On the spectrum of the lattice spin-boson Hamiltonian for any coupling: 1D case

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March 3, 2022

Keywords: spin-boson Hamiltonian, block operator matrix, bosonic Fock space, annihilation and creation operators, Birman-Schwinger principle, essential, point and discrete spectrum

Subject classification: 81Q10, 35P20, 47N50

Abstract: A lattice model of radiative decay (so-called spin-boson model) of a two level atom and at most two photons is considered. The location of the essential spectrum is described. For any coupling constant the finiteness of the number of eigenvalues below the bottom of its essential spectrum is proved. The results are obtained by considering a more general model H for which the lower bound of its essential spectrum is estimated. Conditions which guarantee the finiteness of the number of eigenvalues of H, below the bottom of its essential spectrum are found. It is shown that the discrete spectrum might be infinite if the parameter functions are chosen in a special form.

1 Introduction

Block operator matrices are matrices where the entries are linear operators between Banach or Hilbert spaces [16]. One special class of block operator matrices are Hamiltonians associated with systems of non-conserved number of quasi-particles on a lattice. Their number can be unbounded as in the case of spin-boson models or bounded as in the case of "truncated" spin-boson models. They arise, for example, in the theory of solid-state physics [11], quantum field theory [3] and statistical physics [9, 10].

In a well-known model of radiative decay (the so-called spin-boson model) it is assumed that an atom, which can be in two states – ground state with energy $-\varepsilon$ and excited state with energy ε – emits and absorbs photons, going over from one state to the other [6, 10, 15, 17]. The energy operator of such a system is given by the (formal) expression [6, 10, 15, 17]

$$\mathcal{A} := \varepsilon \sigma_z + \int_{\mathbb{R}^d} w(k) a^*(k) a(k) dk + \alpha \sigma_x \int_{\mathbb{R}^d} v(k) (a^*(k) + a(k)) dk \tag{1.1}$$

and acts in the Hilbert space

$$\mathcal{L} := \mathbb{C}^2 \otimes \mathcal{F}_{s}(L_2(\mathbb{R}^d)), \tag{1.2}$$

where \mathbb{C}^2 is the state of the two-level atom and $\mathcal{F}_s(L_2(\mathbb{R}^d))$ is the symmetric Fock space for bosons. In the following we consider the lattice analog of the standard spin-boson Hamiltonian cf. [11]. In the "algebraic" sense, a lattice spin-boson Hamiltonian is similar to a standard one with only the difference is that \mathcal{A} does not act in the Euclidean space \mathbb{R}^d but on a d-dimensional torus \mathbb{T}^d . This means that we have to replace \mathbb{R}^d by \mathbb{T}^d in formulas (1.1) and (1.2). We write elements F of the space \mathcal{L} in the form

$$F = \{ f_0^{(\sigma)}, f_1^{(\sigma)}(k_1), f_2^{(\sigma)}(k_1, k_2), \dots, f_n^{(\sigma)}(k_1, k_2, \dots, k_n), \dots \}$$

of functions of an increasing number of variables (k_1, \ldots, k_n) , $k_i \in \mathbb{T}^d$, and a discrete variable $\sigma = \pm$; the functions are symmetric with respect to the variables k_i , $i = 1, \ldots, n$, $n \in \mathbb{N}$. The norm in \mathcal{L} is given by

$$||F||^2 := \sum_{\sigma=\pm} |f_0^{(\sigma)}|^2 + \sum_{\sigma,n} \int_{(\mathbb{T}^d)^n} |f_n^{(\sigma)}(k_1,\dots,k_n)|^2 dk_1 \dots dk_n.$$
 (1.3)

In the expression (1.1), the operators $a^*(k)$ and a(k) are "creation and annihilation" operators, $\varepsilon > 0$,

$$\sigma_z := \left(\begin{array}{cc} 1 & 0 \\ 0 & -1 \end{array} \right), \quad \sigma_x := \left(\begin{array}{cc} 0 & 1 \\ 1 & 0 \end{array} \right)$$

are Pauli matrices, w(k) is the dispersion of the free field, $\alpha v(k)$ is the coupling between the atoms and the field modes, $\alpha > 0$ is the coupling constant.

However, the problem of complete spectral description of the operator \mathcal{A} still seems rather difficult. In this connection, it is natural to consider simplified ("truncated") models [6, 10, 17] that differ from the model described above with respect to the number of bosons which is bounded by $N, N \in \mathbb{N}$. The Hilbert state space of each such model is then the space $\mathcal{L}_N := \mathbb{C}^2 \otimes \mathcal{F}_s^{(N)}(L_2(\mathbb{T}^d)) \subset \mathcal{L}$, where

$$\mathcal{F}_{s}^{(N)}(L_{2}(\mathbb{T})) := \mathbb{C} \oplus L_{2}(\mathbb{T}^{d}) \oplus L_{2}^{sym}((\mathbb{T}^{d})^{2}) \oplus \ldots \oplus L_{2}^{sym}((\mathbb{T}^{d})^{N}).$$

Here $L_2^{\text{sym}}((\mathbb{T}^d)^n)$ is the Hilbert space of symmetric functions of n variables, and the norm in \mathcal{L}_N is introduced as in (1.3). Then the truncated Hamiltonian \mathcal{A}_N is given in \mathcal{L}_N by $\mathcal{A}_N := P_{\mathcal{L}_N} \mathcal{A} P_{\mathcal{L}_N}$, where $P_{\mathcal{L}_N}$ is the projection of the space \mathcal{L} onto the subspace \mathcal{L}_N , and \mathcal{A} is the Hamiltonian (1.1).

The standard spin-boson Hamiltonian with N=1,2 was completely studied in [10] for small values of the parameter α . The case N=3 was considered in [17]. The existence of wave operators and their asymptotic completeness were proven there. In [6], the case of arbitrary N was investigated. In particular, using a Mourre type estimate, a complete spectral characterization of the spin-boson Hamiltonian are given for sufficiently small, but nonzero coupling constant.

Let us introduce the corresponding model operator for the case N=2. For simplicity we denote $\mathcal{H}_0:=\mathbb{C},\ \mathcal{H}_1:=L_2(\mathbb{T}^d),\ \mathcal{H}_2:=L_2^{\mathrm{sym}}((\mathbb{T}^d)^2)$ and $\mathcal{H}:=\mathcal{F}_{\mathrm{s}}^{(2)}(L_2(\mathbb{T}^d))$.

In the Hilbert space \mathcal{H} we consider the model operator H that admit an 3×3 tridiagonal block operator matrix representation

$$H := \left(\begin{array}{ccc} H_{00} & H_{01} & 0 \\ H_{01}^* & H_{11} & H_{12} \\ 0 & H_{12}^* & H_{22} \end{array} \right)$$

with the entries $H_{ij}: \mathcal{H}_j \to \mathcal{H}_i, i \leq j, i, j = 0, 1, 2$ defined by

$$H_{00}f_0 = w_0 f_0, \quad H_{01}f_1 = \int_{\mathbb{T}^d} v_0(t) f_1(t) dt,$$

$$(H_{11}f_1)(x) = w_1(x) f_1(x), \quad (H_{12}f_2)(x) = \int_{\mathbb{T}^d} v_1(t) f_2(x,t) dt,$$

$$(H_{22}f_2)(x,y) = w_2(x,y) f_2(x,y), \quad f_i \in \mathcal{H}_i, \quad i = 0, 1, 2;$$

where w_0 is a real number, $w_1(\cdot)$ $v_i(\cdot)$, i = 0, 1, are real-valued analytic functions on \mathbb{T}^d and $w_2(\cdot, \cdot)$ is a real-valued symmetric analytic function on $(\mathbb{T}^d)^2$. Under these assumptions the operator H is bounded and self-adjoint.

An important problem in the spectral theory of such model operators is to study the number of eigenvalues located outside the essential spectrum. We remark that the operator H has been considered before and the following results were obtained: The location of the essential spectrum of H has been described in [7] for any $d \ge 1$. The existence of infinitely many eigenvalues below the bottom of the essential spectrum of H has been announced in [8] for d = 3. Its complete proof was given in [2] for a special parameter functions and in [1] for more general case. An asymptotics of the form $\mathcal{U}_0|\log |\lambda||$ ($0 < \mathcal{U}_0 < \infty$) for the number of eigenvalues on the left of λ , $\lambda < \min \sigma_{\text{ess}}(H)$ was obtained in [1]. The conditions for the finiteness of the discrete spectrum of H was found in [13] for the case d = 3.

In the present paper we consider the case d = 1. We study the relation between the lower bounds of the two-particle and three-particle branches of the essential spectrum of H. Under natural assumptions on the parameters we prove the finiteness of the discrete spectrum of H using the Birman-Schwinger principle. In this analysis the finiteness of the points which gives global minima for the function $w_2(\cdot,\cdot)$ is important. We give a counter-example, which shows the infiniteness of the discrete spectrum if the number of such points is not finite. For this case the exact view of the eigenvalues and eigenvectors are found, their multiplicities are calculated. We notice that a part of the results is typical for d = 1, in fact, they do not have analogues in the case $d \ge 2$.

Using a connection between the operators A_2 and H, and applying obtained results from above we describe the essential spectrum of A_2 , and show that the operator A_2 has finitely many eigenvalues below the bottom of its essential spectrum for any coupling constant α .

The paper is organized as follows. Section 1 is an introduction. In Section 2, the main results for H are formulated. In Section 3, we estimate the lower bound of the essential spectrum of H. In Section 4, we apply the Birman-Schwinger principle to H. Section 5 is devoted to the proof of the finiteness of the discrete spectrum of H. In Section 6 we discuss the case when the discrete spectrum of H is infinite. In Section 7, the finiteness and infiniteness of the discrete spectrum of H is established, when the lower bounds of the two- and three-particle branches of the essential spectrum are coincide. An application to lattice model of radiative decay of a two level atom and at most two photons (truncated spin-boson model on a lattice) illustrates our results.

2 The main results for H

The spectrum, the essential spectrum, the point spectrum and the discrete spectrum of a bounded self-adjoint operator will be denoted by $\sigma(\cdot)$, $\sigma_{\rm ess}(\cdot)$, $\sigma_{\rm p}(\cdot)$ and $\sigma_{\rm disc}(\cdot)$, respectively.

To study the spectral properties of H we introduce a following family of bounded self-adjoint operators (generalized Friedrichs models) h(x), $x \in \mathbb{T}$, which acts in $\mathcal{H}_0 \oplus \mathcal{H}_1$ as

$$h(x) := \begin{pmatrix} h_{00}(x) & h_{01} \\ h_{01}^* & h_{11}(x) \end{pmatrix},$$

where

$$h_{00}(x)f_0 = w_1(x)f_0, \quad h_{01}f_1 = \frac{1}{\sqrt{2}} \int_{\mathbb{T}} v_1(t)f_1(t)dt,$$

 $(h_{11}(x)f_1)(y) = w_2(x,y)f_1(y), \quad f_i \in \mathcal{H}_i, \quad i = 0, 1.$

Let the operator $h_0(x)$, $x \in \mathbb{T}$ act in $\mathcal{H}_0 \oplus \mathcal{H}_1$ as

$$h_0(x) := \left(\begin{array}{cc} 0 & 0 \\ 0 & h_{11}(x) \end{array} \right).$$

The perturbation $h(x) - h_0(x)$ of the operator $h_0(x)$ is a self-adjoint operator of rank 2. Therefore in accordance with the Weyl theorem about the invariance of the essential spectrum under the finite rank perturbations, the essential spectrum of the operator h(x) coincides with the essential spectrum of $h_0(x)$. It is evident that $\sigma_{\text{ess}}(h_0(x)) = [m_x, M_x]$, where the numbers m_x and M_x are defined by

$$m_x := \min_{y \in \mathbb{T}} w_2(x, y)$$
 and $M_x := \max_{y \in \mathbb{T}} w_2(x, y)$.

This yields $\sigma_{\rm ess}(h(x)) = [m_x, M_x].$

For any $x \in \mathbb{T}$ we define an analytic function $\Delta(x;\cdot)$ (the Fredholm determinant associated with the operator h(x)) in $\mathbb{C} \setminus [m_x, M_x]$ by

$$\Delta(x;z) := w_1(x) - z - \frac{1}{2} \int_{\mathbb{T}} \frac{v_1^2(t)dt}{w_2(x,t) - z}.$$

Note that for the discrete spectrum of h(x) the equality

$$\sigma_{\rm disc}(h(x)) = \{ z \in \mathbb{C} \setminus [m_x, M_x] : \Delta(x; z) = 0 \}$$

holds (see Lemma 3.2).

The following theorem [7] describes the location of the essential spectrum of H by the spectrum of the family h(x) of generalized Friedrichs model.

Theorem 2.1 For the essential spectrum of H the following equality holds

$$\sigma_{\mathrm{ess}}(H) = \sigma \cup [m, M], \quad \sigma := \bigcup_{x \in \mathbb{T}} \sigma_{\mathrm{disc}}(h(x))$$

where the numbers m and M are defined by

$$m := \min_{x,y \in \mathbb{T}} w_2(x,y)$$
 and $M := \max_{x,y \in \mathbb{T}} w_2(x,y)$.

The sets σ and [m, M] are called two- and three-particle branches of the essential spectrum of H, respectively.

Throughout this paper we assume that the function $w_2(\cdot,\cdot)$ has a unique non-degenerate global minimum at the point $(0,0) \in \mathbb{T}^2$.

For $\delta > 0$ and $a \in \mathbb{T}$ we set

$$U_{\delta}(a) := \{ x \in \mathbb{T} : |x - a| < \delta \}.$$

We remark that if $v_1(0) = 0$, then from analyticity of $v_1(\cdot)$ on \mathbb{T} it follows that there exist positive numbers C_1, C_2 and δ such that the inequalities

$$C_1|x|^{\alpha} \le |v_1(x)| \le C_2|x|^{\alpha}, \quad x \in U_{\delta}(0)$$
 (2.1)

hold for some $\alpha \in \mathbb{N}$. Since the function $w_2(0,\cdot)$ has a unique non-degenerate global minimum at y=0 (see proof of Lemma 3.3), one can easily seen from the estimate (2.1) that for any $x\in\mathbb{T}$ the integral

$$\int_{\mathbb{T}} \frac{v_1^2(t)dt}{w_2(x,t) - m}$$

is positive and finite. The Lebesgue dominated convergence theorem yields

$$\Delta(0; m) = \lim_{x \to 0} \Delta(x; m),$$

and hence if $v_1(0) = 0$, then the function $\Delta(\cdot; m)$ is continuous on \mathbb{T} .

Let us denote by E_{\min} the lower bound of the essential spectrum of H.

The main results of the present paper as follows.

Theorem 2.2 For the lower bound E_{\min} the following assertions hold:

- (i) If $v_1(0) \neq 0$, then $E_{\min} < m$;
- (ii) If $v_1(0) = 0$ and $\min_{x \in \mathbb{T}} \Delta(x; m) < 0$, then $E_{\min} < m$; (iii) If $v_1(0) = 0$ and $\min_{x \in \mathbb{T}} \Delta(x; m) \ge 0$, then $E_{\min} = m$.

Theorem 2.3 If one of the assertions

- (i) $v_1(0) \neq 0$;
- (ii) $v_1(0) = 0$ and $\min_{x \in \mathbb{R}} \Delta(x; m) < 0$;
- (iii) $v_1(0) = 0$ and $\min_{x \in \mathbb{T}} \Delta(x; m) > 0$,

is satisfied, then the operator H has a finite number of eigenvalues lying below E_{\min} .

Remark 2.4 Since the function $w_2(\cdot,\cdot)$ is continuous on the compact set \mathbb{T}^2 there exist at least one point $(x_0,y_0)\in\mathbb{T}^2$ such that the function $w_2(\cdot,\cdot)$ attains its global maximum at this point. If $v_1(y_0)\neq 0$, then similar arguments show that for the upper bound E_{\max} of the essential spectrum of H we have $E_{\max}>M$ and the operator H has a finite number of eigenvalues greater than E_{\max} .

Remark 2.5 The results can be easily generalized to the case when the function $w_2(\cdot, \cdot)$ has a finite number of non-degenerate global minima at several points of \mathbb{T}^2 . Here finiteness of the number of such points is important. If the number of such points is infinite, then the discrete spectrum of H can be infinite, for a corresponding example see Section 6.

Remark 2.6 The case $v_1(0) = 0$ and $\min_{x \in \mathbb{T}} \Delta(x; m) = 0$ is considered in Section 7, where it is shown that the discrete spectrum of H can be finite or infinite depending on the parameter functions.

3 Lower bound of the essential spectrum of H

In this section first we study the discrete spectrum of h(x) and then Theorem 2.2 will be proven.

Proposition 3.1 The perturbation determinant $\Delta_{h(x)/h_0(x)}(z)$ of the operator $h_0(x)$ by the operator $h(x) - h_0(x)$ has form

$$\Delta_{h(x)/h_0(x)}(z) = -\frac{1}{z}\Delta(x;z), \quad z \in \mathbb{C} \setminus \sigma(h_0(x)).$$

Proof. Since the operator $h(x) - h_0(x)$ is trace class, even of rank 2, the perturbation determinant $\Delta_{h(x)/h_0(x)}(z)$ is well-defined by

$$\Delta_{h(x)/h_0(x)}(z) := \det \left(I + (h(x) - h_0(x))(h_0(x) - z)^{-1} \right).$$

Without loss of generality we can assume that $||v_1|| = 1$. We choose the orthonormal basis $\{\varphi_n\}_n \subset \mathcal{H}_1$ by the following way: $\varphi_1 := v_1$ and $\varphi_j \perp v_1$ for all $j \geq 2$. Introduce

$$\psi_1 := \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ \varphi_1 \end{pmatrix}, \quad \psi_2 := \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ -\varphi_1 \end{pmatrix}, \quad \psi_j := \begin{pmatrix} 0 \\ \varphi_{j-1} \end{pmatrix}, \quad j \ge 3.$$

By the construction $\{\psi_n\}_n \subset \mathcal{H}_0 \oplus \mathcal{H}_1$ is an orthonormal. Set

$$a_{ij}(x;z) := ((h(x) - h_0(x))(h_0(x) - z)^{-1}\psi_i, \psi_j), \quad i, j \in \mathbb{N}.$$

Simple calculation show that

$$a_{11}(x;z) = -\frac{1}{z}w_1(x) + \frac{1}{\sqrt{2}} \int_{\mathbb{T}} \frac{v_1^2(t)dt}{w_2(x,t) - z} - \frac{1}{z} \frac{1}{\sqrt{2}};$$

$$a_{12}(x;z) = -\frac{1}{z}w_1(x) + \frac{1}{\sqrt{2}} \int_{\mathbb{T}} \frac{v_1^2(t)dt}{w_2(x,t) - z} + \frac{1}{z} \frac{1}{\sqrt{2}};$$

$$a_{21}(x;z) = -\frac{1}{z}w_1(x) - \frac{1}{\sqrt{2}} \int_{\mathbb{T}} \frac{v_1^2(t)dt}{w_2(x,t) - z} - \frac{1}{z} \frac{1}{\sqrt{2}};$$

$$a_{22}(x;z) = -\frac{1}{z}w_1(x) - \frac{1}{\sqrt{2}} \int_{\mathbb{T}} \frac{v_1^2(t)dt}{w_2(x,t) - z} + \frac{1}{z} \frac{1}{\sqrt{2}};$$

$$a_{ij}(x;z) = \delta_{ij}, \quad \text{otherwise.}$$

Here δ_{ij} is Kronecker delta. Therefore,

$$\Delta_{h(x)/h_0(x)}(z) = \frac{1}{4} \det \begin{pmatrix} 2 - a_{11}(x;z) & a_{12}(x;z) \\ a_{21}(x;z) & 2 - a_{22}(x;z) \end{pmatrix} = -\frac{1}{z} \Delta(x;z).$$

Proposition is proved.

The following lemma is a simple consequence of the Proposition 3.1 and of [5, chapter IV].

Lemma 3.2 For any fixed $x \in \mathbb{T}$ the operator h(x) has an eigenvalue $z(x) \in \mathbb{C} \setminus [m_x, M_x]$ if and only if $\Delta(x; z(x)) = 0$.

In the next two lemmas we describe the number and location of the eigenvalues of h(x).

Lemma 3.3 If $v_1(0) \neq 0$, then there exists $\delta > 0$ such that for any $x \in U_{\delta}(0)$ the operator h(x) has a unique eigenvalue z(x), lying on the left of m_x .

Proof. Since the function $w_2(\cdot, \cdot)$ has a unique non-degenerate global minimum at the point $(0,0) \in \mathbb{T}^2$, by the implicit function theorem there exist $\delta > 0$ and an analytic function $y_0(\cdot)$ on $U_{\delta}(0)$ such that for any $x \in U_{\delta}(0)$ the point $y_0(x)$ is the unique non-degenerate minimum of the function $w_2(x,\cdot)$ and $y_0(0) = 0$. Therefore, we have $w_2(x,y_0(x)) = m_x$ for any $x \in U_{\delta}(0)$.

Let $\widetilde{w}_2(\cdot,\cdot)$ be the function defined on $U_{\delta}(0) \times \mathbb{T}$ as $\widetilde{w}_2(x,y) := w_2(x,y+y_0(x)) - m_x$. Then for any $x \in U_{\delta}(0)$ the function $\widetilde{w}_2(x,\cdot)$ has a unique non-degenerate zero minimum at the point $0 \in \mathbb{T}$. Now using the equality

$$\int_{\mathbb{T}} \frac{v_1^2(t)dt}{w_2(x,t) - m_x} = \int_{\mathbb{T}} \frac{v_1^2(t + y_0(x))dt}{\widetilde{w}_2(x,t)}, \quad x \in U_{\delta}(0),$$

the continuity of the function $v_1(\cdot)$, the conditions $v_1(0) \neq 0$ and $y_0(0) = 0$ it is easy to see that $\lim_{z \to m_x = 0} \Delta(x; z) = -\infty$ for all $x \in U_{\delta}(0)$.

Since for any $x \in \mathbb{T}$ the function $\Delta(x;\cdot)$ is continuous and monotonically decreasing on $(-\infty, m_x)$ the equality

$$\lim_{z \to -\infty} \Delta(x; z) = \infty \tag{3.1}$$

implies that for any $x \in U_{\delta}(0)$ the function $\Delta(x;\cdot)$ has a unique zero z = z(x), lying in $(-\infty, m_x)$. By Lemma 3.2 the number z(x) is the eigenvalue of h(x).

Lemma 3.4 *Let* $v_1(0) = 0$.

- (i) If $\min_{x \in \mathbb{T}} \Delta(x; m) \geq 0$, then for any $x \in \mathbb{T}$ the operator h(x) has no eigenvalues, lying on the left of m.
- (ii) If $\min_{x \in \mathbb{T}} \Delta(x; m) < 0$, then there exists a non-empty set $G \subset \mathbb{T}$ such that for any $x \in G$ the operator h(x) has a unique eigenvalue z(x), lying on the left of m.

Proof. First we recall that if $v_1(0) = 0$, then the function $\Delta(\cdot; m)$ is a continuous on \mathbb{T} . Let $\min_{x \in \mathbb{T}} \Delta(x; m) \geq 0$. Since for any $x \in \mathbb{T}$ the function $\Delta(x; \cdot)$ is monotonically decreasing on $(-\infty, m)$ we have $\Delta(x; z) > \Delta(x; m) \geq \min_{x \in \mathbb{T}} \Delta(x; m) \geq 0$, that is, $\Delta(x; z) > 0$ for all $x \in \mathbb{T}$ and z < m. Therefore, by Lemma 3.2 for any $x \in \mathbb{T}$ the operator h(x) has no eigenvalues in $(-\infty, m)$.

Now we suppose that $\min_{x \in \mathbb{T}} \Delta(x; m) < 0$ and introduce the following subset of \mathbb{T} :

$$G := \{ x \in \mathbb{T} : \Delta(x; m) < 0 \}.$$

Since $\Delta(\cdot\,;m)$ is a continuous on the compact set \mathbb{T} , there exists a point $x^0\in\mathbb{T}$ such that $\min_{x\in\mathbb{T}}\Delta(x\,;m)=\Delta(x^0\,;m)$, that is, $x^0\in G$. So, the set G is a non-empty. Note that if $\max_{x\in\mathbb{T}}\Delta(x\,;m)<0$, then $\Delta(x\,;m)<0$ for all $x\in\mathbb{T}$ and hence $G=\mathbb{T}$.

Since for any $x \in \mathbb{T}$ the function $\Delta(x;\cdot)$ is a continuous and monotonically decreasing on $(-\infty, m]$ by the equality (3.1) for any $x \in G$ there exists a unique point $z(x) \in (-\infty, m)$ such that $\Delta(x; z(x)) = 0$. By Lemma 3.2 for any $x \in G$ the point z(x) is the unique eigenvalue of h(x).

By the construction of G the inequality $\Delta(x;m) \geq 0$ holds for all $x \in \mathbb{T} \setminus G$. In this case one can see that for any $x \in \mathbb{T} \setminus G$ the operator h(x) has no eigenvalues in $(-\infty, m)$.

Proof of Theorem 2.2. Let $v_1(0) \neq 0$. Then by Lemma 3.3 there exists $\delta > 0$ such that for any $x \in U_{\delta}(0)$ the operator h(x) has a unique eigenvalue z(x), lying on the left of m_x . In particular, $z(0) < m_0$. Since $m = \min_{x \in \mathbb{T}} m_x = m_0$ it follows that $\min \sigma \leq z(0) < m$, that is, $E_{\min} < m$.

Let $v_1(0)=0$. Then two cases are possible: $\min_{x\in\mathbb{T}}\Delta(x\,;m)\geq 0$ or $\min_{x\in\mathbb{T}}\Delta(x\,;m)<0$. In the case $\min_{x\in\mathbb{T}}\Delta(x\,;m)\geq 0$, by the part (i) of Lemma 3.4 for any $x\in\mathbb{T}$ the operator h(x) has no eigenvalues in $(-\infty,m)$, that is, $\min\sigma\geq m$. By Theorem 2.1 it means that $E_{\min}=m$.

For the case $\min_{x \in \mathbb{T}} \Delta(x; m) < 0$, using the part (ii) of Lemma 3.4 we obtain $\min \sigma \leq z(x') < m$ for all $x' \in G$, that is, $E_{\min} < m$.

4 The Birman-Schwinger principle.

For a bounded self-adjoint operator A acting in the Hilbert space \mathcal{R} and for a real number λ , we define [4] the number $n(\lambda, A)$ by the rule

$$n(\lambda, A) := \sup \{ \dim(F) : (Au, u) > \lambda, u \in F \subset \mathcal{R}, ||u|| = 1 \}.$$

The number $n(\lambda, A)$ is equal to infinity if $\lambda < \max \sigma_{\rm ess}(A)$; if $n(\lambda, A)$ is finite, then it is equal to the number of the eigenvalues of A bigger than λ .

Let us denote by N(z) the number of eigenvalues of H on the left of $z, z \leq E_{\min}$. Then we have $N(z) = n(-z, -H), \quad -z > -E_{\min}$.

Since the function $\Delta(\cdot;\cdot)$ is positive on $(x,z) \in \mathbb{T} \times (-\infty, E_{\min})$, there exists a positive square root of $\Delta(x;z)$ for all $x \in \mathbb{T}$ and $z < E_{\min}$.

In our analysis of the discrete spectrum of H the crucial role is played by the self-adjoint compact 2×2 block operator matrix T(z), $z < E_{\min}$ acting on $\mathcal{H}_0 \oplus \mathcal{H}_1$ as

$$T(z) := \begin{pmatrix} T_{00}(z) & T_{01}(z) \\ T_{01}^*(z) & T_{11}(z) \end{pmatrix}$$

with the entries $T_{ij}(z): \mathcal{H}_j \to \mathcal{H}_i, i \leq j, i, j = 0, 1$ defined by

$$T_{00}(z)g_0 = (1+z-w_0)g_0, \quad T_{01}(z)g_1 = -\int_{\mathbb{T}} \frac{v_0(t)g_1(t)dt}{\sqrt{\Delta(t;z)}};$$

$$(T_{11}(z)g_1)(x) = \frac{v_1(x)}{2\sqrt{\Delta(x;z)}} \int_{\mathbb{T}} \frac{v_1(t)g_1(t)dt}{\sqrt{\Delta(t;z)}(w_2(x,t)-z)}.$$

Here $g_i \in \mathcal{H}_i$, i = 0, 1.

The following lemma is a modification of the well-known Birman-Schwinger principle for the operator H (see [1]).

Lemma 4.1 The operator T(z) is compact and continuous in $z < E_{\min}$ and

$$N(z) = n(1, T(z)).$$

For the proof of this lemma see Lemma 5.1 of [1].

5 Finiteness of the number of eigenvalues of H

In this section we prove the finiteness of the number of eigenvalues of H, that is, Theorem 2.3. We have divided the proof into a sequence of lemmas.

Lemma 5.1 There exist positive numbers C_1, C_2, C_3 and δ such that the following inequalities hold

(i)
$$C_1(x^2+y^2) \le w_2(x,y) - m \le C_2(x^2+y^2), x,y \in U_\delta(0);$$

(ii)
$$w_2(x,y) - m \ge C_3$$
, $(x,y) \notin U_{\delta}(0) \times U_{\delta}(0)$.

Proof. Since the function $w_2(\cdot, \cdot)$ is analytic on \mathbb{T}^2 and it has a unique non-degenerate global minimum at the point $(0,0) \in \mathbb{T}^2$, the following decomposition holds

$$w_2(x,y) = m + \frac{1}{2} \left(\frac{\partial^2 w_2(0,0)}{\partial x^2} x^2 + 2 \frac{\partial^2 w_2(0,0)}{\partial x \partial y} xy + \frac{\partial^2 w_2(0,0)}{\partial y^2} y^2 \right) + O(|x|^3 + |y|^3)$$

as $x, y \to 0$. Therefore, there exist positive numbers C_1, C_2, C_3 and δ such that (i) and (ii) hold true.

Lemma 5.2 Let the assumption (iii) of Theorem 2.3 be fulfilled. Then there exists a positive number C_1 such that the inequality $\Delta(x;z) \geq C_1$ holds for all $x \in \mathbb{T}$ and $z \leq m$.

Proof. By assumption (iii) of Theorem 2.3 we have $\min_{x \in \mathbb{T}} \Delta(x; m) > 0$. Since for any $x \in \mathbb{T}$ the function $\Delta(x; \cdot)$ is monotonically decreasing in $(-\infty, m]$, we have

$$\Delta(x;z) \ge \Delta(x;m) \ge \min_{x \in \mathbb{T}} \Delta(x;m) > 0$$

for all $x \in \mathbb{T}$ and $z \leq m$. Now setting $C_1 := \min_{x \in \mathbb{T}} \Delta(x; m)$ we complete the proof.

We recall that by Lemma 3.2 the set σ is equal to the set of all complex numbers $z \in \mathbb{C} \setminus [m_x, M_x]$ such that $\Delta(x; z) = 0$ for some $x \in \mathbb{T}$.

If the condition (i) or (ii) of Theorem 2.3 holds, then by the assertions (i) and (ii) of Theorem 2.2 we have $E_{\min} \in \sigma$, hence there exists $x_1 \in \mathbb{T}$ such that $\Delta(x_1; E_{\min}) = 0$. Since $E_{\min} < m$, the function $\Delta(\cdot; E_{\min})$ is a regular in \mathbb{T} . Therefore, the number of zeros of this function is finite.

Let $\{x \in \mathbb{T} : \Delta(x; E_{\min}) = 0\} = \{x_1, \dots, x_n\}$ and k_j be the multiplicity of x_j for $j \in \{1, \dots, n\}$. The fact $E_{\min} < m$ implies that the difference $w_2(x, y) - z$ is positive for all $x, y \in \mathbb{T}$ and $z \leq E_{\min}$. Hence the function $(w_2(\cdot, \cdot) - z)^{-1}$ is an analytic one on \mathbb{T}^2 for all $z \leq E_{\min}$. Then there exists a number $\delta > 0$ such that for any $i, j \in \{1, \dots, n\}$ and $z \leq E_{\min}$ the following representations are valid

$$\frac{v_1(x)v_1(y)}{2(w_2(x,y)-z)} = \sum_{k=0}^{[k_i/2]} c_{ik}^{(1)}(z;x)(y-x_i)^k + (y-x_i)^{[k_i/2]+1} m_i^{(1)}(z;x,y), \tag{5.1}$$

for all $x \in \mathbb{T}$ and $y \in U_{\delta}(x_i)$;

$$\frac{v_1(x)v_1(y)}{2(w_2(x,y)-z)} = \sum_{k=0}^{[k_i/2]} c_{ik}^{(2)}(z;y)(x-x_i)^k + (x-x_i)^{[k_i/2]+1} m_i^{(2)}(z;x,y), \tag{5.2}$$

for all $x \in U_{\delta}(x_i)$ and $y \in \mathbb{T}$;

$$\frac{v_1(x)v_1(y)}{2(w_2(x,y)-z)} = \sum_{k=0}^{[k_i/2]} \sum_{r=0}^{[k_j/2]} d_{ij}^{kr}(z)(x-x_i)^k (y-x_j)^r
+ \sum_{k=0}^{[k_i/2]} \sum_{r=[k_j/2]+1}^{\infty} d_{ij}^{kr}(z)(x-x_i)^k (y-x_j)^r
+ \sum_{k=[k_i/2]+1}^{\infty} \sum_{r=0}^{[k_j/2]} d_{ij}^{kr}(z)(x-x_i)^k (y-x_j)^r + (x-x_i)^{[k_i/2]+1} (y-x_j)^{[k_j/2]+1} q_{ij}(z;x,y),$$
(5.3)

for all $(x,y) \in U_{\delta}(x_i) \times U_{\delta}(x_j)$. Here $[\cdot]$ is the entire part of a, for any $z \leq E_{\min}$ the numbers $d_{ij}^{kr}(z)$ are some real coefficients, the functions $c_{ik}^{(\alpha)}(z;\cdot)$, $\alpha=1,2;$ $m_i^{(1)}(z;\cdot,\cdot);$ $m_i^{(2)}(z;\cdot,\cdot)$ and $q_{ij}^{kr}(z;\cdot,\cdot)$ are some analytic functions on \mathbb{T} ; $\mathbb{T} \times U_{\delta}(x_i);$ $U_{\delta}(x_i) \times \mathbb{T}$ and $U_{\delta}(x_i) \times U_{\delta}(x_j)$, respectively.

Lemma 5.3 Let the assumption (i) or (ii) of Theorem 2.3 be fulfilled and $j \in \{1, ..., n\}$. Then there exist numbers C > 0 and $\delta > 0$ such that the inequality

$$\frac{|x - x_j|^{[k_j/2]+1}}{\sqrt{\Delta(x;z)}} \le C \tag{5.4}$$

holds for all $x \in U_{\delta}(x_i)$ and $z \leq E_{\min}$.

Proof. If the assumption (i) or (ii) of Theorem 2.3 holds, then by Theorem 2.2 we have $E_{\min} < m$. Since the function $\Delta(x;\cdot)$ is monotonically decreasing on $(-\infty, m)$ we have

$$\frac{|x - x_j|^{[k_j/2]+1}}{\sqrt{\Delta(x;z)}} \le \frac{|x - x_j|^{[k_j/2]+1}}{\sqrt{\Delta(x;E_{\min})}}$$

for all $z \leq E_{\min}$. Taking into account the fact that the number k_j is the multiplicity of the x_j and the function $\Delta(\cdot; E_{\min})$ is analytic on \mathbb{T} we obtain the inequality (5.4).

Lemma 5.4 Let the assumptions of Theorem 2.3 be satisfied. Then for any $z \leq E_{\min}$ the operator T(z) can be represented in the form $T(z) = T_0(z) + T_1(z)$, where the operator-valued function $T_0(\cdot)$ is continuous in the operator-norm in $(-\infty; E_{\min}]$ and $T_1(z)$ is a finite-dimensional operator for all $z \leq E_{\min}$ whose dimension is independent of z.

Proof. Since the operators $T_{00}(z)$, $T_{01}(z)$ and $T_{01}^*(z)$ are of rank one independently of z, it is sufficient to study the operator $T_{11}(z)$.

We denote the kernel of the integral operator $T_{11}(z)$ by $T_{11}(z;x,y)$, that is,

$$T_{11}(z;x,y) := \frac{v_1(x)v_1(y)}{2\sqrt{\Delta(x;z)}(w_2(x,y)-z)\sqrt{\Delta(y;z)}}.$$

First we will prove the statement of lemma under the assumption (i) or (ii) of Theorem 2.3. In this case $E_{\min} < m$ and using the representations (5.1)–(5.4) we obtain $T_{11}(z) = T_{11}^0(z) + T_{11}^1(z)$, where the kernels $T_{11}^0(z; x, y)$ and $T_{11}^1(z; x, y)$ of the integral operators $T_{11}^0(z)$ and $T_{11}^1(z)$ has form

$$T_{11}^{0}(z;x,y) := (1 - \chi_{V_{\delta}}(x))(1 - \chi_{V_{\delta}}(y))T_{11}(z;x,y) + \frac{(1 - \chi_{V_{\delta}}(x))}{\sqrt{\Delta(x;z)}} \sum_{i=1}^{n} \frac{\chi_{V_{\delta}}(y)(y - x_{i})^{[k_{i}/2]+1}}{\sqrt{\Delta(y;z)}} M_{i}^{(1)}(z;x,y) + \frac{(1 - \chi_{V_{\delta}}(y))}{\sqrt{\Delta(y;z)}} \sum_{i=1}^{n} \frac{\chi_{V_{\delta}}(x)(x - x_{i})^{[k_{i}/2]+1}}{\sqrt{\Delta(x;z)}} M_{i}^{(2)}(z;x,y) + \chi_{V_{\delta}}(x)\chi_{V_{\delta}}(y) \sum_{i=1}^{n} \frac{(x - x_{i})^{[k_{i}/2]+1}(y - x_{j})^{[k_{j}/2]+1}}{\sqrt{\Delta(x;z)}\sqrt{\Delta(y;z)}} Q_{ij}(z;x,y);$$

$$T_{11}^{1}(z;x,y) := \frac{(1-\chi_{V_{\delta}}(x))\chi_{V_{\delta}}(y)}{\sqrt{\Delta(x;z)}\sqrt{\Delta(y;z)}} \sum_{i=1}^{n} \sum_{k=0}^{[k_{i}/2]} (y-x_{i})^{k} c_{ik}^{(1)}(z;x)$$

$$+ \frac{\chi_{V_{\delta}}(x)(1-\chi_{V_{\delta}}(y))}{\sqrt{\Delta(x;z)}\sqrt{\Delta(y;z)}} \sum_{i=1}^{n} \sum_{k=0}^{[k_{i}/2]} (x-x_{i})^{k} c_{ik}^{(2)}(z;y)$$

$$+ \frac{\chi_{V_{\delta}}(x)\chi_{V_{\delta}}(y)}{\sqrt{\Delta(x;z)}\sqrt{\Delta(y;z)}} \sum_{i,j=1}^{n} \left(\sum_{k=0}^{[k_{i}/2]} \sum_{r=0}^{[k_{j}/2]} d_{ij}^{kr}(z)(x-x_{i})^{k}(y-x_{j})^{r}\right)$$

$$+ \sum_{k=0}^{[k_{i}/2]} \sum_{r=[k_{i}/2]+1}^{\infty} d_{ij}^{kr}(z)(x-x_{i})^{k}(y-x_{j})^{r} + \sum_{k=[k_{i}/2]+1}^{\infty} \sum_{r=0}^{[k_{j}/2]} d_{ij}^{kr}(z)(x-x_{i})^{k}(y-x_{j})^{r}\right),$$

respectively, where $V_{\delta} := \bigcup_{i=1}^{n} U_{\delta}(x_i), \ \chi_A(\cdot)$ is the characteristic function of the set $A \subset \mathbb{T}$,

$$\begin{split} M_i^{(1)}(z;x,y) &:= \left\{ \begin{array}{l} m_i^{(1)}(z;x,y), \ (x,y) \in \mathbb{T} \times U_{\delta}(x_i), \\ 0, & \text{otherwise}, \end{array} \right. \\ M_i^{(2)}(z;x,y) &:= \left\{ \begin{array}{l} m_i^{(2)}(z;x,y), \ (x,y) \in U_{\delta}(x_i) \times \mathbb{T}, \\ 0, & \text{otherwise}, \end{array} \right. \\ Q_{ij}(z;x,y) &:= \left\{ \begin{array}{l} q_{ij}(z;x,y), \ (x,y) \in U_{\delta}(x_i) \times U_{\delta}(x_j), \\ 0, & \text{otherwise}. \end{array} \right. \end{split}$$

Applying Lemma 5.3 we obtain that the function $T_{11}^0(z;\cdot,\cdot)$, $z \leq E_{\min}$ is square-integrable on \mathbb{T}^2 and converges almost everywhere to $T_{11}^0(E_{\min};\cdot,\cdot)$ as $z \to E_{\min} - 0$. Then by the Lebesgue dominated convergence theorem the operator $T_{11}^0(z)$ converges in the operator-norm to $T_{11}^0(E_{\min})$ as $z \to E_{\min} - 0$. The finite dimensionality of the operator $T_{11}^1(z)$ follows from the definition of $T_{11}^1(z;x,y)$. Now setting

$$T_0(z) := \left(\begin{array}{cc} 0 & 0 \\ 0 & T_{11}^0(z) \end{array} \right), \quad T_1(z) := \left(\begin{array}{cc} T_{00}(z) & T_{01}(z) \\ T_{01}^*(z) & T_{11}^1(z) \end{array} \right)$$

we complete proof of Lemma 5.4 under the assumption (i) or (ii) of Theorem 2.3.

Let the assumption (iii) of Theorem 2.3 be satisfied. Then by Theorem 2.2 we have $E_{\min} = m$. Applying Lemmas 5.1 and 5.2 and as well as inequality (2.1) one can see that the function $|T_{11}(z;\cdot,\cdot)|$ can be estimated by

$$C_1 \left(1 + \frac{|x|^{\alpha}|y|^{\alpha}}{x^2 + y^2} \right)$$

for $z \leq m$ with $\alpha \geq 1$. The latter function is a square-integrable on \mathbb{T}^2 and the function $T_{11}(z;\cdot,\cdot)$ converges almost everywhere to $T_{11}(m;\cdot,\cdot)$ as $z \to m-0$. Then by the Lebesgue dominated convergence theorem the operator $T_{11}(z)$ converges in the norm to $T_{11}(m)$ as $z \to m-0$. Now setting

$$T_0(z) := \begin{pmatrix} 0 & 0 \\ 0 & T_{11}(z) \end{pmatrix}, \quad T_1(z) := \begin{pmatrix} T_{00}(z) & T_{01}(z) \\ T_{01}^*(z) & 0 \end{pmatrix}$$

we complete proof of Lemma 5.4 under the assumption (iii) of Theorem 2.3.

We are now ready for the proof of Theorem 2.3.

Proof of Theorem 2.3. Using the Weyl inequality

$$n(\lambda_1 + \lambda_2, A_1 + A_2) \le n(\lambda_1, A_1) + n(\lambda_2, A_2)$$
 (5.5)

for the sum of compact operators A_1 and A_2 and for any positive numbers λ_1 and λ_2 we have

$$n(1,T(z)) \leq n(2/3,T_0(z)) + n(1/3,T_1(z))$$

$$\leq n(1/3,T_0(z) - T_0(E_{\min})) + n(1/3,T_0(E_{\min})) + n(1/3,T_1(z))$$
 (5.6)

for all $z < E_{\min}$.

By virtue of Lemma 5.4 the operator $T_0(E_{\min})$ is compact and hence $n(1/3, T_0(E_{\min})) < \infty$ and $n(1/3, T_0(z) - T_0(E_{\min}))$ tends to zero as $z \to E_{\min} - 0$. Since $T_1(z)$ is a finite-dimensional operator and its dimension is independent of $z, z < E_{\min}$, there exists a number F such that for all $z < E_{\min}$, we have $n(1/3, T_1(z)) \le F < \infty$. So, by the inequality (5.6) we obtain that the number n(1, T(z)) is finite for all $z < E_{\min}$.

Now Lemma 4.1 implies that N(z) = n(1, T(z)) as $z < E_{\min}$ and hence

$$\lim_{z \to E_{\min} - 0} N(z) = N(E_{\min}) \le n(1/3, T_0(E_{\min})) + n(1/3, T_1(E_{\min})) < \infty.$$

It means that the number of eigenvalues of H lying on the left of E_{\min} is finite.

6 Infiniteness of the number of eigenvalues of H

In this section we consider the case when the parameter functions $v_i(\cdot)$, $i = 1, 2, w_1(\cdot)$ and $w_2(\cdot, \cdot)$ have the special forms:

$$v_0(x) := 0, \quad w_1(x) := a, \quad v_1(x) := b, \quad a, b \in \mathbb{R} \setminus \{0\};$$

 $w_2(x, y) := \varepsilon(x - y), \quad \varepsilon(x) := 1 - \cos x.$

It is obvious that the function $w_2(\cdot,\cdot)$ has non-degenerate minimum at the points of the form (x,x) for any $x \in \mathbb{T}$. Then it is clear that the number $z = w_0$ is an eigenvalue of H with the associated eigenvector $f = (f_0,0,0)$ with $f_0 \neq 0$ and the equality holds $\sigma_{\rm ess}(H) = \{E_{\rm min}\} \cup [0,2] \cup \{E_{\rm max}\}$, where $E_{\rm min}$ and $E_{\rm max}$ are zeros of the function $\Delta(\cdot)$ defined on $\mathbb{C} \setminus [0,2]$ by

$$\Delta(z) := a - z - \frac{b^2}{2} \int_{\mathbb{T}} \frac{dt}{\varepsilon(t) - z}$$

such that $E_{\min} < 0$ and $E_{\max} > 2$.

We define the function $D(\cdot)$ on $\mathbb{C} \setminus \sigma_{\text{ess}}(H)$ as

$$D(z) := \prod_{k=0}^{\infty} D_k(z), \quad D_k(z) := 1 - \frac{1}{2\Delta(z)} d_k(z), \quad d_k(z) := \int_{\mathbb{T}} \frac{\cos(kt)dt}{\varepsilon(t) - z}.$$

The following lemma establishes a connection between eigenvalues of the operator H and zeros of the function $D(\cdot)$.

Lemma 6.1 The number $z \in \mathbb{C} \setminus (\sigma_{ess}(H) \cup \{w_0\})$ is an eigenvalue of H if and only if D(z) = 0. Moreover, if for some $k \in \mathbb{N}$ the number $z_k \in \mathbb{C} \setminus \sigma_{ess}(H)$ is an eigenvalue of H with $D_k(z_k) = 1 - \lambda_k(z_k) = 0$, then the corresponding eigenvector $f^{(k)}$ has the form $f^{(k)} := (0, f_1^{(k)}, f_2^{(k)})$, where the functions $f_1^{(k)}$ and $f_2^{(k)}$ are defined by

$$f_1^{(k)}(x) := \exp(\pm ikx), \quad f_2^{(k)}(x,y) := \frac{b(f_1^{(k)}(x) + f_1^{(k)}(y))}{2(\varepsilon(x-y) - z_k)}.$$
 (6.1)

Proof. Let the number $z \in \mathbb{C} \setminus (\sigma_{\text{ess}}(H) \cup \{w_0\})$ be an eigenvalue of H and $f = (f_0, f_1, f_2) \in \mathcal{H}$ be the corresponding eigenvector. Then f_0 , f_1 and f_2 satisfy the following system of equations

$$(w_0 - z)f_0 = 0;$$

$$(a - z)f_1(x) + b \int_{\mathbb{T}} f_2(x, t)dt = 0;$$

$$\frac{b}{2}(f_1(x) + f_1(y)) + (\varepsilon(x - y) - z)f_2(x, y) = 0.$$
(6.2)

Using the condition $z \neq w_0$ we get from the first equation of (6.3) that $f_0 = 0$. Since $z \notin [0, 2]$, from the third equation of the system (6.3) for f_2 we find

$$f_2(x,y) = -\frac{b(f_1(x) + f_1(y))}{2(\varepsilon(x-y) - z)}.$$
(6.3)

Substituting the expression (6.3) for f_2 into the second equation of the system (6.3) and using the fact that $\Delta(z) \neq 0$ for any $z \in \mathbb{C} \setminus \sigma_{\text{ess}}(H)$ we conclude that the number $z \in \mathbb{C} \setminus (\sigma_{\text{ess}}(H) \cup \{w_0\})$ is an eigenvalue of H if and only if the number 1 is an eigenvalue of the integral operator T(z) in $L_2(\mathbb{T})$ with the kernel

$$\frac{b^2}{2\Delta(z)(\varepsilon(x-y)-z)}.$$

Since the function $(\varepsilon(\cdot) - z)^{-1}$ is continuous on \mathbb{T} and $\Delta(z) \neq 0$ for all $z \in \mathbb{C} \setminus \sigma_{\text{ess}}(H)$, the operator $\widetilde{T}(z)$ is Hilbert-Schmidt and as well trace class. Hence, the determinant $\det(I - \widetilde{T}(z))$ of the operator $I - \widetilde{T}(z)$ exists and is given by the formula (see Theorem XIII.106 of [14])

$$\det(I - \widetilde{T}(z)) = \prod_{k=0}^{\infty} (1 - \lambda_k(z)), \tag{6.4}$$

where I is the identity operator on $L_2(\mathbb{T})$ and the numbers $\{\lambda_k(z)\}$ are the eigenvalues of T(z) counted with their algebraic multiplicities. By Theorem XIII.105 of [14] the number 1 is an eigenvalue of $\widetilde{T}(z)$ if and only if $\det(I - \widetilde{T}(z)) = 0$.

Let φ be the eigenfunction of $\widetilde{T}(z)$ associated with the eigenvalue λ , that is,

$$\lambda \varphi(x) = \frac{b^2}{2\Delta(z)} \int_{\mathbb{T}} \frac{\varphi(t)dt}{\varepsilon(x-t) - z}.$$

By expanding φ into a series with respect to the basis $\{\exp(ikx)\}_{k\in\mathbb{Z}}$ we obtain

$$\lambda c_k \exp(ikx) = \frac{b^2 c_k}{2\Delta(z)} \int_{\mathbb{T}} \frac{\exp(ikt)dt}{\varepsilon(x-t) - z}$$

or $\lambda(z) = b^2 d_k(z)/(2\Delta(z))$. Then for any $k \in \mathbb{Z}$ the eigenvalue $\lambda_k(z)$ of the operator $\widetilde{T}(z)$ in formula (6.4) can be expressed by $\lambda_k(z) = b^2 d_k(z)/(2\Delta(z))$ and the corresponding eigenfunction $\varphi_k(\cdot)$ has the form $\varphi_k(x) := \exp(ikx)$.

For any $k \in \mathbb{Z}$ and $z \in \mathbb{C} \setminus \sigma_{\mathrm{ess}}(H)$ we have $d_k(z) = d_{-k}(z)$. Hence, if the number $1 = \lambda_k(z_k)$, $z_k \in \mathbb{C} \setminus \sigma_{\mathrm{ess}}(H)$ is an eigenvalue of $\widetilde{T}(z_k)$, then $\varphi_k(x) = \exp(\pm ikx)$ is the corresponding eigenfunction of $\widetilde{T}(z_k)$. From here it follows that if the number $z_k \in \mathbb{C} \setminus \sigma_{\mathrm{ess}}(H)$ is an eigenvalue of H with $D_k(z_k) = 1 - \lambda_k(z_k) = 0$, then the corresponding eigenvector $f^{(k)}$ has the form $f^{(k)} := (0, f_1^{(k)}, f_2^{(k)})$, where $f_1^{(k)}$ and $f_2^{(k)}$ are defined by (6.1). \square Let $\mathbb{N}_0 := \mathbb{N} \cup \{0\}$.

Lemma 6.2 For the functions $\Delta(\cdot)$ and $d_k(\cdot)$, $k \in \mathbb{N}_0$, the equalities hold

$$\Delta(z) = a - z - \frac{\pi b^2}{\sqrt{z^2 - 2z}}, \quad d_k(z) = \frac{2\pi \left[1 - z - \sqrt{z^2 - 2z}\right]^k}{\sqrt{z^2 - 2z}}, \quad z < 0;$$

$$\Delta(z) = a - z + \frac{\pi b^2}{\sqrt{z^2 - 2z}}, \quad d_k(z) = \frac{2\pi \left[1 - z + \sqrt{z^2 - 2z}\right]^k}{\sqrt{z^2 - 2z}}, \quad z > 2.$$

Proof. The assertion of lemma for the case z < 0 can be proven similarly to Lemma 10 of [12]. We consider the case z > 2. Using the identity

$$\int_0^{\pi} \frac{\cos(kt)dt}{1 + 2c\cos t + c^2} = \frac{\pi(-c)^k}{1 - c^2},$$

where $k \in \mathbb{N}_0$ and 0 < c < 1, one can show that

$$\int_0^{\pi} \frac{\cos(kt)dt}{1 + (2c/(1+c^2))\cos t} = \frac{\pi(-c)^k}{\sqrt{1 - (2c)^2/(1+c^2)^2}}.$$
 (6.5)

Since the function $\varepsilon(\cdot)$ is an even, the function $d_k(\cdot)$ has form

$$d_k(z) = 2 \int_0^{\pi} \frac{\cos(kt)dt}{1 - \cos t - z} = \frac{2}{1 - z} \int_0^{\pi} \frac{\cos(kt)dt}{1 + \cos t/(z - 1)}.$$

Introducing the notation $c_z := z - 1 - \sqrt{z^2 - 2z}$ we obtain

$$d_k(z) = \frac{2}{1-z} \int_0^{\pi} \frac{\cos(kt)dt}{1 + 2c_z \cos t/(1 + c_z^2)}.$$

It is clear that $c_z \in (0,1)$ for all z > 2 and hence the equality (6.5) completes proof of lemma for the case z > 2.

Now we formulate the result about infiniteness of the discrete spectrum of H.

Theorem 6.3 (i) The operator H has an infinite number of eigenvalues $\{\xi_k^{(\alpha)}\}_0^{\infty}$ with $\alpha=1,2,3$ such that $\{\xi_k^{(1)}\}_0^{\infty} \subset (-\infty, E_{\min}), \ \{\xi_k^{(2)}\}_0^{\infty} \subset (E_{\max}, \infty), \ \{\xi_k^{(3)}\}_0^{\infty} \subset (2, E_{\max}) \ and$

$$\lim_{k \to \infty} \xi_k^{(1)} = E_{\min}, \quad \lim_{k \to \infty} \xi_k^{(2)} = \lim_{k \to \infty} \xi_k^{(3)} = E_{\max}.$$

For $\alpha = 1, 2, 3$ the multiplicity of every eigenvalue $\xi_k^{(\alpha)}$, $k \in \mathbb{N}$ is two, the multiplicity of $\xi_0^{(3)}$ is one or two and $\xi_0^{(1)}$, $\xi_0^{(2)}$ are simple eigenvalues of H. Moreover, the eigenvalues $\xi_k^{(1)}$ resp. $\xi_k^{(2)}$ are solutions of the rational equations

$$\frac{\pi b^2}{\Delta(z)} \frac{\left[1 - z - \sqrt{z^2 - 2z}\right]^k}{\sqrt{z^2 - 2z}} = 1, \quad z < E_{\min};$$

resp.

$$\frac{\pi b^2}{\Delta(z)} \frac{\left[1-z+\sqrt{z^2-2z}\right]^k}{\sqrt{z^2-2z}} = 1, \quad z>E_{\max}.$$

(ii) The operator H has no eigenvalues in $(E_{\min}, 0)$.

Proof. (i) For any fixed $k \in \mathbb{N}_0$ we have

$$\lim_{z \to \pm \infty} D_k(z) = 1, \quad \lim_{z \to E_{\min} - 0} D_k(z) = \lim_{z \to E_{\max} + 0} D_k(z) = -\infty.$$

Since the function $D_k(\cdot)$ is continuous in $(-\infty, E_{\min})$ and (E_{\max}, ∞) there exist numbers $\xi_k^{(1)} \in (-\infty, E_{\min})$ and $\xi_k^{(2)} \in (E_{\max}, \infty)$ such that $D_k(\xi_k^{(\alpha)}) = 0$ for $\alpha = 1, 2$. The equality $\lim_{z \to \pm \infty} D(z) = 1$ and the analyticity of the function $D(\cdot)$ on $\mathbb{C} \setminus (\{E_{\min}\} \cup [0, 2] \cup \{E_{\max}\})$ imply that $\lim_{k \to \infty} \xi_k^{(1)} = E_{\min}$ and $\lim_{k \to \infty} \xi_k^{(2)} = E_{\max}$. By Lemma 6.1 for any $\alpha = 1, 2$ and $k \in \mathbb{N}_0$ the number $\xi_k^{(\alpha)}$ is an eigenvalue of H and the corresponding eigenvector $f^{(k)}$ has the form $f^{(k)} := (0, f_1^{(k)}, f_2^{(k)})$, where $f_1^{(k)}$ and $f_2^{(k)}$ are defined by (6.1) with $z_k = \xi_k^{(\alpha)}$. Moreover, $\xi_0^{(1)}$, $\xi_0^{(2)}$ are simple eigenvalues and for any $k \in \mathbb{N}$ the multiplicities of $\xi_k^{(\alpha)}$ are two.

Note that the function $D_k(\cdot)$ is defined on $(2, E_{\text{max}})$ and for z > 2 we have

$$-1 < 1 - z + \sqrt{z^2 - 2z} < 0.$$

By Lemma 6.2 the function $D_k(\cdot)$ can be rewritten as

$$D_k(z) = 1 - \frac{1}{2\Delta(z)} \frac{\left[1 - z + \sqrt{z^2 - 2z}\right]^k}{\sqrt{z^2 - 2z}}, \quad z \in (2, E_{\text{max}});$$

therefore, for any fixed $z \in (2, E_{\text{max}})$ the equality $\lim_{k \to \infty} D_k(z) = 1$ holds.

It is clear that $\Delta(z) > 0$ for all $z \in (2, E_{\text{max}})$ and hence the inequality $D_{2k+1}(z) > 1$ holds for all $k \in \mathbb{N}_0$ and $z \in (2, E_{\text{max}})$. Since $\lim_{k \to \infty} D_k(z) = 1$ for any fixed $z \in (2, E_{\text{max}})$, there exists a subsequence $\{k_n\} \subset 2\mathbb{N}_0$ such that $D_{k_n}((E_{\text{max}} + 2)/2) > 0$ holds for any $n \in \mathbb{N}_0$.

Now the equality $\lim_{z\to E_{\max}-0} D_{2k}(z) = -\infty$ and the continuity of the function $D_k(\cdot)$ imply that $D_{k_n}(\xi_n^{(3)}) = 0$ for some $\xi_n^{(3)} \in ((E_{\max}+2)/2, E_{\max})$. It follows from the analyticity of $D(\cdot)$ on $\mathbb{C} \setminus (\{E_{\min}\} \cup [0,2] \cup \{E_{\max}\})$ that $\lim_{n\to\infty} \xi_n^{(3)} = E_{\max}$. Now repeated application of Lemma 6.1 implies that the number $\xi_n^{(3)}$ is an eigenvalue of H. Similarly, for any $n \in \mathbb{N}$ the multiplicity of $\xi_n^{(3)}$ is two. If $k_0 = 0$, then $\xi_0^{(3)}$ is a simple, otherwise its multiplicity is two.

(ii) It is clear that $0 < 1 - z - \sqrt{z^2 - 2z} < 1$ for z < 0. Then by Lemma 6.2 the function $D_k(\cdot)$ can be represented as

$$D_k(z) = 1 - \frac{1}{2\Delta(z)} \frac{\left[1 - z - \sqrt{z^2 - 2z}\right]^k}{\sqrt{z^2 - 2z}}, \quad z \in (E_{\min}; 0).$$

Since $\Delta(z) < 0$ for all $z \in (E_{\min}, 0)$ the inequality $D_k(z) > 1$ holds for such z and hence D(z) > 1. By Lemma 6.1 the operator has no eigenvalues in $(E_{\min}, 0)$.

7 The case $v_1(0) = 0$ and $\Delta(0; m) = 0$

In this section we are going to discuss the discrete spectrum of H for the case $v_1(0) = 0$ and $\Delta(0; m) = 0$. In this case the discrete spectrum of H might be finite or infinite depending on the behavior of the parameter functions.

Case I: Infiniteness. Let the parameter functions $v_1(\cdot)$, $w_1(\cdot)$ and $w_2(\cdot,\cdot)$ have the form

$$v_1(x) = \sqrt{\mu}\sin x, \quad \mu > 0; \quad w_1(x) = 1 + \sin^2 x;$$

$$w_2(x, y) = \varepsilon(x) + l\varepsilon(x + y) + \varepsilon(y), \quad \varepsilon(x) := 1 - \cos x, \ l > 0.$$
(7.1)

Then the function $w_2(\cdot,\cdot)$ has a unique non-degenerate zero minimum (m=0) at the point $(0,0) \in \mathbb{T}^2$ and $v_1(0) = 0$. It is easy to see that for

$$\Delta(x;z) = 1 + \sin^2 x - z - \frac{\mu}{2} \int_{\mathbb{T}} \frac{\sin^2 t \, dt}{\varepsilon(x) + l\varepsilon(x+t) + \varepsilon(t) - z}$$

we have $\Delta(0;0) = 0$ if and only if

$$\mu = \mu_0 := (1+l) \left(\int_0^{\pi} \frac{\sin^2 t \, dt}{\varepsilon(t)} \right)^{-1} = \frac{1+l}{\pi}.$$

The following decomposition plays an important role in the proof of the infiniteness of the discrete spectrum of H.

Lemma 7.1 The following decomposition

$$\Delta(x;z) = \Delta(0;0) + \frac{\mu\pi(1+2l-l^2)}{(1+l)^2\sqrt{1+2l}}\sqrt{x^2 - \frac{2(1+l)}{1+2l}z} + O(x^2) + O(\sqrt{|z|})$$

holds as $x \to 0$ and $z \to -0$.

Proof. Let $\delta > 0$ be sufficiently small and $\mathbb{T}_{\delta} := \mathbb{T} \setminus (-\delta, \delta)$. We rewrite the function $\Delta(\cdot; \cdot)$ in the form $\Delta(x; z) = \Delta_1(x; z) + \Delta_2(x; z)$, where

$$\Delta_{1}(x;z): = 1 + \sin^{2} x - z - \frac{\mu}{2} \int_{\mathbb{T}_{\delta}} \frac{\sin^{2} t \, dt}{\varepsilon(x) + l\varepsilon(x+t) + \varepsilon(t) - z},$$

$$\Delta_{2}(x;z): = -\frac{\mu}{2} \int_{-\delta}^{\delta} \frac{\sin^{2} t \, dt}{\varepsilon(x) + l\varepsilon(x+t) + \varepsilon(t) - z}.$$

Since $\Delta_1(\cdot;z)$ is an even analytic function on \mathbb{T} for any $z\leq 0$, we have

$$\Delta_1(x;z) = \Delta_1(0;0) + O(x^2) + O(|z|) \tag{7.2}$$

as $x \to 0$ and $z \to -0$. Using

$$\sin x = x + O(x^3), \quad 1 - \cos x = \frac{1}{2}x^2 + O(x^4), \quad x \to 0$$
 (7.3)

we obtain

$$\Delta_2(x;z) := -\mu \int_{-\delta}^{\delta} \frac{t^2 dt}{(1+l)x^2 + 2lxt + (1+l)t^2 - 2z} + O(x^2) + O(|z|)$$

as $x \to 0$ and $z \to -0$. For the convenience we rewrite the latter integral as

$$\int_{-\delta}^{\delta} \frac{t^2 dt}{(1+l)x^2 + 2lxt + (1+l)t^2 - 2z}$$

$$= \frac{2\delta}{1+l} - \frac{lx}{1+l} \int_{-\delta}^{\delta} \frac{2t dt}{(1+l)x^2 + 2lxt + (1+l)t^2 - 2z}$$

$$-\frac{(1+l)x^2 - 2z}{1+l} \int_{-\delta}^{\delta} \frac{dt}{(1+l)x^2 + 2lxt + (1+l)t^2 - 2z}.$$

Now we study each integral in the last equality. For the integral in the second summand we obtain

$$\int_{-\delta}^{\delta} \frac{2tdt}{(1+l)x^2 + 2lxt + (1+l)t^2 - 2z} = \frac{1}{1+l} \log \left| 1 + \frac{4lx\delta}{(1+l)x^2 - 2lx\delta + (1+l)\delta^2 - 2z} \right| - \frac{2lx}{1+l} \int_{-\delta}^{\delta} \frac{dt}{(1+l)x^2 + 2lxt + (1+l)t^2 - 2z}.$$

Since

$$\log \left| 1 + \frac{4lx\delta}{(1+l)x^2 - 2lx\delta + (1+l)\delta^2 - 2z} \right| = O(x)$$

as $x \to 0$, comparing the last expressions we obtain

$$\begin{split} & \int_{-\delta}^{\delta} \frac{t^2 dt}{(1+l)x^2 + 2lxt + (1+l)t^2 - 2z} = \frac{2\delta}{1+l} \\ & - \left(\frac{1+2l-l^2}{(1+l)^2}x^2 - \frac{2}{1+l}z\right) \int_{-\delta}^{\delta} \frac{dt}{(1+l)x^2 + 2lxt + (1+l)t^2 - 2z} + O(x^2) + O(|z|) \end{split}$$

as $x \to 0$ and $z \to -0$. Using the identity

$$\int_{a}^{b} \frac{dt}{x^2 + t^2} = \frac{1}{|x|} \left(\arctan \frac{b}{|x|} - \arctan \frac{a}{|x|} \right)$$
 (7.4)

we have

$$\int_{-\delta}^{\delta} \frac{dt}{(1+l)x^2 + 2lxt + (1+l)t^2 - 2z} = \frac{1}{1+l} \int_{-\delta}^{\delta} \frac{dt}{(t+\frac{l}{1+l}x)^2 + \frac{1+2l}{(1+l)^2}x^2 - \frac{2}{l+1}z}$$

$$= \frac{1}{(1+l)\sqrt{\frac{1+2l}{(1+l)^2}x^2 - \frac{2}{l+1}z}} \left(\arctan\frac{\delta + \frac{l}{1+l}x}{\sqrt{\frac{1+2l}{(1+l)^2}x^2 - \frac{2}{l+1}z}} + \arctan\frac{\delta - \frac{l}{1+l}x}{\sqrt{\frac{1+2l}{(1+l)^2}x^2 - \frac{2}{l+1}z}}\right).$$

The following properties of the arctan function

$$\arctan y + \arctan \frac{1}{y} = \frac{\pi}{2}, \quad y \ge 0 \quad \text{and} \quad \arctan y = O(y), \quad y \to 0$$
 (7.5)

imply that

$$\begin{split} & \int_{-\delta}^{\delta} \frac{dt}{(1+l)x^2 + 2lxt + (1+l)t^2 - 2z} = \\ & \left(\frac{1+2l-l^2}{(1+l)^2} x^2 - \frac{2}{1+l}z \right) \frac{\pi}{(1+l)\sqrt{\frac{1+2l}{(1+l)^2}x^2 - \frac{2}{l+1}z}} + O(\sqrt{\frac{1+2l}{(1+l)^2}x^2 - \frac{2}{l+1}z}) \end{split}$$

as $x \to 0$ and $z \to -0$. Taking into account

$$\left(\frac{1+2l-l^2}{(1+l)^2}x^2 - \frac{2}{1+l}z\right) \frac{\pi}{(1+l)\sqrt{\frac{1+2l}{(1+l)^2}x^2 - \frac{2}{l+1}z}}$$

$$= \pi \frac{1+2l-l^2}{(1+l)^2\sqrt{1+2l}}\sqrt{x^2 - \frac{2(1+l)}{2l+1}z} + O(\sqrt{-z}),$$

we obtain

$$\Delta_2(x;z) = \Delta_2(0;0) + \frac{\mu\pi(1+2l-l^2)}{(1+l)^2\sqrt{1+2l}}\sqrt{x^2 - \frac{2(1+l)}{2l+1}z} + O(x^2) + O(\sqrt{-z})$$
 (7.6)

as $x \to 0$. The equalities (7.2) and (7.6) give the proof of lemma.

Let $T(\delta; z)$ be the operator in $\mathcal{H}_0 \oplus \mathcal{H}_1$ defined by

$$T(\delta;z) := \begin{pmatrix} 0 & 0 \\ 0 & T_{11}(\delta;z) \end{pmatrix},$$

where $T_{11}(\delta;z)$ is the integral operator on $L_2(\mathbb{T})$ with the kernel

$$\frac{1}{\pi} \frac{(1+l)^2 \sqrt{1+2l}}{1+2l-l^2} \frac{1}{\sqrt[4]{x^2 - \frac{2(1+l)}{2l+1}z}} \frac{\chi_{(-\delta;\delta)}(x)\chi_{(-\delta;\delta)}(y)xy}{(1+l)x^2 + 2lxy + (1+l)y^2 - 2z} \frac{1}{\sqrt[4]{y^2 - \frac{2(1+l)}{2l+1}z}}.$$

Lemma 7.2 Let $\mu = \mu_0$. Then for any $z \leq 0$ the operator $F(z) := T(z) - T(\delta; z)$ is compact and the operator-valued function $F(\cdot)$ is continuous in the operator-norm in $(-\infty, 0]$.

Proof. Denote by $T_{11}(z; x, y)$ and $T_{11}(\delta, z; x, y)$ the kernel of the operator $T_{11}(z)$ and $T_{11}(\delta; z)$, respectively, and set $F(z; x, y) := T_{11}(z; x, y) - T_{11}(\delta, z; x, y)$. We split the function $F(z; \cdot, \cdot)$, z < 0 into four parts

$$F(z; x, y) = F_0(z; x, y) + F_1(z; x, y) + F_2(z; x, y) + F_3(z; x, y),$$

where

$$F_{0}(z;x,y): = (1 - \chi_{(-\delta,\delta)}(x)\chi_{(-\delta,\delta)}(y))T_{11}(z;x,y),$$

$$F_{1}(z;x,y): = \frac{\mu}{2} \frac{\chi_{(-\delta,\delta)}(x)}{\sqrt{\Delta(x;z)}} \frac{\chi_{(-\delta,\delta)}(y)}{\sqrt{\Delta(y;z)}}$$

$$\times \left(\frac{\sin x \sin y}{\varepsilon(x) + l\varepsilon(x+y) + \varepsilon(y) - z} - \frac{2xy}{(1+l)x^{2} + 2lxy + (1+l)y^{2} - 2z}\right),$$

$$F_{2}(z;x,y): = \left(\frac{\mu}{\sqrt{\Delta(x;z)}} - \frac{\mu}{\sqrt{\frac{\mu\pi(1+2l-l^{2})}{(1+l)^{2}\sqrt{1+2l}}}} \sqrt{x^{2} - \frac{2(1+l)}{1+2l}z}}\right) \frac{1}{\sqrt{\Delta(y;z)}}$$

$$\times \frac{\chi_{(-\delta,\delta)}(x)\chi_{(-\delta,\delta)}(y)xy}{(1+l)x^{2} + 2lxy + (1+l)y^{2} - 2z},$$

$$F_{3}(z;x,y): = \frac{1}{\sqrt{\frac{\mu\pi(1+2l-l^{2})}{(1+l)^{2}\sqrt{1+2l}}}} \frac{\chi_{(-\delta,\delta)}(x)\chi_{(-\delta,\delta)}(y)xy}{(1+l)x^{2} + 2lxy + (1+l)y^{2} - 2z}}$$

$$\times \left(\frac{\mu}{\sqrt{\Delta(y;z)}} - \frac{\mu}{\sqrt{\frac{\mu\pi(1+2l-l^{2})}{(1+l)^{2}\sqrt{1+2l}}}} \sqrt{y^{2} - \frac{2(1+l)}{1+2l}z}}\right).$$

We show that the functions $F_i(z;\cdot,\cdot)$, i=0,1,2,3 are square-integrable on \mathbb{T}^2 for any fixed $z \leq 0$. First we note that for any fixed $z \leq 0$ the function $F_0(z;\cdot,\cdot)$ is bounded on \mathbb{T}^2 and hence it is a square-integrable on this set.

Using the decompositions (7.3) we obtain that there exists C > 0 such that for any $z \le 0$ the inequality

$$\left| \frac{\sin x \sin y}{\varepsilon(x) + l\varepsilon(x+y) + \varepsilon(y) - z} - \frac{2xy}{(1+l)x^2 + 2lxy + (1+l)y^2 - 2z} \right| \le C|xy|, \quad x, y \in (-\delta, \delta)$$

holds. Therefore, for any fixed $z \leq 0$ the function $F_1(z;\cdot,\cdot)$ is a square-integrable on \mathbb{T}^2 .

By Lemma 7.1 for any $x \in (-\delta, \delta)$ and $z \in (-\delta, 0)$ we get the estimate

$$\left| \frac{\mu}{\sqrt{\Delta(x\,;z)}} - \frac{\mu}{\sqrt{\frac{\mu\pi(1+2l-l^2)}{(1+l)^2\sqrt{1+2l}}}\sqrt{x^2 - \frac{2(1+l)}{1+2l}z}} \right| \leq \frac{C\sqrt{-z}}{\sqrt[4]{(x^2-z)^3}} + C\sqrt{|x|}.$$

It follows from the last estimate and Lemma 7.1 that

$$|F_2(z;x,y)| \leq \frac{C\sqrt{-z}}{\sqrt[4]{(x^2-z)^3}} \frac{|xy|}{(1+l)x^2 + 2lxy + (1+l)y^2 - 2z} \frac{1}{\sqrt[4]{y^2-z}} + \frac{C|x|^{3/2}|y|}{(1+l)x^2 + 2lxy + (1+l)y^2 - 2z} \frac{1}{\sqrt[4]{y^2-z}}$$

or

$$|F_2(z;x,y)| \le \frac{C|x|^{1/2}|y|^{3/2}\sqrt{-z}}{(1+l)x^2 + 2lxy + (1+l)y^2 - 2z} + \frac{C|x|^{3/2}|y|^{3/2}}{(1+l)x^2 + 2lxy + (1+l)y^2 - 2z}$$

for all $x, y \in (-\delta, \delta)$ with some positive constant C. Since

$$\int_{-\delta}^{\delta} \int_{-\delta}^{\delta} \left(\frac{|x|^{1/2} |y|^{3/2}}{(1+l)x^2 + 2lxy + (1+l)y^2 - 2z} \right)^2 dx dy \le C|\log(-z)|,$$

for any fixed $z \leq 0$ the function $F_2(z;\cdot,\cdot)$ is a square-integrable on \mathbb{T}^2 . By the same way we can show the square-integrability of $F_3(z;\cdot,\cdot)$ on \mathbb{T}^2 for any fixed $z \leq 0$.

Hence, the operator $T_{11}(z) - T_{11}(\delta; z)$ belongs to the Hilbert-Schmidt class for all $z \leq 0$. In combination with the continuity of the kernel of the operator with respect to z < 0, this implies the continuity of $T_{11}(z) - T_{11}(\delta; z)$ with respect to $z \leq 0$.

By the definition the operators $T_{00}(z)$, $T_{01}(z)$ and $T_{01}^*(z)$ are rank 1 operators and they are continuous from the left up to z=0. Consequently the operator F(z) is compact and the operator-valued function $F(\cdot)$ is continuous in the operator-norm in $(-\infty, 0]$.

By the structure of $T(\delta; z)$ we have $\sigma(T(\delta; z)) = \{0\} \cup \sigma(T_{11}(\delta; z))$.

The subspace of functions g having support in $(-\delta, \delta)$ is an invariant subspace for the operator $T_{11}(\delta; z)$. Let $T_{11}^{(0)}(\delta; z)$ be the restriction of the operator $T_{11}(\delta; z)$ to the subspace $L_2(-\delta, \delta)$, that is, the integral operator with the kernel

$$\frac{1}{\pi} \frac{(1+l)^2 \sqrt{1+2l}}{1+2l-l^2} \frac{1}{\sqrt[4]{x^2 - \frac{2(1+l)}{2l+1}z}} \frac{xy}{(1+l)x^2 + 2lxy + (1+l)y^2 - 2z} \frac{1}{\sqrt[4]{y^2 - \frac{2(1+l)}{2l+1}z}},$$

 $x, y \in (-\delta, \delta)$. Then we have $\sigma(T_{11}(\delta; z)) = \sigma(T_{11}^{(0)}(\delta; z)), z < 0$.

Let $L_2^{\rm o}(-\delta,\delta)$ and $L_2^{\rm e}(-\delta,\delta)$ be the spaces of odd and even functions, respectively. It is easily to check that $T_{11}^{(0)}(\delta;z):L_2^{\rm o}(-\delta,\delta)\to L_2^{\rm o}(-\delta,\delta)$ and $T_{11}^{(0)}(\delta;z):L_2^{\rm e}(-\delta,\delta)\to L_2^{\rm e}(-\delta,\delta)$.

Let us consider the unitary operator

$$U_{\rm o}: L_2^{\rm o}(-\delta, \delta) \to L_2(0, \delta), \quad (U_{\rm o}f)(x) = \sqrt{2}f(x).$$

Then

$$U_{\rm o}^{-1}: L_2(0,\delta) \to L_2^{\rm o}(-\delta,\delta), \quad (U_{\rm o}^{-1}f)(x) = \begin{cases} \frac{1}{\sqrt{2}}f(x) & \text{as } x \ge 0\\ -\frac{1}{\sqrt{2}}f(x) & \text{as } x < 0. \end{cases}$$

Let $T_{\rm o}(z) := U_{\rm o} T_{11}^{(0)}(\delta; z) U_{\rm o}^{-1}$. Then $\sigma(T_{11}^{(0)}(\delta; z)) \supset \sigma(T_{\rm o}(z))$, where $T_{\rm o}(z)$ is the integral operator acting on $L_2(0, \delta)$ with the kernel

$$T_{o}(z;x,y): = \frac{1}{\pi} \frac{(1+l)^{2}\sqrt{1+2l}}{1+2l-l^{2}} \frac{1}{\sqrt[4]{x^{2} - \frac{2(1+l)}{2l+1}z}} \left[\frac{xy}{(1+l)x^{2} + 2lxy + (1+l)y^{2} - 2z} + \frac{xy}{(1+l)x^{2} - 2lxy + (1+l)y^{2} - 2z} \right] \frac{1}{\sqrt[4]{y^{2} - \frac{2(1+l)}{2l+1}z}}.$$

Let $T_1(z)$, z < 0 be an integral operator on $L_2(0, \delta)$ with the kernel $\chi_{\Omega(z)}(x)T_0(z; x, y)\chi_{\Omega(z)}(y)$, where $\Omega(z) := (|z|^{1/2}, \delta]$.

Lemma 7.3 Let $\mu = \mu_0$. Then for any $z \in (-\delta, 0]$ the operator $G(z) := T_0(z) - T_3(z)$ is compact and the operator-valued function $G(\cdot)$ is continuous in the operator-norm in $(-\delta, 0]$.

The subspace of functions g having support in $\Omega(z)$ is an invariant subspace for the operator $T_1(z)$. Let $T_2(z)$ be the restriction of the operator $T_1(z)$ to the subspace $L_2(\Omega(z))$, that is, the integral operator with kernel $T_2(z;x,y) := T_0(z;x,y), x,y \in \Omega(z)$. Then we have $\sigma(T_1(z)) = \sigma(T_2(z))$ for $z \in (-\delta, 0]$.

Let us consider the unitary dilation

$$U: L_{2}(\Omega(z)) \to L_{2}(-\pi, \pi),$$

$$(Uf)(x) = \sqrt{\frac{R(z)}{2\pi}} (|z|^{1/2} e^{\frac{R(z)}{2\pi}(x+\pi)})^{1/2} f(|z|^{1/2} e^{\frac{R(z)}{2\pi}(x+\pi)}),$$

$$U^{-1}: L_{2}(-\pi, \pi) \to L_{2}(\Omega(z)),$$

$$(U^{-1}f)(p) = \sqrt{\frac{2\pi}{R(z)}} |x|^{-1/2} f(\frac{2\pi}{R(z)} \log \frac{|x|}{|z|^{1/2}} - \pi), \quad R(z) := -\log \frac{|z|^{1/2}}{\delta}.$$

The operator $T_3(z) := UT_2(z)U^{-1}$ is integral operator on $L_2(-\pi,\pi)$ with the kernel

$$\begin{split} T_3(z;x,y) &:= \frac{1}{\pi} \frac{(1+l)^2 \sqrt{1+2l}}{1+2l-l^2} \frac{R(z)}{2\pi} \frac{e^{\frac{3R(z)}{4\pi}(x+\pi)}}{(e^{\frac{R(z)}{\pi}(x+\pi)} + \frac{2(1+l)}{2l+1})^{1/4}} \\ &\times \left[\frac{1}{(1+l)e^{\frac{R(z)}{\pi}(x+\pi)} + 2le^{\frac{R(z)}{2\pi}(x+\pi)}e^{\frac{R(z)}{2\pi}(y+\pi)} + (1+l)e^{\frac{R(z)}{\pi}(y+\pi)} + 2} \right. \\ &+ \frac{1}{(1+l)e^{\frac{R(z)}{\pi}(x+\pi)} - 2le^{\frac{R(z)}{2\pi}(x+\pi)}e^{\frac{R(z)}{2\pi}(y+\pi)} + (1+l)e^{\frac{R(z)}{\pi}(y+\pi)} + 2} \right] \frac{e^{\frac{3R(z)}{4\pi}(y+\pi)}}{(e^{\frac{R(z)}{4\pi}(y+\pi)} + \frac{2(1+l)}{2l+1})^{1/4}}. \end{split}$$

Lemma 7.4 Let $\mu = \mu_0$. Then for any $z \in (-\delta, 0]$ the operator $G_1(z) := T_3(z) - T_4(z)$ is compact and the operator-valued function $G_1(\cdot)$ is continuous in the operator-norm in $(-\delta, 0]$, where the operator $T_4(z)$ is an integral operator on $L_2(-\pi, \pi)$ with kernel $T_4(z; x)$,

$$T_4(z;x) := \frac{1}{\pi} \frac{(1+l)^2 \sqrt{1+2l}}{1+2l-l^2} \frac{R(z)}{2\pi} \left[\frac{1}{2(1+l)\operatorname{ch} \frac{R(z)}{2\pi}(x) + 2l} + \frac{1}{2(1+l)\operatorname{ch} \frac{R(z)}{2\pi}(x) - 2l} \right].$$

Let us define in $L_2(-\pi,\pi)$ the operator $S(z), z \in (-\delta,0)$ by

$$S(z) := \sum_{k \in \mathbb{Z}} \lambda_k(z)(\varphi_k, \cdot)\varphi_k,$$

$$\lambda_k(z) := \frac{(1+l)^2 \sqrt{1+2l}}{1+2l-l^2} \frac{1}{1+l} \frac{1}{\sin(\arccos\frac{l}{1+l})} \frac{\sinh(\arccos\frac{l}{1+l}\frac{2k\pi}{R(z)}) + \sinh((\pi - \arccos\frac{l}{1+l})\frac{2k\pi}{R(z)})}{\sinh\frac{2k\pi^2}{R(z)}},$$

where $\varphi_0(x) := \frac{1}{2\pi}$ and $\varphi_n(x) := \frac{1}{\sqrt{2\pi}} e^{inx}$ as $n \neq 0$.

Lemma 7.5 Let $\mu = \mu_0$. Then for any $z \in (-\delta, 0]$ the operator $G_2(z) := T_4(z) - S(z)$ is compact and the operator-valued function $G_2(\cdot)$ is continuous in the operator-norm in $(-\delta, 0]$.

Proof. Note that the operator $T_4(z)$ is convolution type. Therefore the eigenvalues of $T_4(z)$ can be found. By the Hilbert-Schmidt theorem the operator $T_4(z)$ can be decomposed as

$$T_4(z) = \sum_{n \in \mathbb{Z}} u_n(z)(\varphi_n, \cdot)\varphi_n,$$

where

$$u_k(z) := \frac{(1+l)^2 \sqrt{1+2l}}{1+2l-l^2} \frac{R(z)}{2\pi^2} \int_{-\pi}^{\pi} \left[\frac{e^{int}}{2(1+l)\operatorname{ch}(\frac{R(z)}{2\pi}t) + 2l} + \frac{e^{int}}{2(1+l)\operatorname{ch}(\frac{R(z)}{2\pi}t) - 2l} \right] dt.$$

We represent $u_k(z)$ as

$$u_k(z) = \tilde{u}_k(z) - O_k(z), \tag{7.7}$$

where

$$\tilde{u}_k(z) := \frac{(1+l)^2 \sqrt{1+2l}}{1+2l-l^2} \frac{1}{\pi} \int_{-\infty}^{\infty} \left[\frac{e^{i\frac{2\pi}{R(z)}nt}}{2(1+l)\text{ch}s+2l} + \frac{e^{i\frac{2\pi}{R(z)}nt}}{2(1+l)\text{ch}s-2l} \right] dt,$$

$$O_k(z) := \frac{(1+l)^2 \sqrt{1+2l}}{1+2l-l^2} \frac{1}{\pi} \int_{|t| > \frac{R(z)}{2}} \left[\frac{e^{i\frac{2\pi}{R(z)}nt}}{2(1+l)\text{ch}s+2l} + \frac{e^{i\frac{2\pi}{R(z)}nt}}{2(1+l)\text{ch}s-2l} \right] dt.$$

Using the equality

$$\frac{1}{2\pi} \int_{-\infty}^{\infty} \frac{e^{itr} \sin a}{\operatorname{ch} t + \cos a} dt = \frac{\operatorname{sh}(ar)}{\operatorname{sh}(\pi r)}, \quad \text{as} \quad |a| < \pi$$

we get the equality $\tilde{u}_k(z) = \lambda_k(z)$. It follows from (7.7) that the difference $G_2(z) := T_4(z) - S(z)$ is a Hilbert-Schmidt operator and continuous up to z = 0.

Lemma 7.6 Let $\mu = \mu_0$. There exist l > 0 and $\rho > 0$ such that $\lim_{z \to -0} n(1 + \rho, S(z)) = \infty$.

Proof. Since

$$\lambda_k(0) := \lim_{z \to -0} \lambda_k(z) = \frac{(1+l)^2 \sqrt{1+2l}}{1+2l-l^2} \frac{1}{1+l} \frac{1}{\sin(\arccos\frac{l}{1+l})} \frac{\arccos\frac{l}{1+l} + \sin(\pi - \arccos\frac{l}{1+l})}{\pi}.$$

It is easy to check that for any $k \in \mathbb{Z}$ it takes place $\lambda_k(0) > 1$ as l = 2.

Main result of this section is the following statement.

Theorem 7.7 Let $\mu = \mu_0$ and the parameter functions $v_1(\cdot)$, $w_1(\cdot)$ and $w_2(\cdot, \cdot)$ be given by (7.1). Then there exists a l > 0 such that the operator H has a infinite number of eigenvalues lying below $E_{\min} = 0$.

Proof. Using the Weyl's inequality (5.5) for any $z \in (-\delta, 0]$ we have the inequalities

$$n(1+\delta,S(z)) \leq n(1+\frac{4\delta}{5},T_4(z)) + n(\frac{\delta}{5},S(z)-T_4(z)),$$

$$n(1+\frac{4\delta}{5},T_4(z)) \leq n(1+\frac{3\delta}{5},T_3(z)) + n(\frac{\delta}{5},T_4(z)-T_3(z)),$$

$$n(1+\frac{3\delta}{5},T_3(z)) = n(1+\frac{3\delta}{5},T_1(z)) \leq n(1+\frac{2\delta}{5},T_1(z)) + n(\frac{\delta}{5},T_1(z)-T_0(z)),$$

$$n(1+\frac{2\delta}{5},T_0(z)) = n(1+\frac{\delta}{5},T_{11}^{(0)}(\delta;z)) + n(\frac{\delta}{5},T_0(z)-T_{11}^{(0)}(\delta;z)),$$

$$n(1+\frac{\delta}{5},T_{11}^{(0)}(\delta;z)) = n(1+\frac{\delta}{5},T(\delta;z)) \leq n(1,T(z)) + n(\frac{\delta}{5},T(\delta;z)-T(z)).$$

According to Lemmas 7.2–7.5 we get the inequalities

$$n(\frac{\delta}{5}, S(z) - T_4(z)) < \infty, \ n(\frac{\delta}{5}, T_4(z) - T_3(z)) < \infty, \ n(\frac{\delta}{5}, T_1(z) - T_0(z)) < \infty,$$

$$n(\frac{\delta}{5}, T_0(z) - T_{11}^{(0)}(\delta; z)) < \infty, \quad n(\frac{\delta}{5}, T(\delta; z) - T(z)) < \infty$$

for any $z \in (-\delta, 0]$. Then $n(1 + \rho, S(z)) \leq C + n(1, T(z))$, where C > 0 does not depend of $z \in (-\delta, 0]$. Hence by Lemma 7.6 we obtain the proof of the theorem.

Case II: Finiteness. Let the parameter functions $v_1(\cdot)$, $w_1(\cdot)$ and $w_2(\cdot,\cdot)$ have the form

$$v_1(x) = \sqrt{\mu(1-\cos x)}, \quad \mu > 0; \quad w_1(x) = 2 - \cos x; \quad w_2(x,y) = 2 - \cos x - \cos y.$$

Then the function $w_2(\cdot,\cdot)$ has a unique non-degenerate global zero minimum (m=0) at the point $(0,0) \in \mathbb{T}^2$ and $v_1(0) = 0$. It is easy to see that for

$$\Delta(x;z) = 2 - \cos x - z - \mu \int_0^{\pi} \frac{(1 - \cos t)^2 dt}{2 - \cos x - \cos t - z}$$

we have $\Delta(0;0) = 0$ if and only if $\mu = 1/\pi$.

Lemma 7.8 Let $\mu = 1/\pi$. Then there exist the numbers $C_1, C_2 > 0$ and $\delta > 0$ such that

$$C_1 x^2 \le \Delta(x;0) \le C_2 x^2, \quad x \in U_{\delta}(0).$$

Proof. Let $\delta > 0$ be sufficiently small. We rewrite the function $\Delta(\cdot; 0)$ in the form $\Delta(x; 0) = \Delta_1(x) + \Delta_2(x)$, where

$$\Delta_1(x) := 2 - \cos x - \mu \int_{\delta}^{\pi} \frac{(1 - \cos t)^2 dt}{2 - \cos x - \cos t - z}, \ \Delta_2(x; z) := -\mu \int_{0}^{\delta} \frac{(1 - \cos t)^2 dt}{2 - \cos x - \cos t - z}.$$

Since $\Delta_1(\cdot)$ is an even analytic function on \mathbb{T} , we have

$$\Delta_1(x) = \Delta_1(0) + O(x^2) \tag{7.8}$$

as $x \to 0$. Using the expansion (7.3) for $1 - \cos x$ we obtain

$$\Delta_2(x) := -\frac{\mu}{4} \int_0^{\delta} \frac{t^4 dt}{x^2 + t^2} + O(x^2)$$

as $x \to 0$. Since

$$\int_0^{\delta} \frac{t^4 dt}{x^2 + t^2} = \frac{\delta^3}{3} - \delta x^2 + x^4 \int_0^{\delta} \frac{dt}{x^2 + t^2},$$

by the properties (7.4) and (7.5) we have

$$\Delta_2(x) = \Delta_2(0) + O(x^2) \tag{7.9}$$

as $x \to 0$. Recall that if $\mu = 1/\pi$, then $\Delta(0;0) = 0$. Now, taking into account the equalities (7.8) and (7.9) we obtain $\Delta(x;0) = O(x^2)$ as $x \to 0$, which implies that there exist $C_1, C_2 > 0$ and $\delta > 0$ such that the assertion of lemma holds.

Lemma 7.9 Let $\mu = 1/\pi$. For any $z \le 0$ the operator T(z) is compact and continuous on the left up to z = 0.

Proof. Let $\mu = 1/\pi$. Denote by Q(z; x, y) the kernel of the integral operator $T_{11}(z)$, z < 0, that is,

$$Q(z; x, y) := \frac{v_1(x)v_1(y)}{2\sqrt{\Delta(x; z)}(w_2(x, y) - z)\sqrt{\Delta(y; z)}}.$$

By virtue of decomposition (7.3) and Lemma 7.8 the kernel Q(z; x, y) is estimated by the square-integrable function

$$C_1\left(1+\frac{\chi_\delta(x)\chi_\delta(y)|x||y|}{x^2+y^2}\right),$$

defined on \mathbb{T}^2 , where $\chi_{\delta}(\cdot)$ is the characteristic function of $(-\delta, \delta)$. Hence for any $z \leq 0$ the operator $T_{11}(z)$ is Hilbert-Schmidt.

The kernel function of $T_{11}(z)$, z < 0 is continuous in $x, y \in \mathbb{T}$. Therefore the continuity of the operator $T_{11}(z)$ from the left up to z = 0 follows from Lebesgue's dominated convergence theorem. Since for all $z \leq 0$ the operators $T_{00}(z)$, $T_{01}(z)$ and $T_{01}^*(z)$ are of rank 1 and continuous from the left up to z = 0 one concludes that T(z) is compact and continuous from the left up to z = 0.

Using Lemma 7.9 we can now proceed analogously to the proof of Theorem 2.3 to show the finiteness of the negative discrete spectrum of H.

8 Application

In this section we investigate the spectrum of A_2 , introduced in Section 1 applying the results for H. We recall that the operator A_2 has a 3×3 tridiagonal block operator matrix representation

$$\mathcal{A}_2 := \left(\begin{array}{ccc} \mathcal{A}_{00} & \mathcal{A}_{01} & 0 \\ \mathcal{A}_{01}^* & \mathcal{A}_{11} & \mathcal{A}_{12} \\ 0 & \mathcal{A}_{12}^* & \mathcal{A}_{22} \end{array} \right),$$

where matrix elements A_{ij} are defined by

$$\mathcal{A}_{00}f_0^{(\sigma)} = \varepsilon \sigma f_0^{(\sigma)}, \quad \mathcal{A}_{01}f_1^{(\sigma)} = \alpha \int_{\mathbb{T}} v(t)f_1^{(-\sigma)}(t)dt,$$

$$(\mathcal{A}_{11}f_1^{(\sigma)})(x) = (\varepsilon \sigma + w(x))f_1^{(\sigma)}(x), \quad (\mathcal{A}_{12}f_2^{(\sigma)})(x) = \alpha \int_{\mathbb{T}} v(t)f_2^{(-\sigma)}(x,t)dt,$$

$$(\mathcal{A}_{22}f_2^{(\sigma)})(x,y) = (\varepsilon \sigma + w(x) + w(y))f_2^{(\sigma)}(x,y), \quad f = \{f_0^{(\sigma)}, f_1^{(\sigma)}, f_2^{(\sigma)}; \sigma = \pm\} \in \mathcal{L}_2.$$

We make the following assumptions: $\varepsilon > 0$; the dispersion $w(\cdot)$ is an analytic on \mathbb{T} and has a unique zero minimum at the point $0 \in \mathbb{T}$; $v(\cdot)$ is a real-valued analytic function on \mathbb{T} ; the coupling constant $\alpha > 0$ is an arbitrary.

Consider the following permutation operator

$$\Phi: \mathcal{L}_2 \to \mathcal{H} \oplus \mathcal{H},$$

$$\Phi: (f_0^{(+)}, f_0^{(-)}, f_1^{(+)}, f_1^{(-)}, f_2^{(+)}, f_2^{(-)}) \to (f_0^{(+)}, f_1^{(-)}, f_2^{(+)}, f_0^{(-)}, f_1^{(+)}, f_2^{(-)}).$$

To investigate the spectral properties of \mathcal{A}_2 we introduce the following two bounded self-adjoint operators $\mathcal{A}_2^{(\sigma)}$, $\sigma=\pm$, which acts in $\mathcal{F}_s^{(2)}(L_2(\mathbb{T}))$ as

$$\mathcal{A}_2^{(\sigma)} := \left(\begin{array}{ccc} \widehat{\mathcal{A}}_{00}^{(\sigma)} & \widehat{\mathcal{A}}_{01} & 0\\ \widehat{\mathcal{A}}_{01}^* & \widehat{\mathcal{A}}_{11}^{(\sigma)} & \widehat{\mathcal{A}}_{12}\\ 0 & \widehat{\mathcal{A}}_{12}^* & \widehat{\mathcal{A}}_{22}^{(\sigma)} \end{array} \right)$$

with the entries

$$\widehat{\mathcal{A}}_{00}^{(\sigma)} f_0 = \varepsilon \sigma f_0, \quad \widehat{\mathcal{A}}_{01} f_1 = \alpha \int_{\mathbb{T}} v(t) f_1(t) dt,$$

$$(\widehat{\mathcal{A}}_{11}^{(\sigma)} f_1)(x) = (-\varepsilon \sigma + w(x)) f_1(x), \quad (\widehat{\mathcal{A}}_{12} f_2)(x) = \alpha \int_{\mathbb{T}} v(t) f_2(x, t) dt,$$

$$(\widehat{\mathcal{A}}_{22}^{(\sigma)} f_2)(x, y) = (\varepsilon \sigma + w(x) + w(y)) f_2(x, y), \quad (f_0, f_1, f_2) \in \mathcal{F}_{s}^{(2)}(L_2(\mathbb{T})).$$

The definitions of the operators A_2 , $A_2^{(\sigma)}$ and Φ imply that

$$\Phi \mathcal{A}_2 \Phi^{-1} = \text{diag}\{\mathcal{A}_2^{(+)}, \mathcal{A}_2^{(-)}\}.$$

The following theorem describes the relation between spectra of A_2 and $A_2^{(\sigma)}$.

Theorem 8.1 The equality $\sigma(A_2) = \sigma(A_2^{(+)}) \cup \sigma(A_2^{(-)})$ holds. Moreover,

$$\sigma_{\mathrm{ess}}(\mathcal{A}_2) = \sigma_{\mathrm{ess}}(\mathcal{A}_2^{(+)}) \cup \sigma_{\mathrm{ess}}(\mathcal{A}_2^{(-)}), \quad \sigma_{\mathrm{p}}(\mathcal{A}_2) = \sigma_{\mathrm{p}}(\mathcal{A}_2^{(+)}) \cup \sigma_{\mathrm{p}}(\mathcal{A}_2^{(-)}).$$

Remark 8.2 Since the part of $\sigma_{disc}(\mathcal{A}_2^{(\sigma)})$ can be located in $\sigma_{ess}(\mathcal{A}_2)$ we have the inclusion

$$\sigma_{\rm disc}(\mathcal{A}_2) \subseteq \sigma_{\rm disc}(\mathcal{A}_2^{(+)}) \cup \sigma_{\rm disc}(\mathcal{A}_2^{(-)}).$$
 (8.1)

To describe the location of the essential spectrum of A_2 we introduce the following two families of bounded self-adjoint operators $h^{(\sigma)}(x)$, $x \in \mathbb{T}$, which acts in $\mathcal{F}_s^{(1)}(L_2(\mathbb{T}))$ as

$$h^{(\sigma)}(x) := \begin{pmatrix} h_{00}^{(\sigma)}(x) & h_{01} \\ h_{01}^* & h_{11}^{(\sigma)}(x) \end{pmatrix},$$

where

$$h_{00}^{(\sigma)}(x)f_0 = (-\sigma\varepsilon + w(x))f_0, \quad h_{01}f_1 = \frac{\alpha}{\sqrt{2}} \int_{\mathbb{T}} v(t)f_1(t)dt,$$

$$(h_{11}^{(\sigma)}(x)f_1)(y) = (\sigma\varepsilon + w(x) + w(y))f_1(y), \quad (f_0, f_1) \in \mathcal{F}_{s}^{(1)}(L_2(\mathbb{T})).$$

By Theorem 2.1 for the essential spectrum of $\mathcal{A}_2^{(\sigma)}$ the following equality holds

$$\sigma_{\mathrm{ess}}(\mathcal{A}_2^{(\sigma)}) = \bigcup_{x \in \mathbb{T}} \sigma_{\mathrm{disc}}(h^{(\sigma)}(x)) \cup [\sigma \varepsilon, 2M_w + \sigma \varepsilon], \quad M_w := \max_{x \in \mathbb{T}} w(x).$$

Now taking into account last equality we obtain from Theorem 8.1 that

$$\sigma_{\rm ess}(\mathcal{A}_2) = \bigcup_{\sigma = \pm} \bigcup_{x \in \mathbb{T}} \sigma_{\rm disc}(h^{(\sigma)}(x)) \cup [-\varepsilon, 2M_w - \varepsilon] \cup [\varepsilon, 2M_w + \varepsilon].$$

To estimate the lower bound of the essential spectrum of \mathcal{A}_2 , for any $x \in \mathbb{T}$ we define the Fredholm determinant $\Delta^{(\sigma)}(x;\cdot)$:

$$\Delta^{(\sigma)}(x\,;z):=-\sigma\varepsilon+w(x)-z-\frac{\alpha^2}{2}\int_{\mathbb{T}}\frac{v^2(t)dt}{\sigma\varepsilon+w(x)+w(t)-z}$$

in $\mathbb{C} \setminus [\sigma \varepsilon + w(x), M_w + \sigma \varepsilon + w(x)]$, associated with the operator $h^{(\sigma)}(x)$. By Lemma 3.2 for the discrete spectrum of $h^{(\sigma)}(x)$ the equality

$$\sigma_{\text{disc}}(h^{(\sigma)}(x)) = \{ z \in \mathbb{C} \setminus [\sigma \varepsilon + w(x), M_w + \sigma \varepsilon + w(x)] : \Delta^{(\sigma)}(x; z) = 0 \}$$

holds. By definition of $\Delta^{(\sigma)}(\cdot;\cdot)$ we have

$$\min_{x \in \mathbb{T}} \Delta^{(\sigma)}(x; z) = -\varepsilon \sigma - z - \frac{\alpha^2}{2} \int_{\mathbb{T}} \frac{v^2(t)dt}{\sigma \varepsilon + w(t) - z}, \quad z \le \sigma \varepsilon.$$

If we set $E_{\min}^{(\sigma)} := \min \sigma_{\text{ess}}(\mathcal{A}_2^{(\sigma)})$, then Theorem 8.1 implies

$$E_{\min} := \min \sigma_{\text{ess}}(\mathcal{A}_2) = \min \{ E_{\min}^{(-)}, E_{\min}^{(+)} \}.$$

It is clear that $\min_{x \in \mathbb{T}} \Delta^{(+)}(x; -\varepsilon) < 0$ for all $\alpha > 0$ and hence $E_{\min} \leq E_{\min}^{(+)} < -\varepsilon$.

Now we study the lower bound of the essential spectrum of $\mathcal{A}_2^{(-)}$. If

$$\int_{\mathbb{T}} \frac{v^2(t)dt}{w(t)} = \infty,$$

then again $\min_{x \in \mathbb{T}} \Delta^{(-)}(x; -\varepsilon) < 0$ for all $\alpha > 0$, that is, $E_{\min} \leq E_{\min}^{(-)} < -\varepsilon$. If

$$\int_{\mathbb{T}} \frac{v^2(t)dt}{w(t)} < \infty,$$

then

$$\min_{x \in \mathbb{T}} \Delta^{(-)}(x; -\varepsilon) < 0 \Leftrightarrow \alpha > \alpha_0 := 2\sqrt{\varepsilon} \left(\int_{\mathbb{T}} \frac{v^2(t)dt}{w(t)} \right)^{-1/2};$$

$$\min_{x \in \mathbb{T}} \Delta^{(-)}(x; -\varepsilon) \ge 0 \Leftrightarrow \alpha \le \alpha_0.$$

So, by Theorem 2.2 we obtain $E_{\min}^{(-)} < -\varepsilon$ for all $\alpha > \alpha_0$ and $E_{\min}^{(-)} = -\varepsilon$ for all $\alpha \leq \alpha_0$.

The above analysis leads to $E_{\min} < -\varepsilon$ for all $\alpha > 0$.

Since the parameter functions of $\mathcal{A}_2^{(+)}$ satisfy the conditions of Theorem 2.3, it has finitely many eigenvalues smaller than $E_{\min}^{(+)}$ for all $\alpha > 0$. Similarly for any $\alpha > \alpha_0$ the operator $\mathcal{A}_2^{(-)}$ has a finitely many eigenvalues smaller than $E_{\min}^{(-)}$. For the case $\alpha \leq \alpha_0$ we have $E_{\min}^{(-)} = -\varepsilon$ and $E_{\min} < E_{\min}^{(-)}$. Hence the operator $\mathcal{A}_2^{(-)}$ has a finitely many eigenvalues smaller than E_{\min} . Now by the inclusion (8.1) we conclude that the operator \mathcal{A}_2 has a finitely many eigenvalues smaller than E_{\min} .

Acknowledgment

This work was supported by the IMU Einstein Foundation Program. T. H. Rasulov wishes to thank the Berlin Mathematical School and Weierstrass Institute for Applied Analysis and Stochastics for the invitation and hospitality.

References

- [1] S. Albeverio, S. N. Lakaev, T. H. Rasulov. On the spectrum of an Hamiltonian in Fock space. Discrete spectrum asymptotics. J. Stat. Phys., **127**:2 (2007), 191–220.
- [2] S. Albeverio, S. N. Lakaev, T. H. Rasulov. The Efimov effect for a model operator associated with the Hamiltonian of a non conserved number of particles. Methods Funct. Anal. Topology, **13**:1 (2007), 1–16.
- [3] K. O. Friedrichs. Perturbation of spectra in Hilbert space. Amer. Math. Soc. Providence, Rhole Island, 1965.

- [4] I. M. Glazman. Direct methods of the qualitative spectral analysis of singular differential operators. J.: IPS Trans., 1965.
- [5] I. C. Gohberg, M. G. Krein. Introduction to the theory of linear nonselfadjoint operators. Amer. Math. Soc. Providence, Rhode Island, 1969.
- [6] M. Hübner, H. Spohn. Spectral properties of the spin-boson Hamiltonian. Ann. Inst. Henri Poincaré, 62:3 (1995), 289–323.
- [7] S. N. Lakaev, T. Kh. Rasulov. A model in the theory of perturbations of the essential spectrum of multi-particle operators. Math. Notes. **73**:4 (2003), 521–528.
- [8] S. N. Lakaev, T. Kh. Rasulov. Efimov's effect in a model of perturbation theory of the essential spectrum. Funct. Anal. Appl. **37** (2003), no. 1, 69–71.
- [9] V. A. Malishev, R. A. Minlos. Linear infinite-particle operators. Translations of Mathematical Monographs. 143, AMS, Providence, RI, 1995.
- [10] R. A. Minlos, H. Spohn. The three-body problem in radioactive decay: the case of one atom and at most two photons. Topics in Statistical and Theoretical Physics. Amer. Math. Soc. Transl., Ser. 2, 177, AMS, Providence, RI, 1996, 159–193.
- [11] A. I. Mogilner. Hamiltonians in solid state physics as multiparticle discrete Schrödinger operators: problems and results. Advances in Sov. Math., 5 (1991), 139–194.
- [12] M. É. Muminov, N. M. Aliev. Spectrum of the three-particle Schrödinger operator on a one-dimensional lattice. Theor. Math. Phys., 171:3 (2012), 754–768.
- [13] T. Kh. Rasulov. Discrete spectrum of a model operator in Fock space. Theor. Math. Phys., **152**:3 (2007), 1313–1321.
- [14] M. Reed, B. Simon. Methods of modern mathematical physics. IV: Analysis of Operators. Academic Press, New York, 1979.
- [15] H. Spohn. Ground states of the spin-boson Hamiltonian. Comm. Math. Phys., 123 (1989), 277–304.
- [16] C. Tretter. Spectral theory of block operator matrices and applications. Imperial College Press, 2008.
- [17] Yu. Zhukov, R. Minlos. Spectrum and scattering in a "spin-boson" model with not more than three photons. Theor. Math. Phys., 103:1 (1995), 398–411.