



NUMERICAL MODELING OF UNSATURATED LAYERED SOIL FOR RAINFALL-INDUCED SHALLOW LANDSLIDES

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Abstract. In this paper, a pioneer study on numerical modeling of rainfall-induced shallow landslides in unsaturated layered soil using the variably saturated flow equation is presented. To model the shallow landslides, the infinite slope stability analysis coupled with the hydrological model with the consideration of the fluctuation of time-dependent pore water pressure and Gardner equation for soil water characteristic curve was developed. A linearization process for the nonlinear Richards equation to deal with groundwater flow in unsaturated layered soil is derived using the Gardner model. To solve one-dimensional flow in the unsaturated zone of layered soil profiles, flux conservation and the continuity of pressure potential at the interface between two consecutive layers are considered in the numerical discretization of the finite difference method. The validity of the proposed model is established in three numerical problems by comparing the results with the analytical and other numerical solutions. Application examples have also been conducted. Obtained results demonstrate that the fluctuation of pore water pressure in unsaturated layered soil dominates slope stability of landslides and the lowest factor of safety may occur at the interface between two consecutive layers. The findings observed in this study are a fundamental contribution to environmental protection engineering for landslides in areas with higher occurrence and vulnerability to extreme precipitation.

Keywords: Variably saturated flow equation, Layered soil, Shallow Landslides, Groundwater flow.

Introduction

Due to global climate change, severe weather phenomena such as extreme precipitation events are becoming much more frequent around the world (Pereira *et al.* 2010; Csete, Buzasi 2016). These affect the stability of slopes and have consequences on landslides. Since global warming is expected to increase the frequency and intensity of severe rainfall events, the number of people exposed to landslide risk may dramatically be increased (Gariano, Guzzetti 2016). Accordingly, it is of importance to understand and measure how climate variables and their variability affect landslide hazards. Landslides commonly observed in the foothill areas involve sliding surfaces which is typically called a shallow landslide with an indication of the high diffusion of shallow landslides phenomena in different environmental and geological contexts, all over the world. In shallow landslides, the sliding surface is located in the soil mantle or regolith typically to depths of a few centimeters to a few meters. Shallow landslides may often occur within the vadose zone under unsaturated soil conditions (Lu, Godt 2008; Yeh, Lee 2013). The assessment

of rainfall-induced shallow landslides has long been an interest of research topic in many branches of science and engineering such as soil science, hydrology, and geotechnical engineering (Montgomery, Dietrich 1994; Baum *et al.* 2002; Frattini *et al.* 2004; Tsai, Yang 2006; Tsai, Chen 2010).

In the past, theoretical models have been developed to predict slope stability assuming a steady or quasi-steady water table associated with an infinite-slope stability analysis (Das 2010). With regard to the influence of rainfall infiltration on slope stability in unsaturated zones, Iverson (2000) assumed that this occurred when soil was nearly saturated, further solving the analytical solution of the Richards equation in the simplified form of a one-dimensional linear diffusion equation, and allowing the development of an analytical model of rainfall-induced shallow landslides. Baum *et al.* (2002) expanded Iverson's work and developed a transient rainfall infiltration and grid-based regional slope-stability analysis (TRIGRS) model. Because of its practicability, the TRIGRS model became popular for assessing shallow landslides (Crosta, Frattini

2003; Keim, Skaugset 2003; Lan *et al.* 2005; D'Odorico *et al.* 2005; Liu *et al.* 2008; Baum *et al.* 2010; Bordoni *et al.* 2015; Schilirò *et al.* 2015; Alvioli, Baum 2016).

In these numerical models, the hillslope soil is typically assumed to be homogeneous. Besides, Iverson's method assumed that nearly saturated soil was a simplified form of the linear diffusion equation of the Richards equation, in which the influence of matric suction of soil and the soil water characteristic curve (SWCC) cannot be comprehensively considered during analysis. A complete theory of groundwater flow when rainfall infiltrates unsaturated zones can be described using either the variably saturated flow equation or the generalized Richards equation. The Richards equation is highly nonlinear and cannot directly provide an analytical solution. The soil water content is a major factor that influences the nonlinear physical relationship of the hydraulic conductivity in unsaturated zones (Van Genuchten 1980; Fredlund *et al.* 1994).

Since the appearance of layered soil is much more common than homogeneous soil in nature, the hydrological process in heterogeneous porous media has drawn much attention and been studied (Fok 1970; Aylor, Parlange 1973; Hachum, Alfaro 1980; Samani *et al.* 1989), using analytical or numerical methods (Hanks Bowers 1962; Whisler, Klute 1966; Romano *et al.* 1988; Moldrup *et al.* 1989; Srivastava, Yeh 1991; Corradini *et al.* 2000; Ku, Tsai 2013). Because the numerical modeling of rainfall-induced shallow landslides in unsaturated layered soil using the variably saturated flow equation has hardly been reported, the study proposed a pioneer work on numerical modeling of rainfall-induced shallow landslides in unsaturated layered soil.

To model the shallow landslides, the infinite slope stability analysis coupled with the hydrological model with the consideration of the fluctuation of time-dependent pore water pressure and the SWCC proposed by Gardner was developed. A linearization process for the nonlinear Richards equation to deal with groundwater flow in unsaturated layered soil is derived using the Gardner model. To solve one-dimensional flow in the unsaturated zone of layered soil profiles, flux conservation and the continuity of pressure potential at the interface between two consecutive layers are considered in the numerical discretization of the finite difference method. In addition, model verification and comparison were performed, and a numerical model of rainfall-induced shallow landslides in unsaturated zones was subsequently developed. The formulation of the proposed method is described as follows.

1. Methods

1.1. Formulation of the variably saturated flow equation

Unsaturated soil is composed through a three-phase mixing process that involves a gas phase (air), a liquid phase

(pore water) and a solid phase (soil particles). Soil layers above the water table are termed as the unsaturated zone. In geotechnical engineering, pore water pressure fluctuation reduces soil suction which may cause soil failure. Generally, continuous rainfall reduces soil suction, affecting the stability of slope. The complete theory of groundwater flow and rainfall infiltration of the unsaturated zone can be expressed using the Richards equation:

$$K\nabla^2 H = [S_s S_w + C(h)] \frac{\partial H}{\partial t}, \quad (1)$$

where H represents the total groundwater head, K represents hydraulic conductivity, S_s represents the specific storage, S_w represents saturation, $C(h)$ represents specific capacity, and t represents time. If porous media are heterogeneous and anisotropic, Eq. (1) becomes:

$$\nabla K \nabla H = [S_s S_w + C(h)] \frac{\partial H}{\partial t}. \quad (2)$$

This is known as the generalized Richards equation or variably saturated flow equation, which can be used to describe the groundwater flow in saturated and unsaturated zones simultaneously (Tóthová *et al.* 2007). Eq. (2) can be rewritten as follows for a complete three-dimensional problem:

$$S_s S_w \frac{\partial H}{\partial t} + C(h) \frac{\partial H}{\partial t} = \frac{\partial}{\partial x} \left(K_x(h) \frac{\partial H}{\partial x} \right) + \frac{\partial}{\partial y} \left(K_y(h) \frac{\partial H}{\partial y} \right) + \frac{\partial}{\partial z} \left(K_z(h) \frac{\partial H}{\partial z} \right), \quad (3)$$

where $K_x(h)$, $K_y(h)$, and $K_z(h)$ represent hydraulic conductivity in the x , y and z axis, which are all functions of the pore water pressure head h when unsaturated. To convert Eq. (3) into an equation that considers infiltration, we assume that the x and y planes represent slope vertical sections, the z axis is perpendicular to the xy plan, and the total water head H can be expressed using the pressure head h and elevation head E as follows:

$$H = h + E. \quad (4)$$

When the slope infiltration problem is considered, the elevation head E varies with the slope angle α , as shown in Figure 1. The elevation head can be presented as follows:

$$E = x \sin(\alpha) + z \cos(\alpha). \quad (5)$$

The governing equation can be obtained by substituting Eqs (4) and (5) into Eq. (3):

$$(S_s S_w + C(h)) \frac{\partial h}{\partial t} = \frac{\partial}{\partial x} \left(K_x(h) \left(\frac{\partial h}{\partial x} - \sin \alpha \right) \right) + \frac{\partial}{\partial y} \left(K_y(h) \frac{\partial h}{\partial y} \right) + \frac{\partial}{\partial z} \left(K_z(h) \left(\frac{\partial h}{\partial z} + \cos \alpha \right) \right). \quad (6)$$

Thus, the majority of the variably saturated flow problem was simplified using Eq. (6). Considering a

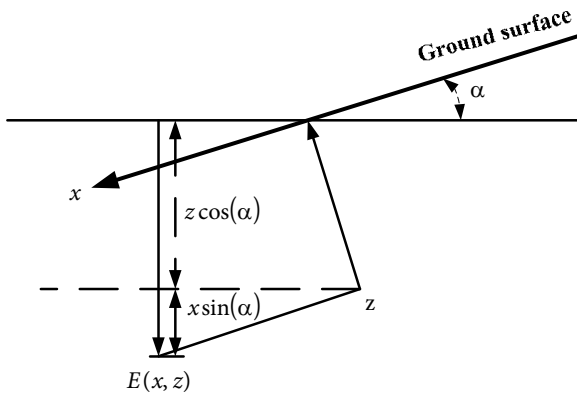


Fig. 1. The definition of the coordinate system in this study

one-dimensional problem for example, rainfall duration was far shorter than the transmission time of pore water pressure in the x and y axis. This consideration can be realistic for slopes, especially steep slopes usually covered by a thin soil mantle. Thus, only vertical water flow variations were considered, which can be presented as a one-dimensional vertical infiltration equation:

$$\left(S_s S_w + C(h)\right) \frac{\partial h}{\partial t} = \frac{\partial}{\partial z} \left(K_z(h) \left(\frac{\partial h}{\partial z} + \cos \alpha \right) \right). \quad (7)$$

When soil is saturated, $S_w = 1$ and $C(h) = 1$. Therefore, the saturated flow equation is a commonly used governing equation for groundwater in the saturated zone, where:

$$S_s \frac{\partial h}{\partial t} = \frac{\partial}{\partial z} \left(K_z \left(\frac{\partial h}{\partial z} + \cos \alpha \right) \right). \quad (8)$$

When soil is not saturated, the groundwater flow equation in unsaturated zone can be expressed as:

$$C(h) \frac{\partial h}{\partial t} = \frac{\partial}{\partial z} \left(K_z(h) \left(\frac{\partial h}{\partial z} + \cos \alpha \right) \right). \quad (9)$$

When the soil layer is not saturated, Eq. (9) can be used to calculate fluctuations in pore water pressure. When the soil layer is gradually saturated by rainfall, Eq. (8) is used for analysis because Eq. (7) primarily describes the pore water pressure changes in saturated and unsaturated zones. This study adopted Eq. (7) as the core governing equation of rainfall-induced pore water pressure fluctuations in unsaturated zones. In 1958, Gardner (Gardner 1958) proposed a simple-one parameter model as follows:

$$S_e = e^{\alpha_g h}, \quad (10)$$

where α_g is the soil pore-size distribution parameter which is related to the pore size distribution of soil. S_e is the effective saturation defined by normalizing volumetric water content with its saturated and residual values as:

$$S_e = \frac{(\theta - \theta_r)}{(\theta_s - \theta_r)}, \quad (11)$$

where θ represents the volumetric water content function,

θ_r represents the residual water content and θ_s represents the saturated water content. Substituting Eq. (11) into Eq. (10), we have

$$\theta = \theta_r + (\theta_s - \theta_r) e^{\alpha_g h}. \quad (12)$$

It is common to normalize the hydraulic conductivity of unsaturated soil with respect to their maximum value. The normalized value, referred to as the relative hydraulic conductivity, K_r , can be expressed as:

$$K_r(h) = \frac{K_z(h)}{K_s}, \quad (13)$$

where K_s is the saturated hydraulic conductivity. $K_r(h)$ is the relative hydraulic conductivity and it is a function of the pressure head which makes the Richards equation highly nonlinear. Rather than using the van Genuchten expression for S_e , a simpler version is used as follows:

$$S_e = K_r. \quad (14)$$

Therefore, $K_r = e^{\alpha_g h}$. Using the Gardner model, we can derive a linearized Richards equation. First of all, we define a new parameter h^* which can be expressed as

$$h^* = e^{\alpha_g h} - \chi, \quad (15)$$

where h^* is the pressure head of the linearized Richards equation and χ is a constant which is a key transformation allowing the solution of the linearized Richards equation using the transform methods (Duffy 1994; Tracy 2006, 2011; Liu *et al.* 2015). The parameter χ has no physical meaning and can be expressed as $\chi = e^{\alpha_g h_d}$. h_d is the pressure head when the soil is dry. Taking the derivative of Eq. (15) with respect to z , we obtain

$$\frac{\partial h^*}{\partial z} = \alpha_g e^{\alpha_g h} \frac{\partial h}{\partial z}. \quad (16)$$

Rearranging Eq. (16), we have

$$\frac{\partial h}{\partial z} = \frac{1}{\alpha_g} e^{-\alpha_g h} \frac{\partial h^*}{\partial z}. \quad (17)$$

Again, substituting Eq. (15) into Eq. (17), we have

$$K_r \frac{\partial h}{\partial z} = e^{\alpha_g h} \left(\frac{1}{\alpha_g} e^{-\alpha_g h} \frac{\partial h^*}{\partial z} \right). \quad (18)$$

Eq. (18) can be written as

$$K_r \frac{\partial h}{\partial z} = \frac{1}{\alpha_g} \frac{\partial h^*}{\partial z}. \quad (19)$$

Taking the derivative of $K_r = e^{\alpha_g h}$ with respect to z , we have

$$\frac{\partial K_r}{\partial z} = \alpha_g e^{\alpha_g h} \frac{\partial h}{\partial z}. \quad (20)$$

Substituting Eq. (17) into Eq. (20), we have

$$\frac{\partial K_r}{\partial z} = \frac{\partial h^*}{\partial z}. \quad (21)$$

Recall Eq. (12), we have

$$\frac{\partial \theta}{\partial t} = (\theta_s - \theta_r) \frac{\partial S_e}{\partial t}. \quad (22)$$

Because $S_e = K_r$, $K_r = e^{\alpha_g h}$, and $h^* = e^{\alpha_g h} - \chi$, we can rewrite the above equation as

$$\frac{\partial \theta}{\partial t} = (\theta_s - \theta_r) \frac{\partial h^*}{\partial t}. \quad (23)$$

Finally, we substitute Eqs. (17), (12), and (23) to Eq. (19). The linearized Richards equation can be obtained as

$$\frac{\partial^2 h^*}{\partial z^2} + \alpha_g \frac{\partial h^*}{\partial z} = \frac{\alpha_g (\theta_s - \theta_r)}{K_s} \cos \alpha \frac{\partial h^*}{\partial t}. \quad (24)$$

For simplicity, the equation can be expressed as

$$K_{\alpha\theta} \frac{\partial^2 h^*}{\partial z^2} + K_\theta \cos \alpha \frac{\partial h^*}{\partial z} = \frac{\partial h^*}{\partial t}, \quad (25)$$

where

$$K_\theta = \frac{K_s}{(\theta_s - \theta_r)} \text{ and } K_{\alpha\theta} = \frac{K_\theta}{\alpha_g}. \quad (26)$$

1.2. The finite difference method

The finite difference method (FDM) is adopted for the numerical discretization. The reason for choosing the FDM is that the FDM can be easily used to integrate the formulation for the discontinuities (or soil interface) which characterize the layers for unsaturated layered soil. In addition, several studies of numerical modeling of shallow landslides have been carried out using the FDM such as Yoshimatsu, Abe (2006), Wang *et al.* (2006), Singh *et al.* (2010), Sarkar *et al.* (2012). The main advantage of the FDM is that it is probably the most straight forward numerical method to formulate complicated partial differential equations such as the Richards equation. In this study, we need to derive the sophisticated mathematical formulation for the unsaturated layered soil. Accordingly, we adopted the FDM as the numerical method. The linearized Richards equation based on the Gardner model is depicted in Eq. (25). To apply the FDM, we have to consider the discretization in spatial and temporal domains. For the spatial discretization, we divide the domain into sections, each of length Δz along the z axis and approximated the first and second derivatives in the linearized Richards equation for each grid point by the central difference formulas. For the time discretization, the implicit scheme is adopted. Accordingly, we obtain the finite difference equation as:

$$K_{\alpha\theta} \left(\frac{h_{i+1}^* - 2h_i^* + h_{i-1}^*}{\Delta z^2} \right) + K_\theta \cos \alpha \left(\frac{h_{i+1}^* - h_{i-1}^*}{2\Delta z} \right) = \left(\frac{h_i^* - h_i^{*n-1}}{\Delta t} \right), \quad (27)$$

where i is the nodal point and Δz is the interval. Δt is the time interval. K_s , α_g , θ_s and θ_r are the same definition as described in the previous section. Considering the steady-state situation, we obtain the finite difference equation of the linearized Richards equation as:

$$K_{\alpha\theta} \left(\frac{h_{i+1}^* - 2h_i^* + h_{i-1}^*}{\Delta z^2} \right) + K_\theta \cos \alpha \left(\frac{h_{i+1}^* - h_{i-1}^*}{2\Delta z} \right) = 0, \quad (28)$$

Eq. (28) can be used to deal with homogenous soil. For unsaturated layered soil, it is necessary to account for the discontinuities of the parameters which characterize the layers. This gives rise to the so-called interface problem for the linearized Richards equation. For the interface between layers, the conservation of water flux and the Darcy's law have to be considered, as shown in Figure 2.

The finite difference equation of the first term, $K_{\alpha\theta} \frac{\partial^2 h^*}{\partial z^2}$ in Eq. (25), can be derived as follows:

$$K_{\alpha\theta} \frac{\partial^2 h^*}{\partial z^2} = \frac{2 \left(\frac{K_{\alpha\theta A}}{K_{\alpha\theta A} + K_{\alpha\theta B}} \right) h_{i-1}^* - 2h_i^* + 2 \left(\frac{K_{\alpha\theta B}}{K_{\alpha\theta A} + K_{\alpha\theta B}} \right) h_{i+1}^*}{\Delta z^2}, \quad (29)$$

where $K_{\alpha\theta A}$ and $K_{\alpha\theta B}$ represent the parameter, $K_{\alpha\theta}$, in layer A and layer B respectively. The finite difference equation of the second term, $K_\theta \frac{\partial h^*}{\partial z}$ in Eq. (25), can be derived as follows:

$$K_\theta \frac{\partial h^*}{\partial z} = K_\theta \frac{\left[\frac{K_{\theta A}}{K_{\theta B}} (h_i^* - h_{i-1}^*) + \frac{K_{\theta B}}{K_{\theta A}} (h_{i+1}^* - h_i^*) \right]}{2\Delta z}, \quad (30)$$

where $K_{\theta A}$ and $K_{\theta B}$ represent the parameter, K_θ , in layer A and layer B respectively. Substituting Eqs. (29) and (30) into Eq. (27) we obtain:

$$K_{\alpha\theta} \frac{\left[2\alpha_A h_{i-1}^{*n} - 2h_i^{*n} + 2\alpha_B h_{i+1}^{*n} \right]}{\Delta z^2} + K_\theta \cos \alpha \frac{\left[\beta_A (h_i^{*n} - h_{i-1}^{*n}) + \beta_B (h_{i+1}^{*n} - h_i^{*n}) \right]}{2\Delta z} = \left(\frac{h_i^{*n} - h_i^{*n-1}}{\Delta t} \right), \quad (31)$$

where $\alpha_A = \frac{K_{\alpha\theta A}}{K_{\alpha\theta A} + K_{\alpha\theta B}}$, $\alpha_B = \frac{K_{\alpha\theta B}}{K_{\alpha\theta A} + K_{\alpha\theta B}}$,

$$\beta_A = \frac{K_{\theta A}}{K_{\theta B}} \text{ and } \beta_B = \frac{K_{\theta B}}{K_{\theta A}}.$$

In Eqs. (28) and (31), $K_{\alpha\theta}$ and K_{θ} at the interface have to be determined. In a stratified soil deposit where the hydraulic conductivity for flow in a given direction changes from layer to layer, an equivalent hydraulic conductivity can be obtained by the following equation:

$$K_{(eqv)} = \frac{H}{\left(\frac{H_1}{K_1}\right) + \left(\frac{H_2}{K_2}\right) + \dots + \left(\frac{H_n}{K_n}\right)}, \quad (32)$$

where $K_{(eqv)}$ is the equivalent hydraulic conductivity, H is the height of the soil, H_i is the thickness of each soil layer, and K_i is the hydraulic conductivity of each soil layer. i is the number of the soil layer. Because we considered only two consecutive layers at the interface in the finite difference discretization. Using Eq. (32), $K_{\alpha\theta}$ and K_{θ} at the interface can be obtained as:

$$K_{\alpha\theta(i)} = \frac{2K_{\alpha\theta A}K_{\alpha\theta B}}{K_{\alpha\theta A} + K_{\alpha\theta B}} \text{ and } K_{\theta(i)} = \frac{2K_{\theta A}K_{\theta B}}{K_{\theta A} + K_{\theta B}}. \quad (33)$$

Accordingly, the finite difference equations of the linearized Richards equation for steady-state and transient conditions can be expressed as:

$$K_{\alpha\theta(i)} \left(\frac{h_{i+1}^* - 2h_i^* + h_{i-1}^*}{\Delta z^2} \right) + K_{\theta(i)} \cos \alpha \left(\frac{h_{i+1}^* - h_{i-1}^*}{2\Delta z} \right) = 0; \quad (34)$$

$$K_{\alpha\theta(i)} \frac{[2\alpha \cdot h_{i-1}^{*n} - 2h_i^{*n} + 2\alpha \cdot h_{i+1}^{*n}]}{\Delta z^2} + K_{\theta(i)} \cos \alpha \frac{[\beta_A (h_i^{*n} - h_{i-1}^{*n}) + \beta_B (h_{i+1}^{*n} - h_i^{*n})]}{2\Delta z} = \left(\frac{h_i^{*n} - h_i^{*(n-1)}}{\Delta t} \right). \quad (35)$$

Furthermore, Eqs (34) and (35) can be expressed as the matrix form and can be written as

$$\mathbf{Ax} = \mathbf{b}, \quad (36)$$

where \mathbf{A} is a $n \times n$ square matrix, \mathbf{x} is a $n \times 1$ vector in which $\mathbf{x} = [h_1^* \dots h_n^*]^T$, and \mathbf{b} is a $n \times 1$ vector. n is the number of unknowns to be solved. Eq. (36) can then be simply solved by $\mathbf{x} = \mathbf{A}^{-1}\mathbf{b}$.

1.3. Slope stability theory of unsaturated soil

The failure mechanism of infinite slopes primarily involves colluvium, weathered rock formations, or shallow failure and planar slides in bedrock alternations located at thin depth below the ground level. The sliding soil thickness is far less than the slope height. In most studies, infinite slopes are used to analyze shallow landslides. The factor of safety (FS) for the infinite slope theory can be expressed as follows:

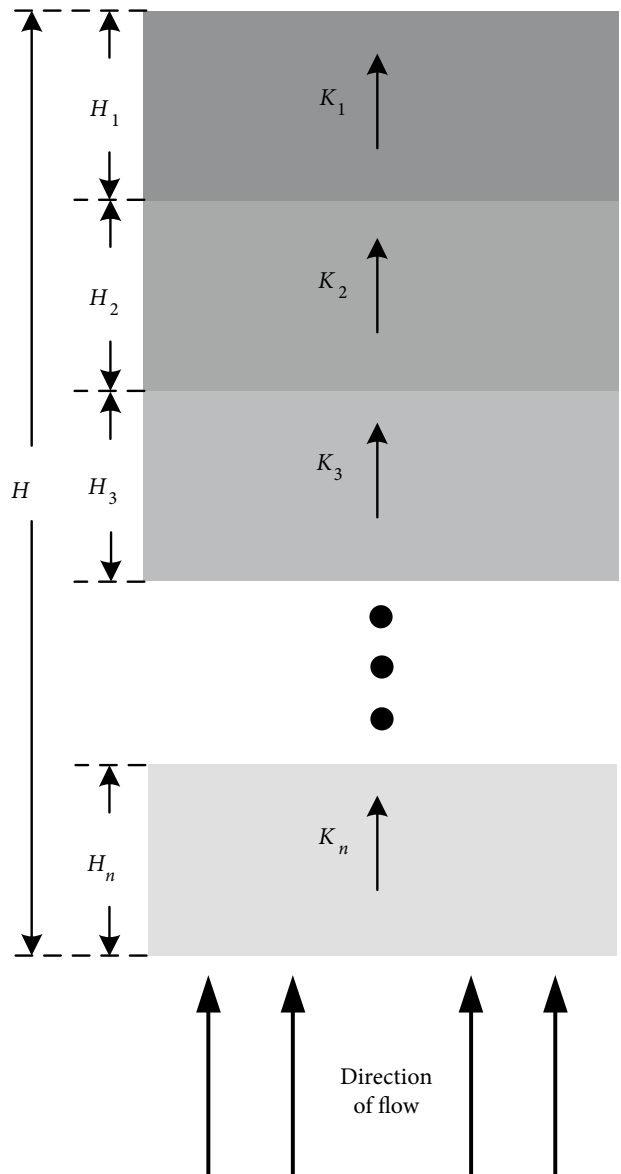


Fig. 2. The configuration of water flow through layered soils in the FDM formulation

$$FS = \frac{c' + (z\gamma_t \cos^2 \alpha - h\gamma_w \cos^2 \alpha) \tan \phi'}{z\gamma_t \cos \alpha \sin \alpha}, \quad (37)$$

where c' represents the effective cohesiveness, ϕ' represents the effective friction angle, z represents the soil thickness, γ_t represents the unit weight of soil and γ_w represents the unit weight of water. Fredlund adopted the perspectives of continuum mechanics and demonstrated that the shear strength equation can be applied to unsaturated soil (Fredlund, Morgenstern 1977). The reason of choosing Fredlund's approach for modeling the FS in unsaturated condition is that its description of the stress state of the unsaturated soil examine the context of multiphase continuum mechanics. Unlike saturated condition, which are two-phase mixing process involved a solid phase and a gas phase or a liquid phase, unsaturated soil is composed through a three-phase mixing process that comprises of

a gas phase, a liquid phase and a solid phase. Therefore, the corresponding pressure of a gas and a liquid phase in unsaturated condition have a direct impact on physical behavior of the shear strength (Lu, Likos 2004). Vanapalli *et al.* proposed the shear strength equation based on the SWCC, which is suitable for unsaturated zones (Vanapalli *et al.* 1996):

$$\tau_f = c' + \sigma \tan \phi' - h\gamma_w \left(\frac{\theta - \theta_r}{\theta_s - \theta_r} \right) \tan \phi', \quad (38)$$

where σ represents the total stress, τ_f represents the shear strength. The shear strength of unsaturated soil was developed to describe the physical consideration to the air-water interphase and these stress state variables (Fredlund *et al.* 1978; Vanapalli *et al.* 1996). The revised shear strength equation was substituted into the original infinite slope equation to obtain the following FS equation for analyzing slope stability in shallow unsaturated zones:

$$FS = \frac{c'}{z\gamma_t \cos \alpha \sin \alpha} + \frac{\tan \phi'}{\tan \alpha} - \frac{h\gamma_w \left(\frac{\theta - \theta_r}{\theta_s - \theta_r} \right) \tan \phi'}{z\gamma_t \cos \alpha \sin \alpha}. \quad (39)$$

2. Verification examples

2.1. Verification example 1

The first example is a one-dimensional transient unsaturated groundwater flow problem in a homogenous soil. A column of soil is initially dry until water begins to infiltrate the soil. A pool of water at the ground surface is then maintained holding the pressure head to zero. The thickness of the soil L is 10 m. The initial condition is described as follows:

$$h(z, t = 0) = -10^5. \quad (40)$$

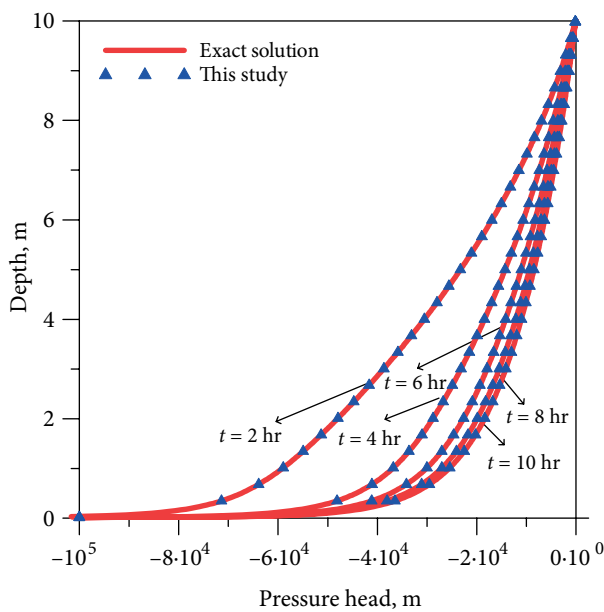


Fig. 3. Result comparison of verification example 1 with the exact solution

The boundary conditions of the top and bottom boundaries are as follows:

$$h(z = 0, t) = -10^5, \quad (41)$$

$$h(z = L, t) = 0. \quad (42)$$

Using Eq. (15), the suction head can be converted to h^* as follows:

$$h^*(z = 0, t) = 0, \quad (43)$$

$$h^*(z = L, t) = 1. \quad (44)$$

In this example, the soil is assumed to be silt which the α_g of the Gardner model is 10^{-4} . The saturated hydraulic conductivity (K_s), saturated water content (θ_s), and residual water content (θ_r) for this example are 9×10^{-5} m/s, 0.50, and 0.11, respectively, as shown in Table 1. For one-dimensional transient unsaturated groundwater flow in a homogenous soil, Tracy proposed the analytical transient solution as follows (Tracy 2011):

$$h^*(z, t) = (1 - \chi) e^{\frac{\alpha_g}{2}(L-z)} \times \left[\frac{\sinh\left(\frac{\alpha_g}{2}z\right)}{\sinh\left(\frac{\alpha_g}{2}L\right)} + \frac{2}{Lc} \sum_{k=1}^{\infty} (-1)^k \frac{\lambda_k}{\mu_k} \sin \lambda_k z e^{-\mu_k t} \right], \quad (45)$$

where $\lambda_k = k\pi/L$, $c = \alpha_g(\theta_s - \theta_r)/K_s$, $\mu_k = (\alpha_g^2/4 + \lambda_k^2)/c$, t is time. L is the thickness of the soil. Figure 3 demonstrates that the numerical results agreed well with the analytical solution for the one-dimensional transient unsaturated groundwater flow in a homogenous soil. It is found that the mean absolute error (MAE) was 1.18×10^{-4} m. The above verification example demonstrates that the proposed numerical scheme for the Richards equation to solve one-dimensional flow processes in the unsaturated zone of soil is validated.

Table 1. Parameters used in the verification example 1

saturated hydraulic conductivity (m/s)	9×10^{-5}
α_g of the Gardner model	10^{-4}
saturated water content	0.50
residual water content	0.11

2.2. Verification example 2

The second example under investigation is also a one-dimensional, transient, unsaturated groundwater flow problem in homogeneous soil. A soil column is initially dry until water begins to infiltrate the soil. A pool of water at the ground surface is maintained, holding the infiltration rate to zero. This is known as one-dimensional Green-Ampt problem (Green, Ampt 1911). The

thickness of the soil L is 10 m. The initial condition is described as follows:

$$h(z, t=0) = -100. \tag{46}$$

The boundary conditions of the top and bottom boundaries are as follows.

$$h(z=0, t) = -100, \tag{47}$$

$$h(z=L, t) = 0. \tag{48}$$

The soil was assumed to be sand, which has saturated hydraulic conductivity (K_s) of 3×10^{-3} m/s. The saturated water content (θ_s) and residual water content (θ_r) are 0.50 and 0.11, respectively. In the Gardner model, there is only one fitting parameter, α_g . The SWCC obtained from

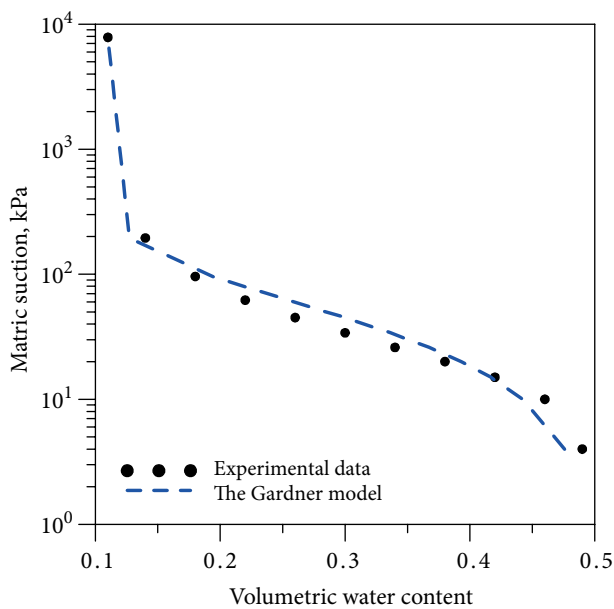


Fig. 4. The fitting curve of the SWCC using the Gardner model

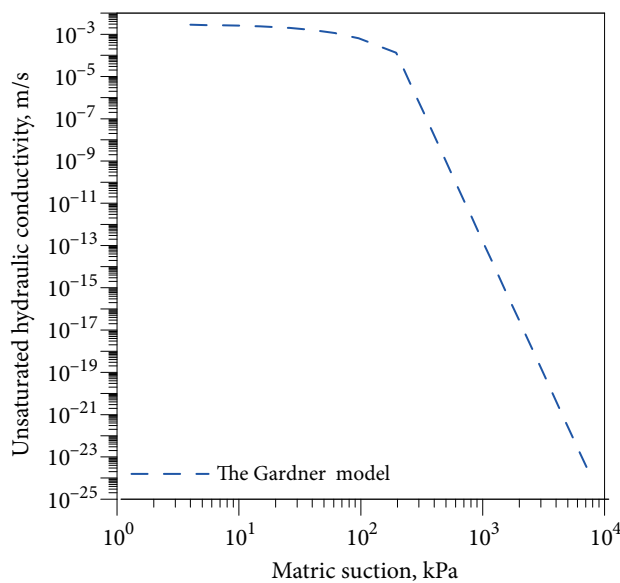


Fig. 5. The relationship between matric suction and unsaturated hydraulic conductivity for the Gardner model

the laboratory is shown in Figure 4. Using the Gardner model, α_g can be obtained by fitting the SWCC with the Gardner model as shown in Figure 4. It is found that the $\alpha_g = 1.6 \times 10^{-2}$ and the MAE was 1.78×10^{-2} .

Using the Gardner model, the nonlinear relationship between the matric suction and unsaturated hydraulic conductivity can be obtained as shown in Figure 5. In this example, the total simulation time is one hour. Figure 6 depicted the comparison of the solution with the numerical solution using the Skeel and Berzins' method (Skeel, Berzins 1990). It is found that the computed results agreed well with those obtained from the Skeel and Berzins' method and the MAE was 2.74×10^{-4} m.

2.3. Verification example 3

The governing equation for unsaturated flow in response to infiltration at ground surface is known as the Richards equation. For one-dimension, it can be written as Eq. (9). Iverson assumed that the Richards equation as a linear diffusion equation and derived the analytical solution (Iverson 2000). Baum *et al.* revised the analytical solution and implemented in TRIGRS (Baum *et al.* 2002). The analytical solution proposed by Baum *et al.* can be written as follows:

$$h(Z, t) = [Z - d] \beta + 2 \sum_{n=1}^N \frac{I_n Z}{K_Z} \times \left[H(t - t_n) \left[D_1(t - t_n) \right]^{\frac{1}{2}} \operatorname{ierfc} \left[\frac{Z}{2 \left[D_1(t - t_n) \right]^{\frac{1}{2}}} \right] \right] - 2 \sum_{n=1}^N \frac{I_n Z}{K_Z} \times \left[H(t - t_{n+1}) \left[D_1(t - t_{n+1}) \right]^{\frac{1}{2}} \operatorname{ierfc} \left[\frac{Z}{2 \left[D_1(t - t_{n+1}) \right]^{\frac{1}{2}}} \right] \right], \tag{49}$$

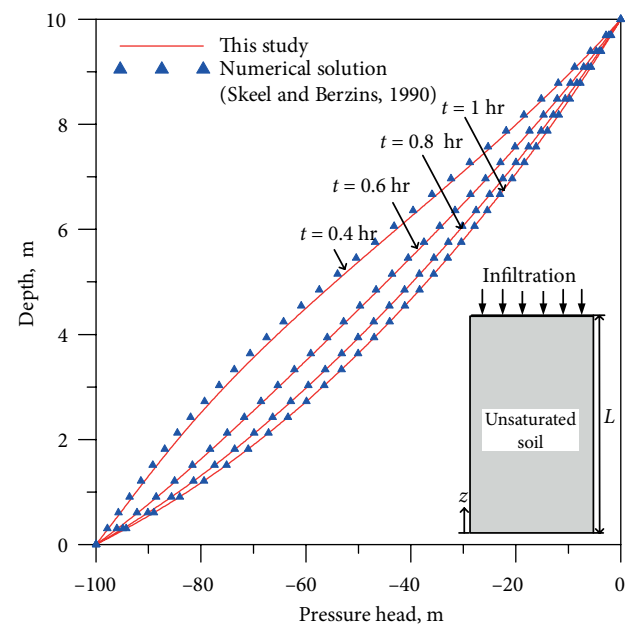


Fig. 6. Result comparison of verification example 2 with the numerical solution

where $Z = z / \cos \alpha$, d is the depth of steady-state water table measured in z direction, $\beta = \lambda \cos \alpha$ and $\lambda = \cos \delta - [I_z / K_z]_{LT}$, where I_z is the long term rainfall flux at ground surface. K_z is the hydraulic conductivity, I_{nz} is the surface flux for the n th time interval, n is the total number of time intervals, D_1 is related to the hydraulic diffusivity D_0 , $H(t - t_{n+1})$ is the Heaviside step function. The $ierfc$ is the complementary error function and can be expressed as follows:

$$ierfc(\eta) = \frac{1}{\sqrt{\pi}} \exp(-\eta^2) - \eta erfc(\eta). \quad (50)$$

We conducted a comparison example with the consideration of unsaturated flow in response to infiltration at ground surface. The input parameters were listed in Table 2. The numerical results of this example were compared to those obtained from the analytical solution of equation. In a scenario set in this study, a soil layer thickness is assumed to be 100 cm and a water table at the junction between the soil base and rock bed. Figure 7 and

Figure 8 depict that the results from the proposed model and the analytical solution. It is interested that the analytical solution from the simplified linear diffusion equation cannot obtain the appropriate result when the value of the hydraulic diffusivity is large as revealed in Figure 8. The

Table 2. Parameters used in the verification example 3

	This study	Baum et al. (2002)
Slope angle (degree)		30
Thickness (m)		1
Infiltration rate (m/hr)		0.2
Infiltration time (hr)		5
Saturated hydraulic conductivity (m/hr)		2.5
Other parameters	$\theta_s = 0.43$ $\theta_r = 0.08$ $\alpha_g = 0.2$	$D_0 = (10 \sim 100) \times K_s$ (m/hr)

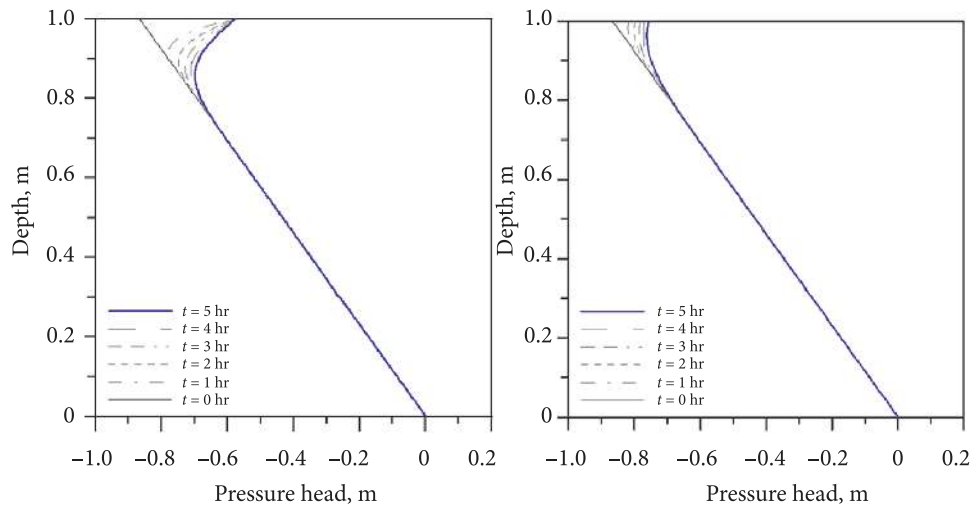


Fig. 7. Comparison of the results for $D_0 = 10 \cdot K_s$ (left: this study, right: TRIGRS model)

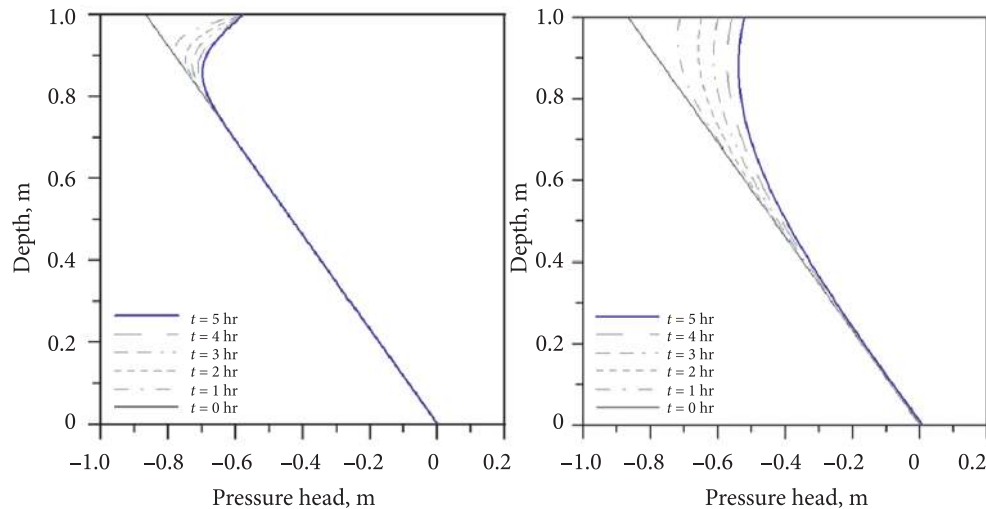


Fig. 8. Comparison of the results for $D_0 = 100 \cdot K_s$ (left: this study, right: TRIGRS model)

results from this example demonstrate that the analytical solution of equation can only be applied on the problems within certain condition and should be used with caution.

3. Application examples

3.1. Application example 1

The first application example under investigation is a one-dimensional, transient, unsaturated groundwater flow problem in a homogeneous soil. The thickness of the soil L is 150 cm and the slope angle is 31 degrees. Since the Richards equation is an initial value problem, an initial condition must be provided to solve the equation. The initial pressure head is given as follows:

$$h(z, t = 0) = (-z - z_w) \cos \alpha, \quad (51)$$

where z_w represents the height of water table. The pore pressure head distribution in unsaturated zones is congruent with the water table function. In this example, we considered $z_w = 50$ cm. The total simulation time is five hours. Considering that a pool of water at the ground surface is maintained holding the infiltration rate, the boundary condition can be defined according to Darcy's law $q = -K_s \nabla H$, and assumes that infiltration rate facilitates complete infiltration and $q = R_f$ as follows:

$$R_f = -K_s \frac{\partial H}{\partial z}. \quad (52)$$

The definition of K_s and H are the same as defined in the previous section. Using Eqs. (4) and (5), the above equation can be written as:

$$R_f = -K_s \left(\frac{\partial h}{\partial z} + \cos \alpha \right). \quad (53)$$

The above equation can then be rewritten as the following equation which is also known as the Neumann boundary condition:

$$\frac{\partial h}{\partial z} = -\cos \alpha - \frac{R_f}{K_s}. \quad (54)$$

If the infiltration rate is greater than saturated hydraulic conductivity, the maximal soil infiltration equals the saturated hydraulic conductivity. The bottom boundary may be assumed to be a no-flow boundary, and is inferred as follows:

$$0 = -K_s \frac{\partial H}{\partial z}. \quad (55)$$

The above equation can be written as:

$$0 = -K_s \left(\frac{\partial h}{\partial z} + \cos \alpha \right). \quad (56)$$

Then, we have

$$\frac{\partial h}{\partial z} = -\cos \alpha. \quad (57)$$

To solve the Richards equation, the boundary conditions must be considered. We adopted $R_f = 5 \times 10^{-5}$ cm/hr.

The soil was assumed to be sand, which has saturated hydraulic conductivity (K_s) of 9×10^{-2} cm/hr. The saturated water content (θ_s) and residual water content (θ_r) are 0.43 and 0.08, respectively. The α_g of the Gardner model is 1×10^{-2} . In this example, the total simulation time is five hours. The effective cohesiveness, effective friction angle and unit weight of soil are 4.6 kPa, 30 degrees and 21.5 kN/m^3 , respectively. Figure 9 and Figure 10 indicate the computed results of pressure head and FS for application example 1.

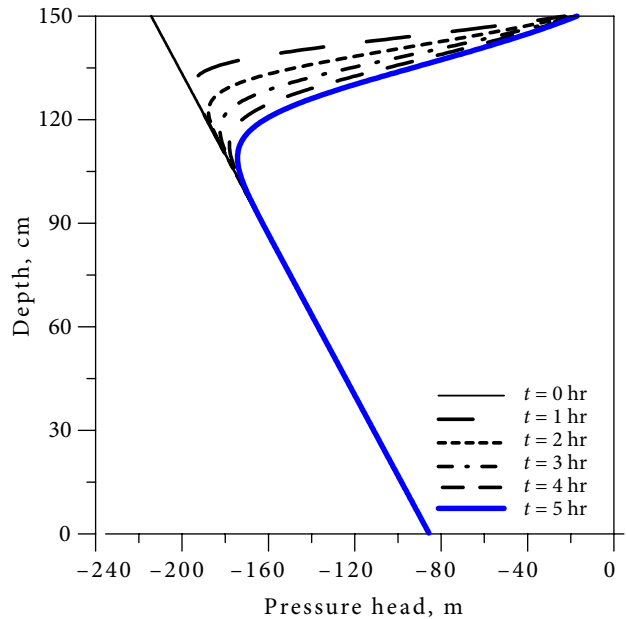


Fig. 9. Computed pressure head of the soil profile for application example 1

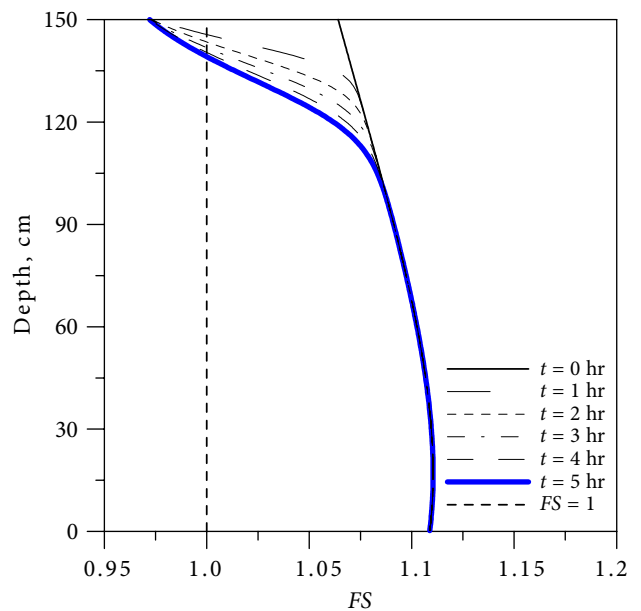


Fig. 10. Computed FS of the soil profile for application example 1

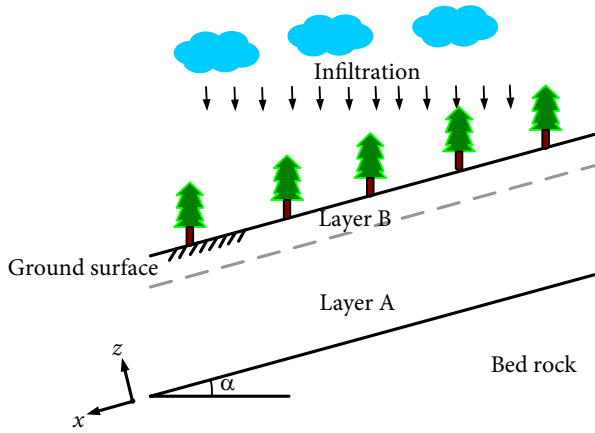


Fig. 11. Schematic illustration of application example 2

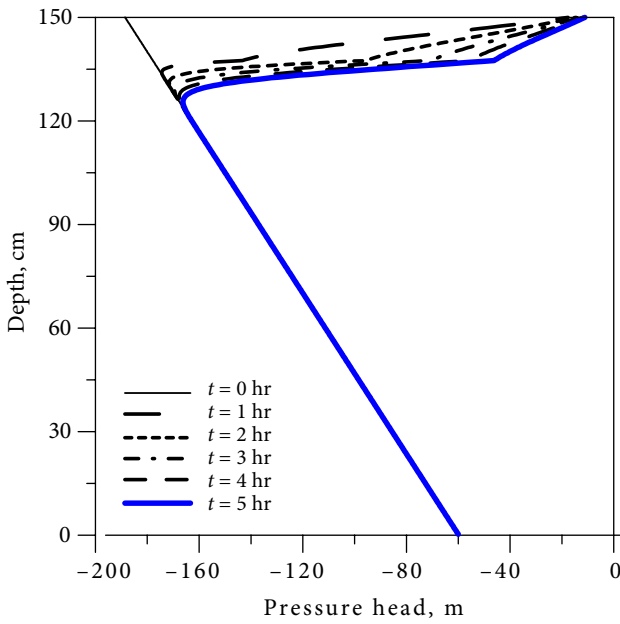


Fig. 12. Computed pressure head of the soil profile for application example 2

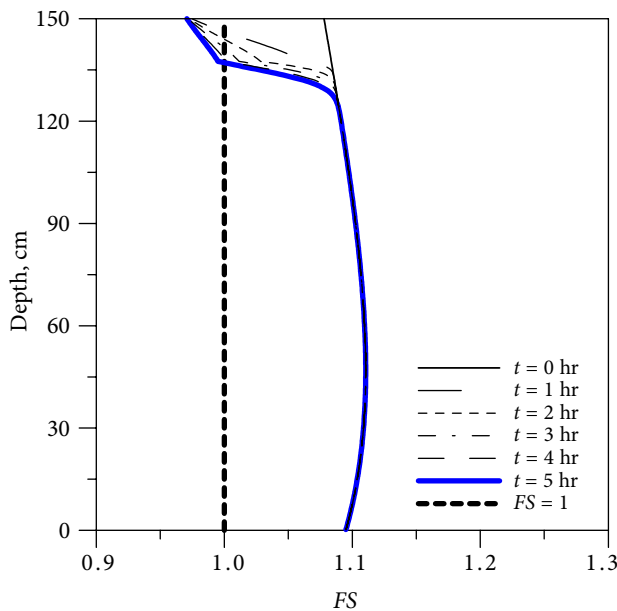


Fig. 13. Computed FS of the soil profile for application example 2

3.2. Application example 2

The second application example under investigation is a one-dimensional, transient, unsaturated groundwater flow problem in two-layered soil. The thickness of the soil L is 150 cm. In the case of two layered soils, layer A and layer B have thickness of 130 cm and 20 cm, respectively, as shown in Figure 11. Table 3 shows the soil parameters for two soil layers. In this example, layer B is more permeable than layer A. The hillslope has a slope angle of 31 degrees. The initial pressure head can be described as Eq. (53). The boundary conditions of the top and bottom boundaries are showed as Eq. (56) and Eq. (59). In this example, we considered $z_w = 80$ cm and $R_f = 5 \times 10^{-5}$ cm/hr. The total simulation time is five hours. The computed pressure heads versus elapsed time for this example are depicted in Figure 12. Figure 12 demonstrates that the fluctuation of pore water pressure in unsaturated layered soil dominates slope stability behavior. Besides, the pressure head at the interface increases quickly and the lowest FS may occur at the interface between two consecutive soil layers during a rainfall even if the top soil layer has the lower hydraulic conductivity. Figure 13 shows the computed FS with respect to time. It shows that the lowest FS occurs at the interface between two consecutive soil layers during a rainfall event. Figure 14 indicates the results of the FS at the depth of 135 cm for this example. It is found that the FS is strongly affected by the rainfall with respect to time. Accordingly, we may conclude that the variation of the pressure head caused by the rainfall event is strongly associated with soil layers and the interface between two consecutive soil layers may play a crucial role for the slope stability.

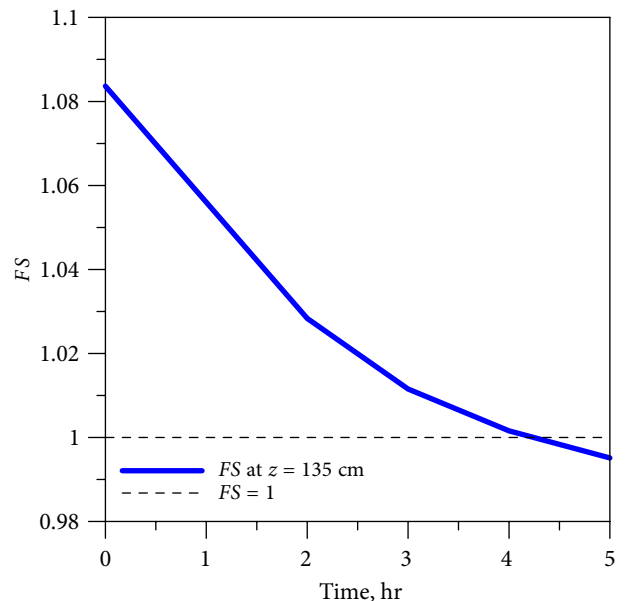


Fig. 14. The computed results of FS at $z = 135$ cm for application example 2

Table 3. Parameters used in the application example 2

Soil	Layer A	Layer B
Hydraulic conductivity K_s (cm/hr)	1×10^{-2}	1×10^{-1}
Saturated water content θ_s	0.43	0.50
Residual water content θ_r	0.08	0.11
α_g of the Gardner model	1×10^{-2}	1×10^{-2}

4. Discussions

In this paper, we carried out three verification examples and two application examples. From the results of verification examples, it is found that the computed numerical results agreed well with the analytical transient solution and those from the Skeel and Berzins' method. Comparison between the proposed numerical model and TRIGRS model (a simplified analytical-based model by assuming the Richards equation as a linear diffusion equation; Iverson 2000; Baum *et al.* 2002) was also conducted. We conducted a comprehensive study in this example with the consideration of several scenarios. It is found that the proposed model in this study can obtain more appropriate results when the value of the hydraulic diffusivity is larger than those from the TRIGRS model.

In the application examples, we further adopted the proposed model to study the slope stability of a landslide with the consideration of the infiltration from the ground surface for homogenous and layered soils. Since the numerical modeling of rainfall-induced shallow landslides in unsaturated layered soil using the variably saturated flow equation is still hardly found, this study demonstrates the advantages of using the proposed model to study the slope stability of landslides for layered soils. From the results of application examples, we also found that the fluctuation of pore water pressure in unsaturated layered soil dominates slope stability of landslides. Accordingly, it can be concluded that the proposed model in this study may have great potential to deal with problems of rainfall-induced shallow landslides in unsaturated layered soil.

The nature of the Richards equation is highly nonlinear and cannot directly provide an analytical solution. The SWCC is a major factor that influences the nonlinear physical relationship. To derive the proposed numerical model of rainfall-induced shallow landslides in unsaturated layered soil, we adopted the Gardner model to formulate the linearized Richards equation. It should be noted that the Gardner model is based on the one-parameter equation with an indication of the rate of desaturation of a soil. Parameters used in mathematical models for the SWCC include fixed points and two or more fitting constants that are usually selected to capture the general shape of the SWCC. Nevertheless, the Gardner model has only

one parameter which may have limitations for fitting the SWCC with complicated shape.

Conclusions

In this study, a numerical model of unsaturated layered soil for rainfall-induced shallow landslides was developed. The infinite slope stability analysis coupled with the hydrological model with the consideration of the SWCC proposed by Gardner was developed to model the shallow landslides triggered by rainfall. The fundamental concepts and the construct of the proposed method are clearly addressed. Findings from this study are drawn as follows:

1. The FDM for the numerical discretization of the variably saturated flow equation based on the Gardner model to deal with groundwater flow in unsaturated layered soil is derived. The proposed method can be used to solve slope stability problems with the appearance of layered soil.

2. The validity of the model is established for a number of test problems by comparing numerical results with the analytical solutions. It is found that the proposed method can be used to model groundwater flow in unsaturated layered soil with high accuracy.

3. The results obtained from this study demonstrate that the slope stability of landslides is strongly dependent to the hydraulic conductivity. It is found that the variation of pore water pressure in unsaturated layered influences the stability of a slope. Besides, the pressure head at the interface increases quickly and the lowest FS may occur at the interface between two consecutive soil layers during a rainfall even if the top soil layer is less permeable.

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Disclosure statement

The authors declare that they don't have any competing financial, professional, or personal interests from any other parties.

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