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Observation of $Z_{\{b\}}(10610)$ and $Z_{\{b\}}(10650)$ Decaying to B Mesons

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Study of $e^+e^- \rightarrow B^{(*)}\bar{B}^{(*)}\pi^\pm$ at $\sqrt{s} = 10.866$ GeV

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We report the analysis of the three-body $e^+e^- \rightarrow B\bar{B}\pi^\pm$, $B\bar{B}^*\pi^\pm$, and $B^*\bar{B}^*\pi^\pm$ processes, including the first observations of the $Z_b^\pm(10610) \rightarrow [B\bar{B}^* + \text{c.c.}]^\pm$ and $Z_b^\pm(10650) \rightarrow [B^*\bar{B}^*]^\pm$ transitions that are found to dominate the corresponding final states. We measure Born cross sections for the three-body production of $\sigma(e^+e^- \rightarrow [B\bar{B}^* + \text{c.c.}]^\pm\pi^\mp) = (17.4 \pm 1.6(\text{stat.}) \pm 1.9(\text{syst.}))$ pb and $\sigma(e^+e^- \rightarrow [B^*\bar{B}^*]^\pm\pi^\mp) = (8.75 \pm 1.15(\text{stat.}) \pm 1.04(\text{syst.}))$ pb and set a 90% C.L. upper limit of $\sigma(e^+e^- \rightarrow [B\bar{B}]^\pm\pi^\mp) < 2.9$ pb. The results are based on a 121.4 fb⁻¹ data sample collected with the Belle detector at a center-of-mass energy near the $\Upsilon(10860)$ peak.

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95 Two new charged bottomonium-like resonances,¹⁰⁷
96 $Z_b(10610)$ and $Z_b(10650)$, have been observed recently¹⁰⁸
97 by the Belle Collaboration in $e^+e^- \rightarrow \Upsilon(nS)\pi^+\pi^-$,
98 $n = 1, 2, 3$ and $e^+e^- \rightarrow h_b(mP)\pi^+\pi^-$, $m = 1, 2$ [1, 2].¹⁰⁹
99 Analysis of the quark composition of the initial and final¹¹⁰
100 states reveals that these hadronic objects have an exotic¹¹¹
101 nature: Z_b should be comprised of (at least) four quarks¹¹²
102 including a $b\bar{b}$ pair. Several models [3] have been pro-¹¹³
103 posed to describe the internal structure of these states. In¹¹⁴
104 Ref. [4], it was suggested that $Z_b(10610)$ and $Z_b(10650)$ ¹¹⁵
105 states might be loosely bound $B\bar{B}^*$ and $B^*\bar{B}^*$ systems,¹¹⁶
106 respectively. If so, it is natural to expect the Z_b states¹¹⁷

to decay to final states with $B^{(*)}$ mesons at substantial rates.

Evidence for the three-body $\Upsilon(10860) \rightarrow B\bar{B}^*\pi$ decay has been reported previously by Belle, based on a data sample of 23.6 fb⁻¹ [5]. In this analysis, we use a data sample with an integrated luminosity of 121.4 fb⁻¹ collected near the peak of the $\Upsilon(10860)$ resonance ($\sqrt{s} = 10.866$ GeV) with the Belle detector [6] at the KEKB asymmetric-energy e^+e^- collider [7]. Note that we reconstruct only three-body $B^{(*)}\bar{B}^{(*)}\pi$ combinations with a charged primary pion. For brevity, we adopt the following notations: the set of $B^+\bar{B}^0\pi^-$ and

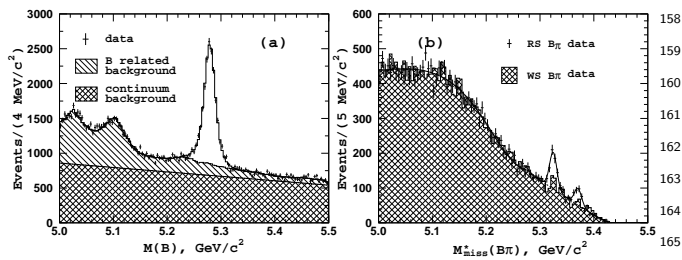


FIG. 1: The (a) invariant mass and (b) $M_{\text{miss}}^*(B\pi)$ distribu-
tion for B candidates in the B signal region. Points with error
bars represent the data. The open histogram in (a) shows the
result of the fit to data. The solid line in (b) shows the result
of the fit to the RS $B\pi$ data; the dashed line represents the
background level.

ing variables that characterize the event topology. Since
the momenta of the two B mesons produced from a three-
body $e^+e^- \rightarrow B^{(*)}B^{(*)}\pi$ decay are low in the center-of-
mass (c.m.) frame (below $0.9 \text{ GeV}/c$), the decay prod-
ucts of different B mesons are essentially uncorrelated so
that the event tends to be spherical. In contrast, hadrons
from continuum events tend to exhibit a back-to-back jet
structure. We use θ_{thr} , the angle between the thrust axis
of the B candidate and that of the rest of the event, to
discriminate between the two cases. The distribution of
 $|\cos\theta_{\text{thr}}|$ is strongly peaked near $|\cos\theta_{\text{thr}}| = 1.0$ for $c\bar{c}$
events and is nearly flat for $B^{(*)}B^{(*)}\pi$ events. We re-
quire $|\cos\theta_{\text{thr}}| < 0.80$ for the $B \rightarrow D^{(*)}\pi$ final states;
this eliminates about 81% of the continuum background
and retains 73% of the signal events.

We identify B candidates by their reconstructed in-
variant mass $M(B)$ and momentum $P(B)$ in the c.m.
frame. We require $P(B) < 1.35 \text{ GeV}/c$ to retain B
mesons produced in both two-body and multibody pro-
cesses. The $M(B)$ distribution for B candidates is shown
in Fig. 1(a). We perform a binned maximum likelihood
fit of the $M(B)$ distribution to the sum of a signal compo-
nent parameterized by a Gaussian function and two back-
ground components: one related to other decay modes of
 B mesons and one due to continuum $e^+e^- \rightarrow q\bar{q}$ pro-
cesses, where $q = u, d, s, c$. The shape of the B -related
background is determined from a large sample of generic
MC; the shape of the $q\bar{q}$ background is parameterized
with a linear function. The parameters of the signal
Gaussian, the normalization of the B -related background
and the parameters of the $q\bar{q}$ background float in the fit.
We find 12263 ± 168 fully reconstructed B mesons. The
 B signal region is defined by requiring $M(B)$ to be within
30 to 40 MeV/c^2 (depending on the B decay mode) of
the nominal B mass.

Reconstructed B^+ or \bar{B}^0 candidates are combined with
 π^- 's — the right-sign (RS) combination — and the
missing mass, $M_{\text{miss}}(B\pi)$, is calculated as $M_{\text{miss}}(B\pi) =$
 $\sqrt{(\sqrt{s} - E_{B\pi})^2/c^4 - P_{B\pi}^2/c^2}$, where $E_{B\pi}$ and $P_{B\pi}$ are
the measured energy and momentum of the reconstructed
 $B\pi$ combination. Signal $e^+e^- \rightarrow BB^*\pi$ events produce
a narrow peak in the $M_{\text{miss}}(B\pi)$ spectrum around the
nominal B^* mass while $e^+e^- \rightarrow B^*B^*\pi$ events produce
a peak at $m_{B^*} + \Delta m_{B^*}$, where $\Delta m_{B^*} = m_{B^*} - m_B$,
due to the missed photon from the $B^* \rightarrow B\gamma$ decay. It
is important to note here that, according to signal MC,
 $BB^*\pi$ events, where the reconstructed B is the one from
the B^* , produce a peak in the $M_{\text{miss}}(B\pi)$ distribution at
virtually the same position as $BB^*\pi$ events, where the
reconstructed B is the primary one. To remove the cor-
relation between $M_{\text{miss}}(B\pi)$ and $M(B)$ and to improve
the resolution, we use $M_{\text{miss}}^* = M_{\text{miss}}(B\pi) + M(B) - m_B$
instead of $M_{\text{miss}}(B\pi)$. The M_{miss}^* distribution for the RS
combinations is shown in Fig. 1(b), where peaks corre-
sponding to the $BB^*\pi$ and $B^*B^*\pi$ signals are evident.
Combinations with π^+ — the wrong sign (WS) combi-

$B^-B^0\pi^+$ final states is referred to as $BB\pi$; the set of
 $B^+B^*0\pi^-$, $B^-B^*0\pi^+$, $B^0B^*\pi^-$ and $\bar{B}^0B^*\pi^-$ final
states is referred to as $BB^*\pi$; and the set of $B^{*+}\bar{B}^*0\pi^-$
and $B^{*-}B^*0\pi^+$ final states is denoted as $B^*B^*\pi$. The in-
clusion of the charge conjugate mode is implied through-
out this report.

We use Monte Carlo (MC) events generated with Evt-
Gen [8] and then processed through a detailed detector
simulation implemented in GEANT3 [9]. The simulated
samples for $e^+e^- \rightarrow q\bar{q}$ ($q = u, d, s, c$, or b) are equiv-
alent to six times the integrated luminosity of the data
and are used to develop criteria to separate signal events
from backgrounds, identify types of background events,
determine the reconstruction efficiency and parameterize
the distributions needed for the extraction of the signal
decays.

B mesons are reconstructed in the following decay
channels: $B^+ \rightarrow J/\psi K^{*+}$, $B^+ \rightarrow \bar{D}^{*0}\pi^+$, $B^0 \rightarrow$
 $J/\psi K^{*0}$, $B^0 \rightarrow D^{*-}\pi^+$. We use Belle standard tech-
niques [10] to reconstruct primary particles such as pho-
tons, pions, kaons, and leptons. The K^{*0} (K^{*+}) is re-
constructed in the $K^+\pi^-$ ($K^0\pi^+$) final state; the invari-
ant mass of the K^* candidate is required to be within
150 MeV/c^2 of the nominal K^* mass [11]. The invari-
ant mass of a $J/\psi \rightarrow \ell^+\ell^-$ candidate is required to
be within 30 (50) MeV/c^2 for $\ell = e$ (μ), of the nom-
inal J/ψ mass. Neutral (charged) D mesons are re-
constructed in the $K^-\pi^+$, $K^-\pi^+\pi^0$, and $K^-\pi^-\pi^+\pi^+$
($K^-\pi^+\pi^+$) modes. To identify D^* candidates, we require
 $|M(D\pi) - M(D) - \Delta m_{D^*}| < 3 \text{ MeV}/c^2$, where $M(D\pi)$
and $M(D)$ are the reconstructed masses of the D^* and
 D candidates, respectively, and $\Delta m_{D^*} = m_{D^*} - m_D$ is the
difference between the nominal D^* and D masses. The
mass windows for narrow states quoted above correspond
to a $\pm 2.5\sigma$ requirement.

The dominant background comes from $e^+e^- \rightarrow c\bar{c}$ con-
tinuum events, where true D mesons produced in e^+e^-
annihilation are combined with random particles to form
a B candidate. This type of background is suppressed us-

nations — are used to evaluate the shape of the combinatorial background. (The $B \rightarrow J/\psi K^0$ mode is not included in the WS sample but both combinations with π^+ and π^- are added to the RS sample.) We apply factor of 1.19 ± 0.01 [12] to the WS distribution to normalize it to the expected number of the background events in the RS sample. There is also a hint for a peaking structure in the WS M_{miss}^* distribution, shown as a hatched histogram in Fig. 1(b). Due to $B^0 - \bar{B}^0$ oscillations, we expect a fraction of the produced B^0 mesons to decay as \bar{B}^0 given by $0.5x_d^2/(1+x_d^2) = 0.1861 \pm 0.0024$, where x_d is the B^0 mixing parameter [11].

Note that the momentum spectrum of B mesons produced in events with initial-state radiation (ISR), $e^+e^- \rightarrow \gamma B\bar{B}$, overlaps significantly with that for B mesons from the three-body $e^+e^- \rightarrow B^{(*)}B^{(*)}\pi$ processes. However, ISR events do not produce peaking structures in the M_{miss}^* distribution.

A binned maximum likelihood fit is performed to fit the M_{miss}^* distribution to the sum of three Gaussian functions to represent three possible signals and two threshold components $A_k(x_k - M_{\text{miss}}^*)^{\alpha_k} \exp\{(M_{\text{miss}}^* - x_k)/\delta_k\}$ ($k = 1, 2$) to parameterize the $q\bar{q}$ and two-body $B^{(*)}B^{(*)}$ backgrounds. The means and widths of the signal Gaussian functions are fixed from the signal MC simulation. The parameters A_k , α_k , δ_k of the background functions are free parameters of the fit; the threshold parameters x_k are fixed from the generic MC. ISR events produce an M_{miss}^* distribution similar to that for $q\bar{q}$ events; these two components are modeled by a single threshold function. The resolution of the signal peaks in Fig. 1(b) is dominated by the c.m. energy spread and is fixed at $6.5 \text{ MeV}/c^2$ and $6.2 \text{ MeV}/c^2$ for the $BB^*\pi$ and $B^*B^*\pi$, respectively as determined from the signal MC. The fit to the RS spectrum yields $N_{BB^*\pi} = 13 \pm 25$, $N_{BB^*\pi} = 357 \pm 30$ and $N_{B^*B^*\pi} = 161 \pm 21$ signal events. The statistical significance of the observed $BB^*\pi$ and $B^*B^*\pi$ signal is 9.3σ and 8.1σ , respectively. The statistical significance is calculated as $\sqrt{-2\ln(\mathcal{L}_0/\mathcal{L}_{\text{sig}})}$, where \mathcal{L}_{sig} and \mathcal{L}_0 denote the likelihood values obtained with the nominal fit and with the signal yield fixed at zero, respectively.

For the subsequent analysis, we require $|M_{\text{miss}}^* - m_{B^*}| < 15 \text{ MeV}/c^2$ to select $BB^*\pi$ signal events and $|M_{\text{miss}}^* - (m_{B^*} + \Delta m_B)| < 12 \text{ MeV}/c^2$, where $\Delta m_B = m_{B^*} - m_B$, to select $B^*B^*\pi$ events. For the selected $B^{(*)}B^*\pi$ candidates, we calculate $M_{\text{miss}}(\pi) = \sqrt{(\sqrt{s} - E_\pi)^2/c^4 - P_\pi^2/c^2}$, where E_π and P_π are the reconstructed energy and momentum, respectively, of the charged pion in the c.m. frame. The $M_{\text{miss}}(\pi)$ distributions are shown in Fig. 2 [13]. We perform a simultaneous binned maximum likelihood fit to the RS and WS samples, assuming the same number (after normalization) and distribution of background events in both samples and known fraction of signal events in the RS sample that leaks to the WS sample due to mixing. To fit the

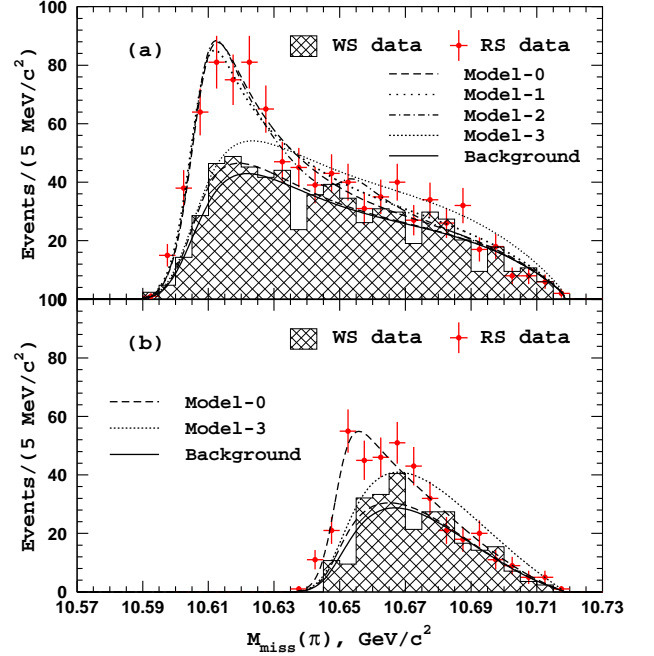


FIG. 2: The $M_{\text{miss}}(\pi)$ distribution for the (a) $BB^*\pi$ and (b) $B^*B^*\pi$ candidate events. Normalization factor is applied for the WS distributions.

$M_{\text{miss}}(\pi)$ spectrum, we use the function

$$F(m) = [f_{\text{sig}}S(m) + B(m)]\epsilon(m)F_{\text{PHSP}}(m), \quad (1)$$

where $m \equiv M_{\text{miss}}(\pi)$; $f_{\text{sig}} = 1.0$ (0.1366 ± 0.0032 , [14]) for the RS (WS) sample; $S(m)$ and $B(m)$ are the signal and background PDFs, respectively; and $F_{\text{PHSP}}(m)$ is the phase space function. To account for the instrumental resolution, we smear the function $F(m)$ with a Gaussian function with $\sigma = 6.0 \text{ MeV}/c^2$ that is dominated by the c.m. energy spread. The reconstruction efficiency is parametrized as $\epsilon(m) \sim \exp((m - m_0)/\Delta)(1 - m/m_0)^{3/4}$, where $m_0 = 10.718 \pm 0.001 \text{ GeV}/c^2$ is an efficiency threshold and $\Delta = 0.094 \pm 0.002 \text{ GeV}/c^2$.

The distribution of background events is parameterized as $B_{B^{(*)}B^*\pi}(m) = b_0 e^{-\beta\delta_m}$, where b_0 and β are fit parameters and $\delta_m = m - (m_{B^{(*)}} + m_{B^*})$. A general form of the signal PDF is written as

$$S(m) = |\mathcal{A}_{Z_b(10610)} + \mathcal{A}_{Z_b(10650)} + \mathcal{A}_{\text{nr}}|^2, \quad (2)$$

where $\mathcal{A}_{\text{nr}} = a_{\text{nr}} e^{i\phi_{\text{nr}}}$ is the non-resonant amplitude parameterized as a complex constant and the two Z_b amplitudes, \mathcal{A}_{Z_b} , are parameterized with Breit-Wigner functions $\mathcal{A}_{Z_b} = a_Z e^{i\phi_Z} / (m^2 - m_Z^2 - i\Gamma_Z m_Z)$. The masses and widths of the Z_b states are fixed at the values obtained from the analyses of $e^+e^- \rightarrow \Upsilon(nS)\pi^+\pi^-$ and $e^+e^- \rightarrow h_b(m_P)\pi^+\pi^-$: $M_{Z_b(10610)} = 10607.2 \pm 2.0 \text{ MeV}/c^2$, $\Gamma_{Z_b(10610)} = 18.4 \pm 2.4 \text{ MeV}$ and $M_{Z_b(10650)} = 10652.2 \pm 1.5 \text{ MeV}/c^2$, $\Gamma_{Z_b(10650)} = 11.5 \pm 2.2 \text{ MeV}$ [1].

TABLE I: Summary of fit results to the $M_{\text{miss}}(\pi)$ distributions for the three-body $BB^*\pi$ and $B^*B^*\pi$ final states.

Mode	Parameter	Model-0	Model-1		Model-2		Model-3
			Solution 1	Solution 2	Solution 1	Solution 2	
$BB^*\pi$	$f_{Z_b(10610)}$	1.0	1.45 ± 0.24	0.64 ± 0.15	1.01 ± 0.13	1.18 ± 0.15	–
	$f_{Z_b(10650)}$	–	–	–	0.05 ± 0.04	0.24 ± 0.11	–
	$\phi_{Z_b(10650)}$, rad.	–	–	–	-0.26 ± 0.68	-1.63 ± 0.14	–
	f_{nr}	–	0.48 ± 0.23	0.41 ± 0.17	–	–	1.0
	ϕ_{nr} , rad.	–	-1.21 ± 0.19	0.95 ± 0.32	–	–	–
	$-2 \log \mathcal{L}$	–304.7	–300.6	–300.5	–301.4	–301.4	–344.5
$B^*B^*\pi$	$f_{Z_b(10650)}$	1.0	1.04 ± 0.15	0.77 ± 0.22	–	–	–
	f_{nr}	–	0.02 ± 0.04	0.24 ± 0.18	–	–	1.0
	ϕ_{nr} , rad.	–	0.29 ± 0.10	1.10 ± 0.44	–	–	–
	$-2 \log \mathcal{L}$	–182.4	–182.4	–182.4	–	–	–209.7

We first analyze the $BB^*\pi$ [$B^*B^*\pi$] data with the simplest hypothesis, Model-0, that includes only the $Z_b(10610)$ [$Z_b(10650)$] amplitude. Results of the fit are shown in Fig. 2; the numerical results are summarized in Table I. The fraction f_X of the total three-body signal attributed to a particular quasi-two-body intermediate state is calculated as $f_X = \int |\mathcal{A}_X|^2 dm / \int S(m) dm$, where \mathcal{A}_X is the amplitude for a particular component X of the three-body amplitude. Next, we extend the hypothesis to include a possible non-resonant component, Model-1, and repeat the fit to the data. Then the $BB^*\pi$ data is fit to a combination of two Z_b amplitudes, Model-2. In both cases, the addition of an extra component to the amplitude does not give a statistically significant improvement in the data description: the likelihood value is only marginally improved (see Table I). The addition of extra components to the amplitude also produces multiple maxima in the likelihood function. As a result, we use Model-0 as our nominal hypothesis. Finally, we fit both samples to a pure non-resonant amplitude (Model-3). In this case, the fit is significantly worse.

If the parameters of the Z_b resonances are allowed to float, the fit to the $BB^*\pi$ data with Model-0 gives 10605 ± 6 MeV/ c^2 and 25 ± 7 MeV for the $Z_b(10610)$ mass and width, respectively, and the fit to the $B^*B^*\pi$ data gives 10648 ± 13 MeV/ c^2 and 23 ± 8 MeV for the $Z_b(10650)$ mass and width, respectively. The large errors here reflect the strong correlation between the resonance parameters.

The three-body Born cross sections are calculated as

$$\sigma(e^+e^- \rightarrow f) = \frac{N_f}{L \mathcal{B}_f \alpha \eta (1 + \delta_{\text{ISR}}) |1 - \Pi|^2}, \quad (3)$$

where N_f is the three-body signal yield and $L = 121.4 \text{ fb}^{-1}$ is the total integrated luminosity. The efficiency-weighted sum of B -meson branching fractions \mathcal{B}_f is determined using both signal MC and two-body $e^+e^- \rightarrow B^{(*)}\bar{B}^{(*)}$ events in data. To avoid the large systematic uncertainties associated with the determination of reconstruction efficiencies for B and D decays to multibody final states, we select a subset of two-body

modes: $B^+ \rightarrow \bar{D}^0[K^+\pi^-]\pi^+$ and $B \rightarrow J/\psi[\ell^+\ell^-]K$, and calculate $\mathcal{B}_f = \mathcal{B}_f^{\text{sel}} \times N_{B^{(*)}\bar{B}^{(*)}}^{\text{all}} / N_{B^{(*)}\bar{B}^{(*)}}^{\text{sel}}$, where the superscripts “sel” and “all” refer to quantities determined for the selected subset of B decay modes and for the full set of modes, respectively. Two-body $e^+e^- \rightarrow B^{(*)}\bar{B}^{(*)}$ events are selected with the requirement $0.90 \text{ GeV}/c < P(B) < 1.35 \text{ GeV}/c$; the B yield is determined from the fit to the $M(B)$ distribution. We find $N_{B^{(*)}\bar{B}^{(*)}}^{\text{all}} = 10131 \pm 152$ and $N_{B^{(*)}\bar{B}^{(*)}}^{\text{sel}} = 2406 \pm 62$. (MC studies show no significant dependence of the reconstruction efficiency on the B momentum.) To account for the non-uniform distribution of signal events over the phase space, we introduce an efficiency correction factor η determined from the MC simulation with signal events generated according to the nominal model. Since we do not observe a signal in the $BB\pi$ final state, no correction is made for this channel. A factor $\alpha = 0.897 \pm 0.002$ is introduced to correct for the effect of neutral B -meson oscillations that is determined using the known B^0 mixing parameter x_d and the yield ratio in data of two-body events with a reconstructed neutral vs. charged B meson. The ISR correction, $1 + \delta_{\text{ISR}}$, for the $B^{(*)}B^*\pi$ final states is calculated using recent results on $\sigma(e^+e^- \rightarrow h_b(mP)\pi^+\pi^-)$ [15] and an observation that the $\Upsilon(5S) \rightarrow h_b(mP)\pi^+\pi^-$ transitions are saturated by the intermediated Z_b production [1]; for the $BB\pi$ final state we assume constant cross section. For the vacuum polarization correction we use $1/|1 - \Pi|^2 = 0.928$ [16]. The results are summarized in Table II.

TABLE II: Summary of results on three-body cross sections. The first (or sole) uncertainty is statistical; the second is systematic.

Parameter	$BB\pi$	$BB^*\pi$	$B^*B^*\pi$
N_f , Events	13 ± 25	357 ± 30	161 ± 21
\mathcal{B}_f , 10^{-6}	293 ± 22	276 ± 21	223 ± 17
η	1.0	1.066	1.182
$1 + \delta_{\text{ISR}}$	0.720 ± 0.017	0.598 ± 0.016	0.594 ± 0.016
σ , pb	< 2.9	$17.4 \pm 1.6 \pm 1.9$	$8.75 \pm 1.15 \pm 1.04$

TABLE III: B branching fractions for the $Z_b^+(10610)$ and $Z_b^+(10650)$ decays. The first quoted uncertainty is statistical; the second is systematic.

Channel	Fraction, %	
	$Z_b(10610)$	$Z_b(10650)$
$\Upsilon(1S)\pi^+$	$0.54^{+0.16+0.11}_{-0.13-0.08}$	$0.17^{+0.07+0.03}_{-0.06-0.02}$
$\Upsilon(2S)\pi^+$	$3.62^{+0.76+0.79}_{-0.59-0.53}$	$1.39^{+0.48+0.34}_{-0.38-0.33}$
$\Upsilon(3S)\pi^+$	$2.15^{+0.55+0.60}_{-0.42-0.43}$	$1.63^{+0.33+0.33}_{-0.42-0.28}$
$h_b(1P)\pi^+$	$3.45^{+0.87+0.86}_{-0.71-0.63}$	$8.41^{+2.43+1.49}_{-2.12-1.06}$
$h_b(2P)\pi^+$	$4.67^{+1.24+1.18}_{-1.00-0.89}$	$14.7^{+3.2+2.8}_{-2.8-2.3}$
$B^+\bar{B}^{*0} + \bar{B}^0B^{*+}$	$85.6^{+1.5+1.5}_{-2.0-2.1}$	—
$B^{*+}\bar{B}^{*0}$	—	$73.7^{+3.4+2.7}_{-4.4-3.5}$

The dominant sources of systematic uncertainties for the three-body production cross sections are the uncertainties in the signal yield extraction (6.9% for $BB^*\pi$ and 8.7% for $B^*B^*\pi$), in the reconstruction efficiency (7.6%) (including secondary branching fractions [11]), in the correction factor α (1%), in the integrated luminosity (1.4%) and in the ISR correction (2.7%). The overall systematic uncertainties for the three-body cross sections are estimated to be 7.9%, 10.8%, and 12.0% for the $BB\pi$, $BB^*\pi$, and $B^*B^*\pi$ final states, respectively.

Using the results of the fit to the $M_{\text{miss}}(\pi)$ spectra with the nominal model (Model-0 in Table I) and the results of the analyses of $e^+e^- \rightarrow \Upsilon(nS)\pi^+\pi^-$ [1] and $e^+e^- \rightarrow h_b(mP)\pi^+\pi^-$ [15, 17], we calculate the ratio of the branching fractions $\mathcal{B}(Z_b^+(10610) \rightarrow \bar{B}^0B^{*+} + B^+\bar{B}^{*0})/\mathcal{B}(Z_b^+(10610) \rightarrow \text{bottomonium}) = 5.93^{+0.99+1.01}_{-0.69-0.73}$ and $\mathcal{B}(Z_b^+(10650) \rightarrow B^{*+}\bar{B}^{*0})/\mathcal{B}(Z_b^+(10650) \rightarrow \text{bottomonium}) = 2.80^{+0.69+0.54}_{-0.40-0.36}$. We also calculate the relative fractions for Z_b decays, assuming that they are saturated by the already observed $\Upsilon(nS)\pi$, $h_b(mP)\pi$, and $B^{(*)}B^*$ channels. The results are presented in Table III.

To summarize, we report the first observations of the three-body $e^+e^- \rightarrow BB^*\pi$ and $e^+e^- \rightarrow B^*B^*\pi$ processes with a statistical significance above 8σ . Measured Born cross sections are $\sigma(e^+e^- \rightarrow [B\bar{B}^* + \text{c.c.}]^\pm\pi^\mp) = (17.4 \pm 1.6 \pm 1.9)$ pb and $\sigma(e^+e^- \rightarrow [B^*\bar{B}^*]^\pm\pi^\mp) = (8.75 \pm 1.15 \pm 1.04)$ pb. For the $e^+e^- \rightarrow BB\pi$ process, we set a 90% confidence level upper limit of $\sigma(e^+e^- \rightarrow [B\bar{B}]^\pm\pi^\mp) < 2.9$ pb. The analysis of the $B^{(*)}B^*$ mass spectra indicates that the total three-body rates are dominated by the intermediate $e^+e^- \rightarrow Z_b(10610)^\mp\pi^\pm$ and $e^+e^- \rightarrow Z_b(10650)^\mp\pi^\pm$ transitions for the $BB^*\pi$ and $B^*B^*\pi$ final states, respectively.

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