heating and heat generation/absorption. Hayat *et al.* [29] investigated the simultaneous characteristics of nonlinear convection and magnetohydrodynamics (MHD) in the Walter-B nanofluid flow induced by a nonlinear stretchable surface in the presence of nonlinear thermal radiation and convective heating. Nonlinear convection and nonlinear thermal radiation impacts on the Maxwell nanofluid by convectively heated stretched sheet are reported by Mahanthesh *et al.* [30].

The intention here is to investigate the feature of convective heating and zero-mass flux conditions in nonlinear convective flow of a magnetic Williamson liquid. The flow is generated due to a slendering surface. The phenomenon of nanoparticles occurs due to consideration of thermophoretic and Brownian motion. Nonlinear thermal radiation, heat generation/absorption and Joule heating are incorporated in the energy transport expression. The governed mathematical expressions are calculated through a numerical algorithm [31–34]. The significance of arising non-dimensional variables is discussed through graphs and tables.

2 Mathematical formulation

Let us scrutinize the two-dimensional (2D) nonlinear convective flow of a magneto Williamson nanomaterial towards a slendering surface at $y = A_1(x + b)^{\frac{1-n}{2}}$. The considered nanoliquid model includes the salient characteristics of thermophoresis and Brownian motion. The nanofluid is conducted electrically through the non-uniform magnetic field ($B(x) = B_0(x + b)^{\frac{n-1}{2}}$) operated perpendicularly to the flow trend (see fig. 1). The supposition of a small magnetic Reynolds number leads to negligible characteristics of electric field. Besides this, the nonlinear thermal radiation, heat

generation/absorption, Joule and convective heating effects are also considered in the heat transfer process. In view of these assumptions the governing expressions for the present situation are [29,35]:

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0, \tag{1}$$

$$u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} = \nu \frac{\partial^2 u}{\partial y^2} + \sqrt{2\nu}\Gamma \left(\frac{\partial u}{\partial y} \right) \frac{\partial^2 u}{\partial y^2} - \frac{\sigma B^2(x)}{\rho_f} u + g \{ \Lambda_1(T - T_\infty) + \Lambda_2(T - T_\infty)^2 \} + g \{ \Lambda_3(C - C_\infty) + \Lambda_4(C - C_\infty)^2 \}, \tag{2}$$

$$u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y} = \alpha^* \frac{\partial^2 T}{\partial y^2} + \frac{\sigma B^2(x)}{(\rho c)_f} u^2 + \tau \left(D_B \frac{\partial C}{\partial y} \frac{\partial T}{\partial y} + \frac{D_T}{T_\infty} \left(\frac{\partial T}{\partial y} \right)^2 \right) + \frac{1}{(\rho c)_f} \frac{16\sigma^*}{3k^*} \frac{\partial}{\partial y} \left(T^3 \frac{\partial T}{\partial y} \right) + \frac{Q(x)}{(\rho c)_f} (T - T_\infty), \tag{3}$$

$$u \frac{\partial C}{\partial x} + v \frac{\partial C}{\partial y} = D_B \frac{\partial^2 C}{\partial y^2} + \frac{D_T}{T_\infty} \frac{\partial^2 T}{\partial y^2}, \tag{4}$$

for conditions [29]

$$u = u_w(x) = a(x + b)^n, \quad v = 0, \quad -k \frac{\partial T}{\partial y} = h_f(T_f - T),$$

$$D_B \frac{\partial C}{\partial y} + \frac{D_T}{T_\infty} \frac{\partial T}{\partial y} = 0 \quad \text{at} \quad y = A_1(x + b)^{\frac{1-n}{2}}, \tag{5}$$

$$u \longrightarrow 0, \quad T \longrightarrow T_\infty, \quad C \longrightarrow C_\infty \quad \text{as} \quad y \longrightarrow \infty, \tag{6}$$

in which (u, v) are the velocity elements along the (x, y) axes, respectively, $\nu = (\frac{\mu}{\rho_f})$ the kinematic viscosity with μ the fluid kinematic viscosity and ρ_f the base fluid density, $\alpha^* = \frac{k}{(\rho c)_f}$ the fluid thermal conductivity, Γ the material parameter of Williamson fluid, B_0 the magnetic field strength, (T, C) the fluid (temperature, concentration), T_f the convective fluid temperature, (T_∞, C_∞) the ambient (temperature, concentration), k the fluid thermal conductivity, $\tau = \frac{(\rho c)_p}{(\rho c)_f}$ the heat capacity ratio, $(\rho c)_f$ the heat capacity of base fluid, $(\rho c)_p$ the effective heat capacity of nanoparticles, (D_T, D_B) the (thermophoretic, Brownian) diffusion coefficients, σ^* the Stefan-Boltzmann constant, k^* the mean absorption coefficient, $Q(x) = Q_0(x + b)^{n-1}$ the non-uniform heat generation/absorption, (Λ_1, Λ_2) the (linear, non-linear) thermal expansion coefficients, (Λ_3, Λ_4) the (linear, nonlinear) concentration expansion coefficients, $u_w(x)$ the stretching velocity, n the power index and $h_f = h(x + b)^{\frac{n-1}{2}}$ the non-uniform heat transfer coefficient and (a, b) the dimensional constants.

Letting [29]

$$\psi = \sqrt{\left(\frac{2}{n+1} \right) a\nu(x + b)^{n+1} F(\xi)}, \quad \xi = y \sqrt{\left(\frac{n+1}{2} \right) \left(\frac{a}{\nu} \right) (x + b)^{n-1}},$$

$$u = a(x + b)^n F'(\xi), \quad v = -\sqrt{\left(\frac{n+1}{2} \right) a\nu(x + b)^{n-1}} \left(F(\xi) + \xi \frac{n-1}{n+1} F'(\xi) \right),$$

$$\Theta(\xi) = \frac{T - T_\infty}{T_f - T_\infty}, \quad \Phi(\xi) = \frac{C - C_\infty}{C_\infty}, \tag{7}$$

eq. (1) is trivially satisfied and other equations are reduced to

$$F''' - \frac{2n}{n+1} F'^2 + FF'' + We \sqrt{\frac{n+1}{2}} F'' F''' - \frac{2}{n+1} M^2 F' + \frac{2}{n+1} \lambda [(1 + \beta_t \Theta)\Theta + N(1 + \beta_c \Phi)\Phi] = 0, \tag{8}$$

$$[1 + R\{1 + (\Theta_f - 1)\Theta\}^3 \Theta']' + Pr \left(F\Theta' + Nb\Theta'\Phi' + Nt\Theta'^2 + \frac{2}{n+1} \delta\Theta + EcM^2 F'^2 \right) = 0, \tag{9}$$

$$\Phi'' + ScF\Phi' + \frac{Nt}{Nb} \Theta'' = 0, \tag{10}$$

$$F(\alpha) = \alpha \left(\frac{1-n}{1+n} \right), \quad F'(\alpha) = 1, \quad \Theta'(\alpha) = -\gamma(1 - \Theta(\alpha)),$$

$$Nb\Phi'(\alpha) + Nt\Theta'(\alpha) = 0 \quad \text{at} \quad \alpha = A_1 \sqrt{\frac{n+1}{2} \frac{a}{\nu}}$$

$$F'(\alpha) \rightarrow 0, \quad \Theta(\alpha) \rightarrow 0, \quad \Phi(\alpha) \rightarrow 1 \quad \text{as} \quad \alpha \rightarrow \infty. \tag{11}$$

Now, setting $F(\xi) = f(\xi - \alpha) = f(\eta)$, $\Theta(\xi) = \theta(\xi - \alpha) = \theta(\eta)$ and $\Phi(\xi) = \phi(\xi - \alpha) = \phi(\eta)$ yields eqs. (8)–(11) in the following forms:

$$f''' - \frac{2n}{n+1}f'^2 + ff'' + We\sqrt{\frac{n+1}{2}}f''f''' - \frac{2}{n+1}M^2f' + \frac{2}{n+1}\lambda[(1 + \beta_t\theta)\theta + N(1 + \beta_c\theta)\theta] = 0, \quad (12)$$

$$[1 + R\{1 + (\theta_f - 1)\theta\}^3\theta']' + \text{Pr}\left(f\theta' + Nb\theta'\phi' + Nt\theta'^2 + \frac{2}{n+1}\delta\theta + EcM^2f'^2\right) = 0, \quad (13)$$

$$\phi'' + Scf\phi' + \frac{Nt}{Nb}\theta'' = 0, \quad (14)$$

$$f(0) = \alpha\left(\frac{1-n}{1+n}\right), \quad f'(0) = 1, \quad \theta'(0) = -\gamma(1 - \theta(0)),$$

$$Nb\phi'(0) + Nt\theta'(0) = 0,$$

$$f'(0) \rightarrow 0, \quad \theta(0) \rightarrow 0, \quad \phi(0) \rightarrow 1. \quad (15)$$

Here the Weissenberg number (We), Hartman number (M), thermal buoyancy parameter (λ), Grashof number in terms of temperature (Gr_x), local Reynolds number (Re_x), nonlinear thermal convection parameter (β_t), nonlinear solutal convection parameter (β_c), ratio of concentration to thermal buoyancy forces (N), Grashof number in terms of concentration (Gr_x^*), temperature ratio parameter (θ_f), radiation parameter (R), Prandtl number (Pr), heat generation ($\delta > 0$), heat absorption ($\delta < 0$), thermophoretic parameter (N_t), Eckert number (Ec), Brownian motion parameter (N_b), Biot number (γ), Schmidt number (Sc) and wall thickness parameter (α) are stated as follows:

$$\begin{aligned} We &= \Gamma\left(\frac{a^3(x+b)^{3n-1}}{\nu}\right)^{1/2}, & M^2 &= \frac{\sigma B_0^2}{\rho_f a}, & \lambda &= \frac{Gr_x}{Re_x^2}, & Gr_x &= \frac{gA_1(T_f - T_\infty)(x+b)^3}{\nu^2}, \\ Re_x &= \frac{u_w}{\nu}(x+b), & \beta_t &= \frac{A_2(T_f - T_\infty)}{A_1}, & \beta_c &= \frac{A_4C_\infty}{A_3}, & N &= \frac{Gr_x^*}{Gr_x}, & Gr_x^* &= \frac{gA_3C_\infty(x+b)^3}{\nu^2}, \\ \theta_f &= \frac{T_f}{T_\infty}, & R &= \frac{16\sigma^*T_\infty^3}{3k^*k}, & \text{Pr} &= \frac{\nu}{\alpha^*}, & \delta &= \frac{Q_0}{(\rho c)_f a}, & N_t &= \frac{\tau D_T(T_f - T_\infty)}{T_\infty\nu}, \\ Ec &= \frac{u_w^2}{c_p(T_f - T_\infty)}, & \gamma &= \frac{h}{k}\sqrt{\frac{\nu}{\alpha^*}}, & N_b &= \frac{\tau D_B C_\infty}{\nu}, & Sc &= \frac{\nu}{D_B}, & \alpha &= A_1\sqrt{\frac{n+1}{2}}\frac{a}{\nu}. \end{aligned} \quad (16)$$

The surface drag force (C_f) and local Nusselt number (Nu) are

$$C_f = \frac{\tau_w}{\rho u_w^2}, \quad Nu = \frac{(x+b)q_w}{k(T_f - T_\infty)}. \quad (17)$$

The following relations will be satisfied for surface shear stress (τ_w) and surface heat flux (q_w), *i.e.*

$$\begin{aligned} \tau_w &= \mu\left(\frac{\partial u}{\partial y} + \frac{\Gamma}{\sqrt{2}}\left(\frac{\partial u}{\partial y}\right)^2\right)\Bigg|_{y=A_1(x+b)^{\frac{1-n}{2}}}, \\ q_w &= -k\left(1 + \frac{16\sigma^*T^3}{3kk^*}\right)\left(\frac{\partial T}{\partial y}\right)\Bigg|_{y=A_1(x+b)^{\frac{1-n}{2}}}. \end{aligned} \quad (18)$$

Replacing eq. (18) in eq. (17) one can obtain

$$\begin{aligned} \sqrt{Re_x}C_f &= \left(\sqrt{\frac{n+1}{2}}f''(0) + We\frac{n+1}{4}(f''(0))^2\right), \\ Nu Re_x^{-1/2} &= -\sqrt{\frac{n+1}{2}}(1 + R\{1 + (\theta_f - 1)\theta(0)\}^3)\theta'(0). \end{aligned} \quad (19)$$

3 Solution methodology

3.1 Computational procedure

This subsection reports the numeric solutions of developed nonlinear systems (12)–(14) subject to boundary conditions (15) through the `bvp4c` technique.

Table 1. Comparative results of surface drag force with [35] through different M , when $n = 1.0$ and $\lambda = 0.0 = \alpha = N = \beta_t = \beta_c = We$.

$-f''(0)$		
M	Ref. [35]	Present results
0.0	1.00000	1.00000
0.2	1.01980	1.01980
0.5	1.11803	1.11803
0.8	1.28063	1.28063
1.0	1.41421	1.41421
1.2	1.56205	1.56205
1.5	1.80303	1.80303

Table 2. Surface drag force ($C_f Re_x^{1/2}$) through distinct values of M , We , λ , β_t and α , when $n = 0.9$, $\delta = 0.1$, $N = N_b = R = \gamma = Ec = 0.5$, $\beta_t = 0.2$, $Sc = \theta_f = 1.1$, $Pr = 1.2$ and $N_t = 0.3$.

M	We	λ	β_t	α	$C_f Re_x^{1/2}$
0.0	1.0	0.5	0.3	0.2	-0.007052794
0.3					-0.05752041
0.5					-0.2758314
0.3	0.0				-0.8824139
	0.4				-0.6697345
	0.8				-0.3299715
	1.0	0.6			-0.3015025
		0.8			-0.2990457
		0.9			-0.2977894
		0.5	0.2		-0.2563464
			0.4		-0.2388347
			0.5		-0.1264708
			0.3	0.0	-0.2882989
				0.1	-0.2395947
				0.5	-0.05607216

3.2 Testing the numerical scheme

The validation of the employed technique with the available published results [36] in a limiting manner is elaborated in this subsection through table 1. It is noticed that our results are in reasonable agreement with [36].

4 Analysis

The salient features of distinct physical variables like Weissenberg number (We), Hartman number (M), thermal buoyancy parameter (λ), nonlinear thermal convection parameter (β_t), nonlinear solutal convection parameter (β_c), ratio of concentration to thermal buoyancy forces (N), temperature ratio parameter (θ_f), radiation parameter (R), Prandtl number (Pr), heat generation parameter ($\delta > 0$), heat absorption parameter ($\delta < 0$), thermophoretic parameter (N_t), Eckert number (Ec), Brownian motion parameter (N_b), Biot number (γ), Schmidt number (Sc) and wall thickness parameter (α) on the non-dimensional temperature (θ), nanoparticle concentration (ϕ), surface drag force ($C_f Re_x^{1/2}$) and Nusselt number ($Nu_x Re_x^{-1/2}$) are illustrated and interpreted in figs. 2–13 and tables 2 and 3.

Figure 2 displays the nature of the temperature field with N_b . Here the strong Brownian diffusion leads to an upsurge in the temperature and the associated thickness of the thermal boundary layer. Small elements in the thermophoresis mechanism are dragged aside from the warm surface to the chilly one. Thus a huge amount of nanoparticles will move apart from the warm surface, which increases the liquid temperature. The radiation impact on the temperature (θ) is exhibited in fig. 3. Clearly the ascending values of the thermal radiation factor (R) enhance the temperature field.

Table 3. Heat transfer rate ($Nu Re_x^{-1/2}$) through distinct values of N_t , N_b , R , θ_f , γ , δ , Pr and Ec , when $n = 0.9$, $We = 1.0$, $\lambda = N = Ec = 0.5$, $\beta_t = M = 0.3$, $\beta_c = \alpha = 0.2$ and $Sc = 1.1$.

N_t	N_b	R	θ_f	γ	δ	Pr	$Nu Re_x^{-1/2}$
0.2	1.0	0.5	0.3	0.2	0.1	1.2	0.201343
0.4							0.207022
0.5							0.210507
0.3	0.2						0.21686
	0.3						0.211506
	0.4						0.206768
	0.5	0.0					0.264135
		0.4					0.227062
		0.8					0.213529
		0.5	1.2				0.179254
			1.3				0.144851
			1.4				0.120746
			1.1	0.3			0.170484
				0.6			0.220432
				0.9			0.232434
				0.5	0.0		0.274422
					0.2		0.337643
					0.4		0.82374
					0.1	1.1	0.196687
						1.3	0.257983
						1.5	0.316771

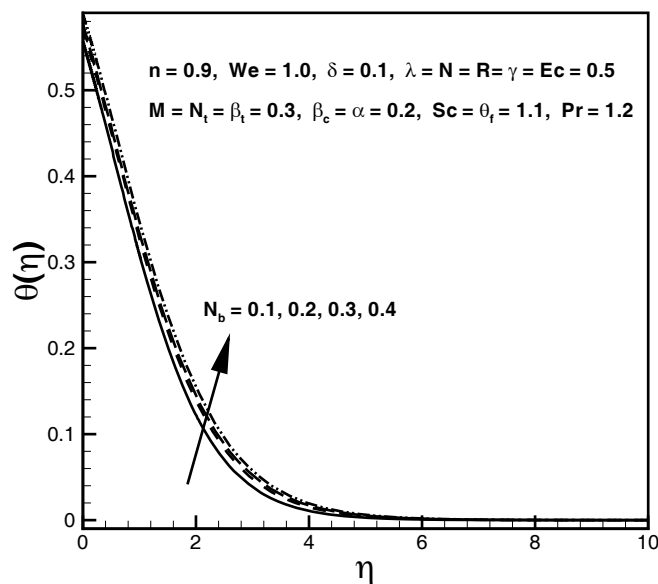


Fig. 2. θ via N_b .

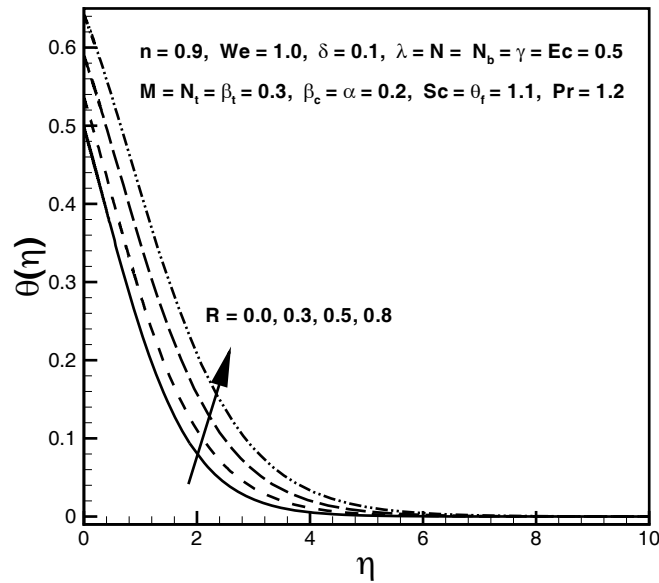


Fig. 3. θ via R .

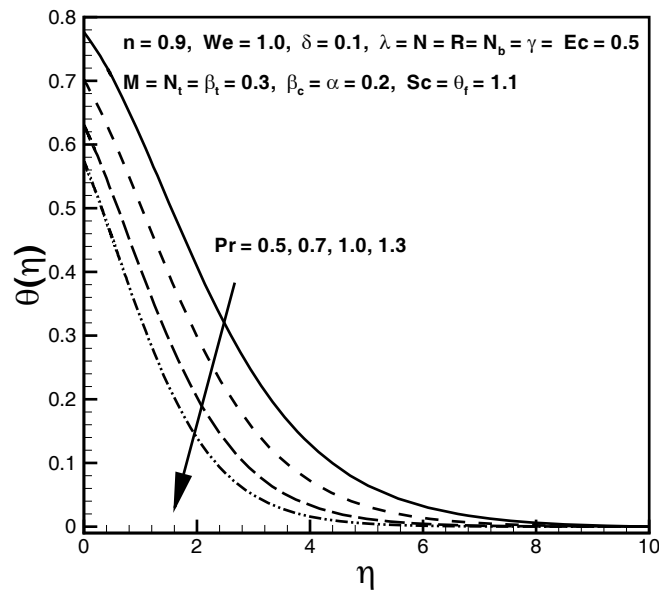


Fig. 4. θ via Pr .

In fact, radiation is the diffusion of heat energy from one section to another. Thermal radiation yields peripheral heat energy to augment the temperature field. Figure 4 shows the effect of the Prandtl number (Pr) on the temperature (θ). Since Pr is contrariwise proportional to the thermal diffusivity, escalating values of Pr yield a weaker thermal diffusivity. Larger Prandtl liquids produce weaker thermal diffusivity and lower Prandtl liquids have stronger thermal diffusivity. Such weaker thermal diffusivity yields a decrease in temperature and its associated boundary layer thickness. The variation of temperature (θ) for different values of θ_f is plotted in fig. 5. It is found that the temperature (θ) and thickness of thermal boundary layer increase through larger θ_f . This happens due to the fact that when we increment θ_f , then T_f upsurges, due to which more heat is transferred to the fluid and so the temperature profile increases. Figure 6 exhibits the effect of γ on the temperature (θ). It is observed that both the temperature and

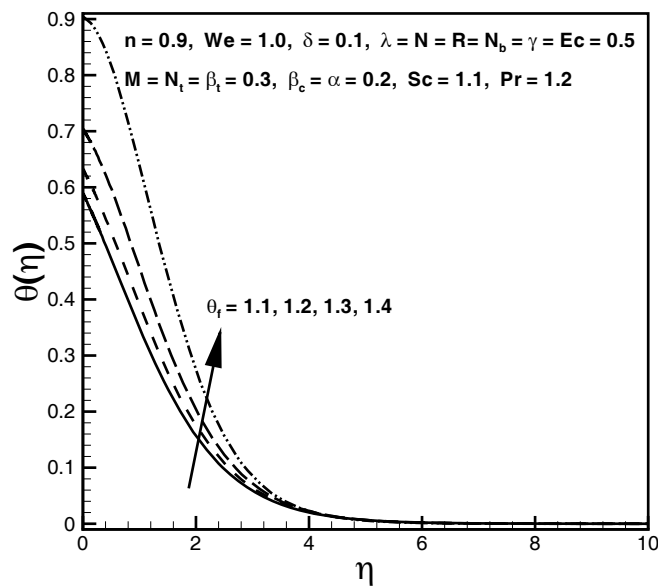


Fig. 5. θ via θ_f .

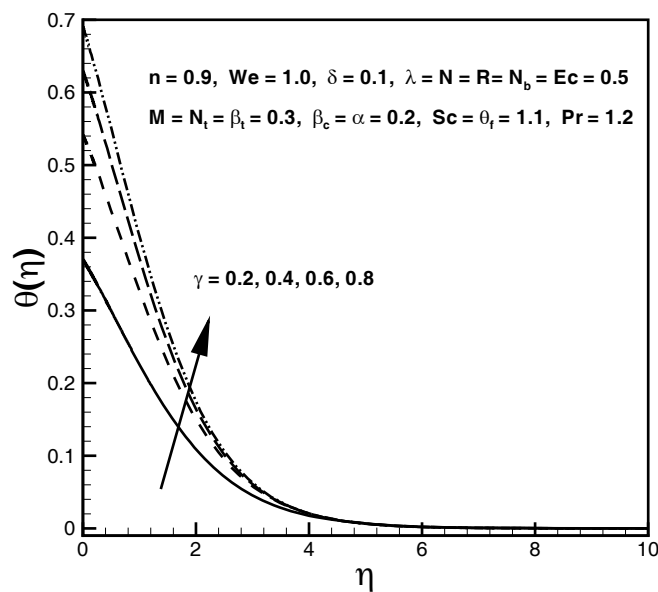


Fig. 6. θ via γ .

associated thickness of thermal layer are higher via larger γ . Physically the heat transfer coefficient is directly related to the Biot number, which is enhanced for larger γ . This enhancement in the heat transfer coefficient leads to higher temperature. Figures 7 and 8 show the influence of heat source or heat generation ($\delta > 0$) and heat sink or heat absorption ($\delta < 0$) on temperature (θ). It is observed that the temperature increases throughout the boundary layer region as ($\delta > 0$) increases; however, a reverse scenario is noticed for heat sink or heat generation ($\delta < 0$). This is because the heat generation or heat source parameter ($\delta > 0$) provides more heat into the fluid, which leads to an intensification of the temperature and associated thermal boundary layer. The variation of Ec on the temperature (θ) is plotted in fig. 9. Here the temperature (θ) and the associated thickness of the thermal layer are increasing functions of Ec . Physically, an increase in Ec means that the heat energy is stored in the fluid due to the frictional

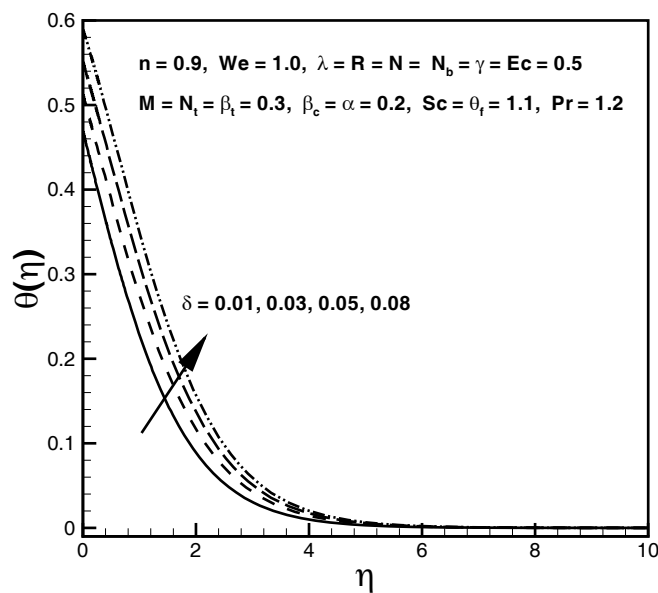


Fig. 7. θ via $\delta > 0$.

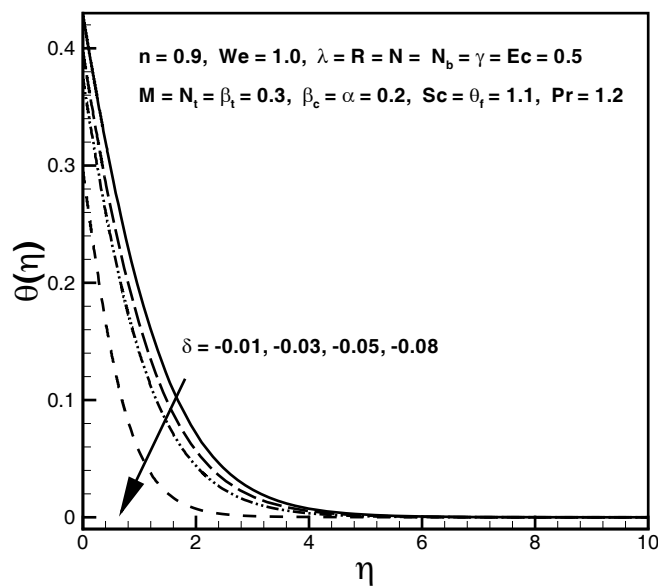


Fig. 8. θ via $\delta < 0$.

or drag forces. As a result, the fluid temperature (θ) enhances. Figure 10 is drawn to have a good knowledge of the Schmidt number (Sc) on nanoparticle concentration (ϕ). As the Schmidt number is inversely proportional to the mass diffusion, so a larger Sc shows a decay in the nanoparticle concentration (ϕ) and its linked thickness of boundary layer. The behavior of Pr versus nanoparticle concentration (ϕ) is addressed through fig. 11. Clearly, a larger Pr leads to a decay in the nanoparticle concentration (ϕ) and the corresponding layer thickness. Figure 12 shows the variation of N_t on nanoparticle concentration (ϕ). Here rising values of N_t yield an increment in nanoparticle concentration (ϕ). Physically, larger N_t creates an additional force and makes nanoparticles pass from the hottest to the coolest region. Therefore nanoparticle concentration (ϕ) and its linked boundary layer are higher for N_t . The variation in the nanoparticle concentration (ϕ) as a result of the changes in N_b is shown in fig. 13. It is found that increasing values of

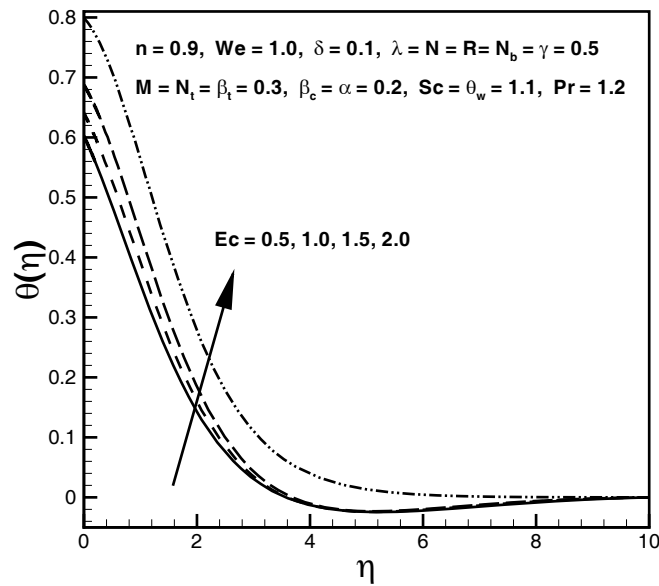


Fig. 9. θ via Ec .

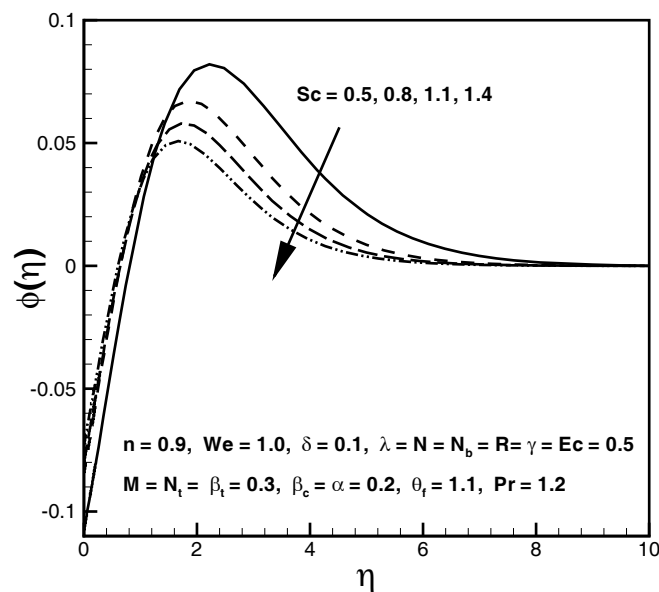


Fig. 10. ϕ via Sc .

N_b reduce the nanoparticle concentration (ϕ). In fact, higher values of the Brownian motion factor (N_b) help to heat the fluid near the boundary. This leads to a decrease in the nanoparticle concentration (ϕ).

Table 2 is built to study the influences of distinct variables, like Ha , We , λ , β_t and α , on the surface drag force ($C_f\sqrt{Re_x}$). Here the surface drag force is an increasing function of Ha , whereas an opposite trend is noted for We , λ , β_t and α . Table 3 elucidates the influence of the Nusselt number for distinct values of N_t , N_b , R , θ_f , γ , δ and Pr . Increasing values of N_t , γ , δ and Pr lead to an enhancement in the Nusselt number, while a reverse behavior is noted for N_b , R and θ_f .

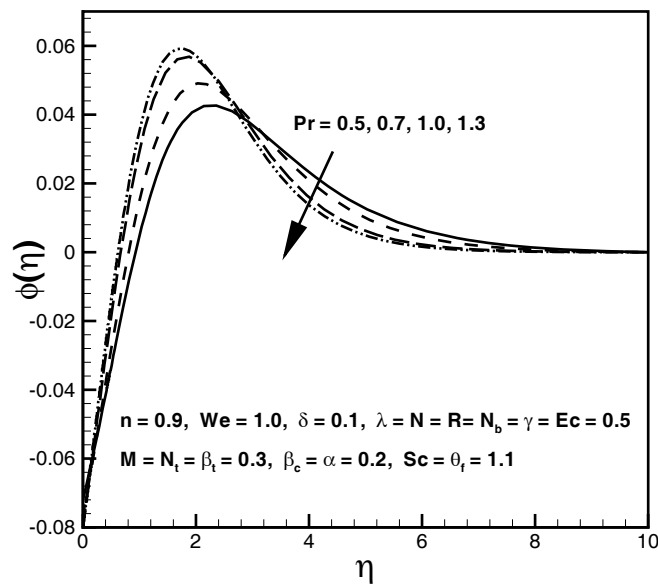


Fig. 11. θ via Pr.

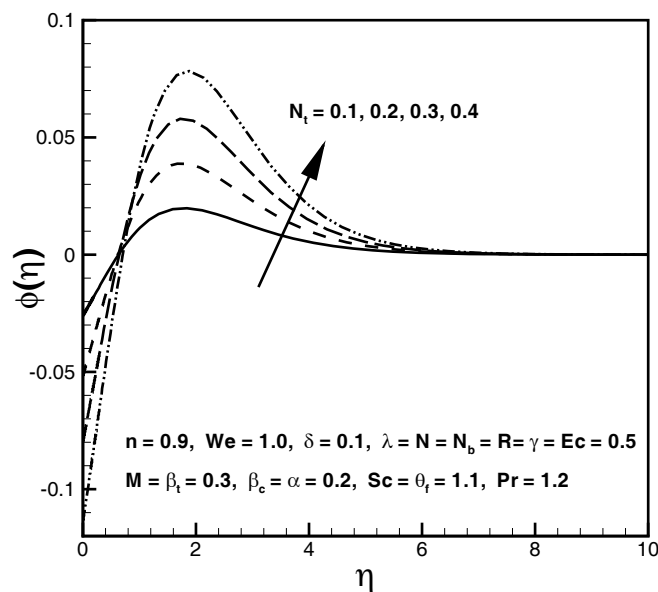


Fig. 12. θ via N_t .

5 Conclusions

Here we explored the interaction of convective heating in a hydromagnetic Williamson nanoliquid induced due to a slendering surface. The nonlinear thermal radiation, heat generation/absorption and Joule heating are also considered in the heat transport expression. The effect of emerging variables on dimensionless quantities is illustrated in detail. The following key points are drawn from this work:

- A larger Pr corresponds to a decay in temperature and thermal boundary layer thicknesses, whereas a reverse situation is noticed via larger γ . This reverse behavior is due to weaker thermal diffusivity, in the case of the Prandtl number, and to a larger heat transfer coefficient, in the case of the Biot number.

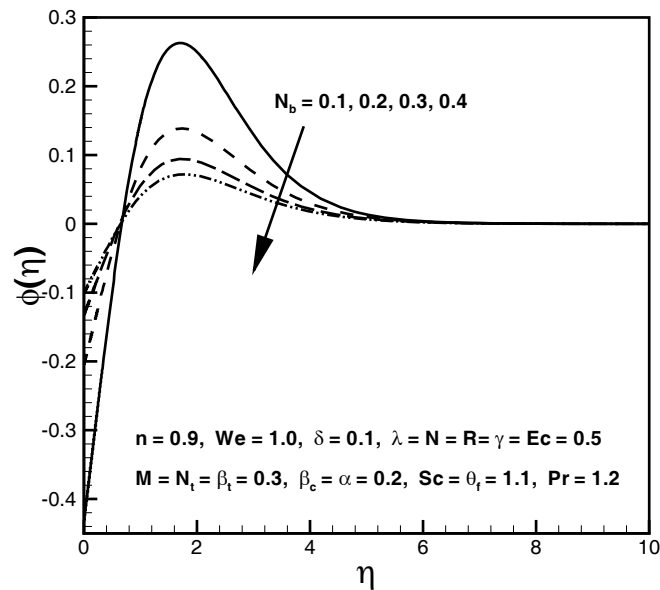


Fig. 13. θ via N_b .

- Both temperature distribution and associated thermal layer are enhanced through larger Eckert number (Ec), heat generation parameter (δ), radiation parameter (R) and temperature ratio parameter (θ_f); however temperature distribution and related thermal layer are lower for heat absorption parameter (δ).
- The increasing thermophoretic parameter increases the profiles of temperature and concentration.
- Consideration of higher Schmidt number (Sc) and Brownian motion factor (N_b) yields lower concentration.
- The values of surface drag force are higher due to the increasing values of the Hartman number (M); however the drag force is reduced for larger Weissenberg number (We), thermal buoyancy parameter (λ), nonlinear thermal convection parameter (β_t) and wall thickness parameter (α).
- Larger values of thermophoretic and Brownian motion parameters lead to an enhancement in the temperature profile due to higher thermal conductivity fluid.
- The local Nusselt number enhances for higher values of Nt , δ and Pr .

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