

# Optical phonons polarized along the $c$ axis of $\text{YBa}_2\text{Cu}_3\text{O}_{6+x}$ , for $x = 0.5 \rightarrow 0.95$

C.C. Homes, T. Timusk, and D.A. Bonn, R. Liang and W.N. Hardy

**Abstract:** The  $c$ -axis polarized phonon spectra of single crystals of  $\text{YBa}_2\text{Cu}_3\text{O}_{6+x}$ , were measured for the doping range  $x = 0.5 \rightarrow 0.95$ , between 10 K and 300 K. The low background electronic conductivity, determined by Kramers-Kronig analysis of the reflectance, leads to a rich phonon structure. With decreased doping the five normally-active  $B_{1u}$  modes broaden and the high-frequency apical oxygen mode splits into two components. We associate the higher of these with the two-fold coordinated copper “sticks”. The  $155 \text{ cm}^{-1}$  low-frequency mode, which involves the apical and the chain-oxygens, splits into at least three components with decreasing doping. Some phonon anomalies which occur near  $T_c$  in the highly-doped material occur well above  $T_c$  in the oxygen-reduced systems. An unusual broad phonon band develops in the normal state at  $\approx 400 \text{ cm}^{-1}$ , which becomes more intense at low doping and low temperatures, borrowing oscillator strength from apical and plane oxygen modes resulting in a major transformation of the phonon spectrum below  $\approx 150 \text{ K}$ .

**Résumé:** Les spectres de phonons polarisés suivant l'axe  $c$ , dans des monocristaux de  $\text{YBa}_2\text{Cu}_3\text{O}_{6+x}$  ont été mesurés, pour l'intervalle de dopage  $x = 0,5$  à  $0,95$ , entre 10 et 300 K. La faible conductivité électronique de fond, déterminée par analyse Kramers-Kronig de la réflectance, donne lieu à une riche structure de phonons. Lorsqu'on diminue le dopage, les cinq modes  $B_{1u}$  normalement actifs s'élargissent, et le mode apical à haute fréquence de l'oxygène se divise en deux composantes. Nous associons la plus haute de ces composantes aux « bâtons » de cuivre à coordination deux. Le mode de basse fréquence, à  $155 \text{ cm}^{-1}$ , qui implique les oxygènes apical et en chaîne, se sépare en au moins trois composantes lorsque le dopage diminue. Certaines anomalies de phonons, présentes au voisinage de  $T_c$  dans les cristaux faiblement dopés, apparaissent à des températures bien plus grandes que  $T_c$  dans les systèmes avec moins de l'oxygène. A environ  $400 \text{ cm}^{-1}$ , il se développe dans l'état normal une bande large insolite qui devient plus intense à faible dopage et à basse température, avec une force d'oscillateur empruntée aux modes apical et plan de l'oxygène, ce qui entraîne une transformation majeure des spectres de phonons au-dessous d'environ 150 K.

[Traduit par la rédaction]

PACS N<sup>os</sup> : 63.20.Dj, 74.25.Gz, 74.72.Bk, 78.30.-j

## 1. Introduction

The role of phonons in high-temperature superconductors is complex. On the one hand there is ample evidence that the cuprates are not conventional electron-phonon BCS superconductors. This evidence includes, for the optimally doped materials a linear resistivity, in the normal state, with a zero-temperature intercept at zero resistance [1, 2], a collapse of transport scattering at the superconducting transition [3–5], and a lack of an isotope effect [6–9].

On the other hand, a large number of studies show that cer-

tain phonons at least are coupled to the carriers, and that this coupling is unconventional [10, 11]. The frequencies and widths of certain phonon lines, again in the optimally-doped materials (precisely those where the isotope effect is small), undergo abrupt changes at  $T_c$ . This effect was first observed by infrared spectroscopy [12–15], then by Raman spectroscopy [15–17], resonant neutron-absorption spectroscopy [18] and neutron scattering [19].

Broad minima in the  $ab$ -plane optical conductivity are seen at frequencies that are in the phonon region. These have been shown to be associated with  $c$ -axis longitudinal optic (LO) phonons [20, 21] that couple to the  $ab$ -plane conductivity by a symmetry-breaking mechanism.

In this paper we focus on the changes in the  $c$ -axis phonon spectra of  $\text{YBa}_2\text{Cu}_3\text{O}_{6+x}$  as we change the oxygen doping. These spectra can be separated from the weak and smooth electronic background conductivity whose intensity rises with doping from a few  $\Omega^{-1}\text{cm}^{-1}$  in the  $x = 0.5$  material<sup>3</sup> to  $450 \Omega^{-1}\text{cm}^{-1}$  in overdoped samples [22]. The background is due to inter-cell hopping of free carriers [23]. Superimposed on this electronic background are five strong phonon lines. These

Received May 8, 1995. Accepted July 14, 1995.

*This article appeared in a special issue of the Canadian Journal of Physics in honour of B.N. Brockhouse.*

**C.C. Homes<sup>1</sup> and T. Timusk<sup>2</sup>** Department of Physics and Astronomy, McMaster University, Hamilton, ON L8S 4M1, Canada. **D.A. Bonn, R. Liang and W.N. Hardy.** Department of Physics and Astronomy, University of British Columbia, Vancouver, BC, V6T 2A6 Canada.

<sup>1</sup>Present Address: Condensed Matter Physics and Materials Science Department, Brookhaven National Laboratory, Upton, NY 11973, USA

<sup>2</sup>Corresponding author (e-mail: timusk@mcmaster.ca).

<sup>3</sup>C.C. Homes, T. Timusk, D.A. Bonn, R. Liang and W.N. Hardy. Unpublished results.

five lines are also seen in the spectra of ceramic superconductors [24, 25], and have been studied by a large number of groups. The extensive literature on ceramic spectra has been summarized by Feile in ref. 26. However, in ceramics, effects due to the background *ab*-plane conductivity distort the spectra, particularly the line strengths, and the *c*-axis polarized phonons are best studied in single crystals. Single crystal studies are difficult due to the small area of the available flux-grown samples. Nevertheless, a number of studies have appeared, with  $\text{YBa}_2\text{Cu}_3\text{O}_{6+x}$  being the most studied material [see footnote 3 and refs. 22, 27–32]. In the present work single crystals with cleaved surfaces are used and the region of doping ranges from the optimally doped, that is the material with the highest  $T_c = 93.5$  K with  $x = 0.95$ , to an underdoped sample with  $x = 0.5$  which is still superconducting with a  $T_c = 53$  K. Preliminary reports of this work have been published [32–34], and the electronic background is the subject of a separate publication [see footnote 3].

## 2. Experimental Details

### 2.1. Experimental

The spectra were obtained by a Kramers-Kronig transformation of *c*-axis polarized reflectance of millimeter-size, high-quality single crystals. The growth of the single crystals and experimental techniques used have been reported on previously [32, 35]. To maximize the signal from the *ac* face of the crystal, an infrared beam larger than the crystal was used. The sample was coated with gold *in situ* and then remeasured at each temperature to obtain a reference spectrum [36].

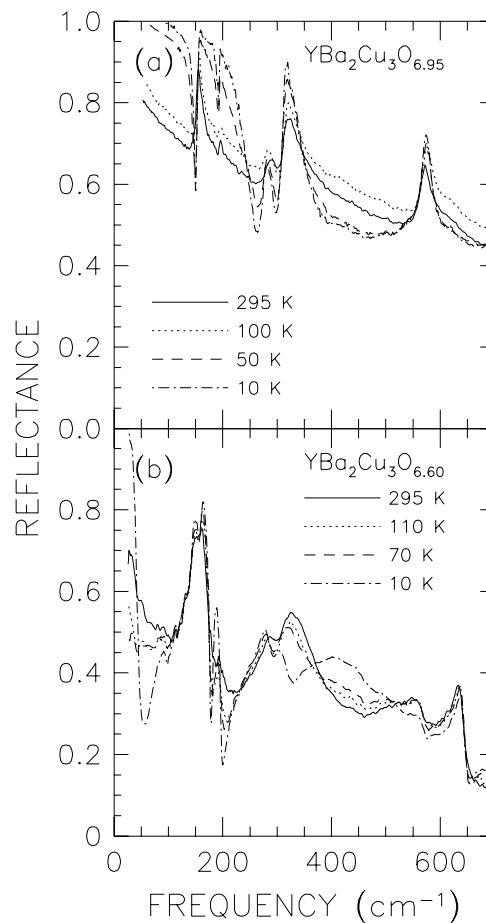
The reflectance of  $\text{YBa}_2\text{Cu}_3\text{O}_{6+x}$ , for two oxygen concentrations,  $x = 0.6$  and  $0.95$ , for light polarized parallel to the *c* axis is shown in the phonon frequency range in Fig. 1 at several temperatures. The reflectance has a continuous background, rising towards unity at low frequency, characteristic of a poor metal, with sharp resonances due to optical phonons. In the superconducting state, the development of the superconducting condensate has the effect of producing a reflectance edge which shifts down to lower frequency as the doping is reduced. Also, as the temperature is lowered, the shapes of the phonon features change due to the changing background conductivity.

Changes in oxygen doping have two principal effects on the phonon spectra: first, as the doping level is lowered, a new phonon feature appears at  $\approx 610$   $\text{cm}^{-1}$ . Secondly, a very broad feature, weak at first, but stronger at low oxygen dopings, appears in the normal state at  $\approx 400$   $\text{cm}^{-1}$ , and becomes stronger at low temperature. These effects are seen more clearly in the optical conductivity, discussed in the following section.

### 2.2. Optical conductivity

A Kramers-Kronig transformation of the reflectance was used to calculate the optical conductivity. The reflectance was extended to regions outside the actual measurements as follows: below the lowest measured frequency Hagen-Rubens frequency dependence:  $[(1-R) \propto \omega^{-1/2}]$  was assumed to hold in the normal state. In the superconducting state, where a plasma edge develops, it was assumed that the reflectance was given by  $(1-R) \propto \omega^{-2}$ . This assumes that the conductivity is finite and constant to zero frequency. There is good evidence

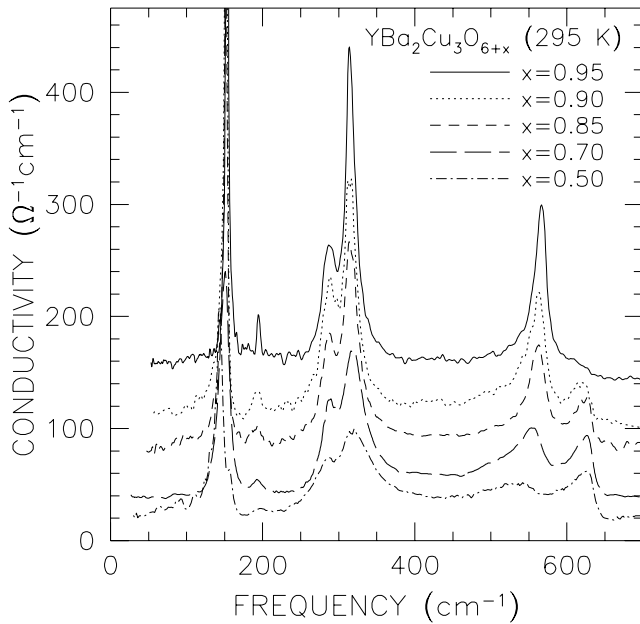
**Fig. 1.** The reflectance of  $\text{YBa}_2\text{Cu}_3\text{O}_{6+x}$  for radiation polarized along the *c*-axis from  $\approx 50$  to  $700$   $\text{cm}^{-1}$  at several temperatures above and below  $T_c$ , for two oxygen dopings  $x$ : (a) fully-doped material with  $x = 0.95$  ( $T_c = 93$  K), and (b) an underdoped crystal with  $x = 0.6$  ( $T_c = 58$  K). The main change in the phonon spectrum in the underdoped material is a new line at  $610$   $\text{cm}^{-1}$  that we associate with the bridging oxygen at the two-fold coordinated copper sites in the chain layer. Another new feature in the  $x = 0.6$  material is a broad band at  $\approx 400$   $\text{cm}^{-1}$  that appears at low temperature.



from comparisons with transport measurements that both assumptions are valid [see footnote 3]. Temperature-dependent reflectance measurements were carried out to  $5000$   $\text{cm}^{-1}$ . The room-temperature measurements of Koch et al. [29] were used from  $5000$   $\text{cm}^{-1}$  to  $3.5 \times 10^5$   $\text{cm}^{-1}$  and above this frequency free-electron behavior was assumed:  $R \propto \omega^{-4}$ .

The  $295$  K conductivity along the *c*-axis is shown in Fig. 2 for a wide range of dopings. The  $10$  K conductivity is shown in Fig. 3. Superimposed on a continuous background that increases in strength with doping are a series of sharp phonon lines. In general, the strength of the phonon bands does not depend strongly on doping, but as the doping is reduced a new phonon appears at  $\approx 610$   $\text{cm}^{-1}$  in the nearly fully-doped material and moves to higher frequency as the oxygen concentration is reduced. We will call this the  $610$   $\text{cm}^{-1}$  line in what follows, but its actual frequency changes with doping and is as

**Fig. 2.** The optical conductivity ( $\sigma_1$ ) at 295 K for  $\text{YBa}_2\text{Cu}_3\text{O}_{6+x}$  for radiation polarized along the  $c$ -axis from  $\approx 50$  to  $700 \text{ cm}^{-1}$  for five oxygen dopings,  $x = 0.5 \rightarrow 0.95$ . As the doping is reduced, the background conductivity is reduced uniformly at all frequencies. At high frequency, in addition to the apical oxygen [O(4)] peak at  $570 \text{ cm}^{-1}$  a new phonon peak grows at  $\approx 610 \text{ cm}^{-1}$ ; this peak is associated with the O(4) mode and the growing density of two-fold coordinated copper sites in the underdoped materials. As the doping is reduced, the O(4) modes acquire an increasingly dispersive line shape.



high as  $635 \text{ cm}^{-1}$  in the  $x = 0.5$  material. The other major effect of doping occurs at low temperature where a new broad band appears at  $400 \text{ cm}^{-1}$  as the doping level is reduced to  $x = 0.7$ . This band seems to increase its oscillator strength at the expense of the two high-frequency modes and the band at  $310 \text{ cm}^{-1}$ . Also, its center frequency appears to decrease as the doping is reduced, reaching a value of  $\approx 375 \text{ cm}^{-1}$  in the  $x = 0.5$  crystal.

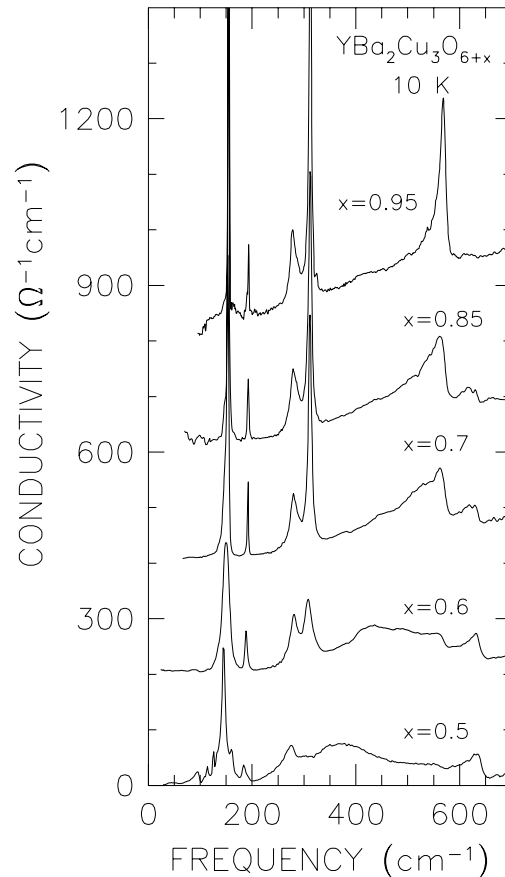
### 3. Results

#### 3.1. Calculation of the phonon parameters

In what follows we refer to the six phonons that we observe by their room temperature frequencies in the fully-doped material, bearing in mind that the actual center frequencies change with both temperature and doping. There are five strong phonons, and at room temperature their frequencies are: 155, 194, 279, 312 and  $570 \text{ cm}^{-1}$ , plus a weaker one at  $610 \text{ cm}^{-1}$ . All the lines have symmetric line shapes except for the two highest ones. The  $570 \text{ cm}^{-1}$  and the  $610 \text{ cm}^{-1}$  lines have Lorentzian shapes at room temperature, but at low temperature they become asymmetric [37]. Such lines are usually described in terms of a Fano line shape

$$\sigma_1(\omega) = A \left[ \frac{(x + q)^2}{(1 + x^2)} \right]$$

**Fig. 3.** The  $c$ -axis optical conductivity ( $\sigma_1$ ) at 10 K for five oxygen dopings,  $x = 0.5 \rightarrow 0.95$ . The curves have been offset vertically in  $200 \text{ } \Omega^{-1}\text{cm}^{-1}$  increments for clarity. There is a major redistribution of spectral weight from the O(4) peaks at  $\approx 570 \text{ cm}^{-1}$  and  $\approx 610 \text{ cm}^{-1}$ , and the in-plane buckling mode at  $310 \text{ cm}^{-1}$ , to a new broad band in the  $\approx 400 \text{ cm}^{-1}$  range. The peak shifts to lower frequency as the doping is reduced, starting as a shoulder on the apical-oxygen line at  $550 \text{ cm}^{-1}$  and ending up centered at  $\approx 375 \text{ cm}^{-1}$  in the  $x = 0.5$  sample.



where  $\sigma_1(\omega)$  is the conductivity,  $A$  is a constant,  $x = (\omega - \omega_i)/\gamma_i$ ,  $\omega_i$  is the phonon frequency,  $\gamma_i$  is the phonon line width and  $q_i$  is a parameter that describes the asymmetry of the  $i$ th phonon [38].

While the Fano line shape describes the  $570 \text{ cm}^{-1}$  phonon at low temperatures, it fails as the line shape becomes more Lorentzian ( $q \rightarrow \infty$ ). A better fit to the data is in terms of a classical Drude-Lorentz model for the dielectric function, for each phonon band, but with an empirical line shape made up of a mixture of the real and imaginary parts of the Lorentzian oscillator:

$$\tilde{\epsilon}(\omega) = \frac{\omega_p^2 e^{i\theta}}{\omega_{\text{TO}}^2 - \omega^2 - i\omega\gamma} \quad (1)$$

where  $\omega_p$  is the effective plasma frequency,  $\omega_{\text{TO}}$  the center frequency and  $\gamma$  the width of the phonon. The parameter  $\theta$  is the phase associated with the asymmetry of the phonon. The

**Table 1.** The parameters used to fit Lorentzian line shapes to the peaks in the optical conductivity of  $\text{YBa}_2\text{Cu}_3\text{O}_{6+x}$ , for  $x = 0.95, 0.85, 0.7$ , and  $0.5$  at 10 K. The two high-frequency copper-oxygen vibrations have been fitted using non-zero rotations.<sup>a</sup>

$x = 0.95$			$x = 0.85$			$x = 0.7$			$x = 0.5$					
$\omega_{\text{TO},i}$	$\gamma_i$	$\omega_{pi}$	$\omega_{\text{TO},i}$	$\gamma_i$	$\omega_{pi}$	$\omega_{\text{TO},i}$	$\gamma_i$	$\omega_{pi}$	$\omega_{\text{TO},i}$	$\gamma_i$	$\omega_{pi}$			
155	2.0	398	{	147	8.4	248	{	136	15	167	{	93		
				153	3.8	351		147	9.8	322		113	5.0	80
								153	4.2	267		125	5.2	117
											145	9.0	363	
												161	5.9	120
194	2.2	128	→	192	2.8	133	→	189	4.2	133	→	184	5.9	110
279	19	357	→	281	19	329	→	280	22	323	→	272	33	299
312	5.2	487	→	312	7.0	417	→	309	18	333	→	301	28	184
570	14	483	{	566	23	321	{	564	30	188	{	551	30	127
				619	23	205		630	24	250		637	24	251

NOTE: All parameters are in  $\text{cm}^{-1}$ .

<sup>a</sup> The two high-frequency apical-oxygen modes have been fit using a Lorentzian with a phase. The angles (in radians) used are for  $x = 0.95$ ,  $\theta = -0.40, 0.0$ , for  $x = 0.85$ ,  $\theta = -0.50, -0.15$ , for  $x = 0.7$ ,  $\theta = -0.62, -0.45$ , for  $x = 0.5$ ,  $\theta = -0.75, -0.80$ , for the low- and high-frequency components, respectively.

line shape is that of a classical oscillator for  $\theta = 0$  but becomes asymmetric for nonzero  $\theta$ .

An advantage of this rotated-Lorentzian line shape over the Fano profile is that one can associate with it an oscillator strength. For a pure Lorentzian, the integrated oscillator strength is defined to be  $\omega_p^2/8$ ; for the rotated Lorentzian, to leading order the oscillator strength is  $\cos(\theta)\omega_p^2/8$ , so that for small rotations, the introduction of an asymmetric line shape does not alter the oscillator sum rule too much. However, for rotations of greater than  $\approx 0.4$  radians, the error introduced is greater than 10%, and the other contributions to the sum rule begin to play an increasingly large role, further compounding this error. The asymmetry of the rotated Lorentzian is related to the Fano parameter  $q$  by  $q^{-1} \propto \tan(\theta/2)$ .

All the phonon lines have been fitted to (1) with a non-linear least-squares method. In addition to the parameters of the modified-Lorentzian oscillator, a linear background was used in the fits. We find that within the accuracy of our experiments, the four low-frequency lines are symmetric and only the  $570 \text{ cm}^{-1}$  line, and the new feature associated with it at  $610 \text{ cm}^{-1}$ , have clearly asymmetric shapes which become more pronounced at low temperatures.

It should be noted that if the Kramers-Kronig analysis is done on data that deviates from the true reflectance, for example by using room-temperature data for the high-frequency extrapolation for a low-temperature data set, there is a tendency for sharp lines to acquire an artificial asymmetry. As the experimental data become more accurate, the lines become more symmetric. The lines that have the highest reflectance are most susceptible to this error, which in our case is the  $155 \text{ cm}^{-1}$  line, which we find accurately symmetric, thus confirming the position of our 100% line. We find that an error of greater than 2% in the unity reflectance line would show an observable asymmetry.

The phonon parameters obtained by the least-squares fitting procedure applied to the rotated-Lorentzian profile with of all the measured spectra are assembled in Table 1. We estimate the error in the parameters to be  $\approx 1\%$  for the center frequencies,  $\approx 5\%$  for the line widths, and  $\approx 5\%$  for the effective plasma frequencies. Where there is a finite rotation angle  $\theta$ , the error could be larger.

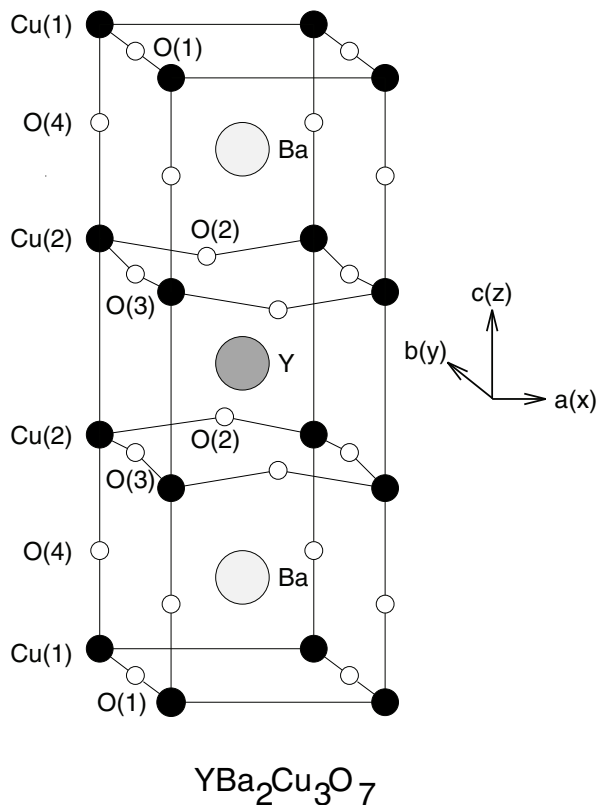
### 3.2. Assignment of phonons

Group theory predicts seven  $B_{1u}$  infrared-active modes polarized in the  $c$ -direction for the fully-oxygenated orthorhombic (Ortho I)  $\text{YBa}_2\text{Cu}_3\text{O}_7$  shown in Fig. 4 [39, 40]. There is crystallographic evidence that the oxygen-reduced material, with a plateau in the  $T_c$  vs  $x$  graph (in the  $T_c \approx 60 \text{ K}$  region), has a doubled unit cell normal to the chain direction, resulting from an alternation of fully-occupied O(1) sites on chains with empty sites. This structure (Ortho II) is expected to have 13 infrared-active modes, but many of the new modes result from extremely small splittings of modes involving atoms far away from the chains. These small splittings may be partially responsible for the systematic broadening of most of the phonon lines at reduced dopings. We expect the largest changes for modes that involve atoms on and near the chains.

The assignment of the infrared-active spectral lines to normal modes is most reliable when the results from inelastic-neutron scattering are combined with a realistic lattice-dynamic model [41]. Such a procedure has been carried out recently for fully-oxygenated  $\text{YBa}_2\text{Cu}_3\text{O}_7$ ,<sup>4</sup> and we make extensive use of the eigenvectors of these models in our discussion of  $c$ -axis phonon spectra. A list of calculated contributions the various modes to the polarizability has been published by Timusk et al. [34]. It is clear from the shell models that infrared modes,

<sup>4</sup>W. Reichardt. Private communication.

**Fig. 4.** The orthorhombic unit cell of  $\text{YBa}_2\text{Cu}_3\text{O}_7$ . The chain oxygen is denoted as O(1), the apical copper as Cu(1), and the O(4) as the apical oxygen atom. When all the O(1) sites are occupied, the Cu(1) atom is four-fold coordinated. Deoxygenation results in the oxygens at the O(1) sites being removed, lowering the coordination of the Cu(1) atom.



unlike Raman modes, which can often be assigned to the motion of a particular ion, usually involve the motion of all the ions in the unit cell. The shell-model eigenvectors suggest that there are some exceptions. For example the  $570\text{ cm}^{-1}$  mode is almost a pure mode of the apical O(4) oxygen (a description of the different copper and oxygen environments is shown in Fig. 4). Similarly, the  $279\text{ cm}^{-1}$  mode is an O(1) chain-oxygen vibration, polarized in the  $c$ -direction and the  $315\text{ cm}^{-1}$  mode an O(2) and O(3) plane-bending vibration. The low-frequency modes involve many more atoms, and while the  $194\text{ cm}^{-1}$  mode has been identified as the yttrium mode, it also involves substantial motion of the O(2) planar-oxygen atoms. The  $155\text{ cm}^{-1}$  mode, often called the barium mode, according to the shell model gets more spectral weight from the O(4) oxygen motion than from the barium, and also involves the motions of the Cu(1) and O(1) chain-copper and -oxygen atoms. This prediction is in agreement with our observation that this modes splits in oxygen-reduced materials.

In the highly-doped  $\text{YBa}_2\text{Cu}_3\text{O}_{6.95}$  material five strong phonons are observed at  $155, 194, 279, 312$  and  $570\text{ cm}^{-1}$  at  $10\text{ K}$ . These frequencies are in excellent agreement with previous measurements on ceramics [24, 42]. Table 2 compares these frequencies with those of the neutron-based shell model. As expected, since the  $k = 0$  frequencies have been fitted to the

**Table 2.** The fitted phonon parameters for the  $c$ -axis  $B_{1u}$  modes in  $\text{YBa}_2\text{Cu}_3\text{O}_{6.95}$  at  $295\text{ K}$  determined from optical and inelastic neutron-scattering measurements. The effective-plasma frequency based on the shell model with formal charges in column 5 [ $Y = +3.0, \text{Ba} = +2.0, \text{Cu}1 = \text{Cu}2 = 2, \text{O}(1) = \text{O}(2) = \text{O}(3) = \text{O}(4) = -2$ ] and fitted charges in column 6, [ $Y = +2.8, \text{Ba} = +2.5, \text{Cu}1 = \text{Cu}2 = 2, \text{O}(1) = -1.8, \text{O}(2) = -1.8, \text{O}(3) = -1.7, \text{O}(4) = -2.6$ ].

Optical <sup>a</sup>			Neutron <sup>b</sup>		
$\omega_{\text{TO}i}$	$\gamma_i$	$\omega_{pi}$	$\omega_i$	$\omega_{pi}^c$	$\omega_{pi}^d$
—	—	—	111	100	115
153	3.6	389	149	283	330
—	—	—	186	22	14
195	3.7	95	198	79	94
287	18.3	316	291	455	400
315	15.6	498	322	493	457
568 <sup>e</sup>	18.4	408	556	256	365

NOTE: All units are in  $\text{cm}^{-1}$ .

<sup>a</sup>Fitted to optical conductivity, this work.

<sup>b</sup>W. Reichardt et al. Private communication.

<sup>c</sup>Effective charges, formal valence charges.

<sup>d</sup>Effective charges, fitted to optical data.

<sup>e</sup>The high-frequency apical-oxygen mode has been fitted using an asymmetric line shape described by a Lorentzian with a rotation angle of  $\theta = -0.27\text{ rad}$ .

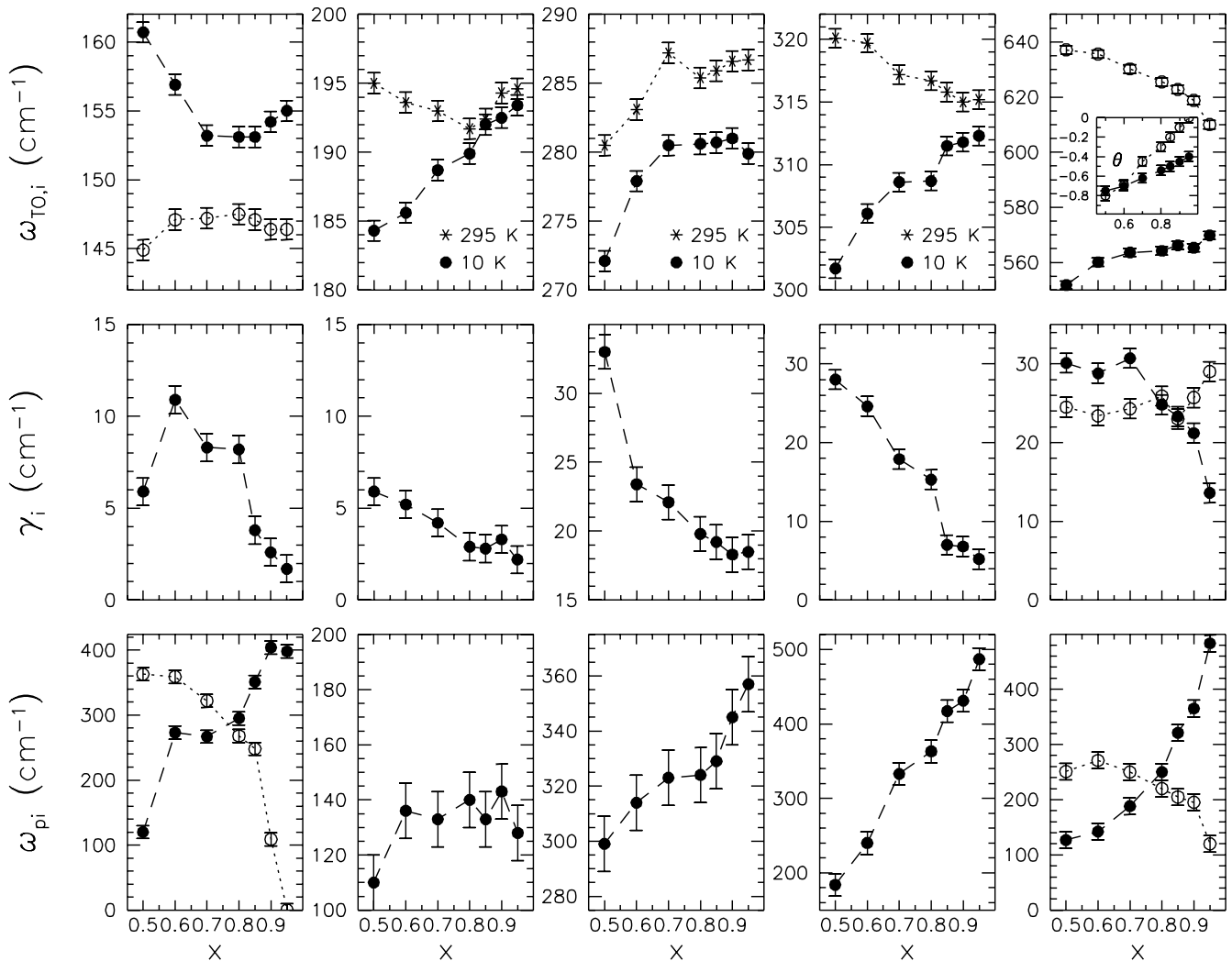
model, the agreement is satisfactory. What is surprising is the excellent agreement between the measured line strengths, and those predicted by the shell-model eigenvectors. In calculating the line strengths we have used two separate approaches: the simplest is to use the formal charges of the ions in all cases. The only case where this seems to be inappropriate is for the  $570\text{ cm}^{-1}$  mode where the observed line strength is larger, and a larger effective charge on the O(4) would give better agreement between the model and the infrared observations. This line is also anomalous in that it develops a strong asymmetry in the oxygen-reduced materials, a sign of strong electron-phonon interaction which may enhance the polarizability of the mode by dynamic charge transfer from the O(4) to the planes.

The second approach is to use a least-squares fit to the effective charges to the observed infrared oscillator strengths. The results are shown in column six of Table 2. The fitted results, as expected, agree better with the optical data. However, it is interesting to note that in order to get this better fit, the main change was the removal of charge from the O(4) ion, and the addition of that charge on the planar oxygen ions.

The two missing lines are predicted by the shell model to be very weak. Of these there is some evidence for the lower one predicted to be at  $111\text{ cm}^{-1}$  in some of our spectra, but they are near the limit of the accuracy of the experiments.

Two modes not predicted by the shell model are seen: the broad band at  $\approx 400\text{ cm}^{-1}$ , and weak peak at  $\approx 610\text{ cm}^{-1}$ . In what follows we will argue that the  $610\text{ cm}^{-1}$  line is associated with the two-fold coordinated copper neighbor of the O(4) oxygen. Such sites are found in oxygen-reduced material where the  $610\text{ cm}^{-1}$  line is much stronger, and its presence in

**Fig. 5.** The variation of the central frequency ( $\omega_{TO,i}$ ), width ( $\gamma_i$ ) and the effective plasma frequency ( $\omega_{pi}$ ) as a function of doping ( $x = 0.5 \rightarrow 0.95$ ) of the five strong phonons observed in the highly-doped materials, at 10 K ( $\bullet$ ). The room-temperature frequencies have been included for the three central panels (\*). Also included in the panels are two phonons that appear only at lower oxygen dopings ( $\circ$ ): in the first panel, a new mode appears at  $\approx 148 \text{ cm}^{-1}$  and gains oscillator strength as the mode observed in the highly-doped materials at  $\approx 155 \text{ cm}^{-1}$  becomes weaker. The second phonon appears at  $\approx 610 \text{ cm}^{-1}$  and, similarly to the  $148 \text{ cm}^{-1}$  phonons, gains oscillator strength with decreasing doping. Inset: rotation angles for the  $\approx 570 \text{ cm}^{-1}$  mode ( $\bullet$ ), and the  $\approx 610 \text{ cm}^{-1}$  mode ( $\circ$ ).



our  $x = 0.95$  sample is an indication that this sample contains a small concentration of broken chain fragments with two-fold coordinated copper.

Overall, at room temperature at least, there is good agreement with the shell model and the observed phonon spectrum and there is no need to invoke special electronic enhancements to explain the intensities of the phonon lines, nor is there any evidence of strong screening by the electronic background.

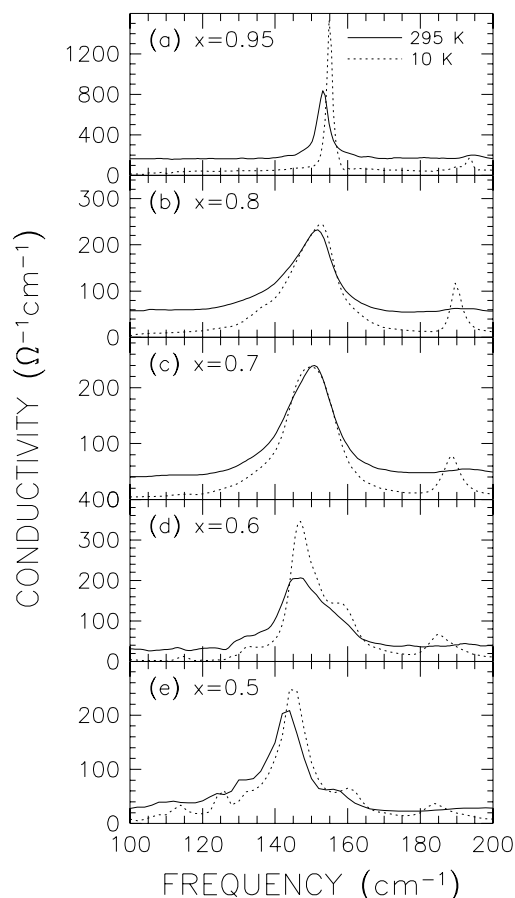
### 3.3. Doping dependence of the phonon parameters

As the doping level is reduced, a number of changes in the phonon parameters take place; these are summarized for the five fundamental phonons observed in the highly-doped material, as well as the two strong splittings associated with the  $155$  and  $570 \text{ cm}^{-1}$  modes at 10 K observed at lower dopings,

in Fig. 5. To a first approximation one expects the phonon frequencies to soften with the expansion of the  $c$ -axis lattice parameter. One generally expects the frequency to vary approximately as the third power of the lattice parameter [43]. Since the  $c$ -axis lattice parameter increases by 0.5% between  $x = 0.93$  to  $x = 0.50$ , one would expect an overall softening of the order of 1.5% as the result of the lattice expansion.

In very rough agreement with this simple picture, many modes soften, but by a larger amount of 3.5% in this doping range. There are also several notable exceptions. First there is a tendency, at room temperature, for several modes to harden as the lattice expands on removal of oxygen, for example the  $312 \text{ cm}^{-1}$  mode. The  $155 \text{ cm}^{-1}$  mode as well as the high-frequency O(4) mode behave in a similar fashion: they split into two components, the higher frequency component hard-

**Fig. 6.** The evolution of the shape of the  $155\text{ cm}^{-1}$  line, at 295 K and 10 K, for five oxygen dopings ( $x = 0.5 \rightarrow 0.95$ ). As the oxygen doping is reduced, the line broadens, first developing shoulders, and then finally splitting into several clearly resolved components.



ens, the lower frequency one softens. As we will show below for the  $570\text{ cm}^{-1}$  mode, and its high-frequency partner the  $610\text{ cm}^{-1}$  mode the higher frequency component corresponds to the O(4) at the two-fold coordinated copper sites whereas the lower frequency corresponds to the O(4) at the four-fold coordinated copper sites. While the shell model eigenvectors suggest that the  $155\text{ cm}^{-1}$  mode also involves a large amount of O(4) motion, many other atoms in the unit cell participate in the motion, and the net result is a behavior that is opposite to that for the  $570\text{ cm}^{-1}$  mode: it is the high-frequency, hard component that loses oscillator strength.

The mode at  $\approx 155\text{ cm}^{-1}$  in the highly-doped material has split into a doublet in the  $x = 0.85$  material, and consists of at least three lines in the  $x = 0.7$  material; the phonon lines are significantly broader than the narrow line in the  $x = 0.95$  material. Fig. 6 shows the evolution of the shape of this line with doping. The multi-component nature of the line and the gradual shift, with reduced doping to lower frequency components is reminiscent of the  $500\text{ cm}^{-1}$  Raman line which has been shown to have many components, associated with chain disorder in underdoped samples [44]. The splitting of the  $\approx 155\text{ cm}^{-1}$  mode is also associated with the loss of O(1)

chain oxygens. Unlike the  $570\text{ cm}^{-1}$  mode, where the O(4)–Cu(1) chain-copper bond is stretched, the  $155\text{ cm}^{-1}$  mode involves the O(4)–Cu(2) bond which is also expected to change its length as the Cu(1) changes from four- to two-fold coordination.

As the doping decreases to  $x \approx 0.7$ , the  $155\text{ cm}^{-1}$  O(4) mode softens and decreases in strength, while the mode at  $148\text{ cm}^{-1}$  hardens slightly, and gains strength. Below  $x \approx 0.7$ , it is clear that the fundamental nature of this mode has changed; the  $155\text{ cm}^{-1}$  feature hardens dramatically, and loses almost all of its oscillator strength; in contrast, the  $148\text{ cm}^{-1}$  mode softens, and its oscillator strength approaches the value of the  $155\text{ cm}^{-1}$  mode in the highly-doped material.

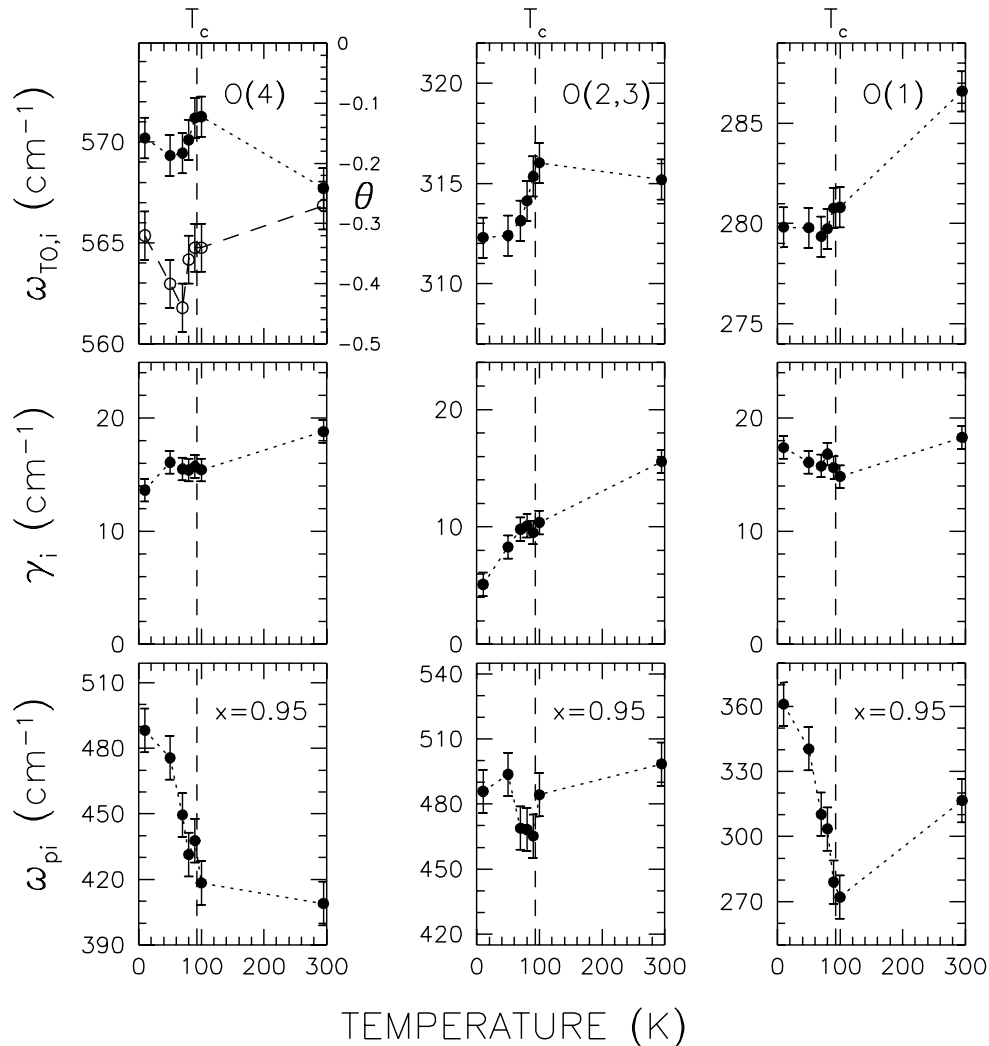
The  $312\text{ cm}^{-1}$  plane-oxygen mode displays drastically different behavior at room temperature and at 10 K. At room temperature, the mode hardens from  $315\text{ cm}^{-1}$  to  $320\text{ cm}^{-1}$  and its oscillator strength remains unchanged as the doping is reduced. However, at 10 K, this mode softens dramatically with decreasing doping, furthermore, at  $x \approx 0.5$  this mode, at low temperature, has lost most of its considerable oscillator strength. The loss of oscillator strength at low temperature in the underdoped materials results from the transfer of their spectral weight to the  $400\text{ cm}^{-1}$  mode, which will be discussed in more detail below.

In the fully-oxygenated material the  $610\text{ cm}^{-1}$  mode can just be seen as a weak side band of the  $570\text{ cm}^{-1}$  mode, but as the doping is reduced the  $610\text{ cm}^{-1}$  grows substantially in spectral weight. The mode hardens substantially to  $645\text{ cm}^{-1}$  in the  $x = 0.5$  material, while the  $570\text{ cm}^{-1}$  mode that it is associated with is steadily losing oscillator strength with decreasing doping. This effect has been attributed by Burns et al. to the shortening of the O(4)–Cu(1) bond by 6% in the same concentration range [45]. In rough agreement with this model the  $610\text{ cm}^{-1}$  line changes its frequency by 6% as well in the same concentration range. In Burns' picture the  $570\text{ cm}^{-1}$  line arising from the O(4) at the four-fold coordinated copper, softens by 3%, a value typical of the overall  $c$ -axis expansion, suggesting a lengthening of that bond. If the interpretation of Burns is correct we would expect a modulation of the O(4) position in the  $a$ -direction by approximately  $0.03\text{ \AA}$  ( $1\text{ \AA} = 10^{-10}\text{ m}$ ) in the Ortho II structure where there is an alternation between two- and four-fold coordinated coppers. There have been reports, by several authors [46–48], of a double peak in the Cu(1)–O(4) bond length distribution measured by XAFS by as much as  $0.1\text{ \AA}$ .

The widths of the phonons increase steadily with decreasing doping, with the exception of the  $155\text{ cm}^{-1}$  mode, which shows some signs of narrowing at low dopings. The steadily increasing widths may be a reflection of the O(1)-site disorder at low dopings. However, if there is some long-range ordering of the O(1) sites at low dopings, as is expected in the Ortho II phase, then there should be some narrowing observed in the lines. While there is some evidence of this in the  $155\text{ cm}^{-1}$  mode, the  $279\text{ cm}^{-1}$  mode shows no signs of narrowing at low-oxygen dopings.

The oscillator strengths are a convenient guide by which to follow the effects of the removal of oxygen from the O(1) sites. In particular, this is true for the  $279\text{ cm}^{-1}$  mode, which is almost exclusively an O(1)-oxygen vibration. In the  $x = 0.5$  material, we expect that the O(1) sites should be half empty. At 295 K, the effective-plasma frequency of this mode decreases

**Fig. 7.** The temperature dependence of the frequency ( $\omega_{\text{TO},i}$ ), line width ( $\gamma_i$ ) and effective plasma frequency ( $\omega_{\text{pi}}$ ) for the three oxygen-related *c*-axis modes of  $\text{YBa}_2\text{Cu}_3\text{O}_{6.95}$  ( $\bullet$ ). The  $\approx 279 \text{ cm}^{-1}$  mode fails to show any discontinuity at  $T_c$ , while the 312 and 570  $\text{cm}^{-1}$  modes both soften below  $T_c$ ; the 570  $\text{cm}^{-1}$  feature also appears to gain oscillator strength below  $T_c$ . The rotation angle used to fit the asymmetric line shape is shown in the O(4) panel ( $\circ$ ), and shows a slightly increasing asymmetry with decreasing temperature.



from  $\approx 310 \text{ cm}^{-1}$  to  $\approx 200 \text{ cm}^{-1}$  in the  $x = 0.5$  material; taking the naïve view that the ratio of the oscillator strengths ( $\propto \omega_p^2$ ) of 0.44 should be a reflection of the half-empty occupancy of the O(1) sites, we find good agreement.

The asymmetric line shapes of the 570 and 610  $\text{cm}^{-1}$  modes require that they be fitted using a rotated Lorentzian. While this line shape generally satisfies oscillator sum rules for small rotations, it has been pointed out that for large rotations ( $\theta > 0.4$  rad) this is no longer true. The values for  $\omega_p$  may generally have increasingly large errors associated with them for large rotations. The results show a decreasing oscillator strength for the 570  $\text{cm}^{-1}$  with decreasing doping, and the 610  $\text{cm}^{-1}$  mode show a steadily increasing oscillator strength with decreasing doping with the two oscillator strengths becoming equal at  $x = 0.7$ . Below  $x \approx 0.6$ , the 610  $\text{cm}^{-1}$  line appears to be decreasing in intensity. This latter result is consistent with the small value for  $\omega_p$  observed in the tetragonal  $\text{YBa}_2\text{Cu}_3\text{O}_6$  material.

#### 3.4. Temperature dependence of the phonon parameters

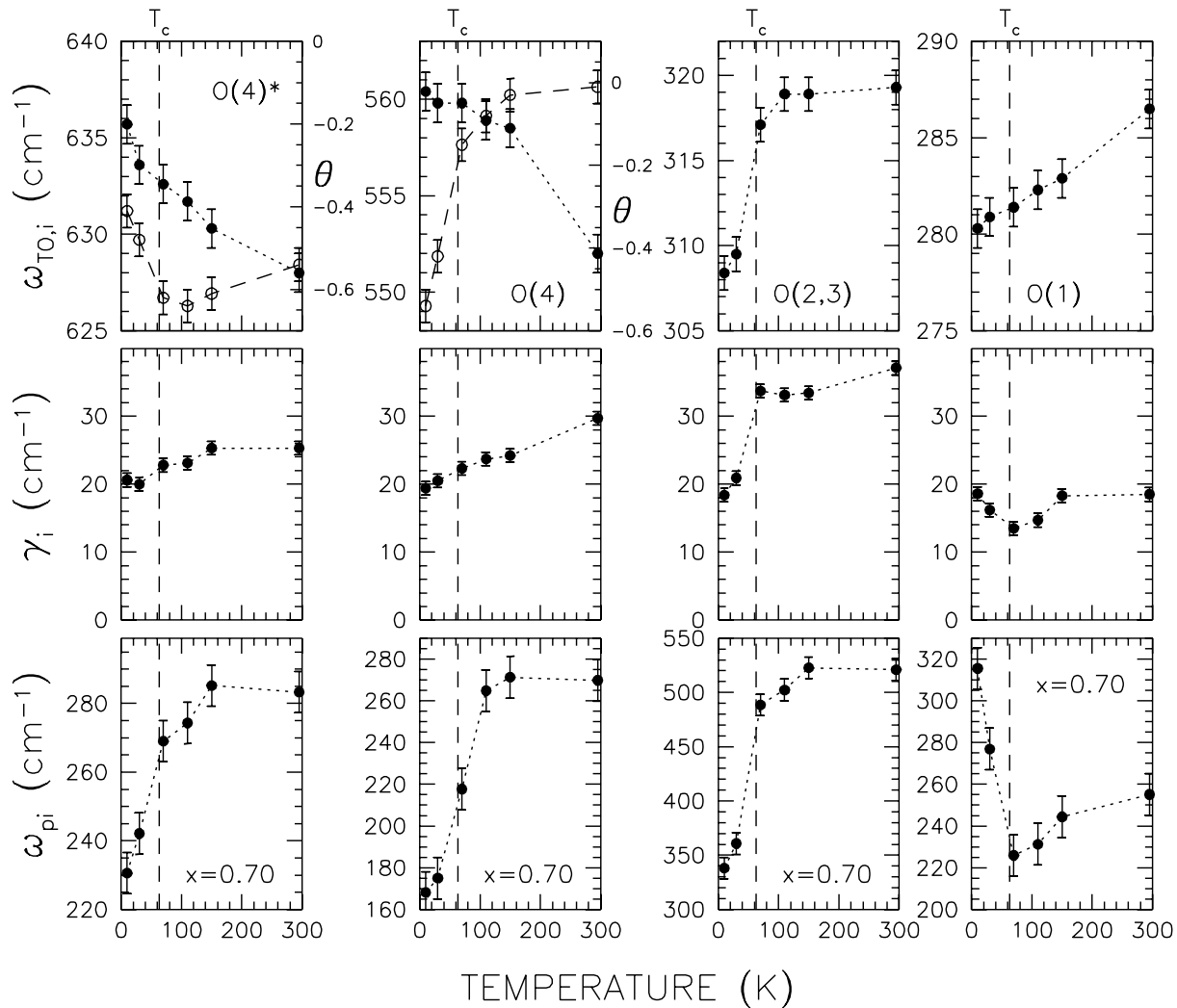
The temperature dependence of the phonon parameters for the three high-frequency modes in  $\text{YBa}_2\text{Cu}_3\text{O}_{6.95}$  is shown in Fig. 7, the results are in general agreement with previous work [12–15]. The temperature dependence for the four high-frequency modes in  $\text{YBa}_2\text{Cu}_3\text{O}_{6.7}$  is shown in Fig. 8.

In the highly-doped material, as the temperature is lowered, the phonon frequencies generally increase due to anharmonic effects, but there are exceptions, most notably the 279  $\text{cm}^{-1}$  chain-oxygen mode, which softens by  $\approx 6 \text{ cm}^{-1}$  between 295 and 10 K. The two low-frequency modes display little temperature dependence. The mode at  $\approx 153 \text{ cm}^{-1}$  hardens slightly from 153  $\text{cm}^{-1}$  at room temperature to 155  $\text{cm}^{-1}$  at 10 K. The feature at  $\approx 194 \text{ cm}^{-1}$  softens to 193  $\text{cm}^{-1}$  at 10 K; this feature shows a slight anomaly at  $T_c$ .

The three remaining modes at 279, 312 and 570  $\text{cm}^{-1}$  are all associated with oxygen vibrations, and show a more dramatic behavior. The 312 and 570  $\text{cm}^{-1}$  modes both harden with de-



**Fig. 8.** The temperature dependence of the frequency ( $\omega_{TO,i}$ ), line width ( $\gamma_i$ ), and effective plasma frequency ( $\omega_{pi}$ ) for the three oxygen-related *c*-axis modes of  $\text{YBa}_2\text{Cu}_3\text{O}_{6.70}$ , and a new fourth feature associated with the change in coordination of the Cu(1) atom, discussed in the text ( $\bullet$ ). The feature at  $\approx 285 \text{ cm}^{-1}$  softens with decreasing temperature, while the two modes at  $\approx 557 \text{ cm}^{-1}$  and  $\approx 630 \text{ cm}^{-1}$  harden smoothly; the feature at  $318 \text{ cm}^{-1}$  softens well above  $T_c$ . The three high-frequency modes begin to loose oscillator strength at  $\approx 150 \text{ K}$ , while the low-frequency mode gains in strength primarily below  $T_c$ . The rotation angles ( $\circ$ ) used to fit the two high-frequency modes are shown in the O(4)\* and O(4) panels: the  $\approx 570 \text{ cm}^{-1}$  mode becomes increasing asymmetric at low temperature, while the mode at  $\approx 610 \text{ cm}^{-1}$  becomes more symmetric.



creasing temperature, but show a discontinuity and soften below  $T_c$ . The peak at  $\approx 279 \text{ cm}^{-1}$ , originally associated with the oxygen atoms in the  $\text{CuO}_2$  planes, has been reassigned as an O(1) chain-oxygen vibration [49]; this mode softens continuously from room temperature to below  $T_c$ , and shows only a weak anomaly at  $T_c$ . However, like the  $570 \text{ cm}^{-1}$  mode, the  $279 \text{ cm}^{-1}$  mode shows some signs of hardening below  $\approx 40 \text{ K}$ . The line widths of the  $279 \text{ cm}^{-1}$  mode shows only a weak temperature dependence, while the line widths of the  $312 \text{ cm}^{-1}$  and  $570 \text{ cm}^{-1}$  modes both narrow with decreasing temperature; interestingly, none show any sign of anomalous behavior at  $T_c$ .

There is substantial variation in the oscillator strengths with temperature. The  $279 \text{ cm}^{-1}$  feature weakens considerably with

decreasing temperature, but then becomes much stronger below  $T_c$ , while the  $312 \text{ cm}^{-1}$  remains relatively constant. The O(4) mode at  $570 \text{ cm}^{-1}$  shows a rapid increase in oscillator strength at  $T_c$ , and initially becomes more asymmetric below  $T_c$ , suggesting the possibility of electron-phonon coupling (all the other modes have been fit using simple Lorentzian oscillators).

The sharpness of the phonons in the highly-doped phase, especially at the low frequency modes at  $155$  and  $279 \text{ cm}^{-1}$ , where anharmonic contribution to the width is small, is a sign of good crystal quality. In previous single crystal work with polished samples, these features were broad and weak, suggesting either poor oxygen homogeneity, or that polishing the sample has induced surface damage or altered the oxygen con-

tent of the material in the surface layers, particularly in the chains. Furthermore, an examination of the reflectance of polished *c*-axis crystals [29, 30] shows that they bear a strong similarity to the reflectance of freshly-annealed ceramics.

The 194  $\text{cm}^{-1}$  mode, according to the shell model, has its eigenvectors confined to plane region of the unit cell, and therefore does not split when alternate chains become deoxygenated, as expected.

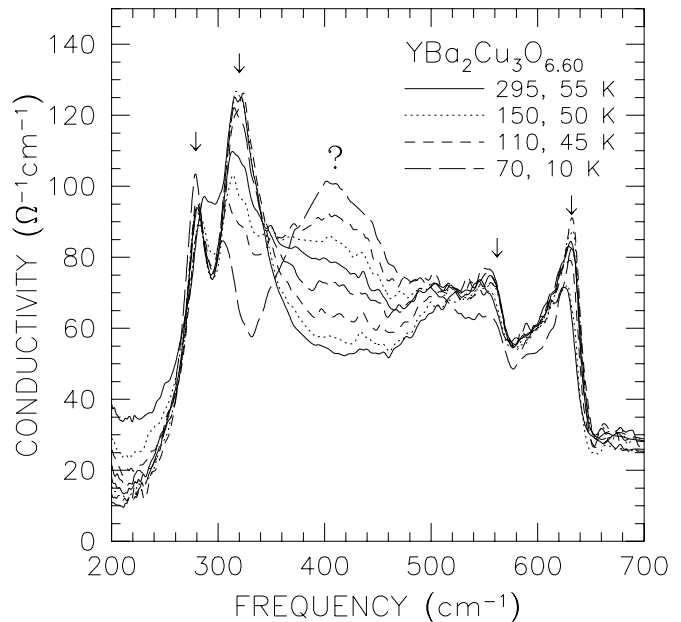
The vibrations at 279 and 312  $\text{cm}^{-1}$  in the highly-doped material harden and soften slightly with decreasing oxygen content, respectively, while also broadening considerably.

The temperature dependence of the phonon parameters in the oxygen-reduced system is dramatically different from the fully-doped case. As Fig. 8 shows for  $\text{YBa}_2\text{Cu}_3\text{O}_{6.70}$  there is now an additional mode at  $\approx 610 \text{ cm}^{-1}$ . The two apical-oxygen modes at 557 and 630  $\text{cm}^{-1}$  harden continuously with decreasing temperature, and there is no sign of a phonon anomaly at the superconducting transition temperature. As noted by Litvinchuck [14] the plane feature at  $\approx 318 \text{ cm}^{-1}$  displays very interesting behavior in that it begins to soften rapidly at  $\approx 100 \text{ K}$ , well above the  $T_c$  of 63 K. The chain feature at  $\approx 285 \text{ cm}^{-1}$  at room temperature softens to  $\approx 279 \text{ cm}^{-1}$  at low temperature, none of these features shows any discontinuity at  $T_c$ .

The behavior of the line widths is similar to that of the oscillator strengths. The strengths of the two high-frequency modes both begin to decrease at  $\approx 140 \text{ K}$ , well above  $T_c$ , as does the strength of the  $\approx 312 \text{ cm}^{-1}$  mode. This is in sharp contrast to the  $\approx 279 \text{ cm}^{-1}$  mode, which loses strength with decreasing temperature until  $T_c$ , at which point it begins to increase very rapidly, this is the only feature which shows any sensitivity to  $T_c$ .

The rapid loss of oscillator strength, particularly in the 312  $\text{cm}^{-1}$  mode in the  $x = 0.7$  material is even more noticeable in the  $x \lesssim 0.6$  materials: in the  $x = 0.5$  material, this mode appears to have completely collapsed at 10 K. The mode at  $\approx 410 \text{ cm}^{-1}$  (or combination of modes), which is very weak in the  $x = 0.95$  material, is gradually becoming stronger as the oxygen content is decreased, so that in the oxygen-reduced  $x = 0.6$  material, it is one of the dominant features. Furthermore, as the other oxygen related vibrations are losing oscillator strength, this feature is becoming correspondingly stronger, as Fig. 9 illustrates. Interestingly, optical sum-rule calculations show no net increase or decrease in the spectral weight in the 250 – 700  $\text{cm}^{-1}$  region, suggesting that the  $\approx 400 \text{ cm}^{-1}$  feature(s) is growing as a result of the transfer of oscillator strength for the other modes, but in particular from the planar copper-oxygen vibrations. The  $\approx 400 \text{ cm}^{-1}$  feature is identifiable well above  $T_c$ , and appears to grow smoothly, and without any sudden discontinuity at  $T_c$ . Furthermore, the 312  $\text{cm}^{-1}$  mode loses a great deal of oscillator strength, and begins to do so well above  $T_c$ . This suggests that while there is a major reorganization of the vibrational energy in the  $\text{CuO}_2$  planes (in the oxygen-reduced materials) as the temperature is lowered; this is a separate mechanism which may affect superconductivity in these materials, but does not appear to be affected by the superconducting transition. Interestingly, similar behavior is observed in the high-frequency copper-oxygen vibrations in cupric oxide ( $\text{CuO}$ ) near the Néel transition [50].

**Fig. 9.** The optical conductivity ( $\sigma_1$ ) of  $\text{YBa}_2\text{Cu}_3\text{O}_{6.60}$  along the *c*-axis as a function of temperature, from 200 to 700  $\text{cm}^{-1}$ , for four temperatures each above and below  $T_c$  ( $\approx 58 \text{ K}$ ). Note the dramatic loss in the strength of the feature at  $\approx 318 \text{ cm}^{-1}$ , and the corresponding increase in the feature (or group of features) centered at  $\approx 400 \text{ cm}^{-1}$ , there is no anomalous behavior near  $T_c$ .



## 4. Discussion

### 4.1. The assignment of the 610 $\text{cm}^{-1}$ phonon

The most striking change at room temperature with doping, is the growth of the 610  $\text{cm}^{-1}$  mode at the expense of the 570  $\text{cm}^{-1}$  mode. This 610  $\text{cm}^{-1}$  feature has been observed in other work and has been the subject of some controversy, since there is lack of agreement in the Raman community on the presence of the corresponding Raman mode at 600  $\text{cm}^{-1}$  [51]. In the tetrahedral  $\text{YBa}_2\text{Cu}_3\text{O}_6$  material, which lacks the O(1) oxygen chains, there is an infrared active  $A_{2u}$  mode that is observed at 645 – 650  $\text{cm}^{-1}$  which, like the 573  $\text{cm}^{-1}$  mode in the orthorhombic phase, has been identified with the O(4) oxygen [52]. In removing the chain oxygens, the four-fold coordinated Cu(1) “squares” become two-fold coordinated “sticks” [51]. The change in coordination number is accompanied by a decrease in the Cu(1)–O(4) distance, and a resulting hardening of the mode causing it to move from 570 to 610  $\text{cm}^{-1}$  [51]. The feature at  $\approx 610 \text{ cm}^{-1}$  therefore appears to be an  $A_{2u}$  mode native to the tetrahedral phase involving the O(4)–Cu(1)–O(4) sticks.

In accord with this picture we find that the oscillator strength of the two modes, at 570 and 610  $\text{cm}^{-1}$ , are nearly identical in the  $x = 0.70$  material. One expects the ideal Ortho II structure to occur at  $x = 0.50$  where there are an equal number of two and four-fold coordinated coppers with every other chain complete and every other chain empty. Also, the temperature dependence of the various mode parameters seen in Fig. 8, is nearly identical for the two modes confirming that they are located on nearly identical structural units.

Some recent Raman work does show intensity variation with

doping that parallels our work with weakening of the  $500\text{ cm}^{-1}$  line as the  $600\text{ cm}^{-1}$  line grows. The two lines reach lines of equal strength at  $500$  and  $600\text{ cm}^{-1}$  in the oxygen reduced,  $x = 0.75$  material.<sup>5</sup> Surprisingly, this is observed for the  $\mathbf{E} \parallel \mathbf{b}$  polarization.

Neutron measurements, too, find a mode in oxygen-reduced material  $x = 0.25$  at  $\mathbf{q} = \mathbf{0}$  at the frequency that is consistent with Burns' model of a high frequency mode for two-fold coordinated Cu(1).<sup>6</sup> It should also be noted that the neutron density of states determined by incoherent scattering shows a high frequency band of modes in the  $x = 0$  phase centered at  $650\text{ cm}^{-1}$ , whereas the fully-doped material has its highest branch at  $\approx 575\text{ cm}^{-1}$  [53], in accord with the notion of a set of high-frequency modes associated with the oxygen-reduced material.

#### 4.2. The $400\text{ cm}^{-1}$ mode

Another trend with decreasing doping is the appearance and growth of a broad band at  $\approx 400\text{ cm}^{-1}$  at low temperature. The band also shifts to lower frequencies as the doping is reduced. As shown in Fig. 3 and in Fig. 9, this unusual feature increases in strength in the normal state as the temperature is lowered, but does not show anomalous behavior at  $T_c$ . At very low dopings, ( $x \lesssim 0.6$ ) where this band is quite strong, it appears to gain oscillator strength at the expense of the  $312$ ,  $570$ , and  $610\text{ cm}^{-1}$  phonons. At room temperature where the phonons are sharp and well defined, the total spectral weight of the in the  $250 - 700\text{ cm}^{-1}$  range, is equal to the total spectral weight, including the  $400\text{ cm}^{-1}$  mode, at low temperature. This conservation of spectral weight is consistent with the idea that the broad peak at  $\approx 400\text{ cm}^{-1}$  is indeed a phonon.

The broad, weak mode at  $\approx 408\text{ cm}^{-1}$  has been observed in ceramic spectra of samples that appear to have been oxygen reduced [12], but this mode does not correspond to any of the calculated  $B_{1u}$  normal modes [39, 40]. A similar mode is seen in  $\text{YBa}_2\text{Cu}_4\text{O}_8$  [54] as well as in  $\text{Pb}_2\text{Sr}_2\text{RCu}_3\text{O}_8$  [55]. The common element in the three compounds is the third copper layer with two-fold coordinated oxygen bonds. In all three systems the mode appears at low temperature, growing in strength at the expense of both in-plane mode at  $312\text{ cm}^{-1}$  as well as the two apical-oxygen modes. This behavior suggests that an unusual transition is taking place that involves several atoms in the unit cell. The transition may be related to the zone-boundary mode reported on by Reichardt et al. [56]; located at approximately  $400\text{ cm}^{-1}$ , it was found to be anomalously broad and involved the "breathing" motions of the ions around the planar copper. This mode is not normally infrared active, but could become so as a result of some symmetry-breaking process. The mode involves the  $c$ -axis motion of the O(4), but not the in-plane motion of the planar oxygens. Its invocation would explain only half of our observations, namely the loss of O(4) spectral weight to the new mode through an electron-phonon process of the type discussed in connection with organic molecules by Rice [57]. In such a process a symmetry-breaking transformation allows totally-symmetric vibrations (here the zone-boundary breathing mode) to become optically active. A necessary ingredient of the process is the presence of low-lying

charge-transfer bands. The excess spectral weight and the asymmetric shape of the O(4) bands suggests that charge is being pumped at the O(4) frequency between chains and planes. This process is already active at high temperature above the formation temperature of the pseudogap. It would then not be unreasonable to anticipate that when the symmetry-breaking transition sets in at  $150\text{ K}$  in the oxygen-reduced materials, that some of this phonon-driven charge is transferred from the O(4) and plane modes, to the newly activated breathing mode. The difficulty with this argument is that it does not apply to the other two compounds where this mode is observed [54, 55].

Possibly related to this unusual phonon, polarized in the  $c$ -direction, is the phenomenon in the  $ab$ -plane conductivity reported on by Reedyk and Timusk where, in  $\text{YBa}_2\text{Cu}_3\text{O}_{6.95}$  the  $c$ -axis longitudinal modes interacted strongly with the  $ab$ -plane conductivity [58, 59]. The strongest of these was a broad phonon with the longitudinal frequency of  $440\text{ cm}^{-1}$ .

The structural transition giving rise to the phonon at  $\approx 400\text{ cm}^{-1}$  band appears to be related to the pseudogap seen at  $\approx 280\text{ cm}^{-1}$  in underdoped materials. It is strong in underdoped materials where the pseudogap is also well developed. In Zn-doped  $\text{YBa}_2\text{Cu}_4\text{O}_8$ , both the pseudogap and the  $400\text{ cm}^{-1}$  mode disappear at a  $1.7\%$  doping level.<sup>7</sup> However, there are some differences. Whereas the pseudogap appears gradually, in case of the  $x = 0.6$  material at a temperature well above  $300\text{ K}$ , the  $400\text{ cm}^{-1}$  band forms quite suddenly at  $\approx 150\text{ K}$ , as shown by the sudden loss of O(4) spectral weight below this temperature.

In summary, at this stage we have only the vaguest understanding of the nature of the phonons giving rise to the band at  $400\text{ cm}^{-1}$ , it is a feature that appears at low temperature in oxygen-reduced materials only, and gets its oscillator strength from plane-buckling and the O(4) modes. Its frequency coincides with the anomalous broad mode seen in neutron scattering, and a broad minimum in the  $ab$ -plane conductivity. Zn doping of  $\text{YBa}_2\text{Cu}_4\text{O}_8$  seems to interfere with the formation of this mode. It should be noted that it is not the result of a simple structural transition, since the *frequencies* of the modes involved change little; there is no soft-mode behavior. It appears that the transition is electronic in nature, the main result being the shift of effective charges of the phonons, not their frequencies.

## 5. Conclusions

The phonons in the oxygen-reduced systems provide a richly detailed structure. At high dopings, there are five strong  $B_{1u}$  modes. At low dopings ( $x \lesssim 0.7$ ), many of the lines split into a number of different components. In particular, the high-frequency apical-oxygen mode splits into two components; the low-frequency component is associated with the four-fold coordinated copper atom, and the high-frequency component with the two-fold coordinated copper. What is most unusual is the development of a broad feature, centered at  $\approx 400\text{ cm}^{-1}$ . This feature, attributed to a (unknown) phonon, becomes more pronounced as oxygen is removed from the system; it grows in strength as the temperature is lowered, appearing to draw oscillator strength from the vibrations at  $\approx 320$ ,  $570$  and  $610\text{ cm}^{-1}$

<sup>5</sup>M.N. Iliev. Private communication.

<sup>6</sup>N. Pyka. Unpublished results quoted in ref. 11.

<sup>7</sup>D.N. Basov, T. Timusk and B. Dabrowski (unpublished)

with decreasing temperature. This behavior develops well above  $T_c$ , and other phonon anomalies are seen in the  $x = 0.7$  material at  $\approx 150$  K.

Since these massive shifts of spectral weight to the  $400\text{ cm}^{-1}$  mode are not seen in other high-temperature superconductors, such as  $\text{La}_{2-x}\text{Sr}_x\text{CuO}_4$  and  $\text{Bi}_2\text{Sr}_2\text{CaCu}_2\text{O}_{8+z}$  [60, 61], they may be related to some special structural property of materials with a third copper layer in the charge reservoir area of the unit cell, and not a fundamental property of the cuprates.

## Acknowledgements

We would like to thank A. Bianconi, J.C. Irwin, W. Reichardt, and J. Röhler for stimulating discussions. This work was supported by the Natural Sciences and Engineering Research Council of Canada and the Canadian Institute for Advanced Research.

## References

1. M. Gurvitch and A.T. Fiory. *Phys. Rev. Lett.* **59**, 1337 (1987).
2. S.M. Martin, A.T. Fiory, R.M. Fleming, L.F. Schneemeyer, and J.V. Waszcak. *Phys. Rev. B: Condens. Matter*, **41**, 846 (1990).
3. D.B. Romero, C.D. Porter, D.B. Tanner, L. Forro, D. Mandrus, L. Mihaly, G.L. Carr, and G.P. Williams. *Phys. Rev. Lett.* **68**, 1590 (1992).
4. D.A. Bonn, R. Liang, P. Dosanjh, and W.N. Hardy. *Phys. Rev. Lett.* **68**, 2390 (1992).
5. R.C. Yu, M.B. Solomon, J.P. Lu, and W.C. Lee. *Phys. Rev. Lett.* **69**, 1431 (1992).
6. B. Batlogg, R.J. Cava, A. Jayaraman et al. *Phys. Rev. Lett.* **58**, 2333 (1987).
7. L.C. Bourne, M.F. Crommie, A. Zettl et al. *Phys. Rev. Lett.* **58**, 2337 (1987).
8. J.P. Franck, J. Jung, G.J. Salomons, W.A. Miner, M.A.-K. Mohammed, J. Chrzanowski, S. Gyax, J.C. Irwin, D.F. Mitchell, and G.I. Sproule. *Physica C (Amsterdam)*, **172**, 90 (1990).
9. M.K. Crawford, M.N. Kunchar, W.E. Farneth, E.M. McCarron III, and S.J. Poon. *Phys. Rev. B: Condens. Matter*, **41**, 282 (1990).
10. A.P. Litvinchuk, C. Thomsen and M. Cardona. *In Physical Properties of High-Temperature Superconductors IV. Edited by D.M. Ginsberg.* World Scientific, Singapore. 1992, pp. 375–470.
11. L. Pintschovius and W. Reichardt. *In Physical Properties of High-Temperature Superconductors IV. Edited by D.M. Ginsberg.* World Scientific, Singapore. 1992, pp. 295–374.
12. D.A. Bonn, J.E. Greedan, C.V. Stager, T. Timusk, M.G. Doss, S.L. Herr, K. Kamarás, D.B. Tanner. *Phys. Rev. Lett.* **58**, 2249 (1987).
13. A. Wittlin, R. Liu, M. Cardona, L. Genzel, M. Bauer, W. König, and E. Garcia-Alvarado. *Solid State Commun.* **64**, 477 (1987).
14. A.P. Litvinchuk, C. Thomsen, and M. Cardona. *Solid State Commun.* **83**, 343 (1992).
15. R.M. Macfarlane, H. Rosen, and H. Seki. *Solid State Commun.* **63**, 831 (1987).
16. B. Friedl, C. Thomsen, and M. Cardona. *Phys. Rev. Lett.* **76**, 1107 (1990).
17. E. Altendorf, X.K. Chen, J.C. Irwin, R. Liang, and W.N. Hardy. *Phys. Rev. B: Condens. Matter*, **47**, 8140 (1993).
18. H.A. Mook, M. Mostoller, J.A. Harvey, N.W. Hill, B.C. Chakoumakos, and B.C. Stales. *Phys. Rev. Lett.* **69**, 2712 (1990).
19. N. Pyka, W. Reichardt, L. Pintschovius, G. Engel, J. Rossat-Mignod, and J.Y. Henry. *Phys. Rev. Lett.* **70**, 1457 (1993).
20. T. Timusk, C.D. Porter and D.B. Tanner. *Phys. Rev. Lett.* **66**, 663 (1991).
21. M. Reedyk and T. Timusk. *Phys. Rev. Lett.* **69**, 2705 (1992).
22. J. Schützmann, S. Tajima, S. Miyamoto, and S. Tanaka. *Phys. Rev. Lett.* **73**, 174 (1994).
23. S.L. Cooper and K.E. Gray. *In Physical Properties of High-Temperature Superconductors IV. Edited by D.M. Ginsberg.* World Scientific, Singapore. 1992, pp. 61–188.
24. D.A. Bonn, A.H. O'Reilly, J.E. Greedan, C.V. Stager, T. Timusk, K. Kamarás, and D.B. Tanner. *Phys. Rev. B: Condens. Matter*, **37**, 1574 (1988).
25. L. Genzel, A. Wittlin, M. Bauer, M. Cardona, E. Schönherr, and A. Simon. *Phys. Rev. B: Condens. Matter*, **40**, 2170 (1989).
26. R. Feile. *Physica C (Amsterdam)* **159**, 1 (1989).
27. I. Bozovic, D. Kirillov, A. Kapitulnik, K. Char, M.R. Hahn, M.R. Beasley, T.H. Geballe, Y.H. Kim, and A.J. Heeger. *Phys. Rev. Lett.* **59**, 2219 (1987).
28. R.T. Collins, Z. Schlesinger, F. Holtzberg, and C. Feild. *Phys. Rev. Lett.* **63**, 422 (1989).
29. B. Koch, M. Dürler, Th. Wolf, G. Roth, G. Zachmann. *Solid State Commun.* **71**, 495 (1990).
30. B. Koch, M. Dürler, H.P. Geserich, Th. Wolf, G. Roth, and G. Zachmann. *In Electronic Properties of High- $T_c$  Superconductors and Related Compounds. Edited by H. Kuzmany, M. Mehrig, J. Fink.* Springer Series in Solid State Sciences, Vol. 99, Springer-Verlag, Berlin-Heidelberg. 1990, p. 290.
31. S.L. Cooper, P. Nyhus, D. Reznik, M.V. Klein, W.C. Lee, D.M. Ginsberg, B.W. Veal, A.P. Paulikas, and B. Dabrowski. *Phys. Rev. Lett.* **70**, 1533 (1993); S.L. Cooper, D. Reznik, A. Kotz, M.A. Karlow, R. Liu, M.V. Klein, W.C. Lee, J. Giapintzakis, D.M. Ginsberg, B.W. Veal, and A.P. Paulikas. *Phys. Rev. B: Condens. Matter*, **47**, 8233 (1993).
32. C.C. Homes, T. Timusk, R. Liang, D.A. Bonn, and W.N. Hardy. *Phys. Rev. Lett.* **71**, 1645 (1993).
33. T. Timusk, C.C. Homes, N. Cao, and D.B. Tanner. *In Proceedings of the Beijing International Conference on High-Temperature Superconductivity (BHTSC '92). Edited by Z.Z. Gan, S.S. Xie, and Z.X. Zhao.* World Scientific, Singapore. 1993, pp. 408–415.
34. T. Timusk, C.C. Homes, and W. Reichardt. *In International Workshop on the Anharmonic Properties of High- $T_c$  Cuprates. Bled, Slovenia. Edited by G. Ruani.* World Scientific, Singapore. 1995.
35. R. Liang, P. Dosanjh, D.A. Bonn, D.J. Baar, J.F. Carolan, and W.N. Hardy. *Physica C (Amsterdam)* **195**, 51 (1992).
36. C.C. Homes, M.A. Reedyk, D.A. Crandles, and T. Timusk. *Appl. Opt.* **32**, 2976 (1993).
37. V.M. Burlkov, K.V. Kraiskaya, A.G. Mitko, E.J. Firsov. *Phys. Lett. A* **142**, 514 (1989); V.M. Burlkov, N.Y. Boldyrev, I. Gaiduk, K.V. Kraiskaya, A.G. Mitko. *Pis'ma Zh. Eksp. Teor. Fiz.* **52**, 733 (1990); *JETP Lett.* **52**, 91 (1990).
38. U. Fano. *Phys. Rev.* **124**, 1866 (1961).
39. F.E. Bates and J.E. Eldridge. *Solid State Commun.* **64**, 1435 (1987).
40. W. Kress, U. Schröder, J. Prade, A.D. Kulkarni, and F.W. de Wette. *Phys. Rev. B: Condens. Matter*, **39**, 2906 (1988).

41. R.A. Cowley, W. Cochran, B.N. Brockhouse, and A.D.B. Woods. *Phys. Rev.* **131**, 1030 (1963).
42. M.K. Crawford, W.E. Farneth, E.M. McCarron III, and R.K. Bordia. *Phys. Rev. B: Condens. Matter*, **38**, 11382 (1988).
43. V.A. Maroni, T.O. Brun, M. Grimsditch, and C.-K. Loong. *Phys. Rev. B: Condens. Matter*, **39**, 4127 (1989).
44. V. Hadjiev, C. Thomsen, J. Kircher, and M. Cardona. *Phys. Rev. B: Condens. Matter*, **47**, 9148 (1993).
45. R.M. Hazen. *In Physical Properties of High-Temperature Superconductors II. Edited by D.M. Ginsberg.* World Scientific, Singapore. 1990, p. 121–198.
46. S.D. Conradson, I. Raistrick, and A.R. Bishop. *Science (Washington)* **248**, 1934 (1990).
47. J. Mustre de Leon, S.D. Conradson, I. Batisitić, and A.R. Bishop. *Phys. Rev. Lett.* **65**, 1675 (1990).
48. E.A. Stern, M. Quian, Y. Yakoby, S.M. Heal, and H. Maeda. *Physica C (Amsterdam)* **209**, 331 (1993).
49. C. Thomsen, A.P. Litvinchuk, E. Schönherr, and M. Cardona. *Phys. Rev. B: Condens. Matter*, **45**, 8145 (1992).
50. C.C. Homes, M. Ziaei, B.P. Clayman, J.C. Irwin, and J.P. Franck. *Phys. Rev. B: Condens. Matter*, **51**, 3140 (1995).
51. G. Burns, F.H. Dacol, C. Feild, and F. Holtzberg. *Physica C (Amsterdam)* **181**, 37 (1991).
52. G. Burns, F.H. Dacol, P.P. Freitas, W. König, and T.S. Plaskett. *Phys. Rev. Lett.* **37**, 5171 (1988).
53. B. Renker, F. Gompf, E. Gering, D. Ewert, H. Rietschek, and A. Dianoux. *Z. Phys. B: Condens. Matter*, **73**, 309 (1988).
54. D.N. Basov, T. Timusk, B. Dabrowski, and J.D. Jorgensen. *Phys. Rev. B: Condens. Matter*, **50**, 3511 (1994).
55. M. Reedyk, T. Timusk, J.S. Xue, and J.E. Greedan. *Phys. Rev. B: Condens. Matter*, **49**, 15984 (1994).
56. W. Reichardt, N. Pyka, L. Pintschovius, B. Hennion, and G. Collin. *Physica C (Amsterdam)* **162–164**, 464 (1989).
57. M.J. Rice, V.M. Yartsev, and C.S. Jacobsen. *Phys. Rev. B: Condens. Matter*, **21**, 3437 (1980).
58. T. Timusk, C.D. Porter and D.B. Tanner. *Phys. Rev. Lett.* **66**, 663 (1991).
59. M. Reedyk and T. Timusk. *Phys. Rev. Lett.* **69**, 2705 (1992).
60. K. Tamasaku, T. Itoh, H. Takagi, and S. Uchida. *Phys. Rev. Lett.* **72**, 3088 (1994).
61. S. Tajima, G.D. Gu, S. Miyamoto, A. Odagawa, and N. Koshizuka. *Phys. Rev. B: Condens. Matter*, **48**, 16164 (1993).