# Photoluminescence Enhancement and Structure Repairing of Monolayer MoSe<sub>2</sub> by Hydrohalic Acid Treatment

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#### ABSTRACT

Atomically thin two-dimensional transition metal dichalcogenides (TMDCs) have attracted much attention recently due to their unique electronic and optical properties for future optoelectronic devices. Chemical vapor deposition (CVD) method is able to generate TMDCs layers with a scalable size and a controllable thickness. However, the TMDC monolayers grown by CVD may incorporate structural defects and it is fundamentally important to understand the relation between photoluminescence and structural defects. In this report, point defects (Se vacancies) and oxidized Se defects in CVD-grown MoSe<sub>2</sub> monolayer are identified by transmission electron microscopy and X-ray photoelectron spectroscopy. These defects can significantly trap free charge carriers and localize excitons, leading to the smearing of free band-to-band exciton emission. Here we report that the simple hydrohalic acid treatment (such as HBr) is able to efficiently suppress the trap state emission and promote the neutral exciton and trion emission in defective MoSe<sub>2</sub> monolayers through the *p*-doping process, where the overall photoluminescence intensity at room temperature can be enhanced by a factor of 30. We show that HBr treatment is able to active distinctive trion and free exciton emissions even from highly defective MoSe<sub>2</sub> layers. Our results suggest that the HBr treatment not only reduces the *n*-doping in MoSe<sub>2</sub> but also reduces the structural defects. The results provide further insights of the control and tailoring the exciton emission from CVD-grown monolayer TMDCs.

## **Table of Content**



Atomically thin transition metal dichalcogenides (TMDCs) can be obtained by mechanical <sup>1, 2</sup> or chemical exfoliation <sup>3</sup> from their bulk crystals. Due to their unique and striking properties, monolayer TMDCs have attracted extensive attention. Semiconducting TMDC monolayers have been demonstrated feasible for various energy related applications, where their electronic properties <sup>1, 4, 5</sup> and high surface areas offer opportunities for applications such as biosensors, <sup>6,7</sup> nano generators, <sup>8</sup> green electronics, <sup>9</sup> electrocatalytic hydrogen generation <sup>10, 11</sup> and energy storage. <sup>12, 13</sup> Semiconducting TMDC monolayers, such as MoS<sub>2</sub>, <sup>14-16</sup> WS<sub>2</sub>, <sup>17-19</sup> MoSe<sub>2</sub> <sup>20, 21</sup> and WSe<sub>2</sub>, <sup>22-25</sup> exhibit direct band gaps and show attractive light emission properties in the visible and near-infrared spectral region, making them suitable for optoelectronic devices. <sup>26</sup> Applications based on TMDCs in light emitting diodes, <sup>27</sup> and photodetectors <sup>28, 29</sup> have been developed recently. Significant efforts have been devoted to the control and tuning of photoluminescence (PL) properties in TMDCs. Several approaches including electrostatic gating. <sup>30</sup> chemical treatment, <sup>31-33</sup> surface plasmonic excitation, <sup>34</sup> photonic crystal cavities <sup>35</sup> and strain engineering <sup>36, 37</sup> have been adopted to modulate the trions and neutral excitons in monolayer TMDCs. CVD-grown TMDC layers could exhibit unique PL property which is absent in their exfoliated TMDC counterparts. For example, extraordinary PL has been observed near the edge of CVD grown WS<sub>2</sub> recently. <sup>17</sup> Edges and point defects in monolayer TMDCs formed during CVD could be engineered to tailor their electronic, optical and catalytic properties. <sup>38-42</sup> Interestingly, the structural defects in TMDCs have been reported to darken and blue shift the PL peak in MoS<sub>2</sub> and WS<sub>2</sub> <sup>38</sup>. Controversially, recent reports have also demonstrated that point defects could activate PL in TMDCs and increase the overall PL intensity. <sup>39</sup> It still demands more experimental efforts to understand the point defects and excitonic characteristics in TMDC layers.

Chemical doping is known as an effective way to modulate the optical and electrical properties of TMDCs. Matsuda *et al.* have reported that the PL intensity of mechanically exfoliated MoS<sub>2</sub> monolayer can be enhanced when it is interacted with a *p*-dopant 2,3,5,6-tetrafluoro-7,7,8,8-tetracyanoquinodimethane (F4TCNQ), where the *p*-doing can switch the recombination of negative trions to the neutral exciton.<sup>32</sup> Yet, the PL intensity enhancement ratio is only around 3 with the F4TCNQ

doping.<sup>32</sup> More recently, Amani *et al.* have reported that PL and minority carrier lifetime in exfoliated MoS<sub>2</sub> monolayers can be uniformly enhanced by an organic super acid treatment. <sup>43</sup> However, the exact mechanism by which the passivation of surface defects is not fully understood. Furthermore, in the previous studies, the modulation of PL properties in TMDCs by surface chemistry and structural defect engineering have been mostly limited to mechanically exfoliated TMDCs. <sup>31, 32, 38, 39</sup> Compared to the exfoliation method, CVD-grown TMDCs are more practical for large scale device applications. <sup>44</sup> Since the CVD-grown TMDCs may exhibit a wide range of defect types and defect density levels, different from the mechanically exfoliated TMDCs, the defect and optical emission relations are yet to be explored.

In this contribution we demonstrate that the PL intensity of monolayer MoSe<sub>2</sub> grown from CVD can be effectively enhanced after exposure to hydrohalic acid vapors, such as HCl, HBr and HI. The overall PL intensity of MoSe<sub>2</sub> dramatically increases for more than 30 times by HBr treatment. Due to the structural defects that commonly exist in CVD grown TMDCs, neutral exciton and trion peaks are rarely observed in low temperature PL measurement. We found the defect trapped exciton can be greatly suppressed by the HBr treatment which allow us to observe the neutral exciton and trion peaks from CVD grown MoSe<sub>2</sub>. Furthermore, our scanning transmission electron microscope (STEM) characterization proves the presence of point defects in MoSe<sub>2</sub> structures. Raman spectrum, PL, X-ray photoelectron spectroscope (XPS) and STEM suggest that the HBr treatment imposes various effects which explain the PL enhancement such as removing impurities, *p*-doping to the MoSe<sub>2</sub>, and reducing the structure of MoSe<sub>2</sub>.

#### **RESULTS AND DISCUSSIONS**

Monolayer MoSe<sub>2</sub> was grown via the CVD process according to our previous reports, <sup>21</sup> where sulfur and MoO<sub>2</sub> powders were used as precursors. Figure 1a illustrates the optical micrograph (OM) of the CVD MoSe<sub>2</sub> monolayers synthesized on sapphire substrates, where the monolayer MoSe<sub>2</sub> exhibits a triangular shape with the lateral size up to 10  $\mu$ m. Here, we use hydrohalic acids as highly efficient agents to tune the exciton PL in MoSe<sub>2</sub> monolayers grown by CVD. Figure 1b schematically illustrates the experimental set-up for the hydrohalic acid treatment. Aqueous solutions of HCl, HBr and HI were bubbled and brought

to the sample by Ar gas, and the AgNO<sub>3</sub> was used to absorb the exhaust gas. Photoluminescence from the as-grown and treated MoSe<sub>2</sub> monolayers excited by a 532 nm laser are shown in Figure S1. Only one high intensity peak with photon energy around 1.53 eV was observed in both pristine and HBr treated MoSe<sub>2</sub>, which is close to the reported band gap of 1.55 eV for monolayer MoSe<sub>2</sub>. Figure 1c compares the PL intensity spatial mappings for as-grown and treated MoSe<sub>2</sub>. As shown in Figure 1d and 1e, PL signals are pronounced at the edge of the as-grown MoSe<sub>2</sub> and become very weak or absent towards the centre of the as-grown monolayer MoSe<sub>2</sub>. A blue-shift was also observed from the centre to the edge. Recently, a similar phenomenon has been observed on CVD-grown WS<sub>2</sub> monolayer by Yu and *et al*, <sup>40</sup> where they suggest that the defects within the crystal act as non-radiative recombination sites and thus quench the intrinsic PL. Consistently, the edge enhanced PL emission has been observed in CVD grown WS<sub>2</sub> monolayer in previous studies, and the darkening of PL in the centre of TMDCs island has be attributed to the charge defect-induced doping.<sup>17, 45</sup> Our result shows that the overall PL intensity of MoSe<sub>2</sub> dramatically increases for more than 30 times after HBr treatment, which suggests the PL tunability of defective monolayer MoSe<sub>2</sub>. Figures 1d and 1e display the PL line scans for PL intensity and photon energy across the sample before and after HBr treatment, respectively. Remarkably, the PL intensity increases significantly and the photon energy blue shifts at the centre of the triangle. For the spatial modulation of PL intensities and positions in MoSe<sub>2</sub>, various factors are responsible, including external electrostatic doping, structural defect and chemical composition change. The darkening of PL in the centre of MoSe<sub>2</sub> islands is attributable to the presence of structural defects such as point defects and dislocations within the metastable nuclei. <sup>38</sup> During the CVD growth procedure, the initial nucleation occurs at the beginning of the growth process followed by the incoming MoSe<sub>2</sub> species coalescing into the nuclei leading to an enlarged MoSe<sub>2</sub> grain. Therefore, the HBr molecules are more likely to adsorb at the nucleation centre of MoSe<sub>2</sub> and results in a more pronounced PL change. In addition to HBr, other hydrohalic acids such as HCl, and HI also show similar PL enhancement effects (see Figure S2). Among the three hydrohalic acids, HBr shows the best PL enhancement performance. In a previous report, <sup>46</sup> the light emission in mechanically exfoliated TMDC flakes were shown to be sensitive to H<sub>2</sub>O molecules.

We have separately examined the effect of moisture and no significant PL enhancement is observed with only H<sub>2</sub>O treatment (details in supporting Figure S3). Therefore, the effect of moisture can be excluded in this study. Furthermore, we have also excluded other possible gaseous adsorption effects, such as O<sub>2</sub>, by comparing the PL intensity change in vacuum for the sample before and after HBr treatment (details in supporting Figure S4). All the results consistently indicate HBr plays a major role in enhancing photoluminescence in MoSe<sub>2</sub>. In the following paragraphs we discuss the HBr treatment effect on the doping level, structural defects, and chemical composition and further reveal the mechanism of optical property modulation in MoSe<sub>2</sub>.

Raman scattering is known to be sensitive to the doping level of 2D materials. We use Raman spectroscopy to investigate the charge-phonon interaction in MoSe<sub>2</sub> layers. For a monolayer MoSe<sub>2</sub> only one Raman active mode (out-of-plane  $A_{1g}$ ) appears. Figure 2a shows the typical Raman spectrum for the MoSe<sub>2</sub> flakes before and after HBr treatment, where a predominate  $A_{1g}$  mode at ~240.4 1/cm is identified for the as-grown MoSe<sub>2</sub>. The out-of-plane vibrational A<sub>1g</sub> mode becomes stronger and shifts to a higher frequency at 241.1 1/cm after HBr treatment. Figure 2b shows the statistical measurements for the Raman frequency and Figure 2c shows the Raman intensity mappings for the as-grown and HBr treated MoSe<sub>2</sub> monolayers. The A1g mode results from the out-of-plane vibration of Se atoms in opposite directions, which couples more strongly with electrons than the in-plane vibrational mode does. Calculation based on first-principles density functional theory shows that electron doping leads to occupation of the bottom of the conduction band at the K point states resulting in a significant change in the electron-phonon coupling of A<sub>1g</sub>.<sup>42</sup> It has also been demonstrated experimentally that *n*-doping typically results in the softening and intensity decrease of the  $A_{1g}$  phonon while *p*-doping causes the hardening and intensity increase of A1g phonon. <sup>31, 47, 48</sup> Therefore, the A1g phonon renormalization can be used to estimate the doping level change in TMDCs. The intensity increase and frequency upshift in A1g mode suggest that HBr treatment imposes *p*-doping to MoSe<sub>2</sub> layers. <sup>49</sup> We provide quantitative analysis of the doping level (details in supporting materials, and Figure S5). The electron carrier density are estimated to be  $8.2 \times 10^{11}$   $cm^{-2}$  for as-grown and  $6.37 \times 10^{10} cm^{-2}$  for HBr treated MoSe<sub>2</sub> respectively, which are within a comparable range with the previous report. <sup>30</sup>

To further reveal the excitonic nature of MoSe<sub>2</sub> monolayer and the effect after HBr treatment, we perform temperature dependence PL measurement down to 10 K for the as-grown and HBr-treated MoSe<sub>2</sub> layers. Figures 3a compares the PL emission from as-grown and HBr treated MoSe<sub>2</sub> at 10K. The neutral exciton of MoSe<sub>2</sub> monolayer has been reported to emit at around 1.66 eV<sup>30</sup>. Figure 3a shows that a strong emission at a lower energy (around 1.56 eV) from our as-grown MoSe<sub>2</sub> with a full width at half-maximum (FWHM) value of 80.6 meV. The peak at 1.56 eV indicates that the radiative excitons are bound to defects which is in good agreement with the previous report<sup>39</sup> since the defect-related, trapped excitons should lead to an emission energy lower than the band-to-band optical transition energy. After HBr treatment, the defect related PL emission MoSe<sub>2</sub> is dramatically suppressed as revealed in Figure 3a. In addition to the broad sub-band-gap emission at around 1.59 eV, we observed two additional peaks at 1.66 and 1.63 eV, assigned as the neutral exciton  $(X^0)$  and the trion peak  $(X^-)$  respectively.<sup>30</sup> Figure 3b shows the evolution of neutral exciton and trion peaks as a function of temperature from 10 K to 300 K. The trion peak becomes negligible when the temperature is higher than 150 K. Figures 3c and 3d show the temperature dependence of the  $X^0$  and  $X^-$  peak position and intensity extracted based on the same fitting method in Figure 3a. Also, a binding energy of 30 meV is extracted, which is in good agreement with previous reports.<sup>30</sup> The temperature dependence of the peak intensity ratio ( $X^{-7} X^{0}$ ) for MoSe<sub>2</sub> monolayers shown in Figure 3d suggests that our MoSe<sub>2</sub> sample exhibits a higher trion to exciton ratio. The appearance of X<sup>0</sup> emission suggests that the highly n-doping feature in pristine MoSe<sub>2</sub> has been reduced considerably after HBr treatments. <sup>30</sup> The temperature dependence measurement of MoSe<sub>2</sub> PL suggests the defects within the as-grown MoSe<sub>2</sub> crystals prohibit the intrinsic exciton emission and the dominate PL peak is mostly from trapped exciton states. While for the HBr-treated MoSe<sub>2</sub>, the trapped exciton state is greatly suppressed, and both exiton and trion peaks are detectable at a low temperature.

Both theoretical <sup>50, 51</sup> and experimental <sup>52, 53</sup> works have demonstrated that the defects such as cation and anion vacancies in TMDCs could induce doping. In order to reveal the as-grown and treated MoSe<sub>2</sub> microscopically, MoSe<sub>2</sub> monolayers are transferred onto a transmission electron microscope (TEM) grid and visualized under TEM and STEM. The transfer process unavoidably induces the folding of MoSe<sub>2</sub> (Top area in Figure 4a). Figure 4b displays the high resolution TEM (HRTEM) image. The inset in Figure 4b corresponds to the Fast Fourier Transform (FFT) of the image Figure 4b, which clearly shows the hexagonal packing of MoSe<sub>2</sub> crystals. In addition to the TEM characterization, a more precise lattice structure measurement was carried out using aberration-corrected STEM. We further visualize the flake via chemical analysis using atomic-resolution Z-contrast imaging with high-angle annular dark field (HAADF) STEM. As the intensity of STEM images directly related to the atomic number (Z-contrast), individual Se atoms can be easily recognized and differentiated from the three-fold coordinated two Se atoms. Figures 4c and 4d are the STEM images of MoSe<sub>2</sub> before and after HBr treatment. By analysing the column-to-column intensity ratio of metal and chalcogen sites in Figures 4c and 4d, the lattice structures have been reconstructed and shown in Figures 4e and 4f. The STEM characterization identifies the presence of large number of single Se vacancies, and we also observe Se point defects in the as-grown MoSe<sub>2</sub> layer (as shown in Figure 4c and 4e up right corner). It was found that the individual Se vacancies are stable under TEM. Meanwhile, we note there are no Mo vacancies observed in the examined locations for both of as-grown and HBr treated samples. The presence of Se vacancies contributes to the *n*-doping and PL weakening in MoSe<sub>2</sub>. We can draw the conclusion that HBr treatment can effectively tailor the localized exciton in MoSe<sub>2</sub> defect sites structured by Se vacancies and release the defect bound excitons. Qualitatively, we note that the Se vacancies in HBr treated samples are relatively less than the as-grown MoSe<sub>2</sub> monolayers.

To obtain more insights of HBr treatment on the lattice structure of  $MoSe_2$ , XPS was used to further reveal the chemical composition change of CVD grown  $MoSe_2$ . As shown in Figure 5a, the binding energies of Mo  $3d_{5/2}$  and Mo  $3d_{3/2}$  are around 229.4 and 232.53 eV respectively, revealing the +IV oxidation chemical state of Mo. The small shoulder peaks at around 233.3 and 236.43 eV in Figure 5a

can be assigned to Mo 3d<sub>5/2</sub> and 3d<sub>3/2</sub> core levels of the hexa-valent state of Mo<sup>54</sup>, originated from not fully selenized MoO<sub>3</sub> precursors used in CVD growth. We note the hexa-valent state becomes barely detectable after HBr treatment, suggesting HBr-treatment can remove these impurities. It is noteworthy, after HBr treatment there is a red shift for the peak positions of Mo<sup>(IV)</sup> 3d<sub>5/2</sub> and 3d<sub>3/2</sub> from 229.6 eV to 229.4 eV which could suggest a reduction of  $Mo^{(IV)}$  oxidation sate or a p-doping<sup>55</sup>, consistent with the conclusion drawn from Raman measurements. Additionally, the full width at half maximum (FWHM) of Mo<sup>(IV)</sup> also decreases from 1.53 to 1.06 eV, indicating a more uniform chemical environment for Mo after the HBr treatment. In Figure 5b, the Se  $3d_{5/2}$  and  $3d_{3/2}$  doublets at (54.8 eV; 55.66 eV) and (55.34 eV; 56.2 eV) can be assigned to terminal Se<sup>2-</sup> and bridging Se<sub>2</sub><sup>2-</sup>, respectively<sup>56</sup>, and these are the two types of covalent bonds observed in both as-grown and HBr treated MoSe<sub>2</sub> layers. The binding energy of 3d<sub>5/2</sub> core levels at 55.9 eV is ascribed to elementary Se<sup>(0)</sup> impurities originated from Se element or the oxidized MoSe<sub>2</sub> (Mo-Se-O) state. The broad peak at around 59.2 eV is attributed to the oxidized Se species such as the SeO<sub>2</sub> impurity <sup>57</sup>. However, these peaks are observable in as-grown but not in HBr-treated MoSe<sub>2</sub> layers. It is known that HBr is able to remove the oxygen in SeO<sub>2</sub> thought the formation of SeBr<sub>4</sub>.<sup>58</sup> We suspect that HBr might be able to remove the oxygen species in oxidized MoSe<sub>2</sub>, where the Mo-Se-O could be transformed to Mo-Se-Br. This postulation agrees well with the observation of FWHM narrowing for Mo<sup>(IV)</sup> peak after HBr treatment. The binding energy evolution for Mo. Se and Br is summarized in Table R3 in Supporting Information. Note that the Br<sup>-</sup> ions are detected at around 68.6 eV (Br 3d<sub>5/2</sub> doublet) after HBr treatment as shown in Figure 5c, which is confirms the reaction between MoSe<sub>2</sub> and HBr. Note that in STEM we could not identify the Br atom since Br has a similar atomic number as Se (35 and 34 respectively). However, we do observe the presence of Br in the XPS and energyfiltered electron energy-loss spectrum (EELS) elementary mapping (shown in Figure S6). This suggests that Br ions adsorb on MoSe<sub>2</sub> as counter ions. Since Se and Br are not distinguishable in STEM, we cannot exclude the possibility that Br ions fill in the Se vacancies as illustrated in Figure 5d.

To further strengthen the arguments of structural repairing by HBr, we calculate the elemental ratio of Mo, Se and Br based on their corresponding XPS peak areas. The atomic ratio of anions versus Mo<sup>(IV)</sup>

is shown in Figure 5e. The ratio of as-grown MoSe<sub>2</sub> is 1.686 suggesting the presence of the high density of Se vacancies and Se<sup>(0)</sup> and SeO<sub>2</sub> impurities. Remarkably, after HBr treatment, the anions to Mo atomic ratio increased to 1.97, this increase is majorly contributed from the incorporation of additional Br anions and the increase in Se<sup>2-</sup>, where the increase is likely from the conversion of the bridging Se<sub>2</sub><sup>2-</sup> to the Se<sup>2-</sup>. We also note the SeO<sub>2</sub> becomes undetectable, but with a dramatically increased percentage of Se<sup>2-</sup>. Revealed by our STEM characterization, the lattice structure defects are mainly the Se vacancies. In brief, all the XPS results suggest that HBr not only *p*-dopes but also removes the impurities Se<sup>(0)</sup> and SeO<sub>2</sub>. Meanwhile, the HBr treatment most likely reduces the Se vacancies by filling Se vacancies using the Se in the bridging form Se<sub>2</sub><sup>2-</sup> to Se<sup>2-</sup>. The passivation of Se vacancies with Br ions might also occur at the same time.

#### CONCLUSIONS

In conclusion, point defects formed by Se vacancies can greatly quench the PL of monolayer MoSe<sub>2</sub> due to the trapping of free charge carriers and non-radiative recombination. Low temperature PL study shows that HBr can effectively suppress the trapped exciton states and populate both exciton and trion emission. Other hydrohalic acids such as HCl, and HI also show similar PL enhancement effects. However, HBr is the most effective chemical with a suitable acidity and thus provides better controllability for tailoring the optical properties of MoSe<sub>2</sub>. The drastic modulation of optical properties by HBr can be attributed to various reasons including removing impurities, *p*-doping to the MoSe<sub>2</sub>, and reducing the structure of MoSe<sub>2</sub>. Our method offers the insights for tuning the properties of 2D TMDCs.

### METHODS

*Growth of MoSe*<sub>2</sub> *layers:* We have developed the CVD method to synthesize monolayer MoSe<sub>2</sub>. The precursor of Mo and Se are MoO<sub>2</sub> powder (Sigma Aldrich,  $\geq$ 99.5% purity) and Se powder (Sigma Aldrich,

 $\geq$ 99.5% purity), respectively. MoO<sub>2</sub> powder (0.3 g) was placed in a ceramic boat at the center of furnace with a 1 cm ×5 cm sapphire substrate placed at downstream of ceramic boat. Selenium powder was heated by another furnace and carried by Ar/H<sub>2</sub> (Ar/H<sub>2</sub> = 65 sccm/5 sccm, 50 torr). The centre of the furnace was gradually heated from room temperature to 800 °C at the rate of 25 °C/min. The temperature of the Se powder at the upstream furnace was raised up to 260 °C when the growth furnace temperature reached 800 °C. In order to growing MoSe<sub>2</sub> triangle domains, the furnace was held at 800 °C for 15 min. After growth, the furnace was naturally cooled down to room temperature.

*Chemical treatment of MoSe<sub>2</sub> layers:* After getting the MoSe<sub>2</sub> triangle domains on sapphire substrate, we used hydrobromic acid (Alfa Aesar, 49% liquid) or other hydrohalic acids as a chemical source. The bubble method was performed to spray the HBr droplets onto the sapphire substrate where the MoSe<sub>2</sub> domains were prepared. After the evaporation of chemical solvent, the MoSe<sub>2</sub> triangle domains were characterized by optical measurements.

*Transfer method for pristine and HBr treated MoSe<sub>2</sub>:* The TMDCs layers were transferred onto target substrates or TEM grids by a poly (methyl methacrylate) (PMMA) (950 PMMA A4, Micro Chem) assisted method. First PMMA thin film was spin-coated on top of the sample, and then the PMMA/MoSe<sub>2</sub>/sapphire was dipped in a 6M HF solution to etch the MoSe<sub>2</sub>/sapphire interface. PMMA/MoSe<sub>2</sub> was lifted from the etching solution and diluted in DI water, and then transferred onto target substrates. The PMMA layer was removed with acetone and isopropanol.

*Characterizations:* The aberration-corrected JEOL-2100F cold field emission gun electron microscope equipped with a DELTA corrector was used to observe the annular dark field (ADF) images in this work. The ADF images were obtained at an accelerating voltage of 60 kV at room temperature. The probe current we used to detect was about 10~15 pA. The scanning rate of the ADF images was set in the region of 8 to 128 µs/pixel. Photoluminescence spectra and Raman spectra were excited by a green laser with 532 nm wavelength in a Witec alpha 300 confocal Raman microscopic system. The spot size of laser beam is about 0.5 µm and a 0.9 N.A. of 100X objective from Carl Zeiss Microscopy GmbH was used to collected the emitted Raman (1800 lines/mm grating) and PL signal (300 lines/mm grating). A higher grating was used to obtain more details Raman line shapes. The Raman wavelength calibration used the

Si Raman peak which is at 520 cm<sup>-1</sup> to be a reference. The sample was cooled down to T = 4 K in a lowvibration cryogen-free cryostat for the low-temperature PL measurement. A 0.42 N.A. objective lens was used in the 4K low-temperature measurement for having long working distance. Chemical configurations were determined by X-ray photoelectron spectroscopy (XPS, Phi V5000). XPS measurements were performed with an Al K $\alpha$  X-ray source on the samples. The energy calibrations were made against the C 1s peak to eliminate the charging of the sample during analysis.

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**Supporting Information Available**: Figures S1-S3 are included. This material is available free of charge *via* the Internet at <u>http://pubs.acs.org</u>.

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