

Optical Frequency Comb Generation Using CMOS Compatible Cascaded Mach-Zehnder Modulators

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Abstract—Optical frequency combs (OFCs) play a key role in a variety of applications and have been implemented primarily using mode-locked lasers, Kerr resonators, and electro-optic (EO) modulation. EO modulation-based OFC generation may not yield as many comb lines nor span as broad a bandwidth as mode-locked lasers or Kerr resonators; however, it offers a high degree of tunability in central frequency and comb spacing. Integrated solutions are of interest as they significantly reduce device footprint and enable large scale system integration. In this paper, we demonstrate flexible on-chip OFC generation using two cascaded EO Mach-Zehnder modulators in silicon photonics. We demonstrate quasi-rectangular OFCs with 9 lines and a comb spacing of up to 10 GHz with an amplitude variation (comb flatness) within 6.5 dB. The corresponding time-domain waveforms have a good fit with sinc-shaped Nyquist pulses having a full-width at half maximum duration as short as 11.4 ps.

Index Terms—Optical frequency comb generation, electro-optic Mach-Zehnder modulators, silicon photonics.

I. INTRODUCTION

OPTICAL frequency combs (OFCs), which consist of a series of equidistant spectral lines with correlated phase [1], have been used extensively in a diverse range of applications including spectroscopy [2, 3], optical frequency metrology [4], precision distance measurement [5], optical and microwave waveform synthesis [6-9], astronomical spectrograph calibration [10, 11], and optical communications [12]. Techniques for OFC generation include, amongst others, the use of mode-lock lasers assisted by an internal or external reference [13-20], Kerr resonators [21-30], and electro-optic (EO) modulation [31-37]. The properties of mode-locked lasers, such as repetition rate, amplitude noise, linewidth, and frequency stability determine the characteristics of the OFC and need to be considered carefully for different applications [18]. By exploiting parametric four-wave mixing, OFC generation based on Kerr resonators may offer octave-spanning, ultrawide bandwidth from a single continuous wave (CW) input, as well

as high frequency and timing stability, which are desirable for telecommunications, spectroscopy, low-noise microwave waveform synthesis, and distance ranging [27]. EO modulation-based OFC generation offers the flexibility to tune readily the center frequency and comb spacing simply by adjusting the input CW wavelength as well as frequency, power, and phase of the RF signal(s) applied to the EO modulators. Cascading a larger number of modulators (e.g., ≥ 3) is often used to increase the number of comb lines and obtain a broadband comb [8], though it may not yield as many comb lines or span as broad a bandwidth as mode-locked lasers or Kerr resonators. In some applications, e.g., the generation of sinc-shaped Nyquist pulses for Nyquist orthogonal time-division multiplexed transmission, the shape of the comb spectrum is more important than obtaining a large number of comb lines. In [35], two modulators are cascaded to produce a rectangular shaped, phase-locked comb with 9 to 10 spectral lines occupying more than 100 GHz bandwidth and corresponding high quality Nyquist optical pulses.

For still other applications in optical communications and microwave photonics, a high degree of system level integration may be more desirable over a large number of comb lines. To date, integrated OFC generators have been reported using various technology platforms such as silica [21, 22], Hydex glass [23], silicon nitride (SiN) [24-27], III-V semiconductors [28, 29], silicon-organic hybrids [36], and silicon-on-insulator (SOI) [32, 37, 38]. Exploiting recent advances in developing high-performance EO modulators in SOI has the potential to yield a highly-tunable integrated comb source.

In this paper, we demonstrate the generation of an EO modulation-based OFC using integrated cascaded push-pull traveling-wave Mach-Zehnder modulators (MZMs) in SOI. By precisely controlling the amplitude and phase of the driving RF signals as well as the bias and thermal heating voltages, a quasi-rectangular OFC with 9 lines, flatness within 6.5 dB, and bandwidths up to 90 GHz have been obtained; the corresponding sinc-shaped Nyquist pulses have a full-width at

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half-maximum (FWHM) duration as short as 11.4 ps.

II. DEVICE DESIGN

Figures 1 (a) and (b) show a schematic of the two cascaded EO MZMs in SOI and a photo of the fabricated device. The modulators are based on the series push-pull design described in [39]; some of the parameters are reviewed here. The devices are fabricated on a 220 nm SOI wafer with a $750 \Omega - cm$ silicon substrate using a CMOS-compatible foundry process available at IME A*STAR. Rib waveguides are doped using three doping levels to form diode waveguides; their cross-sections are shown in Fig. 1(c). The waveguides sit on top of a $3 \mu m$ thick oxide (BOX) and are covered by a $2 \mu m$ thick index matched oxide cladding. The waveguide cross-sections and doping parameters are given in Table 1.

Each modulator has a path imbalance of $130 \mu m$, which introduces an additional phase shift in the lower arm of the MZ interferometer and creates a periodic spectral response. This path imbalance is used to measure the phase shifts associated with bias and thermal tuning, as described below. An n++ doped heater with a length of $215 \mu m$ is overlaid on the upper arm for fine-tuning the phase. The length of each modulator is 4.55 mm and the net p-n junction loading is 4.2 mm. The two modulators are connected in series using 2×2 MMI couplers [40] and in order to characterize each modulator separately, additional taps are used. Vertical grating couplers (VGCs) designed for TE mode operation over the C band [41] serve for input and output coupling and testing. The total size of the device is $1.49 \text{ mm} \times 5.3 \text{ mm}$.

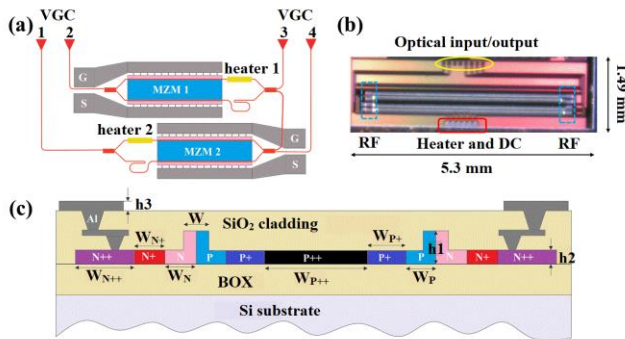


Fig. 1. Integrated OFC generator in SOI: (a) device schematic, (b) photo of fabricated device, and (c) cross-section of push-pull MZM.

TABLE I
PARAMETERS OF THE PUSH-PULL EO MZMS IN SOI

h1	220 nm
h2	90 nm
h3	$2 \mu m$
W	500 nm
W_{N++}	$5.2 \mu m$
W_{N+}	$0.81 \mu m$
W_N	$0.39 \mu m$
W_{P++}	$28 \mu m$
W_{P+}	$0.83 \mu m$
W_P	$0.37 \mu m$

III. EXPERIMENTS AND RESULTS

A. Characterization of the MZMs

First, we characterize each MZM separately. Figure 2(a) shows the transmission spectra of MZM 1 and MZM 2 as a function of reverse (DC) bias voltage. The fiber-to-fiber loss (i.e., static loss in the absence of any applied RF signal) depends on operating wavelength and bias/heating voltages. When measuring the loss of each MZM, we use a wavelength that matches with a peak in the corresponding spectral response without applied voltages. MZM 1 has a fiber-to-fiber insertion loss (i.e., from VGC2 to VGC3) of 18 dB at a wavelength of 1548 nm; this comprises 11 dB coupling loss, 4 dB loss from the MMI, and 3 dB loss from the MZM structure itself. MZM 2 also has a fiber-to-fiber insertion loss (i.e., from VGC4 to VGC1) of 18 dB, but at a wavelength of 1550 nm. The total fiber-to-fiber insertion loss of the cascaded MZMs (i.e., from VGC2 to VGC1) is 24 dB at a wavelength of 1550 nm. When no bias voltage is applied, the extinction ratios at a wavelength of 1550 nm for MZM 1 and MZM 2 are 30 dB and 19.5 dB, respectively. By increasing the bias voltage, the difference in carrier concentration between the two arms of the MZM increases, which induces a higher optical loss imbalance and the degradation in extinction ratio. Both MZMs have a free spectral range of 4.4 nm, which corresponds to the designed path imbalance of $130 \mu m$. The half-wave voltages of both MZMs are $V_\pi = 10 \text{ V}$, indicating $V_\pi L = 4.2 \text{ V} - cm$. The travelling-wave MZM design results in a relatively high value of V_π ; in order to reduce V_π , the doping length must be longer, which will increase the propagation loss and reduce bandwidth. Figure 2(b) shows the transmission spectrum of MZM 1 and MZM 2 as a function of heating voltage (i.e., voltage applied to heater 1 or heater 2). In this case, both MZMs attain a π shift for a heating voltage of 2.5 V. Therefore, we can use the heating voltage and DC bias for coarse and fine phase tuning of the modulator operation.

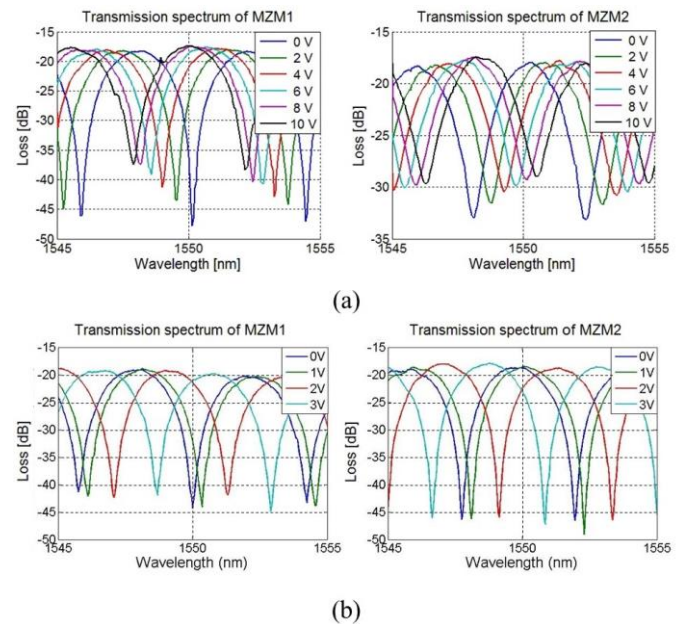


Fig. 2. Measured transmission spectra of MZM 1 and MZM 2 for various (a) reverse bias voltages and (b) heating voltages.

Figure 3 shows the measured S parameters of both MZMs. For a reverse bias of 3 V, the 3 dB (6 dB) bandwidths of MZM 1 and MZM 2 are 13.1 GHz (25 GHz) and 10.9 GHz (22.8 GHz), respectively. Note that the S_{11} value is below -10 dB for frequencies up to 50 GHz, indicating good 50 Ω matching.

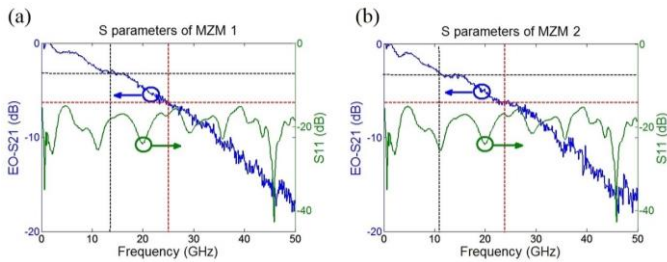


Fig. 3. Measured S_{11} and S_{21} parameters of (a) MZM 1 and (b) MZM 2.

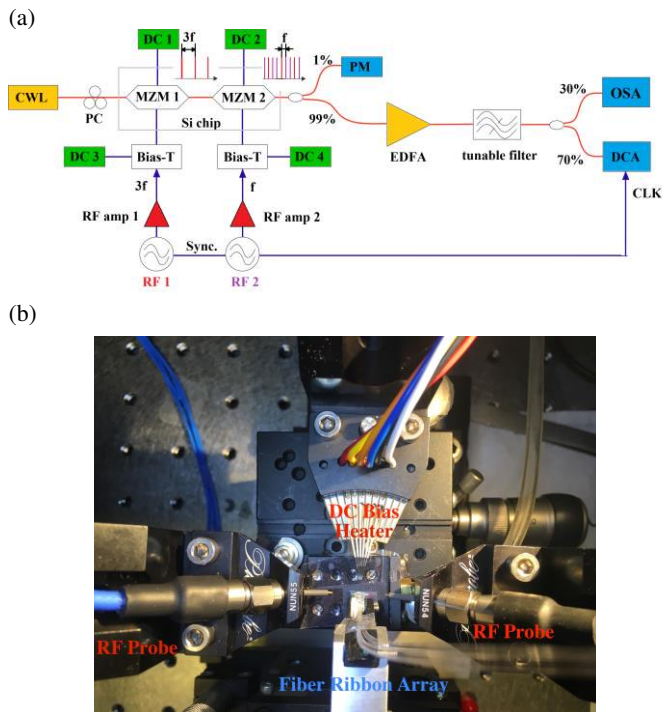


Fig. 4. (a) Schematic and (b) photo of experimental setup for on-chip OFC generation. CWL: continuous wave laser, RF: radio frequency, DC: direct current, PM: power meter, EDFA: erbium-doped fiber amplifier, OSA: optical spectrum analyzer, DCA: digital communications analyzer.

B. OFC generation

Figures 4(a) and 4(b) show a schematic and photo of our experimental setup for on-chip OFC generation. A CW laser at 1550.02 nm with a power of 14.2 dBm is used as a seed carrier for MZM 1. By tuning the amplitude of the RF signal (RF 1), as well as the DC reverse bias and heater voltages of MZM 1, we can obtain two first-order sidebands with approximately the same power as the seed carrier while suppressing the higher-order sidebands. These three spectral lines (which are separated by the frequency of RF 1) are then injected into MZM 2 (driven by a second RF signal, RF 2) as new seed carriers to generate additional first-order sidebands; the final result is a total of 9 comb lines. In order to generate combs with equal frequency spacing, the frequency of RF 2 should be 1/3 that of RF 1.

Moreover, the two RF signals should be synchronized properly by tuning their phases (especially for optimizing the quality of the sinc-shaped Nyquist pulses). The OFC at the output of the chip is amplified using an EDFA and filtered using a tunable filter (Finisar Waveshaper) to suppress out-of-band amplified spontaneous emission noise. The spectra and temporal waveforms are measured using an optical spectrum analyzer (OSA) with a resolution of 10 pm and an optical sampling module connected to a digital sampling oscilloscope with an impulse response time of 6.8 ps. The oscilloscope is operated in sample mode (no averaging) with a persistence time of 100 ms.

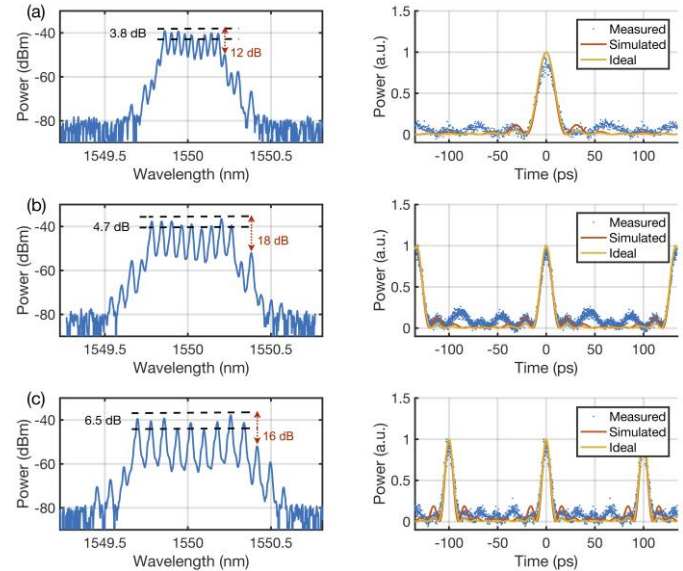


Fig. 5. OFC spectra at the output of the cascaded modulators (left) and corresponding temporal waveforms (right) for a comb spacing of (a) 5 GHz, (b) 7.5 GHz, and (c) 10 GHz. All temporal waveforms are obtained without averaging. The zero crossing pulse durations are 44.3 ps, 28.7 ps, and 25.2 ps, for (a), (b), and (c), respectively, and the corresponding FWHM pulse durations are 17.6 ps, 13.1 ps, and 11.4.

Figure 5 shows the spectra and the corresponding temporal waveforms of the OFC. In the frequency domain, we can control the flatness of the combs by adjusting the RF power and tuning the heater voltages of the MZMs (in our experiments, we set the DC reverse bias voltages to 0 V. The range of the heater voltages used is between 2.8 V to 3.8 V. The power of the RF signals after amplification (i.e., RF amp 1 and RF amp 2) are ~ 11 dBm and 14.5 dBm for RF 1 (higher frequencies) and RF 2 (lower frequencies), respectively. Once the 9 comb lines are generated, the temporal waveforms are optimized further by tuning the phases of the RF signals. The spectra shown in Fig. 5 are for comb spacings of 5 GHz, 7.5 GHz, and 10 GHz and are measured before the EDFA. The corresponding powers at the output of the chip are -28.6 dBm, -26.8 dBm, and -28.5 dBm and the quasi-rectangular shaped combs have a flatness within 3.8 dB, 4.7 dB, and 6.5 dB, and a sideband suppression of 12 dB, 18 dB, and 16 dB. Finally, the optical signal out-of-band noise for the comb lines is ~ 40 dB; future work includes transmission system experiments so that we can assess the performance of our generated OFCs for Nyquist orthogonal

time-division multiplexed and WDM transmission.

To obtain 3 comb lines with roughly equal power (e.g., from the first modulator), we need to suppress intentionally the power of the carrier to match that of the first-order sidebands. Since the 3 dB bandwidths of the two MZMs are in the range of 10.9 GHz to 13.1 GHz, it is easier to obtain 3 comb lines with less power variation at lower frequencies; on the other hand, when we increase the operating frequency, e.g., to 10 GHz which requires on RF signal at 30 GHz), achieving flatness over the 9 comb lines requires higher RF power and greater variation in the heater voltages. Each of the OFC spectra shown in Fig. 5 was optimized separately and this optimization process created the differences in comb flatness. The need to suppress the power of the carrier contributes to a higher dynamic insertion loss (i.e., when an RF signal is applied to the modulators) of ~ 43 dB (compared to the static insertion loss of 18 dB described earlier).

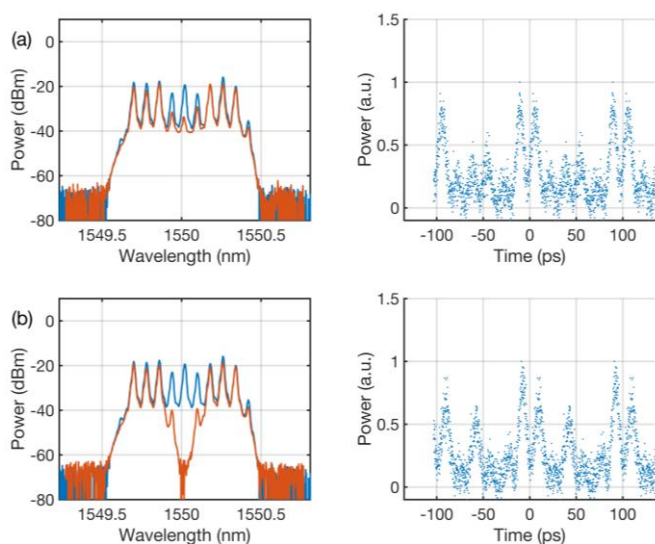


Fig. 6. Waveform shaping: spectra (left) and temporal waveforms (right) when the 3 central comb lines are removed using (a) fine tuning of the heater voltages and (b) a 'notch' filter.

The measured temporal waveforms are in excellent agreement with ideal sinc-shaped Nyquist pulses (assuming a perfect rectangular shaped OFC) and the simulated pulses (using the measured amplitudes of the comb lines). The sinc-shaped Nyquist pulses have FWHM durations of 17.6 ps, 13.1 ps, and 11.4 ps and the corresponding rms timing jitters are 1.2 ps, 1.7 ps, and 1.1 ps (all timing jitter measurements were made using the digital sampling oscilloscope with a precision time base and a persistence time of 100 ms). The RF synthesizers used in our experiments have rms timing jitters below 300 fs. The system (electrical and optical components such as RF amplifiers and EDFA) as well as operating parameters (e.g., synchronization of the applied RF signals to the two MZMs) contribute to an increase in the rms timing jitter of the output Nyquist pulses. The measured rms timing jitters are nevertheless representative of high-quality optical signals, including at 10 Gb/s. Note that due to some phase mismatching between the two RF signals, there is a small difference in the sidelobes between the measured and simulated waveforms; we

believe that we can minimize these sidelobes with further optimization of the RF phases.

We can also 'shape' the comb and corresponding temporal waveform by tuning the heater voltages or using a tunable optical filter. For example, in Fig. 6, we suppress the three middle comb lines by adjusting the heater voltages [6(a)] or using the Waveshaper to implement a notch-like filter [6(b)]. The ability to tailor the comb shape finely allows for the generation of more complex waveforms such as optical pulse bursts, triangular pulses, etc. [6, 42, 43].

IV. SUMMARY AND DISCUSSION

We have demonstrated on-chip generation of quasi-rectangular OFCs using two cascaded push-pull traveling-wave EO MZMs in SOI. We obtained OFCs with a comb spacing of up to 10 GHz (OFC bandwidth of 90 GHz); the flatness varies from 3.8 dB to 6.5 dB and the sideband suppression ranges from 12 dB to 16 dB. The corresponding temporal sinc-shaped Nyquist pulses have FWHM pulse durations as short as 11.4 ps and with rms timing jitters below 1.7 ps.

While OFCs generated using cascaded LiNbO₃ modulators (e.g., [35]) can achieve greater comb flatness (within 0.2 dB), higher sideband suppression (27 dB), and lower timing jitter (< 100 fs), our results demonstrate the potential for what can be obtained with on-chip comb generation and further improvements may be possible. For example, we can increase the available output power (i.e., reducing the dynamic insertion loss) by using VGCs with lower coupling loss [44] as well as having a higher MZM bandwidth and/or reducing V_{π} (both of these will reduce the amount of carrier suppression required to equalize the carrier and sideband power levels). MZM bandwidths of 41 GHz have been demonstrated [39]; thus, it should be possible to obtain 9 comb lines with a spacing of 15 GHz (corresponding to a total bandwidth of 135 GHz). While careful design may reduce V_{π} further, the typical values in SOI are larger compared to LiNbO₃ and as such, it may be difficult to obtain more than 9 lines. However, the integrated OFC generator can be used as a source for Nyquist orthogonal time-division multiplexed transmission [35,45,46] or flexible WDM transmission where flexibility in terms of comb spacing and modulation format is more important than a larger number of comb lines [37].

We believe that our on-chip OFC generator can be used as a building block for developing more advanced integrated subsystems for applications in optical communications, e.g., as an on-chip transmitter, and microwave photonics, e.g., microwave photonic filters and RF channelization.

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REFERENCES

- [1] S. T. Cundiff and J. Ye, "Colloquium: Femtosecond optical frequency combs," *Rev. Mod. Phys.*, vol. 75, no. 1, pp.325-342, Mar 2013.
- [2] M.-G. Suh, Q.-F. Yang, K. Y. Yang, X. Yi, and K. J. Vahala, "Microresonator soliton dual-comb spectroscopy," *Science*, vol. 354, no. 6312, pp. 600-603, Nov 2016.
- [3] C. Gohle, T. Udem, M. Herrmann, J. Rauschenberger, R. Holzwarth, H. A. Schuessler, F. Krausz, and T. W. Hänsch, "A frequency comb in the extreme ultraviolet," *Nature*, vol. 436, no. 7048, pp. 234-237, Jul 2005.
- [4] T. Udem, R. Holzwarth, and T. W. Hänsch, "Optical frequency metrology," *Nature*, vol. 416, no. 6877, pp. 233-237, Mar 2002.
- [5] M.-G. Suh, and K. J. Vahala, "Soliton microcomb range measurement," *Science*, vol. 359, no. 6378, pp. 884-887, Feb 2018.
- [6] C.-B. Huang, Z. Jiang, D. Leaird, J. Caraquitena, and A. M. Weiner, "Spectral line-by-line shaping for optical and microwave arbitrary waveform generations," *Laser Photonics Rev.*, vol. 2, no. 4, pp. 227-248, Aug 2008.
- [7] D. T. Spencer, S. H. Lee, D. Y. Oh, M.-G. Suh, K. Y. Yang, and K. J. Vahala, "An optical-frequency synthesizer using integrated photonics," *Nature*, vol. 557, no. 7703, pp. 81-85, May 2018.
- [8] V. Torres-Company and A. M. Weiner, "Optical frequency comb technology for ultra-broadband radio-frequency photonics," *Laser Photonics Rev.*, vol. 8, no. 3, pp. 368-393, Dec 2013.
- [9] W. Liang, D. Eliyahu, V. S. Ilchenko, A. A. Savchenkov, A. B. Matsko, D. Seidel, and L. Maleki, "High spectral purity Kerr frequency comb radio frequency photonic oscillator," *Nat. Commun.*, vol. 6, pp. 7957, Aug. 2015.
- [10] C.-H. Li, A. J. Benedick, P. Fendel, A. G. Glenday, F. X. Kärtner, D. F. Phillips, D. Sasselov, A. Szentgyorgyi, R. L. Walsworth, "A laser frequency comb that enables radial velocity measurements with a precision of 1 cm s⁻¹," *Nature*, vol. 452, pp. 610-612, Apr 2008.
- [11] T. Steinmetz, T. Wilken, C. Araujo-Hauck, R. Holzwarth, T. W. Hänsch, L. Pasquini, A. Manescau, S. D'Odorico, M. T. Murphy, T. Kentscher, W. Schmidt, and T. Udem, "Laser frequency combs for astronomical observations," *Science*, vol. 321, no. 5894, pp.1335-1337, Sep 2008.
- [12] P. Marin-Palomo, J. N. Kemal, M. Karpov, A. Kordts, J. Pfeifle, M. H. P. Pfeiffer, P. Trocha, S. Wolf, V. Brasch, M. H. Anderson, R. Rosenberger, K. Vijayan, W. Freude, T. J. Kippenberg, and C. Koos, "Microresonator-based solitons for massively parallel coherent optical communications," *Nature*, vol. 546, no. 7657, pp. 274-279, Jun 2017.
- [13] D. J. Jones, S. A. Diddams, J. K. Ranka, A. Stentz, R. S. Windeler, J. L. Hall, and S. T. Cundiff, "Carrier-envelope phase control of femtosecond mode-locked lasers and direct optical frequency synthesis," *Science*, vol. 288, no. 5466, pp.635-639, Apr 2000.
- [14] B. R. Washburn, S. A. Diddams, N. R. Newbury, J. W. Nicholson, M. F. Yan, and C. G. Jørgensen, "Phase-locked, erbium-fiber-laser-based frequency comb in the near infrared," *Opt. Lett.*, vol. 29, no. 3, pp. 250-252, Feb 2004.
- [15] A. Bartels, D. Heinecke, and S. A. Diddams, "10-GHz self-referenced optical frequency comb," *Science*, vol. 326, no. 5953, pp. 681-681, Oct 2009.
- [16] S. A. Diddams, "The evolving optical frequency comb," *J. Opt. Soc. Am. B: Opt. Phys.*, vol. 27, no. 11, pp. B51-B62, Nov 2010.
- [17] D. Burghoff, T.-Y. Kao, N. Han, C. W. I. Chan, X. Cai, Y. Yang, D. J. Hayton, J.-R. Gao, J. L. Reno, and Q. Hu, "Terahertz laser frequency combs," *Nat. Photonics*, vol. 8, no. 6, pp. 462-467, May 2014.
- [18] J. Kim and Y. Song, "Ultralow-noise mode-locked fiber lasers and frequency combs: principles, status, and applications," *Adv. Opt. Photonics*, vol. 8, no. 3, pp. 465-540, 2016.
- [19] S. Uvin, S. Keyvaninia, F. Lelarge, G.-H. Duan, B. Kuyken, and G. Roelkens, "Narrow line width frequency comb source based on an injection-locked III-V-on-silicon mode-locked laser," *Opt. Express*, vol. 24, no. 5, pp. 5277-5286, Mar 2016.
- [20] Z. Wang, K. V. Gasse, V. Moskalenko, S. Latkowski, E. Bente, B. Kuyken, and G. Roelkens, "A III-V-on-Si ultra-dense comb laser," *Light Sci. Appl.*, vol. 6, no. 5, pp. e16260, May 2017.
- [21] J. Li, H. Lee, T. Chen, and K. J. Vahala, "Low-pump-power, low-phase-noise, and microwave to millimeter-wave repetition rate operation in microcombs," *Phys. Rev. Lett.*, vol. 109, no. 23, pp. 233901, Aug 2012.
- [22] P. Del'Haye, A. Schliesser, O. Arcizet, T. Wilken, R. Holzwarth, and T. J. Kippenberg, "Optical frequency comb generation from a monolithic microresonator," *Nature*, vol. 450, no. 7173, pp. 1214-1217, Dec 2007.
- [23] L. Razzari, D. Duchesne, M. Ferrera, R. Morandotti, S. Chu, B. E. Little, and D. J. Moss, "CMOS-compatible integrated optical hyper-parametric oscillator," *Nat. Photonics*, vol. 4, no. 1, pp. 41-45, Dec 2009.
- [24] J. S. Levy, A. Gondarenko, M. A. Foster, A. C. Turner-Foster, A. L. Gaeta, and M. Lipson, "CMOS-compatible multiple-wavelength oscillator for on-chip optical interconnects," *Nat. Photonics*, vol. 4, no. 1, pp. 37-40, Dec 2009.
- [25] T. J. Kippenberg, R. Holzwarth, and S. A. Diddams, "Microresonator-based optical frequency combs," *Science*, vol. 332, no. 6029, pp.555-559, Apr 2011.
- [26] P. Trocha, M. Karpov, D. Ganin, M. H. P. Pfeiffer, A. Kordts, S. Wolf, J. Krockenberger, P. Marin-Palomo, C. Weimann, S. Randel, W. Freude, T. J. Kippenberg, and C. Koos, "Ultrafast optical ranging using microresonator soliton frequency combs," *Science*, vol. 359, no. 6378, pp. 887-891, Feb 2018.
- [27] B. Stern, X. Ji, Y. Okawachi, A. L. Gaeta, and M. Lipson, "Battery-operated integrated frequency comb generator," *Nature*, vol. 562, no. 7727, pp. 401-405, Oct 2018.
- [28] N. Dupuis, C. R. Doerr, L. Zhang, L. Chen, N. J. Sauer, P. Dong, L. L. Buhl, and D. Ahn, "InP-based comb generator for optical OFDM," *J. Lightwave Technol.*, vol. 30, no. 4, pp. 466-472, Oct 2011.
- [29] M. Pu, L. Ottaviano, E. Semenova, and K. Yvind, "Efficient frequency comb generation in AlGaAs-on-insulator," *Optica*, vol. 3, pp. 823-826, Jul 2016.
- [30] B. Kuyken, T. Ideguchi, S. Holzner, M. Yan, T. W. Hänsch, J. V. Campenhout, P. Verheyen, S. Coen, F. Leo, R. Baets, G. Roelkens, and N. Picqué, "An octave-spanning mid-infrared frequency comb generated in a silicon nanophotonic wire waveguide," *Nat. Commun.*, vol. 6, pp. 6310, Feb 2015.
- [31] T. Sakamoto, T. Kawanishi, and M. Izutsu, "Asymptotic formalism for ultraflat optical frequency comb generation using a Mach-Zehnder modulator," *Opt. Lett.*, vol. 32, no. 11, pp. 1515-1517, May 2007.
- [32] I. L. Gheorma, and G. K. Gopalakrishnan, "Flat frequency comb generation with an integrated dual-parallel modulator," *IEEE Photonics Technol. Lett.*, vol. 19, no. 13, pp. 1011-1013, Jun 2007.
- [33] Z. Jiang, C.-B. Huang, D. E. Leaird, and A. M. Weiner, "Optical arbitrary waveform processing of more than 100 spectral comb lines," *Nat. Photonics*, vol. 1, no. 8, pp. 463-467, Aug 2007.
- [34] R. Wu, V. R. Supradeepa, C. M. Long, D. E. Leaird, and A. M. Weiner, "Generation of very flat optical frequency combs from continuous-wave lasers using cascaded intensity and phase modulators driven by tailored radio frequency waveforms," *Opt. Lett.*, vol. 35, no. 19, pp. 3234-3236, Oct 2010.
- [35] M. A. Soto, M. Alem, M. A. Shoaie, A. Vedadi, C.-S. Brès, L. Thévenaz, and T. Schneider, "Optical sinc-shaped Nyquist pulses of exceptional quality," *Nat. Commun.* vol. 4, pp. 2898, Dec 2013.
- [36] C. Weimann, P. C. Schindler, R. Palmer, S. Wolf, D. Bekele, D. Korn, J. Pfeifle, S. Koeber, R. Schmogrow, L. Alloati, D. Elder, H. Yu, W. Bogaerts, L. R. Dalton, W. Freude, J. Leuthold, and C. Koos, "Silicon-organic hybrid (SOH) frequency comb sources for terabit/s data transmission," *Opt. Express*, vol. 22, no. 3, pp. 3629-3637, Feb 2014.
- [37] J. Lin, H. Sepehrian, Y. Xu, L. A. Rusch, and W. Shi, "Frequency comb generation using a CMOS compatible SiP DD-MZM for flexible networks," *IEEE Photonics Technol. Lett.*, vol. 30, no. 17, pp. 1495-1498, Jul 2018.
- [38] I. Demirtzioglou, C. Lacava, K. RH Bottrill, D. J. Thomson, G. T. Reed, D. J. Richardson, and P. Petropoulos, "Frequency comb generation in a silicon ring resonator modulator," *Opt. Express*, vol. 26, no. 2, pp. 790-796, Jan 2018.
- [39] D. Patel, S. Ghosh, M. Chagnon, A. Samani, V. Veerasubramanian, M. Osman, and D. V. Plant, "Design, analysis, and transmission system performance of a 41 GHz silicon photonic modulator," *Opt. Express*, vol. 23, no. 11, pp. 14263-14287, May 2015.
- [40] D. Dai and S. He, "Optimization of ultracompact polarization-insensitive multimode interference couplers based on Si nanowire waveguides," *IEEE Photonics Technol. Lett.*, vol. 18, no. 19, pp. 2017-2019, Oct 2006.
- [41] A. Mekis, S. Gloeckner, G. Masini, A. Narasimha, T. Pinguet, S. Sahni, and P. Dobbelaere, "A grating-coupler-enabled CMOS photonics platform," *IEEE J Sel. Top. Quant.*, vol. 17, no. 3, pp. 597-608, Jun 2011.
- [42] X. Zhou, X. Zheng, H. Wen, H. Zhang, Y. Guo, and B. Zhou, "All optical arbitrary waveform generation by optical frequency comb based on cascading intensity modulation," *Opt. Commun.*, vol. 284, no. 15, pp. 3706-3710, Jul 2011.
- [43] F. Zhang, J. Wu, Y. Li, and J. Lin, "Flat optical frequency comb generation and its application for optical waveform generation," *Opt. Commun.*, vol. 290, pp. 37-42, Mar 2013.

- [44] Y. Ding, C. Peucheret, H. Ou, and K. Yvind, "Fully etched apodized grating coupler on the SOI platform with -0.58 dB coupling efficiency," *Opt. Lett.*, vol. 39, no. 18, pp. 5348-5340, Sep 2014.
- [45] M. Nakazawa, H. Toshihiko, R. Peng, and G. Pengyu, "Ultrahigh-speed "orthogonal" TDM transmission with an optical Nyquist pulse train," *Opt. Express*, vol. 20, no. 2, pp. 1129-1139, Jan 2012.
- [46] H. N. Tan, T. Inoue, and S. Namiki, "Pass-drop operations of 4x172 Gb/s Nyquist OTDM-WDM over cascaded of WSSs using distributed matched filtering," *Prof. Conf. on Opt. Fiber Comm.*, paper JW2A.50, 2013.

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