

1 **Title: Observation of incipient particle formation during flame synthesis by tandem differential**
2 **mobility analysis-mass spectrometry (DMA-MS)**

3

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6

7 **Abstract**

8 While flame aerosol reactor (FLAR) synthesis of nanoparticles is widely used to produce a range of
9 nanomaterials, incipient particle formation by nucleation and vapor condensation is not well understood.
10 This gap in our knowledge of incipient particle formation is caused by limitations in instruments, where,
11 during measurements, the high diffusivity of sub 3 nm particles significantly affects resolution and
12 transport loss. This work used a high resolution Differential Mobility Analyzer (DMA) and an
13 Atmospheric Pressure Interface-Mass Spectrometer (APi-TOF) to observe incipient particle formation
14 during flame synthesis. By tandemly applying these two instruments, Differential Mobility Analysis-
15 Mass Spectrometry (DMA-MS) measured the size and mass of the incipient particles simultaneously,
16 and the effective density of the sub 3 nm particles was estimated. The APi-TOF further provided the
17 chemical compositions of the detected particles based on highly accurate masses and isotope
18 distributions. This study investigated the incipient particle formation in flames with and without the
19 addition of synthesis precursors. Results from FLAR using two types of precursors including tetraethyl
20 orthosilicate (TEOS) and titanium isopropoxide (TTIP) are presented. The effect of the precursor feed
21 rates on incipient particle growth was also investigated.

22

23 **Keywords:** flame aerosol reactor (FLAR), combustion synthesis, incipient particle, differential mobility
24 analysis (DMA), mass spectrometry (MS)

25 **1. Introduction**

26 Flame synthesis is a gas-phase approach for producing nanomaterials on an industrial scale, due to its
27 high reaction temperature and fast reaction rates [1]. Flame-synthesized products, such as metal oxides
28 and carbon-based materials, have been widely applied in catalysis [2], solar energy utilization [3], sensor
29 technology [4], the rubber industry [5], and so on. During flame synthesis, the pyrolysis and oxidation of
30 synthesis precursors, the clustering of vapor molecules, particle nucleation, and particle growth through
31 coagulation, vapor condensation, and sintering, all take place in a single-step manner, resulting in a high
32 yield of nanoparticles [6, 7]. The high temperature and fast reaction rate in flames, on the other hand,
33 make it difficult to analyze detailed particle formation pathways. As the starting of the entire particle
34 formation process, especially below 3 nm, significantly affect the characteristics of the final products,
35 such as size, morphology, crystallinity, and chemical composition. However, research on the early
36 stages of particle formation mechanisms is limited by the performance of existing instruments [8, 9].
37 The *in situ* measurement of particle size distributions in aerosol reactors commonly rely on Differential
38 Mobility Analyzers (DMAs), which can classify a steady and narrowly monomobile stream of charged
39 particles from particles with a continuous spectrum of electrical mobilities [10]. However, the high
40 diffusivity of the sub 3 nm particles greatly decreases the resolution of conventional DMAs and
41 increases the diffusion loss of the measured particles [11]. Molecular Beam Mass Spectrometry (MBMS)
42 is typically used to analyze the composition of flame-generated particles, but the required low-pressure
43 and fuel-rich environment often deviates from the actual operating conditions of flame synthesis [12].
44 These limitations pose difficulties in measuring incipient particles during flame synthesis, hindering a
45 comprehensive understanding of the particle formation mechanisms.

46

47 With the development of a new set of instruments for investigating atmospheric particle nucleation [13]
48 and for analyzing protein properties [14] in the past decade, measuring sub 3 nm particle size and mass

49 at atmospheric conditions has become feasible. DMAs with sheath flow rates of over 100 lpm are used
50 to significantly reduce the residence time and the Brownian diffusion of particles in the instruments,
51 increasing the resolution by orders of magnitude when measuring sub 3 nm particles [15]. The recently
52 developed Atmospheric Pressure Interface Time-Of-Flight Mass Spectrometer (APi-TOF) is able to
53 detect and analyze the chemical composition of atmospheric ions and charged clusters with high
54 transmission and resolution [16]. Selected groups of molecular species were identified as playing an
55 important role in atmospheric particle nucleation and growth [17]. Tandem Differential Mobility
56 Analysis-Mass Spectrometry (DMA-MS) can simultaneously measure particle size and mass, and
57 critical information on particle structure, charging characteristics, and formation mechanisms is obtained
58 [18-20].

59

60 In this study, the incipient particle formation and growth below 3 nm in a premixed flat flame was
61 investigated with the DMA-MS technique, where a high resolution DMA and an APi-TOF were used to
62 counteract the particle Brownian diffusion and loss in the system. Direct measurement with the high
63 resolution DMA coupled with an aerosol electrometer was conducted to provide the size distributions of
64 particles generated during flame synthesis. The APi-TOF was used to determine the compositions of
65 important intermediate particles during the synthesis of SiO₂ and TiO₂. The structure of the incipient
66 flame-synthesized particles was further analyzed with the size and mass data measured by the DMA-MS.

67

68 **2. Methods**

69 *2.1 Experimental Setup*

70 Figure 1 shows the schematic diagram of the experimental setup. The system consisted a premixed flat
71 flame burner, a dilution sampling probe, a high resolution DMA (Herrmann-type [21]), an APi-TOF
72 (TOFWERK AG), and an aerosol electrometer (Model 3068B, TSI Inc.). A premixed flat flame was
73 used in this study due to its uniformity and stability [22]. The stainless steel burner consisted of two
74 concentric tubes with diameters of 0.75 inch and 1 inch, respectively, leaving a gap for passing a stream
75 of N₂ (>99.95%, Linde AG) to protect the flame from the environment. In order to achieve a uniform
76 velocity profile, the head of the burner was capped with a stainless steel honeycomb featuring a mesh
77 size of 0.5 mm. Below the honeycomb cap, the inner tube was filled with 2 mm stainless steel beads for
78 laminarizing the premixed flow streams. The gas mixture was composed of CH₄ (> 99.5%, Linde AG),
79 O₂ (> 99.95%, Linde AG), and N₂, maintained at total flow rates of 1 lpm, 2.85 lpm, and 8 lpm
80 respectively, with the help of mass flow controllers (MKS Inc.). The flame equivalence ratio (ϕ) was
81 calculated to be 0.7. Organometallic precursors for synthesizing nanoparticles were introduced into the
82 flame by bubbling a clean stream of N₂ through liquid precursors of tetraethyl orthosilicate (TEOS, >
83 98 %, Sigma-Aldrich Inc.) or titanium isopropoxide (TTIP, > 97 %, Sigma-Aldrich Inc.), at a
84 temperature of 20 °C. At high temperatures in flames, these synthesis precursors reacted through thermal
85 decomposition and oxidation to generate SiO₂ or TiO₂ nanoparticles, and previous studies demonstrated
86 that a large amount of sub 3 nm particle were formed during these processes [8, 9, 11]. The feed rates of
87 synthesis precursors were calculated according to the materials' saturation pressure data presented by
88 Jang [23] for TEOS and Siefering and Griffin [24] for TTIP. Flame temperature was measured with a
89 type R thermocouple, and the temperature profile is shown as the inset in Figure 1. Due to the low
90 concentration of the synthesis precursors (as shown in Table 1), the effect of adding precursors on flame
91 equivalence ratio and flame temperature were minimal. 5 mm above the head of the burner, a dilution

92 sampling probe was used to introduce the flame-generated particles to downstream instruments while
93 quenching further reactions and particle growth. By considering the thermal expansion of the sampled
94 flow, a dilution ratio of 200 was attained [9]. Note that ionization sources were not applied in the system,
95 implying that the downstream instruments measured the natively charged flame synthesized particles
96 only. Due to a series of reported chemical ionization reactions, the flame acted as a neutral plasma that
97 generated high concentrations of ions, which were sufficient to charge the incipient particles in the flame
98 at various equivalence ratios [12]. It should be also noticed that not all of the flame-generated particles
99 were charged, and the fraction of charged particles was a function of sampling height and flame
100 conditions. To study the properties of neutral particles, well-characterized chemical ionization sources
101 are needed to provide the neutral particles with known charges so that they can be detected by the
102 instruments [25].

103

104 A Herrmann-type DMA was used to classify sub 3 nm particles with high resolution. The DMA was
105 operated in a closed loop to maintain the same flow rate of the aerosol inlet and outlet flows. An inline
106 blower (DOMEL Inc.) provided the recirculating sheath flow (> 500 lpm), and an inline filter and heat
107 exchanger removed the remaining particles and released the heat generated by the blower, respectively.
108 The DMA classifies particles according to the relationship between the drag force and the electrostatic
109 force. When a voltage (V) is applied across the electrodes of the DMA, the classified particles have a
110 uniform mobility (Z) [10], which is further related to the size of a particle by $Z = Cne/3\pi\mu D_p$, where C
111 is the Cunningham slip correction factor, n is the number of charges on the particle, e is the electronic
112 charge, μ is the air viscosity, and D_p is the particle mobility size. For sub 3 nm particles, it is safe to
113 assume that classified particles carry a single charge [26], which was also observed by the mass
114 spectrometer during the experiments, since the isotope peaks that differ by mass units other than 1 were
115 not detected. The potential across the DMA was applied by a high voltage source (Spellman Inc.)

116 controlled by a Labview program. During the measurements, a step voltage of 3 V, a step time of 1 s,
117 and a voltage scanning range of 100 to 1000 V were used to classify particles with size from 0.5 to 2 nm.
118 It should be noted that the mobility size does not necessarily represent the particle physical size,
119 especially for sub 3 nm particles whose chemical composition may significantly determine the structure
120 of the cluster, while the size of non-spherical particles is poorly defined. Existing studies show that the
121 particle mobility size subtracted by 0.3 nm agrees well with the volume size of the sub 3 nm clusters
122 [18]. However, for simplicity, this study used particle mobility size as the indicator of particle physical
123 size. At the same time, the inverse mobility values are marked in the graphs for reference. Before
124 measuring the flame-generated particles, the DMA was first calibrated with the particles generated by
125 electro-spraying a 0.2 mM tetraheptylammonium bromide-methanol solution [15]. The mobility of
126 particles classified at an arbitrary DMA voltage can be accurately determined. We should also note that,
127 although the DMA sheath flow is significantly increased, the high diffusivity of the sub 3 nm particles
128 can still play a role in lowering the resolution of the DMA measurements, which is discussed in more
129 detail in our previous work [11].

130

131 Downstream of the DMA, an aerosol electrometer (EM) collected the classified monomobile particles at
132 a flow rate of 10 lpm. The recorded current was directly proportional to the charged particle
133 concentration if particles carried one unit charge. In the following, the particle size distributions were
134 qualitatively shown with the EM current as a function of particle size, since the data inversion of sub 3
135 nm particle size distributions was difficult to conduct due to the chemistry-dependent charging process
136 [11]. The APi-TOF measured the mass-to-charge ratio (m/z) of the DMA-classified sub 3 nm particles.
137 It can achieve a mass resolving power of 3000 Th/Th and a mass accuracy of 0.002%. The chemical
138 composition of the measured particles was further analyzed by tofTools (a Matlab based set of programs
139 [16, 27]), with the help of the highly accurate atomic mass and isotope distributions. Before

140 measurements, the APi-TOF was calibrated with nitrate ions produced by a chemical ionization source.
141 The DMA-MS technique simultaneously measures the particle size and the mass spectrum of the DMA-
142 classified particles. The mass-size relationship further reveals the structure and effective density of the
143 detected particles. This study investigated the mass-size relationship of the negatively charged particles
144 only, since the positively charged particles were found to be unstable when they were transported from
145 the DMA to the APi-TOF [20, 25, 28]. The obtained data are conveniently represented as contour plots,
146 with the x-axis being the particle size or electrical mobility, and y-axis being m/z. The color of the data
147 point denotes the relative abundance of the signal (black stands for the most intense signal, and white
148 stands for the least intense signal).

149

150 *2.2 Experimental Plan*

151 Five sets of experiments were conducted in this study and tabulated in Table 1. Test 1 studied the
152 properties of charged particles generated from blank flames without the addition of the precursors. Tests
153 2 and 3 were conducted with the addition of different types of precursors to investigate the formation
154 pathways of different types of nanoparticles. Tests 4 and 5, together with Test 2, further examined the
155 influence of synthesis precursor feed rates on the incipient particle formation during flame synthesis.

156

157

158 **3. Results and Discussion**

159 *3.1 Particle formation in blank flames*

160 Figure 2 shows the DMA-measured size distributions of the natively charged particles generated from
161 the blank flame without the addition of the precursors. Charged particles below 1.8 nm in both polarities
162 were observed in large quantities. The chemical ionization reactions in the flame are major sources of
163 these charged particles [12]. The positively charged particles had a larger and broader distribution of
164 sizes, yielding a relatively smaller electrical mobility compared to negatively charged particles. Similar
165 results were observed in commonly used ionization sources such as radioactive neutralizers [20, 29] and
166 corona dischargers [28]. The explanation for this phenomenon may be that, during the ionization process,
167 the relatively large organic molecules act as positive charge carriers, while negative charge carriers are
168 dominated by electrons. The electrons further combine with other molecules to form relatively smaller
169 negatively charged particles with low proton affinities. Based on charge neutrality, the concentration
170 difference between the positively and the negatively charged particles (Fig. 2) also indicated that the
171 remaining negative charge carriers were electrons, whose electrical mobility is too high ($> 1000 \text{ cm}^2/\text{Vs}$
172 [12]) to be measured by the DMA. These charged particles may actively collide with nanoparticles
173 during flame synthesis conditions. According to Fuchs charging theory [25], this difference in the
174 averaged electrical mobility will cause a higher fraction of particles carrying negative charges, which
175 has been observed in previous studies [30, 31].

176

177 Figure 3 displays the mass-size relationship of the negatively charged particles generated from the blank
178 flame. At the inverse mobility of 0.48 Vs/cm^2 (mobility size of 1.01 nm) where the EM detected the
179 highest particle concentration (Fig. 2), the mass spectrum indicated that these particles were mainly
180 composed of nitrate ions (NO_3^- at m/z of 62 Th, and $\text{HNO}_3 \cdot \text{NO}_3^-$ at m/z of 125 Th). The formation of
181 these nitrate ions might be caused by the active NO_x production in the blank flame, especially when the

182 combustible mixture was premixed and the flame was operated in a fuel-lean condition [32]. The NO_x
183 species may further react with water vapor to form nitrate ions. Due to their low proton affinities, these
184 nitrate ions became the dominant negative charge carriers, which was also observed in other types of
185 ionization sources [20, 28, 33]. The existence of two different masses (62 Da and 125 Da) corresponding
186 to a same mobility size (1.01 nm) suggested that the ions might be fragmented when they transported
187 from the DMA to the APi-TOF. Aside from the nitrate ions, negatively charged particles with sizes and
188 masses larger than 1.01 nm and 125 Da respectively were also detected. They were probably generated
189 during the collision between electrons with relatively larger organic molecules. The detection of a
190 spectrum of ions implied the complexity of chemical ionization reactions and the following particle
191 charging process in flames.

192

193 *3.2 Addition of flame synthesis precursors*

194 The sub 3 nm particle size distributions obtained under different synthesis conditions are shown in Fig. 4.
195 When synthesis precursors were added to the flame, the DMA measurements detected particles larger
196 than flame-generated particles, as a result of particle formation and growth. The sizes of these particles
197 were also found to be discrete instead of continuous, which implied that certain stable species might act
198 as important intermediates during particle formation. The average size and concentration of the
199 positively charged particles was still greater than those of the negatively charged particles, where the
200 existence of relatively small charge carriers might play an important role, as explained above.

201

202 The mass-size relationships of the flame-generated sub 3 nm particles during synthesis conditions are
203 shown in Fig. 5. Major negatively charged species, such as $\text{Si}_2\text{O}_4(\text{H}_2\text{O})_3\text{NO}_3^-$ (236 Th),
204 $\text{Si}_3\text{O}_6(\text{H}_2\text{O})_3\text{NO}_3^-$ (296 Th), $\text{TiO}_2(\text{NO}_3)_3^-$ (266 Th), and $\text{Ti}_2\text{O}_8(\text{NO}_3)_2^-$ (348 Th) were detected by the
205 APi-TOF. Table 2 lists the incipient particles with identified compositions and atomic masses. The

206 existence of nitrate ions in the silicon and titanium-containing particles implied a strong interaction
207 between the flame-generated ions and the synthesized particles. Since flame synthesis is conducted with
208 N₂ as the diluting gas on many occasions, these detected ions may act as contaminants for particle
209 growth and crystallization. Hence, further studies investigating the incipient particle formation
210 mechanisms without N₂ participation are needed. Figure 5 also shows that the detected particles fell into
211 different bands with strong signal intensities. A lot of these bands were separated by the atomic masses
212 of N (14 Da), O (18 Da), or H₂O (18 Da), instead of by the atomic mass of SiO₂ (60 Da) or TiO₂ (80 Da).
213 This result suggested the high involvement of blank flame-generated species in particle synthesis. To
214 confirm that the measured charged particles represent the characteristics of those generated in the flame,
215 further experiments using an enhanced condensation particle counter and a charged particle remover was
216 conducted. Results show that under the tested conditions, the charging fraction of sub 3 nm particles is
217 extremely high, where more than 90% of the flame-generated sub 3 nm particles were charged [34]. The
218 high charge fraction also contradicts the classical charging theories, suggesting that further studies on
219 the charging mechanisms of particles in flames are needed.

220

221 Depending on the type of nanoparticles synthesized, the mass-size relationships show different trends in
222 Fig. 5. At a same mobility size of 1.30 nm, the particles generated from TTIP reactions (~ 550 Da) were
223 heavier than the particles generated from TEOS reactions (~ 420 Da), which was possibly caused by the
224 higher atomic mass of titanium. Based on the mass and mobility size values, an effective density (ρ_{eff})

225 could be calculated for these incipient particles following the method introduced below. By assuming
226 that the detected particles were spherical, the effective density satisfied the relationship of

227 $m = \pi D_p^3 \rho_{eff} / 6$, where m is the particle mass and D_p is the particle size. The value of the effective

228 density should guarantee that most of the data points in the contour plot fall on the curve representing

229 the spherical particle mass-size relationship. Further calculation indicated that using particle volume size

230 (mobility size subtracted by 0.3 nm [18]) as D_p provided a better fit when calculating ρ_{eff} , since the
231 mobility size of a particle overestimates the physical size below 3 nm, due to the enhanced interaction
232 between the charged cluster and the dipole it induced in the gas molecules [18]. The particles generated
233 from flames with the addition of TEOS and TTIP had effective densities of 1.42 g/cm³ and 1.75 g/cm³,
234 respectively. These incipient particles were less dense than the SiO₂ and TiO₂ bulk crystals (2.65 g/cm³
235 and 4.23 g/cm³), possibly because these particles were in an amorphous state, where the atoms were not
236 tightly packed. The attachment of the species other than oxides, such as nitrate and organic molecules,
237 on these incipient particles might also lower the density of the detected particles.

238

239 The simultaneously measured mobility and mass of the flame-generated incipient particles could be
240 compared with the data calculated from empirically determined mass-mobility relationships. The most
241 widely used mass-mobility relationship was presented by Kilpatrick [35] and was further fitted with a
242 function of

$$243 \quad Z = \exp[-0.0347 \ln^2(m) - 0.0376 \ln(m) + 1.4662], \quad (1)$$

244 where Z and m represent the particle electrical mobility (unit: cm²/Vs) and atomic mass (unit: Da),
245 respectively [36]. The fitted functions of Kilpatrick's mass-mobility relationship are also displayed in
246 Fig. 5, showing that, at a same electrical mobility, the actual particle mass is higher than the mass
247 predicted by the empirical relationship. This discrepancy can be explained by the fact that the electrical
248 mobility is largely determined by the structure of particles, while particles with similar structures can
249 have different chemical compositions and atomic masses. Due to the existence of relatively heavy
250 species such as silicon and titanium, the flame-generated incipient particles had higher masses. Since
251 researchers often rely on Kilpatrick's relationship to convert the measured mobility to the mass of
252 particles in order to decipher the particle composition, this measured result proves that the existing
253 mass-mobility relationships will be dependent on the type of chemical species. Directly using these

254 relationships may therefore cause errors. To better predict the mass-mobility relationships, numerical
255 methods were used by researchers to consider the physical collision and potential interaction between
256 the molecular clusters and particles, where a desirable agreement was observed between the calculated
257 and the experimentally measured mass and mobility data [37].

258

259 *3.3 Effect of the synthesis precursor feed rates*

260 Figure 6 shows the effect of TEOS feed rates on the DMA-MS measured incipient particle size and mass
261 distributions. The concentration of the particles with larger sizes and masses increased as more
262 precursors were fed to the flame, indicating a stronger particle growth process by vapor condensation
263 and coagulation. The concentration of the smaller charged particles with low mass and size decreased
264 due to coagulation and the scavenging effect caused by the existing larger particles. The calculated
265 effective density remained the same (Fig. 6), demonstrating that the particle formation pathway did not
266 change as a function of precursor feed rates.

267 **4. Conclusions**

268 The incipient particle formation during flame synthesis was investigated for the first time using the
269 tandem DMA-MS technique. A high resolution DMA and an APi-TOF were used to measure the size
270 and mass of the sub 3 nm particles simultaneously in the flames without the addition of synthesis
271 precursors and with the addition of TEOS or TTIP.

272

273 Measurements in a blank flame detected a large number of sub 3 nm particles generated from chemical
274 ionization reactions, and determined that nitrate ions dominated in the negative ions. The formation of
275 nitrate ions may be related with the NO_x formation in flames. Measurements conducted with the addition
276 of synthesis precursors found particles with discrete size distributions, indicating the existence of stable
277 particles as important intermediates during flame synthesis. The blank flame-generated ions played an
278 important role during particle synthesis, since the APi-TOF observed the appearance of nitrate ions in
279 particles containing silicon or titanium. Future work on flame synthesis while manipulating the ion
280 properties, may bring new perspectives on manufacturing functional nanomaterials at high temperatures.
281 The effective densities of the incipient particles were calculated by assuming that the particles were
282 spheres. These particles had lower densities than the bulk materials of SiO₂ and TiO₂, possibly because
283 of the impurities in the particles and their amorphous structures. The commonly used Kilpatrick's mass-
284 mobility relationship was also evaluated in this study, and the difference between the measured data and
285 the Kilpatrick's relationship suggested that particle compositions largely determined particle mass and
286 mobility. As precursor feed rate increased, particles with larger mass and sizes were formed due to
287 enhanced coagulation and vapor condensation. The unchanged particle effective density implied that the
288 particle formation pathway in flames was not a function of precursor feed rates.

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304 **References:**

- 305
- 306 [1] P. Roth, Proc. Combust. Inst. 31 (2007) 1773-1788.
- 307 [2] F. Niu, S. Li, Y. Zong, Q. Yao, J. Phys. Chem. C, 118 (2014) 19165-19171.
- 308 [3] E. Thimsen, N. Rastgar, P. Biswas, J. Phys. Chem. C 112 (2008) 4134-4140.
- 309 [4] S. Bakrania, C. Perez, M. Wooldridge, Proc. Combust. Inst. 31 (2007) 1797-1804.
- 310 [5] P.T. Spicer, C. Artelt, S. Sanders, S.E. Pratsinis, J. Aerosol Sci. 29 (1998) 647-659.
- 311 [6] J. Jiang, D.-R. Chen, P. Biswas, Nanotechnology, 18 (2007) 285603.
- 312 [7] Y. Zhang, L. Shuiqing, S. Deng, Q. Yao, S.D. Tse, J. Aerosol Sci. 44 (2012) 71-82.
- 313 [8] J. Fang, Y. Wang, M. Attoui, T.S. Chadha, J.R. Ray, W.-N. Wang, Y.-S. Jun, P. Biswas, Anal. Chem.
- 314 86 (2014) 7523-7529.
- 315 [9] Y. Wang, P. Liu, J. Fang, W.-N. Wang, P. Biswas, J. Nanopart. Res. 17 (2015) 1-13.
- 316 [10] E. Knutson, K. Whitby, J. Aerosol Sci. 6 (1975) 443-451.
- 317 [11] Y. Wang, J. Fang, M. Attoui, T.S. Chadha, W.-N. Wang, P. Biswas, J. Aerosol Sci. 71 (2014) 52-64.
- 318 [12] A.B. Fialkov, Prog. Energy Combust. Sci. 23 (1997) 399-528.
- 319 [13] M. Kulmala, J. Kontkanen, H. Junninen, K. Lehtipalo, H.E. Manninen, T. Nieminen, T. Petäjä, M.
- 320 Sipilä, S. Schobesberger, P. Rantala, Science 339 (2013) 943-946.
- 321 [14] C.J. Hogan Jr, J.F. de la Mora, J. Am. Soc. Mass. Spectrom. 22 (2011) 158-172.
- 322 [15] J. Fernandez de la Mora, Aerosol Measurement: Principles, Techniques, and Applications, John
- 323 Wiley & Sons, Inc., Hoboken, New Jersey, U.S., 2011, p. 697.
- 324 [16] H. Junninen, M. Ehn, T. Petäjä, L. Luosujärvi, T. Kotiaho, R. Kostianen, U. Rohner, M. Gonin, K.
- 325 Fuhrer, M. Kulmala, Atmos. Meas. Tech. 3 (2010) 1039-1053.
- 326 [17] M. Ehn, J.A. Thornton, E. Kleist, M. Sipilä, H. Junninen, I. Pullinen, M. Springer, F. Rubach, R.
- 327 Tillmann, B. Lee, Nature 506 (2014) 476-479.
- 328 [18] C. Larriba, C.J. Hogan Jr, M. Attoui, R. Borrajo, J.F. Garcia, J.F. de la Mora, Aerosol Sci. Technol.
- 329 45 (2011) 453-467.
- 330 [19] A. Maißer, V. Premnath, A. Ghosh, T.A. Nguyen, M. Attoui, C.J. Hogan, PCCP 13 (2011) 21630-
- 331 21641.
- 332 [20] A. Maißer, J.M. Thomas, C. Larriba-Andaluz, S. He, C.J. Hogan, J. Aerosol Sci. 90 (2015) 36-50.
- 333 [21] W. Herrmann, T. Eichler, N. Bernardo, J. Fernández de la Mora, Annual Conference of the
- 334 American Association for Aerosol Research, St. Louis, MO, U.S., 2000.
- 335 [22] X. Wang, E. Cotter, K.N. Iyer, J. Fang, B.J. Williams, P. Biswas, Proc. Combust. Inst. 35 (2015)
- 336 2347-2354.
- 337 [23] H.D. Jang, Aerosol Sci. Technol. 30 (1999) 477-488.
- 338 [24] K. Siefert, G. Griffin, J. Electrochem. Soc. 137 (1990) 1206-1208.
- 339 [25] J. Kangasluoma, M. Attoui, H. Junninen, K. Lehtipalo, A. Samodurov, F. Korhonen, N. Sarnela, A.
- 340 Schmidt-Ott, D. Worsnop, M. Kulmala, J. Aerosol Sci. 87 (2015) 53-62.
- 341 [26] N. Fuchs, Geofis. pura appl. 56 (1963) 185-193.
- 342 [27] H. Junninen, *Data Cycle in Atmospheric Physics: from Detected Millivolts to Understanding the*
- 343 *Atmosphere*, PhD thesis, University of Helsinki, Helsinki, Finland, 2014.
- 344 [28] H. Manninen, A. Franchin, S. Schobesberger, A. Hirsikko, J. Hakala, A. Skromulis, J.
- 345 Kangasluoma, M. Ehn, H. Junninen, A. Mirme, Atmos. Meas. Tech. Discuss. 4 (2011) 2099-2125.
- 346 [29] P. Kallinger, G. Steiner, W.W. Szymanski, J. Nanopart. Res. 14 (2012) 1-8.
- 347 [30] M.M. Maricq, Aerosol Sci. Tech. 42 (2008) 247-254.
- 348 [31] M. Sahu, J. Park, P. Biswas, J. Nanopart. Res. 14 (2012) 1-11.
- 349 [32] S.R. Turns, An introduction to combustion, McGraw-Hill, New York, U.S., 1996.
- 350 [33] G. Steiner, T. Jokinen, H. Junninen, M. Sipilä, T. Petäjä, D. Worsnop, G.P. Reischl, M. Kulmala,
- 351 Aerosol Sci. Technol. 48 (2014) 261-270.

- 352 [35] Y. Wang, J. Kangasluoma, M. Attoui, H. Junninen, M. Kulmala, T. Petäjä, P. Biswas, *Anal. Chem.*
353 *Submitted.*
- 354 [35] W. Kilpatrick, in: *Proc. Annu. Conf. Mass Spectrosc*, 1971, pp. 320-325.
- 355 [36] J.M. Mäkelä, V. Jokinen, T. Mattila, A. Ukkonen, J. Keskinen, *J. Aerosol Sci.* 27 (1996) 175-190.
- 356 [37] C. Larriba, C.J. Hogan Jr, *J. Phys. Chem. A*, 117 (2013) 3887-3901.
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358 **List of Table Captions**

359 **Table 1.** Experimental plan.

360 **Table 2.** Chemical compositions of major negatively charged particles detected by the APi-TOF under
361 different flame conditions.

362

363 **List of Figure Captions**

364 **Figure 1.** Schematic diagram of the experimental setup for measuring the incipient particles generated
365 during flame synthesis. The Herrmann DMA classified particles with the same electrical mobility. The
366 APi-TOF and the electrometer provided the mass spectrum and the concentration of the classified
367 particles. The inset figure shows the temperature profile along the centerline above the burner.

368 **Figure 2.** Size distributions of sub 3 nm charged particles generated from the blank flame.

369 **Figure 3.** Contour plot showing the abundance of the blank flame-generated negatively charged
370 particles as a function of size and m/z .

371 **Figure 4.** Size distributions of sub 3 nm charged particles under different synthesis conditions. a)
372 positively charged particles; b) negatively charged particles. Note different scales of y-axes.

373 **Figure 5.** Contour plots showing the abundance of negatively charged particles during flame synthesis
374 conditions as a function of size and m/z . a) using TEOS as synthesis precursor; b) using TTIP as
375 synthesis precursor. The mass-size relationships assuming that particles were spherical are displayed as
376 short-dashed lines. The fit to Kilpatrick's mass-mobility relationship is displayed as long-dashed lines.

377 **Figure 6.** Contour plots showing the abundance of negatively charged particles generated at different
378 TEOS feed rates as a function of size and m/z . a) TEOS feed rate of 0.235 mmol/hr; b) TEOS feed rate
379 of 0.353 mmol/hr. The mass-size relationships assuming that particles were spherical are displayed as
380 short-dashed lines.

381

1 **Tables**2 **Table 1.** Experimental plan.

Test	Precursor Type	Feed rates [mmol/hr]
1	N/A	N/A
2	TEOS	0.118
3	TTIP	0.157
4	TEOS	0.235
5	TEOS	0.353

3

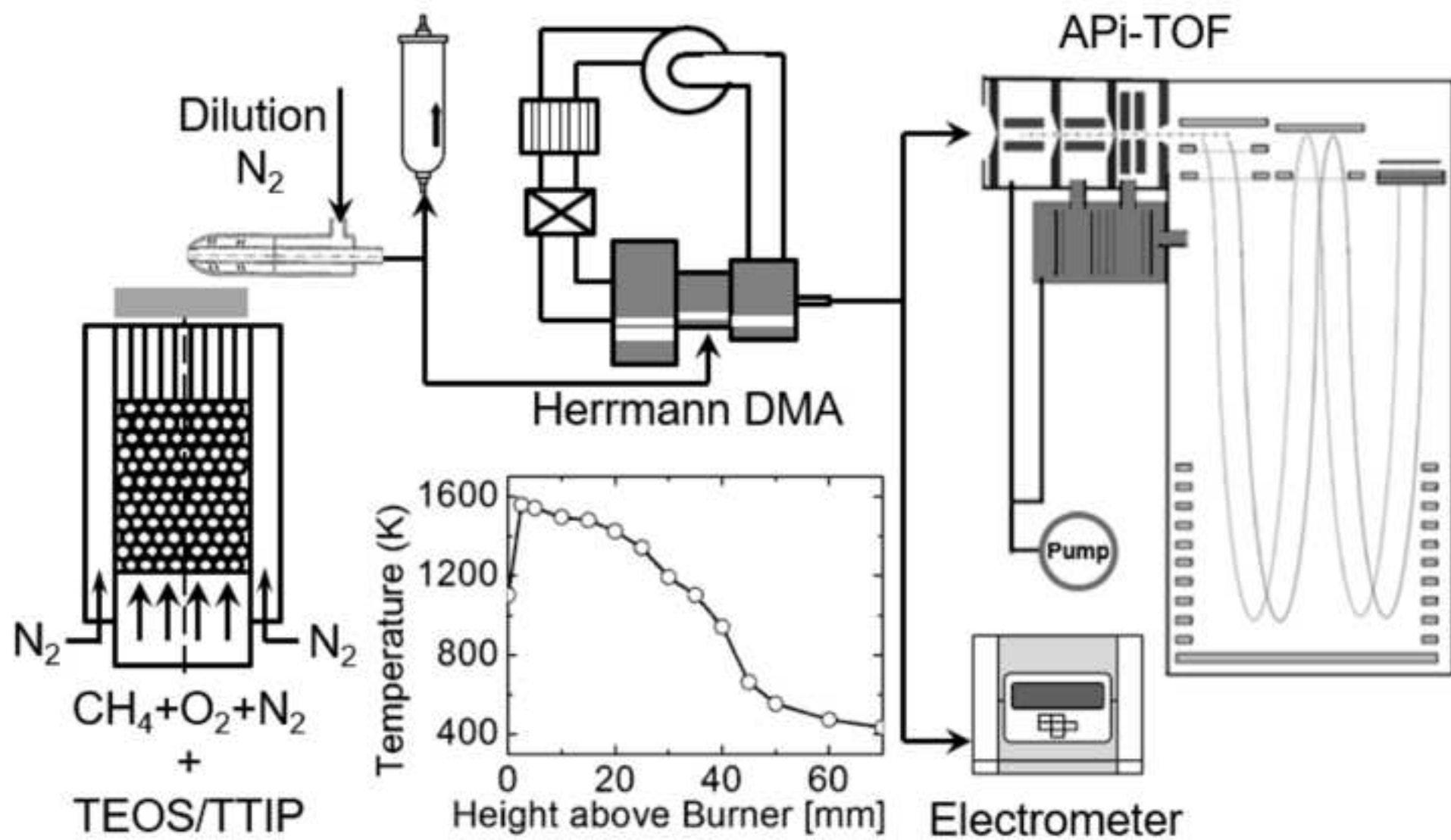
4 **Table 2.** Chemical compositions of major negatively charged particles detected by the APi-TOF under
5 different flame conditions.

Blank flame		TEOS addition		TTIP addition	
Chemical formula	m/z	Chemical formula	m/z	Chemical formula	m/z
NO_3^-	61.9878	NO_3^-	61.9878	NO_3^-	61.9878
$HNO_3 \cdot NO_3^-$	124.9835	$Si_2H_6NO_{10}^-$	235.9530	$TiN_2O_{10}^-$	235.9032
		$Si_3H_6NO_{12}^-$	295.9198	$TiN_3O_{11}^-$	265.9012
		$Si_3H_8NO_{13}^-$	313.9303	$Ti_2NO_{11}^-$	285.8437
		$Si_4H_{10}NO_{16}^-$	319.9077	$Ti_2N_2O_{14}^-$	347.8308

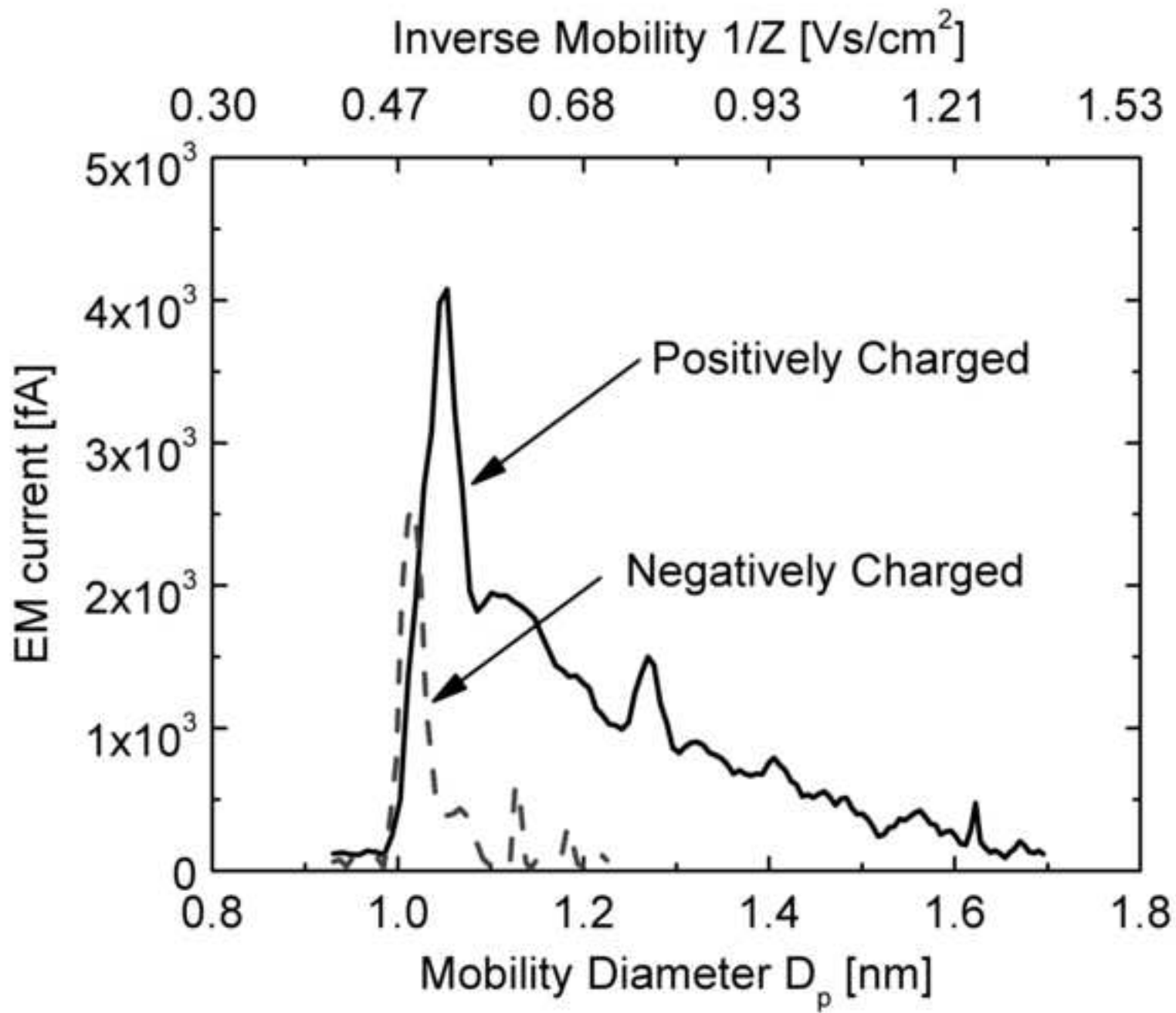
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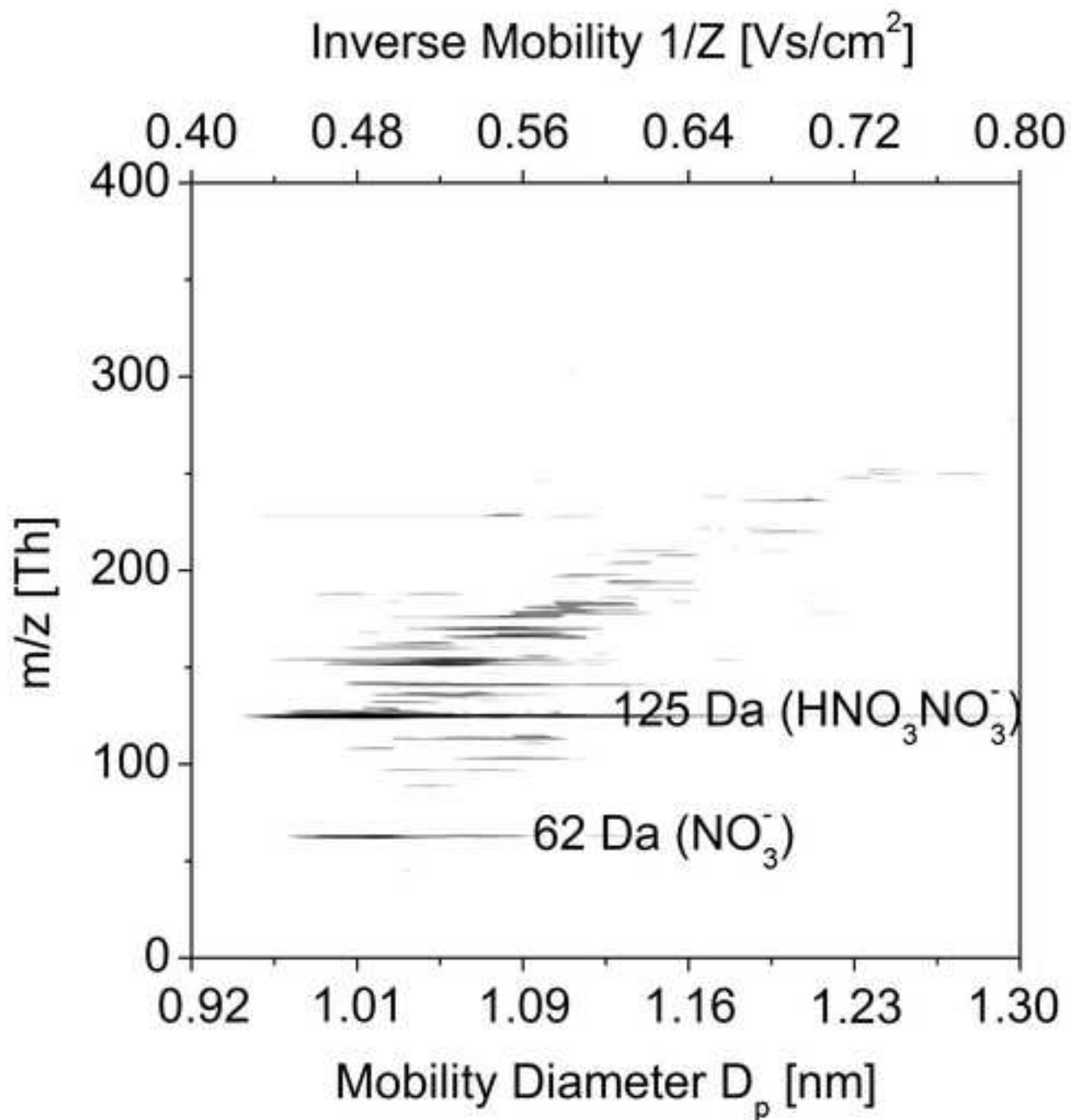
*Figure 1



*Figure 2



*Figure 3



*Figure 4

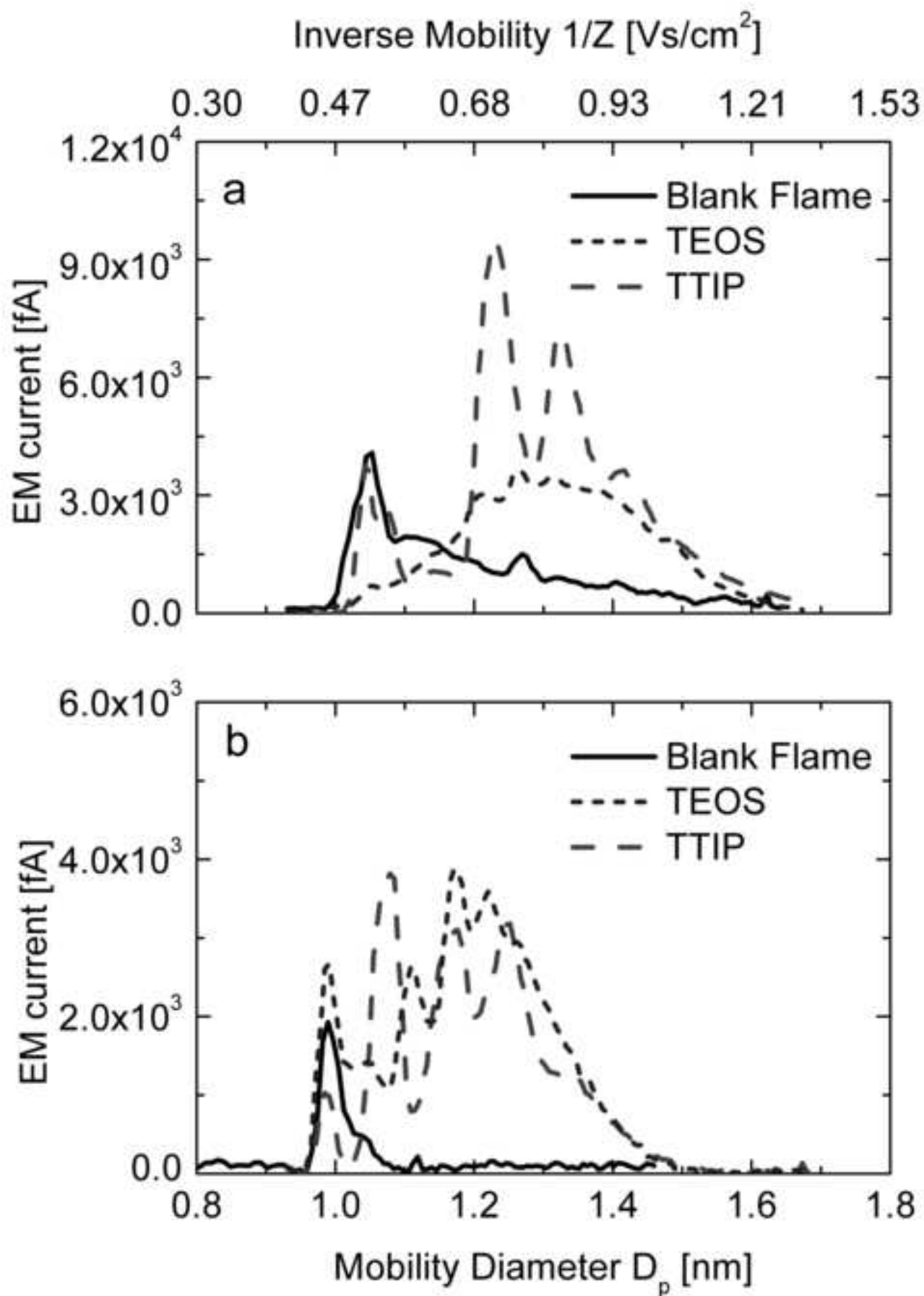


Figure 5

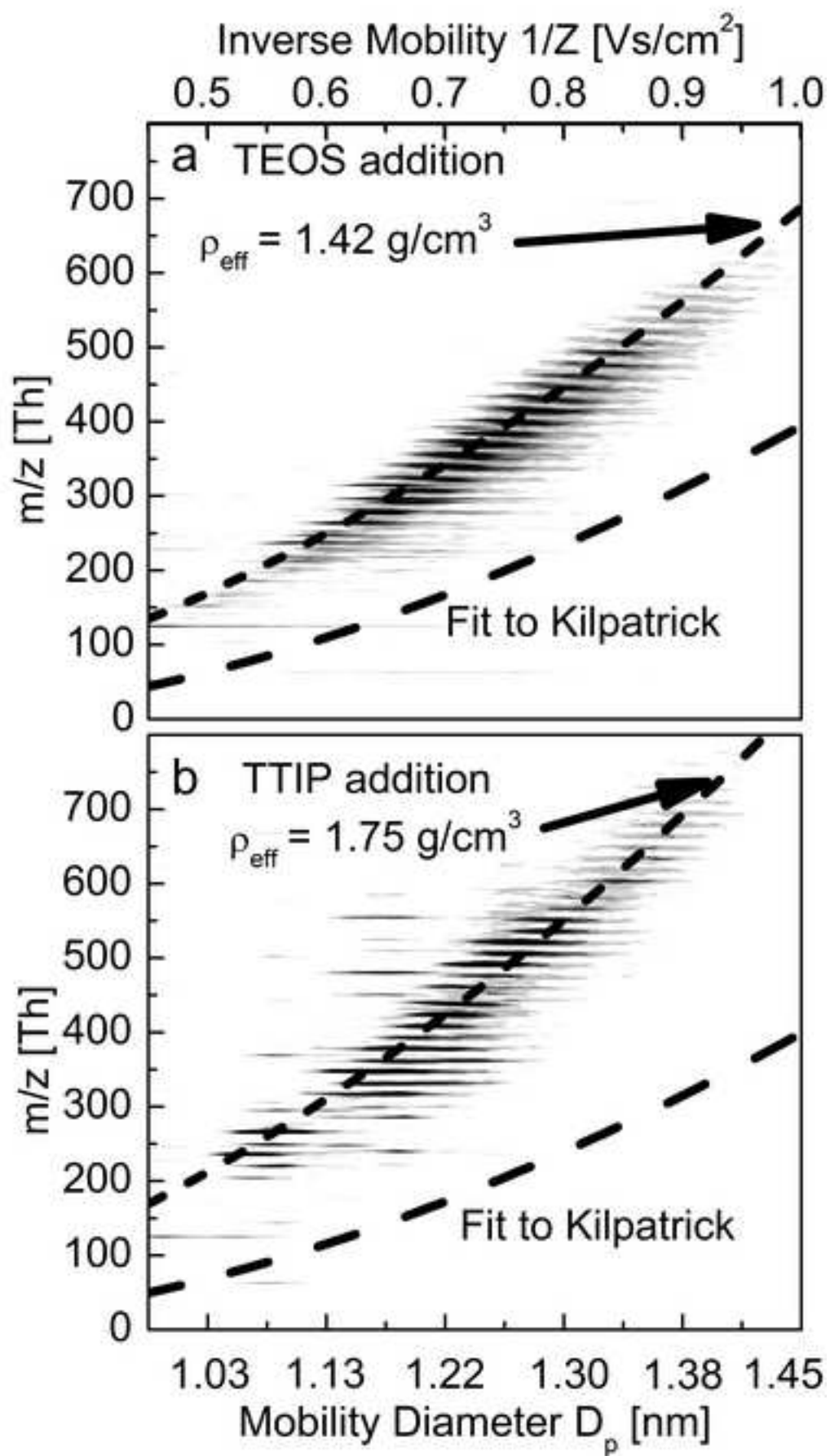


Figure 6

