

On quantum symmetries of ADE graphs

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Abstract

The double triangle algebra (DTA) associated to an ADE graph is considered. A description of its bialgebra structure based on a reconstruction approach is given. This approach takes as initial data the representation theory of the DTA as given by Ocneanu's cell calculus. It is also proved that the resulting DTA has the structure of a weak $*$ -Hopf algebra. As an illustrative example, the case of the graph A_3 is described in detail.

1 Introduction

This paper deals with the correspondence between rational conformal field theories (RCFT) of $SU(2)$ -type and ADE graphs [4, 5, 22]. More precisely we focus on the construction of the so-called double triangle algebra (DTA) [20] associated to an ADE graph. The DTA is a bialgebra and the “algebra of quantum symmetries” describes the tensor product of representations associated with one of its product structures (the same name sometimes denotes the bialgebra itself). Aside from its interest as a mathematical structure, the motivation in considering the DTA stems from the fact that its knowledge makes it possible to construct the modular invariant partition function, as well as other objects, in particular the so-called twisted partition functions [6, 7, 8, 23, 25] associated with the corresponding RCFT – actually [23] one has to take into account existence of boundaries and defects.

The present work contributes to the understanding of the DTA in two respects. First, it provides a precise description of the DTA. To the knowledge of the authors such a description is not available in the literature. Furthermore our approach is constructive in the sense that, given an ADE graph, one can construct the corresponding DTA¹. This is done starting from the calculus of Ocneanu’s cells and connections. More precisely, the approach we employ amounts essentially to take the above mentioned calculus as describing the representation theory of an algebraic structure to be found: the DTA. In this paper², the DTA has a product called \cdot which is determined from the cell calculus. It has also a coproduct Δ . This coproduct determines a product on the dual \widehat{DTA} that corresponds to the composition of endomorphisms. In other references this later product is called “composition product” whereas the product \cdot that we study here is called “convolution product”. We decided to focus the present paper on the convolution product since this operation is the non-trivial one³: the other operation is simply the composition of endomorphisms stemming from the definition of the underlying vector space structure of the DTA (this vector space is defined as $End^{gr}(\mathcal{E})$, the graded vector space \mathcal{E} will be defined later).

The second subject we address is the assertion that the DTA has the structure of a weak $*$ -Hopf algebra (WHA) [3]. This assertion is not new, it is given in ref.[23] that gives arguments based on considering solutions of the

¹The main computational effort is to compute the connections associated with the corresponding Ocneanu cell systems.

²A very simple example is analysed in [9] but notations are not the same (and the point of view is quite different).

³These properties could be discussed in terms of nets of subfactors [1, 2], a notion that we do not use here.

so-called [3, 19] "big pentagon equation" (BPE). In contrast to ref.[23] we do not assume any a priori knowledge of such a solution. In our work, each structural map of the corresponding WHA is constructed in terms of the available data, i.e. connections on cell systems, and all the WHA properties of these maps are proved using properties of the cell calculus (which was introduced in [18] and described for instance in [11] and [24]). In order to make contact with ref.[23], it would be desirable to establish a precise connection between the present results and solutions of the BPE .

The paper is organized as follows. Section II gives generalities about graphs and specify the requirements we want the DTA to fulfill. These requirements deal mainly with the representation theory of a "searched for" C^* -algebra, that we call DTA. The remaining sections uncover the structure maps of the DTA out of the data given in section II. Section III deals with the algebra structure taking advantage of the fact that the DTA is a finite dimensional C^* -algebra. Section IV gives in terms of connections, what we may call "the weak bialgebra structure maps of the DTA", that is, product, coproduct and counit in terms of connections. Section V considers the antipode. The main sections are supplemented by four appendices. All the general results are exemplified in detail for the case of the graph A_3 .

2 The double triangle algebra

2.1 Preliminaries

Let us consider a graph G with n_v vertices. One can characterize a graph by its adjacency matrix M . Its size is $n_v \times n_v$ and its (v_1, v_2) matrix element is an integer n if vertex v_1 is connected to vertex v_2 by n edges. The normalized⁴ eigenvector with maximum eigenvalue β of the adjacency matrix M is called the Perron-Frobenius eigenvector and its components will be denoted by μ_{v_i} , $i = 1, \dots, n_v$.

We can define over G a vector space \mathcal{P} whose elements are paths. An elementary path of length n is a ordered n -uple of contiguous vertices in G . Two vertices are contiguous if there exists an edge connecting them. A path is a linear combination over \mathbb{C} of elementary paths. Therefore these elementary paths provide a preferred basis of \mathcal{P} .

This vector space is graded by the length of paths. There is a subspace \mathcal{E} of \mathcal{P} given by "essential paths", that is paths that are annihilated by all

⁴Set a smallest component to be equal to 1.

the Ocneanu's operators c_k , $k \in \mathbb{N}$. The operator c_k acting on a path of length $n \leq k$ gives zero, otherwise ($n > k$) it is given by,

$$c_k(v_0, v_1, \dots, v_{k-1}, v_k, v_{k+1}, \dots, v_n) = \sqrt{\frac{\mu_{v_k}}{\mu_{v_{k+1}}}} \delta_{v_{k-1}, v_{k+1}}(v_0, v_1, \dots, v_{k-1}, v_{k+2}, \dots, v_n) \quad (2.1)$$

where $(v_0, v_1, v_2, \dots, v_{k-1}, v_k, v_{k+1}, \dots, v_n)$ denotes an elementary path of length n passing through the vertices v_0 to v_n of G . That is the operator c_k eliminates any one-step backtracking sub-path that starts at step k of the path to which the operator is applied and multiplies the result by a number given in terms of the components of the Perron-Frobenius eigenvector.

There is a natural product in \mathcal{P} defined by concatenation of paths. The concatenation product of two paths is zero if the ending vertex of the first path is not equal to the starting vertex of the second path. If the above holds then the product path is simply the extension of the first path by the second. In symbols take $\xi_i = (v_0^i, \dots, v_n^i)$ and $\xi_j = (v_0^j, \dots, v_m^j)$ then the concatenation product $\xi_i \star \xi_j$ of ξ_i and ξ_j is given by,

$$\xi_i \star \xi_j = \delta_{v_n^i v_0^j}(v_0^i, \dots, v_n^i, v_1^j, \dots, v_m^j) \quad . \quad (2.2)$$

We shall restrict to graphs where the dimension of \mathcal{E} is finite. This restriction is very strong and essentially⁵ reduces the family of admissible graphs to those belonging to the ADE series [20].

The basis of elementary paths restricted to the maximum length of essential path will be denoted by $\{\xi_i\}$. We are interested in the length preserving endomorphisms of \mathcal{E} that we denote $End^{gr}(\mathcal{E})$. We denote the dual vector space of \mathcal{E} by $\hat{\mathcal{E}}$ and by $\{\xi^i\}$ the dual basis to the $\{\xi_i\}$. Hence a basis of $End^{gr}(\mathcal{E})$ is given by the objects $\{\xi_i \otimes \xi^j\}$.

Example 2.1 (The case of A_3). *The graph $G = A_3$ and its corresponding adjacency matrix M are,*

$$\begin{array}{c} 0 \quad 1 \quad 2 \\ \bullet \text{---} \bullet \text{---} \bullet \end{array} \quad , \quad M = \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix} \quad (2.3)$$

⁵One could be more accurate, but such a discussion is not needed for our present purpose. Notice that an extension of the results described here to other types of graphs, for example to graphs belonging to higher Coxeter-Dynkin systems related to $SU(N)$ -type RCFTs [12, 13, 21, 26, 27] requires an appropriate modification of the definition of essential paths.

where rows and columns are ordered as $0, 1, 2$ (the values for vertex v can be $0, 1$ or 2). The maximum eigenvalue is $\beta = \sqrt{2}$ and the Perron-Frobenius eigenvector is $(1, \sqrt{2}, 1)$. There are ten essential paths in this graph. Denoting them by the corresponding succession of vertices, they are the following,

$$\begin{aligned} (0) &= 0 & (1) &= 1 & (2) &= 2 \\ (01) &= r_0 & (12) &= r_1 & (10) &= l_1 & (20) &= l_2 \\ (012) &= d & (210) &= g & \frac{(121)-(101)}{\sqrt{2}} &= \gamma \end{aligned} \tag{2.4}$$

where we have also included a shorthand notation. From (2.4) we see that the dimension of $End^{gr}(\mathcal{E})$ is $3^2 + 4^2 + 3^2 = 34$.

2.2 The double triangle algebra(DTA)

As mentioned in the introduction the strategy adopted in this paper is to reconstruct an algebraic structure out of certain requirements that we want to be fulfilled. They will involve information about the representation theory but not directly about the product structure. Indeed one of the results of this analysis will be the product \cdot for this algebra. Requirements:

1. Vector space structure: It is given by $End^{gr}(\mathcal{E})$.

We emphasize that this requirement is only relative to the choice of the underlying vector space structure and not the product itself since the product law that we study on the vector space $End^{gr}(\mathcal{E})$ is (*not*) the composition of endomorphisms but another product that we call \cdot . Rather than giving this product explicitly we shall obtain it by describing all its representations.

2. It is a C^* -algebra⁶.
3. Fundamental ⁷irreps of the product \cdot : These are matrix $*$ -representations

⁶Being a C^* -algebra it is semisimple, and is therefore graded for this product \cdot (simple blocks) but this grading differs from the one given by the length of paths in $End^{gr}(\mathcal{E})$.

⁷Fundamental in the sense that they generate any other irrep by taking adequate linear combinations of tensor products of them. See item (4) for the definition of tensor product representations.

⁸ f whose matrix element $\alpha\beta$ is given by ⁹,

$$\Phi_{\alpha\beta}^f(\xi \otimes \xi') = \alpha \begin{array}{c} \xrightarrow{\xi} \\ \downarrow f \\ \xrightarrow{\xi'} \end{array} \beta \quad (2.5)$$

where the matrix indices α, β label length-one paths on G . If it does not happen that,

$$s(\alpha) = s(\xi) \ , r(\alpha) = s(\xi') \ , s(\beta) = r(\xi) \ , r(\beta) = r(\xi') \quad (2.6)$$

then (2.5) vanishes¹⁰. The symbol on the r.h.s. of (2.5), is a complex number and is by definition the connection associated to that cell in a fundamental representation. These values should satisfy the following conditions,

(0) Zero length paths:

If $\#\xi = \#\xi' = 0$ then,

$$\alpha \begin{array}{c} \xrightarrow{\xi} \\ \downarrow f \\ \xrightarrow{\xi'} \end{array} \alpha = \delta_{s(\alpha)\xi} \delta_{r(\alpha)\xi'} \quad (2.7)$$

(i) Unitarity¹¹:

$$\sum_{\alpha, \xi'} \overline{\alpha \begin{array}{c} \xrightarrow{\xi} \\ \downarrow f \\ \xrightarrow{\xi'} \end{array} \beta} \alpha \begin{array}{c} \xrightarrow{\lambda} \\ \downarrow f \\ \xrightarrow{\xi'} \end{array} \eta = \delta_{\beta\eta} \delta_{\xi\lambda} \quad (2.9)$$

⁸By this we mean that the star operation in the DTA corresponds to hermitian conjugation of matrices in these representations

⁹In the drawing (2.5) the labels for horizontal paths ξ and ξ' denote generic elements of \mathcal{E} and $\hat{\mathcal{E}}$ respectively.

¹⁰In (2.6) we have denoted by $s(\alpha)$ the starting vertex of path α and by $r(\alpha)$ its final vertex.

¹¹Using (2.9) and (2.10) the following relation is obtained,

$$\sum_{\beta, \xi} \overline{\alpha \begin{array}{c} \xrightarrow{\xi} \\ \downarrow f \\ \xrightarrow{\xi'} \end{array} \beta} \alpha' \begin{array}{c} \xrightarrow{\xi} \\ \downarrow f \\ \xrightarrow{\lambda} \end{array} \beta = \delta_{\alpha\alpha'} \delta_{\xi'\lambda} \quad (2.8)$$

conditions (2.9) and (2.8) are called "bi"unitarity in ref.[20].

(ii) Reflection:

$$\begin{array}{c} \xi \\ \bullet \xrightarrow{n} \bullet \\ \alpha \downarrow f \downarrow \beta \\ \bullet \xrightarrow{\xi'} \bullet \end{array} = \sqrt{\frac{\mu_f^\xi \mu_i^{\xi'}}{\mu_i^\xi \mu_f^{\xi'}}} \overline{\begin{array}{c} \tilde{\xi} \\ \bullet \xrightarrow{n} \bullet \\ \beta \downarrow f \downarrow \alpha \\ \bullet \xrightarrow{\tilde{\xi}'} \bullet \end{array}} = \sqrt{\frac{\mu_f^\xi \mu_i^{\xi'}}{\mu_i^\xi \mu_f^{\xi'}}} \overline{\begin{array}{c} \xi' \\ \tilde{\alpha} \downarrow f \downarrow \tilde{\beta} \\ \bullet \xrightarrow{\xi} \bullet \end{array}} \quad (2.10)$$

where $\tilde{\alpha}$ denotes the path that is obtained from α reversing the arrow.

(iii) Concatenation properties:

$$\begin{array}{c} \xi_1 \star \xi_2 \\ \bullet \xrightarrow{n} \bullet \\ \alpha \downarrow f \downarrow \beta \\ \bullet \xrightarrow{\xi_1 \star \xi_2'} \bullet \end{array} = \begin{array}{c} \xi_1 \\ \bullet \xrightarrow{n} \bullet \\ \alpha \downarrow f \downarrow \gamma \\ \bullet \xrightarrow{\xi_1'} \bullet \end{array} \begin{array}{c} \xi_2 \\ \bullet \xrightarrow{n} \bullet \\ \gamma \downarrow f \downarrow \beta \\ \bullet \xrightarrow{\xi_2'} \bullet \end{array} \quad (2.11)$$

where the length one path γ is determined by the ending(starting) vertices of $\xi_1(\xi_2)$ and $\xi_1'(\xi_2')$.

Remark 2.2. It is very important to realize at this stage that, although connections have been defined in eq.(2.5) for elements in $End^{gr}(\mathcal{E})$, property (iii) allows to define them for elements in $End^{gr}(\mathcal{P})$. This is so since by concatenation as in eq.(2.11) it is possible to build any path out of length-zero and length-one paths that are necessarily essential.

4. The tensor product representation is given by,

$$\Phi_{\alpha \star \beta, \alpha' \star \beta'}^{f \otimes f'}(\xi \otimes \xi') = \sum_{\xi_i} \Phi_{\alpha, \alpha'}^f(\xi \otimes \xi^i) \Phi_{\beta, \beta'}^{f'}(\xi_i \otimes \xi') \quad . \quad (2.12)$$

where f, f' can be any of the fundamentals and where the dually paired basis vectors ξ_i and ξ^i have been defined at the end of section 2.1.

Example 2.3 (The case of A_3 .) *In this case there is only one fundamental representation that according to what is shown in appendix A can be chosen as,*

$$\begin{array}{c} 0 \quad r_0 \quad 1 \\ \bullet \xrightarrow{\quad} \bullet \\ r_0 \downarrow f \downarrow l_1 \\ \bullet \xrightarrow{\quad} \bullet \\ 1 \quad l_1 \quad 0 \end{array} = 1 \quad , \quad \begin{array}{c} 0 \quad r_0 \quad 1 \\ \bullet \xrightarrow{\quad} \bullet \\ r_0 \downarrow f \downarrow r_1 \\ \bullet \xrightarrow{\quad} \bullet \\ 1 \quad r_1 \quad 2 \end{array} = 1 \quad , \quad \begin{array}{c} 2 \quad l_2 \quad 1 \\ \bullet \xrightarrow{\quad} \bullet \\ l_2 \downarrow f \downarrow l_1 \\ \bullet \xrightarrow{\quad} \bullet \\ 1 \quad l_1 \quad 0 \end{array} = 1 \quad , \quad \begin{array}{c} 2 \quad l_2 \quad 1 \\ \bullet \xrightarrow{\quad} \bullet \\ l_2 \downarrow f \downarrow r_1 \\ \bullet \xrightarrow{\quad} \bullet \\ 1 \quad r_1 \quad 2 \end{array} = -1 \quad ,$$

$$\begin{array}{c} 1 \quad l_1 \quad 0 \\ \bullet \xrightarrow{\quad} \bullet \\ l_1 \downarrow f \downarrow r_0 \\ \bullet \xrightarrow{\quad} \bullet \\ 0 \quad r_0 \quad 1 \end{array} = \frac{1}{\sqrt{2}} \quad , \quad \begin{array}{c} 1 \quad r_1 \quad 2 \\ \bullet \xrightarrow{\quad} \bullet \\ l_1 \downarrow f \downarrow l_2 \\ \bullet \xrightarrow{\quad} \bullet \\ 0 \quad r_0 \quad 1 \end{array} = \frac{1}{\sqrt{2}} \quad ,$$

$$\begin{array}{ccc}
 \begin{array}{ccc}
 1 & l_1 & 0 \\
 \bullet & \xrightarrow{\quad} & \bullet \\
 r_1 \downarrow & f & \downarrow r_0 \\
 \bullet & \xrightarrow{\quad} & \bullet \\
 2 & l_2 & 1
 \end{array} & = \frac{1}{\sqrt{2}} & , \quad \begin{array}{ccc}
 1 & r_1 & 2 \\
 \bullet & \xrightarrow{\quad} & \bullet \\
 r_1 \downarrow & f & \downarrow l_2 \\
 \bullet & \xrightarrow{\quad} & \bullet \\
 2 & l_2 & 1
 \end{array} = \frac{-1}{\sqrt{2}}
 \end{array} \tag{2.13}$$

where for the sake of completeness we have included the corresponding vertex labels in each cell. Using (iii) it is possible to compute the value of cells for longer horizontal paths, for example,

$$\begin{aligned}
 & \begin{array}{ccc}
 0 & r_0 \star r_1 & 2 \\
 \bullet & \xrightarrow{\quad} & \bullet \\
 r_0 \downarrow & f & \downarrow l_2 \\
 \bullet & \xrightarrow{\quad} & \bullet \\
 1 & \gamma & 1
 \end{array} = \frac{1}{\sqrt{2}} \left(\begin{array}{ccc}
 0 & r_0 \star r_1 & 2 \\
 \bullet & \xrightarrow{\quad} & \bullet \\
 r_0 \downarrow & f & \downarrow l_2 \\
 \bullet & \xrightarrow{\quad} & \bullet \\
 1 & r_1 \star l_2 & 1
 \end{array} - \begin{array}{ccc}
 0 & r_0 \star r_1 & 2 \\
 \bullet & \xrightarrow{\quad} & \bullet \\
 r_0 \downarrow & f & \downarrow l_2 \\
 \bullet & \xrightarrow{\quad} & \bullet \\
 1 & l_1 \star r_0 & 1
 \end{array} \right) \\
 & = \frac{1}{\sqrt{2}} \left(\begin{array}{ccc}
 0 & r_0 & 1 \\
 \bullet & \xrightarrow{\quad} & \bullet \\
 r_0 \downarrow & f & \downarrow r_1 \\
 \bullet & \xrightarrow{\quad} & \bullet \\
 1 & r_1 & 2
 \end{array} \begin{array}{ccc}
 1 & r_1 & 2 \\
 \bullet & \xrightarrow{\quad} & \bullet \\
 r_1 \downarrow & f & \downarrow r_0 \\
 \bullet & \xrightarrow{\quad} & \bullet \\
 2 & l_2 & 1
 \end{array} - \begin{array}{ccc}
 0 & r_0 & 1 \\
 \bullet & \xrightarrow{\quad} & \bullet \\
 r_0 \downarrow & f & \downarrow l_1 \\
 \bullet & \xrightarrow{\quad} & \bullet \\
 1 & l_1 & 0
 \end{array} \begin{array}{ccc}
 1 & r_1 & 2 \\
 \bullet & \xrightarrow{\quad} & \bullet \\
 l_1 \downarrow & f & \downarrow l_2 \\
 \bullet & \xrightarrow{\quad} & \bullet \\
 0 & r_0 & 1
 \end{array} \right) = 1 \tag{2.14}
 \end{aligned}$$

which, as we see, corresponds to a horizontal "concatenation" of basic cells. Using 4. we can compute the value of cells in tensor product representation, for example,

$$\begin{aligned}
 \Phi_{r_0 \star l_1, r_1 \star l_2}^{f \otimes f}(r_0 \otimes r_0) &= r_0 \star l_1 \begin{array}{ccc}
 0 & r_0 & 1 \\
 \bullet & \xrightarrow{\quad} & \bullet \\
 \downarrow & f \otimes f & \downarrow \\
 \bullet & \xrightarrow{\quad} & \bullet \\
 0 & r_0 & 1
 \end{array} r_1 \star l_2 \begin{array}{ccc}
 0 & r_0 & 1 \\
 \bullet & \xrightarrow{\quad} & \bullet \\
 \downarrow & f & \downarrow \\
 \bullet & \xrightarrow{\quad} & \bullet \\
 1 & r_1 & 2
 \end{array} \begin{array}{ccc}
 1 & r_1 & 2 \\
 \bullet & \xrightarrow{\quad} & \bullet \\
 \downarrow & f & \downarrow \\
 \bullet & \xrightarrow{\quad} & \bullet \\
 0 & r_0 & 1
 \end{array} \\
 &= 1/\sqrt{2} = \Phi_{r_0, r_1}^x(r_0 \otimes r_1) \Phi_{l_1, l_2}^y(r_1 \otimes r_0) \tag{2.15}
 \end{aligned}$$

which, as we see, corresponds to a vertical "concatenation" of basic cells.

3 Algebra structure of the DTA

Proposition 3.1. *There exists a basis of the DTA denoted by $E_{\eta\eta'}^x$, where the index x labels the irreducible representations of the DTA and η, η' are indices in the irrep x . In this basis the product is given by,*

$$E_{\eta_1 \eta'_1}^{x_1} \cdot E_{\eta_2 \eta'_2}^{x_2} = \delta_{\eta'_1 \eta_2} \delta_{x_1 x_2} E_{\eta_1 \eta'_2}^{x_1} \tag{3.1}$$

The identity for this product is given by,

$$\mathbb{1} = \sum_{x, \eta} E_{\eta\eta}^x \tag{3.2}$$

The matrix \ast -representations (2.5) are homomorphisms Φ^x in the following way,

$$\Phi_{\alpha\gamma}^x(E_{\eta_1 \eta'_1}^{x_1} \cdot E_{\eta_2 \eta'_2}^{x_2}) = \sum_{\beta} \Phi_{\alpha\beta}^x(E_{\eta_1 \eta'_1}^{x_1}) \Phi_{\beta\gamma}^x(E_{\eta_2 \eta'_2}^{x_2}) \tag{3.3}$$

The star structure takes the following form,

$$(E_{\eta\eta'}^x)^* = E_{\eta'\eta}^x \tag{3.4}$$

The scalar product is given by,

$$\langle E_{\eta_1\eta'_1}^{x_1} | E_{\eta_2\eta'_2}^{x_2} \rangle = \delta_{x_1x_2} \delta_{\eta_1\eta_2} \delta_{\eta'_1\eta'_2} \tag{3.5}$$

Proof. Since we are restricting to graphs where the dimension of \mathcal{E} is finite and using 2. we conclude that the DTA we are considering are finite dimensional C^* -algebras. Hence they are isomorphic to a direct sum of matrix algebras corresponding to their irreps. Since we are dealing with matrix algebras we can take a basis consisting of matrix units,

$$(E_{\eta\eta'}^x)_{\alpha\beta} = \delta_{\eta\alpha} \delta_{\eta'\beta} \tag{3.6}$$

in this basis of matrix units, matrix multiplication corresponds to (3.1). That $\mathbb{1}$ in eq.(3.2) is the identity for this product is very simple to verify. As we mentioned above, the terms of this direct sum decomposition¹² correspond to the irreducible representations of the DTA

$$\Phi_{\alpha\beta}^y(E_{\eta\eta'}^x) = \delta_{xy} \delta_{\alpha\eta} \delta_{\beta\eta'} \tag{3.7}$$

Replacing (3.7) in (3.3) you verify that the later holds. Regarding the star structure using 3. we have,

$$\Phi_{\alpha\beta}^y((E_{\eta\eta'}^x)^*) = \overline{\Phi_{\beta\alpha}^y(E_{\eta\eta'}^x)} = \delta_{xy} \delta_{\beta\eta} \delta_{\alpha\eta'} = \Phi_{\alpha\beta}^y(E_{\eta'\eta}^x) \quad \forall x, y, \eta, \eta', \alpha, \beta \tag{3.8}$$

which leads to (3.4). Regarding the scalar product in the basis $\{E_{\eta\eta'}^x\}$ we note that since there is a unique correspondence between C^* -algebras and operator algebras and since the scalar product (3.5) leads to the operator norm, it must be that one. \square

3.1 Relation between basis

Let us adopt the notation that matrix irreps are represented by similar symbols as in the case of the fundamentals, i.e.,

$$\Phi_{\alpha\beta}^x(\xi \otimes \xi') = \begin{array}{ccc} & \xi & \\ & \bullet \xrightarrow{-n} \bullet & \\ \alpha \downarrow & x & \downarrow \beta \\ & \bullet \xrightarrow{\xi'} \bullet & \end{array} \tag{3.9}$$

¹²The discussion concerning the direct sum decomposition of tensor product representations can be rephrased in graphical terms by using the notion of cell systems associated with representations[10].

where x runs over all the irreps of the DTA(not only the fundamentals) as in proposition (3.1). It is worth noting that condition (2.6) should be fulfilled also in the case where x is not one of the fundamentals. This is so because the fundamentals tensorially generate any other irrep, recalling the definition (2.12) of tensor product representation we see that condition (2.6) is also fulfilled in any tensor product of fundamentals.

We have introduced two basis for the DTA, one in terms of endomorphisms of essential paths and another in terms of matrix units for the product in the DTA, there is a relation between them as stated by the following proposition,

Proposition 3.2. *The two basis $\{\xi \otimes \xi'\}$ and $\{E_{\alpha\beta}^x\}$ are related by,*

$$\xi \otimes \xi' = \sum_{x,\eta,\eta'} \eta \begin{array}{c} \xrightarrow{\xi} \\ \downarrow \eta \\ \xrightarrow{x} \\ \downarrow \eta' \\ \xrightarrow{\xi'} \end{array} E_{\eta\eta'}^x \quad (3.10)$$

where n is the length of the essential paths ξ and ξ' .

Proof. In general we have the following relation between these basis,

$$\xi \otimes \xi' = \sum_{x,\eta,\eta'} B(n, \xi, \xi', x, \eta, \eta') E_{\eta\eta'}^x \quad (3.11)$$

applying $\Phi_{\alpha\beta}^y$ to both sides of (3.11) and using (2.5) and (3.7)we get,

$$B(n, \xi, \xi', x, \eta, \eta') = \eta \begin{array}{c} \xrightarrow{\xi} \\ \downarrow \eta \\ \xrightarrow{x} \\ \downarrow \eta' \\ \xrightarrow{\xi'} \end{array} \eta'. \quad (3.12)$$

□

Defining the inverse cells by,

$$\sum_{x,\alpha,\beta} \eta \begin{array}{c} \xrightarrow{\xi_1} \\ \downarrow \eta \\ \xrightarrow{x^{-1}} \\ \downarrow \eta' \\ \xrightarrow{\xi'_1} \end{array} \eta' \begin{array}{c} \xrightarrow{\xi_2} \\ \downarrow \eta \\ \xrightarrow{x} \\ \downarrow \eta' \\ \xrightarrow{\xi'_2} \end{array} \eta' = \delta_{\xi_1 \xi_2} \delta_{\xi'_1 \xi'_2} \quad (3.13)$$

we have,

$$E_{\eta\eta'}^x = \sum_{\xi,\xi'} \eta \begin{array}{c} \xrightarrow{\xi} \\ \downarrow \eta \\ \xrightarrow{x^{-1}} \\ \downarrow \eta' \\ \xrightarrow{\xi'} \end{array} \eta' \xi \otimes \xi' \quad (3.14)$$

replacing (3.10) in (3.14) the following also holds,

$$\sum_{\xi, \xi'} \begin{array}{c} \xi \\ \bullet \xrightarrow{-n} \bullet \\ \eta \downarrow \quad \uparrow \eta' \\ \bullet \xrightarrow{x^{-1}} \bullet \\ \xi' \end{array} \begin{array}{c} \xi \\ \bullet \xrightarrow{-n} \bullet \\ \rho \downarrow \quad \uparrow \rho' \\ \bullet \xrightarrow{y} \bullet \\ \xi' \end{array} = \delta_{xy} \delta_{\eta\rho} \delta_{\eta'\rho'} \quad (3.15)$$

The result that follows will be very useful for further developments,

Proposition 3.3. *For any $(\xi \otimes \xi')_{\mathcal{P}} \in \text{End}^{gr}(\mathcal{P})$ there exists a unique element in $\text{End}^{gr}(\mathcal{E})$ given by,*

$$(\xi \otimes \xi')_{\mathcal{E}} = \sum_{x, \alpha, \beta} \begin{array}{c} \xi \\ \bullet \xrightarrow{-n} \bullet \\ \alpha \downarrow \quad \uparrow \beta \\ \bullet \xrightarrow{x} \bullet \\ \xi' \end{array} E_{\alpha\beta}^x \quad (3.16)$$

in such a way that¹³,

$$\begin{aligned} \Phi_{\alpha\beta}^x((\xi \otimes \xi')_{\mathcal{P}}) &= \Phi_{\alpha\beta}^x((\xi \otimes \xi')_{\mathcal{E}}) \\ \Phi_{\alpha\beta}^x((\xi_1 \otimes \xi'_1)_{\mathcal{E}} \cdot (\xi_2 \otimes \xi'_2)_{\mathcal{E}}) &= \Phi_{\alpha\beta}^x((\xi_1 \otimes \xi'_1)_{\mathcal{P}} \star (\xi_2 \otimes \xi'_2)_{\mathcal{P}}) \end{aligned} \quad (3.18)$$

Proof. The first equation in (3.18) is a consequence of the definition of $(\xi \otimes \xi')_{\mathcal{E}}$ in (3.16) and the remark, following eq.(2.11), about the extension of the definition of connections for elements in $\text{End}^{gr}(\mathcal{P})$. For the second equation in (3.18) we have,

$$\begin{aligned} &\Phi_{\alpha\beta}^x((\xi_1 \otimes \xi'_1)_{\mathcal{E}} \cdot (\xi_2 \otimes \xi'_2)_{\mathcal{E}}) = \\ &= \sum_{x_1, \alpha_1, \beta_1, x_2, \alpha_2, \beta_2} \begin{array}{c} \xi_1 \\ \bullet \xrightarrow{-n} \bullet \\ \alpha_1 \downarrow \quad \uparrow \beta_1 \\ \bullet \xrightarrow{x_1} \bullet \\ \xi'_1 \end{array} \begin{array}{c} \xi_2 \\ \bullet \xrightarrow{-n} \bullet \\ \alpha_2 \downarrow \quad \uparrow \beta_2 \\ \bullet \xrightarrow{x_2} \bullet \\ \xi'_2 \end{array} = \Phi_{\alpha\beta}^x(E_{\alpha_1\beta_1}^{x_1} \cdot E_{\alpha_2\beta_2}^{x_2}) \\ &= \sum_{\beta_1} \begin{array}{c} \xi_1 \\ \bullet \xrightarrow{-n} \bullet \\ \alpha \downarrow \quad \uparrow \beta_1 \\ \bullet \xrightarrow{x} \bullet \\ \xi'_1 \end{array} \begin{array}{c} \xi_2 \\ \bullet \xrightarrow{-n} \bullet \\ \beta_1 \downarrow \quad \uparrow \beta \\ \bullet \xrightarrow{x} \bullet \\ \xi'_2 \end{array} = \Phi_{\alpha\beta}^x((\xi_1 \otimes \xi'_1)_{\mathcal{P}} \star (\xi_2 \otimes \xi'_2)_{\mathcal{P}}) \end{aligned} \quad (3.19)$$

where we have employed eq.(3.1) and (3.7) to write the second equality and (2.11) for the last equality. \square

¹³In eq. (3.18) we use the following definition for the concatenation product of elements in $\text{End}^{gr}(\mathcal{P})$,

$$(\xi_1 \otimes \xi'_1)_{\mathcal{P}} \star (\xi_2 \otimes \xi'_2)_{\mathcal{P}} = (\xi_1 \star \xi_2)_{\mathcal{P}} \otimes (\xi'_1 \star \xi'_2)_{\mathcal{P}} \quad (3.17)$$

Proposition 3.4. *The following holds,*

$$\langle E_{\eta\eta'}^x | \xi \otimes \xi' \rangle = \begin{array}{c} \xi \\ \bullet \xrightarrow{-n} \bullet \\ \eta \downarrow \quad x \quad \downarrow \eta' \\ \bullet \xrightarrow{\quad} \bullet \\ \xi' \end{array} \quad (3.20)$$

$$\begin{aligned} \langle \rho \otimes \rho' | \xi \otimes \xi' \rangle &= \sum_{x, \eta, \eta'} \langle \rho \otimes \rho' | E_{\eta\eta'}^x \rangle \langle E_{\eta\eta'}^x | \xi \otimes \xi' \rangle = \\ &= \sum_{x, \eta, \eta'} \overline{\begin{array}{c} \rho \\ \bullet \xrightarrow{-n} \bullet \\ \eta \downarrow \quad x \quad \downarrow \eta' \\ \bullet \xrightarrow{\quad} \bullet \\ \rho' \end{array}} \begin{array}{c} \xi \\ \bullet \xrightarrow{-n} \bullet \\ \eta \downarrow \quad x \quad \downarrow \eta' \\ \bullet \xrightarrow{\quad} \bullet \\ \xi' \end{array} \end{aligned} \quad (3.21)$$

Proof. Taking scalar products in both sides of (3.10) and using (3.5) you get the first equality. The other follows from orthogonality and completeness of the $|E_{\eta\eta'}^x\rangle$ basis. \square

Furthermore as shown in appendix B the form of the scalar product is restricted to,

Proposition 3.5.

$$\langle \rho \otimes \rho' | \xi \otimes \xi' \rangle \propto \delta_{\rho\xi} \delta_{\rho'\xi'} \quad (3.22)$$

4 Weak bialgebra structure

4.1 Product

The expression of the product in the basis of endomorphisms of essential paths is given by the following.

Proposition 4.1. *We have,*

$$\xi_1 \otimes \xi'_1 \cdot \xi_2 \otimes \xi'_2 = \sum_{n_3, \xi_3, \xi'_3} P_{\xi'_1 \xi'_2 \xi'_3}^{\xi_1 \xi_2 \xi_3} \xi_3 \otimes \xi'_3 \quad (4.1)$$

where,

$$P_{\xi'_1 \xi'_2 \xi'_3}^{\xi_1 \xi_2 \xi_3} = \sum_{x, \eta_1, \eta_2, \eta_3} \begin{array}{c} \xi_1 \\ \bullet \xrightarrow{-n_1} \bullet \\ \eta_1 \downarrow \quad x \quad \downarrow \eta_2 \\ \bullet \xrightarrow{\quad} \bullet \\ \xi'_1 \end{array} \begin{array}{c} \xi_2 \\ \bullet \xrightarrow{-n_2} \bullet \\ \eta_2 \downarrow \quad x \quad \downarrow \eta_3 \\ \bullet \xrightarrow{\quad} \bullet \\ \xi'_2 \end{array} \begin{array}{c} \xi_3 \\ \bullet \xrightarrow{-n} \bullet \\ \eta_1 \downarrow \quad x^{-1} \quad \downarrow \eta_3 \\ \bullet \xrightarrow{\quad} \bullet \\ \xi'_3 \end{array} \quad (4.2)$$

Proof. Replacing (3.10) in the l.h.s. of (4.1), using (3.1) and employing (3.11) you get the r.h.s. of (4.1) and (4.2). \square

4.2 Coproduct

The definition of a tensor product representation, as given by (2.12), can be rephrased in terms of a coproduct.

Proposition 4.2. *The following coproduct,*

$$\Delta(\xi \otimes \xi') = \sum_{\xi_i} (\xi \otimes \xi^i)_{\mathcal{E}} \otimes (\xi_i \otimes \xi')_{\mathcal{E}} \quad (4.3)$$

where $\xi_i(\xi^i)$ are elements of the (dual) elementary paths basis of $\mathcal{E}(\hat{\mathcal{E}})$ respectively. This coproduct implies (2.12), it is a coassociative algebra morphism and it satisfies,

$$\Delta((\xi \otimes \xi')^*) = \Delta(\xi \otimes \xi')^* \quad (a \otimes b)^* = a^* \otimes b^* \quad (4.4)$$

Proof. Replacing eq.(4.3) in the r.h.s. of,

$$\Phi_{\alpha_1 \circ \alpha_2, \beta_1 \circ \beta_2}^{x \otimes y}(\xi \otimes \xi') = \Phi_{\alpha_1, \beta_1}^{x_1} \otimes \Phi_{\alpha_2, \beta_2}^{x_2}(\Delta(\xi \otimes \xi')) \quad (4.5)$$

using (2.5) and (3.18) you get the same as the r.h.s. of (2.12). Coassociativity is a simple verification. . In order to verify eq. (4.4) note that using (3.10) and (3.4) you obtain,

$$(\xi \otimes \xi')^* = \sum_{x, \eta, \eta'} \overline{\begin{array}{c} \xi \\ \eta \downarrow \begin{array}{c} \bullet \xrightarrow{-n_2} \bullet \\ \bullet \xrightarrow{x} \bullet \\ \bullet \xrightarrow{\xi'} \bullet \end{array} \uparrow \eta' \\ \xi' \end{array}} E_{\eta' \eta}^x \quad (4.6)$$

on the other hand using (3.10) and reflection you have.

$$\begin{aligned} \tilde{\xi} \otimes \tilde{\xi}' &= \sum_{x, \eta', \eta} \eta' \downarrow \begin{array}{c} \tilde{\xi} \\ \bullet \xrightarrow{-n_2} \bullet \\ \bullet \xrightarrow{x} \bullet \\ \bullet \xrightarrow{\tilde{\xi}'} \bullet \end{array} \uparrow \eta E_{\eta' \eta}^x = \sum_{x, \eta', \eta} \sqrt{\frac{\mu_i^\xi \mu_f^{\xi'}}{\mu_f^\xi \mu_i^{\xi'}}} \overline{\begin{array}{c} \xi \\ \eta \downarrow \begin{array}{c} \bullet \xrightarrow{-n_2} \bullet \\ \bullet \xrightarrow{x} \bullet \\ \bullet \xrightarrow{\xi'} \bullet \end{array} \uparrow \eta' \\ \xi' \end{array}} E_{\eta' \eta}^x = \\ &= \sqrt{\frac{\mu_i^\xi \mu_f^{\xi'}}{\mu_f^\xi \mu_i^{\xi'}}} (\xi \otimes \xi')^* \end{aligned} \quad (4.7)$$

where in the last equality we have used (4.6). Using (4.7) it is simple to verify (4.4). The morphism property of the coproduct is dealt with in appendix C. \square

The coproduct can be expressed in the basis $\{E_{\alpha\beta}^x\}$ as stated below.

Proposition 4.3. *We have,*

$$\Delta(E_{\eta\eta'}^x) = \sum_{x_1, \eta_1, \eta'_1, x_2, \eta_2, \eta'_2} \tilde{P}_{\eta\eta' x_1 x_2}^{x \eta_1 \eta'_1 \eta_2 \eta'_2} E_{\eta_1 \eta'_1}^{x_1} \otimes E_{\eta_2 \eta'_2}^{x_2} \quad (4.8)$$

where,

$$\tilde{P}_{\eta\eta' x_1 x_2}^{x \eta_1 \eta'_1 \eta_2 \eta'_2} = \sum_{n, \xi_1, \xi_2, \xi_3} \begin{array}{c} \xi_1 \\ \bullet \xrightarrow{n_1} \bullet \\ \eta \downarrow \quad \downarrow \eta' \\ \bullet \xrightarrow{x^{-1}} \bullet \\ \xi_3 \end{array} \begin{array}{c} \xi_1 \\ \bullet \xrightarrow{n_1} \bullet \\ \eta_1 \downarrow \quad \downarrow \eta'_1 \\ \bullet \xrightarrow{x_1} \bullet \\ \xi_2 \end{array} \begin{array}{c} \xi_2 \\ \bullet \xrightarrow{n_1} \bullet \\ \eta_2 \downarrow \quad \downarrow \eta'_2 \\ \bullet \xrightarrow{x_2} \bullet \\ \xi_3 \end{array} \quad (4.9)$$

Proof. Replacing (3.14) in the l.h.s. of (4.8), employing (4.3) and using (3.10) two times you get the result. \square

4.3 Counit

Regarding the counit the following holds.

Proposition 4.4. *A counit satisfying the property,*

$$(\epsilon \otimes \mathbb{1})\Delta = \mathbb{1} = (\mathbb{1} \otimes \epsilon)\Delta \quad (4.10)$$

is given by,

$$\epsilon(\xi \otimes \xi') = \delta_{\xi\xi'} \quad (4.11)$$

Proof. Just replace (4.11) in (4.10). \square

It is important to note that the relation $\Delta(\mathbb{1}) = \mathbb{1} \otimes \mathbb{1}$ does not hold in the DTA, which is therefore *not* a Hopf algebra. The structure of the DTA appearing up to this part corresponds to a weak bialgebra in the terminology of [17] definition 2.1. As we shall see in the next section there is also an antipode and the DTA is actually a weak Hopf algebra or "quantum groupoid" (this general notion is discussed in [3, 14, 15, 16, 17]).

5 Weak *-Hopf algebra structure. Antipode

According to Theorem 8.7 of [17], we have

Theorem 5.1 (Nill). *Let $(A, \mathbb{1}, \Delta, \epsilon)$ be a weak bialgebra and $S : A \rightarrow A$ be a bialgebra antiautomorphism. If there exists a non-degenerate linear functional $\lambda : A \rightarrow \mathbb{C}$ such that ¹⁴,*

$$a_{(1)}\lambda(ba_{(2)}) = S(b_{(1)})\lambda(b_{(2)}a) \quad , \forall a, b \in A \tag{5.2}$$

then S is an antipode and A is a weak Hopf algebra.

Below we prove the following result,

Theorem 5.2. *The DTA algebra has the structure of a weak $*$ -Hopf algebra with product given by (3.1), unit given by (3.2), coproduct (4.3), counit (4.11) , star (3.4) and antipode,*

$$S(\xi_i \otimes \xi^j) = \tilde{\xi}_j \otimes \tilde{\xi}^i \tag{5.3}$$

Proof. Take,

$$\lambda(E_{\eta\eta'}^x) = \delta_{x\mathbf{0}}\delta_{\eta\eta'} \tag{5.4}$$

where $\mathbf{0}$ denotes the trivial representation as described in appendix D. Let,

$$a = \xi_1 \otimes \xi'_1 \quad , b = \xi_2 \otimes \xi'_2 \tag{5.5}$$

in (5.2). Using (4.3) we obtain,

$$\Delta(a) = \sum_{\rho_1} (\xi_1 \otimes \rho_1) \otimes (\rho_1 \otimes \xi'_1) \quad \Delta(b) = \sum_{\rho_2} (\xi_2 \otimes \rho_2) \otimes (\rho_2 \otimes \xi'_2) \tag{5.6}$$

therefore consider,

$$\begin{aligned} A = & \sum_{\rho_1} (\xi_1 \otimes \rho_1) \lambda((\xi_2 \otimes \xi'_2) \cdot (\rho_1 \otimes \xi'_1)) \\ & - \sum_{\rho_2} S((\xi_2 \otimes \rho_2)) \lambda((\rho_2 \otimes \xi'_2) \cdot (\xi_1 \otimes \xi'_1)) \end{aligned} \tag{5.7}$$

so that eq.(5.2) becomes $A = 0$. Now in order to evaluate the products in the arguments of λ we rewrite (5.7) in terms of the basis $\{E_{\eta\eta'}^x\}$ to obtain,

$$\begin{aligned} A = & \sum_{x_1, \eta_1, \eta'_1, x_2, \eta_2, \eta'_2} \left[\sum_{\rho_1} (\xi_1 \otimes \rho_1) \eta_1 \begin{array}{c} \xrightarrow{\xi_2} \\ \downarrow x_1 \\ \xrightarrow{\xi'_2} \end{array} \eta'_1 \quad \eta_2 \begin{array}{c} \xrightarrow{\rho_1} \\ \downarrow x_2 \\ \xrightarrow{\xi'_1} \end{array} \eta'_2 - \\ & - \sum_{\rho_2} S(\xi_2 \otimes \rho_2) \eta_1 \begin{array}{c} \xrightarrow{\rho_2} \\ \downarrow x_1 \\ \xrightarrow{\xi'_2} \end{array} \eta'_1 \quad \eta_2 \begin{array}{c} \xrightarrow{\xi_1} \\ \downarrow x_2 \\ \xrightarrow{\xi'_1} \end{array} \eta'_2 \right] \lambda(E_{\eta_1 \eta'_1}^x \cdot E_{\eta_2 \eta'_2}^x) \end{aligned} \tag{5.8}$$

¹⁴In eq.(5.2) we employed Sweedler's notation for the coproduct, i.e.,

$$\Delta(a) = a_{(1)} \otimes a_{(2)} \quad \Delta(b) = b_{(1)} \otimes b_{(2)} \tag{5.1}$$

now we use (3.1) and (5.4) to obtain $\lambda(E_{\eta_2\eta'_2}^{x_2} \cdot E_{\eta_1\eta'_1}^{x_1}) = \delta_{x_1\mathbf{0}}\delta_{x_1x_2}\delta_{\eta'_1\eta_2}\delta_{\eta_1\eta'_2}$, replacing in (5.8) and using reflection we obtain,

$$\begin{aligned}
A &= \sum_{\eta_1, \eta_2} \left[\sum_{\rho_1} \sqrt{\frac{\mu_f^{\xi_2} \mu_i^{\xi'_2}}{\mu_i^{\xi_2} \mu_f^{\xi'_2}}} (\xi_1 \otimes \rho_1) \eta_1 \begin{array}{c} \overline{\xi_2} \\ \bullet \xrightarrow{n} \bullet \\ \mathbf{0} \\ \bullet \xrightarrow{\xi'_2} \bullet \\ \xi_2 \end{array} \eta_2 \begin{array}{c} \rho_1 \\ \bullet \xrightarrow{n} \bullet \\ \mathbf{0} \\ \bullet \xrightarrow{\xi'_1} \bullet \\ \xi_1 \end{array} \right. \\
&\quad \left. - \sqrt{\frac{\mu_f^{\rho_2} \mu_i^{\xi'_2}}{\mu_i^{\rho_2} \mu_f^{\xi'_2}}} S((\xi_2 \otimes \rho_2)) \begin{array}{c} \overline{\tilde{\rho}_2} \\ \bullet \xrightarrow{n} \bullet \\ \mathbf{0} \\ \bullet \xrightarrow{\xi'_2} \bullet \\ \xi_2 \end{array} \eta_2 \begin{array}{c} \xi_1 \\ \bullet \xrightarrow{n} \bullet \\ \mathbf{0} \\ \bullet \xrightarrow{\xi'_1} \bullet \\ \xi_1 \end{array} \right] \\
&= \sqrt{\frac{\mu_f^{\xi_2} \mu_i^{\xi'_2}}{\mu_i^{\xi_2} \mu_f^{\xi'_2}}} \sum_{\rho_1} (\xi_1 \otimes \rho_1) \langle \tilde{\xi}_2 \otimes \tilde{\xi}'_2 | \rho_1 \otimes \xi'_1 \rangle \\
&\quad - \sum_{\rho_2} \sqrt{\frac{\mu_f^{\rho_2} \mu_i^{\xi'_2}}{\mu_i^{\rho_2} \mu_f^{\xi'_2}}} S((\xi_2 \otimes \rho_2)) \langle \tilde{\rho}_2 \otimes \tilde{\xi}'_2 | \xi_1 \otimes \xi'_1 \rangle \quad (5.9)
\end{aligned}$$

where we have employed (3.5) and the definition of scalar product in a representation appearing in Appendix B. Using the results in appendix B and noting, as shown in appendix D, that in the representation $\mathbf{0}$ the only non-vanishing cells are those with equal upper and lower horizontal paths we conclude that there is the following factorization in eq.(5.9),

$$A = \sqrt{\frac{\mu_f^{\xi_2} \mu_i^{\xi'_2}}{\mu_i^{\xi_2} \mu_f^{\xi'_2}}} \langle \xi'_1 \otimes \xi'_1 | \xi'_1 \otimes \xi'_1 \rangle (\xi_1 \otimes \tilde{\xi}_2 \delta_{\xi'_1 \tilde{\xi}'_2} - S((\xi_2 \otimes \tilde{\xi}_1)) \delta_{\xi'_1 \tilde{\xi}'_2}) \quad (5.10)$$

hence $A = 0$ implies,

$$S(\xi \otimes \xi') = \tilde{\xi}' \otimes \tilde{\xi} \quad (5.11)$$

Next we show S is a bialgebra antiautomorphism, i.e.,

$$S(a \cdot b) = S(b)S(a) \quad (5.12)$$

In order to do this we first express the antipode in the $\{E_{\eta\eta'}^x\}$ basis. Replacing (3.10) in (5.11) we obtain,

$$\sum_{x, \eta, \eta'} \eta \begin{array}{c} \xi \\ \bullet \xrightarrow{n} \bullet \\ x \\ \bullet \xrightarrow{\xi'} \bullet \\ \xi' \end{array} \eta' S(E_{\eta\eta'}^x) = \sum_{y, \rho, \rho'} \rho \begin{array}{c} \tilde{\xi}' \\ \bullet \xrightarrow{n} \bullet \\ y \\ \bullet \xrightarrow{\tilde{\xi}} \bullet \\ \tilde{\xi} \end{array} \rho' E_{\rho\rho'}^y \quad (5.13)$$

multiplying this equation by,

$$\begin{array}{ccc}
 & \xi & \\
 \lambda \downarrow & \begin{array}{c} \bullet \xrightarrow{-n} \bullet \\ \downarrow z^{-1} \\ \bullet \xrightarrow{\xi'} \bullet \end{array} & \downarrow \lambda' \\
 & \xi' &
 \end{array} \tag{5.14}$$

summing up over ξ and ξ' and using (3.15) we have,

$$S(E_{\lambda\lambda'}^z) = \sum_{y,\rho,\rho'} \sum_{\xi\xi'} \lambda \begin{array}{c} \bullet \xrightarrow{-n} \bullet \\ \downarrow z^{-1} \\ \bullet \xrightarrow{\xi'} \bullet \end{array} \lambda' \begin{array}{c} \bullet \xrightarrow{-n} \bullet \\ \downarrow x \\ \bullet \xrightarrow{\xi} \bullet \end{array} \rho' E_{\rho\rho'}^y \tag{5.15}$$

now using two times reflection we have that,

$$\begin{array}{ccc}
 & \xi' & \\
 \rho \downarrow & \begin{array}{c} \bullet \xrightarrow{-n} \bullet \\ \downarrow x \\ \bullet \xrightarrow{\xi} \bullet \end{array} & \downarrow \rho' \\
 & \xi &
 \end{array} = \begin{array}{ccc}
 & \xi & \\
 \tilde{\rho}' \downarrow & \begin{array}{c} \bullet \xrightarrow{-n} \bullet \\ \downarrow x \\ \bullet \xrightarrow{\xi'} \bullet \end{array} & \downarrow \tilde{\rho} \\
 & \xi' &
 \end{array} \tag{5.16}$$

replacing in (5.15) and using once more (3.15) we obtain,

$$S(E_{\lambda\lambda'}^z) = E_{\lambda'\tilde{\lambda}}^z \tag{5.17}$$

hence,

$$S(E_{\eta_1\eta_1'}^{x_1} \cdot E_{\eta_2\eta_2'}^{x_2}) = \delta_{\eta_1'\eta_2} \delta_{x_1x_2} S(E_{\eta_1\eta_2'}^{x_1}) = \delta_{\eta_1'\eta_2} \delta_{x_1x_2} E_{\tilde{\eta}_2'\tilde{\eta}_1}^{x_1} \tag{5.18}$$

on the other hand,

$$S(E_{\eta_1\eta_1'}^{x_1}) \cdot S(E_{\eta_2\eta_2'}^{x_2}) = E_{\tilde{\eta}_2'\tilde{\eta}_2}^{x_2} \cdot E_{\tilde{\eta}_1'\tilde{\eta}_1}^{x_1} = \delta_{\eta_1'\eta_2} \delta_{x_1x_2} E_{\tilde{\eta}_2'\tilde{\eta}_1}^{x_1} \tag{5.19}$$

that proves (5.12). □

Example 5.3 (The case of A_3). For the case of the fundamental representation, that we denote in this case by $\mathbf{1}$, using (3.10) we can express elements of $End^{gr}(\mathcal{E})$ in terms of matrix units. The connections involved can be calculated from the basic ones and horizontal concatenation of cells. This can be summarized in the following matrix,

$$\begin{array}{l}
 01 \\
 10 \\
 12 \\
 21
 \end{array} \begin{pmatrix}
 01 & r_0l_1 & r_0r_1 & -d\gamma \\
 l_1r_0/\sqrt{2} & 10 & -\gamma d & r_1r_0/\sqrt{2} \\
 l_1l_2/\sqrt{2} & -\gamma g & 12 & -r_1l_2/\sqrt{2} \\
 -g\gamma & l_2l_1 & -l_2r_1 & 21
 \end{pmatrix} \tag{5.20}$$

that should be interpreted in the following way. The matrix representing the element $\xi_i \otimes \xi^j$ in the representation f is the one obtained from (5.20) by

replacing in it $\xi_i \xi^j$ by 1 and all the others by 0. Now we consider the tensor product representation $\mathbf{1} \otimes \mathbf{1}$. Taking into account (2.12) and with the same conventions as in (5.20) you obtain the following matrix,

$$\begin{matrix}
 010 \\
 212 \\
 101 \\
 121 \\
 012 \\
 210
 \end{matrix}
 \begin{pmatrix}
 00 & dd & \frac{1}{\sqrt{2}}r_0r_0 & \frac{1}{\sqrt{2}}r_0r_0 & 0 & 0 \\
 gg & 22 & \frac{1}{\sqrt{2}}l_2l_2 & \frac{1}{\sqrt{2}}l_2l_2 & 0 & 0 \\
 \frac{1}{\sqrt{2}}l_1l_1 & \frac{1}{\sqrt{2}}r_1r_1 & 11 & \gamma\gamma & \frac{1}{\sqrt{2}}l_1r_1 & \frac{1}{\sqrt{2}}r_1l_1 \\
 \frac{1}{\sqrt{2}}l_1l_1 & \frac{1}{\sqrt{2}}r_1r_1 & \gamma\gamma & 11 & -\frac{1}{\sqrt{2}}l_1r_1 & \frac{1}{\sqrt{2}}r_1l_1 \\
 0 & 0 & \frac{1}{\sqrt{2}}r_0l_2 & -\frac{1}{\sqrt{2}}r_0l_2 & 02 & dg \\
 0 & 0 & \frac{1}{\sqrt{2}}l_2r_0 & -\frac{1}{\sqrt{2}}l_2r_0 & gd & 20
 \end{pmatrix}
 \quad (5.21)$$

Note that the labels of files and columns are two step paths (not in general essential) in the graph A_3 . In order to decompose this tensor product representation in direct sum of irreps it is worth noting that,

1. In many instances of elements in $End^{gr}(\mathcal{E})$ say $\xi_i \otimes \xi^j$, the indices corresponding to non-vanishing entries in the irrep under study are completely determined because there is only one possibility for them. For example take the element $0 \otimes 0$ then the only possibility for the step two vertical paths is $\alpha = (010)$ and $\beta = (010)$ That is the only possible non-zero cell is,

$$\begin{array}{ccc}
 0 & \longrightarrow & 0 \\
 (010) \downarrow & f & \downarrow (010) \\
 0 & \longrightarrow & 0
 \end{array}
 = 1
 \quad (5.22)$$

In other words, the end points of the horizontal essential paths give only one possibility for the vertical (not necessarily essential) paths.

2. Note that in the case of A_3 if the two initial or final points of the horizontal paths are 1 then in general the associated vertical paths are not fixed and they can be a linear combination of (101) and (121). This combination should be determined in such a way that one gets a block decomposition of the matrix (5.21).

From the above remarks one reduces in this case the block diagonalization problem to the one of 2×2 matrix. It is a simple calculation to see that the right vertical labels are the following combinations of (101) and (121),

$$\gamma = 1/\sqrt{2}((121) - (101)) \quad , \quad \gamma' = 1/\sqrt{2}((121) + (101))
 \quad (5.23)$$

with these horizontal labels one gets for the matrix (5.21) the following,

$$\begin{array}{l}
 010 \\
 212 \\
 \gamma' \\
 \gamma \\
 012 \\
 210
 \end{array}
 \left(
 \begin{array}{cccccc}
 00 & dd & r_0r_0 & 0 & 0 & 0 \\
 gg & 22 & l_2l_2 & 0 & 0 & 0 \\
 l_1l_1 & r_1r_1 & 11 + \gamma\gamma & 0 & 0 & 0 \\
 0 & 0 & 0 & 11 - \gamma\gamma & -l_1r_1 & -r_1l_1 \\
 0 & 0 & 0 & -r_0l_2 & 02 & dg \\
 0 & 0 & 0 & -l_2r_0 & gd & 20
 \end{array}
 \right) \quad (5.24)$$

the first block in (5.24) corresponds to the trivial representation $\mathbf{0}$ that is described in appendix D. The other block is denoted by $\mathbf{2}$ and the following table of tensor product representations decomposition in irreps is easily obtained by employing analogous methods as in the case of $\mathbf{1} \otimes \mathbf{1}$,

$$\begin{aligned}
 \mathbf{0} \otimes \mathbf{1} &= \mathbf{1} \\
 \mathbf{0} \otimes \mathbf{2} &= \mathbf{2} \\
 \mathbf{1} \otimes \mathbf{2} &= \mathbf{1} \\
 \mathbf{1} \otimes \mathbf{1} &= \mathbf{0} \oplus \mathbf{2}
 \end{aligned} \quad (5.25)$$

this table itself can be represented by a graph. In this graph each vertex corresponds to an irrep of the DTA of A_3 . If a line joins two vertices, it means that one irrep can be obtained from the another one by taking tensor product with the fundamental $\mathbf{1}$. Notice that the trivial representation $\mathbf{0}$ is not 1-dimensional (this is a usual feature of weak Hopf algebras). In this case this graph coincides with the graph A_3 . In the literature this graph is referred to as Ocneanu’s graph of quantum symmetries.

In general, the fusion algebra –that describes the tensor product of representations for the composition product in the dual– has no reason to be isomorphic with the algebra of quantum symmetries – that describes the tensor product of representations for the convolution product \cdot . However, in the case of A_n graphs (and only in such cases), these two algebras are isomorphic. For these particular Dynkin diagrams, the block decomposition of the two products (the convolution product \cdot in the DTA and the composition of endomorphisms in its dual) have the same number of direct summands and these summands have also the same dimensions.

A General form of ADE connections

Proposition A.1. *Given a cell system where the four sides of the cells correspond to length-one paths on the same ADE graph G , the following*

two “basic” connections^{15, 16} (we only give one of the two, but the other is obtained by taking the complex conjugate of the first) satisfy unitarity and reflection,

$$\begin{array}{ccc} v_i & \xrightarrow{\quad} & v_l \\ \downarrow & \text{\scriptsize f} & \downarrow \\ v_k & \xrightarrow{\quad} & v_j \end{array} = \delta_{v_k v_l} \epsilon + \delta_{v_i v_j} \sqrt{\frac{\mu_{v_k} \mu_{v_l}}{\mu_{v_i} \mu_{v_j}}} \bar{\epsilon} \quad (\text{A.1})$$

where $\epsilon = ie^{i\pi/2N}$ with N the Coxeter number of G .

Proof. Reflection is a simple check. The unitarity conditions is,

$$A = \sum_{v_l} \begin{array}{ccc} v_i & \xrightarrow{\quad} & v_l \\ \downarrow & \text{\scriptsize f} & \downarrow \\ v_k & \xrightarrow{\quad} & v_j \end{array} \begin{array}{ccc} \overline{v_i} & \xrightarrow{\quad} & \overline{v_l} \\ \downarrow & \text{\scriptsize f} & \downarrow \\ v_m & \xrightarrow{\quad} & v_j \end{array} = \delta_{v_k v_m} \quad (\text{A.2})$$

Consider now the case when $v_i \neq v_j$. In this case the only term of (A.1) that contributes is the first and replacing in (A.2) you get,

$$A = \sum_{v_l} \delta_{v_k v_l} \epsilon \delta_{v_m v_l} \bar{\epsilon} = \delta_{v_k v_m} \quad (\text{A.3})$$

as claimed. In the other case $v_i = v_j$ you have,

$$\begin{aligned} A &= \sum_l \left(\sqrt{\frac{\mu_{v_k} \mu_{v_l}}{\mu_{v_i}^2}} \bar{\epsilon} + \delta_{v_k v_l} \epsilon \right) \left(\sqrt{\frac{\mu_{v_m} \mu_{v_l}}{\mu_{v_i}^2}} \epsilon + \delta_{v_k v_l} \bar{\epsilon} \right) \\ &= \delta_{v_k v_m} + \sqrt{\frac{\mu_{v_m} \mu_{v_l}}{\mu_{v_i}^2}} [\epsilon^2 + \bar{\epsilon}^2 + \frac{1}{\mu_{v_i}} \sum_{\langle v_l v_i \rangle} \mu_{v_l}] \end{aligned} \quad (\text{A.4})$$

where $\langle v_l v_i \rangle$ in (A.4) indicates that the summation is over the vertices v_l in G that are connected with the vertex v_i . This summation can therefore be expressed in terms of the adjacency matrix M of G and its Perron-Frobenius eigenvector v_{pf} as follows,

$$\sum_{\langle v_l v_i \rangle} \mu_{v_l} = \langle \text{row } v_i \text{ of } M, v_{pf} \rangle = \langle \hat{e}_{v_i} M, v_{pf} \rangle \quad (\text{A.5})$$

where $\{\hat{e}_{v_i}\}$ is a basis vector associated to the vertex v_i of G . The scalar product \langle, \rangle in this vector space \mathbb{C}^{N_v} (where N_v is the number of vertices in G) that appears in (A.5) is the Euclidean one in this basis and the matrix elements $M_{v_i v_j}$ of the adjacency matrix M vanish unless vertex v_i is connected to vertex v_j . Now since M is hermitian we have,

$$\sum_{\langle v_l v_i \rangle} \mu_{v_l} = \langle \hat{e}_{v_i} M, v_{pf} \rangle = \langle \hat{e}_{v_i}, M v_{pf} \rangle = \beta \mu_{v_i} \quad (\text{A.6})$$

¹⁵In the drawing below we omit the path indices.

¹⁶These expressions were given by A. Ocneanu in various seminars (unpublished)

replacing (A.6) in (A.4) and noting that for ADE graphs

$$\beta = 2 \cos \pi/N \tag{A.7}$$

the second term in (A.4) cancels since $\epsilon^2 + \bar{\epsilon}^2 + \beta = 0$ leading to (A.2). \square

Furthermore, regarding the uniqueness of connections we have the following result whose proof is a simple check.

Proposition A.2 (Gauge freedom). *If,*

$$\begin{array}{ccc} v_i & \longrightarrow & v_l \\ \downarrow & f & \downarrow \\ v_k & \longrightarrow & v_j \end{array} \tag{A.8}$$

satisfies unitarity and reflection then,

$$\begin{array}{ccc} v_i & \longrightarrow & v_l \\ \downarrow & f' & \downarrow \\ v_k & \longrightarrow & v_j \end{array} = e^{i(\alpha_{v_i} + \alpha_{v_j} - \alpha_{v_k} - \alpha_{v_l})} \begin{array}{ccc} v_i & \longrightarrow & v_l \\ \downarrow & f & \downarrow \\ v_k & \longrightarrow & v_j \end{array} \tag{A.9}$$

also satisfies them with $\alpha_{v_i}, \alpha_{v_j}, \alpha_{v_k}, \alpha_{v_l}$ real numbers associated to the vertices of G .

Example A.3 (The case of A_3). *From example (2.1) we know that $\beta = \sqrt{2}$ for A_3 hence taking into account (A.7) we have $N = 4$ for this case. Thus $\epsilon = e^{i5\pi/8}$. Using (A.1) and making a gauge transformation as in (A.9) with $\alpha_0 = \frac{15}{16}\pi$, $\alpha_1 = 0$, $\alpha_2 = \frac{7}{16}\pi$ you get the results appearing in (2.13).*

B The scalar product.

We first remark that the scalar product (3.5) can be written as,

$$\langle E_{\alpha\beta}^x | E_{\gamma\delta}^y \rangle = Tr[(E_{\alpha\beta}^x)^* E_{\gamma\delta}^y] \tag{B.1}$$

this can easily be verified by using (3.4), (3.1) and (3.6). Using (B.1) we can write the scalar product (3.21) as,

$$\langle \rho \otimes \rho' | \xi \otimes \xi' \rangle = \sum_{x \text{ irrep.}} \langle \rho \otimes \rho' | \xi \otimes \xi' \rangle_x \tag{B.2}$$

where in general for a representation R irreducible or not,

$$\langle \rho \otimes \rho' | \xi \otimes \xi' \rangle_R = Tr_R[(\Phi^R(\rho \otimes \rho'))^\dagger \Phi^R(\xi \otimes \xi')] \tag{B.3}$$

In order to prove proposition 3.2 we first show that,

Proposition B.1.

$$\langle \rho \otimes \rho' | \xi \otimes \xi' \rangle_{R \propto} \delta_{\rho\xi} \delta_{\rho'\xi'} \quad (\text{B.4})$$

with $R = f_1 \otimes \cdots \otimes f_n$ for any n , where f_1, \dots, f_n can be any of the fundamentals.

Proof. We will use induction. For $R = f$ (B.4) holds. This can be seen by recalling the expression of the scalar product (B.3) in terms of cells,

$$\langle \rho \otimes \rho' | \xi \otimes \xi' \rangle_f = \sum_{\eta, \eta'} \overbrace{\begin{array}{ccc} \rho & & \\ \eta \downarrow & \xrightarrow{n} & \eta' \downarrow \\ \rho' & & \end{array}} \quad \begin{array}{ccc} \xi & & \\ \eta \downarrow & \xrightarrow{n} & \eta' \downarrow \\ \xi' & & \end{array} \quad (\text{B.5})$$

so that given for example ξ and ξ' the side path η and η' are uniquely determined¹⁷ thus determining ρ and ρ' to be equal to ξ and ξ' . The constant of proportionality being the modulus of the cell. The induction hypothesis is,

$$\begin{aligned} & \langle \rho \otimes \rho' | \xi \otimes \xi' \rangle_{f_1 \otimes \cdots \otimes f_{n-1}} = \\ & = \text{Tr}_{f_1 \otimes \cdots \otimes f_{n-1}} [(\Phi^{f_1 \otimes \cdots \otimes f_{n-1}}(\rho \otimes \rho'))^\dagger \Phi^{f_1 \otimes \cdots \otimes f_{n-1}}(\xi \otimes \xi')] \\ & = \delta_{\rho\xi} \delta_{\rho'\xi'} \end{aligned} \quad (\text{B.6})$$

we prove it for the case of $R = f_1 \otimes \cdots \otimes f_n (= f_1 \cdots f_n)$ we have,

$$\begin{aligned} & \langle \rho \otimes \rho' | \xi \otimes \xi' \rangle_{f_1 \cdots f_n} \text{Tr}_{f_1 \cdots f_n} [(\Phi^{f_1 \cdots f_n}(\rho \otimes \rho'))^\dagger \Phi^{f_1 \cdots f_n}(\xi \otimes \xi')] = \\ & = \sum_{\lambda, \omega} \text{Tr}_{f_1 \cdots f_n} [\Phi^{f_1 \cdots f_{n-1}}(\lambda \otimes \rho') \Phi^{f_n}(\rho \otimes \lambda)^\dagger \Phi^{f_1 \cdots f_{n-1}}(\omega \otimes \xi') \Phi^{f_n}(\xi \otimes \omega)] \\ & = \sum_{\lambda, \omega} \text{Tr}_{f_n} [(\Phi^{f_n}(\rho \otimes \lambda))^\dagger \Phi^{f_n}(\xi \otimes \omega)] \times \\ & \quad \times \text{Tr}_{f_1 \cdots f_{n-1}} [(\Phi^{f_1 \cdots f_{n-1}}(\lambda \otimes \rho'))^\dagger \Phi^{f_1 \cdots f_{n-1}}(\omega \otimes \xi')] \\ & = \sum_{\lambda, \omega} \delta_{\rho\xi} \delta_{\lambda\omega} \delta_{\rho'\xi'} \propto \delta_{\rho\xi} \delta_{\rho'\xi'} \end{aligned} \quad (\text{B.7})$$

□

Now since any representation appears in the decomposition of tensor products of the fundamentals, proposition B.1 implies proposition 3.2.

¹⁷Recall that for the fundamentals all the paths involved in the corresponding cells are of length 1.

C Morphism property of the coproduct.

Proposition C.1.

$$\begin{aligned} & \Phi_{\alpha_1\beta_1}^{x_1} \otimes \Phi_{\alpha_2\beta_2}^{x_2} [\Delta((\xi_1 \otimes \xi'_1) \cdot (\xi_2 \otimes \xi'_2))] = \\ & = \Phi_{\alpha_1\beta_1}^{x_1} \otimes \Phi_{\alpha_2\beta_2}^{x_2} [\Delta((\xi_1 \otimes \xi'_1)) \cdot \Delta((\xi_2 \otimes \xi'_2))] \end{aligned} \quad (\text{C.1})$$

Proof. We split the proof in some steps,

(i)

$$\begin{aligned} & \Phi_{\alpha_1\beta_1}^{x_1} \otimes \Phi_{\alpha_2\beta_2}^{x_2} (\Delta((\xi_1 \otimes \xi'_1)) \cdot \Delta((\xi_2 \otimes \xi'_2))) = \\ & = \Phi_{\alpha_1\beta_1}^{x_1} \otimes \Phi_{\alpha_2\beta_2}^{x_2} (\Delta_{\mathcal{P}}((\xi_1 \otimes \xi'_1)) \star \Delta_{\mathcal{P}}((\xi_2 \otimes \xi'_2))) \end{aligned} \quad (\text{C.2})$$

Proof.

$$\begin{aligned} & \Phi_{\alpha_1\beta_1}^{x_1} \otimes \Phi_{\alpha_2\beta_2}^{x_2} (\Delta((\xi_1 \otimes \xi'_1)) \cdot \Delta((\xi_2 \otimes \xi'_2))) = \\ & = \sum_{\xi_i, \xi_j} \Phi_{\alpha_1\beta_1}^{x_1} \otimes \Phi_{\alpha_2\beta_2}^{x_2} (\xi_1 \otimes \xi^i)_{\mathcal{E}} \otimes (\xi_i \otimes \xi'_1)_{\mathcal{E}} \cdot (\xi_2 \otimes \xi^j)_{\mathcal{E}} \otimes (\xi_j \otimes \xi'_2)_{\mathcal{E}} \\ & = \sum_{\xi_i, \xi_j} \Phi_{\alpha_1\beta_1}^{x_1} ((\xi_1 \otimes \xi^i)_{\mathcal{E}} \cdot (\xi_2 \otimes \xi^j)_{\mathcal{E}}) \Phi_{\alpha_2\beta_2}^{x_2} ((\xi_i \otimes \xi'_1)_{\mathcal{E}} \cdot (\xi_j \otimes \xi'_2)_{\mathcal{E}}) \\ & = \sum_{\xi_i, \xi_j} \Phi_{\alpha_1\beta_1}^{x_1} ((\xi_1 \otimes \xi^i)_{\mathcal{P}} \star (\xi_2 \otimes \xi^j)_{\mathcal{P}}) \Phi_{\alpha_2\beta_2}^{x_2} ((\xi_i \otimes \xi'_1)_{\mathcal{P}} \star (\xi_j \otimes \xi'_2)_{\mathcal{P}}) \\ & = \Phi_{\alpha_1\beta_1}^{x_1} \otimes \Phi_{\alpha_2\beta_2}^{x_2} (\Delta_{\mathcal{P}}((\xi_1 \otimes \xi'_1)) \star \Delta_{\mathcal{P}}((\xi_2 \otimes \xi'_2))) \end{aligned} \quad (\text{C.3})$$

where we have defined $\Delta_{\mathcal{P}}$ by,

$$\Delta_{\mathcal{P}}(\xi \otimes \xi') = \sum_{\xi_i} (\xi \otimes \xi^i)_{\mathcal{P}} \otimes (\xi_i \otimes \xi')_{\mathcal{P}} \quad (\text{C.4})$$

□

(ii)

$$\Delta_{\mathcal{P}}((\xi_1 \otimes \xi'_1)) \star \Delta_{\mathcal{P}}((\xi_2 \otimes \xi'_2)) = \Delta_{\mathcal{P}}((\xi_1 \otimes \xi'_1) \star (\xi_2 \otimes \xi'_2)) \quad (\text{C.5})$$

Proof.

$$\begin{aligned} & \Delta_{\mathcal{P}}((\xi_1 \otimes \xi'_1)) \star \Delta_{\mathcal{P}}((\xi_2 \otimes \xi'_2)) = \\ & = \sum_{\xi_i, \xi_j, \#\xi_i = \#\xi_1, \#\xi_j = \#\xi_2} [(\xi_1 \otimes \xi^i)_{\mathcal{P}} \otimes (\xi_i \otimes \xi'_1)_{\mathcal{P}}] \star [(\xi_2 \otimes \xi^j)_{\mathcal{P}} \otimes (\xi_j \otimes \xi'_2)_{\mathcal{P}}] \\ & = \sum_{\xi_i, \xi_j, \#\xi_i = \#\xi_1, \#\xi_j = \#\xi_2} [(\xi_1 \star \xi_2) \otimes (\xi^i \star \xi^j)]_{\mathcal{P}} \otimes [(\xi_i \star \xi_j) \otimes (\xi'_1 \star \xi'_2)]_{\mathcal{P}} \\ & = \Delta_{\mathcal{P}}((\xi_1 \otimes \xi'_1) \star (\xi_2 \otimes \xi'_2)) \end{aligned} \quad (\text{C.6})$$

where in the last equality we have used that a summation over elementary paths of length n with fixed extreme vertex is equal to the summation over any pair of paths of length n_1 and n_2 of the concatenation of them if $n_1 + n_2 = n$ and the starting(ending) vertex of the first(last)path coincide with the ones of the length n path. \square

Finally we have,

(iii)

$$\begin{aligned} & \Phi_{\alpha_1\beta_1}^{x_1} \otimes \Phi_{\alpha_2\beta_2}^{x_2} (\Delta_{\mathcal{P}}((\xi_1 \otimes \xi'_1) \star (\xi_2 \otimes \xi'_2))) = \\ & = \Phi_{\alpha_1\beta_1}^{x_1} \otimes \Phi_{\alpha_2\beta_2}^{x_2} (\Delta((\xi_1 \otimes \xi'_1) \cdot (\xi_2 \otimes \xi'_2))) \end{aligned} \quad (\text{C.7})$$

Proof.

$$\begin{aligned} & \Phi_{\alpha_1\beta_1}^{x_1} \otimes \Phi_{\alpha_2\beta_2}^{x_2} (\Delta_{\mathcal{P}}((\xi_1 \otimes \xi'_1) \star (\xi_2 \otimes \xi'_2))) = \\ & = \Phi_{\alpha_1\beta_1}^{x_1} \otimes \Phi_{\alpha_2\beta_2}^{x_2} (\Delta_{\mathcal{P}} \sum_{x,\alpha,\beta} \alpha \begin{array}{c} \xrightarrow{\xi_1 \star \xi_2} \\ \downarrow f \\ \xrightarrow{\xi'_1 \star \xi'_2} \\ \downarrow \end{array} \beta E_{\alpha\beta}^x) \\ & = \Phi_{\alpha_1\beta_1}^{x_1} \otimes \Phi_{\alpha_2\beta_2}^{x_2} (\sum_{x,\alpha,\beta,\xi,\xi'} \alpha \begin{array}{c} \xrightarrow{\xi_1 \star \xi_2} \\ \downarrow x \\ \xrightarrow{\xi'_1 \star \xi'_2} \\ \downarrow \end{array} \beta \alpha \begin{array}{c} \xrightarrow{\xi} \\ \downarrow x^{-1} \\ \xrightarrow{\xi'} \\ \downarrow \end{array} \beta \Delta_{\mathcal{P}}(\xi \otimes \xi')) \\ & = \sum_{x,\alpha,\beta,\xi,\xi',\xi_i} \alpha \begin{array}{c} \xrightarrow{\xi_1 \star \xi_2} \\ \downarrow x \\ \xrightarrow{\xi'_1 \star \xi'_2} \\ \downarrow \end{array} \beta \alpha \begin{array}{c} \xrightarrow{\xi} \\ \downarrow x^{-1} \\ \xrightarrow{\xi'} \\ \downarrow \end{array} \beta \Phi_{\alpha_1\beta_1}^{x_1}(\xi \otimes \xi^i) \Phi_{\alpha_2\beta_2}^{x_2}(\xi_i \otimes \xi') \\ & = \sum_{\xi,\xi',\xi_i} P_{\xi'_1\xi'_2\xi'}^{\xi_1\xi_2\xi} \alpha_1 \begin{array}{c} \xrightarrow{\xi} \\ \downarrow x_1 \\ \xrightarrow{\xi^i} \\ \downarrow \end{array} \beta_1 \alpha_2 \begin{array}{c} \xrightarrow{\xi_i} \\ \downarrow x_2 \\ \xrightarrow{\xi'} \\ \downarrow \end{array} \beta_2 \\ & = \Phi_{\alpha_1\beta_1}^{x_1} \otimes \Phi_{\alpha_2\beta_2}^{x_2} (\Delta((\xi_1 \otimes \xi'_1) \cdot (\xi_2 \otimes \xi'_2))) \end{aligned} \quad (\text{C.8})$$

\square

\square

D The trivial representation $\mathbf{0}$.

Consider a map $\Phi_{\alpha_0\beta_0}^{\mathbf{0}} : \text{End}^{gr}(\mathcal{E}) \rightarrow \mathbb{C}$ defined such that for any irrep x of the DTA and indices α_x, β_x in x there exist unique subscripts $\alpha_0\beta_0$ in $\mathbf{0}$ such

that,

$$\Phi_{\alpha_0 \star \alpha_x, \beta_0 \star \beta_x}^{\mathbf{0} \otimes x}(\xi \otimes \xi') = \Phi_{\alpha_x, \beta_x}^x(\xi \otimes \xi') \quad (\text{D.1})$$

using eqs. (4.3) and (2.5), eq. (D.1) implies,

$$\sum_{\lambda} \alpha_0 \begin{array}{c} \xrightarrow{\xi} \\ \mathbf{0} \\ \xrightarrow{\lambda} \end{array} \beta_0 = \alpha_x \begin{array}{c} \xrightarrow{\lambda} \\ x \\ \xrightarrow{\xi'} \end{array} \beta_x \alpha_x \begin{array}{c} \xrightarrow{\xi} \\ x \\ \xrightarrow{\xi'} \end{array} \beta_x \quad (\text{D.2})$$

multiplying this last equation by,

$$\alpha_x \begin{array}{c} \xrightarrow{\rho} \\ x^{-1} \\ \xrightarrow{\rho'} \end{array} \beta_x$$

summing up over x, α_x, β_x and using (3.13) we obtain,

$$\alpha_0 \begin{array}{c} \xrightarrow{\xi} \\ \mathbf{0} \\ \xrightarrow{\xi'} \end{array} \beta_0 = \delta_{\xi \xi'} \quad (\text{D.3})$$

Note that this last equation implies that for $\xi_1, \xi'_1, \xi_2, \xi'_2$ such that $\#\xi_1 = \#\xi'_1$, $\#\xi_2 = \#\xi'_2$ and $r(\xi_1) = s(\xi_2)$, $r(\xi'_1) = s(\xi'_2)$ we have,

$$\alpha_0 \begin{array}{c} \xrightarrow{\xi_1 \star \xi_2} \\ \mathbf{0} \\ \xrightarrow{\xi'_1 \star \xi'_2} \end{array} \beta_0 = \alpha_0 \begin{array}{c} \xrightarrow{\xi_1} \\ \mathbf{0} \\ \xrightarrow{\xi'_1} \end{array} \gamma_0 \gamma_0 \begin{array}{c} \xrightarrow{\xi_2} \\ \mathbf{0} \\ \xrightarrow{\xi'_2} \end{array} \beta_0 \quad (\text{D.4})$$

To show that this is indeed a representation note that eq.(3.15) implies,

$$\sum_{\xi, \xi'} \alpha \begin{array}{c} \xrightarrow{\xi} \\ \mathbf{0}^{-1} \\ \xrightarrow{\xi'} \end{array} \beta \alpha_0 \begin{array}{c} \xrightarrow{\xi} \\ \mathbf{0} \\ \xrightarrow{\xi'} \end{array} \beta_0 = \sum_{\xi} \alpha \begin{array}{c} \xrightarrow{\xi} \\ \mathbf{0}^{-1} \\ \xrightarrow{\xi} \end{array} \beta = \delta_{\alpha \alpha_0} \delta_{\beta \beta_0} \quad (\text{D.5})$$

hence,

$$\Phi_{\alpha_0 \beta_0}^{\mathbf{0}}(E_{\alpha \beta}^{\mathbf{0}}) = \sum_{\xi \xi'} \alpha \begin{array}{c} \xrightarrow{\xi} \\ \mathbf{0}^{-1} \\ \xrightarrow{\xi'} \end{array} \beta \Phi_{\alpha_0 \beta_0}^{\mathbf{0}}(\xi \otimes \xi') = \sum_{\xi} \alpha \begin{array}{c} \xrightarrow{\xi} \\ \mathbf{0}^{-1} \\ \xrightarrow{\xi} \end{array} \beta = \delta_{\alpha \alpha_0} \delta_{\beta \beta_0} \quad (\text{D.6})$$

hence $\Phi_{\alpha_0 \beta_0}^{\mathbf{0}}(E_{\alpha \beta}^{\mathbf{0}})$ is of the form (3.7) thus being a homomorphism for the matrix product to which the product in the DTA reduces in this basis.

References

- [1] J. Böckenhauer and D. Evans, *Modular invariants, graphs and α induction for nets of subfactors II*. Commun. Math. Phys. 200, 57-103 (1999).
- [2] J. Böckenhauer and D. Evans, *Modular invariants from subfactors* D. Evans lectures at Bariloche Summer School, Argentina, Jan 2000. AMS Contemporary Mathematics 294, R. Coquereaux, A. García and R. Trinchero eds.
- [3] G. Böhm and K. Szlachányi, *A coassociative C^* -quantum group with non-integral dimensions*, Lett. Math. Phys. 38 (1996), no. 4, 437-456, math.QA/9509008 (1995).
- [4] A. Cappelli, C. Itzykson and J.B.Zuber, Nucl. Phys. **B 280** (1987) 445-465; Comm.Math. Phys. **113** (1987) 1-26.
- [5] A. Cappelli, C. Itzykson et J.-B. Zuber, *The ADE classification of minimal and $A_1^{(1)}$ conformal invariant theories*. Commu. Math. Phys. 13,1 (1987).
- [6] R. Coquereaux, *Notes on the quantum tetrahedron*, Moscow Math. J. vol2, n1, Jan.-March 2002, 1-40, hep-th/0011006.
- [7] R. Coquereaux, G. Schieber, *Twisted partition functions for ADE boundary conformal field theories and Ocneanu algebras of quantum symmetries*, J. of Geom. and Phys. 781 (2002), 1-43, hep-th/0107001.
- [8] R. Coquereaux and M. Huerta, *Torus structure on graphs and twisted partition functions for minimal and affine models*. J. of Geom. and Phys., Vol 48, No4, (2003) 580-634, hep-th/0301215.
- [9] R. Coquereaux, *The A_2 Ocneanu quantum groupoid*, Cocoyoc conference, Mexico, hep-th/0311151.
- [10] R. Coquereaux, G. Schieber and R. Trinchero, in preparation.
- [11] D. Evans and Y. Kawahigashi, *Quantum symmetries and operator algebras*, Clarendon Press, 1998.
- [12] P. Di Francesco and J.-B. Zuber, in *Recent developments in Conformal Field Theory*, Trieste Conference, 1989, S. Randjbar-Daemi, E. Sezgin and J.-B. Zuber eds, World Sci, 1990. P. Di Francesco, Int. J. of Mod. Phys. A7, 92, 407.

- [13] P. Di Francesco and J.-B. Zuber, *SU(N) lattice integrable models associated with graphs*, Nucl. Phys. B338, 602 (1990).
- [14] D. Nikshych and L. Vainerman, *Algebraic versions of a finite-dimensional quantum groupoid*, in Lecture Notes in Pure and Appl. Math. Vol. 209 (2000), 189-220, mathQA/9808054.
- [15] D. Nikshych and L. Vainerman, *A Characterization of depth 2 subfactors of II_1 factors*, Journal of Funct. Analysis, Vol. 171, n.2 (2000), 278-307, math.QA/981002.
- [16] D. Nikshych and L. Vainerman, *Finite quantum groupoids and their applications*, in New directions in Hopf algebras, 211–262, Math. Sci. Res. Inst. Publ., 43, Cambridge Univ. Press, Cambridge, 2002, mathQA/0006057.
- [17] F. Nill, *Axioms for weak bialgebras*, math.QA/9805104.
- [18] A. Ocneanu; *Quantized groups, string algebras and Galois theory for algebras*. Warwick (1987), in Operator algebras and applications, Lond Math Soc Lectures note Series 136, CUP (1988).
- [19] A. Ocneanu, *Quantum cohomology, quantum groupoids and subfactors*, talk presented at the First Caribbean School of Mathematics and Theoretical Physics, Saint François, Guadeloupe (1993).
- [20] A. Ocneanu, *Paths on Coxeter diagrams: from Platonic solids and singularities to minimal models and subfactors*. Notes taken by S. Goto, Fields Institute Monographs, AMS 1999, Rajarama Bhat et al, eds.
- [21] A. Ocneanu, *Quantum symmetries for SU(3) CFT-models* A. Ocneanu, The Classification of subgroups of quantum SU(N), Lectures at Bariloche Summer School, Argentina, Jan.2000, AMS Contemporary Mathematics 294, R. Coquereaux, A. García and R. Trinchero eds.
- [22] V. Pasquier, *Two-dimensional critical systems labelled by Dynkin diagrams* Nucl.Phys. B285, 162 (1987).
- [23] V.B. Petkova and J.-B. Zuber, *The many faces of Ocneanu cells*, Nucl. Phys. B 603(2001) 449, hep-th/0101151.
- [24] Ph. Roche, *Ocneanu cell calculus and integrable lattice models* Commun.Math.Phys. 127, 395-424 (1990).
- [25] G. Schieber, *L'algèbre des symétries quantiques d'Ocneanu et la classification des systèmes conformes à 2D*, PhD thesis (available in French and in Portuguese), UP (Marseille) and UFRJ (Rio de Janeiro). Sept. 2003.

- [26] J.-B. Zuber, *CFT, ADE and all that* Lectures at Bariloche Summer School, Argentina, Jan 2000. AMS Contemporary Mathematics 294, R. Coquereaux, A. García and R. Trinchero eds., hep-th/0006151.
- [27] J.-B. Zuber, *Graphs, algebras, conformal field theories and integrable lattice models*, Nucl. Phys. B (Proc. Suppl.) 18B (1990) 313 - 326, *Generalized Dynkin diagrams, root systems and their folding*, YITP International Workshop on Recent Developments in QCD and Hadron Physics, Kyoto, Japan, 16-18 Dec 1996, hep-th/9707046.