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ON A FUNCTIONAL EQUATION ARISING IN THE STABILITY THEORY OF DIFFERENCE-DIFFERENTIAL EQUATIONS . W. B./Castelan\* E. F Intan Lefschetz Center for Dynamical Systems Division of Applied Mathematics Brown University Providence, Rhode Island 02912 and **JUN 21** Departamento de Matematica Universidade Federal de Santa Catarina Florianopolis, Brasil NOOD14-76-C-0278L /115F-MPS-71-02923 DISTRIBUTION STATEMENT A and Approved for public release; Distribution Unlimited E. F. Infante\*\* Lefschetz Center for Dynamical Systems Division of Applied Mathematics Brown University

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The functional differential equation Abstract:

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squared

B((Q(tau = t)) to the T-th power)  $Q'(t) = AQ(t) + BQ^{T}(\tau-t), -\omega < t < \omega,$ 

where A,B are  $n \times n$  constant matrices,  $T \ge 0$ , Q(t) is a differentiable  $n \times n$  matrix and  $Q^{T}(t)$  is its transpose, is studied. Existence, uniqueness and an algebraic representation of its solutions is given.

This equation, of considerable interest in its own right, naturally arises in the construction of Liapunov functionals of difference differential equations of the type  $\dot{x}(t) = Cx(t) + tau + Dx(t-t)$ , where C,D are constant  $n \times n$  matrices. The role played by the matrix  $\Omega(t)$  is analogous to the one played by a positive definite matrix in the construction of Liapunov functions for ordinary differential equations.

In this paper, we show that, in spite of the functional nature of this equation, the linear vector space of its solutions is  $n^2$ ; moreover, we give a complete algebraic characterization of its solutions and indicate computationally simple methods for obtaining these solutions, which we illustrate through an example. Finally, we briefly indicate how to obtain solutions for the nonhomogeneous problem, through the usual variation of constants method. ON A FUNCTIONAL EQUATION ARISING IN THE STABILITY THEORY OF DIFFERENCE-DIFFERENTIAL EQUATIONS

### I. Introduction

The study of difference-differential equations has received considerable attention in recent times [2,6,7], the overwhelming interest devoted to equations with positive delays.

In this brief paper we wish to study the matrix functional equation

$$Q'(t) = AQ(t) + BQ^{T}(\tau-t), -\infty < t < \infty$$
 (1.1)

where A,B are constant  $n \times n$  matrices and  $\tau \ge 0$ . This equation is neither of the retarded nor advanced type. We show, that, unlike the infinite dimensionality of the vector space of solutions of functional differential equations, the linear vector space of solutions of this equation is  $n^2$ . Moreover, we give a simple algebraic characterization of these  $n^2$  linearly independent solutions which parallels the one for ordinary differential equations, indicate some methods of computation of these solutions and allude to the variation of constants formula for the nonhomogeneous problem.

This equation, of interest in its own right, is particularly important since it arises naturally in the process of constructing Liapunov functionals for retarded differential equations of the form  $x'(t) = Cx(t) + Dx(t-\tau)$ . Datko [4] has encountered, in a somewhat different form, this equation, but has not studied it. Repin [11], in his construction of Liapunov functionals uses this equation, but does so erroneously in replacing, in (1.1), the term  $BQ^{T}(\tau-t)$  by  $BQ(\tau-t)$ , making the analysis trivial.

This equation has been used by Infante and Walker [9] in the construction of the Liapunov functional for a scalar differencedifferential equation. The study presented here arose in the use of the solutions of this equation in a forthcoming paper [8] which treats the construction of Liapunov functionals for matrix difference-differential equations.

# II. Existence, Uniqueness and Algebraic Structure of the Solutions

Consider the equation

$$Q'(t) = AQ(t) + BQ^{T}(\tau-t), -\infty < t < \infty,$$
 (2.1)

with the condition

$$Q\left(\frac{\tau}{2}\right) = K, \qquad (2.2)$$

where K is an arbitrary  $n \times n$  matrix; this equation is intimately related to the differential equation

$$Q'(t) = AQ(t) + BR(t), \qquad (2.3)$$
$$-\infty < t < \infty,$$
$$R'(t) = -Q(t)B^{T} - R(t)A^{T},$$

with the initial conditions

$$Q(\frac{\tau}{2}) = K, R(\frac{\tau}{2}) = K^{T}.$$
 (2.4)

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Moreover, for any two  $n \times n$  matrices P,S, let the  $n^2 \times n^2$ matrix P  $\otimes$  S denote their Kronecker (or direct) product [1,10] and introduce the notation for the  $n \times n$  matrix

$$s = (s_{ij}) = \begin{pmatrix} s_{1*} \\ \vdots \\ s_{n*} \end{pmatrix} = (s_{*1}, \dots, s_{*n}),$$

where  $s_{i*}$  and  $s_{*j}$  are, respectively, the i<sup>th</sup> row and the j<sup>th</sup> column of S; further, let there correspond, to the n × n matrix S, the n<sup>2</sup>-vector  $s = (s_{1*}, \dots, s_{n*})^{T}$ .

With this notation [1,10], equations (2.3) and (2.4) can be rewritten as

$$\begin{bmatrix} q(t) \\ r(t) \end{bmatrix} = \begin{bmatrix} A \otimes I & B \otimes I \\ -I \otimes B & -I \otimes A \end{bmatrix} \begin{bmatrix} q(t) \\ r(t) \end{bmatrix}, \qquad (2.5)$$

and

$$q(\frac{\tau}{2}) = [k_{1*}, \dots, k_{n*}]^{T}, r(\frac{\tau}{2}) = [k_{*1}^{T}, \dots, k_{*n}^{T}]^{T},$$
 (2.6)

which, with the obvious correspondence and for simplicity of notation, we denote as

$$p(t) = Cp(t),$$
 (2.7)

$$p(\frac{\tau}{2}) = p_{\tau} \qquad (2.8)$$

Here, p(t) is an  $2n^2$ -vector and C is a  $2n^2 \times 2n^2$  constant matrix.

The use of the Kronecker product, which has allowed us to reduce (2.3) - (2.4) to (2.7) - (2.3) permits us to prove our first result.

<u>Theorem 1</u>: Equation 2.1 with condition (2.2) has a unique solution Q(t) for  $-\infty < t < \infty$ .

<u>Proof</u>: If equation (2.1) with condition (2.2) has a differentiable solution Q(t), then, defining  $R(t) = Q^{T}(t-\tau)$ , the pair of matrices Q(t), R(t) will satisfy equations (2.3) and (2.4); hence, with the notation introduced above, the pair of vectors q(t)and r(t) will satisfy equations (2.5), (2.6). These remarks, the linearity of all the involved equations, and the uniqueness of the solutions of (2.5) - (2.6) immediately implies that if a solution Q(t) exist, it is unique.

On the other hand, (2.5) - (2.6) has a unique solution defined for  $-\infty < t < \infty$ , and this implies the existence of a unique pair of differentiable matrices Q(t) and R(t) defined for  $-\infty < t < \infty$ and satisfying (2.3) - (2.4). But these last equations can be rewritten as

$$\frac{d}{dt} Q(t) = AQ(t) + BR(t), \qquad (2.9)$$

$$\frac{d}{dt} R^{T}(\tau-t) = AR^{T}(\tau-t) + BQ^{T}(\tau-t),$$

with initial condition

$$Q(\frac{T}{2}) = K = R^{T}(\frac{T}{2}),$$
 (2.10)

from which it follows, from uniqueness, that  $R(t) = Q^{T}(\tau-t)$ , completing the proof.

Examination of the above proof makes it clear that knowledge of the solution of (2.5) - (2.6) immediately yields the solution of (2.1) - (2.2). But (2.5) - (2.6) is a standard initial value problem in ordinary differential equations; the structure of the solutions of such problems is well known [3,5]. Moreover, since the  $2n^2 \times 2n^2$  matrix C has a very special structure, it should be possible to recover the structure of the solutions of equation 2.1.

Let us consider, for the moment, the solutions of equation (2.5). Recall [3,5] that it has  $2n^2$  linearly independent solutions which can be obtained in the following fashion. Let  $\lambda_1, \dots, \lambda_p$ ,  $p = 2n^2$ , be the distinct eigenvalues of the matrix C, that is, solutions of the determinental equation

 $det[\lambda I-C] = det \begin{bmatrix} (\lambda I-A) \otimes I & -B \otimes I \\ I \otimes B & I \otimes (\lambda I+A) \end{bmatrix} = 0; \quad (2.11)$ 

each  $\lambda_j$ ,  $j = 1, \dots, p$  with algebraic multiplicity  $m_j$  and geometric multiplicities  $n_j^r$ ,  $\sum_{r=1}^{s} n_j^r = m_j$ ,  $\sum_{j=1}^{r} m_j = 2n^2$ . Then,  $2n^2$  linearly independent solutions of (2.5) (or (2.7)) are given by

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$$\phi_{j,r}^{q}(t) = e^{\lambda_{j}(t - \frac{\tau}{2})} \sum_{\substack{i=1 \\ i=1}}^{q} \frac{(t - \frac{\tau}{2})^{q-i}}{(q-i)!} e^{i}_{j,r}, \qquad (2.12)$$

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where  $q = 1, ..., n_j^r$ , r = 1, ..., s,  $\sum_{r=1}^{s} n_j^r = m_j$ ,  $\sum_{j=1}^{r} m_j = 2n^2$ , and the  $2n^2$  linearly independent eigenvectors and generalized eigenvectors are given by

$$[\lambda_{j}I-C]e_{j,r}^{i} = -e_{j,r}^{i-1}, \quad e_{j,s}^{0} = 0.$$
 (2.13)

A change of notation, and a return from the vector to the matrix form, shows that  $2n^2$  linearly independent solutions of (2.3) are given by

for  $q = 1, ..., n_j^r$ , r = 1, ..., s, j = 1, ..., p,  $\sum_{n=1}^{s} n_j^r = m_j$ ,  $\sum_{j=1}^{j} m_j = 2n^2$ , and where the generalized eigenmatrix pair  $(L_{j,r}^i, M_{j,r}^i)$  associated with the eigenvalue  $\lambda_j$ , satisfy the equations

$$(\lambda_{j}I-A)L_{j,r}^{i} - BM_{j,r}^{i} = -L_{j,r}^{i-1} ,$$

$$L_{j,r}^{i}B^{T} + M_{j,r}^{i}(\lambda_{j}I+A^{T}) = -M_{j,r}^{i-1} ,$$

$$(2.15)$$

for 
$$i = 1, ..., n_j^r$$
,  $r = 1, ..., s$ ,  $L_{j,s}^0 = M_{j,s}^0 = 0$ .

The structure of these equations is a most particular one; indeed, if they are multiplied by -1, transposed, and written in reverse order, they yield

$$(-\lambda_{j}I-A)M_{j,r}^{i^{T}} - BL_{j,r}^{i^{T}} = M_{j,r}^{i-1^{T}},$$

$$M_{j,r}^{i^{T}}B^{i^{T}} + L_{j,r}^{i^{T}}(-\lambda_{j}I+A^{T}) = L_{j,r}^{i-1^{T}},$$
(2.16)

for  $i = 1, ..., n_j^r$ , r = 1, ..., s,  $L_{j,r}^{0^T} = M_{j,r}^{0^T} = 0$ . But this result demonstrates that if  $\lambda_j$  is a solution of (2.11),  $-\lambda_j$  will also be a solution; moreover  $\lambda_j$  and  $-\lambda_j$  have the same geometric multiplicities and the same algebraic multiplicity. Hence, the distinct eigenvalues always appear in pairs  $(\lambda_j, -\lambda_j)$ , and an examination of equations (2.15) and (2.16) shows that if the generalized eigenmatrix pairs corresponding to  $\lambda_j$  are  $(L_{j,r}^i, M_{j,r}^j)$ , the generalized eigenmatrix pairs corresponding to  $-\lambda_j$  will be  $((-1)^{i+1}M_{j,r}^{i^T}, (-1)^{i+1}L_{j,r}^{i^T})$ .

But these remarks imply that if the solution (2.12) corresponding to  $\lambda_j$  is added to the solution (2.12) corresponding to  $-\lambda_j$ multiplied by (-1)<sup>q+1</sup>, the n<sup>2</sup> linearly independent solutions of (2.3) given by

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$$\begin{bmatrix} \Xi_{j,s}^{q}(t) \\ \Pi_{j,s}^{q}(t) \end{bmatrix} = e^{\lambda_{j}(t - \frac{T}{2})} \underbrace{\substack{q \\ i=1}}_{i=1}^{q} \frac{(t - \frac{T}{2})^{q-i}}{(q-i)!} \begin{bmatrix} L_{j,r}^{i} \\ M_{j,r}^{i} \end{bmatrix}$$
$$+ e^{-\lambda_{j}(t - \frac{T}{2})} \underbrace{\substack{q \\ i=1}}_{i=1}^{q} \frac{(t - \frac{T}{2})^{q-i}}{(q-i)!} (-1)^{q+i} \begin{bmatrix} M_{j,r}^{iT} \\ L_{j,r}^{iT} \end{bmatrix}$$

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satisfy the condition  $\Xi_{j,r}^{q}(\frac{\tau}{2}) = \Pi_{j,r}^{qT}(\frac{\tau}{2})$ . But this is precisely condition (2.4); it therefore follows that

$$\Xi_{j,r}^{q}(t) = \sum_{i=1}^{q} \frac{(t - \frac{\tau}{2})^{q-i}}{(q-i)!} [e^{\lambda_{j}(t - \frac{\tau}{2})} L_{j,r}^{i} + (-1)^{q+i} e^{-\lambda_{j}(t - \frac{\tau}{2})} M_{j,r}^{iT}], \quad (2.17)$$

for  $q = 1, ..., n_j^r$ ,  $\sum_{r=1}^{s} n_j^r = m_j$ ,  $\sum_{2j} m_j = 2n^2$ , are  $n^2$  linearly independent solutions of (2.1). We have, therefore proven

<u>Theorem 2</u>: Equation 2.1 has  $n^2$  linearly independent solutions given by equation 2.17, where the generalized eigenmatrix pairs  $(L_{j,s}^i, M_{j,s}^i)$ satisfy equation 2.17 for one of the elements of the pair  $(\lambda_j, -\lambda_j)$ , each of which is a solution of equation 2.11.

It is interesting to remark that the determinental equation (2.11), involving a  $2n^2 \times 2n^2$  determinant, given the commutativity of its elements, can always be rewritten as

$$det[(\lambda I-A) \otimes (\lambda I+A) + B \otimes B] = 0.$$
(2.18)

Theorem 2 gives the desired algebraic representation of the solutions, a representation which is completely analogous to that for ordinary differential equations. It is surprising that the vector space of solutions of (2.1) has dimension  $n^2$ .

### III. Some Further Characterizations

Theorem 2 of the previous section gives a complete characterization of the solutions of our original functional equation. It is possible, however, to give some further properties of the eigenmatrix pairs in certain particular cases; these further characterizations of the eigenmatrix pairs are very useful from a computational viewpoint, as we demonstrate in the next section.

<u>Lemma 1</u>: If, for a  $\lambda_j$  satisfying equation 2.18, there exists an  $\alpha_j \neq 0$  and corresponding vectors  $x_j, y_j$  such that

$$[\lambda_{j}I - A - \alpha_{j}B]x_{j} = 0,$$

$$[\lambda_{j}I + A + \frac{1}{\alpha_{j}}B]y_{j} = 0,$$
(3.1)

then equation 2.15 is satisfied for i = 1 with  $L_{j,r}^{1} = x_{j}y_{j}^{T}$ ,  $M_{m,r}^{1} = \alpha_{j}L_{j,r}^{1}$ .

<u>Proof</u>: A simple substitution of the indicated result into equation (2.15) shows that this equation is satisfied.

The proofs of the following two lemmas are equally obvious.

Lemma 2: If, for a  $\lambda_j$  satisfying equation 2.18, there exist vectors  $x_j, y_j$  such that

$$Bx_{j} = 0$$
(3.2)
  
 $j^{I+A}y_{j} = 0$ 

then equation (2.15) is satisfied for i = 1 with  $L_{j,r}^{l} = 0$ ,  $M_{j,r}^{l} = x_{j}y_{j}^{T}$ .

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Lemma 3: If, for a  $\lambda_j$  satisfying Equation 2.18 there exist vectors  $x_j, y_j$  such that

$$\lambda_{j} \mathbf{I} - \mathbf{A} \mathbf{x}_{j} = 0$$

$$\mathbf{B} \mathbf{y}_{j} = 0$$
(3.3)

then equation 2.15 is satisfied for i = 1 with

$$L_{j,r}^{1} = x_{j}y_{j}^{T}$$
,  $M_{j,r}^{1} = 0$ .

It is noted that, if the assumptions of any of these three lemmas hold, then the form of our eigenmatrix pairs is dyadic. Moreover, determination of the  $x_j$  depends on simultaneous solutions of the determinental equations, for a given  $\lambda_j$  satisfying (2.18),

$$det[\lambda_j I - \Lambda - \alpha_j B] = 0,$$

$$det[\lambda_j I + \Lambda + \frac{1}{\alpha_j} B] = 0,$$
(3.4)

where, for Lemma 2 we let  $\frac{1}{\alpha_j} = 0$  and, for Lemma 3,  $\alpha_j = 0$ . The computations involved for the determination of  $\alpha_j, x_j$  and  $y_j$  are much simpler than those implied by equation (2.15).

If B and either  $(\lambda_j I - A)$  or  $(\lambda_j I + A)$  are not invertible, then Lemmas 2 and 3 apply. It is not difficult to show that, if such is not the case, there always exists at least one  $\alpha_j \neq 0$ such that (3.4) is satisfied, which implies the applicability of Lemma 1.

Lemma 4: If either B is invertible or  $(\lambda_j I-A)$  and  $(\lambda_j I+A)$ are invertible, then there exist an  $\alpha_j \neq 0$  which satisfies (2.22).

<u>Proof</u>: If B is invertible, for a given j,r and i = 1, equations (2.15) are equivalent to

$$B^{-1}(\lambda_{j}I-A)L_{j,r}^{1}(\lambda_{j}I+A)^{T}B^{T}^{-1} = -L_{j,r}^{1}$$

$$B^{-1}(\lambda_{j}I-A)M_{j,r}^{1}(\lambda_{j}I+A)^{T}B^{T}^{-1} = -M_{j,r}^{1}.$$
(3.5)

But this implies that the matrix  $[B^{-1}(\lambda_j I - A)] \otimes [B^{-1}(\lambda_j I + A)]$ has at least one eigenvalue that equals -1. Since [10] the eigenvalues of a Kronecker product are the products of the eigenvalues of the two matrices appearing in the product, then there exist an  $\alpha_j \neq 0$  and vectors  $x_j$  and  $y_j$  such that

$$B^{-1}(\lambda_{j}I-A)x_{j} = \alpha_{j}x_{j},$$

$$B^{-1}(\lambda_{j}I+A)y_{j} = -\frac{1}{\alpha_{j}}y_{j},$$
(3)

6)

which is equivalent to the assumptions of Lemma 1.

If B is not invertible but  $(\lambda_j I-A)$  and  $(\lambda_j I+A)$  are both invertible, then (2.15), for i = 1, are equivalent to

$$(\lambda_{j}I-A)^{-1}B L_{j,r}B^{T}(\lambda_{j}I+A^{T})^{-1} = -L_{j,r}^{1}$$
  
 $(\lambda_{j}I-A)^{-1}B M_{j,r}^{1}B^{T}(\lambda_{j}I+A^{T})^{-1} = -M_{j,r}^{1}$ 

and a repetition of the above argument leads to the same conclusion. These last four Lemmas imply that, associated with each distinct eigenvalue pair  $(\lambda_j, -\lambda_j)$ , a solution (2.17) to our equation exists in which  $L_{j,r}^1$  and  $M_{j,r}^1$  are for some r dyadic and linearly dependent. It is also worth remarking that Lemmas 1-3 are not exclusive, and that more than one set  $(\alpha_j, x_j, y_j)$  could, in some cases, be obtained from the appropriate equations, yielding a similar result for other values of r.

### IV. An Example

Here we present, as an illustration of the computations involved, the construction of the linearly independent solutions for a  $2 \times 2$  systems of the form

$$Q'(t) = AQ(t) + BQ^{T}(\tau-t), -\infty < t < \infty$$
 (4.1)

where, for illustrative purposes, we have chosen

$$A = \begin{bmatrix} -5 & 0 \\ 0 & -2 \end{bmatrix}, \qquad B = \begin{bmatrix} 4 & 0 \\ 2 & 0 \end{bmatrix}.$$
(4.2)

For the computation of the eigenvalues, the determinental equation (2.18) gives

$$det \begin{vmatrix} (\lambda I - A) & 0 & I & -B & 0 & I \\ & & & \\ I & 0 & B & I & 0 & (\lambda I + A) \end{vmatrix} = det[(\lambda I - A) & 0 & (\lambda I + A) + B & 0 & B] =$$

$$= \det \begin{bmatrix} (\lambda+5)(\lambda-5) + 16 & 0 & 0 & 0 \\ 8 & (\lambda+5)(\lambda-2) & 0 & 0 \\ 8 & 0 & (\lambda+2)(\lambda-5) & 0 \\ 4 & 0 & 0 & (\lambda+2)(\lambda-2) \end{bmatrix} = 0,$$

from which we obtain the four pairs of roots  $(\lambda_1, -\lambda_1) = (5, -5)$ ,  $(\lambda_2, -\lambda_2) = (3, -3)$ ,  $(\lambda_3, -\lambda_3) = (2, -2)$ ,  $(\lambda_4, -\lambda_4) = (2, -2) = (\lambda_3, -\lambda_3)$ .

For the first pair, note that  $\lambda_1 = 5$  is an eigenvalue of -A and that B is not invertible. Application of Lemma 2, equation (3.2) yields that

Bx<sub>1</sub> = 0, or 
$$\begin{bmatrix} 4 & 0 \\ 2 & 0 \end{bmatrix}$$
x<sub>1</sub> = 0,  
( $\lambda_1$ I+A)y<sub>1</sub> = 0, or  $\begin{bmatrix} 0 & 0 \\ 0 & 3 \end{bmatrix}$ y<sub>1</sub> = 0;

hence we have  $L_{1,1}^{1} = 0$  and  $M_{1,1}^{1} = \begin{bmatrix} 0\\1 \end{bmatrix}$  [1,0] yielding the solution

$$\Xi_{1,1}^{1}(t) = e^{-5(t-3)} \begin{bmatrix} 0 & 0 \\ 0 \\ 1 & 0 \end{bmatrix}.$$
 (4.3)

The second pair,  $(\lambda_2, -\lambda_2) = (3, -3)$ , does not consist of eigenvalues of either A or -A. Hence, we search for an  $\alpha$  such that equations (3.4) are satisfied, that is

$$\det \begin{bmatrix} 3 + 5 - \alpha_2 4 & 0 \\ & & \\ -2\alpha_2 & 3 + 2 \end{bmatrix} = 0,$$

and

det 
$$\begin{bmatrix} 3 - 5 + \frac{1}{\alpha_2} & 4 & 0 \\ & & & \\ & + \frac{1}{\alpha_2} & 2 & 3 - 2 \end{bmatrix} = 0,$$

which yields the value  $\alpha_2 = 2$ . Equations (3.1) with  $\lambda_j = 3$  and  $\alpha_j = 2$  now yield

$$\mathbf{x}_2 = \begin{bmatrix} 5 \\ 4 \end{bmatrix}$$
,  $\mathbf{y}_2 = \begin{bmatrix} 1 \\ -1 \end{bmatrix}$ 

and, from our first lemma we obtain the solution

$$\mathbb{E}_{2,1}^{1}(t) = e^{3(t-3)} \begin{bmatrix} 5 & 5 \\ -4 \\ 4 & -4 \end{bmatrix} + 2e^{-3(t-3)} \begin{bmatrix} 5 & 4 \\ -5 & -4 \end{bmatrix}.$$
 (4.4)

The last pair of eigenvalues is of algebraic multiplicity two but it is easy to check that they are of geometric multiplicity one, hence we can attempt to treat them once again through our lemmas.

In the first case,  $\lambda_3 = 2$  is an eigenvalue of -A. Then equations (3.2) become

$$Bx_{3,1} = 0, \text{ or } \begin{bmatrix} 4 & 0 \\ 2 & 2 \end{bmatrix} x_{3,1} = 0,$$
$$\lambda_{3}^{(1+A)}y_{3,1} = 0, \text{ or } \begin{bmatrix} -3 & 0 \\ 0 & 0 \end{bmatrix} y_{3,1} = 0,$$

yielding  $L_{3,1}^1 = 0$  and  $M_{3,1}^1 = \begin{bmatrix} 0 \\ 1 \end{bmatrix}$  [0,1], and therefore, the solution

$$E_{3,1}^{1}(t) = e^{-2(t-s)} \begin{bmatrix} 0 & 0 \\ 0 & 1 \end{bmatrix}.$$
 (4.5)

In the second case,  $\lambda_3 = +2$  is an eigenvalue of A, still, equations (3.4) are satisfied for an  $\alpha_3 \neq 0$ ; indeed, we require

$$det[\lambda_{3}I-A-\alpha_{3}B] = det\begin{bmatrix} 2+5-\alpha_{3}4 & 0\\ & & \\ & -\alpha_{3}2 & 2+2 \end{bmatrix} = 0,$$

$$det[\lambda_{3}I + A + \frac{1}{\alpha_{3}}B] = det\begin{bmatrix}2 - 5 + \frac{4}{\alpha_{3}} & 0\\ & & \\ \frac{2}{\alpha_{3}} & 2 - 2\end{bmatrix} = 0,$$

which are clearly satisfied for

$$\alpha_3=\frac{7}{4}$$
 .

The corresponding vectors are immediately computed as

$$\mathbf{x}_{3,2} = \begin{bmatrix} 8 \\ 7 \end{bmatrix}$$
,  $\mathbf{y}_{3,2} = \begin{bmatrix} 0 \\ 1 \end{bmatrix}$ ,

yielding, through Lemma 1, the solution

$$E_{3,2}^{1}(t) = e^{2(t-3)} \begin{bmatrix} 0 & 8 \\ 0 & 7 \end{bmatrix} + \frac{7}{4} \begin{bmatrix} 0 & 0 \\ 8 & 7 \end{bmatrix} e^{-2(t-3)}$$
(4.6)

It is easily seen that these four solutions are linearly independent.

## V. The Nonhomogeneous Problem

We briefly consider the nonhomogeneous problem

$$Q'(t) = AQ(t) + BQ^{T}(\tau-t) + F(t), -\infty < t < \infty$$
 (5.1)

where A, B are constant  $n \times n$  matrices and F(t) is a continuous

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n × n matrix. We seek a particular solution of this problem.

For simplicity of notation, let the  $n^2$  linearly independent solutions given by (2.17) be relabeled as  $z_k(t) = (z_{i,j}^k(t))$ ,  $k = 1, ..., n^2$ . Then, in a manner completely analogous to that for ordinary differential equations, we obtain

Theorem 3: Equation 5.1 has a particular solution given by

$$\Omega(t) = \sum_{k=1}^{n^2} \left[ \int_{\frac{\tau}{2}}^{t} r_k(s) ds \right] z^k(t), \qquad (5.2)$$

where the  $r_k(s)$  are given below.

<u>Proof</u>: Substitution of (5.2) into (5.1), given that the  $z^{k}(t)$  are solutions of the homogeneous equation, yields that the  $r_{k}(s)$  must satisfy the equation

$$\sum_{k=1}^{n^2} r_k(t) z^k(t) = F(t).$$
 (5.3)

In the notation of Section 2, define the n<sup>2</sup>-vectors  $\zeta_{k}(t) = (z(t)_{1*}^{k}, \dots, z_{n*}^{k}(t))^{T}$ ,  $r(t) = (r_{1}(t), \dots, r_{2}(t)^{T}$ ,  $f(t) = (f_{1*}, \dots, f_{n*})^{T}$ , and the  $n^{2} \times n^{2}$  matrix  $\textcircled{(t)} (t) = (\zeta_{1}, \dots, \zeta_{n^{2}})$ . Then (5.3) is equivalent to

$$\Theta(t)r(t) = f(t).$$

(B(t) is clearly nonvanishing for all t, given the linear independence of the solutions  $z^{k}(t)$ ; hence

$$r(t) = \Theta^{-1}(t)f(t)$$

and this concludes the proof.

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