ON THE ASYMPTOTICS OF THE PRINCIPAL EIGENVALUE FOR A ROBIN PROBLEM WITH A LARGE PARAMETER IN PLANAR DOMAINS

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Let $\Omega \subset \mathbb{R}^2$ be a domain having a compact boundary Σ which is Lipschitz and piecewise C^4 smooth, and let ν denote the inward unit normal vector on Σ . We study the principal eigenvalue $E(\beta)$ of the Laplacian in Ω with the Robin boundary conditions $\partial f/\partial\nu + \beta f = 0$ on Σ , where β is a positive number. Assuming that Σ has no convex corners, we show the estimate $E(\beta) = -\beta^2 - \gamma_{\max}\beta + O\left(\beta^{\frac{2}{3}}\right)$ as $\beta \to +\infty$, where γ_{\max} is the maximal curvature of the boundary.

Keywords: eigenvalue, Laplacian, Robin boundary condition, curvature, asymptotics.

1. Introduction

Let $\Omega \subset \mathbb{R}^2$ be an open connected set having a compact Lipschitz piecewise smooth boundary Σ . For $\beta > 0$ consider the operator H_{β} which is the Laplacian $f \mapsto -\Delta f$ with the Robin boundary conditions,

$$\frac{\partial f}{\partial \nu} + \beta f = 0 \text{ on } \Sigma,$$

where ν is the inward unit normal vector. More precisely, H_{β} is the self-adjoint operator in $L^2(\Omega)$ associated with the sesquilinear form

$$h_{\beta}(f,g) = \iint_{\Omega} \overline{\nabla f} \nabla g \, dx - \beta \int_{\Sigma} \overline{f} \, g \, d\sigma, \quad \text{dom} \, h_{\beta} = H^{1}(\Omega); \tag{1}$$

here σ denotes the one-dimensional Hausdorff measure on Σ . The operator H_{β} is semibounded from below. If Ω is bounded, then H_{β} has a compact resolvent, and we denote by $E_{j}(\beta)$, $j \in \mathbb{N}$, its eigenvalues taken according to their multiplicities and enumerated in the non-decreasing order. If Ω is unbounded, then the essential spectrum of H_{β} coincides with $[0, +\infty)$, and the discrete spectrum consists of a finite number of eigenvalues, which we denote again by $E_{j}(\beta)$, $j \in \{1, \ldots, N_{\beta}\}$, and enumerate them in the non-decreasing order taking into account the multiplicities. In the both cases, the principal eigenvalue $E(\beta) := E_{1}(\beta)$ may be defined through the Rayleigh quotients

$$E(\beta) = \inf_{0 \neq f \in \text{dom } h_{\beta}} \frac{h_{\beta}(f, f)}{\|f\|_{L^{2}(\Omega)}^{2}}.$$

It is easy to check that $E(\beta) < 0$: for bounded Ω one can test on f = 1, and for unbounded Ω , one may use $f(x) = \exp\left(-|x|^{\alpha}/2\right)$ with small $\alpha > 0$.

The study of the principal eigenvalue arises in several applications: work [1] discusses the stochastic meaning of the Robin eigenvalues, paper [2] shows that the eigenvalue problem

appears in the study of long-time dynamics related to some reaction-diffusion processes, and a discussion of an interplay between the eigenvalues and the estimate of the critical temperature in a problem of superconductivity may be found in [3].

In the present note, we are interested in the asymptotic behavior of $E(\beta)$ for large values of β . For bounded Ω , this question was already addressed in numerous papers. It was conjectured and partially proven in [2] that one has the asymptotics

$$E(\beta) = -C_{\Omega}\beta^2 + o(\beta^2) \text{ as } \beta \to +\infty,$$
 (2)

for some constant $C_{\Omega}>0$. It seems that the paper [4] contains the first rigorous proof of the above equality for the case of a C^1 smooth Σ , and in that case one has $C_{\Omega} = 1$, as predicted in [2]. Under the same assumption, it was shown in [5] that the asymptotics $E_i(\beta) = -\beta^2 + o(\beta^2), \ \beta \to +\infty, \ \text{hold for any fixed } j \in \mathbb{N}.$ The paper [6] proved the asymptotics (2) for domains whose boundary is C^{∞} smooth with the possible exception of a finite number of corners. If the corner opening angles are $\alpha_j \in (0,\pi) \cup (\pi,2\pi)$, $j=1,\ldots,m$, and $\theta := \min \alpha_i/2$, then $C_{\Omega} = (\sin \theta)^{-2}$ if $\theta < \pi/2$, otherwise $C_{\Omega} = 1$. We remark that the paper [6] formally deals with bounded domains, but the proofs can be easily adapted to unbounded domains with compact boundaries. It should pointed out that domains with cusps need a specific consideration, and the results are different [6,7]. Various generalizations of the above results and some related questions concerning the spectral theory of the Robin Laplacians were discussed in [7–12]. The aim of the present note is to refine the asymptotics (2) for a class of two-dimensional domains. More precisely, we calculate the next term in the asymptotic expansion for piecewise C^4 smooth domains whose boundary has no convex corners, i.e. we assume that either the boundary is smooth or that all corner opening angles are larger than π ; due to the above cited result of [6] we have $C_{\Omega} = 1$ in the both cases.

Let us formulate the assumptions and the result more carefully. Let Σ_k , $k=1,\ldots,n$, be non-intersecting C^4 smooth connected components of the boundary Σ such that $\Sigma = \bigcup_{k=1}^n \overline{\Sigma}_k$. Denote by ℓ_k the length of Σ_k and consider a parametrization of the closure $\overline{\Sigma}_k$ by the arc length, i.e. let $[0,\ell_k]\ni s\mapsto \Gamma_k(s)\equiv \left(\Gamma_{k,1}(s),\Gamma_{k,2}(s)\right)\in \overline{\Sigma}_k$ be a bijection with $|\Gamma'_k|=1$, such that $\Gamma_k\in C^4\bigl([0,\ell_k],\mathbb{R}^2\bigr)$, and we assume that the orientation of each Γ_k is chosen in such a way that $\nu_k(s):=\left(-\Gamma'_{k,2}(s),\Gamma'_{k,1}(s)\right)$ is the inward unit normal vector at the point $\Gamma_k(s)$ of the boundary. If two components Σ_j , Σ_k meet at some point $P:=\Gamma_j(\ell_j)=\Gamma_k(0)$, then two options are allowed: either $\overline{\Sigma_j}\cup\overline{\Sigma_k}$ is C^4 smooth near P or the corner opening angle at P measured inside Ω belongs to $(\pi,2\pi)$.

We denote by $\gamma_k(s)$ the signed curvature of the boundary at the point $\Gamma_k(s)$ and let γ_{\max} denote its global maximum:

$$\gamma_k(s) := \Gamma'_{k,1}(s)\Gamma''_{k,2}(s) - \Gamma''_{k,1}(s)\Gamma'_{k,2}(s), \quad \gamma_{\max} := \max_{k \in \{1,\dots,n\}} \max_{s \in [0,\ell_k]} \gamma_k(s);$$

note that the decomposition of the boundary Σ into the pieces Σ_k is non-unique, but the value γ_{\max} is uniquely determined. Our result is as follows:

Theorem 1. Under the preceding assumptions there holds

$$E(\beta) = -\beta^2 - \gamma_{\text{max}}\beta + O(\beta^{\frac{2}{3}}) \text{ as } \beta \to +\infty.$$

We believe that it is hard to improve the asymptotics without any additional information on the set at which the curvature attains its maximal value. For example, one may expect that the case of a curvature having isolated maxima and the case of a piecewise constant curvature should give different resolutions of the remainder, and we hope to progress in this direction in subsequent works.

At first sight, the Robin eigenvalue problem may look rather similar to the eigenvalue problem for δ -potentials supported by curves, see e.g. [13–15]. This first impression is wrong, and the result of Theorem 1 concerning the secondary asymptotic term is very different from the one obtained in the papers [13,14] for strong δ -potentials; nevertheless, a part of the machinery of [13] plays an important role in our considerations. On the other hand, the asymptotic behavior of the principal Robin eigenvalue shows some analogy with the lowest eigenvalue of the Neumann magnetic Laplacian studied in the theory of superconductivity [16–18].

2. Dirichlet-Neumann bracketing on thin strips

In this section we introduce and study an auxiliary eigenvalue problem, and the results obtained will be used in the next section to prove theorem 1.

Let $\ell > 0$ and let $\Gamma : [0, \ell] \to \mathbb{R}^2$, $s \mapsto \Gamma(s) = (\Gamma_1(s), \Gamma_2(s)) \in \mathbb{R}^2$, be an injective C^4 map such that $|\Gamma'(s)| = 1$ for all $s \in (0, \ell)$. We denote

$$S := \Gamma((0, \ell)), \quad \kappa(s) := \Gamma'_1(s)\Gamma''_2(s) - \Gamma''_1(s)\Gamma'_2(s), \quad \kappa_{\max} := \max_{s \in [0, \ell]} \kappa(s),$$

$$K := \max_{s \in [0, \ell]} |\kappa(s)| + \max_{s \in [0, \ell]} |\kappa'(s)| + \max_{s \in [0, \ell]} |\kappa''(s)|.$$

Due to $\kappa \in C^2([0,\ell])$, the above quantity K is finite.

For a > 0, consider the map

$$\Phi_a: (0,\ell) \times (0,a) \to \mathbb{R}^2, \quad \Phi_a(s,u) = \begin{pmatrix} \Gamma_1(s) - u\Gamma_2'(s) \\ \Gamma_2(s) + u\Gamma_1'(s) \end{pmatrix}.$$

As shown in [13, Lemma 2.1], for any $a \in (0, a_0)$, $a_0 := (2K)^{-1}$, the map Φ_a defines a diffeomorphism between the domains $\Box_a := (0, \ell) \times (0, a)$ and $\Omega_a := \Phi_a(\Box_a)$. In what follows, we always assume that $a \in (0, a_0)$ and we will work with the usual Sobolev space $H^1(\Omega_a)$ and its part $\widetilde{H}^1_0(\Omega_a) := \{ f \in H^1(\Omega_a) : f \lceil_{\partial \Omega_a \setminus \overline{S}} = 0 \}$; here the symbol \lceil means the trace of the function on the indicated part of the boundary.

Here, we introduce two sesquilinear forms in $L^2(\Omega_a)$. The first one, $h_{\beta}^{N,a}$, is defined on $\mathrm{dom}\,h_{\beta}^{N,a}:=H^1(\Omega_a)$ by the expression

$$h_{\beta}^{N,a}(f,g) = \iint_{\Omega_a} \overline{\nabla f} \nabla g \, \mathrm{d}x - \beta \int_S \overline{f} g \, \mathrm{d}\sigma,$$

and the second one, $h_{\beta}^{D,a}$, is its restriction to $\mathrm{dom}\,h_{\beta}^{D,a}:=\widetilde{H}_{0}^{1}(\Omega_{a})$. Both forms are densely defined, symmetric, closed and semibounded from below, and we denote

$$E_{N/D}(\beta, a) = \inf_{0 \neq f \in \text{dom } h_{\alpha}^{N/D, a}} \frac{h_{\beta}^{N/D, a}(f, f)}{\|f\|_{L^{2}(\Omega_{\alpha})}^{2}}.$$
 (3)

We show the following results:

Lemma 2. There exists $a_1 > 0$ such that for any $a \in (0, a_1)$ one has the estimate $E_{N/D}(\beta, a) = -\beta^2 - \kappa_{\max}\beta + O(\beta^{\frac{2}{3}})$ as $\beta \to +\infty$.

The rest of this section is devoted to the proof of lemma 2. We first introduce a suitable decomposition of Ω_a and then provide two-side eigenvalue estimates using operators with separated variables.

On the asymptotics of the principal eigenvalue for a Robin problem

Define $U_a:L^2(\Omega_a)\to L^2(\square_a)$ by $\big(U_af\big)(s,u)=\sqrt{1-u\kappa(s)}f\big(\Phi_a(s,u)\big)$. Clearly, U_a is a unitary operator, and one has $U_a\big(H^1(\Omega_a)\big)=H^1(\square_a)$ and

$$U_a(\widetilde{H}_0^1(\Omega_a)) = \widetilde{H}_0^1(\square_a) := \{ f \in H^1(\square_a) : f(0,\cdot) = f(\ell,\cdot) = 0 \text{ and } f(\cdot,a) = 0 \},$$

where the restrictions should be again understood as the traces. Using integration by parts, one may easily check that for any $f,g\in H^1(\Omega_a)$, one has $h_{\beta}^{N,a}(f,g)=q_{\beta}^{N,a}(U_af,U_ag)$, where the form $q_{\beta}^{N,a}$ is defined on the domain $\mathrm{dom}\,q_{\beta}^{N,a}:=H^1(\square_a)$ by the expression

$$q_{\beta}^{N,a}(f,g) = \iint_{\square_a} \frac{1}{\left(1 - u\kappa(s)\right)^2} \frac{\overline{\partial f}}{\partial s} \frac{\partial g}{\partial s} \, \mathrm{d}s \, \mathrm{d}u + \iint_{\square_a} \frac{\overline{\partial f}}{\partial u} \frac{\partial g}{\partial u} \, \mathrm{d}s \, \mathrm{d}u$$

$$- \iint_{\square_a} V(s,u) \overline{f} \, g \, \mathrm{d}s \, \mathrm{d}u - \beta \int_0^\ell \overline{f(s,0)} g(s,0) \, \mathrm{d}s$$

$$- \frac{1}{2} \int_0^\ell \kappa(s) \, \overline{f(s,0)} g(s,0) \, \mathrm{d}s + \frac{1}{2} \int_0^\ell \frac{\kappa(s)}{1 - a\kappa(s)} \, \overline{f(s,a)} g(s,a) \, \mathrm{d}s \qquad (4)$$

$$+ \frac{1}{2} \kappa'(\ell) \int_0^a \frac{u}{\left(1 - u\kappa(\ell)\right)^3} \, \overline{f(\ell,u)} g(\ell,u) \, \mathrm{d}u$$

$$- \frac{1}{2} \kappa'(0) \int_0^a \frac{u}{\left(1 - u\kappa(0)\right)^3} \, \overline{f(0,u)} g(0,u) \, \mathrm{d}u$$

with

$$V(s,u) := \frac{u\kappa''(s)}{2(1 - u\kappa(s))^3} + \frac{5u^2\kappa'(s)^2}{4(1 - u\kappa(s))^4} + \frac{\kappa(s)^2}{4(1 - u\kappa(s))^2}.$$

Similarly, for any $f,g\in \widetilde{H}^1_0(\Omega_a)$, one has $h^{D,a}_{\beta}(f,g)=q^{D,a}_{\beta}(U_af,U_ag)$, where $q^{D,a}_{\beta}$ is the restriction of $q^{N,a}_{\beta}$ to the domain $\mathrm{dom}\,q^{D,a}_{\beta}:=\widetilde{H}^1_0(\square_a)$; note that for $f,g\in\mathrm{dom}\,q^{D,a}_{\beta}$ the three last terms on the right-hand side of (4) vanish. Using the unitarity of U_a we may rewrite the equalities (3) in the form:

$$E_{N/D}(\beta, a) = \inf_{0 \neq f \in \text{dom } q_{\beta}^{N/D, a}} \frac{q_{\beta}^{N/D, a}(f, f)}{\|f\|_{L^{2}(\square_{a})}^{2}}.$$
 (5)

We would like to reduce the estimation of these quantities to the study of the eigenvalues of certain one-dimensional operators.

Using the one-dimensional Sobolev inequality on $(0, \ell)$ we see that one can find a constant C > 0 independent of a, such that for all $f \in H^1(\square_a)$, one has

$$\int_0^a |f(0,u)|^2 du + \int_0^a |f(\ell,u)|^2 du \le C \Big(\iint_{\square_a} \left| \frac{\partial f}{\partial s} \right|^2 ds du + \iint_{\square_a} |f|^2 ds du \Big).$$

One can also find a constant v > 0, such that $|V(s,u)| \le v$ for all $(s,u) \in \square_a$ and all $a \in (0,a_0)$. Furthermore, again for $(s,u) \in \square_a$ and any $a \in (0,a_0)$, we have

$$\left| \frac{\kappa(s)}{1 - a\kappa(s)} \right| \le 2K, \quad \frac{2}{3} \le \frac{1}{1 - u\kappa(s)} \le 2.$$

For any $M \in \mathbb{N}$, we denote

$$\delta := \frac{\ell}{M}, \quad I_M^j := (j\delta - \delta, j\delta), \quad \Box_{a,M}^j := I_M^j \times (0, a),$$

$$\kappa_{M,j}^- := \inf_{s \in I_M^j} \kappa(s), \quad \kappa_{M,j}^+ := \sup_{s \in I_M^j} \kappa(s), \quad j = 1, \dots, M,$$

and introduce functions $\kappa_M^\pm:(0,\ell)\to\mathbb{R}$ as follows: $\kappa_M^\pm(s):=\kappa_{M,j}^\pm$ if $s\in I_M^j$, and $\kappa_M^\pm(j\delta):=0$ for $j=1,\ldots,M-1$. In addition, we assume that $0< a<(10KC)^{-1}$. Now, we introduce two new sesquilinear forms which will be used to obtain a two-side estimate for $E_{N/D}(\beta,a)$. The first one, $t_\beta^{-,M,a}$, is defined by

$$\operatorname{dom} t_{\beta}^{-,M,a} = H^{1}\left(\bigcup_{j=1}^{M} \Box_{a,M}^{j}\right) \simeq \bigoplus_{j=1}^{M} H^{1}\left(\Box_{a,M}^{j}\right),$$

$$t_{\beta}^{-,M,a}(f,g) = \left(\frac{4}{9} - 4aKC\right) \iint_{\Box_{a}} \frac{\overline{\partial f}}{\partial s} \frac{\partial g}{\partial s} \, \mathrm{d}s \, \mathrm{d}u + \iint_{\Box_{a}} \frac{\overline{\partial f}}{\partial u} \frac{\partial g}{\partial u} \, \mathrm{d}s \, \mathrm{d}u$$

$$-\left(v + 4aKC\right) \iint_{\Box_{a}} \overline{f}g \, \mathrm{d}s \, \mathrm{d}u - \int_{0}^{\ell} \left(\beta + \frac{\kappa_{M}^{+}(s)}{2}\right) \overline{f(s,0)} g(s,0) \, \mathrm{d}s$$

$$-K \int_{0}^{\ell} \overline{f(s,a)} g(s,a) \, \mathrm{d}s.$$

The second one, $t_{\beta}^{+,M,a}$, is defined on the domain $\dim t_{\beta}^{+,M,a}=\bigoplus_{j=1}^M\widetilde{H}_0^1(\square_{a,M}^j)$,

$$\widetilde{H}^1_0(\square_{a,M}^j):=\big\{f\in H^1(\square_{a,M}^j):\, f(j\delta-\delta,\cdot)=f(j\delta,\cdot)=0 \text{ and } f(\cdot,a)=0\big\},$$

through

$$t_{\beta}^{+,M,a}(f,g) = 4 \iint_{\square_a} \frac{\overline{\partial f}}{\partial s} \frac{\partial g}{\partial s} \, \mathrm{d}s \, \mathrm{d}u + \iint_{\square_a} \frac{\overline{\partial f}}{\partial u} \frac{\partial g}{\partial u} \, \mathrm{d}s \, \mathrm{d}u + \iint_{\square_a} \frac{\overline{\partial f}}{\partial u} \frac{\partial g}{\partial u} \, \mathrm{d}s \, \mathrm{d}u + \iint_{\square_a} \frac{\overline{\partial f}}{\partial u} \frac{\partial g}{\partial u} \, \mathrm{d}s \, \mathrm{d}u - \int_0^\ell \left(\beta + \frac{\kappa_M^-(s)}{2}\right) \overline{f(s,0)} g(s,0) \, \mathrm{d}s.$$

One clearly has the inclusions $\operatorname{dom} t_{\beta}^{+,M,a} \subset \operatorname{dom} q_{\beta}^{D,a} \subset \operatorname{dom} q_{\beta}^{N,a} \subset \operatorname{dom} t_{\beta}^{-,M,a}$ and the inequalities

$$t_{\beta}^{-,M,a}(f,f) \leq q_{\beta}^{N,a}(f,f), \qquad f \in \text{dom } q_{\beta}^{N,a},$$

$$q_{\beta}^{N,a}(f,f) = q_{\beta}^{D,a}(f,f), \qquad f \in \text{dom } q_{\beta}^{D,a},$$

$$q_{\beta}^{D,a}(f,f) \leq t_{\beta}^{+,M,a}(f,f), \qquad f \in \text{dom } t_{\beta}^{+,M,a},$$

which justify the estimates

$$E_M^-(\beta, a) \le E_N(\beta, a) \le E_D(\beta, a) \le E_M^+(\beta, a), \tag{6}$$

where we denote

$$E_M^{\pm}(\beta, a) := \inf_{0 \neq f \in \text{dom } t_{\beta}^{\pm, M, a}} \frac{t_{\beta}^{\pm, M, a}(f, f)}{\|f\|_{L^2(\square_a)}^2}.$$

Now, we are going to estimate $E_M^\pm(\beta,a)$ using separation of variables. Note that the forms $t_\beta^{\pm,M,a}$ are densely defined, semibounded from below and closed in $L^2(\square_a)$, therefore, they define some self-adjoint operators $T_\beta^{\pm,M,a}$ in $L^2(\square_a)$, and $E_M^\pm(\beta,a)=\inf \operatorname{spec} T_\beta^{\pm,M,a}$. On the

other hand, due to the fact that the domains $\Box_{a,M}^j$ are disjoint and isometric to one another , we can identify $T_{\beta}^{\pm,M,a}\simeq\bigoplus_{j=1}^MT_{\beta,j}^{\pm,M,a}$, where $T_{\beta,j}^{\pm,M,a}$ are self-adjoint operators acting in $L^2(\Box_{\delta,a})$, $\Box_{\delta,a}:=(0,\delta)\times(0,a)$, and associated respectively with the sesqulinear forms $t_{\beta,j}^{\pm,M,a}$,

$$\begin{split} t_{\beta,j}^{-,M,a}(f,g) &= \Big(\frac{4}{9} - 4aKC\Big) \int_0^\delta \int_0^a \frac{\overline{\partial f}}{\partial s} \frac{\partial g}{\partial s} \mathrm{d}u \, \mathrm{d}s + \int_0^\delta \int_0^a \frac{\overline{\partial f}}{\partial u} \frac{\partial g}{\partial u} \mathrm{d}u \, \mathrm{d}s \\ &- (v + 4aKC) \int_0^\delta \int_0^a \overline{f}g \, \mathrm{d}u \, \mathrm{d}s - \Big(\beta + \frac{\kappa_{M,j}^+}{2}\Big) \int_0^\delta \overline{f(s,0)}g(s,0) \, \mathrm{d}s \\ &- K \int_0^\delta \overline{f(s,a)}g(s,a) \, \mathrm{d}s, \quad \mathrm{dom} \, t_{\beta,j}^{-,M,a} &= H^1(\square_{\delta,a}), \\ t_{\beta,j}^{+,M,a}(f,g) &= 4 \int_0^\delta \int_0^a \overline{\frac{\partial f}{\partial s}} \frac{\partial g}{\partial s} \mathrm{d}u \, \mathrm{d}s + \int_0^\delta \int_0^a \overline{\frac{\partial f}{\partial u}} \frac{\partial g}{\partial u} \mathrm{d}u \, \mathrm{d}s \\ &+ v \int_0^\delta \int_0^a \overline{f}g \, \mathrm{d}u \, \mathrm{d}s - \Big(\beta + \frac{\kappa_{M,j}^-}{2}\Big) \int_0^\delta \overline{f(s,0)}g(s,0) \, \mathrm{d}s, \\ \mathrm{dom} \, t_{\beta,j}^{+,M,a} &= \big\{ f \in H^1(\square_{\delta,a}) : \, f(0,\cdot) = f(\delta,\cdot) = 0 \, \, \text{and} \, \, f(\cdot,a) = 0 \big\}. \end{split}$$

It is routine to check that $T_{\beta,j}^{\pm,M,a}=Q_M^\pm\otimes 1+1\otimes L_{\beta,a}^{\pm,j}$, where Q_M^\pm are the operators acting in $L^2(0,\delta)$ as follows:

$$\begin{split} Q_M^- f &= -\left(\frac{4}{9} - 4aKC\right) f'' - (v + 4aKC) f, \\ \mathrm{dom}\, Q_M^- &= \Big\{ f \in H^2(0,\delta) : f'(0) = f'(\delta) = 0 \Big\}, \\ Q_M^+ f &= -4f'' + vf, \\ \mathrm{dom}\, Q_M^- &= \Big\{ f \in H^2(0,\delta) : f(0) = f(\delta) = 0 \Big\}, \end{split}$$

and $L_{\beta,a}^{\pm,j}$ are the self-adjoint operators in $L^2(0,a)$ both acting as $L_{\beta,a}^{\pm,j}f=-f''$ on the domains

$$\operatorname{dom} L_{\beta,a}^{-,j} = \left\{ f \in H^2(0,a) : f'(0) + \left(\beta + \frac{\kappa_{M,j}^+}{2}\right) f(0) = 0, \ f'(a) - K f(a) = 0 \right\},
\operatorname{dom} L_{\beta,a}^{+,j} = \left\{ f \in H^2(0,a) : f'(0) + \left(\beta + \frac{\kappa_{M,j}^-}{2}\right) f(0) = 0, \ f(a) = 0 \right\}.$$

The spectra of Q_M^{\pm} can be calculated explicitly; in particular, one has

inf spec
$$Q_M^- = -v - 4aKC$$
, inf spec $Q_M^+ = \frac{4\pi^2}{\delta^2} + v \equiv \frac{4\pi^2M^2}{\ell^2} + v$.

Therefore, denoting $E^{\pm,j}(\beta,a):=\inf \operatorname{spec} L^{\pm,j}_{\beta,a}$, we arrive at

$$E_{M}^{-}(\beta, a) = \min_{j} \left(\inf \operatorname{spec} T_{\beta, j}^{-, M, a} \right) = -v - 4aKC + \min_{j} E^{-, j}(\beta, a),$$

$$E_{M}^{+}(\beta, a) = \min_{j} \left(\inf \operatorname{spec} T_{\beta, j}^{+, M, a} \right) = \frac{4\pi^{2}M^{2}}{\ell^{2}} + v + \min_{j} E^{+, j}(\beta, a).$$
(7)

To study the lowest eigenvalues of $L_{\beta,a}^{\pm,j}$, we prove two auxiliary estimates.

Lemma 3. For $a, \beta, \gamma > 0$, let $\Lambda_{a,\beta,\gamma}$ denote the self-adjoint operator in $L^2(0,a)$ acting as $f \mapsto -f''$ on the functions $f \in H^2(0,a)$ satisfying the boundary conditions

 $f'(0) + \beta f(0) = f'(a) - \gamma f(a) = 0$, and let $E(a, \beta, \gamma)$ be its lowest eigenvalue. Let $\beta > 2\gamma$ and $\beta a > 1$, then $\beta^2 < -E(a, \beta, \gamma) < \beta^2 + 123\beta^2 e^{-2\beta a}$.

Proof. Let k>0. Clearly, $E=-k^2$ is an eigenvalue of $\Lambda_{a,\beta,\gamma}$ if one can find $(C_1,C_2)\in\mathbb{C}^2\setminus\left\{(0,0)\right\}$ such that the function $f:x\mapsto C_1e^{kx}+C_2e^{-kx}$ belongs to the domain of $\Lambda_{a,\beta,\gamma}$. The boundary conditions give

$$0 = f'(0) + \beta f(0) = (\beta + k)C_1 + (\beta - k)C_2,$$

$$0 = f'(a) - \gamma f(a) = (k - \gamma)e^{ka}C_1 - (k + \gamma)e^{-ka}C_2,$$

and one has a non-zero solution if the determinant of this system vanishes, i.e. if k satisfies the equation $(k+\beta)(k+\gamma)e^{-ka}=(k-\beta)(k-\gamma)e^{ka}$. Let us look for solutions $k\in(\beta,+\infty)$. One may rewrite the preceding equation as

$$g(k) = h(k), \quad g(k) = \frac{k+\beta}{k-\beta}, \quad h(k) = \frac{k-\gamma}{k+\gamma}e^{2ka}.$$
 (8)

Both functions g and h are continuous. It is readily seen that the function g is strictly decreasing on $(\beta, +\infty)$ with $g(\beta+) = +\infty$ and $g(+\infty)=1$. Conversely, for $\beta > 2\gamma$, the function h is strictly increasing in $(\beta, +\infty)$, being the product of two strictly increasing positive functions, and we have $h(\beta+) = e^{2\beta a}(\beta-\gamma)/(\beta+\gamma) < +\infty$ and $h(+\infty) = +\infty$. These properties of g and h show that there exists a unique solution $k = k(a, \beta, \gamma) \in (\beta, +\infty)$ of (8) and that $E(a, \beta, \gamma) = -k(a, \beta, \gamma)^2$.

To obtain the required estimate we use again the monotonicity of h on $(\beta, +\infty)$ and the inequality $\beta > 2\gamma$. We have

$$\frac{k+\beta}{k-\beta} = g(k) = h(k) > h(\beta+) = \frac{\beta-\gamma}{\beta+\gamma}e^{2\beta a} \ge \frac{e^{2\beta a}}{3},$$

which gives $(1 - 3e^{-2\beta a})k < (1 + 3e^{-2\beta a})\beta$. The assumption $\beta a > 1$ gives the inequality $3e^{-2\beta a} < 1/2$, and we arrive at

$$k < \frac{1 + 3e^{-2\beta a}}{1 - 3e^{-2\beta a}}\beta < (1 + 3e^{-2\beta a})(1 + 15e^{-2\beta a})\beta < (1 + 41e^{-2\beta a})\beta$$

and $k^2 < (1+41e^{-2\beta a})^2\beta^2 < (1+123e^{-2\beta a})\beta^2$. Together with the inclusion $k \in (\beta, +\infty)$ this gives the result.

Lemma 4. For $a, \beta > 0$, let $\Pi_{a,\beta}$ denote the self-adjoint operator in $L^2(0,a)$ acting as $f \mapsto -f''$ on the functions $f \in H^2(0,a)$ satisfying the boundary conditions $f'(0) + \beta f(0) = f(a) = 0$, and let $E(a,\beta)$ be its lowest eigenvalue. Assume that $\beta a > 4/3$, then $\beta^2 - 4\beta^2 e^{-\beta a} < -E(a,\beta) < \beta^2$.

Proof. Let k>0. Proceeding as in the proof of lemma 3, we see that $E=-k^2$ is an eigenvalue of $\Pi_{a,\beta}$ if k satisfies the equation $(\beta+k)e^{-ka}=(\beta-k)e^{ka}$. As the left side of the equation is strictly positive, the right side must be positive as well, which means that all solutions k belong to $(0,\beta)$. Let us rewrite the equation in the form g(k)=0 with $g(k):=\log(\beta+k)-\log(\beta-k)-2ka$. One has g(0)=0, the function g is strictly decreasing in $(0,k_0)$ and strictly increasing in (k_0,β) , with $k_0:=\sqrt{\beta^2-\beta/a}$. Moreover, $g(\beta-)=+\infty$. Therefore, the equation g(k)=0 has a unique solution in (k_0,β) . It follows from the assumption $\beta a>4/3$ that $k_0>\beta/2$, and we can represent $k=\beta-s$ with some $s\in(0,\beta/2)$. Using again the condition g(k)=0, we arrive at the inequality $\log s=\log(2\beta-s)-2\beta a+2sa<\log(2\beta)-\beta a$, which gives $s<2\beta e^{-\beta a}$ and $k=\beta-s>\beta(1-2e^{-\beta a})$. Finally, $-E(a,\beta)=k^2>\beta^2(1-2e^{-\beta a})^2>\beta^2(1-4e^{-\beta a})$. Together with the first inequality $k<\beta$ this gives the desired estimate .

On the asymptotics of the principal eigenvalue for a Robin problem

Let us complete the proof of lemma 2. Denote $a_1 := \min \left\{ a_0, (10KC)^{-1} \right\}$ and pick any $a \in (0, a_1)$, and let $\beta > 3K + 1 + 4/(3a)$. Applying lemma 3 to each of the operators $L_{\beta,a}^{-,j}$ and lemma 4 to each of the operators $L_{\beta,a}^{+,j}$, we arrive at the estimates

$$E^{-,j}(\beta, a) > -\left(\beta + \frac{\kappa_{M,j}^{+}}{2}\right)^{2} - 123\left(\beta + \frac{\kappa_{M,j}^{+}}{2}\right)^{2} \exp\left[-2a\left(\beta + \frac{\kappa_{M,j}^{+}}{2}\right)\right],$$

$$E^{+,j}(\beta, a) < -\left(\beta + \frac{\kappa_{M,j}^{-}}{2}\right)^{2} + 4\left(\beta + \frac{\kappa_{M,j}^{-}}{2}\right)^{2} \exp\left[-a\left(\beta + \frac{\kappa_{M,j}^{-}}{2}\right)\right].$$

To simplify the form of the remainders, we choose $\beta_a > 0$ sufficiently large such that for $\beta > \beta_a$ we have

$$\left(\beta + \frac{K}{2}\right)^2 \exp\left[-2a\left(\beta - \frac{K}{2}\right)\right] + 4\left(\beta + \frac{K}{2}\right)^2 \exp\left[-a\left(\beta - \frac{K}{2}\right)\right] \le \frac{1}{\beta},$$

then for $\beta > \beta_a + 3K + 1 + 4/(3a)$ and all $j = 1, \dots, M$, we have

$$E^{-,j}(\beta,a) > -\beta^2 - \kappa_{M,j}^+ \beta - \frac{K^2}{4} - \frac{1}{\beta}, \quad E^{+,j}(\beta,a) < -\beta^2 - \kappa_{M,j}^- \beta + \frac{1}{\beta}.$$

Using the inequality $\kappa_{M,j}^+ \leq \kappa_{\max}$, we obtain

$$\min_{j} E^{-,j}(\beta, a) > -\beta^{2} - \kappa_{\max}\beta - \frac{K^{2}}{4} - \frac{1}{\beta}.$$
 (9)

Conversely, let $l \in \{1, \dots, M\}$ be such that $\kappa_{M,l}^+ = \kappa_{\max}$. This means that there exists $s \in \overline{I_M^l}$ such that $\kappa(s) = \kappa_{\max}$. Using the Taylor expansion near s, we obtain

$$\kappa_{M,l}^{-} \ge \kappa_{M,l}^{+} - K\delta = \kappa_{\max} - K\delta \equiv \kappa_{\max} - \frac{K\ell}{M}.$$
 (10)

In the previous considerations the number M was arbitrary, and now we pick $M \in \left[\beta^{\frac{1}{3}}, 2\beta^{\frac{1}{3}}\right] \cap \mathbb{N}$, then

$$\min_{j} E^{+,j}(\beta, a) \leq E^{+,l}(\beta, a) < -\beta^{2} - \kappa_{M,l}^{-}\beta + \frac{1}{\beta}$$

$$= -\beta^{2} - \kappa_{\max}\beta + \frac{K\ell}{M}\beta + \frac{1}{\beta} \leq -\beta^{2} - \kappa_{\max}\beta + K\ell\beta^{\frac{2}{3}} + \frac{1}{\beta}. \quad (11)$$

Substituting the estimates (9) and (11) into (7) we arrive at

$$E_{M}^{+}(\beta, a) \leq -\beta^{2} - \kappa_{\max}\beta + K\ell\beta^{\frac{2}{3}} + \frac{1}{\beta} + \frac{4\pi^{2}M^{2}}{\ell^{2}} + v$$

$$= -\beta^{2} - \kappa_{\max}\beta + \left(K\ell + \frac{16\pi^{2}}{\ell^{2}}\right)\beta^{\frac{2}{3}} + v + \frac{1}{\beta},$$

$$E_{M}^{-}(\beta, a) \geq -\beta^{2} - \kappa_{\max}\beta - \frac{K^{2}}{4} - v - 4aKC - \frac{1}{\beta},$$

and the assertion of lemma 2 follows from the two-side estimates (6).

3. Proof of Theorem 1

We continue using the notation introduced just before theorem 1. For a>0, consider the maps

$$\Phi_{k,a}: (0,\ell_k) \times (0,a) \to \mathbb{R}^2, \quad \Phi_{k,a}(s,u) = \begin{pmatrix} \Gamma_{k,1}(s) - u\Gamma'_{k,2}(s) \\ \Gamma_{k,2}(s) + u\Gamma'_{k,1}(s) \end{pmatrix}, \quad k = 1,\ldots, n.$$

As in section 2, we can find $a_0 > 0$ such that for any $a \in (0, a_0)$ these maps are diffeomorphic between $\square_{k,a} := (0, \ell_k) \times (0, a)$ and $\Omega_{k,a} := \Phi_{k,a}(\square_{k,a})$, that $\Omega_{k,a} \subset \Omega$, and that $\Omega_{j,a} \cap \Omega_{k,a} = \emptyset$ for $j \neq k$. Note that the last property follows from the fact that the opening angles of the boundary corners (if any) are reflex. In addition, we set $\Omega_{0,a} := \Omega \setminus \left(\bigcup_{k=1}^n \overline{\Omega_{k,a}}\right)$.

Denote $\widetilde{H}^1_0(\Omega_{k,a}) := \{ f \in H^1(\Omega_{k,a}) : f \lceil_{\partial \Omega_{k,a} \setminus \overline{\Sigma}_k} = 0 \}, k = 1, \ldots, n, \text{ and introduce two new sesquilinear forms } h^{N/D,a}_{\beta} \text{ in } L^2(\Omega), \text{ both defined by the same expression as } h_{\beta} \text{ on the domains}$

$$\operatorname{dom} h_{\beta}^{N,a} = \bigoplus_{k=0}^{n} H^{1}(\Omega_{k,a}), \quad \operatorname{dom} h_{\beta}^{D,a} = H_{0}^{1}(\Omega_{0,a}) \cup \left(\bigoplus_{k=1}^{n} \widetilde{H}_{0}^{1}(\Omega_{k,a})\right),$$

and define

$$E_{N/D}(\beta, a) := \inf_{0 \neq f \in \text{dom } h_{\beta}^{N/D, a}} \frac{h_{\beta}^{N/D, a}(f, f)}{\|f\|_{L^{2}(\Omega)}^{2}}.$$

Due to the inclusions dom $h_{\beta}^{D,a} \subset \text{dom } h_{\beta} \subset \text{dom } h_{\beta}^{N,a}$, we have the inequalities

$$E_N(\beta, a) \le E(\beta) \le E_D(\beta, a).$$
 (12)

Furthermore, due to the fact that the parts $\Omega_{k,a}$ are disjoint and that the set $\Sigma \cap \partial \Omega_{0,a}$ is finite (this is exactly the set of the corners), we have the equality $E_{N/D}(\beta,a) = \min_{k \in \{0,\dots,n\}} E_{k,N/D}(\beta,a)$, with

$$E_{0,N}(\beta,a) := \inf_{0 \neq f \in H^{1}(\Omega_{0,a})} \frac{\|\nabla f\|_{L^{2}(\Omega_{0,a})}^{2}}{\|f\|_{L^{2}(\Omega_{0,a})}^{2}},$$

$$E_{k,N}(\beta,a) := \inf_{0 \neq f \in H^{1}(\Omega_{k,a})} \frac{\|\nabla f\|_{L^{2}(\Omega_{k,a})}^{2} - \beta \|f\|_{L^{2}(\Sigma_{k})}^{2}}{\|f\|_{L^{2}(\Omega_{k,a})}^{2}}, \quad k = 1, \dots, n,$$

$$E_{0,D}(\beta,a) = \inf_{0 \neq f \in H^{1}_{0}(\Omega_{0,a})} \frac{\|\nabla f\|_{L^{2}(\Omega_{0,a})}^{2}}{\|f\|_{L^{2}(\Omega_{0,a})}^{2}},$$

$$E_{k,D}(\beta,a) := \inf_{0 \neq f \in \widetilde{H}^{1}_{0}(\Omega_{k,a})} \frac{\|\nabla f\|_{L^{2}(\Omega_{k,a})}^{2} - \beta \|f\|_{L^{2}(\Sigma_{k})}^{2}}{\|f\|_{L^{2}(\Omega_{k,a})}^{2}}, \quad k = 1, \dots, n.$$

We have clearly $E_{0,N/D}(\beta,a) \ge 0$. Furthermore, in virtue of lemma 2 we can find a > 0 such that for each $k \in \{1,\ldots,n\}$ for $\beta \to +\infty$ we have

$$E_{k,N/D}(\beta,a) = -\beta^2 - \gamma_{k,\max}\beta + O(\beta^{\frac{2}{3}}), \quad \gamma_{k,\max} := \max_{s \in [0,\ell_k]} \gamma_k(s),$$

which gives $E_{N/D}(\beta, a) = -\beta^2 - \gamma_{\text{max}}\beta + O(\beta^{\frac{2}{3}})$, and the assertion of theorem 1 follows from the two-side estimate (12).

Remark 5. A more detailed asymptotic analysis is beyond the scope of the present note, but we mention one case in which the remainder estimate can be slightly improved with minimal efforts. Namely, assume that one of the following conditions is satisfied:

- the boundary Σ is of class C^4 (i.e. there are no corners),
- the curvature does not attain its maximal value $\gamma_{\rm max}$ at the corners,

then

$$E(\beta) = -\beta^2 - \gamma_{\text{max}}\beta + O(\sqrt{\beta}) \text{ as } \beta \to +\infty.$$
 (13)

Indeed, let us pick any $k \in \{1, ..., n\}$ such that $\gamma_{k, \max} = \gamma_{\max}$ and revise the proof of lemma 2 with $\Gamma := \Gamma_k$, $\kappa := \gamma_k$ and $\ell := \ell_k$. For any $s \in [0, \ell]$ with $\kappa(s) = \kappa_{\max}$ we have then $\kappa'(s) = 0$, and we may replace the inequality (10) with

$$\kappa_{M,l}^- \ge \kappa_{M,l}^+ - K\delta^2 = \kappa_{\max} - K\delta^2 \equiv \kappa_{\max} - \frac{K\ell^2}{M^2},$$

and by choosing $M \in \left[\sqrt[4]{\beta}, 2\sqrt[4]{\beta}\right] \cap \mathbb{N}$ we arrive at the estimate $E_{N/D}(\beta, a) = -\beta^2 - \kappa_{\max}\beta + O(\sqrt{\beta})$ as $\beta \to +\infty$, which in turn gives the asymptotics (13).

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References

- [1] G. Del Grosso, M. Campanino. A construction of the stochastic process associated to heat diffusion in a polygonal domain. *Bolletino Unione Mat. Ital.*, **13**-B, P. 876–895 (1976).
- [2] A.A. Lacey, J.R. Ockendon, J. Sabina. Multidimensional reaction diffusion equations with nonlinear boundary conditions. *SIAM J. Appl. Math.*, **58**, P. 1622–1647 (1998).
- [3] T. Giorgi, R. Smits. Eigenvalue estimates and critical temperature in zero fields for enhanced surface super-conductivity. *Z. Angew. Math. Phys.*, **57**, P. 1–22 (2006).
- [4] Y. Lou, M. Zhu. A singularly perturbed linear eigenvalue problem in C^1 domains. *Pacific J. Math.*, **214**, P. 323–334 (2004).
- [5] D. Daners, J.B. Kennedy. On the asymptotic behaviour of the eigenvalues of a Robin problem. *Differential Integr. Equ.*, **23**, P. 659–669 (2010).
- [6] M. Levitin, L. Parnovski. On the principal eigenvalue of a Robin problem with a large parameter. *Math. Nachr.*, **281**, P. 272–281 (2008).
- [7] H. Kovařík. Eigenvalue asymptotic of Robin Laplace operators on two-dimensional domains with cusps. *J. London Math. Soc.*, **83**, P. 256–271 (2011).
- [8] E. Colorado, J. García-Melián. The behavior of the principal eigenvalue of a mixed elliptic problem with respect to a parameter. *J. Math. Appl.*, **377**, P. 53–69 (2010).
- [9] D.Daners.Principal eigenvalues for generalised indefinite Robin problems. Potent. Anal. 38, P.1047-1069(2013)
- [10] R.L. Frank, L. Geisinger. Semi-classical analysis of the Laplace operator with Robin boundary conditions. *Bull. Math. Sci.*, **2**, P. 281–319 (2012).
- [11] T. Giorgi, R. Smits. Bounds and monotonicity for the generalized Robin problem. *Z. Angew. Math. Phys.*, **58**, P. 1–19 (2007).
- [12] H. Kovařík, A. Laptev. Hardy inequalities for Robin Laplacians. J. Funct. Anal., 262, P. 4972–4985 (2012).
- [13] P. Exner, K. Yoshitomi. Asymptotics of eigenvalues of the Schrödinger operator with a strong δ -interaction on a loop. *J. Geom. Phys.*, **41**, P. 344–358 (2002).
- [14] P. Exner, K. Pankrashkin. Strong coupling asymptotics for a singular Schrödinger operator with an interaction supported by an open arc. Submitted, preprint http://arxiv.org/abs/1207.2271.
- [15] V. Lotoreichik. Note on 2D Schrödinger operators with δ -interactions on angles and crossing lines. *Nanosystems: Phys. Chem. Math.*, **4** (2), P. 166–172 (2013).
- [16] B. Helffer, A. Morame. Magnetic bottles in connection with superconductivity. *J. Funct. Anal.*, **185**, P. 604–680 (2001).
- [17] V. Bonnaillie-Noël, M. Dauge. Asymptotics for the low-lying eigenstates of the Schrödinger operator with magnetic field near corners. *Ann. Henri Poincaré*, 7, P. 899–931 (2006).
- [18] S. Fournais, B. Helffer. Accurate eigenvalue estimates for the magnetic Neumann Laplacian. *Ann. Institut Fourier*, **56**, P. 1–67 (2006).