

Orthogonal Polynomials from Hermitian Matrices

Satoru Odake^a and Ryu Sasaki^b

^a Department of Physics, Shinshu University,
Matsumoto 390-8621, Japan

^b Yukawa Institute for Theoretical Physics,
Kyoto University, Kyoto 606-8502, Japan

Abstract

A unified theory of orthogonal polynomials of a discrete variable is presented through the eigenvalue problem of hermitian matrices of finite or infinite dimensions. It can be considered as a matrix version of exactly solvable Schrödinger equations. The hermitian matrices (factorisable Hamiltonians) are real symmetric tri-diagonal (Jacobi) matrices corresponding to second order difference equations. By solving the eigenvalue problem in two different ways, the duality relation of the eigenpolynomials and their dual polynomials is explicitly established. Through the techniques of exact Heisenberg operator solution and shape invariance, various quantities, the two types of eigenvalues (the eigenvalues and the sinusoidal coordinates), the coefficients of the three term recurrence, the normalisation measures and the normalisation constants etc. are determined explicitly.

1 Introduction

Due to the long history and rich and diverse applications, there are many different ways to introduce orthogonal polynomials [1]. Among many types and kinds of orthogonal polynomials, we will focus in this paper on the so-called orthogonal polynomials of a discrete variable [2, 3, 4, 5], whose orthogonality measures are concentrated on discrete points, either finite or infinite in number. We will present a unified theory of these polynomials based on the

eigenvalue problem (2.1) of a special class of hermitian matrices, which are real symmetric *tri-diagonal* (Jacobi) matrices (2.4), (2.11). Since the spectrum of a Jacobi matrix is simple, the orthogonality of eigenvectors is guaranteed. The eigenvalue problem of a Jacobi matrix can be considered as a difference equation version of the Schrödinger equation for one degree of freedom system, the most basic equation of quantum mechanics [6]. Various classical orthogonal polynomials (the Hermite, Laguerre, Jacobi and their restrictions) have appeared as the eigenfunctions of exactly solvable quantum mechanics [6].

Roughly speaking, our line of arguments is a deformation (discretisation) of the main trend of the twentieth century mathematical physics; to pursue the parallelism between the matrix eigenvalue problem and the ordinary differential equations of the Sturm-Liouville type, with *differential* equations replaced by *difference* equations.

We apply many ideas and methods for solving Schrödinger equations: Crum's theorem [7], the factorisation method [8] (or the so-called supersymmetric quantum mechanics [9]), the method of exact Heisenberg operator solutions and the creation/annihilation operators [10, 11] together with symmetries (shape invariance (4.2) [12, 13] and closure relation (4.31) [10]) to elucidate the universal structure of the orthogonal polynomials of a discrete variable. The Jacobi matrix eigenvalue problem $\mathcal{H}\psi = \mathcal{E}\psi$ (2.1) can be solved in a different way, through the three term recurrence relations for the dual polynomial in \mathcal{E} . Combining the *duality*, which can be stated at many different levels, with the above mentioned solution techniques, various quantities of the orthogonal polynomials, the coefficients of the three term recurrence, the orthogonality measure and the normalisation constants, the difference equations for the dual polynomials, etc. are given explicitly in an elementary manner. We stress that duality is established equally for the finite and infinite dimensional cases. All the examples discussed in this paper are taken from the review of Koekoek and Swarttouw [14] and they are known to satisfy the closure relation (4.31). The orthogonal polynomials with the Jackson integral type measure, *e.g.* the Big q -Jacobi polynomial, etc. will need different treatment and will be discussed elsewhere. As an inverse step to characterise the orthogonal polynomials of a discrete variable, we solve the closure relation algebraically to determine the possible forms of the Jacobi matrices.

It should be stressed that the present duality has far richer contents than the well-known duality of the column and row eigenvectors. That is, in the (hermitian) eigenvalues problems, the orthogonality of the complete set of column eigenvectors automatically implies that of the

row eigenvectors. In the present case, the column eigenvectors are polynomials in *sinusoidal coordinate* (2.28) $\eta(x)$ and the row eigenvectors are polynomials in the *eigenvalues* (2.28) $\mathcal{E}(n)$, and the duality is formulated as the equality of these two polynomials on the integer lattice points (3.10).

The formulation of the unified theory of orthogonal polynomials based on the eigenvalue problems of Jacobi matrices is new. All the quantities and formulas are derived from the two input functions $B(x)$ and $D(x)$ (2.4), (2.5). However, a good part of the explicit formulas for specific polynomials explored in section five is already known individually in the existing theories of orthogonal polynomials based on various settings.

This paper is organised as follows. In section two, the general form of the hermitian matrices (or the Hamiltonians) is given and the solution procedure to arrive at the orthogonal polynomials is outlined. They are polynomials in the *sinusoidal coordinate* (2.28) which plays an important role in exactly solvable quantum mechanics [10, 11]. The orthogonality measure is given by the solution $\phi_0(x)$ of the factorised equation (2.17). The ϕ_0 is the ground state wavefunction. In section three the *dual polynomials* (in the eigenvalue $\mathcal{E}(n)$) are introduced by a different solution method (the three term recurrence (3.2)) of the same eigenvalue problem (2.1). The duality at many different levels, *e.g.* $x \leftrightarrow n$, $\eta(x) \leftrightarrow \mathcal{E}(n)$, is displayed (3.14)–(3.18). Section four explores the underlying symmetry properties of the Hamiltonian, shape invariance §4.1, closure relation §4.2. The general formulas of the coefficients A_n and C_n of the three term recurrence are derived. The forms of the sinusoidal coordinate and the diagonal elements of the Hamiltonian are determined in §4.3 from the closure relation. The dual closure relation is discussed in §4.4. Section five provides the explicit forms of various quantities derived in previous sections for the concrete examples of the orthogonal polynomials in the orders given in the review [14]. Some orthogonal polynomials not listed in the review [14] are also discussed in some detail in §5.3. Appendix A gives the possible forms of the Jacobi matrices (the Hamiltonians) (A.15)–(A.16) as the solutions of the closure relation. The results show the boundedness and unboundedness of the Hamiltonians and eigenvalues clearly. Appendix B provides the collection of the definitions of basic symbols and functions for self-containedness. Throughout this paper we use the parameter q in the range $0 < q < 1$.

2 Hamiltonian

The starting point is the eigenvalue problem for a hermitian matrix \mathcal{H}

$$\mathcal{H}\psi = \mathcal{E}\psi. \quad (2.1)$$

The rows and columns of \mathcal{H} are indexed by non-negative integers x and y , either finite

$$\text{case (I)} \quad x, y = 0, 1, 2, \dots, N, \quad (2.2)$$

or infinite:

$$\text{case (II)} \quad x, y = 0, 1, 2, \dots \quad (2.3)$$

Let us call the hermitian matrix \mathcal{H} *Hamiltonian*, since the eigenvalue problem (2.1) can be considered as the Schrödinger equation which is a *difference* equation, instead of differential in ordinary quantum mechanics. The Hamiltonian we consider in this paper has a general form

$$\mathcal{H} \stackrel{\text{def}}{=} -\sqrt{B(x)} e^{\partial} \sqrt{D(x)} - \sqrt{D(x)} e^{-\partial} \sqrt{B(x)} + B(x) + D(x), \quad (2.4)$$

in which the two functions $B(x)$ and $D(x)$ are real and *positive* but vanish at the boundary:

$$B(x) > 0, \quad D(x) > 0, \quad D(0) = 0; \quad B(N) = 0 \text{ for case (I)}. \quad (2.5)$$

In (2.4) $\partial \stackrel{\text{def}}{=} \frac{d}{dx}$ is the differential operator and its exponentiation gives finite shifts:

$$e^{\pm\partial}\psi(x) = \psi(x \pm 1).$$

Thus, in fact, the eigenvalue problem (2.1) can be written as a difference equation on non-negative integer points:

$$\begin{aligned} (B(x) + D(x))\psi(x) - \sqrt{B(x)D(x+1)}\psi(x+1) - \sqrt{B(x-1)D(x)}\psi(x-1) &= \mathcal{E}\psi(x), \\ x = 0, 1, \dots, (N), \dots \end{aligned} \quad (2.6)$$

The boundary condition $D(0) = 0$ is necessary for the term $\psi(-1)$ does not appear, and $B(N) = 0$ is necessary for the term $\psi(N+1)$ does not appear in the finite dimensional matrix case. It is also easy to see that $\psi(0)$ does not vanish for any eigenvector

$$\psi(0) \neq 0, \quad (\psi(N) \neq 0 \text{ for case (I)}). \quad (2.7)$$

Let us remark that the (q -)Askey-scheme of hypergeometric orthogonal polynomials with absolutely continuous weight functions, for example, the Meixner-Pollaczek, the Wilson and the Askey-Wilson polynomials, have already been well understood as the eigenvalue problem of hermitian (self-adjoint) operators \mathcal{H} [13, 10]. The generic forms of the Hamiltonian \mathcal{H} are:

$$\mathcal{H} \stackrel{\text{def}}{=} \sqrt{V(x)} e^{-i\partial} \sqrt{V(x)^*} + \sqrt{V(x)^*} e^{i\partial} \sqrt{V(x)} - V(x) - V(x)^*, \quad (2.8)$$

$$\mathcal{H} \stackrel{\text{def}}{=} \sqrt{V(x)} q^{-i\partial} \sqrt{V(x)^*} + \sqrt{V(x)^*} q^{i\partial} \sqrt{V(x)} - V(x) - V(x)^*. \quad (2.9)$$

These also give difference Schrödinger equations, but the shift is in the *imaginary* direction, instead of real in the case (2.6).

Although the Hamiltonian \mathcal{H} (2.4) is presented in a difference operator form, it is in fact a real symmetric *tridiagonal* (Jacobi) matrix:

$$\mathcal{H} = (\mathcal{H}_{x,y}), \quad \mathcal{H}_{x,y} = \mathcal{H}_{y,x}, \quad (2.10)$$

$$\mathcal{H}_{x,y} = -\sqrt{B(x)D(x+1)} \delta_{x+1,y} - \sqrt{B(x-1)D(x)} \delta_{x-1,y} + (B(x) + D(x)) \delta_{x,y}. \quad (2.11)$$

It is well-known that the spectrum of a Jacobi matrix is simple; that is, there is no degeneracy of the eigenvalues

$$\mathcal{E}(0) < \mathcal{E}(1) < \mathcal{E}(2) < \dots .$$

The first step for solving the eigenvalue equation (2.1) is to rewrite the Hamiltonian (2.4) in a factorised form:

$$\mathcal{H} = \mathcal{A}^\dagger \mathcal{A}, \quad (2.12)$$

$$\mathcal{A}^\dagger = \sqrt{B(x)} - \sqrt{D(x)} e^{-\partial}, \quad (2.13)$$

$$\mathcal{A} = \sqrt{B(x)} - e^\partial \sqrt{D(x)}, \quad (2.14)$$

with the forward (backward) shift operator \mathcal{A} (\mathcal{A}^\dagger) being hermitian conjugate of each other. Or in the matrix form, the tridiagonal (Jacobi) matrix (2.11) is decomposed into a product of \mathcal{A}^\dagger and \mathcal{A}

$$(\mathcal{A}^\dagger)_{x,y} = \sqrt{B(x)} \delta_{x,y} - \sqrt{D(x)} \delta_{x-1,y}, \quad (2.15)$$

$$\mathcal{A}_{x,y} = \sqrt{B(x)} \delta_{x,y} - \sqrt{D(x+1)} \delta_{x+1,y}, \quad (2.16)$$

with \mathcal{A}^\dagger (\mathcal{A}) having non-vanishing entries on the diagonal and one below (above) the diagonal. Throughout this paper we adopt the (standard) convention that the $(0, 0)$ element of the matrix is at the upper left corner. The factorisation also means that the Hamiltonian (2.4) is semi-positive definite. In fact, the ground state (the lowest eigenvector) $\phi_0(x)$ is annihilated by \mathcal{A} and thus it is a zero-mode of the Hamiltonian:

$$\mathcal{A}\phi_0(x) = 0 \implies \mathcal{H}\phi_0(x) = 0, \quad \mathcal{E}(0) = 0. \quad (2.17)$$

As we will see shortly, $\phi_0(x)^2$ is the orthogonality measure for the eigenpolynomials (2.33). This is the same situation as in the ordinary quantum mechanics and also in the above mentioned theory (2.8)-(2.9) for the Wilson and Askey-Wilson polynomials, etc. [13, 10]. The above equation (2.17), being a two term recurrence relation,

$$\frac{\phi_0(x+1)}{\phi_0(x)} = \sqrt{\frac{B(x)}{D(x+1)}}, \quad (2.18)$$

can be solved elementarily with the boundary (initial) condition

$$\phi_0(0) = 1, \quad (2.19)$$

$$\phi_0(x) = \sqrt{\prod_{y=0}^{x-1} \frac{B(y)}{D(y+1)}}, \quad x = 1, 2, \dots \quad (2.20)$$

With the standard convention $\prod_{k=n}^{n-1} * = 1$, the expression (2.20) is valid for $x = 0$ also. For the infinite matrix case (II), the requirement of the finite ℓ^2 norm of the eigenvectors

$$\sum_{x=0}^{\infty} \phi_0(x)^2 = \sum_{x=0}^{\infty} \prod_{y=0}^{x-1} \frac{B(y)}{D(y+1)} < \infty \quad (2.21)$$

imposes constraints on the asymptotic behaviours of $B(x)$ and $D(x)$.

Next let us determine the other eigenvectors (the excited states) in a factored form

$$\psi(x) = \phi_0(x)v(x). \quad (2.22)$$

The eigenvalue problem (2.1) for ψ is rewritten to that for v

$$\tilde{\mathcal{H}}v(x) = \mathcal{E}v(x), \quad (2.23)$$

in which $\tilde{\mathcal{H}}$ is the similarity transformed Hamiltonian in terms of the ground state wavefunction $\phi_0(x)$:

$$\tilde{\mathcal{H}} \stackrel{\text{def}}{=} \phi_0(x)^{-1} \circ \mathcal{H} \circ \phi_0(x)$$

$$= B(x)(1 - e^\theta) + D(x)(1 - e^{-\theta}). \quad (2.24)$$

It gives rise to a simple difference equation for $v(x)$ (2.23). Obviously a constant is a solution

$$\tilde{\mathcal{H}}1 = 0. \quad (2.25)$$

In the matrix form $\tilde{\mathcal{H}}$ is another tridiagonal matrix

$$\tilde{\mathcal{H}} = (\tilde{\mathcal{H}}_{x,y}), \quad \tilde{\mathcal{H}}_{x,y} = B(x)(\delta_{x,y} - \delta_{x+1,y}) + D(x)(\delta_{x,y} - \delta_{x-1,y}). \quad (2.26)$$

For all the cases discussed in this paper, the similarity transformed Hamiltonian $\tilde{\mathcal{H}}$ is *lower triangular* with respect to the special basis

$$1, \eta(x), \eta(x)^2, \dots, \eta(x)^n, \dots, \quad (2.27)$$

spanned by the *sinusoidal coordinate* $\eta(x)$ [13, 10]:

$$\tilde{\mathcal{H}}\eta(x)^n = \mathcal{E}(n)\eta(x)^n + \text{lower orders in } \eta(x). \quad (2.28)$$

Here $\mathcal{E}(n)$ is the corresponding eigenvalue. This property is preserved under the affine transformation of $\eta(x)$, $\eta(x) \rightarrow a\eta(x) + b$ (a, b : constants). As will be shown shortly, the functional form of the sinusoidal coordinate is characterised by the *closure relation* (4.31), (4.32) up to a multiplicative and an additive constant. We choose the additive constant to achieve

$$\eta(0) = 0. \quad (2.29)$$

The lower triangularity (2.28) implies that the eigenvectors can be obtained as polynomials in $\eta(x)$ up to normalisation:

$$\begin{aligned} \tilde{\mathcal{H}}P_n(\eta(x)) &= \mathcal{E}(n)P_n(\eta(x)), \\ \mathcal{H}\phi_n(x) &= \mathcal{E}(n)\phi_n(x), \quad \phi_n(x) = \phi_0(x)P_n(\eta(x)), \end{aligned} \quad n = 0, 1, 2, \dots, (N), \dots \quad (2.30)$$

We choose the simple normalisation of $P_n(\eta(x))$

$$P_n(0) = 1, \quad n = 0, 1, 2, \dots, (N), \dots, \quad (2.31)$$

which is consistent with (2.7) and (2.29). Note that this also means

$$P_0 = 1. \quad (2.32)$$

Thus the eigenvectors are completely specified and the eigenvalue problem (2.1) is completely solved. The orthogonality of the eigenvectors of the Jacobi matrix \mathcal{H} (2.4), (2.11) implies

$$\sum_x \phi_0(x)^2 P_n(\eta(x)) P_m(\eta(x)) = \frac{1}{d_n^2} \delta_{n,m}, \quad (2.33)$$

in which the normalisation constants $\{d_n > 0\}$ are to be calculated. In terms of the normalised eigenvectors

$$\hat{\phi}_0(x) = d_0 \phi_0(x), \quad \hat{\phi}_n(x) = d_n \phi_0(x) P_n(\eta(x)) = \hat{\phi}_0(x) \frac{d_n}{d_0} P_n(\eta(x)), \quad (2.34)$$

the orthonormality relation reads

$$\sum_x d_n \phi_0(x) P_n(\eta(x)) \cdot d_m \phi_0(x) P_m(\eta(x)) = \delta_{n,m}. \quad (2.35)$$

This in turn implies the completeness relation

$$\sum_n d_n \phi_0(x) P_n(\eta(x)) \cdot d_n \phi_0(y) P_n(\eta(y)) = \delta_{x,y}, \quad (2.36)$$

which states the simplest duality that the row eigenvectors are also orthogonal. When written slightly differently

$$\sum_n d_n^2 P_n(\eta(x)) P_n(\eta(y)) = \frac{1}{\phi_0(x)^2} \delta_{x,y}, \quad (2.37)$$

it displays the duality with (2.33) explicitly.

We have established that $P_n(\eta)$ is an orthogonal polynomial of a discrete variable with the discrete measure $\{\phi_0(x)^2\}$ given explicitly as (2.20), (2.21). They satisfy the three term recurrence relation

$$\eta P_n(\eta) = A_n P_{n+1}(\eta) + B_n P_n(\eta) + C_n P_{n-1}(\eta) \quad (2.38)$$

with coefficients A_n , B_n and C_n ($C_0 = 0$). The boundary condition $P_n(0) = 1$ (2.31) implies

$$B_n = -(A_n + C_n). \quad (2.39)$$

In section 4 we will derive the general formulas of the coefficients A_n and C_n (4.49), (4.50) from the input functions $B(x)$ and $D(x)$ with the help of various quantities implied by the closure relation. This will provide complete specification of the polynomial $P_n(\eta)$.

In the next section, we will show that the explicit form of the normalisation constants $\{d_n\}$ (2.33) can be easily obtained as the components of the ground state wavefunction of the *dual* polynomial as seen clearly from (2.33) and (2.37).

3 Dual Polynomials

The *dual* polynomial is an important concept in the theory of orthogonal polynomials of a discrete variable [2, 3, 4, 5, 14]. Here we show that the dual polynomial arises naturally as the solution of the original eigenvalue problem (2.1) or (2.23) obtained in a different way. Let us rewrite the similarity transformed eigenvalue problem $\tilde{\mathcal{H}}v(x) = \mathcal{E}v(x)$ (2.23) into an explicit matrix form with the change of the notation $v(x) \rightarrow {}^t(Q_0, Q_1, \dots, Q_x, \dots)$

$$\sum_y \tilde{\mathcal{H}}_{x,y} Q_y = \mathcal{E}Q_x, \quad x = 0, 1, \dots, (N), \dots \quad (3.1)$$

Because of the tridiagonality of $\tilde{\mathcal{H}}$, it is in fact a three term recurrence relation for Q_x as a polynomial in \mathcal{E} :

$$\begin{aligned} \mathcal{E}Q_x(\mathcal{E}) &= B(x)(Q_x(\mathcal{E}) - Q_{x+1}(\mathcal{E})) + D(x)(Q_x(\mathcal{E}) - Q_{x-1}(\mathcal{E})), \\ &x = 0, 1, \dots, (N), \dots \end{aligned} \quad (3.2)$$

Starting with the boundary (initial) condition

$$Q_0 = 1, \quad (3.3)$$

which is consistent with (2.7), we determine $Q_x(\mathcal{E})$ as a degree x polynomial in \mathcal{E} . It is easy to see

$$Q_x(0) = 1, \quad x = 0, 1, \dots, (N), \dots \quad (3.4)$$

When \mathcal{E} is replaced by the actual value of the n -th eigenvalue $\mathcal{E}(n)$ (2.28) in $Q_x(\mathcal{E})$, we obtain the explicit form of the eigenvector

$$\sum_y \tilde{\mathcal{H}}_{x,y} Q_y(\mathcal{E}(n)) = \mathcal{E}(n)Q_x(\mathcal{E}(n)), \quad x = 0, 1, \dots, (N), \dots \quad (3.5)$$

In the finite dimensional case (I), $\{Q_0(\mathcal{E}), \dots, Q_N(\mathcal{E})\}$ are determined by (3.2) for $x = 0, \dots, N - 1$. The last equation

$$\mathcal{E}Q_N(\mathcal{E}) = D(N)(Q_N(\mathcal{E}) - Q_{N-1}(\mathcal{E})), \quad (3.6)$$

is the degree $N + 1$ algebraic equation (characteristic equation) for the determination of all the eigenvalues $\{\mathcal{E}(n)\}$. This is a familiar situation encountered in quasi-exactly solvable quantum mechanics [15, 16, 17, 18]. In the next section we will provide further two

independent algebraic methods for the determination of the eigenvalues $\{\mathcal{E}(n)\}$ based on the shape-invariance [13] and the exact Heisenberg operator solution [10], both of which are applicable to the finite (I) as well as the infinite dimensional case (II).

We now have two expressions (polynomials) for the eigenvectors of the problem (2.1) belonging to the eigenvalue $\mathcal{E}(n)$; $P_n(\eta(x))$ and $Q_x(\mathcal{E}(n))$. Due to the simplicity of the spectrum of the Jacobi matrix, they must be equal up to a multiplicative factor α_n ,

$$P_n(\eta(x)) = \alpha_n Q_x(\mathcal{E}(n)), \quad x = 0, 1, \dots, (N), \dots, \quad (3.7)$$

which turns out to be unity because of the boundary (initial) condition at $x = 0$ (2.29), (2.31), (3.3) (or at $n = 0$ (2.17), (2.32), (3.4));

$$P_n(\eta(0)) = P_n(0) = 1 = Q_0(\mathcal{E}(n)), \quad n = 0, 1, \dots, (N), \dots, \quad (3.8)$$

$$P_0(\eta(x)) = 1 = Q_x(0) = Q_x(\mathcal{E}(0)), \quad x = 0, 1, \dots, (N), \dots \quad (3.9)$$

We have established that two polynomials, $\{P_n(\eta)\}$ and its *dual* polynomial $\{Q_x(\mathcal{E})\}$, coincides at the integer lattice points:

$$P_n(\eta(x)) = Q_x(\mathcal{E}(n)), \quad n = 0, 1, \dots, (N), \dots, \quad x = 0, 1, \dots, (N), \dots \quad (3.10)$$

The completeness relation (2.37)

$$\sum_n d_n^2 P_n(\eta(x)) P_n(\eta(y)) = \sum_n d_n^2 Q_x(\mathcal{E}(n)) Q_y(\mathcal{E}(n)) = \frac{1}{\phi_0(x)^2} \delta_{x,y}, \quad (3.11)$$

is now understood as the orthogonality relation of the dual polynomial $Q_x(\mathcal{E})$, and the previous normalisation constant d_n^2 is now the orthogonality measure.

The real symmetric (hermitian) matrix $\mathcal{H}_{x,y}$ (2.11) can be expressed in terms of the complete set of the eigenvalues and the corresponding normalised eigenvectors

$$\mathcal{H}_{x,y} = \sum_n \mathcal{E}(n) \hat{\phi}_n(x) \hat{\phi}_n(y), \quad (3.12)$$

$$\hat{\phi}_n(x) = d_n \phi_0(x) P_n(\eta(x)) = d_n \phi_0(x) Q_x(\mathcal{E}(n)). \quad (3.13)$$

The very fact that it is tridiagonal can be easily verified by using the difference equation for the polynomial $P_n(\eta(x))$ or the three term recurrence relation for $Q_x(\mathcal{E}(n))$.

Here is the list of the dual correspondence:

$$x \leftrightarrow n, \quad (3.14)$$

$$\eta(x) \leftrightarrow \mathcal{E}(n), \quad \eta(0) = 0 \leftrightarrow \mathcal{E}(0) = 0, \quad (3.15)$$

$$B(x) \leftrightarrow -A_n, \quad (3.16)$$

$$D(x) \leftrightarrow -C_n, \quad (3.17)$$

$$\frac{\phi_0(x)}{\phi_0(0)} \leftrightarrow \frac{d_n}{d_0}. \quad (3.18)$$

In the last expression, we inserted $\phi_0(0) = 1$ (2.19) for symmetry. It should be remarked that $B(x)$ and $D(x)$ govern the difference equation for the polynomial $P_n(\eta)$, the solution of which requires the knowledge of the sinusoidal coordinate $\eta(x)$. The same quantities $B(x)$ and $D(x)$ specify the three term recurrence of the dual polynomial $Q_x(\mathcal{E})$ without the knowledge of the spectrum $\mathcal{E}(n)$. It is required for them to be the eigenvectors of the eigenvalue problem (2.1). Likewise, A_n and C_n in (2.38) specify the polynomial $P_n(\eta)$ without the knowledge of the sinusoidal coordinate. As for the dual polynomial $Q_x(\mathcal{E}(n))$, A_n and C_n provide the difference equation (in n), the solution of which needs the explicit form of $\mathcal{E}(n)$. Let us stress that it is the eigenvalue problem (2.1) with the specific Hamiltonian (2.4) that determines the polynomial $P_n(\eta)$ and its dual $Q_x(\mathcal{E})$, the spectrum $\mathcal{E}(n)$, the sinusoidal coordinate $\eta(x)$ and the orthogonality measures $\phi_0(x)^2$ and d_n^2 .

4 Shape Invariance and Closure Relation

Like the Hamiltonians for the (q -)Askey-scheme of hypergeometric orthogonal polynomials with continuous measure, all the Hamiltonians for the orthogonal polynomials of a discrete variable discussed in this paper are endowed with two types of symmetries; shape-invariance [12, 13] and closure relation [10]. The former leads to the determination of the entire spectrum (4.5) and the corresponding eigenvectors (4.6). In other words, it guarantees the exact solvability in the Schrödinger picture. The closure relation (4.31), on the other hand, is essential for the exact Heisenberg operator solution (4.34), which in turn gives the annihilation/creation operators (4.37). They help to determine the coefficients A_n and C_n (4.49), (4.50) in the three term recurrence relation of the polynomial $P_n(\eta)$ as well as the algebraic method (4.42) to determine the eigenvalues $\{\mathcal{E}(n)\}$. In §4.3 we show that the requirement of the closure relation determines the sinusoidal coordinates $\eta(x)$ and $B(x)$ and $D(x)$. In Appendix A, the possible forms of the Hamiltonians are determined (A.15)–(A.16).

4.1 Shape Invariance

Shape invariance dictates the parameter dependence of the Hamiltonian. Let us denote the set of parameters symbolically by $\boldsymbol{\lambda}$:

$$\mathcal{H}(x; \boldsymbol{\lambda}) = \mathcal{A}(x; \boldsymbol{\lambda})^\dagger \mathcal{A}(x; \boldsymbol{\lambda}), \quad \text{or} \quad \mathcal{H}(\boldsymbol{\lambda}) = \mathcal{A}(\boldsymbol{\lambda})^\dagger \mathcal{A}(\boldsymbol{\lambda}). \quad (4.1)$$

Shape invariance simply means

$$\mathcal{A}(\boldsymbol{\lambda})\mathcal{A}(\boldsymbol{\lambda})^\dagger = \kappa\mathcal{A}(\boldsymbol{\lambda} + \boldsymbol{\delta})^\dagger \mathcal{A}(\boldsymbol{\lambda} + \boldsymbol{\delta}) + \mathcal{E}(1; \boldsymbol{\lambda}), \quad (4.2)$$

in which κ is a positive constant and $\mathcal{E}(1; \boldsymbol{\lambda})$ is the eigenvalue of the first excited state $\mathcal{E}(1) > 0$ with the explicit parameter dependence. That is, the original Hamiltonian $\mathcal{H}(\boldsymbol{\lambda})$ and the *associated* Hamiltonian $\mathcal{A}(\boldsymbol{\lambda})\mathcal{A}(\boldsymbol{\lambda})^\dagger$ in Crum's [7] sense (or in the factorisation method [8], or in the so-called super-symmetric quantum mechanics [9]) have the same *shape* with a multiplicative factor κ .¹ The necessary and sufficient condition for (4.2) is

$$\sqrt{B(x+1; \boldsymbol{\lambda})D(x+1; \boldsymbol{\lambda})} = \kappa\sqrt{B(x; \boldsymbol{\lambda} + \boldsymbol{\delta})D(x+1; \boldsymbol{\lambda} + \boldsymbol{\delta})}, \quad (4.3)$$

$$B(x; \boldsymbol{\lambda}) + D(x+1; \boldsymbol{\lambda}) = \kappa(B(x; \boldsymbol{\lambda} + \boldsymbol{\delta}) + D(x; \boldsymbol{\lambda} + \boldsymbol{\delta})) + \mathcal{E}(1; \boldsymbol{\lambda}). \quad (4.4)$$

For the finite dimensional case (I), the new (associated) Hamiltonian $\mathcal{H}(\boldsymbol{\lambda} + \boldsymbol{\delta}) = \mathcal{A}(\boldsymbol{\lambda} + \boldsymbol{\delta})^\dagger \mathcal{A}(\boldsymbol{\lambda} + \boldsymbol{\delta})$ is of N dimensional, and (4.3) holds for $0 \leq x \leq N-1$.

It is rather straightforward to extract the necessary information, as done explicitly for the (q -)Askey-scheme of hypergeometric polynomials with continuous measure [13]. The spectrum is simply generated by $\mathcal{E}(1, \boldsymbol{\lambda})$:

$$\mathcal{E}(n; \boldsymbol{\lambda}) = \sum_{s=0}^{n-1} \kappa^s \mathcal{E}(1; \boldsymbol{\lambda} + s\boldsymbol{\delta}), \quad (4.5)$$

and the corresponding eigenvectors are generated from the known form of the ground state eigenvector ϕ_0 (2.20) together with the multiple action of the successive \mathcal{A}^\dagger operator:

$$\phi_n(x; \boldsymbol{\lambda}) \propto \mathcal{A}(\boldsymbol{\lambda})^\dagger \mathcal{A}(\boldsymbol{\lambda} + \boldsymbol{\delta})^\dagger \mathcal{A}(\boldsymbol{\lambda} + 2\boldsymbol{\delta})^\dagger \cdots \mathcal{A}(\boldsymbol{\lambda} + (n-1)\boldsymbol{\delta})^\dagger \phi_0(x; \boldsymbol{\lambda} + n\boldsymbol{\delta}), \quad (4.6)$$

which is related to a Rodrigues type formula. It should be stressed that shape invariance, (4.5) guarantees the monotonous increasing property (within the valid parameter range)

¹ The continuous q -Hermite polynomial [3, 14] has Hamiltonian (2.9) with $V(x) = 1/((1-z^2)(1-qz^2))$, $z = e^{ix}$, which has no shiftable parameter. The above shape invariance relation gives a simple realisation of the q -oscillator algebra $\mathcal{A}\mathcal{A}^\dagger = q^{-1}\mathcal{A}^\dagger\mathcal{A} + \mathcal{E}_1$, $\mathcal{E}_1 = q^{-1} - 1$, [19].

of $\mathcal{E}(n)$. For all the polynomials discussed in this paper, the spectrum is either linear or quadratic in n or q -quadratic in n :

$$(i) : \mathcal{E}(n) = n, \quad (4.7)$$

$$(ii) : \mathcal{E}(n) = \epsilon n(n + \alpha), \quad \epsilon = \begin{cases} 1 & \text{for } \alpha > -1, \\ -1 & \text{for } \alpha < -N, \end{cases} \quad (4.8)$$

$$(iii) : \mathcal{E}(n) = 1 - q^n, \quad (4.9)$$

$$(iv) : \mathcal{E}(n) = q^{-n} - 1, \quad (4.10)$$

$$(v) : \mathcal{E}(n) = \epsilon(q^{-n} - 1)(1 - \alpha q^n), \quad \epsilon = \begin{cases} 1 & \text{for } \alpha < q^{-1}, \\ -1 & \text{for } \alpha > q^{-N}. \end{cases} \quad (4.11)$$

As we will show shortly in the next subsection, the requirement of closure relation restricts the structure of the sinusoidal coordinate $\eta(x)$, which has the same form as above with n replaced by x , (4.72)–(4.76) reflecting the duality $x \leftrightarrow n$, $\eta(x) \leftrightarrow \mathcal{E}(n)$.

Forward and Backward Shift Operators The counterparts of the \mathcal{A} and \mathcal{A}^\dagger operators in the polynomial space, called the forward and backward shift operators, play an important role in the theory of orthogonal polynomials [14, 3]. Since the shift in x is closely related with the parameter shift, we introduce an auxiliary function $\varphi(x) = \varphi(x; \boldsymbol{\lambda}) > 0$, $\varphi(0) = 1$ to absorb the effects of the parameter shift in the ground state wave function:

$$\varphi(x; \boldsymbol{\lambda}) \stackrel{\text{def}}{=} \sqrt{\frac{B(0; \boldsymbol{\lambda})}{B(x; \boldsymbol{\lambda})} \frac{\phi_0(x; \boldsymbol{\lambda} + \boldsymbol{\delta})}{\phi_0(x; \boldsymbol{\lambda})}}. \quad (4.12)$$

In the finite dimensional case (I), $B(N; \boldsymbol{\lambda}) = 0$ but it is canceled by the zero in $\phi_0(N; \boldsymbol{\lambda} + \boldsymbol{\delta})$ and φ is well defined. With the help of (2.18), it can be rewritten as

$$\varphi(x; \boldsymbol{\lambda}) = \sqrt{\frac{B(0; \boldsymbol{\lambda})}{D(x+1; \boldsymbol{\lambda})} \frac{\phi_0(x; \boldsymbol{\lambda} + \boldsymbol{\delta})}{\phi_0(x+1; \boldsymbol{\lambda})}}. \quad (4.13)$$

Explicitly it reads

$$\varphi(x; \boldsymbol{\lambda}) = \sqrt{\frac{B(0; \boldsymbol{\lambda})}{B(x; \boldsymbol{\lambda})} \prod_{y=0}^{x-1} \frac{B(y; \boldsymbol{\lambda} + \boldsymbol{\delta})}{B(y; \boldsymbol{\lambda})} \frac{D(y+1; \boldsymbol{\lambda})}{D(y+1; \boldsymbol{\lambda} + \boldsymbol{\delta})}}. \quad (4.14)$$

For the Hamiltonians (2.8), (2.9) with a continuous variable x , the corresponding $\varphi(x)$ is given in [10].

Let us introduce the *forward shift operator* $\mathcal{F} = \mathcal{F}(\boldsymbol{\lambda})$, and the *backward shift operator* $\mathcal{B} = \mathcal{B}(\boldsymbol{\lambda})$ as

$$\tilde{\mathcal{H}}(\boldsymbol{\lambda}) = \mathcal{B}(\boldsymbol{\lambda})\mathcal{F}(\boldsymbol{\lambda}), \quad (4.15)$$

$$\mathcal{F}(\boldsymbol{\lambda}) \stackrel{\text{def}}{=} \sqrt{B(0; \boldsymbol{\lambda})} \phi_0(x; \boldsymbol{\lambda} + \boldsymbol{\delta})^{-1} \circ \mathcal{A}(\boldsymbol{\lambda}) \circ \phi_0(x; \boldsymbol{\lambda}), \quad (4.16)$$

$$\mathcal{B}(\boldsymbol{\lambda}) \stackrel{\text{def}}{=} \frac{1}{\sqrt{B(0; \boldsymbol{\lambda})}} \phi_0(x; \boldsymbol{\lambda})^{-1} \circ \mathcal{A}(\boldsymbol{\lambda})^\dagger \circ \phi_0(x; \boldsymbol{\lambda} + \boldsymbol{\delta}). \quad (4.17)$$

With the help of (4.12)–(4.13) we obtain expressions

$$\mathcal{F}(\boldsymbol{\lambda}) = B(0; \boldsymbol{\lambda})\varphi(x; \boldsymbol{\lambda})^{-1}(1 - e^\partial), \quad (4.18)$$

$$\mathcal{B}(\boldsymbol{\lambda}) = \frac{1}{B(0; \boldsymbol{\lambda})} (B(x; \boldsymbol{\lambda}) - D(x; \boldsymbol{\lambda})e^{-\partial})\varphi(x; \boldsymbol{\lambda}). \quad (4.19)$$

With $B(0; \boldsymbol{\lambda})$ appearing in many places, these formulas look rather contrived. However, their action on the polynomial $P_n(\eta) = P_n(\eta(x; \boldsymbol{\lambda}); \boldsymbol{\lambda})$ is rather simple:

$$\mathcal{F}(\boldsymbol{\lambda})P_n(\eta(x; \boldsymbol{\lambda}); \boldsymbol{\lambda}) = f_n(\boldsymbol{\lambda})P_{n-1}(\eta(x; \boldsymbol{\lambda} + \boldsymbol{\delta}); \boldsymbol{\lambda} + \boldsymbol{\delta}), \quad (4.20)$$

$$\mathcal{B}(\boldsymbol{\lambda})P_n(\eta(x; \boldsymbol{\lambda} + \boldsymbol{\delta}); \boldsymbol{\lambda} + \boldsymbol{\delta}) = b_n(\boldsymbol{\lambda})P_{n+1}(\eta(x; \boldsymbol{\lambda}); \boldsymbol{\lambda}), \quad (4.21)$$

in which $f_n(\boldsymbol{\lambda})$ and $b_n(\boldsymbol{\lambda})$ are constants satisfying $\mathcal{E}(n; \boldsymbol{\lambda}) = b_{n-1}(\boldsymbol{\lambda})f_n(\boldsymbol{\lambda})$. In fact, for all the examples given in §5, they take simple values,

$$f_n(\boldsymbol{\lambda}) = \mathcal{E}(n; \boldsymbol{\lambda}), \quad b_n(\boldsymbol{\lambda}) = 1. \quad (4.22)$$

Moreover, the auxiliary function $\varphi(x)$ is related to the sinusoidal coordinate $\eta(x)$ as

$$\varphi(x) = \frac{\eta(x+1) - \eta(x)}{\eta(1)}, \quad (4.23)$$

which gives

$$\frac{\eta(x)}{\eta(1)} = \sum_{y=0}^{x-1} \varphi(y). \quad (4.24)$$

The action of \mathcal{A} , \mathcal{A}^\dagger on the eigenfunction $\phi_n(x)$ is

$$\mathcal{A}(\boldsymbol{\lambda})\phi_n(x; \boldsymbol{\lambda}) = \frac{1}{\sqrt{B(0; \boldsymbol{\lambda})}} f_n(\boldsymbol{\lambda})\phi_{n-1}(x; \boldsymbol{\lambda} + \boldsymbol{\delta}), \quad (4.25)$$

$$\mathcal{A}(\boldsymbol{\lambda})^\dagger\phi_n(x; \boldsymbol{\lambda} + \boldsymbol{\delta}) = \sqrt{B(0; \boldsymbol{\lambda})} b_n(\boldsymbol{\lambda})\phi_{n+1}(x; \boldsymbol{\lambda}). \quad (4.26)$$

In terms of the forward and backward shift operators, the shape invariance condition (4.2) reads

$$\mathcal{F}(\boldsymbol{\lambda})\mathcal{B}(\boldsymbol{\lambda}) = \kappa\mathcal{B}(\boldsymbol{\lambda} + \boldsymbol{\delta})\mathcal{F}(\boldsymbol{\lambda} + \boldsymbol{\delta}) + \mathcal{E}(1; \boldsymbol{\lambda}) \quad (4.27)$$

The necessary and sufficient condition for (4.27) are

$$B(x+1; \boldsymbol{\lambda})\varphi(x+1; \boldsymbol{\lambda}) = \kappa B(x; \boldsymbol{\lambda} + \boldsymbol{\delta})\varphi(x; \boldsymbol{\lambda}), \quad (4.28)$$

$$D(x; \boldsymbol{\lambda})\varphi(x-1; \boldsymbol{\lambda}) = \kappa D(x; \boldsymbol{\lambda} + \boldsymbol{\delta})\varphi(x; \boldsymbol{\lambda}), \quad (4.29)$$

$$B(x; \boldsymbol{\lambda}) + D(x+1; \boldsymbol{\lambda}) = \kappa(B(x; \boldsymbol{\lambda} + \boldsymbol{\delta}) + D(x; \boldsymbol{\lambda} + \boldsymbol{\delta})) + \mathcal{E}(1; \boldsymbol{\lambda}), \quad (4.30)$$

which look simpler than (4.3) and (4.4) thanks to the presence of $\varphi(x)$.

4.2 Closure Relation

The other symmetry of the Hamiltonians for the orthogonal polynomials is called the *closure relation*. It is realised if the $\mathcal{H} = \mathcal{H}(\boldsymbol{\lambda})$ and the *sinusoidal coordinate* $\eta = \eta(x; \boldsymbol{\lambda})$ satisfy

$$[\mathcal{H}, [\mathcal{H}, \eta]] = \eta R_0(\mathcal{H}) + [\mathcal{H}, \eta] R_1(\mathcal{H}) + R_{-1}(\mathcal{H}), \quad (4.31)$$

in which $R_i(z) = R_i(z; \boldsymbol{\lambda})$ is a polynomial in z . The naming is due to the fact that $\eta(x) + R_{-1}(\mathcal{H})/R_0(\mathcal{H})$ undergoes a sinusoidal motion with frequency $\sqrt{R_0(\mathcal{H})}$ in the classical dynamics limit [10]. By similarity transformation in terms of $\phi_0(x; \boldsymbol{\lambda})$, (4.31) reads

$$[\tilde{\mathcal{H}}, [\tilde{\mathcal{H}}, \eta]] = \eta R_0(\tilde{\mathcal{H}}) + [\tilde{\mathcal{H}}, \eta] R_1(\tilde{\mathcal{H}}) + R_{-1}(\tilde{\mathcal{H}}). \quad (4.32)$$

The l.h.s. consists of the operators $e^{2\partial}$, e^∂ , 1 , $e^{-\partial}$, $e^{-2\partial}$ and \mathcal{H} contains $e^{\pm\partial}$. Thus R_i can be parametrised as

$$R_1(z) = r_1^{(1)}z + r_1^{(0)}, \quad R_0(z) = r_0^{(2)}z^2 + r_0^{(1)}z + r_0^{(0)}, \quad R_{-1}(z) = r_{-1}^{(2)}z^2 + r_{-1}^{(1)}z + r_{-1}^{(0)}. \quad (4.33)$$

The coefficients $r_i^{(j)}$ depend on the overall normalisation of the Hamiltonian and the sinusoidal coordinate: $r_1^{(j)} \propto B(0)^{1-j}$, $r_0^{(j)} \propto B(0)^{2-j}$, $r_{-1}^{(j)} \propto \eta(1)B(0)^{2-j}$.

Heisenberg operator solution and Annihilation and Creation operators The closure relation (4.31) enables us to express any multiple commutator

$$[\mathcal{H}, [\mathcal{H}, \dots, [\mathcal{H}, \eta(x)] \dots]]$$

as a linear combination of the operators $\eta(x)$ and $[\mathcal{H}, \eta(x)]$ with coefficients depending on the Hamiltonian \mathcal{H} only. The exact Heisenberg operator solution for the sinusoidal coordinate $\eta(x)$ [10] is given by (t is the time variable):

$$e^{it\mathcal{H}}\eta(x)e^{-it\mathcal{H}} = a^{(+)}e^{i\alpha_+(\mathcal{H})t} + a^{(-)}e^{i\alpha_-(\mathcal{H})t} - R_{-1}(\mathcal{H})R_0(\mathcal{H})^{-1}, \quad (4.34)$$

$$\alpha_{\pm}(z) \stackrel{\text{def}}{=} \frac{1}{2} \left(R_1(z) \pm \sqrt{R_1(z)^2 + 4R_0(z)} \right), \quad (4.35)$$

$$R_1(z) = \alpha_+(z) + \alpha_-(z), \quad R_0(z) = -\alpha_+(z)\alpha_-(z), \quad (4.36)$$

$$a^{(\pm)} \stackrel{\text{def}}{=} \pm \left([\mathcal{H}, \eta(x)] - (\eta(x) + R_{-1}(\mathcal{H})R_0(\mathcal{H})^{-1})\alpha_{\mp}(\mathcal{H}) \right) (\alpha_+(\mathcal{H}) - \alpha_-(\mathcal{H}))^{-1} \quad (4.37)$$

$$= \pm (\alpha_+(\mathcal{H}) - \alpha_-(\mathcal{H}))^{-1} \left([\mathcal{H}, \eta(x)] + \alpha_{\pm}(\mathcal{H})(\eta(x) + R_{-1}(\mathcal{H})R_0(\mathcal{H})^{-1}) \right). \quad (4.38)$$

The positive and negative frequency parts of the Heisenberg operator solution, $a^{(-)}$ and $a^{(+)}$, are the *annihilation* and *creation* operators, which are hermitian conjugate to each other $a^{(+)\dagger} = a^{(-)}$. Applying (4.34) to the eigenvector $\phi_n(x) = \phi_0(x)P_n(\eta(x))$ and using the three term recurrence relation for the polynomial $P_n(\eta)$ (2.38), we obtain [10]

$$a^{(+)}\phi_n(x) = A_n\phi_{n+1}(x), \quad (4.39)$$

$$a^{(-)}\phi_n(x) = C_n\phi_{n-1}(x). \quad (4.40)$$

The constant part in (4.34) is related to B_n :

$$R_{-1}(\mathcal{E}(n))R_0(\mathcal{E}(n))^{-1} = -B_n = A_n + C_n. \quad (4.41)$$

The exact energy eigenvalues $\{\mathcal{E}(n)\}$ can be obtained as the solution of the non-linear equation

$$\alpha_{\pm}(\mathcal{E}(n)) = \mathcal{E}(n \pm 1) - \mathcal{E}(n) \quad (4.42)$$

starting with $\mathcal{E}(0) = 0$ [10]. The above relation simplifies the form of the annihilation-creation operators applied to $\phi_n(x)$

$$\begin{aligned} & a^{(\pm)}\phi_n(x) \quad (4.43) \\ &= \frac{\pm 1}{\mathcal{E}(n+1) - \mathcal{E}(n-1)} \left([\mathcal{H}, \eta(x)] + (\mathcal{E}(n) - \mathcal{E}(n \mp 1))\eta(x) + \frac{R_{-1}(\mathcal{E}(n))}{\mathcal{E}(n \pm 1) - \mathcal{E}(n)} \right) \phi_n(x). \end{aligned}$$

This is the eigenvector version of the raising/lowering operators or the structure relation for orthogonal polynomials [20]. The polynomial version is obtained by the similarity transformation in terms of $\phi_0(x)$, $\tilde{a}^{(\pm)} \stackrel{\text{def}}{=} \phi_0(x)^{-1} \circ a^{(\pm)} \circ \phi_0(x)$,

$$\begin{aligned} & \tilde{a}^{(\pm)}P_n(\eta(x)) \quad (4.44) \\ &= \frac{\pm 1}{\mathcal{E}(n+1) - \mathcal{E}(n-1)} \left([\tilde{\mathcal{H}}, \eta(x)] + (\mathcal{E}(n) - \mathcal{E}(n \mp 1))\eta(x) + \frac{R_{-1}(\mathcal{E}(n))}{\mathcal{E}(n \pm 1) - \mathcal{E}(n)} \right) P_n(\eta(x)). \end{aligned}$$

Determination of A_n and C_n The coefficients A_n and C_n of the three term recurrence relation of the eigenpolynomial $P_n(\eta)$ can be determined elementarily from the above expression for B_n (4.41) and the duality. The starting point is the three term recurrence relations of the polynomial $P_n(\eta)$ and its dual $Q_x(\mathcal{E})$:

$$\eta(x)P_n(\eta(x)) = A_n(P_{n+1}(\eta(x)) - P_n(\eta(x))) + C_n(P_{n-1}(\eta(x)) - P_n(\eta(x))), \quad (4.45)$$

$$\mathcal{E}(n)Q_x(\mathcal{E}(n)) = B(x)(Q_x(\mathcal{E}(n)) - Q_{x+1}(\mathcal{E}(n))) + D(x)(Q_x(\mathcal{E}(n)) - Q_{x-1}(\mathcal{E}(n))). \quad (4.46)$$

By putting $x = 0$ in (4.46) and using $Q_0 = 1$, $D(0) = 0$, we obtain

$$P_n(\eta(1)) = Q_1(\mathcal{E}(n)) = \frac{\mathcal{E}(n) - B(0)}{-B(0)}, \quad (4.47)$$

in which the first equality is due to the duality (3.10). Next, by putting $x = 1$ in (4.45) and using (4.47), we obtain

$$\eta(1)(\mathcal{E}(n) - B(0)) = A_n(\mathcal{E}(n+1) - \mathcal{E}(n)) + C_n(\mathcal{E}(n-1) - \mathcal{E}(n)). \quad (4.48)$$

The two equations (4.41) and (4.48) for $n \geq 1$ give

$$A_n = \frac{\frac{R_{-1}(\mathcal{E}(n))}{R_0(\mathcal{E}(n))}(\mathcal{E}(n) - \mathcal{E}(n-1)) + \eta(1)(\mathcal{E}(n) - B(0))}{\mathcal{E}(n+1) - \mathcal{E}(n-1)}, \quad (4.49)$$

$$C_n = \frac{\frac{R_{-1}(\mathcal{E}(n))}{R_0(\mathcal{E}(n))}(\mathcal{E}(n) - \mathcal{E}(n+1)) + \eta(1)(\mathcal{E}(n) - B(0))}{\mathcal{E}(n-1) - \mathcal{E}(n+1)}, \quad (4.50)$$

and for $n = 0$ we obtain from (4.41)

$$A_0 = R_{-1}(0)R_0(0)^{-1} = r_{-1}^{(0)}/r_0^{(0)}, \quad C_0 = 0. \quad (4.51)$$

They are slightly simplified ($n \geq 1$) in terms of α_{\pm} :

$$A_n = \frac{R_{-1}(\mathcal{E}(n)) + \eta(1)(\mathcal{E}(n) - B(0))\alpha_+(\mathcal{E}(n))}{\alpha_+(\mathcal{E}(n))(\alpha_+(\mathcal{E}(n)) - \alpha_-(\mathcal{E}(n)))}, \quad (4.52)$$

$$C_n = \frac{R_{-1}(\mathcal{E}(n)) + \eta(1)(\mathcal{E}(n) - B(0))\alpha_-(\mathcal{E}(n))}{\alpha_-(\mathcal{E}(n))(\alpha_-(\mathcal{E}(n)) - \alpha_+(\mathcal{E}(n)))}. \quad (4.53)$$

Note that (4.48) for $n = 0$ gives another important relation

$$A_0\mathcal{E}(1) + B(0)\eta(1) = 0. \quad (4.54)$$

This relation ensures that the expression for C_n (4.53) vanishes for $n = 0$. Written differently, the above relation (4.54) means an equality

$$\frac{B(0)}{\mathcal{E}(1)} = \frac{-A_0}{\eta(1)} \quad (4.55)$$

between the two intensive quantities. That is, they are independent of the overall normalisation of the Hamiltonian \mathcal{H} of the polynomial system and of $\mathcal{H}^{\text{dual}}$ of its dual.

Dual Hamiltonian and Determination of d_n/d_0 By the duality (3.10), the three term recurrence relation for $P_n(\eta)$ (4.45) translates to the difference equation in n for the dual polynomial $Q_x(\mathcal{E}(n))$;

$$\eta(x)Q_x(\mathcal{E}(n)) = A_n(Q_x(\mathcal{E}(n+1)) - Q_x(\mathcal{E}(n))) + C_n(Q_x(\mathcal{E}(n-1)) - Q_x(\mathcal{E}(n))). \quad (4.56)$$

With abuse of notation, it can be written as an eigenvalue equation for the operator $\tilde{\mathcal{H}}^{\text{dual}}$ with eigenvalue $\eta(x)$:

$$\tilde{\mathcal{H}}^{\text{dual}} \stackrel{\text{def}}{=} -A_n(1 - e^{\partial_n}) - C_n(1 - e^{-\partial_n}), \quad (4.57)$$

$$\tilde{\mathcal{H}}^{\text{dual}} Q_x(\mathcal{E}(n)) = \eta(x)Q_x(\mathcal{E}(n)). \quad (4.58)$$

The corresponding Hamiltonian matrix is parametrised by $n, m = 0, 1, \dots, (N), \dots$,

$$\mathcal{H}^{\text{dual}} \stackrel{\text{def}}{=} -\sqrt{-A_n} e^{\partial_n} \sqrt{-C_n} - \sqrt{-C_n} e^{-\partial_n} \sqrt{-A_n} - (A_n + C_n), \quad A_n < 0, \quad C_n < 0, \quad (4.59)$$

$$(\mathcal{H}^{\text{dual}})_{n,m} = -\sqrt{A_n C_{n+1}} \delta_{n+1,m} - \sqrt{A_{n-1} C_n} \delta_{n-1,m} - (A_n + C_n) \delta_{n,m}. \quad (4.60)$$

Following the same path that led to $\phi_0(x)$ (2.20), the orthogonality measure for the polynomial $P_n(\eta(x))$, we obtain the formula for the orthogonality measure d_n for the dual polynomial $Q_x(\mathcal{E}(n))$:

$$\frac{d_n}{d_0} = \sqrt{\prod_{m=0}^{n-1} \frac{A_m}{C_{m+1}}}. \quad (4.61)$$

It should be stressed that this same d_n is the normalisation constant of the polynomial $P_n(\eta(x))$ (2.33). With the above formula, one only needs to evaluate d_0 explicitly.

4.3 Determination of $\eta(x)$ and $B(x) + D(x)$

In the previous subsection, consequences of the closure relation are explored. Here we will consider the closure relation as algebraic constraints for unspecified Hamiltonian (*i.e.* $B(x)$)

and $D(x)$) and as yet undetermined sinusoidal coordinate $\eta(x)$. Under very mild assumptions (on top of the general conditions (2.5)), $\eta(0) = 0$, $\eta(x) > \eta(y)$ for $x > y$, $x, y \in \mathbb{Z}_+$, we determine the possible forms of $\eta(x)$ and $B(x) + D(x)$. The characterisation of all the systems satisfying the closure relation (*i.e.* determination of $B(x)$ and $D(x)$ separately) will be relegated to Appendix A.

The necessary and sufficient condition for the closure relation in the polynomial space (4.32) is

$$\eta(x+2) - 2\eta(x+1) + \eta(x) = r_0^{(2)}\eta(x) + r_{-1}^{(2)} + r_1^{(1)}(\eta(x+1) - \eta(x)), \quad (4.62)$$

$$\eta(x-2) - 2\eta(x-1) + \eta(x) = r_0^{(2)}\eta(x) + r_{-1}^{(2)} + r_1^{(1)}(\eta(x-1) - \eta(x)), \quad (4.63)$$

$$\begin{aligned} & (\eta(x+1) - \eta(x))(B(x+1) + D(x+1) - B(x) - D(x)) \\ &= -(r_0^{(2)}\eta(x) + r_{-1}^{(2)})(B(x+1) + D(x+1) + B(x) + D(x)) - (r_0^{(1)}\eta(x) + r_{-1}^{(1)}) \\ & \quad - (\eta(x+1) - \eta(x))(r_1^{(1)}(B(x+1) + D(x+1)) + r_1^{(0)}), \end{aligned} \quad (4.64)$$

$$\begin{aligned} & (\eta(x-1) - \eta(x))(B(x-1) + D(x-1) - B(x) - D(x)) \\ &= -(r_0^{(2)}\eta(x) + r_{-1}^{(2)})(B(x-1) + D(x-1) + B(x) + D(x)) - (r_0^{(1)}\eta(x) + r_{-1}^{(1)}) \\ & \quad - (\eta(x-1) - \eta(x))(r_1^{(1)}(B(x-1) + D(x-1)) + r_1^{(0)}), \end{aligned} \quad (4.65)$$

$$\begin{aligned} & -2(\eta(x+1) - \eta(x))B(x)D(x+1) - 2(\eta(x-1) - \eta(x))D(x)B(x-1) \\ &= (r_0^{(2)}\eta(x) + r_{-1}^{(2)})\left((B(x) + D(x))^2 + B(x)D(x+1) + D(x)B(x-1)\right) \\ & \quad + (r_0^{(1)}\eta(x) + r_{-1}^{(1)})(B(x) + D(x)) + r_0^{(0)}\eta(x) + r_{-1}^{(0)} \\ & \quad + r_1^{(1)}\left((\eta(x+1) - \eta(x))B(x)D(x+1) + (\eta(x-1) - \eta(x))D(x)B(x-1)\right). \end{aligned} \quad (4.66)$$

The first two equations (4.62) and (4.63), two above and below the diagonal of (4.32), determine $\eta(x)$. Then the next two equations (4.64) and (4.65), one above and below the diagonal, give $B(x) + D(x)$ as a function of $\eta(x)$. In Appendix A, we show that the last equation (4.66), the diagonal part, determines $B(x)$ and $D(x)$ separately.

By subtracting (4.62) from (4.63) (with x replaced by $x+2$), we obtain

$$0 = (r_1^{(1)} - r_0^{(2)})(\eta(x+2) - \eta(x)), \quad (4.67)$$

which gives a general constraint on the parameters $r_0^{(2)} = r_1^{(1)}$. By subtracting (4.64) from (4.65) (with x replaced by $x+1$) and using the above constraint, we obtain

$$0 = (2r_1^{(0)} - r_0^{(1)})(\eta(x+1) - \eta(x)). \quad (4.68)$$

With the two constraints

$$r_0^{(2)} = r_1^{(1)}, \quad r_0^{(1)} = 2r_1^{(0)}, \quad (4.69)$$

the first equation (4.62) is reduced to a three term recurrence relation for $\eta(x)$:

$$\eta(x+2) - (2 + r_1^{(1)})\eta(x+1) + \eta(x) = r_{-1}^{(2)}, \quad (4.70)$$

which can be solved by the input $\eta(0) = 0$ and $\eta(1)$. The solutions are classified by the roots of the ‘characteristic equation’:

$$t^2 - (2 + r_1^{(1)})t + 1 = 0.$$

In order to have two positive roots (necessary for the condition $\eta(x) > \eta(y)$ for $x > y$), $r_1^{(1)} \geq 0$ is needed. For the case $r_1^{(1)} = 0$, we have a quadratic solution:

$$\eta(x) = \frac{1}{2}r_{-1}^{(2)}x\left(x - 1 + \frac{2\eta(1)}{r_{-1}^{(2)}}\right) \quad (\text{for } r_{-1}^{(2)} \neq 0); \quad \eta(x) = \eta(1)x \quad (\text{for } r_{-1}^{(2)} = 0). \quad (4.71)$$

For the rest $r_1^{(1)} > 0$, we have a q -quadratic solution. It should be noted that the overall multiplicative factor is irrelevant. After choosing $\eta(1)$ and other parameters properly, we arrive at

$$(i) : \quad \eta(x) = x, \quad (4.72)$$

$$(ii) : \quad \eta(x) = \epsilon'x(x+d), \quad \epsilon' = \begin{cases} 1 & \text{for } d > -1, \\ -1 & \text{for } d < -N, \end{cases} \quad (4.73)$$

$$(iii) : \quad \eta(x) = 1 - q^x, \quad (4.74)$$

$$(iv) : \quad \eta(x) = q^{-x} - 1, \quad (4.75)$$

$$(v) : \quad \eta(x) = \epsilon'(q^{-x} - 1)(1 - dq^x), \quad \epsilon' = \begin{cases} 1 & \text{for } d < q^{-1}, \\ -1 & \text{for } d > q^{-N}. \end{cases} \quad (4.76)$$

Here as usual q is $0 < q < 1$. Similar conclusion was reached in a different context by Atakishiyev-Rahman-Suslov [21]. This result can be considered as a difference equation version of Bochner’s theorem [22, 23, 24]. In fact, the original setting of Bochner’s theorem [22] is rather limited and its generalisation within the framework of ordinary quantum mechanics is given in [10], which contains various sinusoidal coordinates $\eta(x) = x, x^2, \cos x, \sinh x, e^{-x}$, etc.

In these cases the parameters in (4.69) read

$$(i) : \quad r_1^{(1)} = 0, \quad r_{-1}^{(2)} = 0, \quad (4.77)$$

$$(ii) : r_1^{(1)} = 0, \quad r_{-1}^{(2)} = 2\epsilon', \quad (4.78)$$

$$(iii) : r_1^{(1)} = (q^{-\frac{1}{2}} - q^{\frac{1}{2}})^2, \quad r_{-1}^{(2)} = -(q^{-\frac{1}{2}} - q^{\frac{1}{2}})^2, \quad (4.79)$$

$$(iv) : r_1^{(1)} = (q^{-\frac{1}{2}} - q^{\frac{1}{2}})^2, \quad r_{-1}^{(2)} = (q^{-\frac{1}{2}} - q^{\frac{1}{2}})^2, \quad (4.80)$$

$$(v) : r_1^{(1)} = (q^{-\frac{1}{2}} - q^{\frac{1}{2}})^2, \quad r_{-1}^{(2)} = (q^{-\frac{1}{2}} - q^{\frac{1}{2}})^2 \epsilon' (1 + d), \quad (4.81)$$

and $\eta(x \pm 1) - \eta(x)$ are given by

$$(i) : \eta(x \pm 1) - \eta(x) = \pm 1, \quad (4.82)$$

$$(ii) : \eta(x \pm 1) - \eta(x) = \pm \epsilon' (2x \pm 1 + d), \quad (4.83)$$

$$(iii) : \eta(x \pm 1) - \eta(x) = (1 - q^{\pm 1})q^x, \quad (4.84)$$

$$(iv) : \eta(x \pm 1) - \eta(x) = (q^{\mp 1} - 1)q^{-x}, \quad (4.85)$$

$$(v) : \eta(x \pm 1) - \eta(x) = \epsilon' (q^{\mp 1} - 1)q^{-x} (1 - dq^{2x \pm 1}). \quad (4.86)$$

Solving (4.64) and (4.65) is somewhat involved. With the constraints (4.69) incorporated, (4.64) reads,

$$\begin{aligned} & (\eta(x+1) - \eta(x))(a_{x+1} - a_x) \\ &= -(r_1^{(1)}\eta(x) + r_{-1}^{(2)})(a_{x+1} + a_x) - (2r_1^{(0)}\eta(x) + r_{-1}^{(1)}) - (\eta(x+1) - \eta(x))(r_1^{(1)}a_{x+1} + r_1^{(0)}), \end{aligned}$$

which can be rewritten as

$$\begin{aligned} & (\eta(x+1) - \eta(x) + r_1^{(1)}\eta(x+1) + r_{-1}^{(2)})a_{x+1} + (-\eta(x+1) + \eta(x) + r_1^{(1)}\eta(x) + r_{-1}^{(2)})a_x \\ &= -(\eta(x+1) + \eta(x))r_1^{(0)} - r_{-1}^{(1)}. \end{aligned} \quad (4.87)$$

Here we have introduced the abbreviation

$$a_x \stackrel{\text{def}}{=} B(x) + D(x). \quad (4.88)$$

With the help of (4.70), (4.87) can be simplified to

$$(\eta(x+2) - \eta(x+1))a_{x+1} + (\eta(x-1) - \eta(x))a_x = -(\eta(x+1) + \eta(x))r_1^{(0)} - r_{-1}^{(1)}. \quad (4.89)$$

By multiplying $\eta(x) - \eta(x+1)$ to (4.89), we obtain a two term recurrence relation for $a_x = B(x) + D(x)$:

$$(\eta(x+2) - \eta(x+1))(\eta(x) - \eta(x+1))a_{x+1} - (\eta(x+1) - \eta(x))(\eta(x-1) - \eta(x))a_x$$

$$= r_1^{(0)}(\eta(x+1)^2 - \eta(x)^2) + r_{-1}^{(1)}(\eta(x+1) - \eta(x)), \quad (4.90)$$

which can be solved with the boundary (initial) condition

$$a_0 = B(0). \quad (4.91)$$

The result is that the diagonal component of \mathcal{H} and $\tilde{\mathcal{H}}$, $a_x \stackrel{\text{def}}{=} B(x) + D(x)$, multiplied by $(\eta(x+1) - \eta(x))(\eta(x-1) - \eta(x))$ is a quadratic polynomial in $\eta(x)$:

$$(\eta(x+1) - \eta(x))(\eta(x-1) - \eta(x))a_x = r_1^{(0)}\eta(x)^2 + r_{-1}^{(1)}\eta(x) + \eta(1)\eta(-1)a_0, \quad (4.92)$$

in which we have introduced a constant $\eta(-1)$

$$\eta(-1) \stackrel{\text{def}}{=} r_{-1}^{(2)} - \eta(1), \quad (4.93)$$

such that (4.70) is formally satisfied. This is a very important characterisation of $B(x)$ and $D(x)$ as we will see in the concrete examples in section 5. It should be remarked that the factor $(\eta(x+1) - \eta(x))(\eta(x-1) - \eta(x))$ is also a quadratic polynomial in $\eta(x)$,

$$- (\eta(x+1) - \eta(x))(\eta(x-1) - \eta(x)) = r_1^{(1)}\eta(x)^2 + 2r_{-1}^{(2)}\eta(x) - \eta(1)\eta(-1), \quad (4.94)$$

for all the cases listed in (4.72)–(4.76), as can be easily verified from (4.82)–(4.86). Therefore we obtain

$$a_x = - \frac{r_1^{(0)}\eta(x)^2 + r_{-1}^{(1)}\eta(x) + \eta(1)\eta(-1)B(0)}{r_1^{(1)}\eta(x)^2 + 2r_{-1}^{(2)}\eta(x) - \eta(1)\eta(-1)}. \quad (4.95)$$

The analysis of the remaining equation (4.66), which is the diagonal part of the closure relation (4.32) is relegated to Appendix A.

4.4 Dual Closure Relation

Since the sinusoidal coordinate plays the central role in the theory of orthogonal polynomials, understanding its general properties is important. Let us now consider the ‘dual’ closure relation by interchanging $\eta \leftrightarrow \mathcal{H}$ ($\Leftrightarrow \mathcal{E}$) in the closure relation (4.31):

$$[\eta, [\eta, \mathcal{H}]] = \mathcal{H} R_0^{\text{dual}}(\eta) + [\eta, \mathcal{H}] R_1^{\text{dual}}(\eta) + R_{-1}^{\text{dual}}(\eta), \quad (4.96)$$

in which $R_i^{\text{dual}}(\eta)$ are polynomials in η . By similarity transformation in terms of $\phi_0(x)$, we obtain

$$[\eta, [\eta, \tilde{\mathcal{H}}]] = \tilde{\mathcal{H}} R_0^{\text{dual}}(\eta) + [\eta, \tilde{\mathcal{H}}] R_1^{\text{dual}}(\eta) + R_{-1}^{\text{dual}}(\eta), \quad (4.97)$$

which is again a relationship among tridiagonal and diagonal matrices. It is equivalent to the following set of three equations:

$$(\eta(x+1) - \eta(x))^2 = R_0^{\text{dual}}(\eta(x+1)) - (\eta(x+1) - \eta(x))R_1^{\text{dual}}(\eta(x+1)), \quad (4.98)$$

$$(\eta(x-1) - \eta(x))^2 = R_0^{\text{dual}}(\eta(x-1)) - (\eta(x-1) - \eta(x))R_1^{\text{dual}}(\eta(x-1)), \quad (4.99)$$

$$0 = (B(x) + D(x))R_0^{\text{dual}}(\eta(x)) + R_{-1}^{\text{dual}}(\eta(x)). \quad (4.100)$$

These imply

$$R_1^{\text{dual}}(\eta(x)) = (\eta(x+1) - \eta(x)) + (\eta(x-1) - \eta(x)), \quad (4.101)$$

$$R_0^{\text{dual}}(\eta(x)) = -(\eta(x+1) - \eta(x))(\eta(x-1) - \eta(x)), \quad (4.102)$$

$$R_{-1}^{\text{dual}}(\eta(x)) = -(B(x) + D(x))R_0^{\text{dual}}(\eta(x)). \quad (4.103)$$

These expressions are the dual formulas of $\mathcal{E}(n)$, (4.36), (4.41) and (4.42). The right hand sides of (4.101)–(4.103) are in fact polynomials in $\eta(x)$:

$$R_1^{\text{dual}}(z) = r_1^{(1)}z + r_{-1}^{(2)}, \quad (4.104)$$

$$R_0^{\text{dual}}(z) = r_1^{(1)}z^2 + 2r_{-1}^{(2)}z - \eta(1)\eta(-1), \quad (4.105)$$

$$R_{-1}^{\text{dual}}(z) = r_1^{(0)}z^2 + r_{-1}^{(1)}z + \eta(1)\eta(-1)B(0). \quad (4.106)$$

The first two relations are consequences of (4.70) (see also (4.94)), and the last relation is nothing but the main result (4.92) of subsection 4.3. The coefficients of $R_i^{\text{dual}}(z)$ defined like as (4.33) satisfy the same constraints as (4.69). The number of independent parameters in $R_i^{\text{dual}}(z)$ is six, just the same as in $R_i(z)$ with (4.69). Among them, as mentioned after (4.33), the overall normalisation of \mathcal{H} and $\eta(x)$ are immaterial. This reduces the number of essential parameters to four, which corresponds to the most generic cases of the Racah §5.1.1 and the q -Racah polynomials §5.1.5.

5 Concrete Examples

Here we will discuss the properties of various orthogonal polynomials in our scheme. We stick to the standard notation as far as possible; the inevitable deviation is caused by symmetry consideration and our universal normalisation of polynomials (2.31), (3.4). When we use non-standard parametrisation (Racah, q -Racah, Hahn, q -Hahn and their duals) we adopt different symbols so that no confusion will ensue. We are well aware of the shortcomings

of deviating from the standard notation; for example, the connection to other existing concepts/quantities are blurred. However, we strongly believe that presenting the entire theory of orthogonal polynomials of a discrete variable in the hermitian matrix (or operator) form has enough merits of elucidating their common structure to compensate the shortcomings. The orthogonal polynomials whose orthogonality relation is written in terms of a q -integral (Jackson integral), *e.g.* the Big q -Jacobi polynomial, etc. will need different treatment and will be discussed elsewhere. We start with the finite dimensional case (I) and then move to the infinite dimensional case (II). In each case the Askey-scheme of hypergeometric orthogonal polynomials will be followed by the q -scheme polynomials, in the order presented in the review of Koekoek and Swarttouw [14], with the number *e.g.* [KS3.2] attached indicating the subsection there. The subsection is named after the polynomial to be discussed in the subsection. We explain various quantities and concepts in some detail for the first examples, the (q -) Racah polynomials. For the rest, mostly formulas only are given except for those needing some attention and/or remarks. As mentioned earlier, a good part of the explicit formulas for specific polynomials is already known in various contexts.

The only input is the Hamiltonian \mathcal{H} or the two positive functions $B(x)$ and $D(x)$. In order to define them, the parameters must be specified. So we start from the parameters and their shift properties under shape invariance. The other quantities are defined and derived as explained in previous sections. Most basic definitions and notation are recapitulated in Appendix B for self-containedness.

5.1 Finite Dimensional Case (I)

In this subsection x takes a finite range of values

$$x = 0, 1, 2, \dots, N.$$

5.1.1 Racah [KS1.2] (self-dual with different parameters)

As well-known, the Racah polynomial is the most general hypergeometric orthogonal polynomial of a discrete variable. All the other (non- q) polynomials are obtained by restriction or limiting procedure. The set of parameters $\boldsymbol{\lambda}$ and their shifting unit $\boldsymbol{\delta}$ and the multiplication factor κ (4.2) are:

$$\boldsymbol{\lambda} = (a, b, c, d), \quad \boldsymbol{\delta} = (1, 1, 1, 1), \quad \kappa = 1, \quad (5.1)$$

which are different from the standard ones $(\alpha, \beta, \gamma, \delta)$. Let us define \tilde{d} , ϵ and ϵ'

$$\tilde{d} \stackrel{\text{def}}{=} a + b + c - d - 1, \quad (5.2)$$

$$\epsilon \stackrel{\text{def}}{=} \begin{cases} 1 & \text{for } \tilde{d} > -1, \\ -1 & \text{for } \tilde{d} < -N, \end{cases} \quad \epsilon' \stackrel{\text{def}}{=} \begin{cases} 1 & \text{for } d > -1, \\ -1 & \text{for } d < -N, \end{cases} \quad (5.3)$$

and $B(x)$ and $D(x)$ as

$$B(x; \boldsymbol{\lambda}) = -\epsilon \frac{(x+a)(x+b)(x+c)(x+d)}{(2x+d)(2x+1+d)}, \quad (5.4)$$

$$D(x; \boldsymbol{\lambda}) = -\epsilon \frac{(x+d-a)(x+d-b)(x+d-c)x}{(2x-1+d)(2x+d)}. \quad (5.5)$$

The finiteness condition $B(N) = 0$ can be achieved by

$$a = -N \quad \text{or} \quad b = -N \quad \text{or} \quad c = -N. \quad (5.6)$$

Since $B(x)$ and $D(x)$ are symmetric in a, b and c , we assume without loss of generality

$$c = -N, \quad a \geq b. \quad (5.7)$$

Then we restrict our argument to the following parameter ranges in which $B(x)$ and $D(x)$ are positive:

$$d > 0, \quad a > N + d, \quad 0 < b < 1 + d \quad \Rightarrow \quad (\epsilon, \epsilon') = (1, 1), \quad (5.8)$$

$$d < -2N, \quad a > 0, \quad N + d < b < 1 - N \quad \Rightarrow \quad (\epsilon, \epsilon') = (1, -1), \quad (5.9)$$

$$d > 0, \quad 0 < a < 1 + d, \quad b < 1 - N \quad \Rightarrow \quad (\epsilon, \epsilon') = (-1, 1), \quad (5.10)$$

$$d < -2N, \quad N + d < a < 1 - N, \quad b < 1 + d \quad \Rightarrow \quad (\epsilon, \epsilon') = (-1, -1). \quad (5.11)$$

The energy eigenvalue and the sinusoidal coordinate are

$$\mathcal{E}(n; \boldsymbol{\lambda}) = \epsilon n(n + \tilde{d}), \quad \eta(x; \boldsymbol{\lambda}) = \epsilon' x(x + d). \quad (5.12)$$

The polynomial is

$$\begin{aligned} P_n(\eta(x; \boldsymbol{\lambda}); \boldsymbol{\lambda}) &= {}_4F_3 \left(\begin{matrix} -n, n + \tilde{d}, -x, x + d \\ a, b, c \end{matrix} \middle| 1 \right) \\ &= R_n(\epsilon' \eta(x; \boldsymbol{\lambda}); a - 1, \tilde{d} - a, c - 1, d - c), \end{aligned} \quad (5.13)$$

in which ${}_4F_3$ is the standard hypergeometric series (B.3) and R_n is the standard notation for the Racah polynomial [14]. The orthogonality measure (or the squared ground state

wavefunction) $\phi_0(x)^2$ and the normalisation constants d_n^2 (2.33) are obtained from (2.20) and (4.61) (with A_n and C_n given below),

$$\phi_0(x; \boldsymbol{\lambda})^2 = \frac{(a, b, c, d)_x}{(1 + d - a, 1 + d - b, 1 + d - c, 1)_x} \frac{2x + d}{d}, \quad (5.14)$$

$$\begin{aligned} d_n(\boldsymbol{\lambda})^2 &= \frac{(a, b, c, \tilde{d})_n}{(1 + \tilde{d} - a, 1 + \tilde{d} - b, 1 + \tilde{d} - c, 1)_n} \frac{2n + \tilde{d}}{\tilde{d}} \\ &\times \frac{(-1)^N (1 + d - a, 1 + d - b, 1 + d - c)_N}{(\tilde{d} + 1)_N (d + 1)_{2N}}, \end{aligned} \quad (5.15)$$

in which $(a)_n$ is the Pochhammer symbol (B.1). The format for $d_n(\boldsymbol{\lambda})^2$, throughout this section, consists of two parts separated by a \times symbol: $d_n^2 = (d_n^2/d_0^2) \times d_0^2$ reflecting the duality (3.18). The second part d_0^2 satisfies the relation $\sum_x \phi_0(x)^2 = 1/d_0^2$. The coefficients of the three term recurrence (2.38) A_n and C_n for the polynomial $P_n(\eta)$ obtained from the closure relation (4.49)–(4.50) are:

$$A_n(\boldsymbol{\lambda}) = \epsilon' \frac{(n + a)(n + b)(n + c)(n + \tilde{d})}{(2n + \tilde{d})(2n + 1 + \tilde{d})}, \quad (5.16)$$

$$C_n(\boldsymbol{\lambda}) = \epsilon' \frac{(n + \tilde{d} - a)(n + \tilde{d} - b)(n + \tilde{d} - c)n}{(2n - 1 + \tilde{d})(2n + \tilde{d})}. \quad (5.17)$$

The three functions R_1 , R_0 and R_{-1} appearing in the closure relation (4.31) are:

$$R_1(z; \boldsymbol{\lambda}) = 2\epsilon, \quad (5.18)$$

$$R_0(z; \boldsymbol{\lambda}) = 4\epsilon z + \tilde{d}^2 - 1, \quad (5.19)$$

$$R_{-1}(z; \boldsymbol{\lambda}) = 2\epsilon' z^2 + \epsilon\epsilon' (2(ab + bc + ca) - (1 + d)(1 + \tilde{d}))z + \epsilon' abc(\tilde{d} - 1). \quad (5.20)$$

Under the parameter shift $\boldsymbol{\lambda} \rightarrow \boldsymbol{\lambda} + \boldsymbol{\delta}$ the functions $B(x)$ and $D(x)$ behave:

$$B(x; \boldsymbol{\lambda} + \boldsymbol{\delta}) = B(x + 1; \boldsymbol{\lambda}) \frac{2x + 3 + d}{2x + 1 + d}, \quad D(x; \boldsymbol{\lambda} + \boldsymbol{\delta}) = D(x; \boldsymbol{\lambda}) \frac{2x - 1 + d}{2x + 1 + d}. \quad (5.21)$$

From this and (4.14), $\varphi(x)$ is

$$\varphi(x; \boldsymbol{\lambda}) = \frac{2x + d + 1}{d + 1}. \quad (5.22)$$

The constants $f_n(\boldsymbol{\lambda})$ and $b_n(\boldsymbol{\lambda})$ appearing in (4.20)–(4.21) are

$$f_n(\boldsymbol{\lambda}) = \mathcal{E}(n; \boldsymbol{\lambda}), \quad b_n(\boldsymbol{\lambda}) = 1. \quad (5.23)$$

By comparing $B(x)$, $D(x)$ with A_n and C_n , $\phi_0(x)^2$ with d_n^2 , it is clear that the theory is self-dual with the parameter correspondence $(a, b, c, d, \epsilon, \epsilon') \leftrightarrow (a, b, c, \tilde{d}, \epsilon', \epsilon)$.

5.1.2 Hahn [KS1.5] (dual to dual Hahn §5.1.3)

The parameters (a, b) are very slightly different from the standard ones (α, β) ($\alpha + 1 = a, \beta + 1 = b$) for the Hahn polynomial. For obvious reasons, we adopt the same parameters for the Hahn and dual Hahn polynomials:

$$\boldsymbol{\lambda} = (a, b, N), \quad \boldsymbol{\delta} = (1, 1, -1), \quad \kappa = 1, \quad (5.24)$$

$$\epsilon = \begin{cases} 1 & \text{for } a, b > 0, \\ -1 & \text{for } a, b < 1 - N, \end{cases} \quad (5.25)$$

$$B(x; \boldsymbol{\lambda}) = \epsilon(x + a)(N - x), \quad D(x; \boldsymbol{\lambda}) = \epsilon x(b + N - x), \quad (5.26)$$

$$\mathcal{E}(n; \boldsymbol{\lambda}) = \epsilon n(n + a + b - 1), \quad \eta(x; \boldsymbol{\lambda}) = x, \quad (5.27)$$

$$P_n(\eta(x; \boldsymbol{\lambda}); \boldsymbol{\lambda}) = {}_3F_2 \left(\begin{matrix} -n, n + a + b - 1, -x \\ a, -N \end{matrix} \middle| 1 \right) = Q_n(\eta(x; \boldsymbol{\lambda}); a - 1, b - 1, N), \quad (5.28)$$

$$\phi_0(x; \boldsymbol{\lambda})^2 = \frac{N!}{x!(N - x)!} \frac{(a)_x (b)_{N-x}}{(b)_N}, \quad (5.29)$$

$$d_n(\boldsymbol{\lambda})^2 = \frac{N!}{n!(N - n)!} \frac{(a)_n (2n + a + b - 1)(a + b)_N}{(b)_n (n + a + b - 1)_{N+1}} \times \frac{(b)_N}{(a + b)_N}, \quad (5.30)$$

$$A_n(\boldsymbol{\lambda}) = -\frac{(n + a)(n + a + b - 1)(N - n)}{(2n - 1 + a + b)(2n + a + b)}, \quad (5.31)$$

$$C_n(\boldsymbol{\lambda}) = -\frac{n(n + b - 1)(n + a + b + N - 1)}{(2n - 2 + a + b)(2n - 1 + a + b)}, \quad (5.32)$$

$$R_1(z; \boldsymbol{\lambda}) = 2\epsilon, \quad (5.33)$$

$$R_0(z; \boldsymbol{\lambda}) = 4\epsilon z + (a + b - 2)(a + b), \quad (5.34)$$

$$R_{-1}(z; \boldsymbol{\lambda}) = -\epsilon(2N - a + b)z - a(a + b - 2)N, \quad (5.35)$$

$$B(x; \boldsymbol{\lambda} + \boldsymbol{\delta}) = B(x + 1; \boldsymbol{\lambda}), \quad D(x; \boldsymbol{\lambda} + \boldsymbol{\delta}) = D(x; \boldsymbol{\lambda}), \quad (5.36)$$

$$\varphi(x; \boldsymbol{\lambda}) = 1, \quad f_n(\boldsymbol{\lambda}) = \mathcal{E}(n; \boldsymbol{\lambda}), \quad b_n(\boldsymbol{\lambda}) = 1. \quad (5.37)$$

Obviously the Hahn and dual Hahn polynomials are dual to each other. The standard parameters for the latter are (γ, δ) .

5.1.3 dual Hahn [KS1.6] (dual to §5.1.2)

The parameters (a, b) are very slightly different from the standard ones (γ, δ) ($\gamma + 1 = a, \delta + 1 = b$) for the dual Hahn polynomial:

$$\boldsymbol{\lambda} = (a, b, N), \quad \boldsymbol{\delta} = (1, 0, -1), \quad \kappa = 1, \quad (5.38)$$

$$\epsilon = \begin{cases} 1 & \text{for } a, b > 0, \\ -1 & \text{for } a, b < 1 - N, \end{cases} \quad (5.39)$$

$$B(x; \boldsymbol{\lambda}) = \frac{(x+a)(x+a+b-1)(N-x)}{(2x-1+a+b)(2x+a+b)}, \quad (5.40)$$

$$D(x; \boldsymbol{\lambda}) = \frac{x(x+b-1)(x+a+b+N-1)}{(2x-2+a+b)(2x-1+a+b)}, \quad (5.41)$$

$$\mathcal{E}(n; \boldsymbol{\lambda}) = n, \quad \eta(x; \boldsymbol{\lambda}) = \epsilon x(x+a+b-1), \quad (5.42)$$

$$P_n(\eta(x; \boldsymbol{\lambda}); \boldsymbol{\lambda}) = {}_3F_2 \left(\begin{matrix} -n, x+a+b-1, -x \\ a, -N \end{matrix} \middle| 1 \right) = R_n(\epsilon \eta(x; \boldsymbol{\lambda}); a-1, b-1, N), \quad (5.43)$$

$$\phi_0(x; \boldsymbol{\lambda})^2 = \frac{N!}{x!(N-x)!} \frac{(a)_x (2x+a+b-1)(a+b)_N}{(b)_x (x+a+b-1)_{N+1}}, \quad (5.44)$$

$$d_n(\boldsymbol{\lambda})^2 = \frac{N!}{n!(N-n)!} \frac{(a)_n (b)_{N-n}}{(b)_N} \times \frac{(b)_N}{(a+b)_N}, \quad (5.45)$$

$$A_n(\boldsymbol{\lambda}) = -\epsilon(n+a)(N-n), \quad C_n(\boldsymbol{\lambda}) = -\epsilon n(b+N-n), \quad (5.46)$$

$$R_1(z; \boldsymbol{\lambda}) = 0, \quad R_0(z; \boldsymbol{\lambda}) = 1, \quad R_{-1}(z; \boldsymbol{\lambda}) = 2\epsilon z^2 - \epsilon(2N-a+b)z - \epsilon aN, \quad (5.47)$$

$$B(x; \boldsymbol{\lambda} + \boldsymbol{\delta}) = B(x+1; \boldsymbol{\lambda}) \frac{2x+2+a+b}{2x+a+b}, \quad D(x; \boldsymbol{\lambda} + \boldsymbol{\delta}) = D(x; \boldsymbol{\lambda}) \frac{2x-2+a+b}{2x+a+b}, \quad (5.48)$$

$$\varphi(x; \boldsymbol{\lambda}) = \frac{2x+a+b}{a+b}, \quad f_n(\boldsymbol{\lambda}) = \mathcal{E}(n; \boldsymbol{\lambda}), \quad b_n(\boldsymbol{\lambda}) = 1. \quad (5.49)$$

With the present parametrisation, the following duality (5.28) \leftrightarrow (5.43), (5.27) \leftrightarrow (5.42), (5.26) \leftrightarrow (5.46), (5.31)–(5.32) \leftrightarrow (5.40)–(5.41), (5.29) \leftrightarrow (5.45), (5.30) \leftrightarrow (5.44), is obvious.

5.1.4 Krawtchouk [KS1.10] (self-dual)

$$\boldsymbol{\lambda} = (p, N), \quad \boldsymbol{\delta} = (0, -1), \quad \kappa = 1, \quad (5.50)$$

$$0 < p < 1, \quad (5.51)$$

$$B(x; \boldsymbol{\lambda}) = p(N-x), \quad D(x; \boldsymbol{\lambda}) = (1-p)x, \quad (5.52)$$

$$\mathcal{E}(n; \boldsymbol{\lambda}) = n, \quad \eta(x; \boldsymbol{\lambda}) = x, \quad (5.53)$$

$$P_n(\eta(x; \boldsymbol{\lambda}); \boldsymbol{\lambda}) = {}_2F_1 \left(\begin{matrix} -n, -x \\ -N \end{matrix} \middle| p^{-1} \right) = K_n(\eta(x; \boldsymbol{\lambda}); p, N), \quad (5.54)$$

$$\phi_0(x; \boldsymbol{\lambda}) = \frac{N!}{x!(N-x)!} \left(\frac{p}{1-p} \right)^x, \quad (5.55)$$

$$d_n(\boldsymbol{\lambda})^2 = \frac{N!}{n!(N-n)!} \left(\frac{p}{1-p} \right)^n \times (1-p)^N, \quad (5.56)$$

$$A_n(\boldsymbol{\lambda}) = -p(N-n), \quad C_n(\boldsymbol{\lambda}) = -(1-p)n, \quad (5.57)$$

$$R_1(z; \boldsymbol{\lambda}) = 0, \quad R_0(z; \boldsymbol{\lambda}) = 1, \quad R_{-1}(z; \boldsymbol{\lambda}) = (2p-1)z - pN, \quad (5.58)$$

$$B(x; \boldsymbol{\lambda} + \boldsymbol{\delta}) = B(x+1; \boldsymbol{\lambda}), \quad D(x; \boldsymbol{\lambda} + \boldsymbol{\delta}) = D(x; \boldsymbol{\lambda}), \quad (5.59)$$

$$\varphi(x; \boldsymbol{\lambda}) = 1, \quad f_n(\boldsymbol{\lambda}) = \mathcal{E}(n; \boldsymbol{\lambda}), \quad b_n(\boldsymbol{\lambda}) = 1. \quad (5.60)$$

5.1.5 q -Racah [KS3.2] (self-dual with different parameters)

This is the first example of the q -scheme of the orthogonal polynomials. Among them the q -Racah polynomial is the most general. The set of parameters $\boldsymbol{\lambda}$ is different from the standard one $(\alpha, \beta, \gamma, \delta)$ in the same manner as for the Racah polynomial. The shifting unit $\boldsymbol{\delta}$ and the multiplication factor κ (4.2) are:

$$q^\lambda = (a, b, c, d), \quad \boldsymbol{\delta} = (1, 1, 1, 1), \quad \kappa = q^{-1}. \quad (5.61)$$

Here q^λ stands for $q^{(\lambda_1, \lambda_2, \dots)} = (q^{\lambda_1}, q^{\lambda_2}, \dots)$. They are shifted multiplicatively. By definition the parameter q is not shifted. Let us introduce $\tilde{d}, \epsilon, \epsilon'$

$$\tilde{d} = abcd^{-1}q^{-1}, \quad (5.62)$$

$$\epsilon = \begin{cases} 1 & \text{for } \tilde{d} < q^{-1}, \\ -1 & \text{for } \tilde{d} > q^{-1}, \end{cases} \quad \epsilon' = \begin{cases} 1 & \text{for } d < q^{-1}, \\ -1 & \text{for } d > q^{-1}. \end{cases} \quad (5.63)$$

The functions $B(x)$ and $D(x)$ are

$$B(x; \boldsymbol{\lambda}) = -\epsilon \frac{(1 - aq^x)(1 - bq^x)(1 - cq^x)(1 - dq^x)}{(1 - dq^{2x})(1 - dq^{2x+1})}, \quad (5.64)$$

$$D(x; \boldsymbol{\lambda}) = -\epsilon \tilde{d} \frac{(1 - a^{-1}dq^x)(1 - b^{-1}dq^x)(1 - c^{-1}dq^x)(1 - q^x)}{(1 - dq^{2x-1})(1 - dq^{2x})}. \quad (5.65)$$

The finiteness condition $B(N) = 0$ can be realised by one of the following choices:

$$a = q^{-N} \quad \text{or} \quad b = q^{-N} \quad \text{or} \quad c = q^{-N}. \quad (5.66)$$

Thanks to the symmetry in a, b and c , we choose for simplicity

$$c = q^{-N}, \quad 0 < a \leq b, \quad 0 < d, \quad (5.67)$$

and consider the following parameter ranges in which $B(x)$ and $D(x)$ are positive:

$$0 < d < 1, \quad 0 < a < q^N d, \quad qd < b < 1 \quad \Rightarrow \quad (\epsilon, \epsilon') = (1, 1), \quad (5.68)$$

$$d > q^{-2N}, \quad 0 < a < 1, \quad q^{1-N} < b < q^N d \quad \Rightarrow \quad (\epsilon, \epsilon') = (1, -1), \quad (5.69)$$

$$0 < d < 1, \quad qd < a < 1, \quad b > q^{1-N} \quad \Rightarrow \quad (\epsilon, \epsilon') = (-1, 1), \quad (5.70)$$

$$d > q^{-2N}, \quad q^{1-N} < a < q^N d, \quad b > qd \quad \Rightarrow \quad (\epsilon, \epsilon') = (-1, -1). \quad (5.71)$$

In case some of the parameters a, b and d are non-positive, the situation is slightly complicated but can be treated in a similar way. The energy eigenvalue, the sinusoidal coordinate and the q -polynomial are:

$$\mathcal{E}(n; \boldsymbol{\lambda}) = \epsilon(q^{-n} - 1)(1 - \tilde{d}q^n), \quad \eta(x; \boldsymbol{\lambda}) = \epsilon'(q^{-x} - 1)(1 - dq^x), \quad (5.72)$$

$$\begin{aligned}
P_n(\eta(x; \boldsymbol{\lambda}); \boldsymbol{\lambda}) &= {}_4\phi_3\left(\begin{matrix} q^{-n}, \tilde{d}q^n, q^{-x}, dq^x \\ a, b, c \end{matrix} \middle| q; q\right) \\
&= R_n(\epsilon'(1+d+\eta(x; \boldsymbol{\lambda})); aq^{-1}, \tilde{d}a^{-1}, cq^{-1}, dc^{-1}), \tag{5.73}
\end{aligned}$$

in which ${}_4\phi_3$ is the basic hypergeometric series (B.4). The orthogonality measure $\phi_0(x)^2$ and the normalisation constants d_n^2 (2.33) are obtained from (2.20) and (4.61) (with A_n and C_n given below),

$$\phi_0(x; \boldsymbol{\lambda})^2 = \frac{(a, b, c, d; q)_x}{(a^{-1}dq, b^{-1}dq, c^{-1}dq, q; q)_x \tilde{d}^x} \frac{1 - dq^{2x}}{1 - d}, \tag{5.74}$$

$$\begin{aligned}
d_n(\boldsymbol{\lambda})^2 &= \frac{(a, b, c, \tilde{d}; q)_n}{(a^{-1}\tilde{d}q, b^{-1}\tilde{d}q, c^{-1}\tilde{d}q, q; q)_n d^n} \frac{1 - \tilde{d}q^{2n}}{1 - \tilde{d}} \\
&\times \frac{(-1)^N (a^{-1}dq, b^{-1}dq, c^{-1}dq; q)_N \tilde{d}^N q^{\frac{1}{2}N(N+1)}}{(\tilde{d}q; q)_N (dq; q)_{2N}}, \tag{5.75}
\end{aligned}$$

in which $(a; q)_n$ is the q -Pochhammer symbol (B.2). The coefficients of the three term recurrence (2.38) A_n and C_n for the polynomial $P_n(\eta)$ obtained from the closure relation (4.49)–(4.50) are:

$$A_n(\boldsymbol{\lambda}) = \epsilon' \frac{(1 - aq^n)(1 - bq^n)(1 - cq^n)(1 - \tilde{d}q^n)}{(1 - \tilde{d}q^{2n})(1 - \tilde{d}q^{2n+1})}, \tag{5.76}$$

$$C_n(\boldsymbol{\lambda}) = \epsilon' d \frac{(1 - a^{-1}\tilde{d}q^n)(1 - b^{-1}\tilde{d}q^n)(1 - c^{-1}\tilde{d}q^n)(1 - q^n)}{(1 - \tilde{d}q^{2n-1})(1 - \tilde{d}q^{2n})}. \tag{5.77}$$

The three functions R_1 , R_0 and R_{-1} appearing in the closure relation (4.31) are:

$$R_1(z; \boldsymbol{\lambda}) = (q^{-\frac{1}{2}} - q^{\frac{1}{2}})^2 z', \quad z' \stackrel{\text{def}}{=} z + \epsilon(1 + \tilde{d}), \tag{5.78}$$

$$R_0(z; \boldsymbol{\lambda}) = (q^{-\frac{1}{2}} - q^{\frac{1}{2}})^2 (z'^2 - (q^{-\frac{1}{2}} + q^{\frac{1}{2}})^2 \tilde{d}), \tag{5.79}$$

$$\begin{aligned}
R_{-1}(z; \boldsymbol{\lambda}) &= (q^{-\frac{1}{2}} - q^{\frac{1}{2}})^2 \left(\epsilon'(1+d)z'^2 - \epsilon'(a+b+c+d+\tilde{d} + (ab+bc+ca)q^{-1})z' \right. \\
&\quad \left. + \epsilon'((1-a)(1-b)(1-c)(1-\tilde{d}q^{-1}) \right. \\
&\quad \left. + (a+b+c-1-d\tilde{d} + (ab+bc+ca)q^{-1}))(1+\tilde{d}) \right). \tag{5.80}
\end{aligned}$$

For the shifted parameters the functions B and D are:

$$B(x; \boldsymbol{\lambda} + \boldsymbol{\delta}) = B(x+1; \boldsymbol{\lambda}) \frac{1 - dq^{2x+3}}{1 - dq^{2x+1}}, \quad D(x; \boldsymbol{\lambda} + \boldsymbol{\delta}) = q^2 D(x; \boldsymbol{\lambda}) \frac{1 - dq^{2x-1}}{1 - dq^{2x+1}}. \tag{5.81}$$

From this and (4.14), $\varphi(x)$ is

$$\varphi(x; \boldsymbol{\lambda}) = \frac{q^{-x} - dq^{x+1}}{1 - dq}. \tag{5.82}$$

The constants $f_n(\boldsymbol{\lambda})$ and $b_n(\boldsymbol{\lambda})$ appearing in (4.20)-(4.21) are

$$f_n(\boldsymbol{\lambda}) = \mathcal{E}(n; \boldsymbol{\lambda}), \quad b_n(\boldsymbol{\lambda}) = 1. \quad (5.83)$$

By comparing $B(x)$, $D(x)$ with A_n and C_n , $\phi_0(x)^2$ with d_n^2 , it is evident that the theory is self-dual with the parameter correspondence $(a, b, c, d, \epsilon, \epsilon') \leftrightarrow (a, b, c, \tilde{d}, \epsilon', \epsilon)$.

5.1.6 q -Hahn [KS3.6] (dual to §5.1.7)

Here and in the next subsection (dual q -Hahn) we will consider, for simplicity, the positive parameter range of a and b , which are very slightly different from the standard ones (α, β) ($\alpha q = a, \beta q = b$) for the Hahn and dual Hahn polynomials. The standard parameters for the latter are (γ, δ) . These are the multiplicative counterparts of the parameter change for the Hahn polynomials. The parameter range could well be enlarged.

$$q^\lambda = (a, b, q^N), \quad \boldsymbol{\delta} = (1, 1, -1), \quad \kappa = q^{-1}, \quad (5.84)$$

$$\epsilon = \begin{cases} 1 & \text{for } 0 < a < 1, \quad 0 < b < 1, \\ -1 & \text{for } a > q^{1-N}, \quad b > q^{1-N}, \end{cases} \quad (5.85)$$

$$B(x; \boldsymbol{\lambda}) = \epsilon(1 - aq^x)(q^{x-N} - 1), \quad D(x; \boldsymbol{\lambda}) = \epsilon a q^{-1}(1 - q^x)(q^{x-N} - b), \quad (5.86)$$

$$\mathcal{E}(n; \boldsymbol{\lambda}) = \epsilon(q^{-n} - 1)(1 - abq^{n-1}), \quad \eta(x; \boldsymbol{\lambda}) = q^{-x} - 1, \quad (5.87)$$

$$P_n(\eta(x; \boldsymbol{\lambda}); \boldsymbol{\lambda}) = {}_3\phi_2\left(\begin{matrix} q^{-n}, abq^{n-1}, q^{-x} \\ a, q^{-N} \end{matrix} \middle| q; q\right) = Q_n(1 + \eta(x; \boldsymbol{\lambda}); aq^{-1}, bq^{-1}, N|q), \quad (5.88)$$

$$\phi_0(x; \boldsymbol{\lambda})^2 = \frac{(q; q)_N}{(q; q)_x (q; q)_{N-x}} \frac{(a; q)_x (b; q)_{N-x}}{(b; q)_N a^x}, \quad (5.89)$$

$$d_n(\boldsymbol{\lambda})^2 = \frac{(q; q)_N}{(q; q)_n (q; q)_{N-n}} \frac{(a, abq^{-1}; q)_n}{(abq^N, b; q)_n} \frac{1 - abq^{2n-1}}{1 - abq^{-1}} \times \frac{(b; q)_N a^N}{(ab; q)_N}, \quad (5.90)$$

$$A_n(\boldsymbol{\lambda}) = -\frac{(q^{n-N} - 1)(1 - aq^n)(1 - abq^{n-1})}{(1 - abq^{2n-1})(1 - abq^{2n})}, \quad (5.91)$$

$$C_n(\boldsymbol{\lambda}) = -aq^{n-N-1} \frac{(1 - q^n)(1 - abq^{n+N-1})(1 - bq^{n-1})}{(1 - abq^{2n-2})(1 - abq^{2n-1})}, \quad (5.92)$$

$$R_1(z; \boldsymbol{\lambda}) = (q^{-\frac{1}{2}} - q^{\frac{1}{2}})^2 z', \quad z' \stackrel{\text{def}}{=} z + \epsilon(1 + abq^{-1}), \quad (5.93)$$

$$R_0(z; \boldsymbol{\lambda}) = (q^{-\frac{1}{2}} - q^{\frac{1}{2}})^2 (z'^2 - ab(1 + q^{-1})^2), \quad (5.94)$$

$$R_{-1}(z; \boldsymbol{\lambda}) = (q^{-\frac{1}{2}} - q^{\frac{1}{2}})^2 \left(z'^2 - (a(1 + bq^{-1}) + (1 + aq^{-1})q^{-N})\epsilon z' \right. \\ \left. + a(1 + q^{-1})((a - 1)bq^{-1} + (1 + bq^{-1})q^{-N}) \right), \quad (5.95)$$

$$B(x; \boldsymbol{\lambda} + \boldsymbol{\delta}) = B(x + 1; \boldsymbol{\lambda}), \quad D(x; \boldsymbol{\lambda} + \boldsymbol{\delta}) = q^2 D(x; \boldsymbol{\lambda}), \quad (5.96)$$

$$\varphi(x; \boldsymbol{\lambda}) = q^{-x}, \quad f_n(\boldsymbol{\lambda}) = \mathcal{E}(n; \boldsymbol{\lambda}), \quad b_n(\boldsymbol{\lambda}) = 1. \quad (5.97)$$

Obviously the q -Hahn and dual q -Hahn are dual to each other.

5.1.7 dual q -Hahn [KS3.7] (dual to §5.1.6)

For obvious reasons, we adopt the same parameters (a, b) for the q -Hahn and dual q -Hahn polynomials:

$$q^\lambda = (a, b, q^N), \quad \delta = (1, 0, -1), \quad \kappa = q^{-1}, \quad (5.98)$$

$$\epsilon = \begin{cases} 1 & \text{for } 0 < a < 1, \quad 0 < b < 1, \\ -1 & \text{for } a > q^{1-N}, \quad b > q^{1-N}, \end{cases} \quad (5.99)$$

$$B(x; \lambda) = \frac{(q^{x-N} - 1)(1 - aq^x)(1 - abq^{x-1})}{(1 - abq^{2x-1})(1 - abq^{2x})}, \quad (5.100)$$

$$D(x; \lambda) = aq^{x-N-1} \frac{(1 - q^x)(1 - abq^{x+N-1})(1 - bq^{x-1})}{(1 - abq^{2x-2})(1 - abq^{2x-1})}, \quad (5.101)$$

$$\mathcal{E}(n; \lambda) = q^{-n} - 1, \quad \eta(x; \lambda) = \epsilon(q^{-x} - 1)(1 - abq^{x-1}), \quad (5.102)$$

$$\begin{aligned} P_n(\eta(x; \lambda); \lambda) &= {}_3\phi_2 \left(\begin{matrix} q^{-n}, abq^{x-1}, q^{-x} \\ a, q^{-N} \end{matrix} \middle| q; q \right) \\ &= R_n(\epsilon(1 + abq^{-1} + \eta(x; \lambda)); aq^{-1}, bq^{-1}, N|q), \end{aligned} \quad (5.103)$$

$$\phi_0(x; \lambda)^2 = \frac{(q; q)_N}{(q; q)_x (q; q)_{N-x}} \frac{(a, abq^{-1}; q)_x}{(abq^N, b; q)_x} \frac{1 - abq^{2x-1}}{1 - abq^{-1}}, \quad (5.104)$$

$$d_n(\lambda)^2 = \frac{(q; q)_N}{(q; q)_n (q; q)_{N-n}} \frac{(a; q)_n (b; q)_{N-n}}{(b; q)_N a^n} \times \frac{(b; q)_N a^N}{(ab; q)_N}, \quad (5.105)$$

$$A_n(\lambda) = -\epsilon(1 - aq^n)(q^{n-N} - 1), \quad C_n(\lambda) = -\epsilon aq^{-1}(1 - q^n)(q^{n-N} - b), \quad (5.106)$$

$$R_1(z; \lambda) = (q^{-\frac{1}{2}} - q^{\frac{1}{2}})^2 z', \quad z' \stackrel{\text{def}}{=} z + 1, \quad (5.107)$$

$$R_0(z; \lambda) = (q^{-\frac{1}{2}} - q^{\frac{1}{2}})^2 z'^2, \quad (5.108)$$

$$\begin{aligned} R_{-1}(z; \lambda) &= (q^{-\frac{1}{2}} - q^{\frac{1}{2}})^2 \left(\epsilon(1 + abq^{-1})z'^2 - \epsilon(a(1 + bq^{-1}) + (1 + aq^{-1})q^{-N})z' \right. \\ &\quad \left. + \epsilon a(1 + q^{-1})q^{-N} \right), \end{aligned} \quad (5.109)$$

$$B(x; \lambda + \delta) = B(x + 1; \lambda) \frac{1 - abq^{2x+2}}{1 - abq^{2x}}, \quad D(x; \lambda + \delta) = q^2 D(x; \lambda) \frac{1 - abq^{2x-2}}{1 - abq^{2x}}, \quad (5.110)$$

$$\varphi(x; \lambda) = \frac{q^{-x} - abq^x}{1 - ab}, \quad f_n(\lambda) = \mathcal{E}(n; \lambda), \quad b_n(\lambda) = 1. \quad (5.111)$$

5.1.8 quantum q -Krawtchouk [KS3.14]

$$q^\lambda = (p, q^N), \quad \delta = (1, -1), \quad \kappa = q, \quad p > q^{-N}, \quad (5.112)$$

$$B(x; \lambda) = p^{-1} q^x (q^{x-N} - 1), \quad D(x; \lambda) = (1 - q^x)(1 - p^{-1} q^{x-N-1}), \quad (5.113)$$

$$\mathcal{E}(n; \boldsymbol{\lambda}) = 1 - q^n, \quad \eta(x; \boldsymbol{\lambda}) = q^{-x} - 1, \quad (5.114)$$

$$P_n(\eta(x; \boldsymbol{\lambda}); \boldsymbol{\lambda}) = {}_2\phi_1\left(\begin{matrix} q^{-n}, q^{-x} \\ q^{-N} \end{matrix} \middle| q; pq^{n+1}\right) = K_n^{\text{qtm}}(1 + \eta(x; \boldsymbol{\lambda}); p, N; q), \quad (5.115)$$

$$\phi_0(x; \boldsymbol{\lambda})^2 = \frac{(q; q)_N}{(q; q)_x (q; q)_{N-x}} \frac{p^{-x} q^{x(x-1-N)}}{(p^{-1} q^{-N}; q)_x}, \quad (5.116)$$

$$d_n(\boldsymbol{\lambda})^2 = \frac{(q; q)_N}{(q; q)_n (q; q)_{N-n}} \frac{p^{-n} q^{-Nn}}{(p^{-1} q^{-n}; q)_n} \times (p^{-1} q^{-N}; q)_N, \quad (5.117)$$

$$A_n(\boldsymbol{\lambda}) = -p^{-1} q^{-n-N-1} (1 - q^{N-n}), \quad C_n(\boldsymbol{\lambda}) = -(q^{-n} - 1)(1 - p^{-1} q^{-n}), \quad (5.118)$$

$$R_1(z; \boldsymbol{\lambda}) = (q^{-\frac{1}{2}} - q^{\frac{1}{2}})^2 z', \quad z' \stackrel{\text{def}}{=} z - 1, \quad (5.119)$$

$$R_0(z; \boldsymbol{\lambda}) = (q^{-\frac{1}{2}} - q^{\frac{1}{2}})^2 z'^2, \quad (5.120)$$

$$R_{-1}(z; \boldsymbol{\lambda}) = (q^{-\frac{1}{2}} - q^{\frac{1}{2}})^2 (z'^2 + p^{-1}(1 + p + q^{-N-1})z' + p^{-1}(1 + q^{-1})), \quad (5.121)$$

$$B(x; \boldsymbol{\lambda} + \boldsymbol{\delta}) = q^{-1} B(x + 1; \boldsymbol{\lambda}), \quad D(x; \boldsymbol{\lambda} + \boldsymbol{\delta}) = q D(x; \boldsymbol{\lambda}), \quad (5.122)$$

$$\varphi(x; \boldsymbol{\lambda}) = q^{-x}, \quad f_n(\boldsymbol{\lambda}) = \mathcal{E}(n; \boldsymbol{\lambda}), \quad b_n(\boldsymbol{\lambda}) = 1. \quad (5.123)$$

The dual quantum q -Krawtchouk polynomial with the same parameter and

$$B(x; \boldsymbol{\lambda}) = p^{-1} q^{-x-N-1} (1 - q^{N-x}), \quad D(x; \boldsymbol{\lambda}) = (q^{-x} - 1)(1 - p^{-1} q^{-x}), \quad (5.124)$$

$$\mathcal{E}(n; \boldsymbol{\lambda}) = q^{-n} - 1, \quad \eta(x; \boldsymbol{\lambda}) = 1 - q^x, \quad (5.125)$$

$$A_n(\boldsymbol{\lambda}) = -p^{-1} q^n (q^{n-N} - 1), \quad C_n(\boldsymbol{\lambda}) = -(1 - q^n)(1 - p^{-1} q^{n-N-1}), \quad (5.126)$$

has not been reported in Koekoek-Swarttouw [14]. It is interesting to note that the functions $B(x)$ and $D(x)$ in (5.124) are related to those of the affine q -Krawtchouk polynomial §5.1.10 with the change of variable $x \leftrightarrow N - x$, $B \leftrightarrow D$:

$$B(x) = D^{\text{affine } q\text{-Krawtchouk}}(N - x), \quad D(x) = B^{\text{affine } q\text{-Krawtchouk}}(N - x), \quad (5.127)$$

with the change of the parameter $p^{\text{affine } q\text{-Krawtchouk}} \rightarrow p^{-1} q^{-N-1}$.

5.1.9 q -Krawtchouk [KS3.15] and dual q -Krawtchouk [KS3.17]

$$q^\lambda = (p, q^N), \quad \boldsymbol{\delta} = (2, -1), \quad \kappa = q^{-1}, \quad p > 0, \quad (5.128)$$

$$B(x; \boldsymbol{\lambda}) = q^{x-N} - 1, \quad D(x; \boldsymbol{\lambda}) = p(1 - q^x), \quad (5.129)$$

$$\mathcal{E}(n; \boldsymbol{\lambda}) = (q^{-n} - 1)(1 + pq^n), \quad \eta(x; \boldsymbol{\lambda}) = q^{-x} - 1, \quad (5.130)$$

$$P_n(\eta(x; \boldsymbol{\lambda}); \boldsymbol{\lambda}) = {}_3\phi_2\left(\begin{matrix} q^{-n}, q^{-x}, -pq^n \\ q^{-N}, 0 \end{matrix} \middle| q; q\right) = K_n(1 + \eta(x; \boldsymbol{\lambda}); p, N; q), \quad (5.131)$$

$$\phi_0(x; \boldsymbol{\lambda})^2 = \frac{(q; q)_N}{(q; q)_x (q; q)_{N-x}} p^{-x} q^{\frac{1}{2}x(x-1) - xN}, \quad (5.132)$$

$$d_n(\boldsymbol{\lambda})^2 = \frac{(q; q)_N}{(q; q)_n (q; q)_{N-n}} \frac{(-p; q)_n}{(-pq^{N+1}; q)_n} \frac{1 + pq^{2n}}{1 + p} \times \frac{p^N q^{\frac{1}{2}N(N+1)}}{(-pq; q)_N}, \quad (5.133)$$

$$A_n(\boldsymbol{\lambda}) = -\frac{(q^{n-N} - 1)(1 + pq^n)}{(1 + pq^{2n})(1 + pq^{2n+1})}, \quad C_n(\boldsymbol{\lambda}) = -pq^{2n-N-1} \frac{(1 - q^n)(1 + pq^{n+N})}{(1 + pq^{2n-1})(1 + pq^{2n})}, \quad (5.134)$$

$$R_1(z; \boldsymbol{\lambda}) = (q^{-\frac{1}{2}} - q^{\frac{1}{2}})^2 z', \quad z' \stackrel{\text{def}}{=} z + 1 - p, \quad (5.135)$$

$$R_0(z; \boldsymbol{\lambda}) = (q^{-\frac{1}{2}} - q^{\frac{1}{2}})^2 (z'^2 + p(q^{-\frac{1}{2}} + q^{\frac{1}{2}})^2), \quad (5.136)$$

$$R_{-1}(z; \boldsymbol{\lambda}) = (q^{-\frac{1}{2}} - q^{\frac{1}{2}})^2 (z'^2 + (p - q^{-N})z' + p(1 + q^{-1})(1 - q^{-N})), \quad (5.137)$$

$$B(x; \boldsymbol{\lambda} + \boldsymbol{\delta}) = B(x + 1; \boldsymbol{\lambda}), \quad D(x; \boldsymbol{\lambda} + \boldsymbol{\delta}) = q^2 D(x; \boldsymbol{\lambda}), \quad (5.138)$$

$$\varphi(x; \boldsymbol{\lambda}) = q^{-x}, \quad f_n(\boldsymbol{\lambda}) = \mathcal{E}(n; \boldsymbol{\lambda}), \quad b_n(\boldsymbol{\lambda}) = 1. \quad (5.139)$$

The dual q -Krawtchouk polynomial in this parametrisation has

$$B(x; \boldsymbol{\lambda}) = \frac{(q^{x-N} - 1)(1 + pq^x)}{(1 + pq^{2x})(1 + pq^{2x+1})}, \quad D(x; \boldsymbol{\lambda}) = pq^{2x-N-1} \frac{(1 - q^x)(1 + pq^{x+N})}{(1 + pq^{2x-1})(1 + pq^{2x})}, \quad (5.140)$$

$$\mathcal{E}(n; \boldsymbol{\lambda}) = q^{-n} - 1, \quad \eta(x; \boldsymbol{\lambda}) = (q^{-x} - 1)(1 + pq^x), \quad (5.141)$$

$$A_n(\boldsymbol{\lambda}) = -(q^{n-N} - 1), \quad C_n(\boldsymbol{\lambda}) = -p(1 - q^n). \quad (5.142)$$

These are to be compared with the standard parametrisation of the dual q -Krawtchouk [KS3.17], in which the parameter c is identified as

$$c = -pq^N. \quad (5.143)$$

dual q -Krawtchouk [KS3.17] in the standard parametrisation

$$q^\lambda = (c, q^N), \quad \boldsymbol{\delta} = (0, -1), \quad \kappa = q^{-1}, \quad c < 0, \quad (5.144)$$

$$B(x; \boldsymbol{\lambda}) = \frac{(q^{x-N} - 1)(1 - cq^{x-N})}{(1 - cq^{2x-N})(1 - cq^{2x+1-N})}, \quad (5.145)$$

$$D(x; \boldsymbol{\lambda}) = -cq^{2x-2N-1} \frac{(1 - q^x)(1 - cq^x)}{(1 - cq^{2x-1-N})(1 - cq^{2x-N})}, \quad (5.146)$$

$$\mathcal{E}(n; \boldsymbol{\lambda}) = q^{-n} - 1, \quad \eta(x; \boldsymbol{\lambda}) = (q^{-x} - 1)(1 - cq^{x-N}), \quad (5.147)$$

$$P_n(\eta(x; \boldsymbol{\lambda}); \boldsymbol{\lambda}) = {}_3\phi_2 \left(\begin{matrix} q^{-n}, q^{-x}, cq^{x-N} \\ q^{-N}, 0 \end{matrix} \middle| q; q \right) = K_n(1 + cq^{-N} + \eta(x; \boldsymbol{\lambda}); c, N|q), \quad (5.148)$$

$$\phi_0(x; \boldsymbol{\lambda})^2 = \frac{(q; q)_N}{(q; q)_x (q; q)_{N-x}} \frac{(cq^{-N}; q)_x q^{Nx - \frac{1}{2}x(x+1)}}{(cq; q)_x (-c)^x} \frac{1 - cq^{2x-N}}{1 - cq^{-N}}, \quad (5.149)$$

$$d_n(\boldsymbol{\lambda})^2 = \frac{(q; q)_N}{(q; q)_n (q; q)_{N-n}} (-c)^{-n} q^{\frac{1}{2}n(n-1)} \times \frac{1}{(c^{-1}; q)_N}, \quad (5.150)$$

$$A_n(\boldsymbol{\lambda}) = -(q^{n-N} - 1), \quad C_n(\boldsymbol{\lambda}) = cq^{-N}(1 - q^n), \quad (5.151)$$

$$R_1(z; \boldsymbol{\lambda}) = (q^{-\frac{1}{2}} - q^{\frac{1}{2}})^2 z', \quad z' \stackrel{\text{def}}{=} z + 1, \quad (5.152)$$

$$R_0(z; \boldsymbol{\lambda}) = (q^{-\frac{1}{2}} - q^{\frac{1}{2}})^2 z'^2, \quad (5.153)$$

$$R_{-1}(z; \boldsymbol{\lambda}) = (q^{-\frac{1}{2}} - q^{\frac{1}{2}})^2 ((1 + cq^{-N})z'^2 - (1 + c)q^{-N}z'), \quad (5.154)$$

$$B(x; \boldsymbol{\lambda} + \boldsymbol{\delta}) = B(x + 1; \boldsymbol{\lambda}) \frac{1 - cq^{2x+3-N}}{1 - cq^{2x+1-N}}, \quad D(x; \boldsymbol{\lambda} + \boldsymbol{\delta}) = q^2 D(x; \boldsymbol{\lambda}) \frac{1 - cq^{2x-1-N}}{1 - cq^{2x+1-N}}, \quad (5.155)$$

$$\varphi(x; \boldsymbol{\lambda}) = \frac{q^{-x} - cq^{1-N}q^x}{1 - cq^{1-N}}, \quad f_n(\boldsymbol{\lambda}) = \mathcal{E}(n; \boldsymbol{\lambda}), \quad b_n(\boldsymbol{\lambda}) = 1. \quad (5.156)$$

5.1.10 affine q -Krawtchouk [KS3.16] (self-dual)

$$q^\lambda = (p, q^N), \quad \boldsymbol{\delta} = (1, -1), \quad \kappa = q^{-1}, \quad 0 < p < q^{-1}, \quad (5.157)$$

$$B(x; \boldsymbol{\lambda}) = (q^{x-N} - 1)(1 - pq^{x+1}), \quad D(x; \boldsymbol{\lambda}) = pq^{x-N}(1 - q^x), \quad (5.158)$$

$$\mathcal{E}(n; \boldsymbol{\lambda}) = q^{-n} - 1, \quad \eta(x; \boldsymbol{\lambda}) = q^{-x} - 1, \quad (5.159)$$

$$P_n(\eta(x; \boldsymbol{\lambda}); \boldsymbol{\lambda}) = {}_3\phi_2 \left(\begin{matrix} q^{-n}, q^{-x}, 0 \\ pq, q^{-N} \end{matrix} \middle| q; q \right) = K_n^{\text{aff}}(1 + \eta(x; \boldsymbol{\lambda}); p, N; q), \quad (5.160)$$

$$\phi_0(x; \boldsymbol{\lambda})^2 = \frac{(q; q)_N}{(q; q)_x (q; q)_{N-x}} \frac{(pq; q)_x}{(pq)^x}, \quad (5.161)$$

$$d_n(\boldsymbol{\lambda})^2 = \frac{(q; q)_N}{(q; q)_n (q; q)_{N-n}} \frac{(pq; q)_n}{(pq)^n} \times (pq)^N, \quad (5.162)$$

$$A_n(\boldsymbol{\lambda}) = -(q^{n-N} - 1)(1 - pq^{n+1}), \quad C_n(\boldsymbol{\lambda}) = -pq^{n-N}(1 - q^n), \quad (5.163)$$

$$R_1(z; \boldsymbol{\lambda}) = (q^{-\frac{1}{2}} - q^{\frac{1}{2}})^2 z', \quad z' \stackrel{\text{def}}{=} z + 1, \quad (5.164)$$

$$R_0(z; \boldsymbol{\lambda}) = (q^{-\frac{1}{2}} - q^{\frac{1}{2}})^2 z'^2, \quad (5.165)$$

$$R_{-1}(z; \boldsymbol{\lambda}) = (q^{-\frac{1}{2}} - q^{\frac{1}{2}})^2 (z'^2 - (pq + (1 + p)q^{-N})z' + p(1 + q)q^{-N}), \quad (5.166)$$

$$B(x; \boldsymbol{\lambda} + \boldsymbol{\delta}) = B(x + 1; \boldsymbol{\lambda}), \quad D(x; \boldsymbol{\lambda} + \boldsymbol{\delta}) = q^2 D(x; \boldsymbol{\lambda}), \quad (5.167)$$

$$\varphi(x; \boldsymbol{\lambda}) = q^{-x}, \quad f_n(\boldsymbol{\lambda}) = \mathcal{E}(n; \boldsymbol{\lambda}), \quad b_n(\boldsymbol{\lambda}) = 1. \quad (5.168)$$

5.2 Infinite Dimensional Case (II)

In this subsection x takes an infinite range of values

$$x = 0, 1, 2, \dots,$$

In contrast to the finite dimensional case, the structure of the polynomials is severely constrained by the asymptotic forms of the functions $B(x)$ and $D(x)$, which are determined in

Appendix A. It is easy to see for $\eta(x) = x$ (i) (4.72), a quadratic spectrum $\mathcal{E}(n) \sim n(n + \alpha)$ is not possible. In order this to happen in (2.28), $B(x)$ and $D(x)$ must have the same coefficient β with the different sign for the leading quadratic term, $B(x) \sim \beta x^2$, $D(x) \sim -\beta x^2$, which violates the positivity of $B(x)$ and $D(x)$. It is also easy to see $\eta(x) = \epsilon'x(x + d)$ (ii) (4.73) is not possible. If the leading power of $B(x)$ and $D(x)$ is quadratic, the coefficient must be the same with the same sign, $B(x) \sim \beta x^2$, $D(x) \sim \beta x^2$, $\beta > 0$. Then for large n (2.28) gives a negative leading term

$$\tilde{\mathcal{H}}\eta(x)^n \sim -\beta n^2 \eta(x)^n + \dots, \quad (5.169)$$

which cannot give a normalisable eigenvector for a positive definite Hamiltonian. If the leading power of $B(x)$ and $D(x)$ is linear, the results in Appendix A (A.12) and (A.14) tell that $r_1^{(0)} = 0$ and $B(x)$ and $D(x)$ have the opposite sign leading terms

$$B(x) = \frac{r_{-1}^{(0)}\eta(x)}{\eta(1)B(0)(\eta(x+1) - \eta(x-1))} + \text{lower order}, \quad (5.170)$$

$$D(x) = -\frac{r_{-1}^{(0)}\eta(x)}{\eta(1)B(0)(\eta(x+1) - \eta(x-1))} + \text{lower order}. \quad (5.171)$$

This simply contradicts the positivity of $B(x)$ and $D(x)$. If the leading power of $B(x)$ and $D(x)$ is a constant or a negative power in x , then eigenpolynomials in $\eta(x)$ of $\tilde{\mathcal{H}}$ do not exist. Thus we have two self-dual polynomials in x , the Meixner §5.2.1 and Charlier §5.2.2.

As for the q -polynomials, the consequences of the asymptotic behaviours of $B(x)$ and $D(x)$ are easy to see. If these functions are *bounded* as in the cases of the sinusoidal coordinates $\eta(x) = q^{-x} - 1$ (iv) (4.75), and $\eta(x) = \epsilon'(q^{-x} - 1)(1 - dq^x)$ (v) (4.76), the Hamiltonians are *bounded* and eigenvalues cannot take *unbounded* forms $\mathcal{E}(n) = q^{-n} - 1$, $(q^{-n} - 1)(1 - \alpha q^n)$ and vice versa. As mentioned in Appendix A, there is no self-dual q -polynomial of infinite dimension.

5.2.1 Meixner [KS1.9] (self-dual)

This is the first example of the self-dual polynomial (5.175), which is symmetric under interchange $x \leftrightarrow n$, with $\mathcal{E}(n) = n$ and $\eta(x) = x$ (5.174). Compare (5.173)↔(5.177), (5.176):

$$\boldsymbol{\lambda} = (\beta, c), \quad \boldsymbol{\delta} = (1, 0), \quad \kappa = 1, \quad \beta > 0, \quad 0 < c < 1, \quad (5.172)$$

$$B(x; \boldsymbol{\lambda}) = \frac{c}{1-c}(x + \beta), \quad D(x; \boldsymbol{\lambda}) = \frac{1}{1-c}x, \quad (5.173)$$

$$\mathcal{E}(n; \boldsymbol{\lambda}) = n, \quad \eta(x; \boldsymbol{\lambda}) = x, \quad (5.174)$$

$$P_n(\eta(x; \boldsymbol{\lambda}); \boldsymbol{\lambda}) = {}_2F_1\left(\begin{matrix} -n, -x \\ \beta \end{matrix} \middle| 1 - c^{-1}\right) = M_n(\eta(x; \boldsymbol{\lambda}); \beta, c), \quad (5.175)$$

$$\phi_0(x; \boldsymbol{\lambda})^2 = \frac{(\beta)_x c^x}{x!}, \quad d_n(\boldsymbol{\lambda})^2 = \frac{(\beta)_n c^n}{n!} \times (1 - c)^\beta, \quad (5.176)$$

$$A_n(\boldsymbol{\lambda}) = -\frac{c}{1-c}(n + \beta), \quad C_n(\boldsymbol{\lambda}) = -\frac{1}{1-c}n, \quad (5.177)$$

$$R_1(z; \boldsymbol{\lambda}) = 0, \quad R_0(z; \boldsymbol{\lambda}) = 1, \quad R_{-1}(z; \boldsymbol{\lambda}) = -\frac{1+c}{1-c}z - \frac{\beta c}{1-c}, \quad (5.178)$$

$$B(x; \boldsymbol{\lambda} + \boldsymbol{\delta}) = B(x + 1; \boldsymbol{\lambda}), \quad D(x; \boldsymbol{\lambda} + \boldsymbol{\delta}) = D(x; \boldsymbol{\lambda}), \quad (5.179)$$

$$\varphi(x; \boldsymbol{\lambda}) = 1, \quad f_n(\boldsymbol{\lambda}) = \mathcal{E}(n; \boldsymbol{\lambda}), \quad b_n(\boldsymbol{\lambda}) = 1. \quad (5.180)$$

5.2.2 Charlier [KS1.12] (self-dual)

$$\boldsymbol{\lambda} = a, \quad \boldsymbol{\delta} = 0, \quad \kappa = 1, \quad a > 0, \quad (5.181)$$

$$B(x; \boldsymbol{\lambda}) = a, \quad D(x; \boldsymbol{\lambda}) = x, \quad (5.182)$$

$$\mathcal{E}(n; \boldsymbol{\lambda}) = n, \quad \eta(x; \boldsymbol{\lambda}) = x, \quad (5.183)$$

$$P_n(\eta(x; \boldsymbol{\lambda}); \boldsymbol{\lambda}) = {}_2F_0\left(\begin{matrix} -n, -x \\ - \end{matrix} \middle| -a^{-1}\right) = C_n(\eta(x; \boldsymbol{\lambda}); a), \quad (5.184)$$

$$\phi_0(x; \boldsymbol{\lambda})^2 = \frac{a^x}{x!}, \quad d_n(\boldsymbol{\lambda})^2 = \frac{a^n}{n!} \times e^{-a}, \quad (5.185)$$

$$A_n(\boldsymbol{\lambda}) = -a, \quad C_n(\boldsymbol{\lambda}) = -n, \quad (5.186)$$

$$R_1(z; \boldsymbol{\lambda}) = 0, \quad R_0(z; \boldsymbol{\lambda}) = 1, \quad R_{-1}(z; \boldsymbol{\lambda}) = -z - a, \quad (5.187)$$

$$B(x; \boldsymbol{\lambda} + \boldsymbol{\delta}) = B(x + 1; \boldsymbol{\lambda}), \quad D(x; \boldsymbol{\lambda} + \boldsymbol{\delta}) = D(x; \boldsymbol{\lambda}), \quad (5.188)$$

$$\varphi(x; \boldsymbol{\lambda}) = 1, \quad f_n(\boldsymbol{\lambda}) = \mathcal{E}(n; \boldsymbol{\lambda}), \quad b_n(\boldsymbol{\lambda}) = 1. \quad (5.189)$$

5.2.3 little q -Jacobi [KS3.12]

The universal normalisation (2.31) differs from the standard normalisation of the little q -Jacobi polynomial p_n as shown explicitly in (5.193). This does not affect the normalisation measure $\phi_0(x)^2$ (5.194) but the normalisation constants d_n^2 (5.195) and the coefficients of the three term recurrence A_n and C_n (5.196)–(5.197) are different from the standard ones:

$$q^\boldsymbol{\lambda} = (a, b), \quad \boldsymbol{\delta} = (1, 1), \quad \kappa = q^{-1}, \quad 0 < a < q^{-1}, \quad b < q^{-1}, \quad (5.190)$$

$$B(x; \boldsymbol{\lambda}) = a(q^{-x} - bq), \quad D(x; \boldsymbol{\lambda}) = q^{-x} - 1, \quad (5.191)$$

$$\mathcal{E}(n; \boldsymbol{\lambda}) = (q^{-n} - 1)(1 - abq^{n+1}), \quad \eta(x; \boldsymbol{\lambda}) = 1 - q^x, \quad (5.192)$$

$$\begin{aligned}
P_n(\eta(x; \boldsymbol{\lambda}); \boldsymbol{\lambda}) &= (-a)^{-n} q^{-\frac{1}{2}n(n+1)} \frac{(aq; q)_n}{(bq; q)_n} {}_2\phi_1 \left(\begin{matrix} q^{-n}, abq^{n+1} \\ aq \end{matrix} \middle| q; q^{x+1} \right) \\
&= (-a)^{-n} q^{-\frac{1}{2}n(n+1)} \frac{(aq; q)_n}{(bq; q)_n} p_n(1 - \eta(x; \boldsymbol{\lambda}); a, b|q), \tag{5.193}
\end{aligned}$$

$$\phi_0(x; \boldsymbol{\lambda})^2 = \frac{(bq; q)_x}{(q; q)_x} (aq)^x, \tag{5.194}$$

$$d_n(\boldsymbol{\lambda})^2 = \frac{(bq, abq; q)_n a^n q^{n^2}}{(q, aq; q)_n} \frac{1 - abq^{2n+1}}{1 - abq} \times \frac{(aq; q)_\infty}{(abq^2; q)_\infty}, \tag{5.195}$$

$$A_n(\boldsymbol{\lambda}) = -aq^{2n+1} \frac{(1 - bq^{n+1})(1 - abq^{n+1})}{(1 - abq^{2n+1})(1 - abq^{2n+2})}, \tag{5.196}$$

$$C_n(\boldsymbol{\lambda}) = -\frac{(1 - q^n)(1 - aq^n)}{(1 - abq^{2n})(1 - abq^{2n+1})}, \tag{5.197}$$

$$R_1(z; \boldsymbol{\lambda}) = (q^{-\frac{1}{2}} - q^{\frac{1}{2}})^2 z', \quad z' \stackrel{\text{def}}{=} z + 1 + abq, \tag{5.198}$$

$$R_0(z; \boldsymbol{\lambda}) = (q^{-\frac{1}{2}} - q^{\frac{1}{2}})^2 (z'^2 - ab(1 + q)^2), \tag{5.199}$$

$$R_{-1}(z; \boldsymbol{\lambda}) = (q^{-\frac{1}{2}} - q^{\frac{1}{2}})^2 (-z'^2 + (1 + a)z' - a(1 + q)(1 - bq)), \tag{5.200}$$

$$B(x; \boldsymbol{\lambda} + \boldsymbol{\delta}) = q^2 B(x + 1; \boldsymbol{\lambda}), \quad D(x; \boldsymbol{\lambda} + \boldsymbol{\delta}) = D(x; \boldsymbol{\lambda}), \tag{5.201}$$

$$\varphi(x; \boldsymbol{\lambda}) = q^x, \quad f_n(\boldsymbol{\lambda}) = \mathcal{E}(n; \boldsymbol{\lambda}), \quad b_n(\boldsymbol{\lambda}) = 1. \tag{5.202}$$

The dual little q -Jacobi polynomial with the same parameter and

$$B(x; \boldsymbol{\lambda}) = aq^{2x+1} \frac{(1 - bq^{x+1})(1 - abq^{x+1})}{(1 - abq^{2x+1})(1 - abq^{2x+2})}, \quad D(x; \boldsymbol{\lambda}) = \frac{(1 - q^x)(1 - aq^x)}{(1 - abq^{2x})(1 - abq^{2x+1})}, \tag{5.203}$$

$$\mathcal{E}(n; \boldsymbol{\lambda}) = 1 - q^n, \quad \eta(x; \boldsymbol{\lambda}) = (q^{-x} - 1)(1 - abq^{x+1}), \tag{5.204}$$

$$A_n(\boldsymbol{\lambda}) = -a(q^{-n} - bq), \quad C_n(\boldsymbol{\lambda}) = -(q^{-n} - 1), \tag{5.205}$$

has not been reported in Koekoek-Swarttouw [14]. The dual little q -Jacobi polynomial was introduced by Atakishiyev and Klimyk [25], which has the same three term recurrence relation and the orthogonality measure as above. The difference equation was not mentioned there.

5.2.4 q -Meixner [KS3.13]

$$q^\lambda = (b, c), \quad \boldsymbol{\delta} = (1, -1), \quad \kappa = q, \quad 0 < b < q^{-1}, \quad c > 0, \tag{5.206}$$

$$B(x; \boldsymbol{\lambda}) = cq^x(1 - bq^{x+1}), \quad D(x; \boldsymbol{\lambda}) = (1 - q^x)(1 + bcq^x), \tag{5.207}$$

$$\mathcal{E}(n; \boldsymbol{\lambda}) = 1 - q^n, \quad \eta(x; \boldsymbol{\lambda}) = q^{-x} - 1, \tag{5.208}$$

$$P_n(\eta(x; \boldsymbol{\lambda}); \boldsymbol{\lambda}) = {}_2\phi_1 \left(\begin{matrix} q^{-n}, q^{-x} \\ bq \end{matrix} \middle| q; -c^{-1}q^{n+1} \right) = M_n(1 + \eta(x; \boldsymbol{\lambda}); b, c; q), \tag{5.209}$$

$$\phi_0(x; \boldsymbol{\lambda})^2 = \frac{(bq; q)_x}{(q, -bcq; q)_x} c^x q^{\frac{1}{2}x(x-1)}, \quad d_n(\boldsymbol{\lambda})^2 = \frac{(bq; q)_n}{(q, -c^{-1}q; q)_n} \times \frac{(-bcq; q)_\infty}{(-c; q)_\infty}, \quad (5.210)$$

$$A_n(\boldsymbol{\lambda}) = -cq^{-n-1}(q^{-n} - bq), \quad C_n(\boldsymbol{\lambda}) = -(q^{-n} - 1)(1 + cq^{-n}), \quad (5.211)$$

$$R_1(z; \boldsymbol{\lambda}) = (q^{-\frac{1}{2}} - q^{\frac{1}{2}})^2 z', \quad z' \stackrel{\text{def}}{=} z - 1, \quad (5.212)$$

$$R_0(z; \boldsymbol{\lambda}) = (q^{-\frac{1}{2}} - q^{\frac{1}{2}})^2 z'^2, \quad (5.213)$$

$$R_{-1}(z; \boldsymbol{\lambda}) = (q^{-\frac{1}{2}} - q^{\frac{1}{2}})^2 (z'^2 + (1 - c - bc)z' - c(1 + q^{-1})), \quad (5.214)$$

$$B(x; \boldsymbol{\lambda} + \boldsymbol{\delta}) = q^{-2}B(x + 1; \boldsymbol{\lambda}), \quad D(x; \boldsymbol{\lambda} + \boldsymbol{\delta}) = D(x; \boldsymbol{\lambda}), \quad (5.215)$$

$$\varphi(x; \boldsymbol{\lambda}) = q^{-x}, \quad f_n(\boldsymbol{\lambda}) = \mathcal{E}(n; \boldsymbol{\lambda}), \quad b_n(\boldsymbol{\lambda}) = 1. \quad (5.216)$$

The dual q -Meixner polynomial with the same parameter and

$$B(x; \boldsymbol{\lambda}) = cq^{-x-1}(q^{-x} - bq), \quad D(x; \boldsymbol{\lambda}) = (q^{-x} - 1)(1 + cq^{-x}), \quad (5.217)$$

$$\mathcal{E}(n; \boldsymbol{\lambda}) = q^{-n} - 1, \quad \eta(x; \boldsymbol{\lambda}) = 1 - q^x, \quad (5.218)$$

$$A_n(\boldsymbol{\lambda}) = -cq^n(1 - bq^{n+1}), \quad C_n(\boldsymbol{\lambda}) = -(1 - q^n)(1 + bcq^n), \quad (5.219)$$

has not been reported in Koekoek-Swarttouw [14]. This can be considered another q -version of the Meixner polynomial §5.2.1 with the sinusoidal coordinate $\eta(x) = 1 - q^x$.

5.2.5 little q -Laguerre/Wall [KS3.20] (dual to §5.2.6)

The universal normalisation (2.31) differs from the standard normalisation of the little q -Laguerre/Wall polynomial p_n as shown explicitly in (5.223). As other examples, d_n^2 , A_n and C_n are different from the standard ones:

$$q^\lambda = a, \quad \boldsymbol{\delta} = 1, \quad \kappa = q^{-1}, \quad 0 < a < q^{-1}, \quad (5.220)$$

$$B(x; \boldsymbol{\lambda}) = aq^{-x}, \quad D(x; \boldsymbol{\lambda}) = q^{-x} - 1, \quad (5.221)$$

$$\mathcal{E}(n; \boldsymbol{\lambda}) = q^{-n} - 1, \quad \eta(x; \boldsymbol{\lambda}) = 1 - q^x, \quad (5.222)$$

$$P_n(\eta(x; \boldsymbol{\lambda}); \boldsymbol{\lambda}) = {}_2\phi_0\left(\begin{matrix} q^{-n}, q^{-x} \\ - \end{matrix} \middle| q; a^{-1}q^x\right) = (a^{-1}q^{-n}; q)_n p_n(1 - \eta(x; \boldsymbol{\lambda}); a|q), \quad (5.223)$$

$$\phi_0(x; \boldsymbol{\lambda})^2 = \frac{(aq)^x}{(q; q)_x}, \quad d_n(\boldsymbol{\lambda})^2 = \frac{a^n q^{n^2}}{(q, aq; q)_n} \times (aq; q)_\infty, \quad (5.224)$$

$$A_n(\boldsymbol{\lambda}) = -aq^{2n+1}, \quad C_n(\boldsymbol{\lambda}) = -(1 - q^n)(1 - aq^n), \quad (5.225)$$

$$R_1(z; \boldsymbol{\lambda}) = (q^{-\frac{1}{2}} - q^{\frac{1}{2}})^2 z', \quad z' \stackrel{\text{def}}{=} z + 1, \quad (5.226)$$

$$R_0(z; \boldsymbol{\lambda}) = (q^{-\frac{1}{2}} - q^{\frac{1}{2}})^2 z'^2, \quad (5.227)$$

$$R_{-1}(z; \boldsymbol{\lambda}) = (q^{-\frac{1}{2}} - q^{\frac{1}{2}})^2 (-z'^2 + (1 + a)z' - a(1 + q)), \quad (5.228)$$

$$B(x; \boldsymbol{\lambda} + \boldsymbol{\delta}) = q^2 B(x + 1; \boldsymbol{\lambda}), \quad D(x; \boldsymbol{\lambda} + \boldsymbol{\delta}) = D(x; \boldsymbol{\lambda}), \quad (5.229)$$

$$\varphi(x; \boldsymbol{\lambda}) = q^x, \quad f_n(\boldsymbol{\lambda}) = \mathcal{E}(n; \boldsymbol{\lambda}), \quad b_n(\boldsymbol{\lambda}) = 1. \quad (5.230)$$

The dual little q -Laguerre/Wall polynomial with the same parameter and

$$B(x; \boldsymbol{\lambda}) = aq^{2x+1}, \quad D(x; \boldsymbol{\lambda}) = (1 - q^x)(1 - aq^x), \quad (5.231)$$

$$\mathcal{E}(n; \boldsymbol{\lambda}) = 1 - q^n, \quad \eta(x; \boldsymbol{\lambda}) = q^{-x} - 1, \quad (5.232)$$

$$A_n(\boldsymbol{\lambda}) = -aq^{-n}, \quad C_n(\boldsymbol{\lambda}) = -(q^{-n} - 1), \quad (5.233)$$

is Al-Salam Carlitz II polynomial as seen below. This was reported by Atakishiyev and Klimyk [25].

5.2.6 Al-Salam-Carlitz II [KS3.25] (dual to §5.2.5)

The universal normalisation (2.31) differs from the standard normalisation of the Al-Salam-Carlitz II polynomial $V_n^{(a)}$ as shown explicitly in (5.237). As other examples, d_n^2 , A_n and C_n are different from the standard ones:

$$q^\lambda = a, \quad \boldsymbol{\delta} = 0, \quad \kappa = q, \quad 0 < a < q^{-1}, \quad (5.234)$$

$$B(x; \boldsymbol{\lambda}) = aq^{2x+1}, \quad D(x; \boldsymbol{\lambda}) = (1 - q^x)(1 - aq^x), \quad (5.235)$$

$$\mathcal{E}(n; \boldsymbol{\lambda}) = 1 - q^n, \quad \eta(x; \boldsymbol{\lambda}) = q^{-x} - 1, \quad (5.236)$$

$$P_n(\eta(x; \boldsymbol{\lambda}); \boldsymbol{\lambda}) = {}_2\phi_0 \left(\begin{matrix} q^{-n}, q^{-x} \\ - \end{matrix} \middle| q; a^{-1}q^n \right) = (-a^{-1}q^{\frac{1}{2}(n-1)})^n V_n^{(a)}(1 + \eta(x; \boldsymbol{\lambda}); q), \quad (5.237)$$

$$\phi_0(x; \boldsymbol{\lambda})^2 = \frac{a^x q^{x^2}}{(q, aq; q)_x}, \quad d_n(\boldsymbol{\lambda})^2 = \frac{(aq)^n}{(q; q)_n} \times (aq; q)_\infty, \quad (5.238)$$

$$A_n(\boldsymbol{\lambda}) = -aq^{-n}, \quad C_n(\boldsymbol{\lambda}) = -(q^{-n} - 1), \quad (5.239)$$

$$R_1(z; \boldsymbol{\lambda}) = (q^{-\frac{1}{2}} - q^{\frac{1}{2}})^2 z', \quad z' \stackrel{\text{def}}{=} z - 1, \quad (5.240)$$

$$R_0(z; \boldsymbol{\lambda}) = (q^{-\frac{1}{2}} - q^{\frac{1}{2}})^2 z'^2, \quad (5.241)$$

$$R_{-1}(z; \boldsymbol{\lambda}) = (q^{-\frac{1}{2}} - q^{\frac{1}{2}})^2 (z'^2 + (1 + a)z'), \quad (5.242)$$

$$B(x; \boldsymbol{\lambda} + \boldsymbol{\delta}) = q^{-2} B(x + 1; \boldsymbol{\lambda}), \quad D(x; \boldsymbol{\lambda} + \boldsymbol{\delta}) = D(x; \boldsymbol{\lambda}), \quad (5.243)$$

$$\varphi(x; \boldsymbol{\lambda}) = q^{-x}, \quad f_n(\boldsymbol{\lambda}) = \mathcal{E}(n; \boldsymbol{\lambda}), \quad b_n(\boldsymbol{\lambda}) = 1. \quad (5.244)$$

5.2.7 alternative q -Charlier [KS3.22]

The universal normalisation (2.31) differs from the standard normalisation of the alternative q -Charlier polynomial K_n as shown explicitly in (5.248). As other examples, d_n^2 , A_n and C_n

are different from the standard ones:

$$q^\lambda = a, \quad \delta = 2, \quad \kappa = q^{-1}, \quad a > 0, \quad (5.245)$$

$$B(x; \lambda) = a, \quad D(x; \lambda) = q^{-x} - 1, \quad (5.246)$$

$$\mathcal{E}(n; \lambda) = (q^{-n} - 1)(1 + aq^n), \quad \eta(x; \lambda) = 1 - q^x, \quad (5.247)$$

$$\begin{aligned} P_n(\eta(x; \lambda); \lambda) &= q^{nx} {}_2\phi_1 \left(\begin{matrix} q^{-n}, q^{-x} \\ 0 \end{matrix} \middle| q; -a^{-1}q^{-n+1} \right) \\ &= (-aq^n)^{-n} K_n(1 - \eta(x; \lambda); a; q), \end{aligned} \quad (5.248)$$

$$\phi_0(x; \lambda)^2 = \frac{a^x q^{\frac{1}{2}x(x+1)}}{(q; q)_x}, \quad d_n(\lambda)^2 = \frac{a^n q^{\frac{1}{2}n(3n-1)}}{(q; q)_n} \frac{(-a; q)_\infty}{(-aq^n; q)_\infty} \frac{1 + aq^{2n}}{1 + a} \times \frac{1}{(-aq; q)_\infty}, \quad (5.249)$$

$$A_n(\lambda) = -aq^{3n+1} \frac{1 + aq^n}{(1 + aq^{2n})(1 + aq^{2n+1})}, \quad C_n(\lambda) = -\frac{1 - q^n}{(1 + aq^{2n-1})(1 + aq^{2n})}, \quad (5.250)$$

$$R_1(z; \lambda) = (q^{-\frac{1}{2}} - q^{\frac{1}{2}})^2 z', \quad z' \stackrel{\text{def}}{=} z + 1 - a, \quad (5.251)$$

$$R_0(z; \lambda) = (q^{-\frac{1}{2}} - q^{\frac{1}{2}})^2 (z'^2 + a(q^{-\frac{1}{2}} + q^{\frac{1}{2}})^2), \quad (5.252)$$

$$R_{-1}(z; \lambda) = (q^{-\frac{1}{2}} - q^{\frac{1}{2}})^2 (-z'^2 + z' - a(1 + q)), \quad (5.253)$$

$$B(x; \lambda + \delta) = q^2 B(x + 1; \lambda), \quad D(x; \lambda + \delta) = D(x; \lambda), \quad (5.254)$$

$$\varphi(x; \lambda) = q^x, \quad f_n(\lambda) = \mathcal{E}(n; \lambda), \quad b_n(\lambda) = 1. \quad (5.255)$$

The dual alternative q -Charlier polynomial with the same parameter and

$$B(x; \lambda) = aq^{3x+1} \frac{1 + aq^x}{(1 + aq^{2x})(1 + aq^{2x+1})}, \quad D(x; \lambda) = \frac{1 - q^x}{(1 + aq^{2x-1})(1 + aq^{2x})}, \quad (5.256)$$

$$\mathcal{E}(n; \lambda) = 1 - q^n, \quad \eta(x; \lambda) = (q^{-x} - 1)(1 + aq^x), \quad (5.257)$$

$$A_n(\lambda) = -a, \quad C_n(\lambda) = -(q^{-n} - 1), \quad (5.258)$$

has not been reported in Koekoek-Swarttouw [14]. The dual alternative q -Charlier polynomial was introduced by Atakishiyev and Klimyk [25], which has the same three term recurrence relation and the orthogonality measure as above. The difference equation was not mentioned there.

5.2.8 q -Charlier [KS3.23]

$$q^\lambda = a, \quad \delta = -1, \quad \kappa = q, \quad a > 0, \quad (5.259)$$

$$B(x; \lambda) = aq^x, \quad D(x; \lambda) = 1 - q^x, \quad (5.260)$$

$$\mathcal{E}(n; \lambda) = 1 - q^n, \quad \eta(x; \lambda) = q^{-x} - 1, \quad (5.261)$$

$$P_n(\eta(x; \boldsymbol{\lambda}); \boldsymbol{\lambda}) = {}_2\phi_1\left(\begin{matrix} q^{-n}, q^{-x} \\ 0 \end{matrix} \middle| q; -a^{-1}q^{n+1}\right) = C_n(1 + \eta(x; \boldsymbol{\lambda}); a; q), \quad (5.262)$$

$$\phi_0(x; \boldsymbol{\lambda})^2 = \frac{a^x q^{\frac{1}{2}x(x-1)}}{(q; q)_x}, \quad d_n(\boldsymbol{\lambda})^2 = \frac{q^n}{(-a^{-1}q, q; q)_n} \times \frac{1}{(-a; q)_\infty}, \quad (5.263)$$

$$A_n(\boldsymbol{\lambda}) = -aq^{-2n-1}, \quad C_n(\boldsymbol{\lambda}) = -(q^{-n} - 1)(1 + aq^{-n}), \quad (5.264)$$

$$R_1(z; \boldsymbol{\lambda}) = (q^{-\frac{1}{2}} - q^{\frac{1}{2}})^2 z', \quad z' \stackrel{\text{def}}{=} z - 1, \quad (5.265)$$

$$R_0(z; \boldsymbol{\lambda}) = (q^{-\frac{1}{2}} - q^{\frac{1}{2}})^2 z'^2, \quad (5.266)$$

$$R_{-1}(z; \boldsymbol{\lambda}) = (q^{-\frac{1}{2}} - q^{\frac{1}{2}})^2 (z'^2 + (1 - a)z' - a(1 + q^{-1})), \quad (5.267)$$

$$B(x; \boldsymbol{\lambda} + \boldsymbol{\delta}) = q^{-2}B(x + 1; \boldsymbol{\lambda}), \quad D(x; \boldsymbol{\lambda} + \boldsymbol{\delta}) = D(x; \boldsymbol{\lambda}), \quad (5.268)$$

$$\varphi(x; \boldsymbol{\lambda}) = q^{-x}, \quad f_n(\boldsymbol{\lambda}) = \mathcal{E}(n; \boldsymbol{\lambda}), \quad b_n(\boldsymbol{\lambda}) = 1. \quad (5.269)$$

The dual q -Charlier polynomial with the same parameter and

$$B(x; \boldsymbol{\lambda}) = aq^{-2x-1}, \quad D(x; \boldsymbol{\lambda}) = (q^{-x} - 1)(1 + aq^{-x}), \quad (5.270)$$

$$\mathcal{E}(n; \boldsymbol{\lambda}) = q^{-n} - 1, \quad \eta(x; \boldsymbol{\lambda}) = 1 - q^x, \quad (5.271)$$

$$A_n(\boldsymbol{\lambda}) = -aq^n, \quad C_n(\boldsymbol{\lambda}) = -(1 - q^n), \quad (5.272)$$

has not been reported in Koekoek-Swarttouw [14]. This can be considered another q -version of the Charlier polynomial §5.2.2 with sinusoidal coordinate $\eta(x) = 1 - q^x$. It should be emphasised that these two different q -versions of the Charlier polynomial have very different characters. The q -Charlier has a bounded spectrum and its dual has an unbounded spectrum. This is always the case for the infinite dimensional q -polynomials and their duals.

5.3 Other Polynomials

Here we will discuss orthogonal polynomials which are not listed in Koekoek-Swarttouw's review [14]. The naming of the polynomials is very tentative.

5.3.1 alternative q -Hahn

The q -Hahn polynomial §5.1.6 is a q -version of the Hahn polynomial §5.1.2 with the sinusoidal coordinate $\eta(x) = q^{-x} - 1$. The other q -version with $\eta(x) = 1 - q^x$ and the same spectrum $\mathcal{E}(n) = \epsilon(q^{-n} - 1)(1 - abq^{n-1})$ as the q -Hahn polynomial is as follows (the parameter range could well be enlarged):

$$q^\lambda = (a, b, q^N), \quad \boldsymbol{\delta} = (1, 1, -1), \quad \kappa = q^{-1}, \quad (5.273)$$

$$\epsilon = \begin{cases} 1 & \text{for } 0 < a < 1, \quad 0 < b < 1, \\ -1 & \text{for } a > q^{1-N}, \quad b > q^{1-N}, \end{cases} \quad (5.274)$$

$$B(x; \boldsymbol{\lambda}) = \epsilon a q^{-1} (1 - q^{N-x})(q^{-x} - b), \quad D(x; \boldsymbol{\lambda}) = \epsilon (1 - a q^{N-x})(q^{-x} - 1), \quad (5.275)$$

$$\mathcal{E}(n; \boldsymbol{\lambda}) = \epsilon (q^{-n} - 1)(1 - a b q^{n-1}), \quad \eta(x; \boldsymbol{\lambda}) = 1 - q^x, \quad (5.276)$$

$$P_n(\eta(x; \boldsymbol{\lambda}); \boldsymbol{\lambda}) = {}_3\phi_2 \left(\begin{matrix} q^{-n}, a b q^{n-1}, q^{-x} \\ b, q^{-N} \end{matrix} \middle| q; a^{-1} q^{x+1-N} \right), \quad (5.277)$$

$$\phi_0(x; \boldsymbol{\lambda})^2 = \frac{(q; q)_N}{(q; q)_x (q; q)_{N-x}} \frac{a^x (a; q)_{N-x} (b; q)_x}{(a; q)_N}, \quad (5.278)$$

$$d_n(\boldsymbol{\lambda})^2 = \frac{(q; q)_N}{(q; q)_n (q; q)_{N-n}} \frac{(b, a b q^{-1}; q)_n a^n q^{n(n-1)}}{(a, a b q^N; q)_n} \frac{1 - a b q^{2n-1}}{1 - a b q^{-1}} \times \frac{(a; q)_N}{(a b; q)_N}, \quad (5.279)$$

$$A_n(\boldsymbol{\lambda}) = -a q^{n+N} \frac{(q^{n-N} - 1)(1 - b q^n)(1 - a b q^{n-1})}{(1 - a b q^{2n-1})(1 - a b q^{2n})}, \quad (5.280)$$

$$C_n(\boldsymbol{\lambda}) = -\frac{(1 - q^n)(1 - a q^{n-1})(1 - a b q^{n+N-1})}{(1 - a b q^{2n-2})(1 - a b q^{2n-1})}, \quad (5.281)$$

$$R_1(z; \boldsymbol{\lambda}) = (q^{-\frac{1}{2}} - q^{\frac{1}{2}})^2 z', \quad z' \stackrel{\text{def}}{=} z + \epsilon (1 + a b q^{-1}), \quad (5.282)$$

$$R_0(z; \boldsymbol{\lambda}) = (q^{-\frac{1}{2}} - q^{\frac{1}{2}})^2 (z'^2 - a b (1 + q^{-1})^2), \quad (5.283)$$

$$R_{-1}(z; \boldsymbol{\lambda}) = (q^{-\frac{1}{2}} - q^{\frac{1}{2}})^2 \left(-z'^2 + (1 + a q^{-1} + a(1 + b q^{-1}) q^N) \epsilon z' - a(1 + q^{-1})(1 - b + b(1 + a q^{-1}) q^N) \right), \quad (5.284)$$

$$B(x; \boldsymbol{\lambda} + \boldsymbol{\delta}) = q^2 B(x + 1; \boldsymbol{\lambda}), \quad D(x; \boldsymbol{\lambda} + \boldsymbol{\delta}) = D(x; \boldsymbol{\lambda}), \quad (5.285)$$

$$\varphi(x; \boldsymbol{\lambda}) = q^x, \quad f_n(\boldsymbol{\lambda}) = \mathcal{E}(n; \boldsymbol{\lambda}), \quad b_n(\boldsymbol{\lambda}) = 1. \quad (5.286)$$

This is the most generic form of the functions $B(x)$ and $D(x)$ for the sinusoidal coordinate $\eta(x) = 1 - q^x$ (quadratic polynomials in q^{-x}) as mentioned (A.8) in Appendix A. It is interesting to note that the functions $B(x)$ and $D(x)$ are related to those of the q -Hahn polynomial with the change of variable $x \leftrightarrow N - x$, $B \leftrightarrow D$:

$$B(x) = D^{q\text{-Hahn}}(N - x), \quad D(x) = B^{q\text{-Hahn}}(N - x). \quad (5.287)$$

5.3.2 alternative q -Krawtchouk

The q -Krawtchouk polynomial §5.1.9 is a q -version of the Krawtchouk polynomial §5.1.4 with the sinusoidal coordinate $\eta(x) = q^{-x} - 1$. The other q -version with $\eta(x) = 1 - q^x$ and the same spectrum $\mathcal{E}(n) = (q^{-n} - 1)(1 + p q^n)$ is as follows:

$$q^\lambda = (p, q^N), \quad \boldsymbol{\delta} = (2, -1), \quad \kappa = q^{-1}, \quad p > 0, \quad (5.288)$$

$$B(x; \boldsymbol{\lambda}) = p(1 - q^{N-x}), \quad D(x; \boldsymbol{\lambda}) = q^{-x} - 1, \quad (5.289)$$

$$\mathcal{E}(n; \boldsymbol{\lambda}) = (q^{-n} - 1)(1 + pq^n), \quad \eta(x; \boldsymbol{\lambda}) = 1 - q^x, \quad (5.290)$$

$$P_n(\eta(x; \boldsymbol{\lambda}); \boldsymbol{\lambda}) = {}_3\phi_1\left(\begin{matrix} q^{-n}, -pq^n, q^{-x} \\ q^{-N} \end{matrix} \middle| q; -p^{-1}q^{x-N}\right), \quad (5.291)$$

$$\phi_0(x; \boldsymbol{\lambda})^2 = \frac{(q; q)_N}{(q; q)_x (q; q)_{N-x}} p^x q^{\frac{1}{2}x(x+1)}, \quad (5.292)$$

$$d_n(\boldsymbol{\lambda})^2 = \frac{(q; q)_N}{(q; q)_n (q; q)_{N-n}} \frac{(-p; q)_n p^n q^{\frac{1}{2}n(3n-1)}}{(-pq^{N+1}; q)_n} \frac{1 + pq^{2n}}{1 + p} \times \frac{1}{(-pq; q)_N}, \quad (5.293)$$

$$A_n(\boldsymbol{\lambda}) = -pq^{2n+N+1} \frac{(q^{n-N} - 1)(1 + pq^n)}{(1 + pq^{2n})(1 + pq^{2n+1})}, \quad C_n(\boldsymbol{\lambda}) = -\frac{(1 - q^n)(1 + pq^{n+N})}{(1 + pq^{2n-1})(1 + pq^{2n})}, \quad (5.294)$$

$$R_1(z; \boldsymbol{\lambda}) = (q^{-\frac{1}{2}} - q^{\frac{1}{2}})^2 z', \quad z' \stackrel{\text{def}}{=} z + 1 - p, \quad (5.295)$$

$$R_0(z; \boldsymbol{\lambda}) = (q^{-\frac{1}{2}} - q^{\frac{1}{2}})^2 (z'^2 + p(q^{-\frac{1}{2}} + q^{\frac{1}{2}})^2), \quad (5.296)$$

$$R_{-1}(z; \boldsymbol{\lambda}) = (q^{-\frac{1}{2}} - q^{\frac{1}{2}})^2 (-z'^2 + (1 - pq^N)z' - p(1 + q)(1 - q^N)), \quad (5.297)$$

$$B(x; \boldsymbol{\lambda} + \boldsymbol{\delta}) = q^2 B(x + 1; \boldsymbol{\lambda}), \quad D(x; \boldsymbol{\lambda} + \boldsymbol{\delta}) = D(x; \boldsymbol{\lambda}), \quad (5.298)$$

$$\varphi(x; \boldsymbol{\lambda}) = q^x, \quad f_n(\boldsymbol{\lambda}) = \mathcal{E}(n; \boldsymbol{\lambda}), \quad b_n(\boldsymbol{\lambda}) = 1. \quad (5.299)$$

The functions $B(x)$ and $D(x)$ are related to those of the q -Krawtchouk polynomial with the change of variable $x \leftrightarrow N - x$, $B \leftrightarrow D$:

$$B(x) = D^{q\text{-Krawtchouk}}(N - x), \quad D(x) = B^{q\text{-Krawtchouk}}(N - x). \quad (5.300)$$

5.3.3 alternative affine q -Krawtchouk (self-dual)

The affine q -Krawtchouk §5.1.10 is a self-dual polynomial with $\eta(x) = q^{-x} - 1$ and $\mathcal{E}(n) = q^{-n} - 1$. There is another self-dual polynomial with $\eta(x) = 1 - q^x$ and $\mathcal{E}(n) = 1 - q^n$:

$$q^\lambda = (p, q^N), \quad \boldsymbol{\delta} = (1, -1), \quad \kappa = q, \quad p > q^{-N}, \quad (5.301)$$

$$B(x; \boldsymbol{\lambda}) = (1 - q^{N-x})(1 - p^{-1}q^{-x-1}), \quad D(x; \boldsymbol{\lambda}) = p^{-1}q^{N-x}(q^{-x} - 1), \quad (5.302)$$

$$\mathcal{E}(n; \boldsymbol{\lambda}) = 1 - q^n, \quad \eta(x; \boldsymbol{\lambda}) = 1 - q^x, \quad (5.303)$$

$$P_n(\eta(x; \boldsymbol{\lambda}); \boldsymbol{\lambda}) = {}_2\phi_2\left(\begin{matrix} q^{-n}, q^{-x} \\ pq, q^{-N} \end{matrix} \middle| q; pq^{x+n+1-N}\right), \quad (5.304)$$

$$\phi_0(x; \boldsymbol{\lambda})^2 = \frac{(q; q)_N}{(q; q)_x (q; q)_{N-x}} \frac{(p^{-1}q^{-N}; q)_N}{(p^{-1}q^{-N}; q)_{N-x}} p^x q^{x(x+1-N)}, \quad (5.305)$$

$$d_n(\boldsymbol{\lambda})^2 = \frac{(q; q)_N}{(q; q)_n (q; q)_{N-n}} \frac{(p^{-1}q^{-N}; q)_N}{(p^{-1}q^{-N}; q)_{N-n}} p^n q^{n(n+1-N)} \times (pq)^{-N}, \quad (5.306)$$

$$A_n(\boldsymbol{\lambda}) = -(1 - p^{-1}q^{-n-1})(1 - q^{N-n}), \quad C_n(\boldsymbol{\lambda}) = -p^{-1}q^{N-n}(q^{-n} - 1), \quad (5.307)$$

$$R_1(z; \boldsymbol{\lambda}) = (q^{-\frac{1}{2}} - q^{\frac{1}{2}})^2 z', \quad z' \stackrel{\text{def}}{=} z - 1, \quad (5.308)$$

$$R_0(z; \boldsymbol{\lambda}) = (q^{-\frac{1}{2}} - q^{\frac{1}{2}})^2 z'^2, \quad (5.309)$$

$$R_{-1}(z; \boldsymbol{\lambda}) = (q^{-\frac{1}{2}} - q^{\frac{1}{2}})^2 \left(-z'^2 - (p^{-1}q^{-1} + (1 + p^{-1})q^N)z' - p^{-1}(1 + q^{-1})q^N \right), \quad (5.310)$$

$$B(x; \boldsymbol{\lambda} + \boldsymbol{\delta}) = B(x + 1; \boldsymbol{\lambda}), \quad D(x; \boldsymbol{\lambda} + \boldsymbol{\delta}) = q^{-2}D(x; \boldsymbol{\lambda}), \quad (5.311)$$

$$\varphi(x; \boldsymbol{\lambda}) = q^x, \quad f_n(\boldsymbol{\lambda}) = \mathcal{E}(n; \boldsymbol{\lambda}), \quad b_n(\boldsymbol{\lambda}) = 1. \quad (5.312)$$

The functions $B(x)$ and $D(x)$ are related to those of the quantum q -Krawtchouk polynomial with the change of variable $x \leftrightarrow N - x$, $B \leftrightarrow D$:

$$B(x) = D^{\text{quantum } q\text{-Krawtchouk}}(N - x), \quad D(x) = B^{\text{quantum } q\text{-Krawtchouk}}(N - x). \quad (5.313)$$

6 Summary and Comments

A unified theory of orthogonal polynomials of a discrete variable is presented through the eigenvalue problem of hermitian matrices of finite or infinite dimensions. The hermitian matrices (\mathcal{H}) are real and tridiagonal (Jacobi) matrices (2.4) that can be factorised $\mathcal{H} = \mathcal{A}^\dagger \mathcal{A}$ and thus are positive semi-definite. The orthogonality measure $\phi_0(x)^2$ is obtained as a solution of the factored equation $\mathcal{A}\phi_0(x) = 0$, as in the ordinary quantum mechanics. Then the eigenvalue problem can be solved in two different ways; the first as a difference equation and the second through the three term recurrence relations. The former gives the eigenvector polynomial $P_n(\eta(x))$ and the latter provides the dual polynomial $Q_x(\mathcal{E}(n))$ satisfying the relation $P_n(\eta(x)) = Q_x(\mathcal{E}(n))$ through the universal normalisation condition $P_n(0) = 1 = Q_x(0)$. Here $\mathcal{E}(n)$ is the eigenvalue of \mathcal{H} and $\eta(x)$ is the sinusoidal coordinate satisfying the closure relation $[\mathcal{H}, [\mathcal{H}, \eta(x)]] = \eta(x)R_0(\mathcal{H}) + [\mathcal{H}, \eta(x)]R_1(\mathcal{H}) + R_{-1}(\mathcal{H})$. Thanks to the closure relation and the shape invariance, the entire spectrum, the coefficients A_n and C_n of the three term recurrence relation of $P_n(\eta)$, its normalisation constants d_n^2 are determined algebraically. As byproducts, some as yet unexplored (not reported in the review of Koekoek-Swartzouw [14]) dual polynomials with explicit forms of the difference equation, three term recurrence, the normalisation measure/constants are presented in §5.3. On top of them we mentioned the dual quantum q -Krawtchouk (5.124)–(5.126), the dual q -Meixner (5.217)–(5.219) and the dual q -Charlier (5.270)–(5.272).

After completing the main part of the paper, we became aware of the work of Terwilliger [26] on the correspondence between the Leonard pair [27] and a class of orthogonal polynomials. Although the general setting of the problem and the methods are markedly different,

Terwilliger's results have some overlap with the finite dimensional case (I) of the present paper.

Acknowledgements

This work is supported in part by Grants-in-Aid for Scientific Research from the Ministry of Education, Culture, Sports, Science and Technology, No.18340061 and No.19540179. This work was also supported in part by the Italian MIUR (Internazionalizzazione Program) within the joint SISSA–YITP research project on “Fundamental Interactions and the Early Universe.”

Appendix A: Possible Forms of the Hamiltonians

Here we continue from §4.3 the algebraic analysis of the closure relation to determine the possible function forms of $B(x)$ and $D(x)$. The remaining equation to be solved is (4.66), which is the diagonal part of the closure relation (4.32).

With (4.69) and (4.70), the equation to be analysed (4.66) reads

$$\begin{aligned} & (\eta(x) - \eta(x+2))B(x)D(x+1) + (\eta(x) - \eta(x-2))B(x-1)D(x) \\ &= (\eta(x+1) - 2\eta(x) + \eta(x-1))a_x^2 + (2r_1^{(0)}\eta(x) + r_{-1}^{(1)})a_x + r_0^{(0)}\eta(x) + r_{-1}^{(0)}. \end{aligned} \quad (\text{A.1})$$

Before proceeding, let us stress again that the input of the eigenvalue problem (2.1) is $B(x)$ and $D(x)$. The quantities $r_i^{(j)}$ ($i = 1, 0, -1$; $j = 0, 1, 2$) are determined by the consistency of the closure relation (4.31). In Appendix A, however, we turn the logic and regard $r_i^{(j)}$ as input data and determine the possible forms of $B(x)$ and $D(x)$ which satisfy (4.62)–(4.66). In this approach the positivity of $B(x)$ and $D(x)$ must be verified at the last step, since it is not guaranteed automatically. By multiplying $\eta(x-1) - \eta(x+1)$ to (A.1), we obtain

$$\begin{aligned} & (\eta(x+2) - \eta(x))(\eta(x+1) - \eta(x-1))B(x)D(x+1) - (x \rightarrow x-1) \\ &= \left((\eta(x-1) - \eta(x))^2 - (\eta(x+1) - \eta(x))^2 \right) a_x^2 \\ &+ (\eta(x-1) - \eta(x+1))(2r_1^{(0)}\eta(x) + r_{-1}^{(1)})a_x + (\eta(x-1) - \eta(x+1))(r_0^{(0)}\eta(x) + r_{-1}^{(0)}) \\ &= \frac{(r_1^{(0)}\eta(x+1)\eta(x) + r_{-1}^{(1)}\eta(x+1) + a'_0)(r_1^{(0)}\eta(x+1)\eta(x) + r_{-1}^{(1)}\eta(x) + a'_0)}{(\eta(x+1) - \eta(x))^2} \\ &- r_0^{(0)}\eta(x+1)\eta(x) - r_{-1}^{(0)}(\eta(x+1) + \eta(x)) - (x \rightarrow x-1), \end{aligned} \quad (\text{A.2})$$

where $a'_0 = \eta(1)\eta(-1)a_0$ and $a_0 = B(0)$. For the second equality, the main result of §4.3, (4.92) is used. The above equation is simplified as

$$\begin{aligned} & (\eta(x+2) - \eta(x))(\eta(x+1) - \eta(x-1))B(x)D(x+1) \\ &= \frac{(r_1^{(0)}\eta(x+1)\eta(x) + r_{-1}^{(1)}\eta(x+1) + a'_0)(r_1^{(0)}\eta(x+1)\eta(x) + r_{-1}^{(1)}\eta(x) + a'_0)}{(\eta(x+1) - \eta(x))^2} \\ & \quad - \eta(1)a_0(r_{-1}^{(1)} + \eta(1)a_0) - r_0^{(0)}\eta(x+1)\eta(x) - r_{-1}^{(0)}(\eta(x+1) + \eta(x) - \eta(-1)). \end{aligned} \quad (\text{A.3})$$

For each case of $\eta(x)$, (4.72)–(4.76), we will solve $B(x)$ from the above equation (A.3), with the result of $a_x = B(x) + D(x)$, (4.92). Let us parametrise $B(x)$ as

$$B(x) = \frac{\tilde{B}(x)}{(\eta(x+1) - \eta(x))(\eta(x+1) - \eta(x-1))}. \quad (\text{A.4})$$

Note that the denominator is the dual counterpart of $\alpha_+(\mathcal{E}(n))(\alpha_+(\mathcal{E}(n)) - \alpha_-(\mathcal{E}(n)))$ in the denominator of A_n in (4.52).

We eliminate $D(x+1)$ from (A.3) (multiplied by $(\eta(x+2) - \eta(x+1))(\eta(x+1) - \eta(x))^2$) by using (4.92):

$$\begin{aligned} & -\tilde{B}(x) \left((r_1^{(0)}\eta(x+1)^2 + r_{-1}^{(1)}\eta(x+1) + a'_0)(\eta(x+2) - \eta(x)) + \tilde{B}(x+1)(\eta(x+1) - \eta(x)) \right) \\ &= (\eta(x+2) - \eta(x+1)) \\ & \quad \times \left((r_1^{(0)}\eta(x+1)\eta(x) + r_{-1}^{(1)}\eta(x+1) + a'_0)(r_1^{(0)}\eta(x+1)\eta(x) + r_{-1}^{(1)}\eta(x) + a'_0) \right. \\ & \quad \left. - (\eta(x+1) - \eta(x))^2 \left(\eta(1)a_0(r_{-1}^{(1)} + \eta(1)a_0) + r_0^{(0)}\eta(x+1)\eta(x) \right. \right. \\ & \quad \left. \left. + r_{-1}^{(0)}(\eta(x+1) + \eta(x) - \eta(-1)) \right) \right). \end{aligned} \quad (\text{A.5})$$

This equation determines $\tilde{B}(x)$ uniquely from the boundary (initial) value $\tilde{B}(0) = \eta(1)(\eta(1) - \eta(-1))B(0)$.

We have carried out this last step by assuming that $\tilde{B}(x)$ is a polynomial in x , $\tilde{B}(x) = \sum_{j=0}^M \beta_j x^j$ for the cases (i) and (ii) and that $\tilde{B}(x)$ is a Laurent polynomial in q^x , $\tilde{B}(x) = \sum_{j=M_1}^{M_2} \beta_j q^{jx}$ for the cases (iii) – (v). The results are as follows.

case (i), $\eta(x) = x$:

A quadratic polynomial ($M = 2$) solution $\tilde{B}(x)$ is obtained and $B(x)$ and $D(x)$ are:

$$B(x) = (\text{quadratic in } x), \quad D(x) = x(\text{linear in } x), \quad (\text{A.6})$$

which contains the Hahn polynomial §5.1.2 as the most generic case.

case (ii), $\eta(x) = \epsilon'x(x+d)$:

A quartic polynomial ($M = 4$) solution $\tilde{B}(x)$ is obtained and $B(x)$ and $D(x)$ are:

$$B(x) = \frac{(x+d)(\text{cubic in } x)}{(2x+d)(2x+1+d)}, \quad D(x) = \frac{x(\text{cubic in } x)}{(2x+d)(2x-1+d)}, \quad (\text{A.7})$$

which contains the Racah polynomial §5.1.1 as the most generic case.

case (iii), $\eta(x) = 1 - q^x$:

A quadratic polynomial ($M_1 = 0, M_2 = 2$) solution $\tilde{B}(x)$ in q^x is obtained and $B(x)$ and $D(x)$ are:

$$B(x) = (\text{quadratic in } q^{-x}), \quad D(x) = (q^{-x} - 1)(\text{linear in } q^{-x}). \quad (\text{A.8})$$

The most generic one in this category reported in the Koekoek-Swarttouw's review [14] is the little q -Jacobi polynomial §5.2.3. The dual quantum q -Krawtchouk polynomial (5.124)–(5.126) also belongs to this case. There is a more general polynomial with one more parameter as tentatively called alternative q -Hahn polynomial §5.3.1. For the infinite dimensional case (II), the functions $B(x)$ and $D(x)$ grow exponentially. Therefore the corresponding spectrum is unbounded, *i.e.* type (iv) (4.10) or (v) (4.11). In other words, there is no self-dual $\eta(x) = 1 - q^x$, $\mathcal{E}(n) = 1 - q^n$ polynomial in the infinite dimensional case (II).

case (iv), $\eta(x) = q^{-x} - 1$:

A quadratic polynomial ($M_1 = -2, M_2 = 0$) solution $\tilde{B}(x)$ in q^{-x} is obtained and $B(x)$ and $D(x)$ are:

$$B(x) = (\text{quadratic in } q^x), \quad D(x) = (1 - q^x)(\text{linear in } q^x), \quad (\text{A.9})$$

which contains the q -Hahn polynomial §5.1.6 as the most generic one. In contrast to the previous case the functions $B(x)$ and $D(x)$ are all bounded in the the infinite dimensional case (II). Therefore the corresponding spectrum is also bounded, *i.e.* type (iii) (4.9) only. In other words, there is no self-dual $\eta(x) = q^{-x} - 1$, $\mathcal{E}(n) = q^{-n} - 1$ polynomial in the infinite dimensional case (II).

case (v), $\eta(x) = \epsilon'(q^{-x} - 1)(1 - dq^x)$:

A quadratic polynomial ($M_1 = -2, M_2 = 2$) solution $\tilde{B}(x)$ in q^x and q^{-x} is obtained and $B(x)$ and $D(x)$ are:

$$B(x) = \frac{(1 - dq^x)(\text{cubic in } q^x)}{(1 - dq^{2x})(1 - dq^{2x+1})}, \quad D(x) = \frac{(1 - q^x)(\text{cubic in } q^x)}{(1 - dq^{2x})(1 - dq^{2x-1})}, \quad (\text{A.10})$$

which contains the most generic case of the q -Racah polynomial §5.1.5. The functions $B(x)$ and $D(x)$ are all bounded in the the infinite dimensional case (II). The corresponding spectrum is also bounded, *i.e.* type (iii) (4.9) only. The case (v) does not have self-dual polynomial, either. Thus the most generic functional forms of $B(x)$ and $D(x)$ are determined for each case of the five different sinusoidal coordinates (i)–(v).

Determination of polynomials satisfying certain forms of difference equations for given $\eta(x)$ (quadratic or q -quadratic) has a long history [23, 24, 27, 26]. Bochner's theorem [22] on Sturm-Liouville polynomials is a precursor. The present characterisation in terms of the closure relation (4.32) is consistent with the existing ones.

It is interesting to note that in all the known cases the following relation is satisfied:

$$\frac{r_0^{(0)}}{B(0)^2} + \frac{r_1^{(0)}}{B(0)} \frac{r_{-1}^{(0)}}{\eta(1)B(0)^2} - \left(\frac{r_{-1}^{(0)}}{\eta(1)B(0)^2} \right)^2 = 0. \quad (\text{A.11})$$

When this condition is met, the solution of (A.5) for all the cases (i)–(v) has a simple expression in terms of $\eta(x)$:

$$\begin{aligned} \tilde{B}(x) &= -r_1^{(0)}\eta(x)\eta(x+1) + \frac{r_{-1}^{(0)}}{\eta(1)B(0)}\eta(x)(\eta(x+1) - \eta(x)) \\ &\quad - r_{-1}^{(1)}\eta(x) + \eta(1)B(0)(\eta(x+1) - \eta(x) - \eta(-1)). \end{aligned} \quad (\text{A.12})$$

For the other function $D(x)$ we parametrise

$$D(x) = \frac{\tilde{D}(x)}{(\eta(x-1) - \eta(x))(\eta(x-1) - \eta(x+1))}. \quad (\text{A.13})$$

The corresponding expression for $\tilde{D}(x)$ is

$$\begin{aligned} \tilde{D}(x) &= -r_1^{(0)}\eta(x)\eta(x-1) + \frac{r_{-1}^{(0)}}{\eta(1)B(0)}\eta(x)(\eta(x-1) - \eta(x)) \\ &\quad - r_{-1}^{(1)}\eta(x) + \eta(1)B(0)(\eta(x-1) - \eta(x) - \eta(-1)). \end{aligned} \quad (\text{A.14})$$

The above relation (A.11) is a consequence of (4.54), which can be regarded as a consistency condition. From (4.35), (4.42) and (4.51) we have $\mathcal{E}(1) = \alpha_+(\mathcal{E}(0)) = \frac{1}{2}(r_1^{(0)} + \sqrt{r_1^{(0)2} + 4r_0^{(0)}})$ and $A_0 = r_{-1}^{(0)}/r_0^{(0)}$. Substituting these into (4.54) (multiplied by $r_0^{(0)}/\mathcal{E}(1)$), we obtain

$$r_{-1}^{(0)} + \frac{1}{2} \left(\sqrt{r_1^{(0)2} + 4r_0^{(0)}} - r_1^{(0)} \right) \eta(1)B(0) = 0,$$

namely

$$\frac{1}{2}\sqrt{r_1^{(0)2} + 4r_0^{(0)}} = \frac{1}{2}r_1^{(0)} - \frac{r_{-1}^{(0)}}{\eta(1)B(0)}.$$

Squaring this gives (A.11).

By using (4.54), the expressions for $B(x)$ and $D(x)$, (A.4) with (A.12) and (A.13) with (A.14), are cast in a similar form as A_n (4.52) and C_n (4.53):

$$B(x) = -\frac{R_{-1}^{\text{dual}}(\eta(x)) + \mathcal{E}(1)(\eta(x) + A_0)\alpha_+^{\text{dual}}(\eta(x))}{\alpha_+^{\text{dual}}(\eta(x))(\alpha_+^{\text{dual}}(\eta(x)) - \alpha_-^{\text{dual}}(\eta(x)))}, \quad (\text{A.15})$$

$$D(x) = -\frac{R_{-1}^{\text{dual}}(\eta(x)) + \mathcal{E}(1)(\eta(x) + A_0)\alpha_-^{\text{dual}}(\eta(x))}{\alpha_-^{\text{dual}}(\eta(x))(\alpha_-^{\text{dual}}(\eta(x)) - \alpha_+^{\text{dual}}(\eta(x)))}. \quad (\text{A.16})$$

Here $R_i^{\text{dual}}(z)$ and $\alpha_{\pm}^{\text{dual}}(z)$ are functions appearing in the dual closure relation §4.4. The functions $R_i^{\text{dual}}(z)$ are given in (4.104)–(4.106) and $\alpha_{\pm}^{\text{dual}}(z)$ are defined by (c.f. (4.35), (4.42))

$$\alpha_{\pm}^{\text{dual}}(z) \stackrel{\text{def}}{=} \frac{1}{2}\left(R_1^{\text{dual}}(z) \pm \sqrt{R_1^{\text{dual}}(z)^2 + 4R_0^{\text{dual}}(z)}\right), \quad (\text{A.17})$$

which satisfy

$$\alpha_{\pm}^{\text{dual}}(\eta(x)) = \eta(x \pm 1) - \eta(x). \quad (\text{A.18})$$

Similarly A_n and C_n can be cast into the forms like $B(x)$ in (A.4) with (A.12) and $D(x)$ in (A.13) with (A.14).

Appendix B: Some definitions related to the hypergeometric and q -hypergeometric functions

For self-containedness we collect several definitions related to the (q -)hypergeometric functions [14].

◦ Pochhammer symbol $(a)_n$:

$$(a)_n \stackrel{\text{def}}{=} \prod_{k=1}^n (a+k-1) = a(a+1)\cdots(a+n-1) = \frac{\Gamma(a+n)}{\Gamma(a)}. \quad (\text{B.1})$$

◦ q -Pochhammer symbol $(a; q)_n$:

$$(a; q)_n \stackrel{\text{def}}{=} \prod_{k=1}^n (1 - aq^{k-1}) = (1-a)(1-aq)\cdots(1-aq^{n-1}). \quad (\text{B.2})$$

◦ hypergeometric series ${}_rF_s$:

$${}_rF_s \left(\begin{matrix} a_1, \dots, a_r \\ b_1, \dots, b_s \end{matrix} \middle| z \right) \stackrel{\text{def}}{=} \sum_{n=0}^{\infty} \frac{(a_1, \dots, a_r)_n}{(b_1, \dots, b_s)_n} \frac{z^n}{n!}, \quad (\text{B.3})$$

where $(a_1, \dots, a_r)_n \stackrel{\text{def}}{=} \prod_{j=1}^r (a_j)_n = (a_1)_n \cdots (a_r)_n$.

○ q -hypergeometric series (the basic hypergeometric series) ${}_r\phi_s$:

$${}_r\phi_s \left(\begin{matrix} a_1, \dots, a_r \\ b_1, \dots, b_s \end{matrix} \middle| q; z \right) \stackrel{\text{def}}{=} \sum_{n=0}^{\infty} \frac{(a_1, \dots, a_r; q)_n}{(b_1, \dots, b_s; q)_n} (-1)^{(1+s-r)n} q^{(1+s-r)n(n-1)/2} \frac{z^n}{(q; q)_n}, \quad (\text{B.4})$$

where $(a_1, \dots, a_r; q)_n \stackrel{\text{def}}{=} \prod_{j=1}^r (a_j; q)_n = (a_1; q)_n \cdots (a_r; q)_n$.

References

- [1] G. Szegő, *Orthogonal polynomials*, 4th edition, Amer. Math. Soc. Colloq. Pub. **23** Providence, R.I. (1975).
- [2] A. F. Nikiforov, S. K. Suslov, and V. B. Uvarov, *Classical Orthogonal Polynomials of a Discrete Variable*, Springer, Berlin, (1991).
- [3] G. E. Andrews, R. Askey and R. Roy, *Special Functions*, Encyclopedia of mathematics and its applications, Cambridge, (1999).
- [4] G. Gasper and M. Rahman, *Basic hypergeometric series*, 2nd ed. Encyclopedia of mathematics and its applications, Cambridge, (2004).
- [5] M. E. H. Ismail, *Classical and quantum orthogonal polynomials in one variable*, Encyclopedia of mathematics and its applications, Cambridge, (2005).
- [6] S. Flügge, *Practical Quantum Mechanics*, Springer, Berlin, (1994).
- [7] M. M. Crum, “Associated Sturm-Liouville systems,” Quart. J. Math. Oxford Ser. (2) **6** (1955) 121-127, [arXiv:physics/9908019](#).
- [8] L. Infeld and T. E. Hull, “The factorization method,” Rev. Mod. Phys. **23** (1951) 21-68.
- [9] See, for example, a review: F. Cooper, A. Khare and U. Sukhatme, “Supersymmetry and quantum mechanics,” Phys. Rep. **251** (1995) 267-385.
- [10] S. Odake and R. Sasaki, “Unified Theory of Annihilation-Creation Operators for Solvable (‘Discrete’) Quantum Mechanics,” J. Math. Phys. **47** (2006) 102102 (33pages), [arXiv:quant-ph/0605215](#); “Exact solution in the Heisenberg picture and annihilation-creation operators,” Phys. Lett. **B641** (2006) 112-117, [arXiv:quant-ph/0605221](#).

- [11] S. Odake and R. Sasaki, “Exact Heisenberg operator solutions for multi-particle quantum mechanics,” *J. Math. Phys.* **48** (2007) 082106, [arXiv:0706.0768 \[quant-ph\]](#).
- [12] L. E. Gendenshtein, “Derivation of exact spectra of the Schrodinger equation by means of supersymmetry,” *JETP Lett.* **38** (1983) 356-359.
- [13] S. Odake and R. Sasaki, “Shape Invariant Potentials in “Discrete Quantum Mechanics,” *J. Nonlinear Math. Phys.* **12** Suppl. 1 (2005) 507-521, [arXiv:hep-th/0410102](#); “Equilibrium Positions, Shape Invariance and Askey-Wilson Polynomials,” *J. Math. Phys.* **46** (2005) 063513 (10pages), [arXiv:hep-th/0410109](#); “Calogero-Sutherland-Moser Systems, Ruijsenaars-Schneider-van Diejen Systems and Orthogonal Polynomials,” *Prog. Theor. Phys.* **114** (2005) 1245-1260, [arXiv:hep-th/0512155](#).
- [14] R. Koekoek and R.F. Swarttouw, “The Askey-scheme of hypergeometric orthogonal polynomials and its q -analogue,” [arXiv:math.CA/9602214](#).
- [15] A. G. Ushveridze, “Exact solutions of one- and multi-dimensional Schrödinger equations,” *Sov. Phys.-Lebedev Inst. Rep.* **2**, (1988) 50, 54-58; *Quasi-exactly solvable models in quantum mechanics* (IOP, Bristol, 1994); A. Y. Morozov, A. M. Perelomov, A. A. Rosly, M. A. Shifman and A. V. Turbiner, “Quasiexactly solvable quantal problems: one-dimensional analog of rational conformal field theories,” *Int. J. Mod. Phys.* **A5** (1990) 803-832.
- [16] A. V. Turbiner, “Quasi-Exactly-Solvable Problems and $sl(2)$ Algebra,” *Comm. Math. Phys.* **118** (1988) 467-474.
- [17] C. M. Bender and G. V. Dunne, “Quasi-Exactly Solvable Systems and Orthogonal Polynomials,” *J. Math. Phys.* **37** (1996) 6-11.
- [18] R. Sasaki, “Quasi Exactly Solvable Difference Equations,” *J. Math. Phys.* **48** (2007) 122104 (11pages), [arXiv:0708.0702 \[nlin.SI\]](#); “New Quasi Exactly Solvable Difference Equation,” [arXiv:0712.2616 \[nlin.SI\]](#); S. Odake and R. Sasaki, “Multi-Particle Quasi Exactly Solvable Difference Equations,” *J. Math. Phys.* **48** (2007) 122105 (8pages), [arXiv:0708.0716 \[nlin.SI\]](#).
- [19] S. Odake and R. Sasaki, “ q -oscillator from the q -Hermite Polynomial,” [arXiv:0710.2209 \[hep-th\]](#).

- [20] T.H. Koornwinder, “The structure relation for Askey-Wilson polynomials,” [arXiv:math.CA/0601303](#).
- [21] N.M. Atakishiyev, M. Rahman and S.K. Suslov, “On classical orthogonal polynomials,” *Constr. Approx.* **11** (1995) 181–226;
- [22] S. Bochner, “Über Sturm-Liouvillesche Polynomsysteme,” *Math. Zeit.* **29** (1929) 730–736.
- [23] R. Askey and J. Wilson, “Some basic hypergeometric orthogonal polynomials that generalize Jacobi polynomials,” *Mem. Amer. Math. Soc.* **54** No.319 (1985) 1-55; F. Alberto Grünbaum and L. Haine, “The q -version of a theorem of Bochner,” *J. Comput. Appl. Math.* **68** (1996) 103-114; M. E. H. Ismail, “A generalization of a theorem of Bochner,” *J. Comput. Appl. Math.* **159** (2003) 319324.
- [24] L. Vinet and A. Zhedanov, “Generalized Bochner theorem: characterization of the Askey-Wilson polynomials,” [arXiv:0712.0069 \[math.CA\]](#).
- [25] N.M. Atakishiyev and A. U. Klimyk, “On q -orthogonal polynomials, dual to little and big q -Jacobi polynomials,” [arXiv:math.CA/0307250](#); “A set of orthogonal polynomials, dual to alternative q -Charlier polynomials,” [arXiv:math.CA/0312312](#); “Duality of q -polynomials, orthogonal on countable sets of points,” [arXiv:math.CA/0411249](#).
- [26] P. Terwilliger, “Leonard pairs and the q -Racah polynomials,” *Linear Algebra Appl.* **387** (2004) 235–276, [arXiv:math.QA/0306301](#); “Two linear transformations each tridiagonal with respect to an eigenbasis of the other; an algebraic approach to the Askey scheme of orthogonal polynomials,” [arXiv:math.QA/0408390](#).
- [27] D. Leonard, “Orthogonal polynomials, duality, and association schemes,” *SIAM J. Math. Anal.* **13** (1982) 656–663.