

Polarization splitting directional coupler using tilted subwavelength gratings

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On-chip polarization splitters are key elements for coherent optical communication systems and polarization diversity circuits. These devices are often implemented with directional couplers that are symmetric for one polarization and strongly asymmetric for the other polarization. To achieve this asymmetry, highly dissimilar waveguides are used in each coupler arm, often requiring additional material layers or etch steps. Here we demonstrate polarization splitting with a directional coupler composed of two fully etched subwavelength waveguides, which only differ in the tilt angle of the silicon segments. Our device exhibits deep-UV compatible feature sizes, is only 14 μm long, and covers a 72 nm bandwidth with insertion losses below 1 dB and an extinction ratio in excess of 15 dB.

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Silicon-on-Insulator (SOI) offers multiple benefits as a photonics platform, including CMOS compatible fabrication, and a high index contrast that enables dense integration [1]. The high contrast enables waveguiding at sub-micron scale and, consequently, very high integration densities, albeit at the cost of a high polarization dependence, which is why polarization diversity schemes are required for many applications [2–5]. While larger micron-scale waveguides can significantly reduce this polarization dependence [6,7], they do so at the expense of integration density. Polarization beam splitters (PBS) are key building blocks for on-chip polarization diversity architectures. For applications in optical communications, such as polarization multiplexed coherent receivers, PBS with an extinction ratio of 16 dB are deemed almost ideal as digital signal processing is used at the receiver end. Moreover, insertion losses have to be kept as low as possible [5].

PBS can be implemented with a variety of approaches, including 1D or 2D grating couplers [8,9], multimode interferometers [10], slotted waveguides [11], mode evolution based devices [12,13] or devices designed via inverse/non-linear algorithms [14,15]. Among those approaches, mode evolution based PBS traditionally show the better performance at expense of relatively large footprints. On the contrary, inverse/non-linear algorithm designs achieve extremely short devices at expense of a complex design process and a limited low-loss bandwidth. However, arguably the most widely used mechanism for polarization splitting in integrated photonic chips is directional coupling (DC-PBS) [16–26]. These splitters can be classified in two categories, symmetric and asymmetric. Symmetric DC-PBS comprise two identical coupled waveguides where both polarizations satisfy the phase matching condition. The device length L is designed to verify $L = mL_{\pi}^{\text{TE}} = nL_{\pi}^{\text{TM}}$, where L_{π} is the coupling length of the TE/TM modes, and m and n are integers with $|m - n|$ being an odd number. However, symmetric DC-PBS are often long because the above condition usually requires high values for m and n . Moreover, the operational bandwidth and robustness of symmetric DC-PBS is limited because they exhibit a wavelength dependent behavior for both polarizations caused by the narrowband behavior of the phase matching condition. On the contrary, in an asymmetric DC-PBS only one of the polarizations satisfies the phase matching condition, thus coupling to the second waveguide, while the other polarization remains in the original waveguide. The device length is then optimized for the polarization that couples to the second waveguide, while the other polarization does not impose additional restrictions. As a result, asymmetric configurations are typically more compact and exhibit a better performance and tolerance to fabrication errors. On the downside, achieving the required asymmetry complicates the design and fabrication process, requiring additional materials [17,22,25], or two etch steps [26].

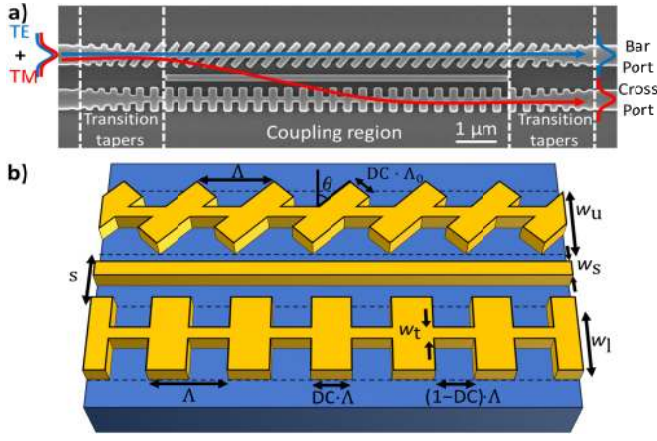


Fig. 1. (a) Scanning electron image (SEM) of the directional coupler based polarization beam splitter (before SiO₂ cladding deposition). The tilted SWG segments in the top waveguide break the structural symmetry for the TE mode, which remains in this waveguide. The TM mode is unaffected by the tilt and couples to the cross port. (b) Schematic of the coupling region of the polarization beam splitter (the cladding is not shown for the sake of clarity).

Subwavelength metamaterials, which can be fabricated in a single etch step [27–29], offer the promising capability to implement lithographically tailorable metamaterials modeled as homogeneous uniaxial crystals [30]. These structures can be exploited to implement asymmetric directional couplers by combining conventional and subwavelength waveguides [31,32]. Recently it has been shown that tilting the subwavelength segments significantly changes the propagation constant of TE polarized modes [33,34], but does not substantially affect the TM modes. In this paper we leverage this property to demonstrate for the first time a high-performance asymmetric DC-PBS based on two subwavelength waveguides which only differs in the tilting angle of the silicon segments, as shown in Fig. 1(a). This tilt breaks the phase matching for TE polarization, without affecting the cross-coupling for TM polarization. Our device exhibits a compact footprint of $2 \times 14 \mu\text{m}^2$. The measurements demonstrate a 72 nm bandwidth with a polarization extinction ratio in excess of 15 dB and, critically, sub-decibel insertion loss. Our device, shown in Fig. 1(b), is based on two subwavelength grating (SWG) waveguides separated by a distance $s = 0.5 \mu\text{m}$, to allow power coupling between them. The strip between the arms is a performance optimization that is detailed below. The SWG waveguide widths are initially set to $w_u = w_l = 0.6 \mu\text{m}$ to sufficiently confine the field, thus avoiding leakage losses [35], while being insensible to fabrication jitter [36]. The period of the structure is $\Lambda_0 = 275 \text{ nm}$ when $\theta = 0^\circ$, with a duty cycle $DC = 0.5$ to facilitate fabrication. The tip width $w_t = 100 \text{ nm}$ provides a smooth transition between the input waveguide and the SWG structure. In this case, we have chosen a typical 100 nm tip, however, this value could be slightly increased to facilitate fabrication but would result in a longer device. The grating segments of the upper waveguide are tilted by an angle, θ . We enlarge the period of both structures to $\Lambda = \Lambda_0 / \cos(\theta)$, thereby maintaining the distance between the SWG segments. The smallest feature of the device is w_t at the homogeneous-to-periodic taper tip. We use a standard 220-nm-thick silicon layer with a silicon dioxide top cladding with refractive indexes $n_{\text{Si}} = 3.476$, $n_{\text{SiO}_2} = 1.444$, at a central wavelength of $\lambda_0 = 1.55 \mu\text{m}$. To illustrate the principle of operation of our device, we first consider a simplified structure

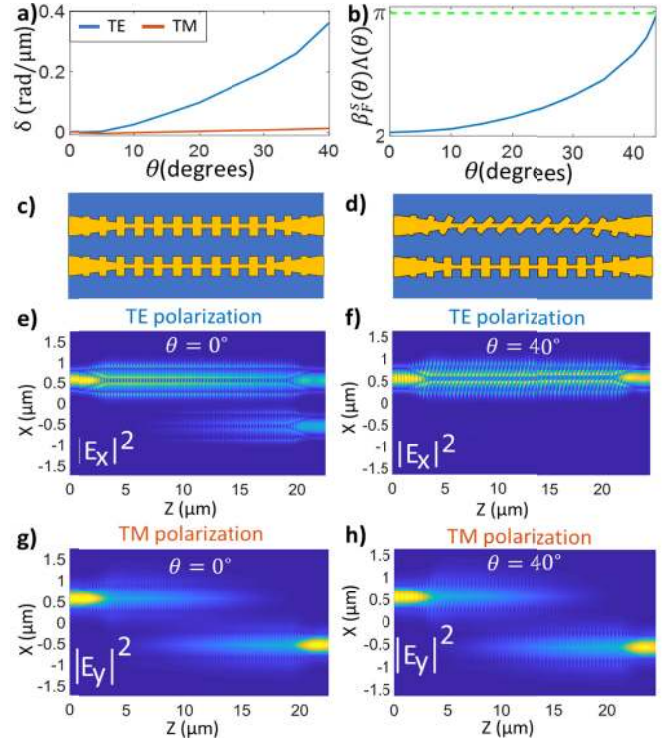


Fig. 2. (a) TE and TM phase matching detuning between the upper and lower waveguides with respect to the tilt angle θ . (b) Phase length of a period as a function of the tilt angle for the fundamental TE supermode. The Bragg condition is reached when $\beta_F^z(\theta) \cdot \Lambda(\theta) = \pi$, where β_F^z is the phase constant of the fundamental TE supermode and $\Lambda(\theta) = \Lambda_0 / \cos(\theta)$. (c), (d) Schematic representation of the proposed directional coupler for a tilt angle of (c) $\theta = 0^\circ$ and (d) $\theta = 40^\circ$. (e), (f), (g), (h) 3D-FDTD simulations of the squared magnitude of the main component of the propagating electric field for (e), (f) TE and (g), (h) TM polarization for a tilt angle of (e), (g) $\theta = 0^\circ$ and (f), (h) $\theta = 40^\circ$.

without the silicon strip between the directional coupler arms [See Fig. 2(c) and Fig. 2(d)]. According to coupled mode theory [37], in a directional coupler the maximum power transferred to the cross port is inversely proportional to the detuning δ between the effective indexes of the modes of the individual waveguides with $\delta = \pi \Delta n_{\text{eff}} / \lambda$. Figure 2(a) shows the detuning between the waveguides as a function of the tilt angle for both TE and TM polarization. The effective indexes of the modes were calculated using 3D-FDTD simulations [38], following the procedure outlined in Ref. [39]. As expected, increasing the tilt angle results in a strong detuning for the fundamental TE modes, while the effect on the fundamental TM modes is approximately 25 times smaller for the 40° tilt structure. While higher tilt angles would result in higher detuning between the individual TE waveguide modes, the fundamental TE supermode reaches the Bragg condition $\beta_F^z(\theta) \Lambda(\theta) = \pi$ for a tilt angle of $\theta \sim 45^\circ$ [See Fig. 2(b)], due to the effective increase of the period with the tilt angle. Here β_F^z is the phase constant of the fundamental TE supermode. We compensate the minor detuning for TM polarization at a 40° tilt angle by slightly changing the width of the lower waveguide from $0.6 \mu\text{m}$ to $0.65 \mu\text{m}$. The effect of the asymmetry introduced by the tilt is illustrated in Figs. 2(e), (f) and Figs. 2(g), (h) for tilt angles $\theta = 0^\circ$ and $\theta = 40^\circ$. The propagation of the electromagnetic field reveals that while TM polarization completely couples to the cross port in both cases, the 40° tilt helps completely isolate the waveguides for TE polarization.

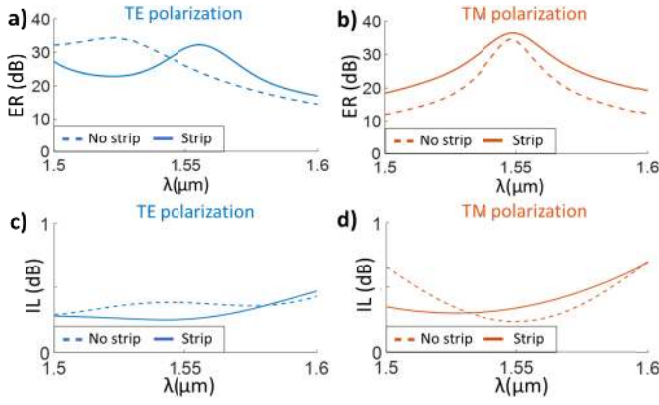


Fig. 3. Simulated extinction ratio for the (a) TE and (b) TM polarizations, and insertion loss for the (c) TE and (d) TM polarizations, of the tilted SWG directional coupler with (solid line) and without (dashed lines) the central silicon strip.

To evaluate the performance of this preliminary design, we use the insertion loss and extinction ratio, defined as:

$$IL^{TE} = P_{\text{bar}}^{TE}/P_{\text{in}}^{TE}; \quad IL^{TM} = P_{\text{cross}}^{TM}/P_{\text{in}}^{TM}; \quad (1)$$

$$ER^{TE} = P_{\text{cross}}^{TE}/P_{\text{bar}}^{TE}; \quad ER^{TM} = P_{\text{bar}}^{TM}/P_{\text{cross}}^{TM}; \quad (2)$$

where $P_{\text{bar}}^{TE/TM}$ and $P_{\text{cross}}^{TE/TM}$ is the power in the fundamental TE/TM mode in the bar and cross ports, and $P_{\text{in}}^{TE/TM}$ is the power of the fundamental TE/TM mode the input port. To calculate these parameters, we analyze the structure with 3D-FDTD simulations for different wavelengths. The dashed lines in Fig. 3 show the simulated IL and ER of the tilted SWG directional coupler, achieving an ER > 20 dB over a 40 nm bandwidth with negligible IL, being mainly limited by the wavelength dependent phase matching condition for TM polarization.

To improve the bandwidth of the device, we introduce the central silicon strip (see Fig. 1), transforming the device into a bridged directional coupler [32,40,41]. From an intuitive point of view, this silicon central strip only affects the even TM supermode, thereby allowing us to engineer the coupling length, defined as:

$$L_{\pi}^{TM}(\lambda) = \frac{\lambda}{2[n_1^{TM}(\lambda) - n_2^{TM}(\lambda)]}, \quad (3)$$

where λ is the free-space wavelength and $n_{1,2}^{TM}$ are the effective indices of the even and the odd TM polarized supermodes. In Fig. 4(a) we show the resulting coupling lengths for TM polarization for the proposed directional coupler with a tilt angle of 40° when we change the width of the silicon strip. It is observed that as the strip is widened, the coupling length is reduced. Moreover, as the variation of the beat length with wavelength is, to first order, proportional to the square of the beat length, the device also becomes less wavelength dependent. Increasing the width w_s over 150 nm has an undesired influence on the losses of the device [See Fig. 4(b)], which is why we choose the 150-nm-wide silicon strip which yields an almost flat wavelength response with a small penalty in insertion losses. The solid lines in Fig. 3 show the performance of the optimized PBS design with a tilt angle $\theta = 40^\circ$ and a silicon strip width $w_s = 150$ nm, revealing an ER higher than 20 dB over an 86 nm bandwidth for both polarizations with IL lower than 1 dB. The input tapers and the coupling region are 8 periods and 27 periods long, respectively, resulting in a device length of only $14 \mu\text{m}$.

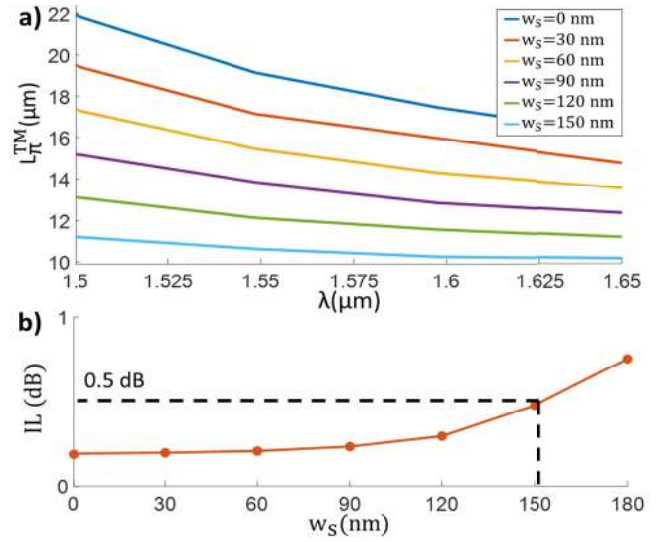


Fig. 4. TM polarization (a) coupling length and (b) insertion loss, of the tilted SWG directional couplers with different silicon strip widths.

The device has been fabricated in an open-access commercial 220-nm silicon-on-insulator platform via e-beam lithography with a 2- μm -thick buried oxide (BOX) and a 2.2- μm -thick SiO₂ upper cladding [42]. In our measurement setup light from a tunable laser is coupled into a lensed polarization maintaining fiber. The fiber is mounted on a rotary stage, which, together with a Glan-Thompson polarizer, is used to achieve horizontal (TE) and vertical (TM) polarization states with a high extinction ratio. The light from the lensed fiber is coupled into the chip by using an SWG edge coupler with high and broadband coupling efficiency for both polarizations [43]. At the output, light is focused onto a photodetector via a microscope objective, with negligible crosstalk between the output ports which are separated by $30 \mu\text{m}$. The fabricated device is characterized by measuring the power in both output ports (bar and cross) when the device is excited through the input port. Each device is excited first with TE polarized waveguide mode and then with TM mode. To determine the insertion loss, we use a reference waveguide to calibrate the setup and fiber-to-chip coupling losses for both polarizations. We then calculate the polarization extinction ratio by dividing the power in both output ports [see Eq. (2)]. In Figs. 5(a) and 5(b), we show the measured polarization extinction ratio and insertion loss of the fabricated device for both polarizations after applying the minimum phase technique [44]. The fabricated nominal device achieves an extinction ratio in excess of 15 dB and an insertion loss below 1 dB for both polarizations over a 72 nm bandwidth (1505 nm – 1577 nm). Figure 5(a) indicates that for long wavelengths the device performance is limited by the extinction ratio for TE polarization. Operation at longer wavelengths can be achieved by increasing the duty-cycle of the SWG structure, to enhance the modal confinement, which in turn suppresses the coupling of the TE mode to the cross port. Note that, for TM polarization, this increase in duty-cycle will only result in a minor detuning of the phase matching condition. In Figs. 5(c) and 5(d), we show the measured device performance when the duty cycle is increased from 50% to 60%, i.e. the silicon blocks are 28 nm longer, and the overall device length is enlarged two periods. This biased device this yields insertion loss below 1 dB and extinction ratio in excess of 15 dB at longer wavelengths, from 1534 nm to 1601 nm, i.e. a 67 nm bandwidth. Note that for applications incoherent optical

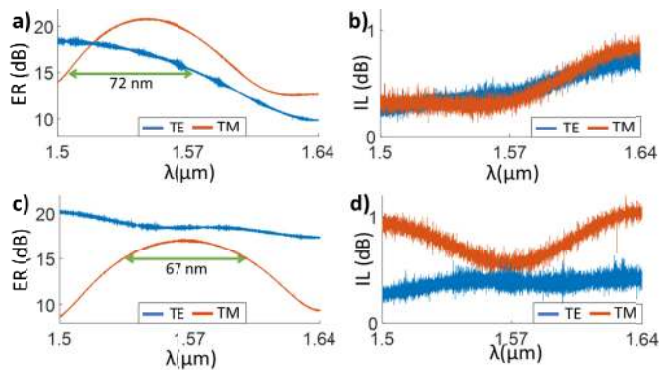


Fig. 5. Measured (a) extinction ratio and (b) insertion loss of the nominal device. The IL is below 1 dB and the ER is better than 15 dB in bandwidth of 72 nm. Measured (c) extinction ratio and (d) insertion loss for a device with a duty cycle increased by 10%. The IL are <1 dB and the extinction ratio is >15 dB in a bandwidth of 67 nm.

receivers these specifications result in a power penalty of less than 1.5 dB [5].

In conclusion, we have proposed and experimentally demonstrated a new type of directional coupler based polarization beam splitter, which uses tilted subwavelength segments to control the symmetry between the coupler arms. The device exhibits a comparatively large minimum feature size of 100 nm and can be fabricated in a single etch step. The low insertion loss of the device, lower than 1 dB over the 143 nm measured bandwidth, place the device as an ideal component for polarization multiplexed coherent receivers. We believe that this new concept of refractive index symmetry engineering using tilted subwavelength gratings will open promising venues for polarization control in integrated optical devices.

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