

## PHONON-ASSISTED AND PHONON-FREE RADIATIVE TRANSITIONS IN SULFUR-DOPED GaAs DIFFUSED JUNCTIONS

(electroluminescence, quantum  
efficiency, 4.2–77°K, E)

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In the course of electroluminescence studies of GaAs junctions we investigated the emission spectra of sulfur-doped zinc-diffused diodes. The peaks in the spectra correspond to three main transitions and their associated phonon emission satellites.

The diodes were obtained by a diffusion of Zn from a surface concentration of  $5.0 \times 10^{18} \text{ cm}^{-3}$  into a sulfur-doped crystal with an electron concentration of  $2.0 \times 10^{16} \text{ cm}^{-3}$ . The net impurity profile was shown by capacitance measurements to be abrupt with a zero-bias width of  $3340 \text{ \AA}$  at  $300^\circ\text{K}$ . The areas of the junctions were  $\sim 10^{-3} \text{ cm}^2$ .

The diode series resistance decreased with an increase in temperature, from  $12 \text{ } \Omega$  at  $4.2^\circ\text{K}$ ,  $6 \text{ } \Omega$  at  $20^\circ\text{K}$  and to  $1.7 \text{ } \Omega$  at  $77^\circ\text{K}$ , a behavior possibly due to impurity band conduction at low temperatures in the *p*-type side of the diode.

The average external efficiency at  $4.2^\circ\text{K}$ , estimated from the measured external efficiency at  $77^\circ\text{K}$  and the ratio of the integrated spectral intensities at the two temperatures, was 30%. The external efficiencies at  $77^\circ\text{K}$  varied from 0.4 to 0.8% and were measured with a silicon solar cell whose detection efficiency was assumed to be 0.5. Figure 1 shows the variation of external efficiency with current at  $4.2^\circ\text{K}$  for a diode whose ratio of integrated intensities was 80. The maximum value of 30% suggests the plausibility of a quantum efficiency equal to unity in the high current region.

The most intense emission peak shifts with applied voltage. A plot of integrated light intensity vs photon energy of the peak (Fig. 2) at 4.2 and  $20^\circ\text{K}$  displays two regions. In region A,  $h\nu < 1.48 \text{ eV}$ ,

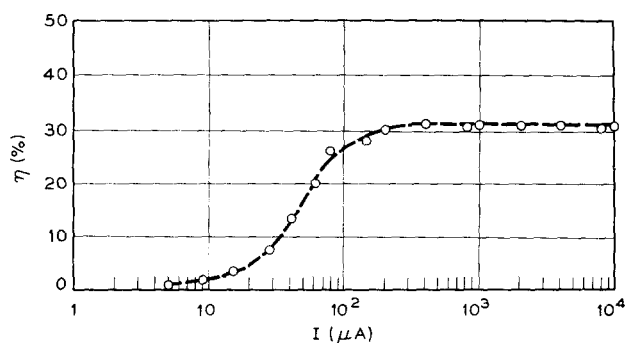


Fig. 1. Variation of external efficiency with current.

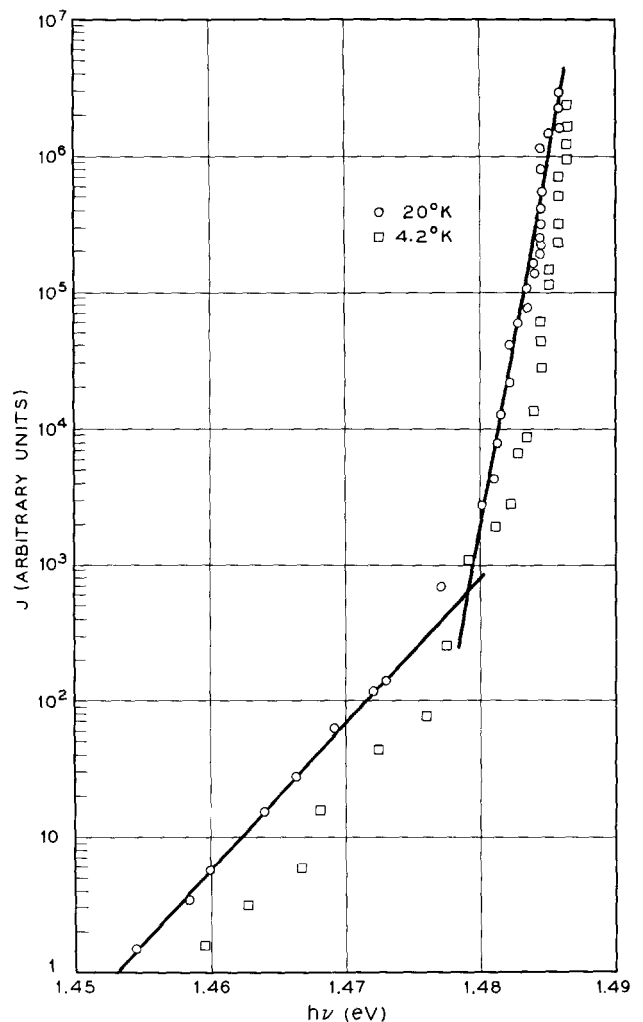


Fig. 2. Plot of the integrated light intensity vs photon energy at 4.2 and  $20^\circ\text{K}$ , for the main emission ( $\nu_1$  in Fig. 3).

the intensity is given by  $J = J_0 \exp[S_A h\nu]$  where  $S_A = 247 \text{ (eV)}^{-1}$ . In region B,  $h\nu > 1.48 \text{ eV}$  the slope changes to  $S_B \approx 1400 \text{ (eV)}^{-1}$ .  $S_B$  remains constant up to  $77^\circ\text{K}$ , suggesting a tunneling mechanism.<sup>3</sup> Because of the decrease of quantum efficiency at  $77^\circ\text{K}$ , it was impossible to obtain data in region A at this temperature.

Figure 3A and 3B shows the electroluminescence spectra at  $4.2^\circ\text{K}$  and  $77^\circ\text{K}$  for a characteristic

current of 25 mA. The lesser peaks shift with applied voltage in the same manner as the most intense peak. The spectrum at 20°K is the same as the one at 4.2°K, except for being displaced to lower energies by 2.5 meV as can be noted from Fig. 2. A peak at 1.486 eV is seen at very low currents (Fig. 3C). Below 15  $\mu$ A it is the only emission present; at this point the main peak appears and the 1.486-eV peak, which grows superlinearly with current, saturates drastically at 20  $\mu$ A. With further increase of the current the main peak shifts toward higher energies and the 1.486-eV peak is then masked completely; its energy does not shift with current, before or after saturation. It is quite similar to one which appears at low current densities in Te-doped junctions<sup>4</sup> and saturates in the same fashion at somewhat higher current densities. This suggests that the transition does not involve sulfur atoms, but some residual impurity; its energy and saturation suggest a bound exciton annihilation, but further work remains to be done.

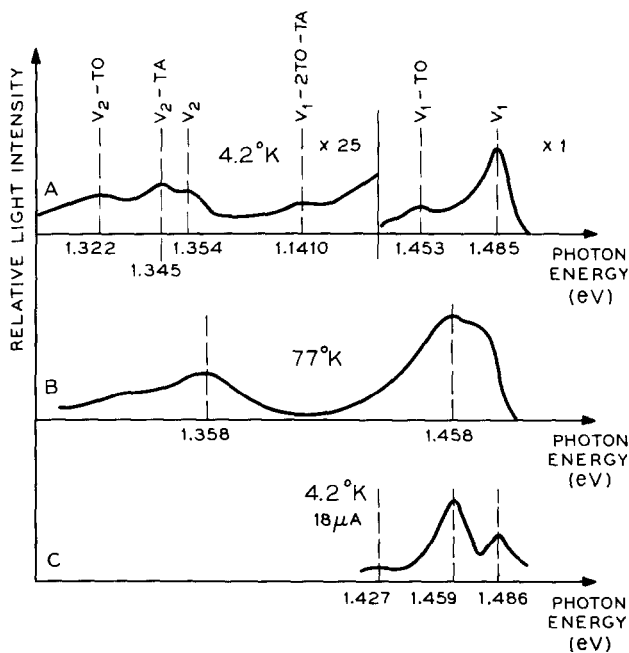


Fig. 3. Electroluminescence spectra at 4.2 and 77°K. A and B, 25 mA characteristic current; C, 18  $\mu$ A.

The rest of the spectrum at 4.2°K (Fig. 3A) consists of two main transitions;  $\nu_1$  at 1.485 eV and  $\nu_2$  at 1.354 eV. The remaining peaks can be explained as phonon satellites involving the emission of transverse optical (TO) phonons of 32-meV energy and of transverse acoustic (TA) phonons of  $\sim$ 9 meV as shown in Fig. 3A. The TO phonons have been reported by Cochran and others<sup>5</sup> and Spitzer.<sup>6</sup> The energy of the TA phonon has been uncertain; Cochran and others inferred a value of 9 to 8 meV from phonon combination bands while Spitzer explained the same data by a different phonon assignment where the lowest energy phonon is TA of 15.7 meV. Our direct observation of the 9-meV TA phonon lends support to the five phonon assignments of Cochran and others. As the spectrum in Fig. 3A was not seen in Te-doped diodes,<sup>4</sup> one is tempted to relate those lines to transitions involving sulfur atoms, the two main emissions  $\nu_1$  and  $\nu_2$  concerning sulfur atoms in different lattice positions. The second line  $\nu_2$ , which shifts with applied voltage, might also be due to an extraneous impurity in the S-doped diodes. The observation that the main emission occurs at a frequency lower than the 1.486-eV impurity line suggests that the strongest peak is not due to band-to-band transitions. The same is true for diodes with Te concentrations up to  $5 \times 10^{17}$  cm<sup>-3</sup> where the 1.486-eV peak was observed at low currents.

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