

## ON THE QUENCHING RATE ESTIMATE

By

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**§1. Introduction.** In this paper, we consider the problem

$$\begin{aligned} u_t - u_{xx} &= (1 - u)^{-\beta}, & |x| < l, t > 0, \\ u(\pm l, t) &= 0, & t > 0, \\ u(x, 0) &= u_0(x), & |x| \leq l, \end{aligned} \tag{1.1}$$

where  $\beta > 0$ ,  $l > 0$ ,  $0 \leq u_0 < 1$  is smooth. The solution  $u$  of (1.1) is said to be quenching if  $u$  reaches 1 in finite time  $T$ . Note that in this case  $u_t$  blows up at the same time  $T$ . This phenomenon has been studied by many authors (see, for example, the references cited in [8] and [11]). In particular, for any  $\beta > 0$  there exists a positive constant  $l_* = l_*(\beta)$  such that  $u$  quenches for any  $u_0$  if  $l > l_*$ . Hereafter we shall assume that  $u$  quenches and that  $u_0$  satisfies

$$u_0'' + (1 - u_0)^{-\beta} \geq 0. \tag{1.2}$$

Let  $T \in (0, \infty)$  be the quenching time for  $u$ . We say that  $a$  is a quenching point for  $u$  if there exists a sequence  $\{(x_n, t_n)\}$  such that  $x_n \rightarrow a$ ,  $t_n \uparrow T$ , and  $u(x_n, t_n) \rightarrow 1$  as  $n \rightarrow \infty$ . The set of all such points (for the same  $T$ ) is called the quenching set.

In [8] we first proved that the quenching set consists of finite points which remain a positive distance from  $x = \pm l$ . Then we studied the quenching rate of the solution near any quenching point. Let  $\gamma = (\beta + 1)^{-1}$  and  $k = \gamma^{-\gamma}$ . We obtained the following quenching rate estimate.

**THEOREM A.** If  $a$  is a quenching point, then

$$\lim_{t \uparrow T} (1 - u(x, t))(T - t)^{-\gamma} = k \tag{1.3}$$

uniformly for  $|x - a| \leq C\sqrt{T - t}$  for any positive constant  $C$ .

But, there we only proved this theorem for  $\beta \geq 3$ . In [11], Levine conjectured that this theorem should hold for all  $\beta > 0$ . Recently, Fila and Hulshof [3] improved this result to any  $\beta \geq 1$  using a convexity argument of [5]. The purpose of this paper is to complete this quenching rate estimate for any  $\beta > 0$ .

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Recall the similarity variables

$$y = (x - a)/\sqrt{T - t}, \quad s = -\ln(T - t),$$

$$w(y, s) = (1 - u(x, t))(T - t)^{-\gamma}.$$

Then  $w$  satisfies the equation

$$w_s = w_{yy} - \frac{1}{2}yw_y + \gamma w - w^{-\beta} \tag{1.4}$$

in the set  $W = \{(y, s) : |y \exp(-s/2) + a| < l, s > -\ln T\}$ . The key step for the proof of Theorem A is to show that the only positive global solution to

$$w'' - \frac{1}{2}yw' + \gamma w - w^{-\beta} = 0, \quad y \in \mathbf{R}, \tag{1.5}$$

under some growth condition at  $|y| = \infty$ , is the constant solution  $w \equiv k$ .

We note that the quenching rate estimate for the corresponding higher-dimensional radial case for  $\beta > 1$  was obtained by the author in [9]. Recently, Fila, Hulshof, and Quittner [4] have completed this result for all  $\beta > 0$ .

This paper is organized as follows. In Sec. 2, we give the asymptotic behavior of solutions  $w$  of (1.5) at  $|y| = \infty$  by using the method of [9](see also [2]). Then we apply this result to obtain the quenching rate estimate (1.3) for  $0 < \beta < 1$  in Sec. 3.

**§2. Asymptotic behaviors.** In this section, we let  $w = w(y)$  be any positive global solution of (1.5). The main result of this section is as follows.

**THEOREM 2.1.** If  $w(y)$  is not identically equal to  $k$ , then  $w(y)$  behaves either as  $|y|^{2\gamma}$  or as  $|y|^{-(1+2\gamma)} \exp(y^2/4)$  at  $|y| = \infty$ .

Let

$$f(w) = \gamma w - w^{-\beta} \quad \text{and} \quad F(w) = \int_k^w f(s)ds, \quad w > 0.$$

Note that  $f(w) > 0$  if  $w > k$ ,  $< 0$  if  $w < k$ ; and  $F(w) > 0$  for  $w \neq k$ . Rewrite (1.5) as

$$[w'^2/2 + F(w)]'(y) = yw'^2(y)/2. \tag{2.1}$$

Since the right-hand side of (2.1) is nonnegative for  $y \geq 0$ , the limit

$$\lim_{y \rightarrow \infty} [w'^2/2 + F(w)](y) = l_+$$

exists and is nonnegative. We claim that  $l_+ > 0$  if  $w \neq k$ . Indeed, if  $l_+ = 0$ , then  $w'(y) \rightarrow 0$  and  $w(y) \rightarrow k$  as  $y \rightarrow \infty$ . By integrating (2.1) from  $y$  to  $\infty$ , we obtain that  $w'(y) \equiv 0$  and  $w \equiv k$  for  $y \geq 0$ . This contradiction leads to the conclusion  $l_+ > 0$ . Let  $v = w'$ . Then  $v' = (y/2)v - f(w)$ .

**LEMMA 2.2.** Let  $\alpha$  be a positive constant. Then the region

$$A_\alpha \equiv \{(w, v) : w \geq k, v \geq \alpha w\}$$

is a positively invariant region, i.e., there exists  $y_0 = y_0(\alpha)$  such that if  $y_1 \geq y_0$  and  $(w(y_1), v(y_1)) \in A_\alpha$  then  $(w(y), v(y)) \in A_\alpha, \forall y \geq y_1$ .

*Proof.* Take  $y_0 = 2(\alpha + \gamma/\alpha)$ . Since the vector field  $(w'(y), v'(y))$  is always pointing inward to the region  $A_\alpha$  for  $y \geq y_0$ , the lemma follows.  $\square$

**LEMMA 2.3.** Suppose that  $w(y) \not\equiv k$ . Then  $w$  cannot assume the value  $k$  at infinitely many points as  $y \rightarrow \infty$ .

*Proof.* Suppose that there is a sequence  $y_m \rightarrow \infty$  as  $m \rightarrow \infty$  such that  $w(y_m) = k, \forall m$ . Without loss of generality we may assume that  $w'(y_m) > 0$  for all  $m$ . Recall that  $l_+ > 0$ . Hence

$$w'(y_m) \rightarrow \sqrt{2l_+} \text{ as } m \rightarrow \infty.$$

Take any number  $\alpha \in (0, \sqrt{2l_+})$  and consider the positively invariant set  $A = A_{\alpha/k}$ . Let  $y_0 = y_0(\alpha/k)$  be the constant obtained in Lemma 2.2. Then there is an  $m_0$  sufficiently large such that  $w'(y_{m_0}) \geq \alpha$  and  $y_{m_0} \geq y_0$ . Since  $(w(y_{m_0}), v(y_{m_0})) \in A$ , from Lemma 2.2 it follows that  $(w(y), v(y)) \in A$  for all  $y \geq y_{m_0}$ . Thus  $v(y) > 0, \forall y \geq y_{m_0}$ , which is a contradiction and the lemma is proved.  $\square$

Let  $w(y)$  be a nonconstant positive global solution of (1.5). From Lemma 2.3 it follows that either  $w(y) > k, \forall y \geq \bar{y}$ , or  $w(y) < k, \forall y \geq \bar{y}$ , for some  $\bar{y} > 0$ . From the differential equation (1.5), we observe that any critical point  $y$  of  $w$  is a local maximum point if  $w(y) > k$ , and is a local minimum point if  $w(y) < k$ . Moreover, by the local uniqueness of solutions of ordinary differential equations, there cannot exist a point  $y$  with  $w(y) = k$  and  $w'(y) = 0$  except when  $w \equiv k$ . Therefore, there is  $y_0 \geq \bar{y}$  such that either

$$w'(y) > 0, \quad \forall y \geq y_0, \quad \text{or} \quad w'(y) < 0, \quad \forall y \geq y_0.$$

Let

$$L = \lim_{y \rightarrow \infty} w(y).$$

We claim that  $L > k$ , if  $w'(y) > 0$  for all  $y \geq y_0$ . Indeed, if  $L \leq k$ , then from the equation

$$w'' = yw'/2 - f(w)$$

we obtain that  $w''(y) > 0$  for  $y \geq y_0$ . Thus

$$w(y) = w(y_0) + \int_{y_0}^y w'(\xi)d\xi \geq w'(y_0)(y - y_0) \rightarrow \infty$$

as  $y \rightarrow \infty$ , a contradiction. Hence  $L > k$ , if  $w' > 0, \forall y \geq y_0$ . Similarly,  $L < k$ , if  $w'(y) < 0$  for all  $y \geq y_0$ .

**LEMMA 2.4.** If  $w \not\equiv k$ , then  $w$  must be strictly increasing to  $+\infty$  as  $y \rightarrow \infty$ .

*Proof.* Suppose that  $L < \infty$ . Since the integral

$$\int_{y_0}^{\infty} w'(y)dy = L - w(y_0)$$

is finite, there is a sequence  $y_m \rightarrow \infty$  such that  $w'(y_m) \rightarrow 0$  as  $m \rightarrow \infty$ . Rewrite (1.5) as

$$w''/y - w'/2 = -f(w)/y$$

and integrate it from  $y_0$  to  $y_m$ . Using an integration by parts for the first term, the integral on the left-hand side remains bounded as  $m \rightarrow \infty$ . But,

$$\left| \int_{y_0}^{y_m} \frac{f(w(y))}{y} dy \right| \rightarrow \infty$$

as  $m \rightarrow \infty$ . This contradiction leads to the conclusion of the lemma.  $\square$

*Proof of Theorem 2.1.* From Lemma 2.4 we can easily obtain the asymptotic behavior of  $w(y)$  at  $y = \infty$  by the method used in [9]. The motivation is from Remark 2 in [2] and the proof is based on using L'Hôpital's rule (cf. [2] and [9]). For the reader's convenience, we outline the proof here.

First, using Lemma 2.2 and 2.4, we can show that the limit

$$\alpha \equiv \lim_{y \rightarrow \infty} \frac{w'(y)}{w(y)}$$

exists and that  $\alpha$  is either 0 or  $\infty$ . Suppose that  $\alpha = 0$ . Then we have

$$\lim_{y \rightarrow \infty} \frac{yw'(y)}{w(y)} = 2\gamma,$$

by applying L'Hôpital's rule and using the formula

$$\frac{w'(y)}{w(y)} = \exp(y^2/4) \int_y^\infty [\gamma + a(s)] \exp(-s^2/4) ds,$$

where  $a(s) \rightarrow 0$  as  $s \rightarrow \infty$ . Thus, from

$$\lim_{y \rightarrow \infty} y^\delta \left[ \frac{yw'(y)}{w(y)} - 2\gamma \right] = 0$$

for  $\delta \in (0, 2)$ , we conclude that for any  $\delta \in (0, 2)$  there exists a positive constant  $C_\delta$  such that

$$w(y) = C_\delta y^{2\gamma} [1 + o(y^{-\delta})]$$

as  $y \rightarrow \infty$ . The case for  $\alpha = \infty$  is similar.

The asymptotic behavior of  $w(y)$  at  $y = -\infty$  follows by a similar argument and the theorem is proved.  $\square$

In the sequel, we shall call a nonconstant positive global solution of (1.5), which behaves as  $|y|^{2\gamma}$  both at  $y = \infty$  and at  $y = -\infty$ , as a *slow orbit*.

**REMARK 2.5.** For any  $\beta > 0$ , using the Sturm comparison theorem (cf. [10]), we can show that every nonconstant positive solution of (1.5) must be strictly convex for all  $y$  sufficiently large and/or for all  $-y$  sufficiently large. Hence there is no slow orbit for  $\beta \geq 1$ .

The proof of Remark 2.5 is quite similar to that of [10, Theorem 2]. Here we compare the function  $w''$  with the function  $w^{-\beta} - w''$  and compare the function  $w^{-\beta} - w''$  with  $w'$ . Then, for a nonconstant positive solution  $w(y)$  of (1.5), if  $y_0 \geq 0$  (which we may assume without loss of generality) is such that  $w'(y_0) = 0$  and  $w'(y) > 0$  for all  $y > y_0$ , we have  $w''(y) > 0$  for all  $y \geq y_0$ .

For  $0 < \beta < 1$ , let

$$\phi(y) = K|y|^{2\gamma}, \quad K = [2\gamma(2\gamma - 1)]^{-\gamma}.$$

Note that  $K$  is positive, since  $\gamma > \frac{1}{2}$  in this case. Also,  $\phi(y)$  is a continuously differentiable function satisfying (1.5) for all  $y \neq 0$ . Similar to [9, Theorem 3.4], we have the following result.

**THEOREM 2.6.** Any slow orbit  $w(y)$  must intersect the function  $\phi(y)$  at least twice in  $y > 0$  and/or in  $y < 0$ .

*Proof.* Since the proof is quite similar to that of [9, Theorem 3.4], we only sketch the proof. To begin with, we take the minimum point  $y_0$  of  $w$ . Without loss of generality we may assume that  $y_0 \geq 0$ . First, let

$$g(y) = \gamma w(y) - \frac{1}{2}y w'(y), \quad h(y) = w''(y), \quad \text{and} \quad V(y) = g(y)h'(y) - g'(y)h(y).$$

Then

$$V(y) = \rho^{-1}(y)\{\rho(y_0)V(y_0) + \int_{y_0}^y \rho(t)\beta(\beta + 1)w^{-(\beta+2)}(t)[w'(t)]^2 g(t) dt\}$$

where  $\rho(y) = \exp(-y^2/4)$ , and

$$h(y) = \frac{h(y_0)}{g(y_0)}g(y) + g(y) \int_{y_0}^y \frac{V(t)}{g^2(t)} dt.$$

Recall that  $g(y) \rightarrow 0$  and  $g'(y) \rightarrow 0$  as  $y \rightarrow \infty$  for any slow orbit  $w(y)$ . By the choice of  $y_0$ , we have  $V(y_0) \geq 0$  and  $g(y_0) > 0$ . If  $g(y) > 0$  for all  $y \geq y_0$ , then we will have  $V(y) \rightarrow \infty$  as  $y \rightarrow \infty$ . Hence  $h(y) \rightarrow \infty$  as  $y \rightarrow \infty$ , a contradiction. Therefore,  $g(y)$  must have a zero in  $y > y_0$ .

Next, let

$$U(y) = w(y)\phi'(y) - w'(y)\phi(y).$$

Then  $U(0) = 0$  and  $U$  satisfies

$$\begin{aligned} U(y) &= 2Ky^{2\gamma-1}g(y), \\ U' - \frac{y}{2}U &= w\phi[\phi^{-(\beta+1)} - w^{-(\beta+1)}], \quad y \neq 0. \end{aligned}$$

If  $w(y_0) < \phi(y_0)$ , then clearly  $w$  intersects  $\phi$  at least once in  $(0, y_0)$ . From here and proceeding as in the proof of Lemma 3.6 of [9], we obtain that  $w$  intersects  $\phi$  at least twice in  $y > 0$ . Hence the theorem follows.  $\square$

**§3. The quenching rate.** We assume that  $\beta \in (0, 1)$  and hence  $2\gamma > 1$ . Let  $a$  be a quenching point for  $u$ . Without loss of generality we may assume that  $a = 0$ . First, we recall that  $w(y, s) \geq 1/B$  in  $W$  for some positive constant  $B$ . Applying the maximum principle (cf. [6]) to the function

$$J(x, t) = \frac{1}{2}u_x^2 - \frac{C}{1-\beta}(1-u)^{1-\beta}$$

for some constant  $C \geq 1$ , we obtain that

$$|u_x| \leq \sqrt{\frac{C}{1-\beta}}(1-u)^{(1-\beta)/2} \quad \text{in } (-l, l) \times (0, T). \tag{3.1}$$

From (3.1) it follows that

$$w(y, s) \leq C(|y|^{2\gamma} + 1) \quad \text{in } W. \quad (3.2)$$

Now, applying the energy method of [7] (for details see [8]), we can show that  $w(y, s)$  tends to a positive global solution of (1.5) as  $s \rightarrow \infty$ . By (3.2), this limit function must be a slow orbit, if it is not identically equal to the constant  $k$ .

We claim that this limit function intersects  $\phi(y)$  at most once both in  $y > 0$  and in  $y < 0$ . First, consider the case  $y > 0$ . If  $w(y, s) > \phi(y)$  in  $W \cap \{y > 0\}$ , then we are done. Otherwise, we choose  $s_0 < \infty$  such that  $w(y_0, s_0) \leq \phi(y_0)$  for some  $y_0 > 0$  with  $(y_0, s_0) \in W$ . Then proceeding as in the proofs of [1, Sec. 3] there is a  $\delta > 0$  such that  $w(y, s)$  intersects  $\phi(y)$  exactly once in  $(0, \delta \exp(s/2))$  for all  $s > s_0$ . The case for  $y < 0$  is similar. Therefore, the assertion follows. By Theorem 2.6, this limit function must be identically equal to the constant  $k$ . Hence Theorem A is proved for  $0 < \beta < 1$ .

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