

NONLOCAL PROBLEM FOR A MIXED TYPE FOURTH-ORDER DIFFERENTIAL EQUATION WITH HILFER FRACTIONAL OPERATOR

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Abstract: In this paper, we consider a non-self-adjoint boundary value problem for a fourth-order differential equation of mixed type with Hilfer operator of fractional integro-differentiation in a positive rectangular domain and with spectral parameter in a negative rectangular domain. The mixed type differential equation under consideration is a fourth order differential equation with respect to the second variable. Regarding the first variable, this equation is a fractional differential equation in the positive part of the segment, and is a second-order differential equation with spectral parameter in the negative part of this segment. A rational method of solving a nonlocal problem with respect to the Hilfer operator is proposed. Using the spectral method of separation of variables, the solution of the problem is constructed in the form of Fourier series. Theorems on the existence and uniqueness of the problem are proved for regular values of the spectral parameter. For sufficiently large positive integers in unique determination of the integration constants in solving countable systems of differential equations, the problem of small denominators arises. Therefore, to justify the unique solvability of this problem, it is necessary to show the existence of values of the spectral parameter such that the quantity we need is separated from zero for sufficiently large n . For irregular values of the spectral parameter, an infinite number of solutions in the form of Fourier series are constructed. Illustrative examples are provided.

Keywords: Mixed type equation, Non-self-adjoint boundary value problem, Hilfer operator, Mittag-Leffler function, Spectral parameter, Solvability.

1. Problem statement

In a rectangular domain $\Omega = \{(t, x) : -a < t < b, 0 < x < 1\}$, we consider the partial differential equation of mixed type

$$0 = \begin{cases} \left(D^{\alpha, \gamma} + \frac{\partial^4}{\partial x^4}\right)U(t, x), & (t, x) \in \Omega_1, \\ \left(\frac{\partial^2}{\partial t^2} + \omega^2 \frac{\partial^4}{\partial x^4}\right)U(t, x), & (t, x) \in \Omega_2, \end{cases} \quad (1.1)$$

where $\Omega_1 = \Omega \cap (t > 0)$, $\Omega_2 = \Omega \cap (t < 0)$, ω is positive spectral parameter, a and b are positive real numbers,

$$D^{\alpha, \gamma} = J_{0+}^{\gamma-\alpha} \frac{d}{dt} J_{0+}^{1-\gamma} \quad (0 < \alpha \leq \gamma \leq 1)$$

is the Hilfer operator, and

$$I_{0+}^{\nu} \varphi(t) = \frac{1}{\Gamma(\nu)} \int_0^t (t-\tau)^{\nu-1} \varphi(\tau) d\tau, \quad \nu > 0$$

is the Riemann–Liouville integral operator [2, pp. 112, 113].

Nonlocal problem. It is required to find a function $U(t, x)$, which belongs to the class

$$t^{1-\gamma} \frac{\partial^k U}{\partial x^k} \in C(\overline{\Omega}_1), \quad \frac{\partial^k U}{\partial x^k} \in C(\overline{\Omega}_2), \quad D^{\alpha, \gamma} U \in C(\Omega_1), \quad U_{tt} \in C(\Omega_2), \quad U_{xxxx} \in C(\Omega_1 \cup \Omega_2), \quad (1.2)$$

$k = \overline{0, 3}$ and satisfies the homogeneous equation (1.1) in the domain $\Omega_1 \cup \Omega_2$, the homogeneous boundary value conditions

$$U|_{x=0} = \frac{\partial^2 U}{\partial x^2}|_{x=1} = 0, \quad \frac{\partial^k U}{\partial x^k}|_{x=0} = \frac{\partial^k U}{\partial x^k}|_{x=1}, \quad k = 1, 3, \quad t \neq 0, \quad (1.3)$$

the nonlocal condition

$$U(-a, x) = U(b, x) + \varphi(x), \quad 0 \leq x \leq 1, \quad (1.4)$$

and the gluing conditions

$$\lim_{t \rightarrow +0} J_{0+}^{1-\gamma} U(t, x) = \lim_{t \rightarrow -0} U(t, x), \quad \lim_{t \rightarrow +0} J_{0+}^{1-\alpha} \frac{d}{dt} J_{0+}^{1-\gamma} U(t, x) = \lim_{t \rightarrow -0} U_t(t, x), \quad (1.5)$$

where $\varphi(x)$ is a given sufficiently smooth function.

Let $(t_0; b) \subset \mathbb{R}^+ \equiv [0; \infty)$ be a finite interval, and let $\alpha > 0$. The Riemann–Liouville α -order fractional integral of a function f is defined as follows:

$$I_{t_0+}^{\alpha} f(t) = \frac{1}{\Gamma(\alpha)} \int_{t_0}^t (t-s)^{\alpha-1} f(s) ds, \quad t \in (t_0; b),$$

where $\Gamma(\alpha)$ is the Gamma function [2, p. 112].

Let $n-1 < \alpha \leq n$, $n \in \mathbb{N}$. The Riemann–Liouville α -order fractional derivative of a function f is defined as follows [9, Vol. 1, p. 27]:

$$D_{t_0+}^{\alpha} f(t) = \frac{d^n}{dt^n} I_{t_0+}^{n-\alpha} f(t), \quad t \in (t_0; b).$$

The Caputo α -order fractional derivative of a function f is defined [9, Vol. 1, p. 34] by

$${}^*D_{t_0+}^{\alpha} f(t) = I_{t_0+}^{n-\alpha} f^{(n)}(t) = \frac{1}{\Gamma(n-\alpha)} \int_{t_0}^t \frac{f^{(n)}(s) ds}{(t-s)^{\alpha-n+1}}.$$

Both the derivatives are reduced to the n th order derivatives for $\alpha = n \in \mathbb{N}$ [9, Vol. 1, pp. 27, 34]:

$$D_{t_0+}^n f(t) = {}^*D_{t_0+}^n f(t) = \frac{d^n f}{dt^n}.$$

The so-called generalized Riemann–Liouville fractional derivative (referred to as the Hilfer fractional derivative) of order α , $n-1 < \alpha \leq n$, $n \in \mathbb{N}$, and type β , $0 \leq \beta \leq 1$, is defined by the following composition of three operators: [2, p. 113]:

$$D_{t_0+}^{\alpha, \beta} f(t) = I_{t_0+}^{\beta(n-\alpha)} \frac{d^n}{dt^n} I_{t_0+}^{(1-\beta)(n-\alpha)} f(t).$$

For $\beta = 0$, this operator is reduced to the Riemann–Liouville fractional derivative ($D_{t_0+}^{\alpha,0} = D_{t_0+}^{\alpha}$) and the case $\beta = 1$ corresponds to the Caputo fractional derivative: $D_{t_0+}^{\alpha,1} = {}_*D_{t_0+}^{\alpha}$.

Let $t_0 = 0$ and $\gamma = \alpha + \beta n - \alpha\beta$. It is easy to see that $\alpha \leq \gamma \leq n$. Then it is convenient to use another notation for the operator $D_{0+}^{\alpha,\beta} f(t)$:

$$D^{\alpha,\gamma} f(t) = D_{0+}^{\alpha,\beta} f(t). \quad (1.6)$$

For the first time, the generalized Riemann–Liouville operator was introduced in [2] by R. Hilfer on the basis of fractional time evolutions that arise during the transition from the microscopic scale to the macroscopic time scale. Using the integral transforms, he investigated the Cauchy problem for the generalized diffusion equation, the solution of which is presented in the form of the Fox H-function. We also note [10, 11], where the generalized Riemann–Liouville operator was used in studying dielectric relaxation in glass-forming liquids with different chemical compositions.

In [23], boundary value problems for a fractional diffusion equation with the Hilfer fractional derivative in finite and infinite domains were studied. In the finite domain, the spectral method and the Laplace transform method were used for solving the problem. In the domain infinite with respect to the spatial variable, the Cauchy problem was solved by the Fourier–Laplace integral transform method.

In [12], the properties of the generalized Riemann–Liouville operator were investigated in a special functional space, and an operational method was developed for solving fractional differential equations with this operator. Based on the results of [12], the authors of [15] have developed an operational method for solving fractional differential equations containing a finite linear combination of the generalized Riemann–Liouville operators with various parameters. In [17], the problem of source identification was studied for the generalized diffusion equation with the operator $D^{\alpha,\gamma}$. We also note the work [4], in which inverse problems were investigated for a generalized fourth-order parabolic equation with the operator $D^{\alpha,\gamma}$.

The construction of various models of theoretical physics problems by the aid of fractional calculus is described in [9, Vols. 4, 5], [16, 26]. A specific physical interpretation of the Hilfer fractional derivative, describing the random motion of a particle moving on the real line at Poisson paced times with finite velocity is given in [25]. A detailed review of the application of fractional calculus in solving applied problems is given in [9, Vols. 6–8], [19]. More detailed information as well as a bibliography related to the theory of fractional integro-differentiation, including the Hilfer fractional derivative, can be found in the recently published monograph [24]. In [7], the boundary value problems for the generalized modified moisture transfer equation and difference methods for their numerical implementation were considered.

Nonlocal problems can arise in studying various problems of mathematical biology, predicting soil moisture, problems of plasma. Note that nonlocal conditions of the type (1.3) take place in modeling the problems of the flow around a profile by a subsonic velocity stream with a supersonic zone [20]. More detailed information on nonlocal problems can be found in the monograph [18]. We would like to note some works [14, 30–32], where nonlocal problems for partial differential and integro-differential equations with derivatives of integer or fractional orders were studied.

As for the equations of mixed type, we note the work [8], where I.M. Gel'fand considered an example of gas motion in a channel surrounded by a porous medium, and the gas motion in a channel was described by a wave equation, while the diffusion equation was posed outside the channel. Ya. S. Uflyand considered a problem on the propagation of electric oscillations in compound lines when the losses on a semi-infinite line were neglected and the rest of the line was treated as a cable with no leaks [28]. He reduced this problem to a mixed parabolic-hyperbolic type equation. In [27], a hyperbolic-parabolic system arising in pulse combustion was investigated.

Nonlocal problems for partial differential equations of mixed type were studied by many authors, in particular, in [13, 21, 22, 29, 33]. We would like to note also the results on nonlocal problems

for parabolic-hyperbolic type equations with fractional order derivatives [1, 3]. But these listed works relate mainly to nonlocal problems for fractional mixed type equations of second order. As for mixed fourth-order equations with derivatives of integer or fractional orders, nonlocal problems in such formulation have not been previously studied.

In this paper, we consider a non-self-adjoint boundary value problem for a mixed type fourth-order differential equation with Hilfer operator of fractional integro-differentiation. The spectral method of separation of variables is used taking into account the features of the fractional integro-differentiation operator. We study the solvability of the nonlocal problem (1.1)–(1.5) for various values of the spectral parameter. This work is a further development and generalization of the results of [5, 6, 20].

2. Ordinary differential equation with Hilfer operator

We consider the Cauchy problem for a differential equation of fractional order with the operator $D^{\alpha, \gamma}$

$$\begin{cases} D^{\alpha, \gamma} u(t) = \lambda u(t) + f(t), & t \in (0, \ell), \\ \lim_{t \rightarrow +0} J_{0+}^{1-\gamma} u(t) = u_0, \end{cases} \quad (2.1)$$

where $f(t)$ is a given continuous function and $u_0 = \text{const}$.

Note that the Laplace method was used for solving this problem in [4]. In [15], a solution was found by the operational calculus for a problem more general than (2.1) in a specially constructed functional space. In our work, in contrast to these studies, we use a more rational way to solve problem (2.1), which allows us to obtain an explicit solution.

We prove the following Lemma.

Lemma 1. *Assume that $f(t) \in C(0; \ell] \cap L_1(0; \ell)$. Then a solution of problem (2.1) $u(t) \in C(0; \ell] \cap L_1(0; \ell)$ is representable as follows:*

$$u(t) = u_0 t^{\gamma-1} E_{\alpha, \gamma}(\lambda t^\alpha) + \int_0^t (t-\tau)^{\alpha-1} E_{\alpha, \alpha}(\lambda(t-\tau)^\alpha) f(\tau) d\tau, \quad (2.2)$$

where

$$E_{\alpha, \beta}(z) = \sum_{k=0}^{\infty} \frac{z^k}{\Gamma(\alpha k + \beta)}, \quad z, \alpha, \beta \in \mathbb{C}, \quad \text{Re}(\alpha) > 0$$

is the Mittag-Leffler function [9, Vol. 1, pp. 269–295].

P r o o f. By virtue of the formula (1.6), we rewrite the differential equation of problem (2.1) in the form

$$J_{0+}^{\gamma-\alpha} D_{0+}^{\gamma} u(t) = \lambda u(t) + f(t).$$

Further, applying the operator J_{0+}^{α} to both sides of this equation and taking into account the linearity of this operator and the following formula [15]:

$$J_{0+}^{\beta} D_{0+}^{\beta} u(t) = u(t) - \frac{1}{\Gamma(\gamma)} J_{0+}^{1-\beta} u(t)|_{t=0} t^{\beta-1},$$

we obtain

$$u(t) = \frac{u_0}{\Gamma(\gamma)} t^{\gamma-1} + J_{0+}^{\alpha} f(t) + \lambda J_{0+}^{\alpha} u(t). \quad (2.3)$$

Using the lemma from [6, p. 123], we represent the solution of equation (2.3) as

$$u(t) = \frac{u_0}{\Gamma(\gamma)} t^{\gamma-1} + J_{0+}^\alpha f(t) + \lambda \int_0^t (t-\tau)^{\alpha-1} E_{\alpha,\alpha}(\lambda(t-\tau)^\alpha) \left[\frac{u_0}{\Gamma(\gamma)} \tau^{\gamma-1} + J_{0+}^\alpha f(\tau) \right] d\tau. \tag{2.4}$$

We rewrite representation (2.4) as the sum of two expressions $u(t) = I_1(t) + I_2(t)$, where

$$I_1(t) = u_0 \left[\frac{t^{\gamma-1}}{\Gamma(\gamma)} + \frac{\lambda}{\Gamma(\gamma)} \int_0^t (t-\tau)^{\alpha-1} E_{\alpha,\alpha}(\lambda(t-\tau)^\alpha) \tau^{\gamma-1} d\tau \right], \tag{2.5}$$

$$I_2(t) = J_{0+}^\alpha f(t) + \lambda \int_0^t (t-\tau)^{\alpha-1} E_{\alpha,\alpha}(\lambda(t-\tau)^\alpha) J_{0+}^\alpha f(\tau) d\tau. \tag{2.6}$$

We make the change of variables $s = t - \tau$ in formula (2.5) and use the following formulas [9, Vol. 1, pp. 269–295]:

$$E_{\alpha,\mu}(z) = \frac{1}{\Gamma(\mu)} + z E_{\alpha,\mu+\alpha}(t), \quad \alpha > 0, \quad \mu > 0, \tag{2.7}$$

$$\frac{1}{\Gamma(\nu)} \int_0^z (z-t)^{\nu-1} E_{\alpha,\beta}(\lambda t^\alpha) t^{\beta-1} dt = z^{\beta+\nu-1} E_{\alpha,\beta+\nu}(\lambda z^\alpha), \quad \nu > 0, \quad \beta > 0. \tag{2.8}$$

Then we obtain the following representation for integral (2.5):

$$I_1(t) = u_0 t^{\gamma-1} E_{\alpha,\gamma}(\lambda t^\alpha). \tag{2.9}$$

The integral in the formula (2.6) is transformed as follows:

$$\begin{aligned} & \int_0^t (t-\tau)^{\alpha-1} E_{\alpha,\alpha}(\lambda(t-\tau)^\alpha) J_{0+}^\alpha f(\tau) d\tau = \\ &= \frac{1}{\Gamma(\alpha)} \int_0^t (t-\tau)^{\alpha-1} E_{\alpha,\alpha}(\lambda(t-\tau)^\alpha) d\tau \int_0^\tau (\tau-s)^{\alpha-1} f(s) ds = \\ &= \frac{1}{\Gamma(\alpha)} \int_0^t f(s) ds \int_s^t (t-\tau)^{\alpha-1} (\tau-s)^{\alpha-1} E_{\alpha,\alpha}(\lambda(t-\tau)^\alpha) d\tau. \end{aligned} \tag{2.10}$$

In view of (2.8), the second integral in the latter equality of formula (2.10) can be written as

$$\int_s^t (t-\tau)^{\alpha-1} (\tau-s)^{\alpha-1} E_{\alpha,\alpha}(\lambda(t-\tau)^\alpha) d\tau = \Gamma(\alpha) (t-\tau)^{2\alpha-1} E_{\alpha,2\alpha}(\lambda(t-\tau)^\alpha).$$

Then, taking into account (2.7), we represent formula (2.6) in the following form:

$$I_2(t) = \int_0^t (t-\tau)^{\alpha-1} E_{\alpha,\alpha}(\lambda(t-\tau)^\alpha) f(\tau) d\tau. \tag{2.11}$$

Substituting (2.9) and (2.11) into the sum $u(t) = I_1(t) + I_2(t)$, we obtain formula (2.2). The lemma is proved. □

3. Uniqueness of solution of the nonlocal problem

We study this problem by the spectral method of separating variables and seek particular solutions of the nonlocal problem in the form of a product of two functions $U(t, x) = u(t) \cdot \vartheta(x)$. From equation (1.1) and boundary value conditions (1.3), we arrive at the following spectral problem:

$$\vartheta^{IV}(x) - \lambda^4 \vartheta(x) = 0, \quad \vartheta(0) = \vartheta''(1) = 0, \quad \vartheta'(1) = \vartheta'(1), \quad \vartheta'''(1) = \vartheta'''(1),$$

where λ^4 is the constant of separation, $0 < \lambda = \text{const}$.

As follows from the results of [5], this spectral problem is non-self-adjoint and has a complete system of eigenfunctions of the following form in the space $L_2(0; 1)$:

$$\begin{aligned} \vartheta_0(x) = 2x, \quad \vartheta_{n1}(x) = 2 \sin \lambda_n x, \quad \vartheta_{n2}(x) = \frac{e^{\lambda_n x} - e^{\lambda_n(1-x)}}{e^{\lambda_n} - 1} + \cos \lambda_n x, \\ \lambda_n = 2\pi n, \quad n \in \mathbb{N}. \end{aligned} \quad (3.1)$$

System (3.1) forms a Riesz basis in $L_2(0; 1)$. In [5], it was also proved that there exists a biorthogonal system of functions with (3.1):

$$\eta_0(x) = 1, \quad \eta_{n1}(x) = \frac{e^{\lambda_n x} + e^{\lambda_n(1-x)}}{e^{\lambda_n} - 1} + \sin 2\pi n x, \quad \eta_{n2}(x) = 2 \cos \lambda_n x. \quad (3.2)$$

System (3.2) also forms a Riesz basis in $L_2(0; 1)$.

Let $U(t, x)$ be a solution of the nonlocal problem. We consider the functions

$$u_0^+(t) = \int_0^1 U(t, x) dx, \quad u_{ni}^+(t) = \int_0^1 U(t, x) \eta_{ni}(x) dx, \quad t > 0, \quad (3.3)$$

$$u_0^-(t) = \int_0^1 U(t, x) dx, \quad u_{ni}^-(t) = \int_0^1 U(t, x) \eta_{ni}(x) dx, \quad i = 1, 2, \quad t < 0, \quad (3.4)$$

where the functions $\eta_0(x)$ and $\eta_{ni}(x)$, $i = 1, 2$, are defined in (3.2).

Applying the operator $D^{\alpha, \gamma}$ with respect to t to both sides of equality (3.3), differentiating (3.4) twice with respect to t , and taking into account equation (1.1), we obtain differential equations with respect to the functions $u_0^\pm(t)$ and $u_{ni}^\pm(t)$, $i = 1, 2$:

$$D^{\alpha, \gamma} u_0^+(t) = 0, \quad D^{\alpha, \gamma} u_{ni}^+(t) + \lambda_n^4 u_{ni}^+(t) = 0, \quad i = 1, 2, \quad t > 0, \quad (3.5)$$

$$\frac{d^2}{dt^2} u_0^-(t) = 0, \quad \frac{d^2}{dt^2} u_{ni}^-(t) + \lambda_n^4 \omega^2 u_{ni}^-(t) = 0, \quad i = 1, 2, \quad t < 0. \quad (3.6)$$

The general solutions of these differential equations (3.5) and (3.6) have the form

$$u_0^\pm(t) = \begin{cases} \frac{A_0}{\Gamma(\gamma)} t^{\gamma-1}, & t > 0, \\ B_0 t + C_0, & t < 0, \end{cases} \quad u_{ni}^\pm(t) = \begin{cases} A_{ni} t^{\gamma-1} E_{\alpha, \gamma}(-\lambda_n^4 t^\alpha), & t > 0, \\ B_{ni} \sin \lambda_n^2 \omega t + C_{ni} \cos \lambda_n^2 \omega t, & t < 0, \end{cases} \quad (3.7)$$

where A_0 , B_0 , C_0 , A_{ni} , B_{ni} , and C_{ni} are arbitrary constants, $i = 1, 2$, $n = 1, 2, \dots$.

Taking into account conditions (1.4) and (1.5), we conclude from (3.3) and (3.4) that the functions $u_0^\pm(t)$ and $u_{ni}^\pm(t)$, $i = 1, 2$, in (3.7) must satisfy the following conditions:

$$\lim_{t \rightarrow +0} J_{0+}^{1-\gamma} u_0^+(t) = \lim_{t \rightarrow -0} u_0^-(t), \quad \lim_{t \rightarrow +0} J_{0+}^{1-\alpha} \left(\frac{d}{dt} J_{0+}^{1-\gamma} u_0^+(t) \right) = \lim_{t \rightarrow -0} \frac{du_0^-(t)}{dt}, \quad (3.8)$$

$$\lim_{t \rightarrow +0} J_{0+}^{1-\gamma} u_{ni}^+(t) = \lim_{t \rightarrow -0} u_{ni}^-(t), \quad \lim_{t \rightarrow +0} J_{0+}^{1-\alpha} \left(\frac{d}{dt} J_{0+}^{1-\gamma} u_{ni}^+(t) \right) = \lim_{t \rightarrow -0} \frac{d u_{ni}^-(t)}{dt}, \quad (3.9)$$

$$u_0^-(-a) = u_0^+(b) + \varphi_0, \quad u_{ni}^-(-a) = u_{ni}^+(b) + \varphi_{ni}, \quad i = 1, 2, \quad (3.10)$$

where

$$\varphi_0 = \int_0^1 \varphi(x) dx, \quad \varphi_{ni} = \int_0^1 \varphi(x) \eta_{ni}(x) dx, \quad i = 1, 2, \quad n = 1, 2, \dots$$

Therefore, we obtain the following systems of algebraic equations:

$$\begin{cases} A_0 = C_0, & B_0 = 0, \\ -B_0 a + C_0 = \frac{A_0}{\Gamma(\gamma)} b^{\gamma-1} + \varphi_0, \end{cases} \quad (3.11)$$

$$\begin{cases} A_{ni} = C_{ni}, & \omega B_{ni} = -\lambda_n^2 A_{ni}, \\ -B_{ni} \sin \lambda_n^2 \omega a + C_{ni} \cos \lambda_n^2 \omega a - A_{ni} b^{\gamma-1} E_{\alpha, \gamma}(-\lambda_n^4 b^\alpha) = \varphi_{ni}. \end{cases} \quad (3.12)$$

Each of systems (3.11) and (3.12) has a unique solution

$$C_0 = A_0, \quad B_0 = 0, \quad A_0 = \frac{\varphi_0}{\Delta_0}, \quad C_{ni} = A_{ni} = \frac{\varphi_{ni}}{\Delta_n(\omega)}, \quad B_{ni} = -\frac{\lambda_n^2}{\omega} \frac{\varphi_{ni}}{\Delta_n(\omega)}, \quad (3.13)$$

if the following condition holds for all $n \in \mathbb{N}_0 = \mathbb{N} \cup \{0\}$:

$$\Delta_n(\omega) = \lambda_n^2 \omega \sin \lambda_n^2 \omega a + \cos \lambda_n^2 \omega a - b^{\gamma-1} E_{\alpha, \gamma}(-\lambda_n^4 b^\alpha) \neq 0. \quad (3.14)$$

Substituting (3.13) into (3.7), we obtain the representation

$$u_0^\pm(t) = \begin{cases} \frac{\varphi_0}{\Gamma(\gamma) \Delta_0} t^{\gamma-1}, & t > 0, \\ \frac{\varphi_0}{\Delta_0}, & t \leq 0, \end{cases} \quad (3.15)$$

$$u_{ni}^\pm(t) = \begin{cases} \frac{\varphi_{ni}}{\Delta_n(\omega)} t^{\gamma-1} E_{\alpha, \gamma}(-\lambda_n^4 t^\alpha), & t > 0, \\ \frac{\varphi_{ni}}{\Delta_n(\omega)} \left(\cos \lambda_n^2 \omega t - \frac{\lambda_n^2}{\omega} \sin \lambda_n^2 \omega t \right), & t \leq 0. \end{cases} \quad (3.16)$$

We show the uniqueness of the solution of the nonlocal problem under condition (3.14). Suppose the opposite. Let the nonlocal problem have two different solutions $U_1(t, x)$ and $U_2(t, x)$, and let $U(t, x) = U_1(t, x) - U_2(t, x)$. It is not difficult to see that $U(t, x)$ is a solution of the homogeneous nonlocal problem ($\varphi(x) = 0$). This is why one only needs to prove that the homogeneous problem has only the trivial solution.

Suppose that condition (3.14) holds and $\varphi(x) \equiv 0$. Then $\varphi_0 = 0$, $\varphi_{ni} = 0$, $i = 1, 2$, and the representations (3.3), (3.4) and (3.15), (3.16) yield

$$\int_0^1 t^{1-\gamma} U(t, x) dx = 0, \quad \int_0^1 t^{1-\gamma} U(t, x) \eta_{ni}(x) dx = 0, \quad t \in [0; b],$$

$$\int_0^1 U(t, x) dx = 0, \quad \int_0^1 U(t, x) \eta_{ni}(x) dx = 0, \quad t \in [-a; 0], \quad i = 1, 2.$$

Further, taking into account the completeness of system (3.2) in the space $L_2(0; 1)$, we conclude that $U(t, x) = 0$ almost everywhere on $[0; 1]$ for all $t \in [-a; b]$. Since $t^{1-\gamma}U(t, x) \in C(\overline{\Omega}_1)$ and $U(t, x) \in C(\overline{\Omega}_2)$, we have $t^{1-\gamma}U(t, x) \equiv 0$ in the domain $\overline{\Omega}$. Therefore, the solution of the nonlocal problem is unique in the domain $\overline{\Omega}$.

Thus, we have proved the following theorem.

Theorem 1. *Suppose that there exists a solution of the nonlocal problem. This solution is unique if condition (3.14) holds for all $n \in \mathbb{N}_0$.*

4. Existence of a solution of the nonlocal problem

Now we consider the case when condition (3.14) is violated. Let $\Delta_m(\omega) = 0$ for all $\omega, \gamma \in (0; 1)$ and $n = m$. Then the homogeneous nonlocal problem ($\varphi(x) \equiv 0$) has a nontrivial solution

$$V_{mi}^\pm(t, x) = v_m^\pm(t) \vartheta_{mi}(x), \quad i = 1, 2, \quad (4.1)$$

where

$$v_m^\pm(t) = \begin{cases} t^{\gamma-1} E_{\alpha, \gamma}(-\lambda_m^4 t^\alpha), & t > 0, \\ \cos \lambda_m^2 \omega t - \frac{\lambda_m^2}{\omega} \sin \lambda_m^2 \omega t, & t < 0. \end{cases}$$

It is easy to verify that, for $\gamma = 1$, the function $V(t, x) = x$ is also a nontrivial solution of the homogeneous nonlocal problem.

From $\Delta_n(\omega) = 0$, we come to the trigonometric equation

$$\sqrt{1 + \omega^2 \lambda_n^4} \sin(\lambda_n^2 \omega a + \rho_n) - b^{\gamma-1} E_{\alpha, \gamma}(-\lambda_n^4 b^\alpha) = 0, \quad (4.2)$$

where $\rho_n = \arcsin\left(1/\sqrt{1 + \omega^2 \lambda_n^4}\right)$ and $\rho_n \rightarrow 0$ as $n \rightarrow +\infty$. Hence, we conclude that the expression $\Delta_n(\omega)$ is zero only if

$$\omega = \frac{1}{\lambda_n^2 a} \left[(-1)^k \arcsin \frac{b^{\gamma-1} E_{\alpha, \gamma}(-\lambda_n^4 b^\alpha)}{\sqrt{1 + \omega^2 \lambda_n^4}} + \pi k - \rho_n \right], \quad k = 1, 2, \dots$$

The set \mathfrak{S} of positive solutions of trigonometric equation (4.2) is called the set of irregular values of the spectral parameter ω .

The set of remaining values of the spectral parameter $\aleph = (0; \infty) \setminus \mathfrak{S}$ is called the set of regular values of the spectral parameter ω .

Since $\Delta_n(\omega)$ is the denominator of a fraction and its values can become quite small for sufficiently large n , the problem of “small denominators” arises. Therefore, in order to justify the unique solvability of the nonlocal problem for regular values of the spectral parameter ω , it is necessary to show that the quantity $\Delta_n(\omega)$ is separated from zero for sufficiently large n .

Lemma 2. *Suppose that $\gamma \in (0; 1]$, a and b are arbitrary positive real numbers, and ω is such that the product $\pi \omega a$ is a rational number. Then, for large n , there exists a positive constant M_0 such that the following estimate holds:*

$$|\Delta_n(\omega)| \geq M_0 > 0. \quad (4.3)$$

P r o o f. I. We set $\omega = p/\pi a$, $p \in \mathbb{N}$. Then we derive from (4.2) that, for all n and $a, b > 0$,

$$|\Delta_n(\omega)| \geq \left| \pm \sqrt{1 + 16n^4 \pi^2 \frac{p^2}{a^2}} - b^{\gamma-1} E_{\alpha, \gamma}(-16n^4 \pi^4 b^\alpha) \right| \geq$$

$$\geq |1 - b^{\gamma-1} E_{\alpha,\gamma}(-16 n^4 \pi^4 b^\alpha)| \geq 1 - b^{\gamma-1} E_{\alpha,\gamma}(-16 n^4 \pi^4 b^\alpha).$$

We use the following properties of the Mittag-Leffler function [9, Vol. 1, pp. 269–295].

(1) For all $\lambda > 0$, $\alpha, \gamma \in (0; 1]$, $\alpha \leq \gamma$, and $t > 0$, the function $t^{\alpha-1} E_{\alpha,\gamma}(-\lambda t^\alpha)$ is completely monotone, i.e.,

$$(-1)^n [t^{\gamma-1} E_{\alpha,\gamma}(-\lambda t^\alpha)]^{(n)} \geq 0, \quad n = 0, 1, 2, \dots \tag{4.4}$$

(2) The following estimate is true for all $\alpha \in (0; 2)$, $\gamma \in \mathbb{R}$, and $\arg z = \pi$:

$$|E_{\alpha,\gamma}(z)| \leq \frac{M}{1 + |z|}, \tag{4.5}$$

where $0 < M = \text{const}$ is independent of z .

Then, (4.4) implies that there exists a number $n_0 \in \mathbb{N}$ such that for all $n > n_0$ we have

$$1 - b^{\gamma-1} E_{\alpha,\gamma}(-16 n^4 \pi^4 b^\alpha) = M_1 > 0.$$

Consequently, $\Delta_n(\omega) \geq M_1 > 0$.

II. Now we set

$$\frac{p}{q} = 4\pi\omega a \in \mathbb{Q} \Leftrightarrow \omega = \frac{1}{4\pi a} \frac{p}{q},$$

where $p, q \in \mathbb{N}$, $(p, q) = 1$. We divide $n^2 p$ by q with a remainder: $n^2 p = sq + r$, $s \in \mathbb{N}$, $0 \leq r < q$. Then from (4.1), we obtain

$$|\Delta_n(\omega)| = \left| \sqrt{1 + \left[\frac{\pi}{a} \left(s + \frac{r}{q}\right)\right]^2} (-1)^s \sin\left(\frac{\pi r}{q} + \rho_n\right) - b^{\gamma-1} E_{\alpha,\gamma}(-16 n^4 \pi^4 b^\alpha) \right|.$$

If $r = 0$, then this case reduces to case I.

Suppose that $r > 0$. Since $\rho_n \rightarrow 0$ as $n \rightarrow +\infty$, there exists a number $n_1 > 0$ such that $\rho_n < \pi/(2q)$ for all $n > n_1$. Thus, we obtain the lower estimate

$$\begin{aligned} |\Delta_n(\omega)| &\geq \left| \sqrt{1 + \left[\frac{\pi}{a} \left(s + \frac{r}{q}\right)\right]^2} \sin\left(\frac{\pi r}{q} + \rho_n\right) - b^{\gamma-1} E_{\alpha,\gamma}(-16 n^4 \pi^4 b^\alpha) \right| \geq \\ &\geq \sqrt{1 + \left[\frac{\pi}{a} \left(s + \frac{r}{q}\right)\right]^2} \left| \sin\left(\frac{\pi r}{q} + \rho_n\right) \right| - b^{\gamma-1} E_{\alpha,\gamma}(-16 n^4 \pi^4 b^\alpha) > \\ &> \frac{\pi}{a} \left(s + \frac{r}{q}\right) \left| \sin\left(\frac{\pi(q-1)}{q} + \frac{\pi}{2q}\right) \right| - 1 = \frac{\pi}{a} \left(s + \frac{r}{q}\right) \sin \frac{\pi}{2q} - 1 = M_2 > 0 \end{aligned}$$

for

$$n_2 \geq \left[a q \left(\pi p \sin \frac{\pi}{2q} \right)^{-1} \right]^{1/2}.$$

Setting $M_0 > \max\{M_1, M_2\}$ and $n > \max\{n_0, n_1, n_2\}$, we complete the proof of the lemma. Lemma 2 is proved. \square

We call the solution of the nonlocal problem (1.1)–(1.5) for regular values of the spectral parameter ω a regular solution of the nonlocal problem. Estimates (4.3) and (4.5) imply the following lemma.

Lemma 3. *The following estimates hold for regular values of the spectral parameter ω :*

$$\begin{aligned} t^{1-\gamma} |u_0^+(t)| &\leq C_1 |\varphi_0|, & t^{1-\gamma} |u_{ni}^+(t)| &\leq C_2 |\varphi_{ni}|, \\ t^{1-\gamma} \left| D^{\alpha, \beta} u_{ni}^+(t) \right| &\leq C_3 n^4 |\varphi_{ni}|, & i = 1, 2, & t \in [0; b]; \\ |u_0^-(t)| &\leq C_4 |\varphi_0|, & |u_{ni}^-(t)| &\leq C_5 n^2 |\varphi_{ni}|, \\ \left| \frac{du_{ni}^-(t)}{dt} \right| &\leq C_6 n^4 |\varphi_{ni}|, & \left| \frac{d^2 u_{ni}^-(t)}{dt^2} \right| &\leq C_7 n^6 |\varphi_{ni}|, & i = 1, 2, & t \in [-a; 0], \end{aligned}$$

where C_k , $k = \overline{1, 7}$, are positive constants.

Since system (3.1) is complete and forms a Riesz basis in $L_2(0; 1)$, we write the solution of the nonlocal problem for regular values of the spectral parameter ω as

$$U(t, x) = \begin{cases} u_0^+(t) \vartheta_0(x) + \sum_{n=1}^{\infty} \sum_{i=1}^2 u_{ni}^+(t) \vartheta_{ni}(x), & (t, x) \in \Omega_1, \\ u_0^-(t) \vartheta_0(x) + \sum_{n=1}^{\infty} \sum_{i=1}^2 u_{ni}^-(t) \vartheta_{ni}(x), & (t, x) \in \Omega_2, \end{cases} \quad (4.6)$$

where $u_0^\pm(t)$, $u_{n1}^\pm(t)$, and $u_{n2}^\pm(t)$ are defined in (3.15) and (3.16).

Indeed, substituting function (4.6) into the mixed equation (1.1) and satisfying conditions (1.3)–(1.5), we obtain problems (3.5), (3.6), (3.8)–(3.10) with respect to the desired functions. The solutions of these problems can be represented as functions (3.15) and (3.16).

Now formally differentiating term-by-term the series (4.6) the required number of times, we obtain the series

$$D^{\alpha, \gamma} U(t, x) = \sum_{n=1}^{\infty} \sum_{i=1}^2 D^{\alpha, \gamma} u_{ni}^+(t) \vartheta_{ni}(x), \quad t > 0, \quad (4.7)$$

$$\frac{\partial^k U(t, x)}{\partial x^k} = u_0^+(t) \frac{d^k \vartheta_0(x)}{dx^k} + \sum_{n=1}^{\infty} \sum_{i=1}^2 u_{ni}^+(t) \frac{d^k \vartheta_{ni}(x)}{dx^k}, \quad k = \overline{1, 4}, \quad t > 0, \quad (4.8)$$

$$\frac{\partial^2 U(t, x)}{\partial t^2} = \sum_{n=1}^{\infty} \sum_{i=0}^2 \frac{d^2 u_{ni}^-(t)}{dt^2} \vartheta_{ni}(x), \quad t < 0, \quad (4.9)$$

$$\frac{\partial^k U(t, x)}{\partial x^k} = u_0^-(t) \frac{d^k \vartheta_0(x)}{dx^k} + \sum_{n=1}^{\infty} \sum_{i=1}^2 u_{n2}^-(t) \frac{d^k \vartheta_{n2}(x)}{dx^k}, \quad k = \overline{0, 4}, \quad t < 0. \quad (4.10)$$

By virtue of Lemma 2 and Lemma 3, we conclude that series (4.9) and (4.10) are majorized by the following sum of series:

$$\sum_{n=1}^{\infty} n^6 |\varphi_{n1}| + \sum_{n=1}^{\infty} n^6 |\varphi_{n2}|. \quad (4.11)$$

Multiplying series (4.7) and (4.8) term-by-term by $t^{1-\gamma}$, we obtain the series

$$\sum_{n=1}^{\infty} \sum_{i=1}^2 t^{1-\gamma} D^{\alpha, \gamma} u_{ni}^+(t) \vartheta_{ni}(x), \quad \sum_{n=1}^{\infty} \sum_{i=1}^2 t^{1-\gamma} u_{ni}^+(t) \frac{d^k \vartheta_{ni}(x)}{dx^k}, \quad k = \overline{0, 4}, \quad t > 0. \quad (4.12)$$

The series in (4.12) are also majorized by the series (4.11). Taking into account the fact that the function $\varphi(x)$ is sufficiently smooth and integrating by parts

$$\varphi_{ni} = \int_0^1 \varphi(x) \eta_{ni}(x) dx, \quad i = 1, 2,$$

we derive

$$\begin{aligned} \varphi_{n1} &= -\frac{1}{(2\pi n)^\gamma} \varphi_{n1}^{(\gamma)} = -\frac{1}{(2\pi n)^\gamma} \left(\varphi^{(\gamma)}(x), \vartheta_{n2}(x) \right), \\ \varphi_{n2} &= \frac{1}{(2\pi n)^\gamma} \varphi_{n2}^{(\gamma)} = \frac{1}{(2\pi n)^\gamma} \left(\varphi^{(\gamma)}(x), \vartheta_{n1}(x) \right). \end{aligned}$$

By virtue of these representations, we apply the Cauchy–Schwartz inequality and Bessel inequality to (4.11)

$$\sum_{n=1}^{\infty} n^6 |\varphi_{ni}| \leq \sum_{n=1}^{\infty} \frac{1}{n} |\varphi_{ni}^{(\gamma)}| \leq \left(\sum_{n=1}^{\infty} \frac{1}{n^2} \right)^{1/2} \left(\sum_{n=1}^{\infty} |\varphi_{ni}^{(\gamma)}|^2 \right)^{1/2} \leq C \|\varphi^{(\gamma)}(x)\|_{L_2(0,1)} < \infty, \quad i = 1, 2.$$

This estimate implies that series (4.9) and (4.10) converge absolutely and uniformly in the domains $\overline{\Omega}_1$ and $\overline{\Omega}_2$, respectively. Therefore, the function $U(t, x)$, represented by series (4.6), possesses properties (1.2) and satisfies conditions (1.3)–(1.5).

We note that $\Delta_n(\omega) = 0$ for irregular values of the spectral parameter ω and $n = k_1, \dots, k_s$, $1 \leq k_1 < k_2 < \dots < k_s$, $s \in \mathbb{N}$ ($\gamma \neq 1$). Then, the following orthogonality conditions are necessary and sufficient for the solvability of systems (3.11) and (3.12):

$$\varphi_{ni} = \int_0^1 \varphi(x) \eta_{ni} dx = 0, \quad i = 1, 2, \quad n = k_1, \dots, k_s. \tag{4.13}$$

In this case, the solutions of the nonlocal problem are representable as a sum of series

$$\begin{aligned} U(t, x) &= u_0^\pm(t) \vartheta_0(x) + \\ &+ \left[\sum_{n=1}^{k_1-1} + \sum_{n=k_1+1}^{k_2-1} + \dots + \sum_{n=k_s+1}^{\infty} \right] \sum_{i=1}^2 u_{ni}^\pm(t) \vartheta_{ni}(x) + \sum_m \sum_{i=1}^2 C_{mi} V_{mi}^\pm(t), \end{aligned} \tag{4.14}$$

where $m = k_1, \dots, k_s$, C_{mi} are arbitrary constants, and the functions $V_{mi}^\pm(t)$, $i = 1, 2$, are defined in (4.1). Note that, in the case $\gamma = 1$, we replace the function $u_0^\pm(t)$ in (4.14) with a constant C_0 ; moreover, the orthogonality condition

$$\varphi_0 = \int_0^1 \varphi(x) dx = 0 \tag{4.15}$$

is added to formula (4.13).

Thus, the following theorem is proved.

Theorem 2. *Suppose that the following conditions are fulfilled:*

$$\begin{aligned} \varphi(x) &\in C^6[0; 1], \quad \varphi^{(\gamma)}(x) \in L_2(0; 1), \quad \varphi^{(2k)}(0) = 0, \\ \varphi^{(2(k+1))}(1) &= 0, \quad k = \overline{0, 2}, \quad \varphi^{(k)}(0) = \varphi^{(k)}(1), \quad k = 1, 3, 5. \end{aligned}$$

Then the nonlocal boundary value problem is uniquely solvable for regular values of the spectral parameter ω , and this solution is represented in the form of the Fourier series (4.6) in the domain Ω .

For irregular values of the spectral parameter ω and some $n = k_1, \dots, k_s$, the nonlocal problem has an infinite number of solutions in the form of series (4.14).

For $\gamma < 1$, the solvability condition has the form (4.13). For $\gamma = 1$ in (4.14), the function $u_0^\pm(t)$ is replaced with a constant C_0 and conditions (4.13), and (4.15) are the solvability conditions.

5. Stability of solution of the nonlocal problem

For regular values of the spectral parameter ω , we consider the question of the stability of the solution of the nonlocal problem with respect to the function $\varphi(x)$ from condition (1.4). To this end, we introduce the norm in the space of continuous functions as follows:

$$\begin{aligned} \|U(t, x)\|_{C(\bar{\Omega})} &= \|t^{1-\gamma}U(t, x)\|_{C(\bar{\Omega}_1)} + \|U(t, x)\|_{C(\bar{\Omega}_2)} = \\ &= \max_{(t, x) \in \bar{\Omega}_1} |t^{1-\gamma}U(t, x)| + \max_{(t, x) \in \bar{\Omega}_1} |U(t, x)|. \end{aligned}$$

Theorem 3. *Suppose that all the conditions of Theorem 2 are fulfilled. Then the following estimate holds for the solution of the nonlocal problem with regular values of the spectral parameter ω :*

$$\|U(t, x)\|_{C(\bar{\Omega})} \leq C \|\varphi'''(x)\|_{C[0;1]}, \quad (5.1)$$

where $0 < C = \text{const}$ is independent of $\varphi(x)$ and $\|f(x)\|_{C[0;1]} = \max_{[0;1]} |f(x)|$.

P r o o f. Let (t, x) be an arbitrary point of the domain $\bar{\Omega}_2$. Then we have the representations

$$\begin{aligned} \varphi_{n1} &= -\frac{1}{\lambda_n^3} \varphi_{n1}^{(3)}, \quad \varphi_{n1}^{(3)} = \int_0^1 \varphi'''(x) \vartheta_{n2}(x) dx, \\ \varphi_{n2} &= \frac{1}{\lambda_n^3} \varphi_{n2}^{(3)}, \quad \varphi_{n2}^{(3)} = \int_0^1 \varphi'''(x) \vartheta_{n1}(x) dx. \end{aligned}$$

Applying Lemma 3 and the Cauchy–Schwarz inequality to (4.6), we obtain

$$\begin{aligned} \|U(t, x)\|_{C(\bar{\Omega}_2)} &\leq 2C_4|\varphi_0| + C_5 \sum_{n=1}^{\infty} \frac{1}{n} (|\varphi_{n1}^{(3)}| + |\varphi_{n2}^{(3)}|) \leq \\ &\leq 2C_4|\varphi_0| + C_5 \left(\sum_{n=1}^{\infty} \frac{1}{n^2} \right)^{1/2} \left(\sum_{n=1}^{\infty} (|\varphi_{n1}^{(3)}| + |\varphi_{n2}^{(3)}|)^2 \right)^{1/2}. \end{aligned}$$

It is well known that the former series converges. Applying the inequality $(|a|+|b|)^2 \leq 2(|a|^2 + |b|^2)$ and the Bessel inequality to the latter series, we obtain

$$\sum_{n=1}^{\infty} (|\varphi_{n1}^{(3)}| + |\varphi_{n2}^{(3)}|)^2 \leq 2 \sum_{n=1}^{\infty} (|\varphi_{n1}^{(3)}|^2 + |\varphi_{n2}^{(3)}|^2) \leq C_{11} \|\varphi'''(x)\|_{L_2(0;1)}^2, \quad 0 < C_{11} = \text{const}. \quad (5.2)$$

Similarly, we can find for all $(t, x) \in \bar{\Omega}_1$ that

$$\|t^{1-\gamma}U(t, x)\|_{C(\bar{\Omega}_1)} \leq C_{12} \|\varphi'''(x)\|_{L_2(0;1)}^2, \quad 0 < C_{12} = \text{const}. \quad (5.3)$$

Estimates (6.1) and (6.2) imply estimate (5.1), where $C = C_{11} + C_{12}$. If we assume that $\|\varphi'''(x)\|_{L_2(0;1)}^2 < \delta$, then the estimate $\|U(t, x)\|_{C(\bar{\Omega})} < \varepsilon$ is true for all $\varepsilon = C \cdot \delta$. The theorem is proved. \square

6. Illustrative examples

Example 1. Consider the nonlocal problem for $\gamma = 1$. Then we have $D^{\alpha,\gamma} = D^{\alpha,1} = {}_C D^\alpha$ and equation (1.1) takes the form

$$0 = \begin{cases} {}_C D^\alpha U(t, x) + \frac{\partial^4 U(t, x)}{\partial x^4}, & t > 0, \\ \frac{\partial^2 U(t, x)}{\partial t^2} + \omega^2 \frac{\partial^4 U(t, x)}{\partial x^4}, & t < 0. \end{cases} \tag{6.1}$$

Equations (6.1) is a mixed type differential equation with the Caputo operator in a positive rectangular domain. We consider it under conditions (1.3)–(1.5). From (3.14), we obtain $A_0 = \varphi_0 = 0$, i.e., we arrive at condition (4.15). The solution of this problem with regular values of the spectral parameter ω can be represented as

$$U(t, x) = \begin{cases} \sum_{n=1}^{\infty} \sum_{i=1}^2 \frac{\varphi_{ni}}{\Delta_n(\omega)} E_{\alpha,1}(-\lambda_n^4 t^\alpha) \vartheta_{ni}(x) + C_{01} x, & (t, x) \in \Omega_1, \\ \sum_{n=1}^{\infty} \sum_{i=1}^2 \frac{\varphi_{ni}}{\Delta_n(\omega)} \left(\cos \lambda_n^2 \omega t - \frac{\lambda_n^2}{\omega} \sin \lambda_n^2 \omega t \right) \vartheta_{ni}(x) + C_{02} x, & (t, x) \in \Omega_2, \end{cases}$$

where $C_{0i} = \text{const}$, $i = 1, 2$.

Example 2. Consider the nonlocal problem for $\gamma = \alpha < 1$. Then we have $D^{\alpha,\gamma} = D^{\alpha,\alpha} = {}_{RL} D^\alpha$ and equation (1.1) takes the form

$$0 = \begin{cases} {}_{RL} D^\alpha U(t, x) + \frac{\partial^4 U(t, x)}{\partial x^4}, & t > 0, \\ \frac{\partial^2 U(t, x)}{\partial t^2} + \omega^2 \frac{\partial^4 U(t, x)}{\partial x^4}, & t < 0. \end{cases} \tag{6.2}$$

Equation (6.2) is a mixed type differential equation with the Riemann–Liouville operator in a positive rectangular domain. We consider it under conditions (1.3)–(1.5). A solution of this problem with regular values of the spectral parameter ω exists and is unique. This solution has a representation coinciding with (4.6) for $\gamma = \alpha < 1$.

Example 3. Consider the case $\gamma = \alpha = 1$. Then we have $D^{\alpha,\gamma} = D^{1,1} = d/dt$ and equation (1.1) takes the form

$$0 = \begin{cases} \frac{\partial U(t, x)}{\partial t} + \frac{\partial^4 U(t, x)}{\partial x^4}, & t > 0, \\ \frac{\partial^2 U(t, x)}{\partial t^2} + \omega^2 \frac{\partial^4 U(t, x)}{\partial x^4}, & t < 0. \end{cases}$$

We obtained a mixed type differential equation of integer order, which is a particular case of equation (6.1) and, therefore, the solvability condition for this problem coincides with condition (4.15), and the solution of the nonlocal problem is represented as

$$U(t, x) = \begin{cases} \sum_{n=1}^{\infty} \sum_{i=1}^2 \frac{\varphi_{ni}}{\Delta_n(\omega)} e^{-\lambda_n^4 t} \vartheta_{ni}(x) + Ax, & (t, x) \in \Omega_1, \\ \sum_{n=1}^{\infty} \sum_{i=1}^2 \frac{\varphi_{ni}}{\Delta_n(\omega)} \left(\cos \lambda_n^2 \omega t - \frac{\lambda_n^2}{\omega} \sin \lambda_n^2 \omega t \right) \vartheta_{ni}(x) + Ax, & (t, x) \in \Omega_2, \end{cases}$$

where $A = \text{const}$.

7. Conclusion

We established a criterion for the existence and uniqueness of the regular solution of the nonlocal problem for a fourth-order differential equation of mixed type with Hilfer operator in a positive rectangular domain and with spectral parameter in a negative rectangular domain. We use the spectral method of separation of variables, which helps us to construct the solution of the nonlocal problem (1.1)–(1.5) in the form of Fourier series. Theorems on the existence and uniqueness of the problem are proved for regular values of the spectral parameter ω . We study also the case of irregular values of spectral parameter ω . Our theorem proving methods are based on expanding the regular solution using a biorthogonal set of functions. The stability of the regular solution of the nonlocal problem with respect to the data is proved.

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