# ON THE $H^{1}-L^{1}$ BOUNDEDNESS OF OPERATORS 

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#### Abstract

We prove that if $q$ is in $(1, \infty), Y$ is a Banach space, and $T$ is a linear operator defined on the space of finite linear combinations of $(1, q)$-atoms in $\mathbb{R}^{n}$ with the property that $$
\sup \{\|T a\| Y: a \text { is a }(1, q) \text {-atom }\}<\infty
$$ then $T$ admits a (unique) continuous extension to a bounded linear operator from $H^{1}\left(\mathbb{R}^{n}\right)$ to $Y$. We show that the same is true if we replace $(1, q)$-atoms by continuous $(1, \infty)$-atoms. This is known to be false for $(1, \infty)$-atoms.


## 1. Introduction

In a recent paper, M. Bownik [3] showed that there exists a linear functional $F$ defined on finite linear combinations of $(1, \infty)$-atoms in $\mathbb{R}^{n}$ with the property that

$$
\sup \{|F(a)|: a \text { is a }(1, \infty) \text {-atom }\}<\infty
$$

but which does not admit a continuous extension to $H^{1}\left(\mathbb{R}^{n}\right)$. If $v$ is a fixed function in $L^{1}\left(\mathbb{R}^{n}\right) \backslash\{0\}$, then the operator $B$, defined on finite linear combinations of $(1, \infty)$ atoms by $B f=F(f) v$, satisfies

$$
\sup \left\{\|B a\|_{L^{1}\left(\mathbb{R}^{n}\right)}: a \text { is a }(1, \infty) \text {-atom }\right\}<\infty
$$

but does not admit an extension to a bounded operator from $H^{1}\left(\mathbb{R}^{n}\right)$ to $L^{1}\left(\mathbb{R}^{n}\right)$. This shows that the argument "the operator $T$ maps $(1, \infty)$-atoms uniformly into $L^{1}\left(\mathbb{R}^{n}\right)$, and hence it extends to a bounded operator from $H^{1}\left(\mathbb{R}^{n}\right)$ to $L^{1}\left(\mathbb{R}^{n}\right)$ " is fallacious.

Fortunately, if $T$ is a Calderón-Zygmund operator, then the uniform boundedness of $T$ on $(1, \infty)$-atoms implies the boundedness from $H^{1}\left(\mathbb{R}^{n}\right)$ to $L^{1}\left(\mathbb{R}^{n}\right)$ (see, for instance, [11, Ch. 7.3, Lemma 1], [2, Ch. 1.9], [7, Ch. III.7] and [8, Thm 6.7.1]).

The purpose of this paper is to show that the operator $B$ constructed above is, to a certain extent, pathological. Indeed, we prove that if $q$ is in $(1, \infty), Y$ is a Banach space, and $T$ is a linear operator defined on finite linear combinations of $(1, q)$-atoms in $\mathbb{R}^{n}$ with the property that

$$
\begin{equation*}
\sup \left\{\|T a\|_{Y}: a \text { is a }(1, q) \text {-atom }\right\}<\infty \tag{1.1}
\end{equation*}
$$

then $T$ admits a unique continuous extension to a bounded linear operator from $H^{1}\left(\mathbb{R}^{n}\right)$ to $Y$. The same conclusion holds if we assume that $T$ is a linear operator

[^0]on finite linear combinations of continuous $(1, \infty)$-atoms in $\mathbb{R}^{n}$ with the property that
\[

$$
\begin{equation*}
\sup \left\{\|T a\|_{Y}: a \text { is a continuous }(1, \infty) \text {-atom }\right\}<\infty \tag{1.2}
\end{equation*}
$$

\]

Note that this does not contradict Bownik's example. Indeed, the restriction of the operator $B$ to continuous $(1, \infty)$-atoms extends to a bounded operator $\widetilde{B}$ from $H^{1}\left(\mathbb{R}^{n}\right)$ to $L^{1}\left(\mathbb{R}^{n}\right)$. However, $B$ and $\widetilde{B}$ will agree on continuous $(1, \infty)$-atoms but not on all $(1, \infty)$-atoms.

To explain the idea of the proofs of these results, we need more notation. Suppose that $q$ is in $(1, \infty]$, and denote by $H_{\mathrm{fin}}^{1, q}\left(\mathbb{R}^{n}\right)$ the vector space of all finite linear combinations of $(1, q)$-atoms. Notice that $H_{\text {fin }}^{1, q}\left(\mathbb{R}^{n}\right)$ consists of all $L^{q}\left(\mathbb{R}^{n}\right)$ functions with compact support and integral 0. Clearly, $H_{\text {fin }}^{1, q}\left(\mathbb{R}^{n}\right)$ is a dense subspace of $H^{1}\left(\mathbb{R}^{n}\right)$. We may define a norm on $H_{\mathrm{fin}}^{1, q}\left(\mathbb{R}^{n}\right)$ as follows:

$$
\|f\|_{H_{\mathrm{fin}}^{1, q}\left(\mathbb{R}^{n}\right)}=\inf \left\{\sum_{j=1}^{N}\left|\lambda_{j}\right|: f=\sum_{j=1}^{N} \lambda_{j} a_{j}, a_{j} \text { is a }(1, q) \text {-atom, } N \in \mathbb{N}\right\} .
$$

Obviously $\|f\|_{H^{1}\left(\mathbb{R}^{n}\right)} \leq\|f\|_{H_{\mathrm{fin}}^{1, q}\left(\mathbb{R}^{n}\right)}$ for every $f$ in $H_{\mathrm{fin}}^{1, q}\left(\mathbb{R}^{n}\right)$. An example due to Y. Meyer (see [12, p. 513], Bownik's paper [3] or [7, p. 370]) shows that $\|\cdot\|_{H^{1}\left(\mathbb{R}^{n}\right)}$ and $\|\cdot\|_{H_{\text {fin }}^{1, \infty}\left(\mathbb{R}^{n}\right)}$ are inequivalent norms on $H_{\text {fin }}^{1, \infty}\left(\mathbb{R}^{n}\right)$. This is the starting point of Bownik's construction.

We prove that Meyer's example itself is somewhat exceptional. Indeed, by using the maximal characterisation of $H^{1}\left(\mathbb{R}^{n}\right)$, we show that if $q<\infty$, then $\|\cdot\|_{H^{1}\left(\mathbb{R}^{n}\right)}$ and $\|\cdot\|_{H_{\text {fin }}^{1, q}\left(\mathbb{R}^{n}\right)}$ are equivalent norms on $H_{\text {fin }}^{1, q}\left(\mathbb{R}^{n}\right)$ (see Section 3). Similarly, we prove that $\|\cdot\|_{H^{1}\left(\mathbb{R}^{n}\right)}$ and $\|\cdot\|_{H_{\text {fin }}^{1, \infty}\left(\mathbb{R}^{n}\right)}$ are equivalent norms on $H_{\text {fin }}^{1, \infty}\left(\mathbb{R}^{n}\right) \cap C\left(\mathbb{R}^{n}\right)$. This immediately implies that operators defined on $H_{\mathrm{fin}}^{1, q}\left(\mathbb{R}^{n}\right)$ which have either property (1.1) or property (1.2) automatically extend to bounded operators from $H^{1}\left(\mathbb{R}^{n}\right)$ to $L^{1}\left(\mathbb{R}^{n}\right)$.

As discussed briefly in Section 3, this equivalence of norms remains true for $H^{p}\left(\mathbb{R}^{n}\right)$ with $0<p<1$ and $(p, q)$-atoms.

The extension property for operators was also proved, by different methods, for $0<p \leq 1$ and $(p, 2)$-atoms and operators taking values in quasi-Banach spaces, by D. Yang and Y. Zhou [17].

A theory of Hardy spaces has been developed in spaces of homogeneous type; see R.R. Coifman and G. Weiss [5]. It is, however, not evident whether our results extend to this case in general. Nevertheless, let $M$ be such a space. By a simple functional analysis argument, we show that if $q$ is in $(1, \infty)$ and $T$ is an operator defined on $H_{\text {fin }}^{1, q}(M)$ satisfying the analogue of (1.1), then $T$ automatically extends to a bounded operator from $H^{1}(M)$ to $L^{1}(M)$ (see Section (4). It may be worth noticing that the proof of this result also applies to certain metric measured spaces ( $M, \rho, \mu$ ) where $\mu$ is only "locally doubling" [10, 4], and [16].

For so-called RD-spaces, which are spaces of homogeneous type having "dimension $n "$ in a certain sense, our complete results were recently extended in the paper [9] by L. Grafakos, L. Liu and Yang. These authors consider $n /(n+1)<p \leq 1$ and quasi-Banach-valued operators.

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## 2. Notation and terminology

Suppose that $(M, \rho, \mu)$ is a space of homogeneous type in the sense of Coifman and Weiss [5] and that $\mu$ is a $\sigma$-finite measure. For the sake of simplicity, we shall assume that $\mu(M)$ is infinite.

Suppose that $q$ is in $(1, \infty]$. For each closed ball $B$ in $M$, we denote by $L_{0}^{q}(B)$ the space of all functions in $L^{q}(M)$ which are supported in $B$ and have integral 0 . Clearly $L_{0}^{q}(B)$ is a closed subspace of $L^{q}(M)$. The union of all spaces $L_{0}^{q}(B)$ as $B$ varies over all balls coincides with the space $L_{c, 0}^{q}(M)$ of all functions in $L^{q}(M)$ with compact support and integral 0 . Fix a reference point $o$ in $M$ and for each positive integer $k$ denote by $B_{k}$ the ball centred at $o$ with radius $k$. A convenient way of topologising $L_{c, 0}^{q}(M)$ is to interpret $L_{c, 0}^{q}(M)$ as the strict inductive limit of the spaces $L_{c, 0}^{q}\left(B_{k}\right)$ (see [1, II, p. 33] for the definition of the strict inductive limit topology). We denote by $X^{q}$ the space $L_{c, 0}^{q}(M)$ with this topology, and write $X_{k}^{q}$ for $L_{c, 0}^{q}\left(B_{k}\right)$.

We recall the basic definitions and results concerning the atomic Hardy space $H^{1}(M)$. The reader is referred to [5] and the references therein for this and more on Hardy spaces defined on spaces of homogeneous type. Suppose that $q$ is in $(1, \infty]$. A $(1, q)$-atom is a function $a$ in $L^{q}(M)$ supported in a ball $B$, with mean value 0 and such that

$$
\left(\frac{1}{\mu(B)} \int_{B}|a|^{q} \mathrm{~d} \mu\right)^{1 / q} \leq \mu(B)^{-1}
$$

if $q$ is finite, and $\|a\|_{\infty} \leq \mu(B)^{-1}$ if $q=\infty$. We denote by $H^{1, q}(M)$ the space of all functions $g$ in $L^{1}(M)$ which admit a decomposition of the form $g=\sum_{j} \lambda_{j} a_{j}$, where the $a_{j}$ are $(1, q)$-atoms and the $\lambda_{j}$ are complex numbers such that $\sum_{j}\left|\lambda_{j}\right|<$ $\infty$. The norm $\|g\|_{H^{1, q}}$ of $g$ in $H^{1, q}(M)$ is the infimum of $\sum_{j}\left|\lambda_{j}\right|$ over all such decompositions. It is well known that all the spaces $H^{1, q}(M)$ with $q \in(1, \infty)$ coincide with $H^{1, \infty}(M)$, and we denote them all by $H^{1}(M)$. Clearly, the vector space $H_{\mathrm{fin}}^{1, q}(M)$ of all finite linear combinations of $(1, q)$-atoms is dense in $H^{1}(M)$ with respect to the norm of $H^{1}(M)$, for $q$ in $(1, \infty]$. Observe also that $H_{\text {fin }}^{1, q}(M)$ and $L_{c, 0}^{q}(M)$ agree as vector spaces, and so do the space of finite linear combinations of continuous $(1, \infty)$-atoms and $H_{\text {fin }}^{1, \infty}(M) \cap C\left(\mathbb{R}^{n}\right)$.

For each ball $B$ and each locally integrable function $f$, we denote by $f_{B}$ the average of $f$ on $B$. Recall that $B M O$ is the Banach space of all locally integrable functions $f$, defined modulo constants, such that

$$
\|f\|_{B M O}=\sup _{B} \frac{1}{\mu(B)} \int_{B}\left|f-f_{B}\right| \mathrm{d} \mu<\infty
$$

The dual of $H^{1}(M)$ may be identified with $B M O$.
There are several characterisations of the space $H^{1}\left(\mathbb{R}^{n}\right)$. We shall make use of the so-called maximal characterisation, which we briefly recall. Suppose that $m$ is an integer with $m>n$, and denote by $\mathcal{A}_{m}$ the set of all functions $\varphi$ in the Schwartz space $\mathcal{S}\left(\mathbb{R}^{n}\right)$ such that

$$
\sup _{|\beta| \leq m} \sup _{x \in \mathbb{R}^{n}}(1+|x|)^{m}\left|D^{\beta} \varphi(x)\right| \leq 1
$$

where $|\beta|$ denotes the length of the multi-index $\beta$. For $\varphi$ in $\mathcal{S}\left(\mathbb{R}^{n}\right)$ denote by $\varphi_{t}$ the function $t^{-n} \varphi(\cdot / t)$. Given $f$ in $L^{1}\left(\mathbb{R}^{n}\right)$, define the "grand maximal function"
$\mathcal{M}_{m} f$ by

$$
\mathcal{M}_{m} f=\sup _{\varphi \in \mathcal{A}_{m}} \sup _{t>0}\left|\varphi_{t} * f\right| .
$$

The following result is classical [6, [13, 7, and 15].
Theorem 2.1. Suppose that $f$ is in $L^{1}\left(\mathbb{R}^{n}\right)$. The following are equivalent:
(i) $f$ is in $H^{1}\left(\mathbb{R}^{n}\right)$;
(ii) the grand maximal function $\mathcal{M}_{m} f$ is in $L^{1}\left(\mathbb{R}^{n}\right)$.

Furthermore, $f \mapsto\left\|\mathcal{M}_{m} f\right\|_{L^{1}\left(\mathbb{R}^{n}\right)}$ is an equivalent norm on $H^{1}\left(\mathbb{R}^{n}\right)$.
The letter $C$ will denote a positive constant, which need not be the same at different occurrences. Given two positive quantities $A$ and $B$, we shall mean by $A \sim B$ that there exists a constant $C$ such that $1 / C \leq A / B \leq C$.

## 3. The Euclidean case

In this section we work in the classical setting of $\mathbb{R}^{n}$.
Theorem 3.1. The following hold:
(i) if $q<\infty$, then $\|\cdot\|_{H_{\text {fin }}^{1, q}\left(\mathbb{R}^{n}\right)}$ and $\|\cdot\|_{H^{1}\left(\mathbb{R}^{n}\right)}$ are equivalent norms on $H_{\text {fin }}^{1, q}\left(\mathbb{R}^{n}\right)$;
(ii) the two norms $\|\cdot\|_{H_{\text {fin }}^{1, \infty}\left(\mathbb{R}^{n}\right)}$ and $\|\cdot\|_{H^{1}\left(\mathbb{R}^{n}\right)}$ are equivalent on $H_{\text {fin }}^{1, \infty}\left(\mathbb{R}^{n}\right) \cap$ $C\left(\mathbb{R}^{n}\right)$.
Proof. Clearly, $\|f\|_{H^{1}\left(\mathbb{R}^{n}\right)} \leq\|f\|_{H_{\text {fin }}^{1, q}\left(\mathbb{R}^{n}\right)}$ for $f$ in $H_{\text {fin }}^{1, q}\left(\mathbb{R}^{n}\right)$ and for $q$ in $(1, \infty]$. Thus, we have to show that for every $q$ in $(1, \infty)$ there exists a constant $C$ such that

$$
\|f\|_{H_{\mathrm{fin}}^{1, q}\left(\mathbb{R}^{n}\right)} \leq C\|f\|_{H^{1}\left(\mathbb{R}^{n}\right)} \quad \forall f \in H_{\mathrm{fin}}^{1, q}\left(\mathbb{R}^{n}\right)
$$

and that a similar estimate holds for $q=\infty$ and all $f$ in $H_{\text {fin }}^{1, \infty}\left(\mathbb{R}^{n}\right) \cap C\left(\mathbb{R}^{n}\right)$.
Suppose that $q$ is in $(1, \infty]$ and that $f$ is in $H_{\mathrm{fin}}^{1, q}\left(\mathbb{R}^{n}\right)$ with $\|f\|_{H^{1}\left(\mathbb{R}^{n}\right)}=1$. By the translation invariance of Lebesgue measure, we may assume that the support of $f$ is contained in the closed ball $B=B(0, R)$ centred at 0 with radius $R$. For each $k$ in $\mathbb{Z}$, denote by $\Omega_{k}$ the level set $\left\{x \in \mathbb{R}^{n}: \mathcal{M}_{m} f(x)>2^{k}\right\}$ of the grand maximal function $\mathcal{M}_{m} f$ of $f$. We choose Whitney cubes $Q_{i}^{k}, i \in \mathbb{N}$, with disjoint interiors satisfying $\Omega_{k}=\bigcup_{i} Q_{i}^{k}$ and

$$
\begin{equation*}
\operatorname{diam}\left(Q_{i}^{k}\right) \leq \eta \operatorname{dist}\left(Q_{i}^{k}, \Omega_{k}^{c}\right) \leq 4 \operatorname{diam}\left(Q_{i}^{k}\right) \tag{3.1}
\end{equation*}
$$

where $\eta$ is a suitable constant in $(0,1)$. Except for the factor $\eta$, this is Theorem VI. 1 of [14, p. 167]. The only modification needed in the proof of [14] concerns the choice of the constant denoted $c$.

By following closely the proof of [15, Theorem III.2, p. 107] or [13, Theorem 3.5, pp. 12-18], we produce an atomic decomposition of $f$ of the form

$$
\begin{equation*}
f=\sum_{i, k} \lambda_{i}^{k} a_{i}^{k} \tag{3.2}
\end{equation*}
$$

such that the following hold:
(a) $\left|\lambda_{i}^{k} a_{i}^{k}\right| \leq C 2^{k}$ for every $k$ in $\mathbb{Z}$;
(b) for each $k$ in $\mathbb{Z}$, the atoms $a_{i}^{k}$ are supported in balls $B_{i}^{k}$ concentric with the $Q_{i}^{k}$ and contained in $\Omega_{k}$. By choosing the constant $\eta$ in (3.1) small enough, depending on the dimension, we can also ensure that the family $\left\{B_{i}^{k}\right\}_{i}$ has the bounded overlap property, uniformly with respect to $k$;
(c) there exists a constant $C$ independent of $f$ such that

$$
\sum_{i, k}\left|\lambda_{i}^{k}\right| \leq C\|f\|_{H^{1}\left(\mathbb{R}^{n}\right)}=C
$$

We write $2 B$ for the closed ball concentric with $B$ whose radius is twice as large. For $\varphi$ in $\mathcal{A}_{m}$ and $x$ in $\mathbb{R}^{n} \backslash(2 B)$ one then has

$$
\begin{aligned}
\left|\varphi_{t} * f(x)\right| & \leq t^{-n} \sup _{y \in B^{c}}|\varphi(y / t)|\|f\|_{L^{1}\left(\mathbb{R}^{n}\right)} \\
& \leq t^{-n}(1+R / t)^{-m}\|f\|_{L^{1}\left(\mathbb{R}^{n}\right)} \quad \forall t \in \mathbb{R}^{+},
\end{aligned}
$$

so that

$$
\mathcal{M}_{m} f(x)=\sup _{\varphi \in \mathcal{A}_{m}} \sup _{t>R}\left|\varphi_{t} * f(x)\right| \leq R^{-n}
$$

since $m>n$. Now, if $x$ is in $\Omega_{k} \backslash(2 B)$, the above inequality and the definition of $\Omega_{k}$ force $2^{k}<R^{-n}$; denote by $k^{\prime}$ the largest integer $k$ such that $2^{k}<R^{-n}$. Then $\overline{\Omega_{k}}$ is contained in $2 B$ for $k>k^{\prime}$.

Next we define the functions $h$ and $\ell$ by

$$
\begin{equation*}
h=\sum_{k \leq k^{\prime}} \sum_{i} \lambda_{i}^{k} a_{i}^{k} \quad \text { and } \quad \ell=\sum_{k>k^{\prime}} \sum_{i} \lambda_{i}^{k} a_{i}^{k} \tag{3.3}
\end{equation*}
$$

Observe that both these series converge in $L^{1}\left(\mathbb{R}^{n}\right)$, simply because $\sum_{i, k}\left|\lambda_{i}^{k}\right|<\infty$, so that $h$ and $\ell$ have integral 0 . Clearly, $f=h+\ell$. Furthermore, the support of $\ell$ is contained in $2 B$, because it is contained in $\bar{\Omega}_{k}$ by (b) above, and $\bar{\Omega}_{k}$ is contained in $2 B$ for all $k>k^{\prime}$. Therefore $h=f=0$ in $(2 B)^{c}$.

To estimate the size of $h$ in $2 B$, we use (a) above and the bounded overlap property of (b), getting

$$
|h| \leq C \sum_{k \leq k^{\prime}} 2^{k} \leq C 2^{k^{\prime}} \leq C|2 B|^{-1}
$$

This proves that $h / C$ is a $(1, \infty)$-atom, where $C$ is independent of $f$.
Now we assume that $q<\infty$ and conclude the proof of (i). Observe that $\ell$ is in $L^{q}\left(\mathbb{R}^{n}\right)$, because $\ell=f-h$, and both $f$ and $h$ are in $L^{q}\left(\mathbb{R}^{n}\right)$.

We claim that the series $\sum_{k>k^{\prime}} \sum_{i} \lambda_{i}^{k} a_{i}^{k}$ converges to $\ell$ in $L^{q}\left(\mathbb{R}^{n}\right)$.
Fixing $s$ in $\mathbb{Z}$, we shall estimate $\sum_{k>k^{\prime}} \sum_{i}\left|\lambda_{i}^{k} a_{i}^{k}\right|$ in $\Omega_{s} \backslash \Omega_{s+1}$. First observe that all terms with $k>s$ vanish outside $\Omega_{s+1}$. Then apply (a) and (b) to get the pointwise bound

$$
\sum_{k>k^{\prime}} \sum_{i}\left|\lambda_{i}^{k} a_{i}^{k}\right| \leq C \sum_{k \leq s} 2^{k} \leq C 2^{s} \leq C \mathcal{M}_{m} f
$$

The constants $C$ above are independent of $f$ and $s$, so that

$$
\sum_{k>k^{\prime}} \sum_{i}\left|\lambda_{i}^{k} a_{i}^{k}\right| \leq C \mathcal{M}_{m} f
$$

in all of $\mathbb{R}^{n}$, with $C$ independent of $f$. Note that $\mathcal{M}_{m} f$ is in $L^{q}\left(\mathbb{R}^{n}\right)$, since $f$ is. This implies that the series defining $\ell$ converges almost everywhere and the limit must coincide with the $L^{1}$ limit $\ell$. The Lebesgue dominated convergence theorem now implies that $\sum_{k>k^{\prime}} \sum_{i} \lambda_{i}^{k} a_{i}^{k}$ converges to $\ell$ in $L^{q}\left(\mathbb{R}^{n}\right)$, and the claim is proved.

Finally, for each positive integer $N$ we denote by $F_{N}$ the finite set of all pairs of integers $(i, k)$ such that $k>k^{\prime}$ and $|i|+|k| \leq N$, and by $\ell_{N}$ the function $\sum_{(i, k) \in F_{N}} \lambda_{i}^{k} a_{i}^{k}$. The function $\ell_{N}$ is in $H_{\text {fin }}^{1, q}\left(\mathbb{R}^{n}\right)$, and $f=h+\ell_{N}+\left(\ell-\ell_{N}\right)$.

Observe that $\ell-\ell_{N}$ will be a small multiple of a $(1, q)$-atom for large $N$. Indeed, by taking $N$ large enough, we can make the corresponding coefficient less than any given $\varepsilon$ in $\mathbb{R}^{+}$. Then

$$
\|f\|_{H_{\mathrm{fin}}^{1, q}\left(\mathbb{R}^{n}\right)} \leq C+\sum_{(i, k) \in F_{N}}\left|\lambda_{i}^{k}\right|+\varepsilon
$$

so that

$$
\|f\|_{H_{\mathrm{fin}}^{1, q}\left(\mathbb{R}^{n}\right)} \leq C+\sum_{(i, k) \in F_{N}}\left|\lambda_{i}^{k}\right| \leq C
$$

by property (c) above, as required to conclude the proof of (i).
Now we finish the proof of (ii). Assume that $f$ is a continuous function in $H_{\text {fin }}^{1, \infty}\left(\mathbb{R}^{n}\right)$. A careful examination of the proof of [15, Theorem III.2, pp. 1078] or [13, Theorem 3.5, pp. 12-18] shows that the atoms $a_{i}^{k}$ that appear in the decomposition (3.2) are then continuous. Furthermore, we see that for each $k$ and $i$ the function $\lambda_{i}^{k} a_{i}^{k}$ depends only on the restriction of $f$ to a ball $\tilde{B}_{i}^{k}$ which is a concentric enlargement of the ball $B_{i}^{k}$ from (b) above, by a fixed scaling factor. It is straightforward to check that if $f$ is constant in $\tilde{B}_{i}^{k}$, then $\lambda_{i}^{k} a_{i}^{k}=0$ and that there exists an absolute constant $C$ such that if $|f|<\varepsilon$ in $\tilde{B}_{i}^{k}$, then $\left|\lambda_{i}^{k} a_{i}^{k}\right|<C \varepsilon$.

Since trivially $\mathcal{M}_{m} f \leq C_{n}\|f\|_{\infty}$, where the constant $C_{n}$ depends only on $n$, the level set $\Omega_{k}$ is empty for all $k$ such that $2^{k} \geq C_{n}\|f\|_{\infty}$. We denote by $k^{\prime \prime}$ the largest integer for which the last inequality does not hold. Then the index $k$ in the sum defining $\ell$ in (3.3) will run only over $k^{\prime}<k \leq k^{\prime \prime}$.

Let $\varepsilon$ be positive. Since $f$ is uniformly continuous, there exists a positive $\delta$ such that $|x-y|<\delta$ implies

$$
|f(x)-f(y)|<\varepsilon
$$

Write $\ell=\ell_{1}^{\varepsilon}+\ell_{2}^{\varepsilon}$ with

$$
\ell_{1}^{\varepsilon}=\sum_{(i, k) \in F_{1}} \lambda_{i}^{k} a_{i}^{k} \quad \text { and } \quad \ell_{2}^{\varepsilon}=\sum_{(i, k) \in F_{2}} \lambda_{i}^{k} a_{i}^{k}
$$

where $F_{1}=\left\{(i, k): \operatorname{diam}\left(\tilde{B}_{i}^{k}\right) \geq \delta, k^{\prime}<k \leq k^{\prime \prime}\right\}$ and $F_{2}=\left\{(i, k): \operatorname{diam}\left(\tilde{B}_{i}^{k}\right)<\right.$ $\left.\delta, k^{\prime}<k \leq k^{\prime \prime}\right\}$. Since $F_{1}$ is a finite set, $\ell_{1}^{\varepsilon}$ is continuous.

To estimate $\ell_{2}^{\varepsilon}$, we denote by $x_{i}^{k}$ the centre of the ball $B_{i}^{k}$ and write for $(i, k)$ in $F_{2}$

$$
f(x)=f\left(x_{i}^{k}\right)+f(x)-f\left(x_{i}^{k}\right)
$$

Then $\left|\lambda_{i}^{k} a_{i}^{k}\right|<C \varepsilon$, because $\left|f(x)-f\left(x_{i}^{k}\right)\right|<\varepsilon$ for $x$ in $\tilde{B}_{i}^{k}$. For fixed $k$ the balls $\left\{B_{i}^{k}\right\}_{i}$ have uniformly bounded overlap, so there exists an absolute constant $C$ such that

$$
\left|\ell_{2}^{\varepsilon}\right| \leq C \sum_{k^{\prime}<k \leq k^{\prime \prime}} \varepsilon \leq C\left(k^{\prime \prime}-k^{\prime}\right) \varepsilon
$$

Since $\varepsilon$ is arbitrary, we can thus split $\ell$ into a continuous part and a part that is uniformly arbitarily small. It follows that $\ell$ is continuous. But then $h=f-\ell$ is also continuous, so that $h$ is a continuous $(1, \infty)$-atom, multiplied by a factor $C$.

To find a finite atomic decomposition of $\ell$, we again use the splitting $\ell=\ell_{1}^{\varepsilon}+\ell_{2}^{\varepsilon}$. Clearly $\ell_{1}^{\varepsilon}$ is for each $\varepsilon$ a finite linear combination of continuous $(1, \infty)$-atoms, and the $\ell^{1}$ norm of the coefficients is controlled by $\|f\|_{H^{1}}$, in view of (c). Observe that $\ell_{2}^{\varepsilon}=\ell-\ell_{1}^{\varepsilon}$ is continuous. Further, $\ell_{2}^{\varepsilon}$ is supported in $2 B$, has integral 0 and satisfies $\left|\ell_{2}^{\varepsilon}\right| \leq C\left(k^{\prime \prime}-k^{\prime}\right) \varepsilon$. Choosing $\varepsilon$, we can thus make $\ell_{2}^{\varepsilon}$ into an arbitrarily small multiple of a continuous $(1, \infty)$-atom.

To sum up, $f=h+\ell_{1}^{\varepsilon}+\ell_{2}^{\varepsilon}$ gives the desired finite atomic decomposition of $f$, with coefficients controlled by $\|f\|_{H^{1}}$.

We have completed the proof of (ii) and that of the theorem.
Remark 3.2. Theorem 3.1 (ii) implies that any function $f$ in $H_{\text {fin }}^{1, \infty}\left(\mathbb{R}^{n}\right) \cap C\left(\mathbb{R}^{n}\right)$ admits a finite decomposition in $(1, \infty)$-atoms such that the sum of the corresponding coefficients is $\leq C\|f\|_{H^{1}\left(\mathbb{R}^{n}\right)}$. Actually, the proof of Theorem 3.1 (ii) shows that we can construct this finite decomposition in such a way that it involves only continuous ( $1, \infty$ )-atoms.

Remark 3.3. Theorem 3.1 extends to $H^{p}\left(\mathbb{R}^{n}\right)$ with $0<p<1$ and $(p, q)$-atoms, where one can now have $1 \leq q \leq \infty$. The proof is rather similar to the one given above, so we only briefly describe the modifications needed for part (i). Thus let $1 \leq q<\infty$. Given $f \in H_{\mathrm{fin}}^{p, q}\left(\mathbb{R}^{n}\right)$ supported in a ball $B_{R}$, the first step is the inequality $\mathcal{M}_{m} f \leq C R^{-n / p}\|f\|_{H^{p}\left(\mathbb{R}^{n}\right)}$, valid outside a larger ball $B_{C R}$. One proves this by comparing the values of $\mathcal{M}_{m} f$ at different points and using the fact that $\left\|\mathcal{M}_{m} f\right\|_{L^{p}\left(\mathbb{R}^{n}\right)} \sim\|f\|_{H^{p}\left(\mathbb{R}^{n}\right)}$. Then the $\Omega_{k}$ and the decompositions $f=\sum \lambda_{i}^{k} a_{i}^{k}=$ $h+\ell$ are introduced as above. The sum $\ell$ now converges in $\mathcal{S}^{\prime}$ and is dominated by $\mathcal{M}_{m} f$. If $q>1$, we have $\mathcal{M}_{m} f \in L^{q}\left(\mathbb{R}^{n}\right)$ and conclude as before that $\ell$ converges in $L^{q}\left(\mathbb{R}^{n}\right)$. For $q=1$, the tail sum $S_{\kappa}=\sum_{k \geq \kappa} \sum_{i} \lambda_{i}^{k} a_{i}^{k}$ tends to 0 in $L^{1}\left(\mathbb{R}^{n}\right)$ as $\kappa \rightarrow+\infty$, because $S_{\kappa}$ is nonzero only in $\Omega_{\kappa}$ and not larger than $|f|+C 2^{\kappa}$ there, and $\left|\Omega_{\kappa}\right|=o\left(2^{-\kappa}\right)$ as $\kappa \rightarrow+\infty$. The rest of the proof proceeds as before. See also [9, Theorem 5.6].

Corollary 3.4. Suppose that $Y$ is a Banach space and that one of the following holds:
(i) $q$ is in $(1, \infty)$ and $T: H_{\mathrm{fin}}^{1, q}\left(\mathbb{R}^{n}\right) \rightarrow Y$ is a linear operator such that

$$
A:=\sup \left\{\|T a\|_{Y}: a \text { is a }(1, q) \text {-atom }\right\}<\infty ;
$$

(ii) $T$ is a $Y$-valued linear operator defined on continuous $(1, \infty)$-atoms such that

$$
A:=\sup \left\{\|T a\|_{Y}: a \text { is a continuous }(1, \infty) \text {-atom }\right\}<\infty .
$$

Then there exists a unique bounded linear operator $\widetilde{T}$ from $H^{1}\left(\mathbb{R}^{n}\right)$ to $Y$ which extends $T$.

Proof. We consider the case (i). Suppose that $f$ is in $H_{\mathrm{fin}}^{1, q}\left(\mathbb{R}^{n}\right), f=\sum_{j=1}^{N} \lambda_{j} a_{j}$ say, where $a_{j}$ are $(1, q)$-atoms. Then the assumption and the triangle inequality give

$$
\|T f\|_{Y} \leq A \sum_{j=1}^{N}\left|\lambda_{j}\right|
$$

By taking the infimum of the right-hand side with respect to all decompositions of $f$ as a finite sum of $(1, q)$-atoms, we obtain

$$
\|T f\|_{Y} \leq A\|f\|_{H_{\mathrm{fin}}^{1, q}\left(\mathbb{R}^{n}\right)}
$$

Now, Theorem3.1(i) implies that the right-hand side is dominated by $C A\|f\|_{H^{1}\left(\mathbb{R}^{n}\right)}$, where $C$ does not depend on $f$, and a density argument completes the proof of the corollary.

The case (ii) is similar.

Remark 3.5. The statement of Corollary 3.4 (i) becomes false if we replace $q$ by $\infty$. A counterexample is given by the operator $B$ defined in the Introduction. Note also that Corollary 3.4 applies to linear functionals.

## 4. Results on spaces of homogeneous type

In this section, we work in a space of homogeneous type $(M, \rho, \mu)$. Recall that we assume that $\mu$ is $\sigma$-finite and that $\mu(M)$ is infinite.
Theorem 4.1. Suppose that $q$ is in $(1, \infty)$ and that $T$ is a linear operator defined on $H_{\text {fin }}^{1, q}(M)$ with the property that

$$
A:=\sup \left\{\|T a\|_{L^{1}(M)}: a \text { is } a(1, q) \text {-atom }\right\}<\infty .
$$

Then there exists a unique bounded linear operator $\widetilde{T}$ from $H^{1}(M)$ to $L^{1}(M)$ which extends $T$.

Proof. We prove the result in the case where $q=2$. The proof in the other cases is similar.

Suppose that $B$ is a ball. For each $f$ in $L_{0}^{2}(B)$ such that $\|f\|_{L^{2}(M)}=1$, the function $\mu(B)^{-1 / 2} f$ is a $(1,2)$-atom, so that

$$
\|T f\|_{L^{1}(M)} \leq A \mu(B)^{1 / 2} \quad \forall f \in L_{0}^{2}(B)
$$

by the assumption. In particular, the restriction of $T$ to $X_{k}^{2}$ is bounded from $X_{k}^{2}$ to $L^{1}(M)$ for each $k$. Thus, $T$ is bounded from $X^{2}$ to $L^{1}(M)$. It follows that $T^{*}$ is bounded from $L^{\infty}(M)$ to the dual of $X^{2}$. But the dual of $X^{2}$ is the quotient space $L_{\text {loc }}^{2}(M) / \mathbb{C}$, since that of $L_{c, 0}^{2}\left(B_{k}\right)$ is $L^{2}\left(B_{k}\right) / \mathbb{C}$. Now, for every $f$ in $L^{\infty}(M)$ and for every (1,2)-atom $a$,

$$
\langle T a, f\rangle=\left\langle a, T^{*} f\right\rangle=\int_{M} a T^{*} f \mathrm{~d} \mu
$$

so that

$$
\left|\int_{M} a T^{*} f \mathrm{~d} \mu\right|=|\langle T a, f\rangle| \leq A\|f\|_{\infty}
$$

A standard argument then shows that $T^{*} f$ belongs to $B M O(M)$ and that

$$
\begin{equation*}
\left\|T^{*} f\right\|_{B M O(M)} \leq 2 A\|f\|_{\infty} \quad \forall f \in L^{\infty}(M) \tag{4.1}
\end{equation*}
$$

We give the details for the reader's convenience. Suppose that $B$ is a ball and observe that

$$
\left[\int_{B}\left|T^{*} f-\left(T^{*} f\right)_{B}\right|^{2} \mathrm{~d} \mu\right]^{1 / 2}=\sup _{\|\varphi\|_{L^{2}(B)}=1}\left|\int_{B} \varphi\left(T^{*} f-\left(T^{*} f\right)_{B}\right) \mathrm{d} \mu\right|
$$

But

$$
\begin{aligned}
\int_{B} \varphi\left(T^{*} f-\left(T^{*} f\right)_{B}\right) \mathrm{d} \mu & =\int_{B}\left(\varphi-\varphi_{B}\right)\left(T^{*} f-\left(T^{*} f\right)_{B}\right) \mathrm{d} \mu \\
& =\int_{B}\left(\varphi-\varphi_{B}\right) T^{*} f \mathrm{~d} \mu
\end{aligned}
$$

and since $\|\varphi\|_{L^{2}(B)}=1$,

$$
\left|\varphi_{B}\right| \leq\left[\frac{1}{\mu(B)} \int_{B}|\varphi|^{2} \mathrm{~d} \mu\right]^{1 / 2} \leq \mu(B)^{-1 / 2}
$$

Write $\psi$ instead of $\varphi-\varphi_{B}$. Then

$$
\|\psi\|_{L^{2}(B)} \leq\|\varphi\|_{L^{2}(B)}+\left|\varphi_{B}\right| \mu(B)^{1 / 2} \leq 2
$$

so that $\psi /\left(2 \mu(B)^{1 / 2}\right)$ is a $(1,2)$-atom. Therefore

$$
\left|\int_{B} \psi T^{*} f \mathrm{~d} \mu\right| \leq 2 A \mu(B)^{1 / 2}\|f\|_{\infty}
$$

Combining the above, we conclude that for every ball $B$

$$
\left[\frac{1}{\mu(B)} \int_{B}\left|T^{*} f-\left(T^{*} f\right)_{B}\right|^{2} \mathrm{~d} \mu\right]^{1 / 2} \leq 2 A\|f\|_{\infty}
$$

and (4.1) follows.
Now we show that $T$ extends to a bounded operator from $H^{1}(M)$ to $L^{1}(M)$ with norm at most $2 A$. Observe that $X^{2}$ and $H_{\text {fin }}^{1,2}(M)$ coincide as vector spaces. For every $g$ in $H_{\text {fin }}^{1,2}(M)$ and for every $f$ in $L^{\infty}(M)$

$$
\begin{aligned}
|\langle T g, f\rangle| & =\left|\left\langle g, T^{*} f\right\rangle\right| \\
& \leq\|g\|_{H^{1}(M)}\left\|T^{*} f\right\|_{B M O(M)} \\
& \leq 2 A\|g\|_{H^{1}(M)}\|f\|_{L^{\infty}(M)} .
\end{aligned}
$$

By taking the supremum of both sides over all functions $f$ in $L^{\infty}(M)$ with $\|f\|_{L^{\infty}(M)}$ =1, we obtain that

$$
\|T g\|_{L^{1}(M)} \leq 2 A\|g\|_{H^{1}(M)} \quad \forall g \in H_{\mathrm{fin}}^{1,2}(M)
$$

Finally we observe that $H_{\text {fin }}^{1,2}(M)$ is dense in $H^{1}(M)$ (with respect to the norm of $H^{1}(M)$ ), and the required conclusion follows by a density argument.

Quite often one encounters the following situation. Suppose that $T$ is a bounded linear operator on $L^{2}(M)$. Then $T$ is automatically defined on $H_{\text {fin }}^{1,2}(M)$. Assume that

$$
A:=\sup \left\{\|T a\|_{L^{1}(M)}: a \text { is a }(1,2) \text {-atom }\right\}<\infty .
$$

By the previous result, the restriction of $T$ to $H_{\text {fin }}^{1,2}(M)$ has a unique extension to a bounded linear operator $\widetilde{T}$ from $H^{1}(M)$ to $L^{1}(M)$. The question is whether the operators $T$ and $\widetilde{T}$ are consistent, i.e., whether they coincide on the intersection $H^{1}(M) \cap L^{2}(M)$ of their domains. The answer to this question is in the affirmative, as the following proposition shows.

Proposition 4.2. Suppose that $T$ is bounded on $L^{2}(M)$ and that

$$
A:=\sup \left\{\|T a\|_{L^{1}(M)}: a \text { is a }(1,2) \text {-atom }\right\}<\infty
$$

Denote by $\widetilde{T}$ the unique continuous linear extension of the restriction of $T$ to $H_{\mathrm{fin}}^{1,2}(M)$ to an operator from $H^{1}(M)$ to $L^{1}(M)$. Then the operators $T$ and $\widetilde{T}$ agree on $H^{1}(M) \cap L^{2}(M)$.

Proof. Suppose that $f$ is in $L^{2}(M) \cap L^{\infty}(M)$ and that $g$ is in $L_{c, 0}^{2}(M)$. Denote by $T^{*}$ the transpose operator of $T$ (as an operator on $\left.L^{2}(M)\right)$. Then

$$
\begin{equation*}
\int_{M} g T^{*} f \mathrm{~d} \mu=\int_{M} T g f \mathrm{~d} \mu . \tag{4.2}
\end{equation*}
$$

Since $g$ is in $H_{\mathrm{fin}}^{1,2}(M)$ and the operators $T$ and $\widetilde{T}$ agree on $H_{\mathrm{fin}}^{1,2}(M)$, we see that

$$
\begin{align*}
\int_{M} T g f \mathrm{~d} \mu & =\int_{M} \widetilde{T} g f \mathrm{~d} \mu  \tag{4.3}\\
& =\left\langle g,(\widetilde{T})^{*} f\right\rangle
\end{align*}
$$

where $(\widetilde{T})^{*}$ denotes the transpose of the operator $\widetilde{T}$ from $H^{1}(M)$ to $L^{1}(M)$. Note that $(\widetilde{T})^{*} f$ is in $B M O(M)$ and $g$ is a multiple of an atom. Thus the above scalar product $\left\langle g,(\widetilde{T})^{*} f\right\rangle$ (with respect to the duality between $H^{1}(M)$ and $B M O(M)$ ) may be written as $\int_{M} g(\widetilde{T})^{*} f \mathrm{~d} \mu$. Therefore, (4.2) and (4.3) imply that

$$
\int_{M} g\left[T^{*} f-(\widetilde{T})^{*} f\right] \mathrm{d} \mu=0 \quad \forall g \in L_{c, 0}^{2}(M)
$$

i.e., for all $g$ in $X^{2}$. Therefore $T^{*} f-(\widetilde{T})^{*} f=0$ in the dual space of $X^{2}$, i.e., in $L_{\text {loc }}^{2}(M) / \mathbb{C}$. This implies that $T^{*} f-(\widetilde{T})^{*} f$ is constant.

Now, suppose that $g$ is in $H^{1}(M) \cap L^{2}(M)$ and that $f$ is in $L^{2}(M) \cap L^{\infty}(M)$. Then

$$
\begin{align*}
\int_{M} T g f \mathrm{~d} \mu & =\int_{M} g T^{*} f \mathrm{~d} \mu \\
& =\int_{M} g(\widetilde{T})^{*} f \mathrm{~d} \mu  \tag{4.4}\\
& =\int_{M} \widetilde{T} g f \mathrm{~d} \mu .
\end{align*}
$$

Since $f$ is an arbitrary function in $L^{2}(M) \cap L^{\infty}(M), T g-\widetilde{T} g=0$ almost everywhere, as required.

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