DOI (Digital Object Identifier) 10.1007/s102310100004

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On a class of Einstein hypersurfaces immersed in a Riemannian manifold

Received: April 7, 1999; in final form: January 7, 2000 Published online: May 10, 2001 – © Springer-Verlag 2001

Abstract. Let $x : M \to \tilde{M}$ be an isometric immersion of a hypersurface M into an (n + 1)-dimensional Riemannian manifold \tilde{M} and let ρ_i $(i \in \{1, ..., n\})$ be the principal curvatures of M. We denote by E and P the distinguished vector field and the curvature vector field of M, respectively, in the sense of [8].

If M is structured by a P-parallel connection [7], then it is Einsteinian. In this case, all the curvature 2-forms are exact and other properties induced by E and P are stated.

The principal curvatures ρ_i are isoparametric functions and the set $(\rho_1, ..., \rho_n)$ defines an isoparametric system [10].

In the last section, we assume that, in addition, M is endowed with an almost symplectic structure. Then, the dual 1-form $\pi = P^{\flat}$ of P is symplectic harmonic. If M is compact, then its 2nd Betti number $b_2 \ge 1$.

Mathematics Subject Classification (2000). 53C25, 53B21, 53D15

1. Preliminaries

Let (M, g) be an *n*-dimensional Riemannian C^{∞} -manifold and ∇ the Levi–Civita connection with respect to *g*. We denote by $\Gamma(TM)$ the set of sections of the tangent bundle *TM* and by $\flat : TM \to T^*M$ the *musical isomorphism* defined by *g* and by $\sharp : T^*M \to TM$ its inverse.

A function $\rho : \mathbf{R}^n \to \mathbf{R}$ is called *isoparametric* [10] if $\|\nabla \rho\|^2$ and div $(\nabla \rho)$ are functions of ρ . Recall $\Delta \rho = -\operatorname{div}(\nabla \rho)$.

More generally, a system of functions $F = (\rho_1, ..., \rho_n)$ such that

(1.1)
$$\langle \nabla \rho_i, \nabla \rho_j \rangle = A_{ij}(F),$$

(1.2)
$$\operatorname{div}(\nabla \rho_i) = B_i(F),$$

(1.3)
$$[\nabla \rho_i, \nabla \rho_j] = \sum_{k=1}^n C_{ij}^k \nabla \rho_k,$$

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 $(i, j, k \in \{1, ..., n\})$, where A_{ij} , B_i and C_{ij}^k are smooth functions, is said to be an *isoparametric system* [10].

Let $\mathcal{O} = \{e_1, ..., e_n\}$ be a local field of adapted vectorial frames over M and let $\mathcal{O}^* = \{\omega^1, ..., \omega^n\}$ be its associated coframe.

Then, the *soldering form dp* (i.e. the canonical vector valued 1-form associated with O) is expressed by

(1.4)
$$dp = \omega^i \otimes e_i$$

and E. Cartan's structure equations, written in indexless form, are

(1.5)
$$\nabla e = \theta \otimes e$$

$$(1.6) d\omega = -\theta \wedge \omega$$

(1.7)
$$d\theta = -\theta \wedge \theta + \Theta.$$

In the above equations, θ and Θ , respectively, are the local connection forms in the tangent bundle *TM* and the curvature 2-forms of *M*.

2. Parallel connection

Let $x : (M, g) \to (\tilde{M}, \tilde{g})$ be the immersion of a C^{∞} -hypersurface M into an (n+1)-dimensional Riemannian manifold \tilde{M} .

Let $\mathcal{B}(\tilde{M})$ be the bundle of orthonormal frames $\mathcal{O} = \{e_A; A = 1, ..., n + 1\}$ over \tilde{M} and $\tilde{\nabla}$ the Levi–Civita connection with respect to \tilde{g} .

We will denote the induced elements on M by suppressing $\tilde{}$ and put $N = e_{n+1}$ for the unit normal vector field of M. Then, the second fundamental form l associated with x is expressed by

(2.1)
$$l = - \langle dp, \nabla N \rangle = \sum_{i=1}^{n} \theta_i^{n+1} \otimes \omega^i.$$

We assume (this is always possible) that *l* is diagonal, i.e. each vector e_i (i = 1, ..., n) is an eigenvector of *l*. Then, we may write

(2.2)
$$\theta_i^{n+1} = \rho_i \omega^i,$$

(no summation over *i*), where $\rho_i \in C^{\infty}(M)$ are the *principal curvatures* of *M*.

In [8], $P_0 = E = \sum_{i=1}^{n} e_i$ and $P_1 = P = \sum_{i=1}^{n} \rho_i e_i$ have been called, respectively, the *distinguished vector* field and the *curvature vector* field associated with x. More generally, $P_r = \sum_{i=1}^{n} \rho_i^r e_i$ is called the *r*-th curvature vector field associated with x, and following [9],

$$h_r = \frac{1}{n} \sum_{i=1}^n \rho_i^r = \frac{1}{n} g(E, P_r)$$

is the *r*-th associated mean curvature of M.

By (2.2) and the structure equation (1.5), one gets

(2.3)
$$\nabla N = -\sum_{i=1}^{n} \rho_i \omega^i \otimes e_i,$$

and it is seen that if all the principal curvatures ρ_i are equal, N is an umbilical section.

Denote by $L: T_p M \to T_p M$ the Weingarten map, i.e. $LZ = \nabla_Z N$ (*L* is linear and self adjoint).

One has

$$\nabla_{P_{r-1}}N=-P_r,$$

which shows that all the curvature vector fields P_r are (up to the sign) related by the Weingarten map.

We assume in this paper that *M* is structured by a *P*-parallel connection, in the sense of [6] and [7]. Consequently, the connection forms θ_i^j satisfy

(2.4)
$$\theta_i^j = \langle P, e_i \wedge e_j \rangle = \rho_i \omega^j - \rho_j \omega^i,$$

where \wedge denotes the wedge product of vector fields.

On behalf of the structure equations (1.6), we easily get, by (2.4), that

$$(2.5) d\omega^i = \pi \wedge \omega^i,$$

where $\pi = P^{\flat}$ denotes the dual form of *P*.

Since all the 1-forms ω^i are exterior recurrent [3], having the same recurrence form π , it follows that π is closed, i.e. $d\pi = 0$, or equivalently,

(2.6)
$$\sum_{i=1}^{n} d\rho_i \wedge \omega^i = 0.$$

On the other hand, by the structure equations (1.5) and (2.4), one gets

(2.7)
$$\nabla \rho_i = \rho_i dp - \omega^i \otimes P + \rho_i \omega^i \otimes N,$$

(no summation over *i*), and since $\nabla_P e_i = \rho_i^2 N$, it follows that all the vector fields e_i , (i = 1, ..., n) are *P*-parallel on *TM*.

Theorem 1. Let $x : M \to \tilde{M}$ be an isometric immersion of a hypersurface M into an (n + 1)-dimensional Riemannian manifold \tilde{M} and let ρ_i (i = 1, ..., n) be the principal curvatures associated with x.

Let E and P, respectively, be the distinguished vector field and the curvature vector field of M.

If M is structured by a P-parallel connection, then M is an Einstein hypersurface of scalar curvature $2(n-1)\rho^2$ ($\rho^2 = ||P||^2 = constant$) and P is an invariant section for the canonical vector-valued 1-form ω_2 of the set of second-order frames on M. In addition, the following properties hold:

- *i) E* and *P* commute and ρ_i represent, up to (n-1), the divergences of the vector basis $\mathcal{O} = \{e_i\}$ on *M* (the indices correspond) and are first integrals of *P*;
- *ii)* all the curvature forms on *M* are exact 2-forms and *P* is a recurrent dynamical system (in the sense of *H*. Poincaré) on *M*;
- iii) the retracting space corresponding to $\Theta_i^j, \omega^i, \omega^j$ is of dimension 3;
- *iv)* the brackets $[e_i, e_j]$ define infinitesimal conformal transformations of the curvature forms Θ_i^j (the indices correspond).

Proof. i) Taking the covariant differential of P, one finds, by (2.7),

(2.8)
$$\nabla P = \sum_{i=1}^{n} d\rho_i \otimes e_i + \rho^2 dp - \pi \otimes P + v \otimes N,$$

where $\rho^2 = \|P\|^2$ and

(2.9)
$$v = \sum_{i=1}^{n} \rho_i^2 \omega^i = P_2^{\flat}.$$

Hence, in order that *P* be also parallel in *TM*, it is necessary and sufficient that the differential of the principal curvatures ρ_i satisfy

(2.10)
$$d\rho_i = \rho_i \pi - \rho^2 \omega^i.$$

One derives that $\rho^2 = \text{constant}$ and

(2.11)
$$\nabla P = v \otimes N.$$

Since, in general, tr $\nabla Z = \text{div } Z$, one derives, by (2.7),

div
$$e_i = \sum_{j=1}^{n} g(\nabla_{e_j} e_i, e_j) = (n-1)\rho_i$$

i.e. the components of *P* represent, up to (n - 1), the divergences of the vector basis $\mathcal{O} = \{e_i\}$.

Taking the covariant differential of the distinguished vector field E, one has, by (2.7),

(2.12)
$$\nabla E = nh_1dp - \omega \otimes P + \pi \otimes N,$$

and remarking that $v(E) = \pi(P) = \rho^2$, we get, by (2.11),

$$(2.13) [E, P] = 0,$$

that is, *E* and *P* commute.

Since, by (2.5), we may write

$$\omega^i \wedge \omega^j \wedge d\omega^i = 0,$$

it follows that all the 2-planes $\{e_i, e_j\}$ are integrable.

In addition, one easily finds $\mathcal{L}_P \rho_i = 0$, i.e. all the principal curvatures ρ_i are *first integrals* of *P*.

ii) Further, taking the exterior differential of the connection forms θ_i^j , one obtains, by (2.5) and (2.10),

(2.14)
$$d\theta_i^j = 2\pi \wedge \theta_i^j - 2\rho^2 \omega^i \wedge \omega^j$$

Now, making use of the structure equations (1.7), after a standard calculation, we get

(2.15)
$$\Theta_i^j = \frac{1}{2} d\theta_i^j.$$

Thus, all the curvature forms are exact.

With respect to the local Riemannian volume element η of M and the local star operator, *, determined by a local orientation of M, one infers, by (2.10),

*div
$$P = (\text{div } P)\eta = \mathcal{L}_P \eta = 0$$
,

where \mathcal{L} denotes the Lie derivative.

Hence, we may say that P defines a *recurrent dynamical system* (Poincaré's theorem of non-return).

We recall that the Ricci tensor field \mathcal{R} and the scalar curvature S are defined by

$$\mathcal{R}(Z, Z') = \sum_{i=1}^{n} g\big(R(e_i, Z)Z', e_i\big),$$

(*R* is the curvature tensor of ∇),

$$S = \sum_{i=1}^{n} \mathcal{R}(e_i, e_i) = \sum_{i=1}^{n} R_{ii},$$

respectively.

By (2.14) and (2.15), one derives $R_{ii} = 2(\rho^2 - \rho_i^2)$. Consequently, the scalar curvature of the hypersurface *M* is given by

$$S = 2(n-1)\rho^2 = \text{constant.}$$

This proves that M is an *Einstein* hypersurface. By (2.4), (2.5), (2.10) and (2.14), one derives, after some calculations,

(2.16)
$$\mathcal{L}_P \omega^i = 0, \quad \mathcal{L}_P \theta^j_i = 0.$$

Then, if $\{e_i^k; i, k = 1, ..., n\}$ is a basis of the Lie algebra $gl(n, \mathbf{R})$ acting on \mathbf{R}^n , the last equations reveal that *P* is an *invariant section* for the canonical vector-valued 1-form

$$\omega_2 = \omega^i \otimes e_i + \theta^i_k \otimes e^k_i$$

of the set of 2-frames $\mathcal{O}^2 M$, i.e. the frames of second order (see also [4]) on M.

iii) We recall the following definition: let \mathcal{I} be an ideal spanned by the linearly independent elements $\omega^i \in \mathcal{O}^*$ and a 2-form $\psi \in \wedge^2 M$. Then, following [2], if q is the smallest integer such that

$$\psi^{q+1} \wedge \omega^1 \wedge ... \wedge \omega^s = 0$$

the retracting space $H(\mathcal{I})$ corresponding to \mathcal{I} is of dimension 2q + s and has the Grassmannian coordinate vector $\phi^q \wedge \omega^1 \wedge ... \wedge \omega^s$.

Coming back to the case under discussion, it follows, on behalf of (2.14) and

(2.15), that the retracting space $H(\mathcal{I})$ corresponding to $\Theta_i^j, \omega^i, \omega^j$ is of dimension 3.

iv) On the other hand, by (2.7), one finds

$$(2.17) \qquad \qquad [e_i, e_j] = \rho_j e_i - \rho_i e_j,$$

and taking account of (2.10) and of the fact that Θ_i^j are closed, one derives

(2.18)
$$\mathcal{L}_{[e_i,e_j]}\Theta_i^j = (\rho_j - \rho_i)\Theta_i^j.$$

Hence, the brackets $[e_i, e_j]$ define infinitesimal conformal transformations of the curvature forms Θ_i^j $(i, j \in \{1, ..., n\})$.

Theorem 2. The vector field $P = \pi^{\sharp}$ is harmonic and all its components ρ_i (the principal curvatures of M) are eigenfunctions of Δ and have, as spectrum, $(n-1)\rho^2$.

Moreover, ρ_i are isoparametric functions and the set $(\rho_1, ..., \rho_n)$ defines an isoparametric system on M.

Proof. Taking the star operator, *, of $\pi = P^{\flat} = \rho_i \omega^i$, one has by a standard formula

(2.19)
$$*\pi = \sum_{i=1}^{n} (-1)^{i-1} \rho_i \omega^1 \wedge \ldots \wedge \hat{\omega}^i \wedge \ldots \wedge \omega^n.$$

Then, using (2.4), one immediatly gets

(2.20)
$$d(*\pi) = 0 \implies \delta\pi = 0,$$

(δ is the adjoint of *d*), which shows that π is a harmonic 1-form, or that *P* is a harmonic vector field.

Further, since, by (2.10), one has

(2.21)
$$\nabla \rho_i = \rho_i P - \rho^2 e_i,$$

and recalling that div $e_i = (n - 1)\rho_i$, one derives

(2.22)
$$\Delta \rho_i = (n-1)\rho^2 \rho_i,$$

which shows that all the principal curvatures ρ_i of *M* are *eigenfunctions* of Δ corresponding to the same *eigenvalue* $(n-1)\rho^2$.

Hence, following a known definition, $(n-1)\rho^2$ is in the *spectrum* of Δ on $\wedge^1 M$ and $\{\rho_i\}$ form the *eigenspace* of the eigenvalue $(n-1)\rho^2$.

Further, one has

$$||\nabla \rho_i||^2 = \rho^2 (\rho^2 - \rho_i^2)$$

Therefore, according to [10], it follows that ρ_i are isoparametric functions.

Moreover, since $\pi = P^{\flat}$, one gets by (2.21) that $\pi(\nabla \rho_i) = 0$, from which we derive

$$<\nabla\rho_i, \nabla\rho_j>=-\rho^2\rho_i\rho_j$$

and

$$[\nabla \rho_i, \nabla \rho_j] = \rho^2 (\rho_i \nabla \rho_j - \rho_j \nabla \rho_i) + \rho_i \rho_j (\rho_i^3 - \rho_j^3) N.$$

Accordingly, by references to (1.1)–(1.3), one may say that the principal curvatures define an isoparametric system on M.

3. Almost symplectic structure

We assume in this section that *M* is of even dimension, say n = 2m, and consider the almost symplectic form

$$\Omega = \sum_{a=1}^{m} \omega^a \wedge \omega^{a^*}, \ a^* = a + m.$$

Taking the exterior differential of Ω , one finds, by (2.5),

$$d\Omega = 2\pi \wedge \Omega.$$

This proves that *M* is endowed with a *conformal symplectic structure* $CSp(m; \mathbf{R})$ having π as a covector of Lee (see also [6], [7]).

Denoting now, as usual, by Ω^{\flat} : $TM \to T^*M$, $Z \mapsto -i_Z \Omega = {}^{\flat}Z$, $Z \in \Gamma(TM)$, the *symplectic isomorphism* and, in order to simplify, setting ${}^{\flat}P = \gamma$, one derives

(3.2)
$$d\gamma = 2\pi \wedge \gamma + 2\rho^2 \Omega.$$

Next, after a short calculation, we find

(3.3)
$$d(\mathcal{L}_P\Omega) = 2\rho^2 \pi \wedge \Omega.$$

Hence, by reference to [6], [7], *P* defines a relative conformal transformation of Ω .

It should also be noted that, by (3.1) and (3.3) and taking account that $\rho^2 =$ constant, one may state that $\mathcal{L}_P \Omega$ and $\rho^2 \Omega$ are homologous.

Therefore, if *M* is compact, then its Betti number of order $2, b_2 \ge 1$.

Since the eigenvalue corresponding to the eigenfunctions ρ_i is positive, we may assume that *M* is compact.

More generally, let $X = \sum_{A=1}^{2m} X^A e_A \in \Gamma(TM)$ be any tangent vector field and denote $\alpha = X^{\flat}$ and $\beta = -{}^{\flat}X$. By (2.7), we have

$$\nabla X = dX^A \otimes e_A + g(X, P)dp - \alpha \otimes P + X^A \rho_A \omega^A \otimes N.$$

If we assume that the differential of the components of X satisfy

$$\begin{cases} dX^a = x_a \omega^{a^*} + t\omega^a, \\ dX^{a^*} = x_{a^*} \omega^a + t\omega^{a^*}, \end{cases}$$

 $x_a, x_{a^*}, t \in C^{\infty}(M)$, then, since

$$\beta = i_X \Omega = \sum_{a=1}^m (X^a \omega^{a^*} - X^{a^*} \omega^a),$$

one derives, by (2.5),

$$d\beta = \pi \wedge \beta + 2t\Omega,$$

which infers $dt + t\pi = 0$, i.e. the existence of such X implies that π is an exact form.

By a straightforward calculation, we find

$$\mathcal{L}_X \Omega = 2(g(P, X) + t)\Omega - \pi \wedge \beta.$$

By exterior differentiation, one gets

$$d(\mathcal{L}_X\Omega) = (df + 2(t+f)\pi) \wedge \Omega, \ f = g(P, X);$$

the above equation shows that the vector field X defines a relative conformal transformation of Ω .

Moreover, since π is an exact form, then, by reference to [1], *X* also defines a *weak automorphism* of Ω .

Further, recall that the Poisson bracket $(u, v)_{\mathcal{P}}$ of two Pfaffians u and v is defined by

$$(u, v)_{\mathcal{P}} = i_{[u^{\sharp}, v^{\sharp}]} \Omega$$

 $[u^{\sharp}, v^{\sharp}]$ is called the associated vector field of $(u, v)_{\mathcal{P}}$.

Coming back to the case under discussion, one finds, by (2.7),

$$(\omega^a, \omega^{a^*})_{\mathcal{P}} = i_{[e_a, e_{a^*}]} \Omega = \rho_a \omega^a + \rho_{a^*} \omega^{a^*},$$

(no summation), and taking the Lie derivative, one derives, by (2.5) and (2.10),

$$\mathcal{L}_{[e_a,e_{a^*}]}\Omega=0.$$

Hence, the structure 2-form Ω is invariant by the vector fields associated with the Poisson brackets $(\omega^a, \omega^{a^*})_{\mathcal{P}}$.

Finally, we recall that if $u \in \wedge^1 M$ is any 1-form and $\tilde{*}$ and $\tilde{\delta}$, respectively, denote the *symplectic adjoint operator* and the *symplectic differentiation operator*, then one has

$$\tilde{*}u = \frac{1}{(2m-1)!}u \wedge \Omega^{2m-1}, \ \tilde{\delta} = \tilde{*}d\tilde{*}u.$$

In the case under discussion, we quickly find $\delta \pi = 0$.

This shows that π is also a *symplectic harmonic* form.

Summing up, we state the following:

Theorem 3. If the Einstein hypersurface M defined in Sect. 2 is endowed with an almost symplectic structure defined by the 2-form Ω , then this structure is necessarily conformal symplectic, having $\pi = P^{\flat}$ as the associated Lee covector. This structure induces the following properties:

- *i)* the curvature vector field P defines a relative conformal transformation of Ω and $\mathcal{L}_P \Omega$ and $\rho^2 \Omega$ are homologous;
- *ii)* if X is any vector field such that its components X^a , X^{a^*} satisfy

$$dX^a = x_a \omega^{a^*} + t\omega^a, \ dX^{a^*} = x_{a^*} \omega^a + t\omega^{a^*},$$

 $(x_a, x_{a^*}, t \in C^{\infty}M, a \in \{1, ..., m\}, a^* = a + m)$, then π is an exact form and X defines a weak automorphism of Ω ;

- iii) Ω is invariant by the associated vector fields of the Poisson brackets $(\omega^a, \omega^{a^*})_{\mathcal{P}}$;
- iv) π is symplectic harmonic.

Acknowledgements. The authors would like to thank the referee for his valuable comments.

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