# RESEARCH



# On the asymptotic behavior of fourth-order functional differential equations

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# Abstract

The aim of this work is to study asymptotic properties of a class of fourth-order delay differential equations. Our results in this paper not only generalize some previous results, but also improve the earlier ones. Examples are considered to elucidate the main results.

Keywords: asymptotic behavior; fourth-order; delay differential equations

# 1 Introduction

This paper is concerned with the oscillatory behavior of solutions of nonlinear fourthorder differential equations of the type

$$\left(r(\tau)\left(x^{\prime\prime\prime\prime}(\tau)\right)^{\alpha}\right)^{\prime} + \int_{a}^{b} q(\tau,\xi) f\left(x\left(g(\tau,\xi)\right)\right) d\sigma(\xi) = 0, \tag{1.1}$$

where the following conditions are satisfied:

- (A<sub>1</sub>)  $r \in C([\tau_0, \infty), (0, \infty)), r'(\tau) > 0$  and  $\alpha$  is a quotient of odd positive integers;
- (A<sub>2</sub>)  $q, g \in C([\tau_0, \infty) \times [a, b], \mathbb{R}), q(\tau, \xi) \ge 0, q(\tau, \xi)$  is not zero on any half line  $[\tau_{\lambda}, \infty) \times [a, b], \tau_{\lambda} \ge \tau_0, g(\tau, \xi) \le \tau$  for  $\tau \ge \tau_0$  and  $\xi \in [a, b], g(\tau, \xi)$  is continuous, nondecreasing with respect to  $\xi$  and  $\lim_{\tau \to \infty} g(\tau, \xi) = \infty$ ;
- (A<sub>3</sub>)  $\sigma \in C([a, b], \mathbb{R})$ ,  $\sigma$  is nondecreasing and the integral of equation (1.1) is in the Riemann-Stieltjes sense;

and the function  $f \in C(\mathbb{R}, \mathbb{R})$  satisfies one of the following conditions:

$$\begin{array}{l} (\mathrm{S}_1) \ f(x)/x^{\alpha} \geq k_1 > 0 \ \text{for} \ x \neq 0; \\ (\mathrm{S}_2) \ f'(x)/|f(x)|^{\frac{1-\alpha}{\alpha}} \geq k_2 > 0 \ \text{for} \ x \neq 0 \ \text{and} \ f(uv) \geq u^{\alpha}f(v) \ \text{for} \ uv > 0. \end{array}$$

By a solution of equation (1.1), we mean a function  $x(\tau) \in C[\tau_x, \infty)$ ,  $\tau_x \ge \tau_0$  such that  $r(\tau)(x'''(\tau))^{\alpha}$  is continuously differentiable for all  $\tau \ge \tau_x$  and satisfies equation (1.1) for all  $\tau \in [\tau_x, \infty)$ . Here, we consider only proper solutions  $x(\tau)$  to equation (1.1) with property  $\sup\{|x(\tau)|: \tau \ge \tau\} > 0$  for any  $\tau \ge \tau_x$ . A solution of equation (1.1) is called oscillatory if it has arbitrary large zeros, otherwise it is called nonoscillatory.

In recent years there has been much research activity concerning the oscillation behavior of solutions of nonlinear differential equations (see [1-21]). In the last few years, many



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papers have appeared on the oscillatory theory of fourth-order differential equations (see [2, 16, 22–25]).

The aim of this paper is to study the oscillatory behavior of the solutions of nonlinear fourth-order differential equations (1.1) under the assumption

$$\int_{\tau_0}^{\infty} \frac{1}{r^{1/\alpha}(s)} \, ds < \infty,\tag{1.2}$$

and we consider the function f with and without monotonicity. The results obtained essentially generalize the results from Zhang [24] and also improve some results from Baculykova [2]. Examples are provided to illustrate new results.

In order to discuss our main results, we need the following lemmas.

**Lemma 1.1** ([15]) If the function y satisfies  $y^{(i)}(\tau) > 0$ , i = 0, 1, ..., n, and  $y^{(n+1)}(\tau) < 0$ , then

$$\frac{y(\tau)}{\tau^n/n!} \geq \frac{y'(\tau)}{\tau^{n-1}/(n-1)!}.$$

**Lemma 1.2** ([1]) Let  $y \in C^n([\tau_0, \infty), (0, \infty))$ . Assume that  $y^{(n)}(\tau)$  is of fixed sign and not identically zero on  $[\tau_0, \infty)$  and that there exists  $\tau_1 \ge \tau_0$  such that  $y^{(n-1)}(\tau)y^{(n)}(\tau) \le 0$  for all  $\tau \ge \tau_1$ . If  $\lim_{\tau \to \infty} y(\tau) \ne 0$ , then for every  $\mu \in (0, 1)$  there exists  $\tau_\mu \ge \tau_1$  such that

$$y(\tau) \geq \frac{\mu}{(n-1)!} \tau^{n-1} \left| y^{(n-1)}(\tau) \right| \quad for \ \tau \geq \tau_{\mu}.$$

## 2 Main results

In this section, we establish new oscillation criteria for solutions of equation (1.1). For the sake of convenience, we insert the following notation:

$$\begin{split} R_1(\tau) &= \int_{\tau}^{\infty} \frac{1}{r^{1/\alpha}(s)} \, ds, \qquad R_i(\tau) = \int_{\tau}^{\infty} R_{i-1}(u) \, du, \quad i = 2, 3, \\ Q(\tau) &= \int_{a}^{b} q(\tau, \xi) \, d\sigma(\xi), \end{split}$$

and  $F_{+}(\tau) = \max\{0, F(\tau)\}.$ 

**Lemma 2.1** If  $x(\tau)$  is an eventually positive three times continuously differentiable function such that  $r(\tau)x'''(\tau)$  is continuously differentiable and  $(r(\tau)x'''(\tau))' \leq 0$  for large t, then one of the following cases holds for large t:

 $\begin{array}{ll} (C_1) & x'(\tau) > 0, & x''(\tau) > 0, & x'''(\tau) > 0, \\ (C_2) & x'(\tau) > 0, & x''(\tau) < 0, & x'''(\tau) > 0, \\ (C_3) & x'(\tau) < 0, & x''(\tau) > 0, & x'''(\tau) < 0, \\ (C_4) & x'(\tau) > 0, & x''(\tau) > 0, & x'''(\tau) < 0. \end{array}$ 

The proof is immediate and hence is omitted.

**Theorem 2.1** Assume that (1.2) and (S<sub>1</sub>) hold. If there exist continuously differentiable functions  $\rho, \vartheta \in C([\tau_0, \infty), (0, \infty))$  such that

$$\int_{\tau_0}^{\infty} \left( k_1 \rho(s) Q(s) \left( \frac{g(s,a)}{s} \right)^{3\alpha} - \frac{2^{\alpha}}{(\alpha+1)^{\alpha+1}} \frac{r(s)(\rho'_+(s))^{\alpha+1}}{(\mu_1 s^2 \rho(s))^{\alpha}} \right) ds = \infty,$$
(2.1)

$$\int_{\tau_0}^{\infty} \left( \vartheta\left(u\right) \int_u^{\infty} \left[ \frac{k_1}{r(v)} \int_v^{\infty} \int_a^b q(s,\xi) \frac{g^{\alpha}(s,\xi)}{s^{\alpha}} \, d\sigma\left(\xi\right) \, ds \right]^{1/\alpha} \, dv - \frac{\vartheta_+^{\prime 2}(u)}{4\vartheta\left(u\right)} \right) du = \infty, \quad (2.2)$$

$$\int_{\tau_0}^{\infty} \left( k_1 R_3^{\alpha}(s) Q(s) - \left(\frac{\alpha}{\alpha+1}\right)^{\alpha+1} \frac{R_2(s)}{R_3(s)} \right) ds = \infty$$
(2.3)

and

$$\int_{\tau_0}^{\infty} \left(\frac{k_1 \mu_2^{\alpha}}{2^{\alpha}} R_1^{\alpha}(s) g^{2\alpha}(s, a) Q(s) - \left(\frac{\alpha}{\alpha+1}\right)^{\alpha+1} \frac{1}{r^{1/\alpha}(s) R_1(s)}\right) ds = \infty$$
(2.4)

for some  $\mu_1, \mu_2 \in (0, 1)$ , then every solution of (1.1) is oscillatory.

*Proof* Let *x* be a nonoscillatory solution of equation (1.1) on the interval  $[\tau_0, \infty)$ . Without loss of generality, we may assume that  $x(\tau) > 0$ . From Lemma 2.1, there exists  $\tau_1 \ge \tau_0$  such that  $x(\tau)$  has one of the four cases (C<sub>1</sub>)-(C<sub>4</sub>) for  $\tau \ge \tau_1$ . *For Case* (C<sub>1</sub>), we define

$$\omega(\tau) = \rho(\tau) \frac{r(\tau)(x''(\tau))^{\alpha}}{x^{\alpha}(\tau)}.$$

Then  $\omega(\tau) > 0$ . By differentiating, we obtain

$$\omega'(\tau) = \frac{\rho'(\tau)}{\rho(\tau)}\omega(\tau) + \rho(\tau)\frac{(r(\tau)(x'''(\tau))^{\alpha})'}{x^{\alpha}(\tau)} - \alpha\rho(\tau)\frac{r(\tau)(x'''(\tau))^{\alpha}}{x^{\alpha+1}(\tau)}x'(\tau).$$
(2.5)

It follows from Lemma 1.2 that

$$x'(\tau) \ge \frac{\mu}{2} \tau^2 x'''(\tau)$$
 (2.6)

for all  $\mu \in (0, 1)$  and every sufficiently large  $\tau$ . From (1.1), (A<sub>2</sub>) and (S<sub>1</sub>), we see that

$$(r(\tau)(x'''(\tau))^{\alpha})' = -\int_{a}^{b} q(\tau,\xi)f(x(g(\tau,\xi))) d\sigma(\xi)$$
  
$$\leq -k_{1}Q(\tau)x^{\alpha}(g(\tau,a)).$$
(2.7)

Thus, by (2.5), (2.6) and (2.7), we get

$$\omega'(\tau) \le \frac{\rho'(\tau)}{\rho(\tau)}\omega(\tau) - k_1\rho(\tau)Q(\tau)\frac{x^{\alpha}(g(\tau,a))}{x^{\alpha}(\tau)} - \frac{\alpha\mu}{2}\tau^2\rho(\tau)r(\tau)\left(\frac{x'''(\tau)}{x(\tau)}\right)^{\alpha+1}.$$
(2.8)

From Lemma 1.1, we have that

$$x(\tau)\geq \frac{\tau}{3}x'(\tau).$$

Integrating this inequality from  $g(\tau, a)$  to  $\tau$ , we get

$$\frac{x(g(\tau,a))}{x(\tau)} \ge \frac{g^3(\tau,a)}{\tau^3},\tag{2.9}$$

which with (2.8) gives

$$\omega'(\tau) \le \frac{\rho'(\tau)}{\rho(\tau)}\omega(\tau) - k_1\rho(\tau)Q(\tau) \left(\frac{g(\tau,a)}{\tau}\right)^{3\alpha} - \frac{\alpha\mu}{2} \frac{\tau^2}{(\rho(\tau)r(\tau))^{1/\alpha}} \omega^{\frac{\alpha+1}{\alpha}}(\tau).$$
(2.10)

By using the inequality

$$Bz - Az^{\frac{1+\alpha}{\alpha}} \le \frac{\alpha^{\alpha}}{(\alpha+1)^{\alpha+1}} B^{\alpha+1} A^{-\alpha} \quad \text{for } A, B > 0 \text{ and } z \ge 0,$$
(2.11)

with  $A = \frac{\alpha \mu}{2} \frac{\tau^2}{(\rho(\tau)r(\tau))^{1/\alpha}}$ ,  $B = \frac{\rho'(\tau)}{\rho(\tau)}$  and  $z = \omega$ , we get

$$\omega'(\tau) \leq -k_1 \rho(\tau) Q(\tau) \left(\frac{g(\tau,a)}{\tau}\right)^{3\alpha} + \frac{2^{\alpha}}{(\alpha+1)^{\alpha+1}} \frac{r(\tau)(\rho'_+(\tau))^{\alpha+1}}{(\mu\tau^2\rho(\tau))^{\alpha}}.$$

This implies that

$$\int_{\tau_1}^{\tau} \left( k_1 \rho(s) Q(s) \left( \frac{g(s,a)}{s} \right)^{3\alpha} - \frac{2^{\alpha}}{(\alpha+1)^{\alpha+1}} \frac{r(s)(\rho'_+(\tau))^{\alpha+1}}{(\mu s^2 \rho(s))^{\alpha}} \right) ds \le \omega(\tau_1)$$

for every  $\mu \in (0,1)$  and all sufficiently large  $\tau$ , which contradicts (2.1).

*Consider Case* (C<sub>2</sub>) *holds.* From Lemma 1.1, we get that  $x(\tau) \ge \tau x'(\tau)$ , by integrating this inequality from  $g(\tau, \xi)$  to  $\tau$ , we get

$$x(g(\tau,\xi)) \geq \frac{g(\tau,\xi)}{\tau}x(\tau).$$

Hence, from  $(S_1)$ , we have

$$f(x(g(\tau,\xi))) \ge k_1 \frac{g^{\alpha}(s,\xi)}{s^{\alpha}} x^{\alpha}(\tau).$$
(2.12)

Integrating (1.1) from  $\tau$  to u and using  $x'(\tau) > 0$ , we obtain

$$r(u)(x'''(u))^{\alpha} - r(\tau)(x'''(\tau))^{\alpha} = -\int_{\tau}^{u} \int_{a}^{b} q(s,\xi) f(x(g(s,\xi))) d\sigma(\xi) ds$$
$$\leq -k_{1}x^{\alpha}(\tau) \int_{\tau}^{u} \int_{a}^{b} q(s,\xi) \frac{g^{\alpha}(s,\xi)}{s^{\alpha}} d\sigma(\xi) ds.$$

Letting  $u \to \infty$ , we see that

$$r(\tau)(x^{\prime\prime\prime}(\tau))^{\alpha} \geq k_1 x^{\alpha}(\tau) \int_{\tau}^{\infty} \int_{a}^{b} q(s,\xi) \frac{g^{\alpha}(s,\xi)}{s^{\alpha}} d\sigma(\xi) ds,$$

and so,

$$x'''(\tau) \ge x(\tau) \left[ \frac{k_1}{r(\tau)} \int_{\tau}^{\infty} \int_{a}^{b} q(s,\xi) \frac{g^{\alpha}(s,\xi)}{s^{\alpha}} d\sigma(\xi) ds \right]^{1/\alpha}.$$

Integrating again from  $\tau$  to  $\infty$ , we get

$$x''(\tau) \le -x(\tau) \int_{\tau}^{\infty} \left[ \frac{k_1}{r(\nu)} \int_{\nu}^{\infty} \int_{a}^{b} q(s,\xi) \frac{g^{\alpha}(s,\xi)}{s^{\alpha}} d\sigma(\xi) ds \right]^{1/\alpha} d\nu.$$
(2.13)

Now, we define

$$w(\tau)=\vartheta(\tau)\frac{x'(\tau)}{x(\tau)}.$$

Then  $w(\tau) > 0$  for  $\tau \ge \tau_1$ . By differentiating and using (2.13), we find

$$w'(\tau) = \frac{\vartheta'(\tau)}{\vartheta(\tau)}w(\tau) + \vartheta(\tau)\frac{x''(\tau)}{x(\tau)} - \vartheta(\tau)\frac{(x'(\tau))^2}{x^2(\tau)}$$
  
$$\leq -\vartheta(\tau)\int_{\tau}^{\infty} \left[\frac{k}{r(\nu)}\int_{\nu}^{\infty}\int_{a}^{b}q(s,\xi)\frac{g^{\alpha}(s,\xi)}{s^{\alpha}}d\sigma(\xi)\,ds\right]^{1/\alpha}d\nu$$
  
$$+\frac{\vartheta'(\tau)}{\vartheta(\tau)}w(\tau) - \frac{1}{\vartheta(\tau)}w^2(\tau).$$
(2.14)

Thus, we obtain

$$w'(\tau) \le -\vartheta(\tau) \int_{\tau}^{\infty} \left[ \frac{k}{r(\nu)} \int_{\nu}^{\infty} \int_{a}^{b} q(s,\xi) \frac{g^{\alpha}(s,\xi)}{s^{\alpha}} d\sigma(\xi) ds \right]^{1/\alpha} d\nu + \frac{(\vartheta'_{+}(\tau))^{2}}{4\vartheta(\tau)}.$$
 (2.15)

Then we get

$$\int_{\tau_1}^{\tau} \left( \vartheta(u) \int_u^{\infty} \left[ \frac{k}{r(v)} \int_v^{\infty} \int_a^b q(s,\xi) \frac{g^{\alpha}(s,\xi)}{s^{\alpha}} d\sigma(\xi) ds \right]^{1/\alpha} dv - \frac{\vartheta_+^{\prime 2}(u)}{4\vartheta(u)} \right) du$$
  
  $\leq w(\tau_1).$ 

This contradicts (2.2).

Assume that Case (C<sub>3</sub>) holds. Since  $r(\tau)(x'''(\tau))^{\alpha}$  is nonincreasing, we have that  $r(s)(x'''(s))^{\alpha} \leq r(\tau)(x'''(\tau))^{\alpha}$  for all  $s \geq \tau \geq \tau_1$ . This yields

$$x'''(s) \le [r(\tau)(x'''(\tau))^{\alpha}]^{1/\alpha} \frac{1}{r^{1/\alpha}(s)}.$$

Integrating this inequality from  $\tau$  to u, we get

$$x^{\prime\prime}(u)-x^{\prime\prime}(\tau)\leq \left[r(\tau)\big(x^{\prime\prime\prime}(\tau)\big)^{\alpha}\right]^{1/\alpha}\int_{\tau}^{u}\frac{1}{r^{1/\alpha}(s)}\,ds.$$

Letting  $u \to \infty$ , we see that

$$-x''(\tau) \le \left[r(\tau) \left(x'''(\tau)\right)^{\alpha}\right]^{1/\alpha} R_1(\tau).$$
(2.16)

By integrating the last inequality from  $\tau$  to  $\infty$ , we obtain

$$x'(\tau) \le \left[r(\tau) \left(x'''(\tau)\right)^{\alpha}\right]^{1/\alpha} R_2(\tau).$$
 (2.17)

Integrating again from  $\tau$  to  $\infty$ , we find

$$x(\tau) \ge -\left[r(\tau) \left(x'''(\tau)\right)^{\alpha}\right]^{1/\alpha} R_3(\tau).$$
(2.18)

Next, we define

$$\psi(\tau)=\frac{r(\tau)(x^{\prime\prime\prime}(\tau))^{\alpha}}{x^{\alpha}(\tau)}.$$

Thus, we see that  $\psi(\tau) < 0$  and satisfies

$$\psi'(\tau) = \frac{[r(\tau)(x'''(\tau))^{\alpha}]'}{x^{\alpha}(\tau)} - \alpha \frac{r(\tau)(x'''(\tau))^{\alpha}}{x^{\alpha+1}(\tau)} x'(\tau).$$

Hence, from (1.1), (2.17) and  $(S_1)$ , we have

$$\psi'(\tau) \leq -k_1 \int_a^b q(\tau,\xi) \frac{x^{\alpha}(g(\tau,\xi))}{x^{\alpha}(\tau)} \, d\sigma(\xi) - \alpha R_2(\tau) \psi^{\frac{1+\alpha}{\alpha}}(\tau).$$

.

Since  $g(\tau,\xi) \le \tau$  and  $x'(\tau) < 0$ , we have that  $x(g(\tau,\xi)) \ge x(\tau)$ . Therefore, we get

$$\psi'(\tau) \le -k_1 Q(\tau) - \alpha R_2(\tau) \psi^{\frac{1+\alpha}{\alpha}}(\tau).$$
(2.19)

From (2.18), we have

$$R_3^{\alpha}(\tau)\psi(\tau) \ge -1. \tag{2.20}$$

Multiplying (2.19) by  $R_3^{\alpha}(\tau)$  and integrating from  $\tau_1$  to  $\tau$ , we obtain

$$R_{3}^{\alpha}(\tau)\psi(\tau) - R_{3}^{\alpha}(\tau_{1})\psi(\tau_{1}) + \alpha \int_{\tau_{1}}^{\tau} R_{2}(s)R_{3}^{\alpha-1}(s)\psi(s) ds$$
$$+ k_{1}\int_{\tau_{1}}^{\tau} R_{3}^{\alpha}(s)Q(s) ds + \alpha \int_{\tau_{1}}^{\tau} R_{2}(s)R_{3}^{\alpha}(s)\psi^{\frac{1+\alpha}{\alpha}}(s) ds \leq 0,$$

which with (2.20) gives

$$1 + R_{3}^{\alpha}(\tau_{1})\psi(\tau_{1}) \ge k_{1} \int_{\tau_{1}}^{\tau} R_{3}^{\alpha}(s)Q(s) ds + \alpha \int_{\tau_{1}}^{\tau} R_{2}(s)R_{3}^{\alpha-1}(s) [\psi(s) + R_{3}(s)\psi^{\frac{1+\alpha}{\alpha}}(s)] ds.$$

Using inequality (2.11) with  $A = R_3$ , B = 1 and  $z = -\psi$ , we get

$$\psi(s) + R_3(s)\psi^{\frac{1+\alpha}{\alpha}}(s) \geq -\frac{\alpha^{\alpha}}{(\alpha+1)^{\alpha+1}}R_3^{-\alpha}(s).$$

It follows that

$$\int_{\tau_1}^{\tau} \left( k_1 R_3^{\alpha}(s) Q(s) - \left(\frac{\alpha}{\alpha+1}\right)^{\alpha+1} \frac{R_2(s)}{R_3(s)} \right) ds \le 1 + R_3^{\alpha}(\tau_1) \psi(\tau_1),$$

but this contradicts (2.3).

In Case (C<sub>4</sub>). In view of the proof of Case (C<sub>3</sub>), we have (2.16) holds. From Lemma 1.2, we have that  $x(\tau) \ge \frac{\mu}{2}\tau^2 x''(\tau)$  for all  $\mu \in (0, 1)$  and every sufficiently large  $\tau$ . Thus, from (A<sub>2</sub>), there exists  $\tau_2 \ge \tau_1$  such that

$$\frac{x(g(\tau,a))}{x''(g(\tau,a))} \ge \frac{\mu}{2}g^2(\tau,a)$$
(2.21)

for  $\tau \geq \tau_2$ . Next, we define

$$\varphi(\tau) = \frac{r(\tau)(x''(\tau))^{\alpha}}{(x''(\tau))^{\alpha}}.$$
(2.22)

We note that  $\varphi(\tau) < 0$  for  $\tau \ge \tau_1$ . By differentiating and using (1.1), (A<sub>3</sub>) and (S<sub>1</sub>), we obtain

$$\varphi'(\tau) \leq -k_1 Q(\tau) \frac{x^{\alpha}(g(\tau,a))}{(x''(g(\tau,a)))^{\alpha}} \frac{(x''(g(\tau,a)))^{\alpha}}{(x''(\tau))^{\alpha}} - \alpha \frac{1}{r^{1/\alpha}(\tau)} \varphi^{\frac{\alpha+1}{\alpha}}(\tau).$$

Hence, (2.21) yields

$$\varphi'(\tau) \le -\frac{k_1 \mu^{\alpha}}{2^{\alpha}} g^{2\alpha}(\tau, a) Q(\tau) - \alpha \frac{1}{r^{1/\alpha}(\tau)} \varphi^{\frac{\alpha+1}{\alpha}}(\tau).$$
(2.23)

From (2.16), we get

 $R_1^{\alpha}(\tau)\varphi(\tau) \geq -1.$ 

Multiplying (2.23) by  $R_1^{\alpha}(\tau)$  and integrating from  $\tau_2$  to  $\tau$ , we obtain

$$1 + R_1^{\alpha}(\tau_2)\varphi(\tau_2) \ge \frac{k_1\mu^{\alpha}}{2^{\alpha}} \int_{\tau_2}^{\tau} R_1^{\alpha}(s)g^{2\alpha}(s,a)Q(s)\,ds$$
$$+ \alpha \int_{\tau_2}^{\tau} \frac{R_1^{\alpha-1}(s)}{r^{1/\alpha}(s)} \big(\varphi(s) + R_1(s)\varphi^{\frac{1+\alpha}{\alpha}}(s)\big)\,ds.$$

By following the same steps as in Case  $(C_3)$ , we get that

$$\begin{split} &\int_{\tau_2}^{\tau} \left(\frac{k_1 \mu^{\alpha}}{2^{\alpha}} R_1^{\alpha}(s) g^{2\alpha}(s, \alpha) Q(s) - \left(\frac{\alpha}{\alpha+1}\right)^{\alpha+1} \frac{1}{r^{1/\alpha}(s) R_1(s)}\right) ds \\ &\leq 1 + R_1^{\alpha}(\tau_2) \varphi(\tau_2), \end{split}$$

which contradicts (2.4). This contradiction completes the proof of Theorem 2.1.  $\Box$ 

**Theorem 2.2** Assume that (1.2) and (S<sub>2</sub>) hold, and let  $g(\tau, \xi)$  have a positive partial derivative on  $I \times [a, b]$  with respect to  $\tau$ . If there exist continuously differentiable functions  $\rho, \vartheta \in C([\tau_0, \infty), (0, \infty))$  such that

$$\int_{\tau_0}^{\infty} \left( \rho(s)Q(s) - \frac{(2\alpha)^{\alpha}}{(\alpha+1)^{\alpha+1}} \frac{r(s)(\rho'_+(s))^{\alpha+1}}{(\mu k_2 s^2 \rho(s)g'(s,a))^{\alpha}} \right) ds = \infty,$$
(2.24)

$$\int_{\tau_0}^{\infty} \left(\vartheta\left(u\right) \int_u^{\infty} \left(\frac{1}{r(v)} \int_v^{\infty} Q(s) \, ds\right)^{1/\alpha} dv - \frac{\alpha}{4k_2} \frac{\vartheta_+^{\prime 2}(u)}{g'(u,a)\vartheta(u)}\right) du = \infty,\tag{2.25}$$

$$\int_{\tau_0}^{\infty} \left( f\left(R_3(s)\right) Q(s) - \frac{\alpha^{\alpha}}{(\alpha+1)^{\alpha+1}} R_2(s) \frac{(f'(R_3(s)))^{\alpha+1}}{k_2^{\alpha} f^{\alpha}(R_3(s))} \right) ds = \infty,$$
(2.26)

and

$$\int_{\tau_0}^{\infty} \left( R_1^{\alpha}(s) f\left(\frac{\mu}{2} g^2(\tau, a)\right) Q(s) - \left(\frac{\alpha}{\alpha+1}\right)^{\alpha+1} \frac{1}{r^{1/\alpha}(s) R_1(s)} \right) ds = \infty$$
(2.27)

for some  $\mu_1, \mu_2 \in (0, 1)$ , then every solution of (1.1) is oscillatory.

*Proof* Let *x* be a nonoscillatory solution of equation (1.1). Without loss of generality, we may assume that  $x(\tau) > 0$ . By Lemma 2.1, there exists  $\tau_1 \ge \tau_0$  such that  $x(\tau)$  has one of the four cases  $(C_1)$ - $(C_4)$  for  $\tau \ge \tau_1$ . *For Case*  $(C_1)$ , since  $g(\tau, \xi)$  is nondecreasing with respect to  $\xi$ ,  $x'(\tau) > 0$  and f'(x) > 0, we have that  $f(x(g(\tau, a))) \le f(x(g(\tau, \xi)))$ . Thus, from (1.1), we get

$$\frac{(r(\tau)(x'''(\tau))^{\alpha})'}{f(x(g(\tau,a)))} \leq -Q(\tau).$$

Now, we define

$$\omega(\tau) = \rho(\tau) \frac{r(\tau)(x'''(\tau))^{\alpha}}{f(x(g(\tau,a)))}.$$

By differentiating and using  $(S_2)$ , we get

$$\omega'(\tau) \le \frac{\rho'(\tau)}{\rho(\tau)} \omega(\tau) - \rho(\tau) Q(\tau) - k_2 \rho(\tau) \frac{r(\tau) (x'''(\tau))^{\alpha}}{f^{1+1/\alpha} (x(g(\tau, a)))} x'(g(\tau, a)) g'(\tau, a).$$
(2.28)

From (A<sub>2</sub>), there exists  $\tau_2 \ge \tau_1$  such that  $g(\tau, a) \ge \tau_1$  for  $\tau \ge \tau_2$ . Hence, from Lemma 1.2 and  $x^{(4)} < 0$ , we obtain

$$x'(g(\tau,a)) \geq \frac{\mu}{2}\tau^2 x'''(g(\tau,a)) \geq \frac{\mu}{2}\tau^2 x'''(\tau)$$

for all  $\mu \in (0, 1)$  and  $\tau \ge \tau_2$ . Therefore, (2.28) yields

$$\omega'(\tau) \le \frac{\rho'(\tau)}{\rho(\tau)}\omega(\tau) - \rho(\tau)Q(\tau) - \frac{\mu k_2}{2} \frac{\tau^2 g'(\tau, a)}{(\rho(\tau)r(\tau))^{1/\alpha}} \omega^{\frac{\alpha+1}{\alpha}}(\tau).$$

$$(2.29)$$

By following the same steps as in Case  $(C_1)$  of the proof of Theorem 2.1, we get a contradiction with (2.24).

For Case  $(C_2)$ . From (1.1),  $(S_2)$  and  $(A_3)$ , we obtain

$$(r(\tau)(x'''(\tau))^{\alpha})' \leq -f(x(g(\tau,a)))Q(\tau)$$

By integrating this inequality from  $\tau$  to  $\infty$ , we obtain

$$r(\tau)(x'''(\tau))^{\alpha} \geq \int_{\tau}^{\infty} f(x(g(s,a)))Q(s)\,ds.$$

Since f'(x) > 0, we get

$$x^{\prime\prime\prime}(\tau) \geq f^{1/\alpha} \left( x(g(\tau,a)) \right) \left( \frac{1}{r(\tau)} \int_{\tau}^{\infty} Q(s) \, ds \right)^{1/\alpha}.$$

Integrating again from  $\tau$  to  $\infty$ , we have

$$x''(\tau) \leq -f^{1/\alpha}\left(x(g(\tau,a))\right) \int_{\tau}^{\infty} \left(\frac{1}{r(\nu)} \int_{\nu}^{\infty} Q(s) \, ds\right)^{1/\alpha} d\nu.$$

Next, we define

$$w(\tau) = \vartheta(\tau) \frac{x'(\tau)}{f^{1/\alpha}(x(g(\tau, a)))}.$$

Then  $w(\tau) > 0$  for  $\tau \ge \tau_1$ . By differentiating and using (2.13), we find

$$\begin{split} w'(\tau) &\leq \frac{\vartheta'(\tau)}{\vartheta(\tau)} w(\tau) - \vartheta(\tau) \int_{\tau}^{\infty} \left( \frac{1}{r(\nu)} \int_{\nu}^{\infty} Q(s) \, ds \right)^{1/\alpha} d\nu \\ &- \frac{1}{\alpha} \vartheta(\tau) \left( \frac{x'(\tau)}{f^{1/\alpha}(x(g(\tau,a)))} \right)^2 \frac{f'(x(g(\tau,a)))}{f^{1-\frac{1}{\alpha}}(x(g(\tau,a)))} x'(g(\tau,a)) g'(\tau,a). \end{split}$$

Since  $x''(\tau) < 0$ , we see that  $x'(g(\tau, a)) > x'(\tau)$ 

$$w'(\tau) \leq \frac{\vartheta'(\tau)}{\vartheta(\tau)}w(\tau) - \vartheta(\tau)\int_{\tau}^{\infty} \left(\frac{1}{r(\nu)}\int_{\nu}^{\infty}Q(s)\,ds\right)^{1/\alpha}d\nu - \frac{k_2}{\alpha}\frac{g'(\tau,a)}{\vartheta(\tau)}w^2(\tau).$$

Then we get

$$w'(\tau) \leq -\vartheta(\tau) \int_{\tau}^{\infty} \left(\frac{1}{r(\nu)} \int_{\nu}^{\infty} Q(s) \, ds\right)^{1/\alpha} d\nu + \frac{\alpha}{4k_2} \frac{\vartheta'^2(\tau)}{g'(\tau, a)\vartheta(\tau)}.$$

Integrating again from  $\tau_2$  to  $\tau$ , we have

$$\int_{\tau_2}^{\tau} \left( \vartheta(u) \int_u^{\infty} \left( \frac{1}{r(v)} \int_v^{\infty} Q(s) \, ds \right)^{1/\alpha} dv - \frac{\alpha}{4k_2} \frac{\vartheta'^2(u)}{g'(u,a)\vartheta(u)} \right) du \leq w(\tau_2),$$

which contradicts (2.25).

*If Case*  $(C_3)$  *holds*. As in the proof of Case  $(C_3)$  of Theorem 2.1, we have that (2.16), (2.17) and (2.18) hold. Then we define

$$\psi(\tau) = \frac{r(\tau)(x'''(\tau))^{\alpha}}{f(x(\tau))}.$$

Thus, we see that  $\psi(\tau) < 0$  and satisfies

$$\psi'(\tau) = \frac{[r(\tau)(x'''(\tau))^{\alpha}]'}{f(x(\tau))} - \frac{r(\tau)(x'''(\tau))^{\alpha}}{f^2(x(\tau))} f'(x(\tau)) x'(\tau).$$

Hence, from (1.1), (2.17) and  $(S_2)$ , we have

$$\psi'(\tau) \leq -\int_a^b q(\tau,\xi) \frac{f(x(g(\tau,\xi)))}{f(x(\tau))} \, d\sigma(\xi) - k_2 R_2(\tau) \psi^{\frac{1+\alpha}{\alpha}}(\tau).$$

Since  $x'(\tau) < 0$ , we get  $f(x(g(\tau, \xi))) \ge f(x(\tau))$ . Therefore, we obtain

$$\psi'(\tau) \leq -Q(\tau) - k_2 R_2(\tau) \psi^{\frac{1+\alpha}{\alpha}}(\tau).$$

From (2.18) and  $(S_2)$ , we have

$$f(R_3(\tau))\psi(\tau) \ge -1. \tag{2.30}$$

Multiplying (2.19) by  $f(R_3(\tau))$  and integrating from  $\tau_1$  to  $\tau$ , we obtain

$$f(R_{3}(\tau))\psi(\tau) - f(R_{3}(\tau_{1}))\psi(\tau_{1}) + \int_{\tau_{1}}^{\tau} R_{2}(s)f'(R_{3}(s))\psi(s) ds + \int_{\tau_{1}}^{\tau} f(R_{3}(s))Q(s) ds + k_{2}\int_{\tau_{1}}^{\tau} R_{2}(s)f(R_{3}(s))\psi^{\frac{1+\alpha}{\alpha}}(s) ds \leq 0.$$

Using inequality (2.11) with  $A = k_2 f(R_3(s))$ ,  $B = f'(R_3(s))$  and  $z = -\psi$ , we get

$$\int_{\tau_1}^{\tau} \left( f(R_3(s)) Q(s) - \frac{\alpha^{\alpha}}{(\alpha+1)^{\alpha+1}} R_2(s) \frac{(f'(R_3(s)))^{\alpha+1}}{k_2^{\alpha} f^{\alpha}(R_3(s))} \right) ds$$
  

$$\leq 1 + f(R_3(\tau_1)) \psi(\tau_1),$$

but this contradicts (2.26).

In Case (C<sub>4</sub>). In view of the proof of Case (C<sub>4</sub>) of Theorem 2.1, we have (2.16) and (2.21) hold. By defining  $\varphi(\tau)$  as the form (2.22), we note that  $\varphi(\tau) < 0$  for  $\tau \ge \tau_1$ . Thus, from (1.1) and (A<sub>2</sub>), we get

$$\varphi'(\tau) \le -Q(\tau) \frac{f(x(g(\tau, a)))}{(x''(g(\tau, a)))^{\alpha}} \frac{(x''(g(\tau, a)))^{\alpha}}{(x''(\tau))^{\alpha}} - \alpha \frac{1}{r^{1/\alpha}(\tau)} \varphi^{\frac{\alpha+1}{\alpha}}(\tau).$$
(2.31)

From (2.21) and  $(S_2)$ , we see that

$$f(x(g(\tau,a))) \ge f\left(\frac{\mu}{2}g^2(\tau,a)\right)(x''(g(\tau,a)))^{\alpha}$$

Hence, (2.31) yields

$$\varphi'(\tau) \leq -f\left(\frac{\mu}{2}g^2(\tau,a)\right)Q(\tau) - \alpha \frac{1}{r^{1/\alpha}(\tau)}\varphi^{\frac{\alpha+1}{\alpha}}(\tau).$$

By following the same steps as in Case  $(C_3)$ , we get that

$$\int_{\tau_1}^{\tau} \left( R_1^{\alpha}(s) f\left(\frac{\mu}{2} g^2(\tau, a)\right) Q(s) - \left(\frac{\alpha}{\alpha+1}\right)^{\alpha+1} \frac{1}{r^{1/\alpha}(s) R_1(s)} \right) ds \leq 1 + R_1^{\alpha}(\tau_2) \varphi(\tau_2),$$

which contradicts (2.27). This contradiction completes the proof of Theorem 2.2.  $\hfill \Box$ 

**Theorem 2.3** Assume that (1.2) and  $(S_1)$  hold. If the differential equations

$$\left(\frac{r(\tau)}{\tau^{2\alpha}} \left(x'(\tau)\right)^{\alpha}\right)' + k_1 Q(\tau) \left(\frac{\mu_1 g^3(\tau, a)}{2\tau^3}\right)^{\alpha} x^{\alpha}(\tau) = 0,$$
(2.32)

$$x''(\tau) + x(\tau) \int_{\tau}^{\infty} \left[ \frac{k_1}{r(\nu)} \int_{\nu}^{\infty} \int_{a}^{b} q(s,\xi) \frac{g^{\alpha}(s,\xi)}{s^{\alpha}} d\sigma(\xi) ds \right]^{1/\alpha} d\nu = 0,$$
(2.33)

$$\left(\frac{1}{R_2^{\alpha}(\tau)} \left(x'(\tau)\right)^{\alpha}\right)' + k_1 Q(\tau) x^{\alpha}(\tau) = 0,$$
(2.34)

and

$$\left(r(\tau)(x'(\tau))^{\alpha}\right)' + \frac{k_1 \mu_2^{\alpha}}{2^{\alpha}} g^{2\alpha}(\tau, a) Q(\tau) x^{\alpha}(\tau) = 0$$
(2.35)

are oscillatory for some  $\mu_1, \mu_2 \in (0, 1)$ , then every solution of (1.1) is oscillatory.

*Proof* Proceeding as in the proof of Theorem 2.1, for Case (C<sub>1</sub>), we have that (2.10) holds. Then, if  $\rho(\tau) = 1$ , we get

$$\omega'(\tau) + k_1 Q(\tau) \left(\frac{g(\tau, a)}{\tau}\right)^{3\alpha} + \frac{\alpha \mu}{2} \frac{\tau^2}{r^{1/\alpha}(\tau)} \omega^{\frac{\alpha+1}{\alpha}}(\tau) \le 0$$
(2.36)

for all  $\mu \in (0, 1)$ . Hence, from [1], we see that (2.32) has a nonoscillatory solution for every  $\mu \in (0, 1)$ , which is a contradiction.

The rest of the proof is the same, and hence is omitted.

From Corollary 1 in Dzurina [3], if

$$\int_{\tau_0}^\infty \frac{1}{\widehat{r}^{1/\alpha}(s)}\,ds = \infty$$

and

$$\liminf_{\tau\to\infty}\widehat{r}^{1/\alpha}(\tau)\left[\int_{\tau_0}^{\tau}\widehat{r}^{-1/\alpha}(s)\,ds\right]^{\alpha+1}\widehat{q}(\tau)>\frac{\alpha}{4},$$

then equation

$$\left(\widehat{r}(\tau)\left(u'(\tau)\right)^{\alpha}\right)' + \widehat{q}(\tau)u^{\alpha}(\tau) = 0$$
(2.37)

is oscillatory. In the following theorem, by using the results of Dzurina [3], we will establish new oscillation criteria for solutions of equation (1.1) under the conditions

$$\int_{\tau_0}^{\infty} \frac{s^2}{r^{1/\alpha}(s)} \, ds = \infty \quad \text{and} \quad \int_{\tau_0}^{\infty} R_2(s) \, ds = \infty.$$
(2.38)

**Theorem 2.4** Assume that (1.2), (2.38) and  $(S_1)$  hold, and let (2.4) hold for some  $\mu_2 \in (0, 1)$ . If

$$\liminf_{\tau \to \infty} \frac{r^{1/\alpha}(\tau)}{\tau^2} \left(\frac{\mu_1 g^3(\tau, a)}{2\tau^3}\right)^{\alpha} \left(\int_{\tau_0}^{\tau} \frac{s^2}{r^{1/\alpha}(s)} \, ds\right)^{\alpha+1} Q(\tau) > \frac{\alpha}{4k_1},\tag{2.39}$$

$$\liminf_{\tau \to \infty} (\tau - \tau_0)^2 \int_{\tau}^{\infty} \left[ \frac{1}{r(\nu)} \int_{\nu}^{\infty} \int_{a}^{b} q(s,\xi) \frac{g^{\alpha}(s,\xi)}{s^{\alpha}} d\sigma(\xi) ds \right]^{1/\alpha} d\nu > \frac{1}{4k_1},$$
(2.40)

and

$$\liminf_{\tau \to \infty} \frac{Q(\tau)}{R_2(\tau)} \left( \int_{\tau_0}^{\tau} R_2(s) \, ds \right)^{\alpha+1} > \frac{\alpha}{4k_1} \tag{2.41}$$

for some  $\mu_1 \in (0,1)$ , then every solution of (1.1) is oscillatory.

Example 2.1 Consider the fourth-order differential equation

$$\left(e^{3\tau}\left(x'''(\tau)\right)^{3}\right)' + \int_{0}^{1} \frac{\delta\tau e^{\tau(\xi+3)}}{e^{\tau} - 1} x^{3}(\tau-\xi) \, d\xi, \qquad (2.42)$$

where  $\delta > 0$  is a constant. We note that

$$R_i(\tau) = e^{-\tau}, \quad i = 1, 2, 3 \text{ and } Q(\tau) = \delta e^{3\tau}.$$

If we choose  $\rho(\tau) = \vartheta(\tau) = 1$  and  $k_1 = 1$ , then it easy to see that conditions (2.1), (2.2), (2.3) and (2.4) hold for  $\delta > \frac{81}{256}$ . Thus, from Theorem 2.1, every solution of equation (2.42) is oscillatory for  $\delta > \frac{81}{256}$ .

Example 2.2 Consider the delay differential equation

$$\left(e^{\tau}x'''(\tau)\right)' + be^{\tau}x(\tau-1) = 0, \quad \tau \ge 1,$$
(2.43)

where b > 0. According to Corollary 4 in [2], equation (2.43) is oscillatory if  $b > \frac{2^5}{e}$ . If we choose  $\rho(\tau) = \vartheta(\tau) = 1$  and  $k_1 = 1$ , then we conclude that (2.1) and (2.2) are satisfied and (2.3) and (2.4) hold for  $b > \frac{1}{4}$ . Hence, by Theorem 2.1, every solution of equation (2.43) is oscillatory for  $b > \frac{1}{4}$ . Then our results supplement and improve some results obtained in [2]. In particular, we consider the equation

$$\left(e^{\tau} x'''(\tau)\right)' + 25\sqrt{2}e^{\tau+2\gamma} x(\tau-\gamma) = 0, \quad \tau \ge \gamma,$$
(2.44)

where  $\gamma = \sin^{-1} \frac{7}{5\sqrt{2}}$ . Since  $b = 25\sqrt{2}e^{2\gamma} > \frac{1}{4}$ , every solution of equation (2.44) is oscillatory. For example,  $x(\tau) = e^{2\tau} \sin(\tau)$  is a solution of equation (2.44). On the other hand, [24] showed that every nonoscillatory solution of

$$(e^{\tau}x'''(\tau))' + \frac{e^{\tau-1/2}}{16}x(\tau-1) = 0, \quad \tau \ge 1,$$

tends to zero as  $\tau \to \infty$ , and we note that  $b = \frac{e^{-1/2}}{16} < \frac{1}{4}$ .

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The authors declare that they have no competing interests.

#### Authors' contributions

All authors read and approved the final manuscript.

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