One-Step Piecewise Polynomial Galerkin Methods for Initial Value Problems*

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Abstract. A new approach to the numerical solution of systems of first-order ordinary differential equations is given by finding local Galerkin approximations on each subinterval of a given mesh of size h. One step at a time, a piecewise polynomial, of degree n and class C^0 , is constructed, which yields an approximation of order $O(h^{2n})$ at the mesh points and $O(h^{n+1})$ between mesh points. In addition, the jth derivatives of the approximation on each subinterval have errors of order $O(h^{n-j+1})$, $1 \le j \le n$. The methods are related to collocation schemes and to implicit Runge-Kutta schemes based on Gauss-Legendre quadrature, from which it follows that the Galerkin methods are A-stable.

1. Introduction. In this paper, we show how Galerkin's method can be employed to devise one-step methods for systems of nonlinear first-order ordinary differential equations. The basic idea is to find local *n*th degree polynomial Galerkin approximations on each subinterval of a given mesh and to match them together continuously, but not smoothly.

For each $n \ge 1$, a method is defined (Section 2) which uses an *n*-point Gauss-Legendre quadrature formula to evaluate certain inner products in the Galerkin equations. For sufficiently small step size h, a unique numerical solution exists and may be found by successive substitution (Section 3). After showing that these Galerkin methods are also collocation methods (Section 4) and implicit Runge-Kutta methods (Section 5), we show that the mesh point errors are of the order $O(h^{2n})$, and the global errors are of the order $O(h^{n+1})$ for the approximate solution and $O(h^{n-i+1})$, $1 \le j \le n$, for its *j*th derivatives (Section 6). A proof of the *A*-stability of the methods is given in Section 7, and numerical results are presented in Section 8.

Discrete one-step methods based on quadrature, other than the classical Runge-Kutta methods, have been studied by several authors, including the explicit schemes in [12, p. 101], [13], [14], [22] and the implicit schemes in [1], [2], [3], [6, Chapters 4, 9], [10], [12, p. 159]. Also, discrete block implicit methods are given in [21], [24], [25]. The methods of this paper, however, yield continuous piecewise polynomial approximations with the inherent benefit of derivative approximations. Earlier uses of piecewise polynomials may be found in [4], [5], [11], [15], [16], [17], [26].

Finally, we remark that recent "semidiscrete" Galerkin methods [7], [9], [18], [19], [23] reduce initial-boundary value problems to systems of ordinary differential equations. When combined with such methods, our techniques open the possibility of "fully discrete" Galerkin methods for these problems.

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2. Piecewise Polynomial Galerkin Methods. We consider the numerical solution of only a single nonlinear ordinary differential equation

$$(1) u'(t) = f(t, u(t)), t_0 \leq t,$$

$$(2) u(t_0) = u_0$$

on a finite interval $[t_0, t_N]$, although the results carry over to systems of such equations. We assume that $f(t, x) \in C^{2n}$ in $[t_0, t_N] \times (-\infty, \infty)$, so that the exact solution $u(t) \in C^{2n+1}[t_0, t_N]$, $n \ge 1$, and we also assume that f has a Lipschitz constant L in this same region.

Let π : $t_i = t_0 + ih$, $0 \le i \le N$, be a uniform mesh for the sake of simplicity. (It will be seen that our arguments do not depend crucially on this assumption since our method is a one-step method and step size changes are easy.) Then, we may approximate u(t) on each subinterval by an nth degree polynomial

$$y(t) = \sum_{j=1}^{n+1} b_i^{(i)} \varphi_{i+j}(t), \qquad t_i \leq t \leq t_{i+1}, 0 \leq i \leq N-1,$$

where $\varphi_{i+j}(t)$ are basis functions which are *n*th degree polynomials on each $[t_i, t_{i+1}]$. For example, $\{\varphi_k\}_{k=1}^{N+n}$ might be the *n*th degree *B*-spline basis functions of Schoenberg [20] or some other piecewise polynomial basis. Since the b_i may change from one subinterval to the next, y(t) need not be as smooth as the $\varphi_k(t)$.

We require that y(t) be continuous on $[t_0, t_N]$ and that it provide a local Galerkin approximation on each subinterval $[t_i, t_{i+1}]$, $0 \le i \le N - 1$. Accordingly, on each subinterval, we write the following n + 1 equations (one linear, n nonlinear) for the $b_i^{(i)}$, $1 \le j \le n + 1$,

(4)
$$y_{i+} = y_{i-}, \quad i \ge 1,$$

= $u_0, \quad i = 0,$

(5)
$$(y' - f(t, y), \varphi_{i+k})_i = 0, \quad 2 \le k \le n+1, 0 \le i \le N-1,$$

using the notation

$$(v, w)_i = \int_{t_i}^{t_{i+1}} v(t)w(t) dt.$$

To obtain a computational form of (4)–(5), we assume that the $(\varphi'_{i+j}, \varphi_{i+k})_i$ in (5) are computed exactly, i.e., analytically or by an exact quadrature formula, while the inner products $(f, \varphi_{i+k})_i$ are replaced by the *n-point Gauss-Legendre quadrature* formula having the form

(6)
$$\int_{t_i}^{t_{i+1}} v(t) dt = h \sum_{k=1}^n w_k v(\sigma_{i,k}) + O(h^{2n+1}),$$

(7)
$$\sigma_{i,k} = t_i + \theta_k h, \quad 1 \leq k \leq n,$$

where $w_k > 0$ and θ_k are the weights and abscissae for [0, 1]. The result is that (4)-(5) are replaced by the following set of N systems of n + 1 nonlinear equations to be solved in succession

(8)
$$Ab^{(i)} = c^{(i)}(b^{(i)}), \quad 0 \le i \le N-1,$$

where

(9)
$$\mathbf{b}^{(i)} = \{b_1^{(i)}, b_2^{(i)}, \cdots, b_{n+1}^{(i)}\}^T,$$

(10)
$$A_{k,j} = \varphi_{i+j}(t_i), \qquad k = 1,$$

$$= (\varphi_{i+k}, \varphi'_{i+j})_i, \qquad 2 \le k \le n+1, 1 \le j \le n+1,$$

$$c_k^{(i)}(\mathbf{b}^{(i)}) = y_i, \qquad k = 1,$$

(11)
$$= h \sum_{m=1}^{n} w_m f \left(\sigma_{i,m}, \sum_{j=1}^{n+1} b_j^{(i)} \varphi_{i+j}(\sigma_{i,m}) \right) \varphi_{i+k}(\sigma_{i,m}), \qquad 2 \leq k \leq n+1.$$

We consider only the cases where A is nonsingular. Certainly, A will be nonsingular when $\{\varphi_{i+k}\}_{k=2}^{n+1}$ span \mathcal{O}_{n-1} , the class of (n-1)st degree polynomials. For then, Ab⁽ⁱ⁾ = 0 implies $y(t_i) = 0$ and $(y', \varphi_{i+k})_i = 0$, $2 \le k \le n+1$, which, in turn, imply y' = 0, y = 0 and b⁽ⁱ⁾ = 0. However, this condition is not necessary, since A is nonsingular in the case of the cubic (n=3) B-spline basis functions used for the computations given in Section 8, but $\{\varphi_{i+k}\}_{k=2}^4$ do not span \mathcal{O}_2 . Since we may multiply (8) by A^{-1} , our numerical method depends on the solution of

(12)
$$\mathbf{b}^{(i)} = \mathbf{A}^{-1} \mathbf{c}^{(i)} (\mathbf{b}^{(i)}), \qquad 0 \le i \le N - 1.$$

3. Existence and Uniqueness of the Numerical Solution. Having let L denote the Lipschitz constant for f in $[t_0, t_N] \times R$, where $R \equiv (-\infty, \infty)$, we use the l_{∞} -norm to show that the right side of (12) is a contraction mapping on R^{n+1} when h is sufficiently small. Since

$$||\mathbf{A}^{-1}\mathbf{c}^{(i)}(\mathbf{b}) - \mathbf{A}^{-1}\mathbf{c}^{(i)}(\mathbf{b}^*)||_{\infty} \leq ||\mathbf{A}^{-1}||_{\infty} ||\mathbf{c}^{(i)}(\mathbf{b}) - \mathbf{c}^{(i)}(\mathbf{b}^*)||_{\infty}$$

and

$$||c^{(i)}(b) - c^{(i)}(b^*)||_{\infty} \leq hQ_1L ||b - b^*||_{\infty},$$

where

(13)
$$Q_1 = \max_{2 \le k \le n+1} \sum_{m=1}^n w_m |\varphi_{i+k}(\sigma_{i,m})| \sum_{i=1}^{n+1} |\varphi_{i+i}(\sigma_{i,m})|,$$

it is clear that

$$||\mathbf{A}^{-1}\mathbf{c}^{(i)}(\mathbf{b}) - \mathbf{A}^{-1}\mathbf{c}^{(i)}(\mathbf{b}^*)||_{\infty} \leq hQ_2 ||\mathbf{b} - \mathbf{b}^*||_{\infty},$$

where

(14)
$$Q_2 = Q_1 L ||\mathbf{A}^{-1}||_{\infty}.$$

Thus, we have a contraction mapping, and (12) has a unique solution which may be found by successive substitution when

$$(15) h < Q_2^{-1}.$$

4. The Galerkin Method as a Collocation Method. We show here that the approximate solution y(t) satisfies (1) at the quadrature points in each subinterval. Using (11), we may write (12) as

(16)
$$b_{i}^{(i)} = A_{i,1}^{-1} y_{i} + \sum_{m=1}^{n} \gamma_{i,m} f(\sigma_{i,m}, y(\sigma_{i,m})), \qquad 1 \leq j \leq n+1,$$

where

$$\gamma_{i,m} = h w_m \sum_{k=2}^{n+1} A_{i,k}^{-1} \varphi_{i+k}(\sigma_{i,m}).$$

Then, from substituting (16) into (3), we have at the quadrature points

(17)
$$y'(\sigma_{i,k}) = \alpha_k y_i + \sum_{m=1}^{n} \beta_{m,k} f(\sigma_{i,m}, y(\sigma_{i,m})), \qquad 1 \leq k \leq n,$$

where

$$\alpha_k = \sum_{i=1}^{n+1} A_{i,1}^{-1} \varphi'_{i+i}(\sigma_{i,k})$$

and

$$\beta_{m,k} = \sum_{i=1}^{n+1} \gamma_{i,m} \varphi'_{i+i}(\sigma_{i,k}).$$

In the following, we make use of the fact that whenever f is independent of u and $f \in \mathcal{O}_{n-1}$, the exact solution $u \in \mathcal{O}_n$ and $y \equiv u$. This follows because the quadrature (6) is exact for $v \in \mathcal{O}_{2n-1}$, in this case $f\varphi_{i+k} \in \mathcal{O}_{2n-1}$, and the exact computation of $(f, \varphi_{i+k})_i$ means (8) is equivalent to (4)–(5). Since u satisfies (4)–(5) and v satisfies (8), they satisfy equivalent equations in this case and, by uniqueness, $u \equiv v$.

Let $q(t) \in \mathcal{O}_n$ be defined by $q(t_i) = 1$, $q'(\sigma_{i,k}) = 0$, $1 \le k \le n$, and let f = q' so that u' = f, $u(t_i) = 1$ leads to u = q = y on $[t_i, t_{i+1}]$. Substituting y = q and f = q' into (17) yields

$$\alpha_k = 0, \qquad 1 \le k \le n.$$

Now for each r, $1 \le r \le n$, let $q_r(t) \in \mathcal{O}_n$ be defined by $q_r(t_i) = 0$, $q'_r(\sigma_{i,k}) = \delta_{r,k}$, $1 \le k \le n$, and let $f = q'_r$ and $u(t_i) = 0$ so that $u = q_r = y$. This time, substituting $y = q_r$ and $f = q'_r$ into (17) shows that

$$\beta_{r,k} = \delta_{r,k}, \qquad 1 \le r, k \le n.$$

Consequently, (17) becomes the collocation equation

$$(20) y'(\sigma_{i,k}) = f(\sigma_{i,k}, y(\sigma_{i,k})), 1 \leq k \leq n,$$

showing that one-step collocation to (1) at the quadrature points by means of a continuous piecewise nth degree polynomial is equivalent to the Galerkin method.

Notice that the proof of this collocation property depends on the use of exactly n points in a quadrature formula (6) which is exact for $v \in \mathcal{O}_{2n-1}$. The proof would break down if (6) had more than n points or different weights and abscissae.

5. The Galerkin Method as an Implicit Runge-Kutta Method. Wright [26] has shown that any one-step collocation method is equivalent to some implicit Runge-Kutta method. Having already shown that the Galerkin method is equivalent to a

certain one-step collocation method, we now derive the *particular* implicit Runge-Kutta method to which they are both equivalent. Of course, the Galerkin and collocation methods yield continuous approximations, so "equivalent" here means "matches the discrete values" of the Runge-Kutta approximation.

From (3) and (16), we have

(21)
$$y_{i+1} = \bar{\alpha}y_i + \sum_{n=1}^{n} \bar{\beta}_m f(\sigma_{i,m}, y(\sigma_{i,m})),$$

where

$$\bar{\alpha} = \sum_{i=1}^{n+1} A_{i,1}^{-1} \varphi_{i+i}(t_{i+1})$$

and

$$\bar{\beta}_m = \sum_{i=1}^{n+1} \gamma_{i,m} \varphi_{i+j}(t_{i+1}).$$

If we let f = 0, $u(t_i) = 1$ so that u = 1 = y, then substituting y = 1 and f = 0 into (21) produces

$$\bar{\alpha} = 1.$$

Next, for each r, $1 \le r \le n$, let $q_r(t) \in \mathcal{O}_n$ be defined as in Section 4. Now, substitution of $y = q_r$ and $f = q'_r$ into (21) leads to

$$a_r(t_{i+1}) = \bar{\beta}_r$$

Since the *n*-point Gauss-Legendre formula (6) is exact for elements of \mathcal{O}_{n-1} , we also have

$$q_r(t_{i+1}) = \int_{t_i}^{t_{i+1}} q_r'(t) dt = h \sum_{k=1}^n w_k q_r'(\sigma_{i,k}) = h w_r,$$

from which it follows that

$$\bar{\beta}_r = h w_r, \qquad 1 \le r \le n.$$

Together, (21)-(23) imply

(24)
$$y_{i+1} = y_i + h \sum_{m=1}^{n} w_m f(\sigma_{i,m}, y(\sigma_{i,m})),$$

and this is simply the implicit Runge-Kutta method based on the n-point Gauss-Legendre formula (6). Again, the proof of (24) depends on the fact that (6) is a Gauss-Legendre formula with exactly n points.

Thus, each of Butcher's implicit Runge-Kutta methods based on *n*-point Gauss-Legendre quadrature [2] has a corresponding "equivalent" Galerkin method using *n*th degree piecewise polynomials.

6. Error Bounds. In the following, a technique similar to that used by Shampine and Watts [21], [25] is employed to obtain asymptotic error bounds for the *discrete values* given by an implicit Runge-Kutta method. We view the Galerkin method as

a discrete one-step method and use Henrici's theory [12, Chapter 2] of such methods. Continuous error bounds are then obtained from the discrete ones.

First, we need to define an increment function. Since, from (20), y'(t) interpolates f(t, y(t)) at $\sigma_{i,k}$, $1 \le k \le n$, the Lagrangian representation for y'(t) is

(25)
$$y'(t) = \sum_{k=1}^{n} l_k(t) f(\sigma_{i,k}, y(\sigma_{i,k})), \quad t_i \leq t \leq t_{i+1},$$

where

$$l_k(t) = \prod_{j=1, j \neq k}^{n} \frac{(t - \sigma_{i,j})}{(\sigma_{i,k} - \sigma_{i,j})}, \qquad 1 \leq k \leq n.$$

Integrating (25) leads to

(26)
$$y(t) = y_i + \sum_{k=1}^{n} f(\sigma_{i,k}, y(\sigma_{i,k})) \int_{t_i}^{t} l_k(s) ds, \quad t_i \leq t \leq t_{i+1}.$$

Using (26), we now may write the Runge-Kutta form of the Galerkin method (24), in terms of an increment function Φ ,

$$(27) y_{i+1} = y_i + h\Phi(t_i, y_i; h), 0 \le i \le N-1,$$

where Φ satisfies

(28)
$$\Phi(t_i, y_i; h) = \sum_{m=1}^{n} w_m g_m(t_i, y_i; h)$$

and

$$g_m(t_i, y_i; h) = f(\sigma_{i,m}, y(\sigma_{i,m}))$$

(29)
$$= f\left(t_i + \theta_m h, y_i + \sum_{k=1}^n g_k(t_i, y_i; h) \int_{t_i}^{t_i + \theta_m h} l_k(s) ds\right), \quad 1 \leq m \leq n.$$

In order for Henrici's theory to apply, we must show that Φ is Lipschitz continuous with respect to y in $\Omega \equiv [t_0, t_N] \times R \times [0, h_0]$. If, for any $i, 0 \le i \le N - 1$, and any $y_i^* \in R$, $y^*(t)$ is the Galerkin approximate solution to u' = f(t, u), $u(t_i) = y_i^*$, $t_i \le t \le t_{i+1}$, then (26) holds for y^*

$$(26') y^*(t) = y_i^* + \sum_{k=1}^n f(\sigma_{i,k}, y^*(\sigma_{i,k})) \int_{t_i}^t l_k(s) ds, t_i \leq t \leq t_{i+1}.$$

Letting B_0 be a constant such that

(30)
$$\sum_{k=1}^{n} \max_{t_i \leq t \leq (i,k)} \left| \int_{t_i}^{t} l_k(s) ds \right| \leq hB_0, \quad 0 \leq i \leq N-1,$$

and subtracting (26) and (26') leads to

(31)
$$\max_{t \leq t \leq t \leq t+1} |y(t) - y^*(t)| \leq \frac{1}{1 - h_0 B_0 L} |y_i - y_i^*|, \quad 0 \leq i \leq N - 1,$$

where $0 \le h \le h_0 < (B_0L)^{-1}$. The Lipschitz condition then follows from (28), (29) and (31) since, for $0 \le h \le h_0$ and $0 \le i \le N - 1$,

$$|\Phi(t_{i}, y_{i}; h) - \Phi(t_{i}, y_{i}^{*}; h)| \leq \sum_{m=1}^{n} w_{m} |g_{m}(t_{i}, y_{i}; h) - g_{m}(t_{i}, y_{i}^{*}; h)|$$

$$\leq L \sum_{m=1}^{n} w_{m} |y(\sigma_{i,m}) - y^{*}(\sigma_{i,m})|$$

$$\leq \frac{L}{1 - h_{0}B_{0}L} |y_{i} - y_{i}^{*}|,$$

where $\sum_{m=1}^{n} w_m = 1$.

Now, we may prove

THEOREM 1. Assume that $f(t, x) \in C^{2n}$ in $[t_0, t_N] \times R$ so that $u(t) \in C^{2n+1}[t_0, t_N]$, and denote by L the Lipschitz constant for f in this region. Let the Galerkin method be defined as in Section 2 for some piecewise polynomial basis functions of degree $n \ge 1$ and the n-point Gauss-Legendre quadrature formula (6). If Q_2 and B_0 are defined by (14) and (30), respectively, and $0 < h \le h_0$ where $0 < h_0 < \min(Q_2^{-1}, (B_0 L)^{-1})$, then there exists a constant M such that

$$|u_i - y_i| \leq Mh^{2n}, \quad 0 \leq i \leq N.$$

Proof. The local truncation error τ_i is defined from (24) by

$$u_{i+1} = u_i + h \sum_{m=1}^{n} w_m f(\sigma_{i,m}, u(\sigma_{i,m})) + \tau_i.$$

Thus,

$$\tau_{i} = \int_{t_{i}}^{t_{i+1}} f(t, u(t)) dt - h \sum_{m=1}^{n} w_{m} f(\sigma_{i,m}, u(\sigma_{i,m}))$$

and, from (6), $|\tau_i| \leq Kh^{2n+1}$, where K is a constant depending on the maximum of $u^{(2n+1)}(t)$ on $[t_0, t_N]$. The bound (33) follows immediately from Henrici's Theorem 2.2 [12]. Q.E.D.

The discrete error bounds (33) agree with those for Butcher's methods [2].

We obtain continuous error bounds in

THEOREM 2. Let the hypotheses of Theorem 1 hold. Then there exist constants E_i , $0 \le j \le n$, such that

(34)
$$\max_{t \le t \le t} |u(t) - y(t)| \le E_0 h^{n+1},$$

and

(35)
$$\max_{t \leq t \leq t_{i+1}} |u^{(i)}(t) - y^{(i)}(t)| \leq E_i h^{n-j+1}, \quad 1 \leq j \leq n, \, 0 \leq i \leq N-1.$$

Proof. We write u(t) in the same form as y(t) in (26) by using the *n*-point Lagrangian quadrature found there

(36)
$$u(t) = u_i + \int_{t_i}^t f(s, u(s)) ds$$

$$= u_i + \sum_{k=1}^n f(\sigma_{i,k}, u(\sigma_{i,k})) \int_{t_i}^t l_k(s) ds + R_n(t), \qquad t_i \leq t \leq t_{i+1},$$

where $R_n(t) = O(h^{n+1})$. Subtracting from (26), we find that

$$\max_{t \leq t \leq t_{i+1}} |u(t) - y(t)| \leq \frac{1}{1 - h_0 B_0 L} |u_i - y_i| + O(h^{n+1}), \quad 0 \leq i \leq N - 1;$$

and (34) follows from (33). If we differentiate (26) and (36) j times using $R_n^{(i)}(t) = O(h^{n-i+1})$ and subtract, we can show that

$$\max_{\substack{i \leq i \leq i+1\\ i \leq j \leq i+1}} |u^{(i)}(t) - y^{(i)}(t)| \leq LB_i \max_{\substack{1 \leq k \leq n\\ 1 \leq k \leq n}} |u(\sigma_{i,k}) - y(\sigma_{i,k})| + O(h^{n-j+1}),$$

for $1 \le j \le n$, $0 \le i \le N-1$, where

$$\sum_{k=1}^{n} \max_{t \leq t \leq t \leq t+1} |l_k^{(i-1)}(t)| \leq B_i.$$

Then (35) follows from (34). Q.E.D.

7. A-Stability of the Galerkin Methods. Dahlquist [8] defines A-stability as follows.

Definition. A k-step method is called A-stable, if all its solutions tend to zero, as $i \to \infty$, when the method is applied with fixed positive h to any differential equation of the form $u' = \lambda u$, where λ is a complex constant with negative real part.

Butcher's implicit Runge-Kutta methods based on Gauss-Legendre quadrature [2] have been shown by Ehle [10] to be A-stable. Ehle observed that the n-stage method, applied to $u' = \lambda u$, yields $y_{i+1} = P_{nn}(\lambda h)y_i$, where $P_{nn}(\lambda h)$ is the nth diagonal Padé rational approximation to exp (λh) . A-stability follows from the fact that $|P_{nn}(\lambda h)| < 1$ for Re $(\lambda h) < 0$. Our Galerkin methods, which from (24) give discrete values y_i identical to those of Butcher's methods [2], are therefore A-stable.

We should remark that Axelsson [1] has used similar properties of subdiagonal and diagonal Padé rational approximations to prove A-stability for implicit Runge-Kutta methods based on Radau and Lobatto quadratures. It is natural then to ask whether a Galerkin method which uses these quadratures rather than Gauss-Legendre would yield corresponding "equivalent" methods. The answer is no. If (6) were an n-point Radau formula with $\sigma_{i,n} = t_{i+1}$, it would be exact only for $v \in \mathcal{O}_{2n-2}$. The quadratures for $(f, \varphi_{i+k})_i$ would not be exact for $f \in \mathcal{O}_{n-1}$, y would not be exact for $u \in \mathcal{O}_n$, (24) would not hold, and the order of the Galerkin method would be $O(h^{n-1})$, whereas Axelsson [1] and Butcher [3] have shown that an n-stage implicit Runge-Kutta method based on Radau quadrature has the order $O(h^{2n-1})$. Similar results are true of Lobatto quadrature;

8. Numerical Examples. In this section, we give numerical results of an A-stable piecewise cubic (n = 3) Galerkin scheme of order 6. We have employed Schoenberg's [20] cubic B-spline basis functions where φ_{i+j} has its support on $[t_{i+j-4}, t_{i+j}]$. The calculations were performed on a CDC 6600, which has about 14 decimal digits, using a successive substitution iteration at each step to solve (12).

First, we consider problems for single equations.

Problem 1.
$$u' = -2tu^2$$
, $u(0) = 1$, $u(t) = 1/(1 + t^2)$, $0 \le t \le 1$.
Problem 2. $u' = 1/(1 + \tan^2 u)$, $u(0) = 0$, $u(t) = \arctan t$, $0 \le t \le 1$.

Problem 3. u' = u - (2t/u), u(0) = 1, $u(t) = (2t+1)^{1/2}$, $0 \le t \le 1$. Problem 4. u' = u, u(0) = 1, $u(t) = e^t$, $0 \le t \le 10$. Several uniform meshes are used for each problem. Tables 1-4 are designed to

Table 1
Error Norms for Problem 1

e(t; h) '	e'(t;h) '	e''(t;h) '	e'''(t;h) '
55(10) ⁻⁴	7.94(10)-2	8.20(10)-1	3.49(10)°
			$4.10(10)^{\circ}(-0.23)$
		. , , ,	$2.71(10)^{\circ} (0.60)$
		` , , ,	$1.46(10)^{\circ}$ (0.89)
		• • • • •	$7.45(10)^{-1}(0.97)$
		` ` ` ` ` `	$3.74(10)^{-1}(0.99)$
			$1.87(10)^{-1}(1.00)$
	54(10) ⁻⁶ (5.38) 18(10) ⁻⁷ (6.18) 79(10) ⁻⁹ (6.04) 81(10) ⁻¹¹ (6.00) 45(10) ⁻¹³ (6.35)	7.94(10) ⁻² 55(10) ⁻⁴ 7.94(10) ⁻² 54(10) ⁻⁶ (5.38) 1.30(10) ⁻² (2.61) 18(10) ⁻⁷ (6.18) 2.65(10) ⁻³ (2.30) 79(10) ⁻⁹ (6.04) 3.75(10) ⁻⁴ (2.82) 81(10) ⁻¹¹ (6.00) 4.83(10) ⁻⁵ (2.96) 45(10) ⁻¹³ (6.35) 6.09(10) ⁻⁶ (2.99) 85(10) ⁻¹³ (0.90) 7.62(10) ⁻⁷ (3.00)	$54(10)^{-6}$ (5.38) $1.30(10)^{-2}$ (2.61) $3.36(10)^{-1}$ (1.29) $18(10)^{-7}$ (6.18) $2.65(10)^{-3}$ (2.30) $1.29(10)^{-1}$ (1.38) $79(10)^{-6}$ (6.04) $3.75(10)^{-4}$ (2.82) $3.61(10)^{-2}$ (1.84) $81(10)^{-11}$ (6.00) $4.83(10)^{-5}$ (2.96) $9.29(10)^{-3}$ (1.96) $45(10)^{-13}$ (6.35) $6.09(10)^{-6}$ (2.99) $2.34(10)^{-3}$ (1.99)

TABLE 2
Error Norms for Problem 2

h	e(t; h)	l'	e'(t;h) '	$ e^{\prime\prime}(t;h) ^{\prime}$	$ e^{\prime\prime\prime}(t;h) ^{\prime}$
$2^{-1} 2^{-2} 2^{-3} 2^{-4} 2^{-5}$	2.48(10) ⁻⁵ 1.28(10) ⁻⁷ 5.79(10) ⁻⁹ 8.62(10) ⁻¹¹ 1.34(10) ⁻¹² 9.24(10) ⁻¹⁴ 1.49(10) ⁻¹³ (40)	(4.46) (6.07) (6.01) (3.86)	4.76(10) ⁻³ (2.67) 5.88(10) ⁻⁴ (3.02) 7.64(10) ⁻⁵ (2.94) 9.48(10) ⁻⁶ (3.01) 1.19(10) ⁻⁶ (2.99)	3.62(10) ⁻¹ 1.11(10) ⁻¹ (1.70) 2.84(10) ⁻² (1.97) 7.30(10) ⁻³ (1.96) 1.82(10) ⁻³ (2.01) 4.56(10) ⁻⁴ (1.99) 1.14(10) ⁻⁴ (2.00)	6.61(10) ⁻¹ (1.29) 5.70(10) ⁻¹ (0.21) 2.85(10) ⁻¹ (1.00) 1.46(10) ⁻¹ (0.97) 7.29(10) ⁻² (1.00)

Table 3

Error Norms for Problem 3

h	e(t; h) '	e'(t;h) '	$ e^{\prime\prime}(t;h) ^{\prime}$	e'''(t; h) '
2^{-1} 2^{-2} 2^{-3} 2^{-4} 2^{-5}	7.08(10) ⁻⁴ 2.22(10) ⁻⁵ (4.99) 4.67(10) ⁻⁷ (5.57) 8.05(10) ⁻⁹ (5.86) 1.30(10) ⁻¹⁰ (5.96) 2.73(10) ⁻¹² (5.57) 1.48(10) ⁻¹² (0.88)	6.03(10) ⁻³ (2.09) 1.14(10) ⁻³ (2.40) 1.83(10) ⁻⁴ (2.64) 2.63(10) ⁻⁵ (2.80) 3.53(10) ⁻⁶ (2.89)	$3.24(10)^{-1}$ $1.50(10)^{-1}(1.11)$ $5.58(10)^{-2}(1.42)$ $1.77(10)^{-2}(1.65)$ $5.07(10)^{-3}(1.81)$ $1.36(10)^{-3}(1.90)$ $3.53(10)^{-4}(1.95)$	2.43(10)° 1.87(10)° (0.37) 1.26(10)° (0.57) 7.55(10)⁻¹ (0.74) 4.19(10)⁻¹ (0.85) 2.21(10)⁻¹ (0.92) 1.14(10)⁻¹ (0.96)

Table 4				
Error Norms for Problem	4			

h	e(t; h) '	e(t; h) ' $ e'(t; h) '$	$ e^{\prime\prime}(t;h) ^{\prime}$	$ e^{\prime\prime\prime}(t;h) ^{\prime}$
$2^{-1} 2^{-2} 2^{-3} 2^{-4} 2^{-5}$	5.35(10) ⁻⁴ (6.01) 8.31(10) ⁻⁶ (6.01) 9.32(10) ⁻⁸ (6.48)	$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	1.48(10) ³ 4.49(10) ² (1.72) 1.24(10) ² (1.86) 3.27(10) ¹ (1.93) 8.38(10) ⁰ (1.96) 2.12(10) ⁰ (1.98) 5.34(10) ⁻¹ (1.99)	5.59(10) ³ 3.90(10) ³ (0.52) 2.31(10) ³ (0.75) 1.26(10) ³ (0.88) 6.59(10) ² (0.94) 3.37(10) ² (0.97) 1.70(10) ² (0.98)

illustrate the $O(h^6)$ mesh point accuracy of Theorem 1 as well as the $O(h^3)$, $O(h^2)$ and O(h) accuracies of the first three derivatives predicted by Theorem 2. The tables give the discrete error norms for y(t; h) and its first three derivatives

(37)
$$||e^{(i)}(t;h)||' = \max_{0 \le i \le N} |e^{(i)}(t_{i+1};h)|, \qquad 0 \le j \le 3,$$

where e = u - y and also in parentheses the *computed orders of accuracy*, based on successive mesh sizes h_1 , h_2 ,

(38)
$$\omega_i = \frac{\log[||e^{(i)}(t; h_1)||'/||e^{(i)}(t; h_2)||']}{\log(h_1/h_2)},$$

i.e., $||e^{(i)}(t; h)||' \approx O(h^{\omega_i}), 0 \le j \le 3$.

Next, we present in Table 5 absolute errors e(t; h) and relative errors e(t; h)/u(t)

Table 5

Absolute and Relative Errors for Problem 5

t	e(t; 1)	e(t; 1)/u(t)	e(t; 0.5)	e(t; 0.5)/u(t)
1	3.79(10)-6	1.03(10)-5	5.76(10) ⁻⁸	1.57(10) ⁻⁷
10	$4.68(10)^{-9}$	$1.03(10)^{-4}$	$7.11(10)^{-11}$	$1.57(10)^{-6}$
20	$4.25(10)^{-13}$	$2.06(10)^{-4}$	$6.45(10)^{-15}$	$3.13(10)^{-6}$
40	$1.75(10)^{-21}$	4.12(10)-4	$2.67(10)^{-23}$	$6.26(10)^{-6}$
60	$5.42(10)^{-30}$	6.19(10)-4	$8.22(10)^{-32}$	$9.39(10)^{-6}$
80	$1.49(10)^{-38}$	$8.25(10)^{-4}$	$2.26(10)^{-40}$	$1.25(10)^{-5}$
100	$3.83(10)^{-47}$	$1.03(10)^{-3}$	$5.83(10)^{-49}$	$1.57(10)^{-5}$

at selected points t_i for h = 1 and 0.5 in

Problem 5. u' = -u, u(0) = 1, $u(t) = e^{-t}$, $0 \le t \le 100$, in order to illustrate the stability of the method.

Finally, we give in Tables 6 and 7 the results of the application of our method to

Table 6						
Error Norms for e ₁ (t; h) of Problem 6	6					

h	$ e_1(t;h) $	'	$ e_i'(t;h) '$	$ e_1^{\prime\prime}(t;h) ^{\prime}$	$ e_1^{\prime\prime\prime}(t;h) ^{\prime}$
2^{-1} 2^{-2} 2^{-3} 2^{-4} 2^{-5}	4.50(10) ⁻⁷ 7.01(10) ⁻⁹ 1.08(10) ⁻¹⁰ 2.19(10) ⁻¹²	(6.08) (6.02) (6.00) (6.02) (5.62)	2.53(10) ⁻³ (2.85) 3.33(10) ⁻⁴ (2.92) 4.29(10) ⁻⁵ (2.96) 5.45(10) ⁻⁶ (2.98) 6.86(10) ⁻⁷ (2.99)	1.77(10) ⁻¹ 5.43(10) ⁻² (1.71) 1.52(10) ⁻² (1.84) 4.01(10) ⁻³ (1.92) 1.03(10) ⁻³ (1.96) 2.62(10) ⁻⁴ (1.98) 6.59(10) ⁻⁵ (1.99)	4.95(10) ⁻¹ (0.59) 2.89(10) ⁻¹ (0.77) 1.57(10) ⁻¹ (0.88) 8.16(10) ⁻² (0.94) 4.16(10) ⁻² (0.97)

Table 7

Error Norms for $e_2(t; h)$ of Problem 6

h	$ e_2(t;h) $) ' 	$ e_2'(t;h) '$	$ e_2^{\prime\prime}(t;h) ^{\prime}$	$ e_2^{\prime\prime\prime}(t;h) ^{\prime}$
$2^{-1} 2^{-2} 2^{-3} 2^{-4} 2^{-5}$	2.58(10) ⁻⁴ 3.82(10) ⁻⁶ 5.91(10) ⁻⁸ 9.20(10) ⁻¹⁰ 1.42(10) ⁻¹¹ 2.13(10) ⁻¹³ 7.27(10) ⁻¹³ (-	(6.02) (6.00) (6.02) (6.06)	9.40(10) ⁻⁴ (2.95) 1.23(10) ⁻⁴ (2.93) 1.58(10) ⁻⁵ (2.96) 2.00(10) ⁻⁶ (2.98) 2.52(10) ⁻⁷ (2.99)	7.69(10) ⁻² 2.17(10) ⁻² (1.82) 5.82(10) ⁻³ (1.90) 1.51(10) ⁻³ (1.95) 3.84(10) ⁻⁴ (1.97) 9.68(10) ⁻⁵ (1.99) 2.43(10) ⁻⁵ (1.99)	2.21(10) ⁻¹ (0.84) 1.18(10) ⁻¹ (0.91) 6.06(10) ⁻² (0.96) 3.08(10) ⁻² (0.98) 1.55(10) ⁻² (0.99)

the system of equations in

Problem 6. $u'_1 = u_1^2 u_2$, $u'_2 = -1/u_1$, $u_1(0) = 1$, $u_2(0) = 1$, $u_1 = e^t u_2 = e^{-t}$, $0 \le t \le 1$.

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