

Onsager-Machlup Theory for Nonequilibrium Steady States and Fluctuation Theorems

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1. Introduction

1.1. Past : Fluctuation Theories in Nonequilibrium Physics

- **Fluctuation-Dissipation Theorem**

(Near equilibrium: Transport coefficients and fluctuation correlations)

- **Onsager-Machlup Fluctuation Theory**

(Near equilibrium: Relaxation process → average decay of a fluctuation away from equilibrium follows linear macroscopic law; Onsager's principle of minimum energy dissipation; the most probable path)

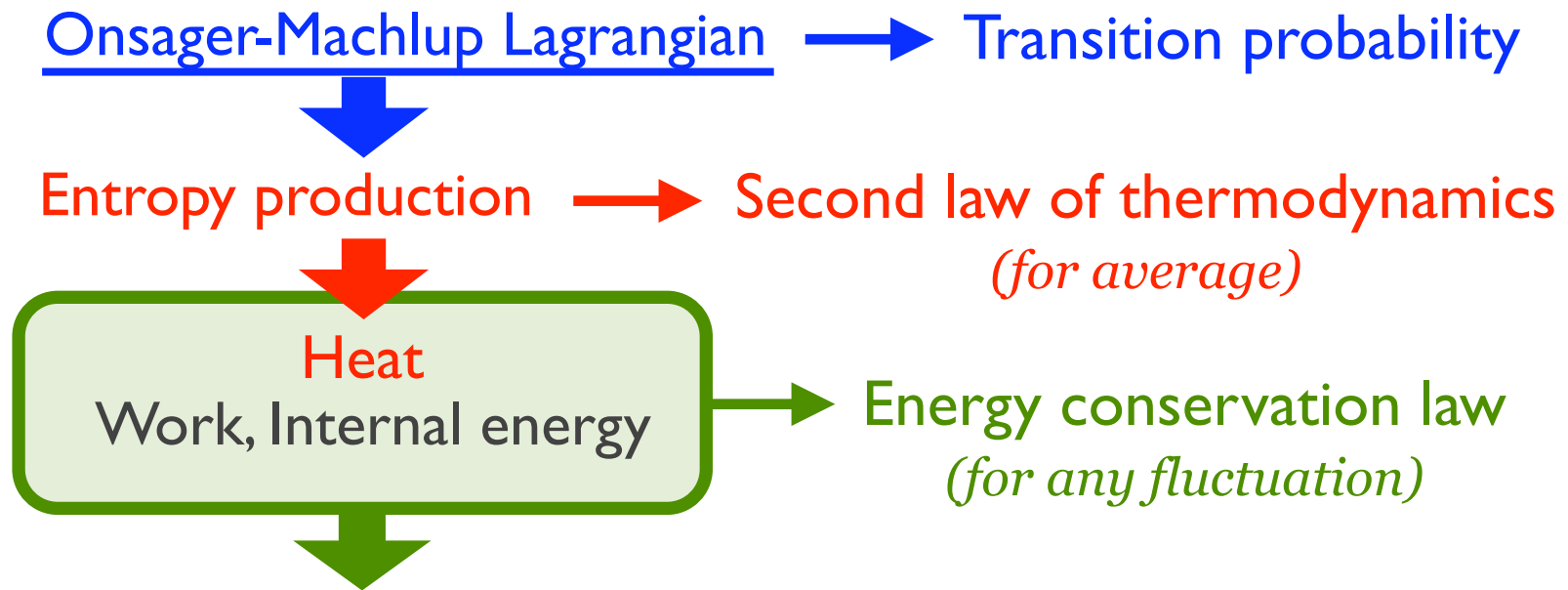
- **Fluctuation Theorems**

(“Far” from equilibrium: Asymmetric property of probability distribution functions of fluctuations)

?

1.2. New : Contents of This Talk

1. Generalization of Onsager-Machlup Fluctuation Theory to Nonequilibrium Steady States



2. Fluctuation Theorems by a Functional Integral Approach

- Nonequilibrium detailed balance relations
- Fluctuation theorems for work and friction
- Extended fluctuation theorem for heat



1.3. Model: Dragged Brownian Particle

(for a nonequilibrium steady state)

- Langevin Equation

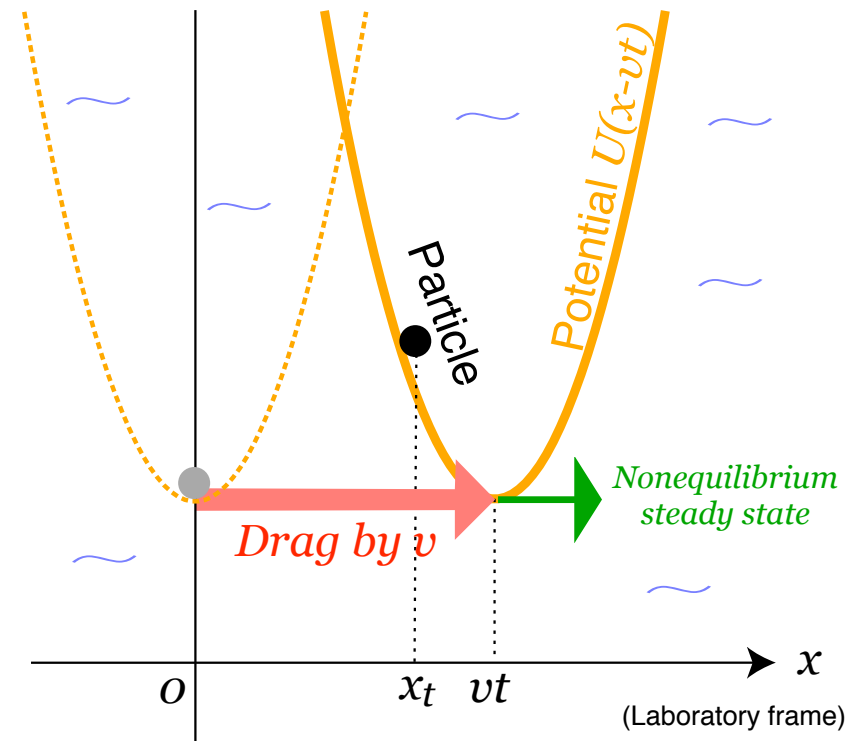
$$m \frac{d^2 x_t}{dt^2} = \underbrace{-\alpha \frac{dx_t}{dt}}_{\text{Friction force}} - \underbrace{\kappa (x_t - vt)}_{\text{Harmonic force}} + \underbrace{\zeta_t}_{\text{Gaussian-white noise}}$$

Harmonic force

$$U(x) = \frac{1}{2} \kappa x^2$$

Gaussian-white noise

$$\begin{aligned} \langle \zeta_t \rangle &= 0, \\ \langle \zeta_{t_1} \zeta_{t_2} \rangle &= \frac{2\alpha}{\beta} \delta(t_1 - t_2) \end{aligned}$$



- **Over-Damped Assumption**

Neglect inertial effect: $m d^2x_t/dt^2 \approx 0$ (or simply $m \approx 0$)

$$\frac{dx_t}{dt} = -\frac{1}{\tau_r} (x_t - vt) + \frac{1}{\alpha} \zeta_t, \quad \tau_r \equiv \frac{\alpha}{\kappa}$$

- **Comoving Frame y : Frame Moving with Velocity v**

$$y_t \equiv x_t - vt$$

$$\frac{dy_t}{dt} = -\frac{1}{\tau_r} y_t - v + \frac{1}{\alpha} \zeta_t$$

Nonequilibrium effect

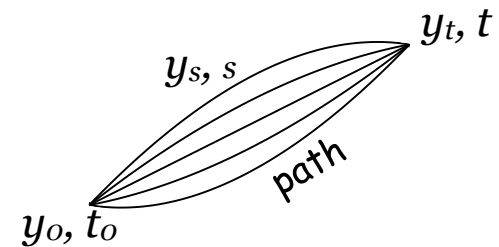
Laboratory experiments (e.g. a Brownian particle captured in an optical trap, or an electric circuit consisting of a resistor and capacitor, etc.)

2. Onsager-Machlup Theory for Nonequilibrium Steady States

- Transition Probability of Particle Position in Time

$$F \left(\begin{matrix} y_t \\ t \end{matrix} \middle| \begin{matrix} y_0 \\ t_0 \end{matrix} \right) = \int_{y_0}^{y_t} \mathcal{D}y_s \exp \left[\int_{t_0}^t ds L^{(v)}(\dot{y}_s, y_s) \right]$$

Functional integral
Probability functional for a path $\{y_s\}$



- Onsager-Machlup Lagrangian function

$$L^{(v)}(\dot{y}_s, y_s) \equiv -\frac{1}{4D} \left(\dot{y}_s + v + \frac{y_s}{\tau_r} \right)^2$$

$$D \equiv \frac{k_B T}{\alpha} \quad \text{Einstein relation}$$

$$= -\frac{1}{2k_B} \left[\frac{\alpha}{2T} (\dot{y}_s + v)^2 + \frac{\alpha}{2T} \left(\frac{y_s}{\tau_r} \right)^2 - \dot{\mathcal{S}}^{(v)}(\dot{y}_s, y_s) \right]$$

- Connection with Thermodynamics

$$\dot{\mathcal{S}}^{(v)}(\dot{y}_s, y_s) \equiv -\frac{1}{T} \kappa y_s (\dot{y}_s + v) \quad \text{Entropy production rate, because:}$$

[i] Second Law of Thermodynamics (*holds for average*)

$$\dot{S}^{(v)}(\langle \dot{y}_t \rangle, \langle y_t \rangle) \geq 0$$

[ii] Energy Conservation Law (*holds for any fluctuation*)

$$Q_t(\{y_s\}) = \mathcal{W}_t^{(v)}(\{y_s\}) - \Delta U(y_t, y_0)$$

Nonequilibrium effect
(zero in equilibrium: $v=0$)

Heat

$$Q_t(\{y_s\}) \equiv T \int_{t_0}^t ds \dot{S}^{(v)}(\dot{y}_s, y_s)$$

Functional

Entropy production rate

Work

$$\mathcal{W}_t^{(v)}(\{y_s\}) \equiv \int_{t_0}^t ds (-\kappa y_s) v$$

Harmonic Force

Internal Energy
Difference

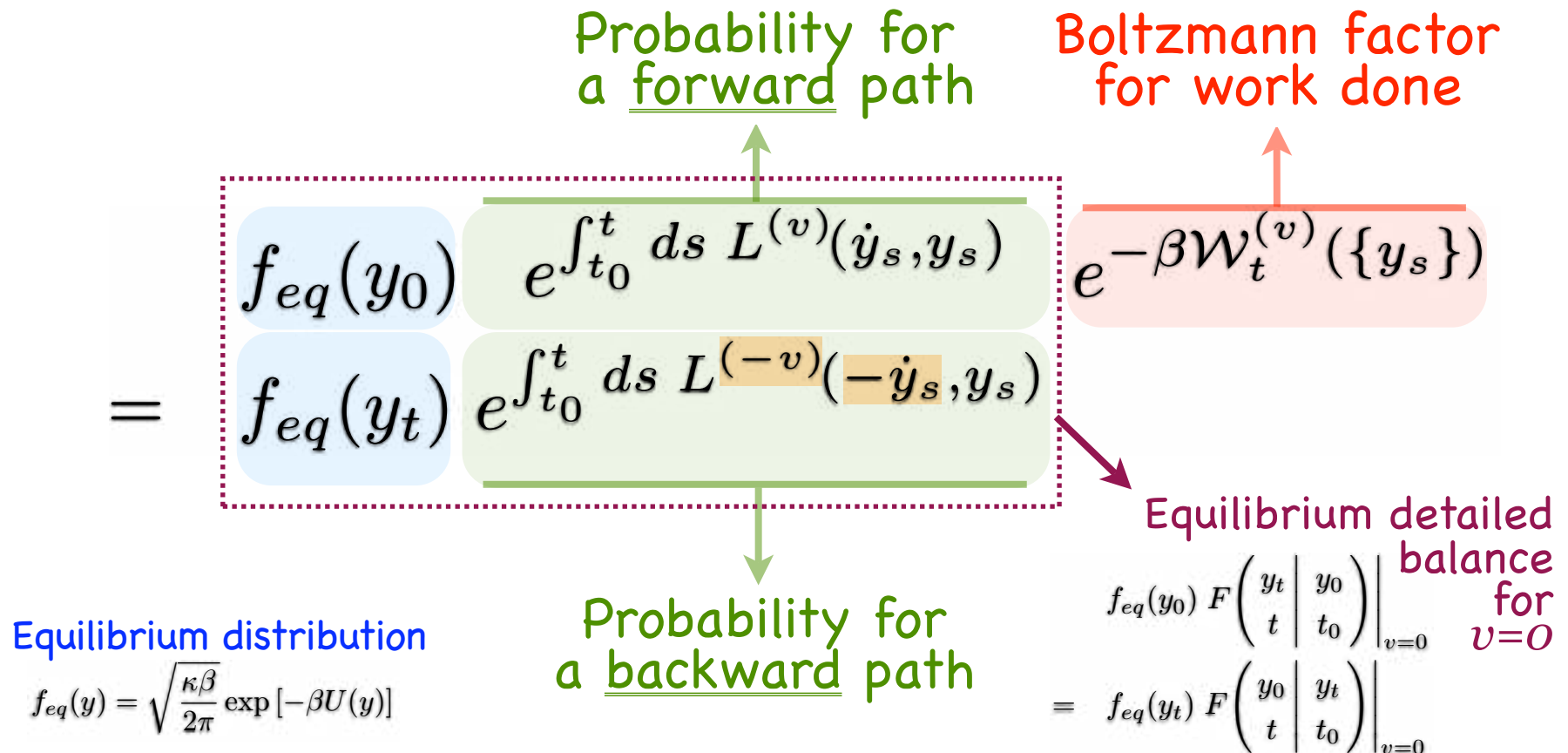
$$\Delta U(y_t, y_0) \equiv U(y_t) - U(y_0)$$

Potential $U(y) = \frac{1}{2} \kappa y^2$

3. Fluctuation Theorems

- Nonequilibrium Detailed Balance (I)

[Due to non-zero v : nonequilibrium effect]



3.1. Fluctuation Theorem for Work

- Distribution function of (dimensionless) work

$$P_w(W, t) = \left\langle\left\langle \delta \left(W - \beta \mathcal{W}_t^{(v)}(\{y_s\}) \right) \right\rangle\right\rangle_t$$

Functional average

$$\left\langle\left\langle X(\{y_s\}) \right\rangle\right\rangle_t \equiv \int dy_t \int_{y_0}^{y_t} \mathcal{D}y_s \int dy_0 e^{\int_{t_0}^t ds L^{(v)}(\dot{y}_s, y_s)} f(y_0, t_0) X(\{y_s\})$$

Distribution of y_0 at t_0

- Work fluctuation theorem

$$\lim_{t \rightarrow +\infty} \frac{P_w(W, t)}{P_w(-W, t)} = \exp(W)$$

for any $f(y_0, t_0)$

Nonequilibrium
detailed balance

3.2. Fluctuation Theorem for Friction

- Energy loss by friction

$$\mathcal{R}_t^{(v)}(y_t, y_0) = \int_{t_0}^t ds \overbrace{(-\alpha \dot{y}_s)}^{\text{Friction force}} v$$

Boltzmann factor
for energy loss
by friction

- Nonequilibrium detailed balance (II)

$$f_{eq}(y_0) e^{\int_{t_0}^t ds L^{(v)}(\dot{y}_s, y_s)} e^{-\beta \mathcal{R}_t^{(v)}(y_t, y_0)} = f_{eq}(y_t) e^{\int_{t_0}^t ds L^{(v)}(-\dot{y}_s, y_s)}$$

No change of
the sign of v

- Distribution function of (dimensionless) friction

$$P_r(R, t) = \left\langle \left\langle \delta \left(R - \beta \mathcal{R}_t^{(v)}(y_t, y_0) \right) \right\rangle \right\rangle_t$$

- Friction fluctuation theorem

$$\frac{P_r(R, t)}{P_r(-R, t)} = \exp(R) \quad \text{for} \quad f(y_0, t_0) = f_{eq}(y_0)$$

(The friction FT is not correct for the steady state initial condition even for $t \rightarrow \infty$.)

3.3. Extended Fluctuation Theorem for Heat (in the long time limit)

- Heat distribution function

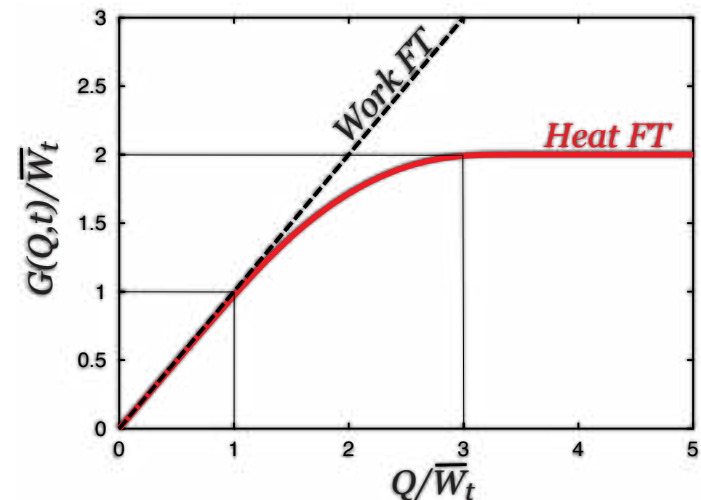
$$\begin{aligned}
 P_q(Q, t) &\stackrel{t \rightarrow +\infty}{\sim} \int dW \int d\Delta U \underbrace{P_w(W, t)}_{\text{Work distribution (Gaussian)}} \underbrace{P_{\Delta e}(\Delta U, t)}_{\text{Energy-difference distribution (Exponential)}} \underbrace{\delta(Q - W + \Delta U)}_{\text{Energy conservation}} \\
 &= \left[e^{-Q+2\bar{W}_t} \text{erfc} \left(-\frac{Q - 3\bar{W}_t}{2\sqrt{\bar{W}_t}} \right) + e^Q \text{erfc} \left(\frac{Q + \bar{W}_t}{2\sqrt{\bar{W}_t}} \right) \right] \\
 &\quad \text{Exponential tail}
 \end{aligned}$$

$$\bar{W}_t \stackrel{t \rightarrow +\infty}{\sim} \alpha \beta v^2 t, \quad \text{erfc}(x) \equiv (2/\sqrt{\pi}) \int_x^{+\infty} dz \exp(-z^2)$$

- Heat fluctuation theorem (scaled)

$$G(Q, t) \equiv \ln \frac{P_q(Q, t)}{P_q(-Q, t)}$$

Experimental check of the heat FT using an electric circuit (Garnier and Ciliberto, 2005)



4. Conclusion

- **Generalization of Onsager-Machlup Theory to Nonequilibrium Steady States**

- Thermodynamics and fluctuations from the Onsager-Machlup Lagrangian function (the second law of thermodynamics, the energy conservation law, etc.)

- **Fluctuation Theorems using a Functional Integral Approach**

- Usage of nonequilibrium detailed balance relations for derivation of fluctuation theorems for work and friction
- Simple argument for the extended fluctuation theorem for heat



Onsager-Machlup
Lagrangian function

Reference: T. Taniguchi and E. G. D. Cohen, e-print cond-mat/0605548

Appendix: Notations in This Talk

- m : mass
- α : friction coefficient
- κ : spring constant
- T : temperature
- k_B : Boltzmann constant
- $\beta=1/(k_B T)$: inverse temperature
- v : velocity to drag the particle
- $\tau_r = \alpha/\kappa$: relaxation time
- ζ_t : Gaussian white random force
- $\langle \dots \rangle$: ensemble average
- $D=1/(\alpha\beta)$: diffusion constant
- $U(y)$: harmonic potential
- $L^{(v)}(\dot{y}, y)$: Lagrangian function
- $\dot{S}^{(v)}(\dot{y}, y)$: entropy production rate
- Q : heat
- W : work
- ΔU : internal energy difference
- $f_{eq}(y)$: equilibrium distribution function
- $f(y, t)$: distribution of position y at time t
- $P_w(W, t)$: work distribution
- $P_r(R, t)$: friction distribution
- $P_q(Q, t)$: heat distribution
- $\langle\langle \dots \rangle\rangle_t$: functional average