

# Probing dark matter couplings to top and bottom at the LHC

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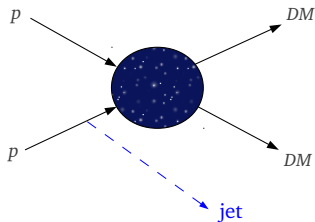
with Liantao Wang and Rocky Kolb, arXiv:1303.6638

April 17, 2013

# Dark matter at colliders

“Mono” searches

- **Monojets:** DM production ( $E_T^{miss}$ ) plus one or two jets
- Can arise from initial state radiation of quarks, gluons

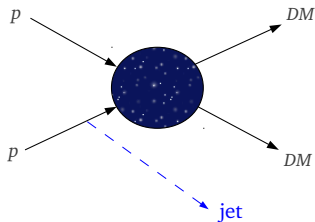


- Useful collider probe for light dark matter, and operators with SD scattering
- Related: monophoton, mono- $Z$ , mono-leptons, **mono- $b$**

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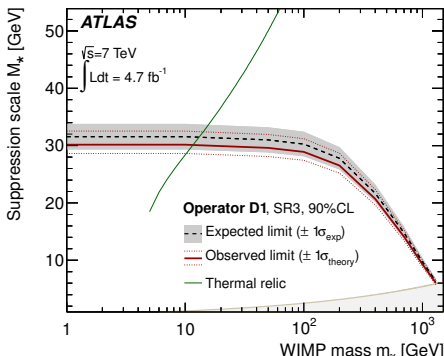
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## Constraints on operators

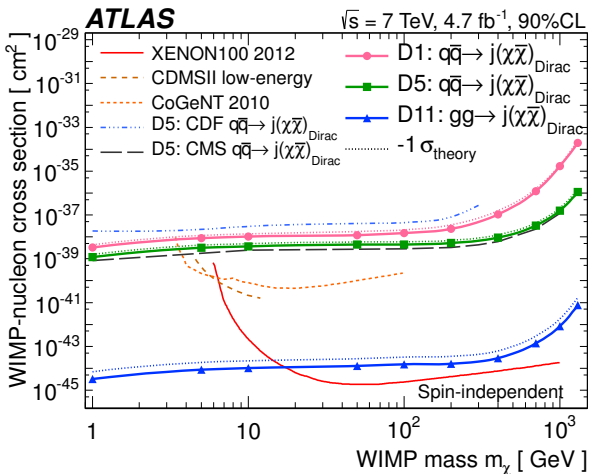
Model-independent effective field theory (EFT) approach: simple connection to direct detection and relic abundance (annihilation cross section)

Single scale  $M_*$  parameterizes size of interactions:

$$\text{D1/Scalar} : \mathcal{O} = \frac{m_q}{M_*^3} \bar{q}q \bar{X}X$$

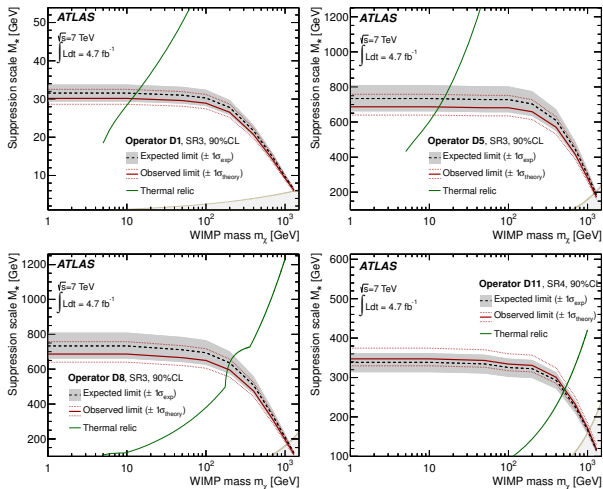


# Map to direct detection



# Constraints on operators

Constraints on  $M_*$  much worse for scalar operator:



# Scalar operator

A closer look at scalar operators:

$$\mathcal{O}_S \equiv \frac{m_q}{M_*^3} \bar{q}q \bar{X}X \sim \frac{y_q g_X}{m_\phi^2} \bar{q}q \bar{X}X$$

- Why  $m_q$ ? Flavor constraints  $\rightarrow$  Minimal Flavor Violation
- Suppressed couplings to light quarks means suppressed monojet signal from gluons and light quarks
- ATLAS published result (2012) - only consider up to  $m_c$

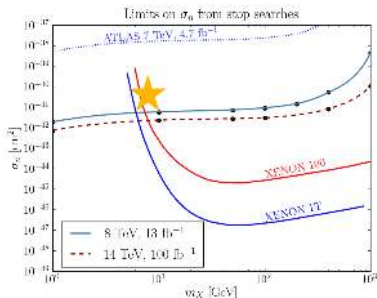
This operator gives rise to spin-independent scattering:

$$\sigma_n = \frac{(0.38m_n)^2 \mu_X^2}{\pi M_*^6} \approx 2 \times 10^{-38} \text{cm}^2 \left( \frac{30 \text{ GeV}}{M_*} \right)^6$$

(Other operators like  $\bar{q}\gamma^5 q \bar{X}\gamma^5 X$  give similar collider constraints but velocity-suppressed direct detection cross sections.)

# Searches with bottoms and tops

- Include heavy quarks in monojet signal
- Significant improvement in constraints on scalar interactions with  $b$ -tagging and  $t\bar{t}$  final states
  - heavy flavor production:  $y_q^2$  wins over pdf suppression
  - $b$ -enriched final states from bottoms and tops
- Direct search for couplings to third generation

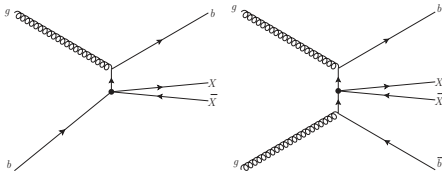




# DM production

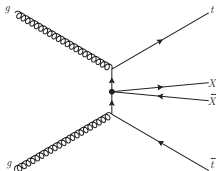
Direct  $b$  production:

- $bg \rightarrow bX\bar{X}$ ,  $gg \rightarrow b\bar{b}X\bar{X}$



DM production in association with tops:

- $bg \rightarrow bX\bar{X}$ ,  $gg \rightarrow b\bar{b}X\bar{X}$



Factor of  $\sim 10$  larger cross section than direct  $b$  production.

## *b*-tagging on monojets

### Monojet:

- $E_T^{miss} > 350\text{GeV}$ ,  $p_T^1 > 100\text{ GeV}$ , No lepton
- No more than 2 jets with  $p_T > 50\text{ GeV}$

### Mono-*b*:

- ***b*-tag on leading jet**

*b*-tagging efficiency (calibrated on tops)  $\sim 60\%$ , mistag  $\sim 0.2\%$  for light quarks at ATLAS

Dominant SM backgrounds for monojet searches:  $Z$ +jets ( $W$  + jets) cut from 500 fb to around 15 fb with *b*-tag.

## Mono- $b$

	Process	Monojet	$b$ -tag	$b$ -tag on $j_1$
Signal	$\bar{X}X + \text{jets}$	11 fb	0.9 fb	0.7 fb
	$\bar{X}X + b + \text{jets}$	65 fb	40 fb	33 fb
	$\bar{X}X + t\bar{t}$	120 fb	63 fb	41 fb

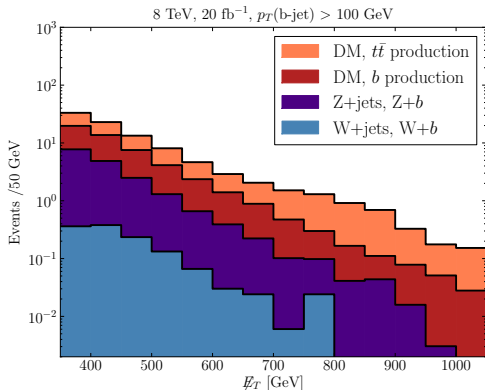
Table: **Inclusive mono- $b$  search at 8 TeV:** Cross sections after cuts applied. For the signal we take  $M^* = 50$  GeV.

**Improvement in bounds from: (1) overall monojet production rate is bigger (2) monojets are  $b$ -enriched**

Simulation of Events:

- MadGraph5 + PYTHIA6 + Delphes
- NLO cross sections from MCFM-Dark (Fox and Williams)

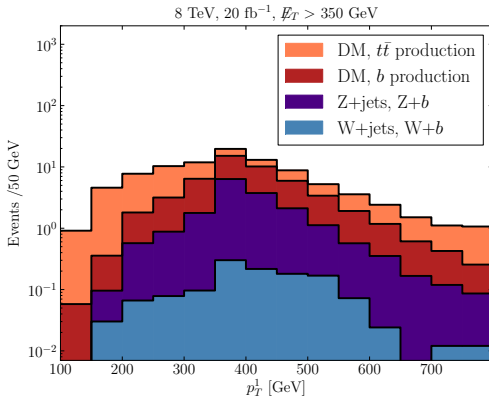
# Mono- $b$ MET distribution



$$(M_* = 50 \text{ GeV})$$

- Direct  $b$ -production isn't significantly harder
- $Z$ +jets often initiated by at least one valence quark
- $t\bar{t}$  has different kinematics and a harder spectrum; it isn't coming from ISR

# Mono- $b$ $p_T$ distribution

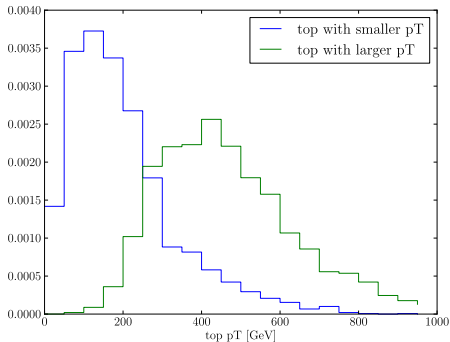


( $M_* = 50$  GeV)

- For  $b$  production, typically just one hard  $b$ -jet
- For top production, second jet may be harder,  $b$ -tagged 30% of the time

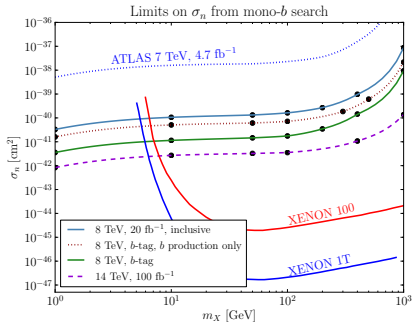
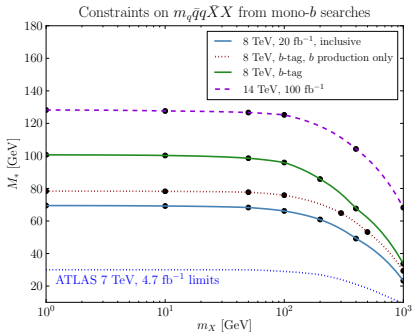
# Top kinematics: events passing mono- $b$ cuts

Associated production of  $t\bar{t}$  can look like “mono- $b$ ” cuts if one top is boosted and the other is not too hard.



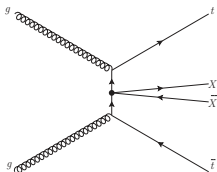
Efficiency WRT veto on 3rd jet: 0.17

# Mono-b limits



## $t\bar{t}$ final state

A direct search for  $t\bar{t} + \text{MET}$  would do better in this case:



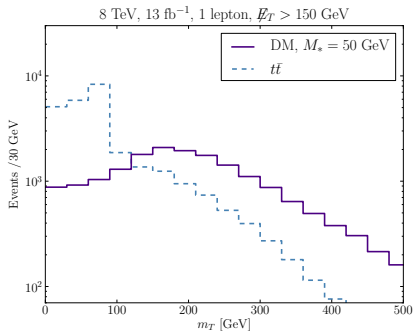
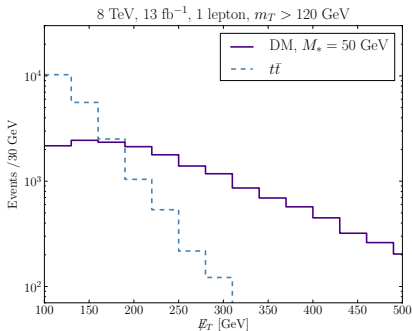
→ SUSY top-partner searches

We apply the ATLAS 8 TeV,  $13 \text{ fb}^{-1}$  study with 1 lepton in the final state:

- 4 jets with  $p_T > (80, 60, 40, 25)$  GeV with 1  $b$ -tag
- We use signal region with  $E_T^{\text{miss}} > 225$  GeV,  $m_T > 130$  GeV

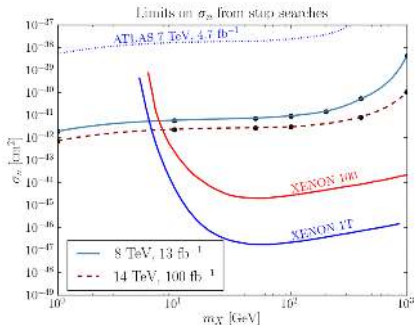
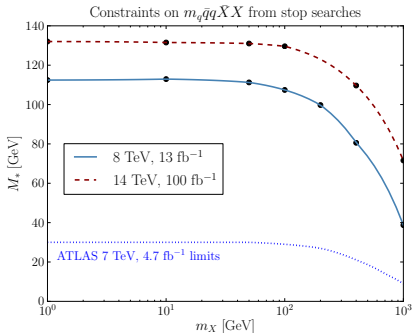


# $t\bar{t}$ final state



Dominant SM background is  $t\bar{t}$  dileptonic decay, where one lepton is either missed or a  $\tau$ .

# Limits from recasting stop search

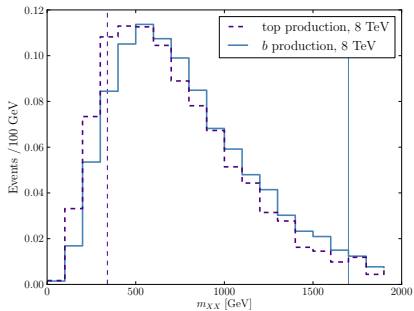
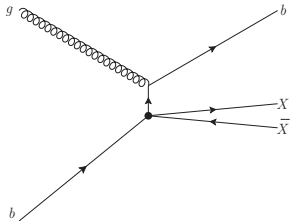


$\sim 20\%$  improvement in  $M_*$  constraint compared to mono- $b$ .  
Another  $10\%$  possible with all-hadronic final state.

# Unitarity

Bounds on  $M_*$  still low - what do these constraints mean?  
(Shoemaker and Vecchi; Fox, Harnik, Primulando, Yu):

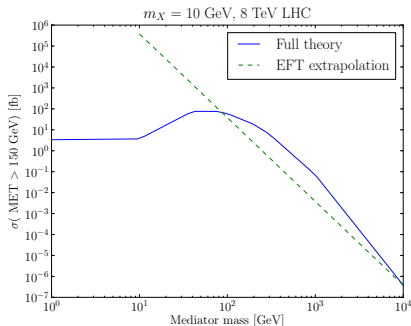
$$\frac{\sqrt{3}}{8\pi} \frac{m_q}{M_*^3} m_{\chi\bar{\chi}}^2 \lesssim \frac{1}{2}$$



$b$ -production seems safe, for top production it depends...

## Example: neutral scalar mediator

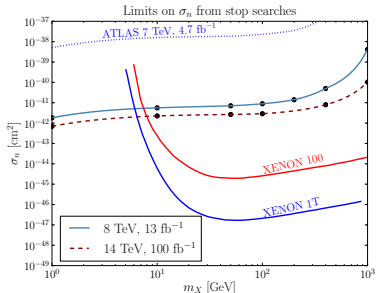
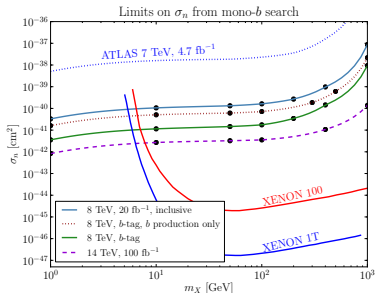
- Interactions through a mediator of mass  $M_\phi$  can go through a resonance:



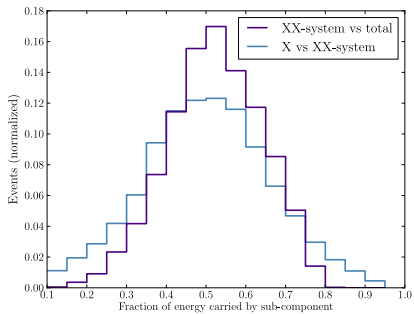
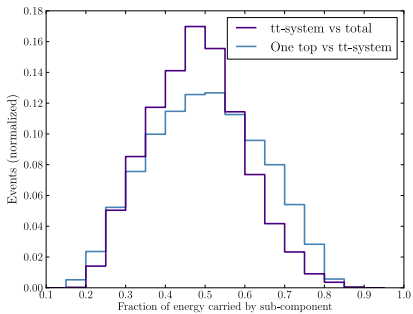
- Other constraints can also apply depending on the model.
- If the scalar is the Higgs, the constraints on Higgs invisible decay -  $Br(H \rightarrow \chi\chi)$  - are already stronger.

# Conclusions

- Monojets searches useful to study DM-SM interactions
- For scalar operators, final states are  $b$  enriched. “Mono- $b$ ” gives significant improvement in bounds.
- Top production even more important: direct  $t\bar{t} + E_T^{miss}$  search
- Direct probe of couplings to third generation

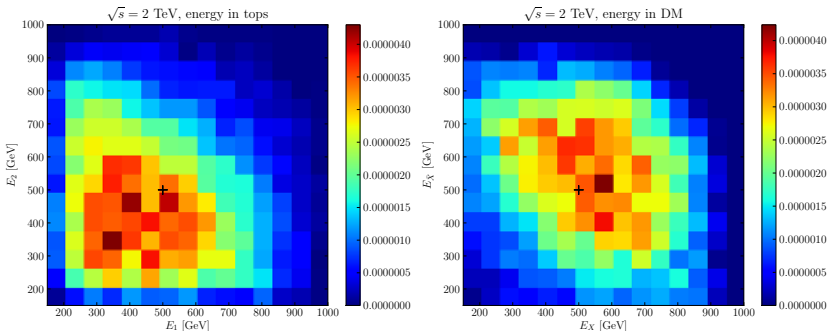


# Top kinematics



# Top kinematics

Compared to  $t\bar{t}$  or stop pair production, in this case each particle carries off about  $\sim 1/4$  of the energy in roughly uncorrelated directions  $\rightarrow$  large  $E_T^{miss}$ .



2D parton level distributions for energies of tops (left) and DM (right)

# Loops

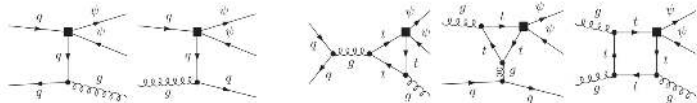


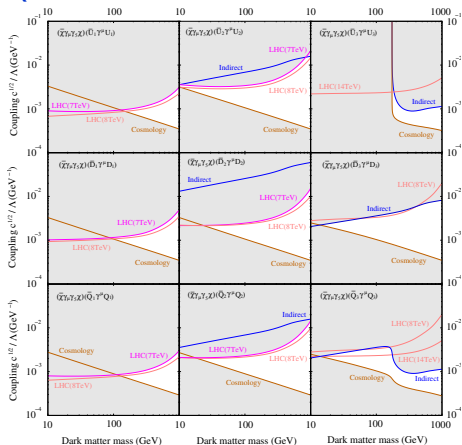
FIG. 1. Typical tree-level (left) and loop-level (right) diagrams leading to minijet events. The black squares denote insertions of four-fermion operators.

Haisch, Kahlhoefer, Unwin:  $M_*$  limit raised to 150 GeV.



# Scalar operator constraints

- Recent paper by Bhattacharjee et al. [1212.5013], “Model Independent Analysis of Interactions between Dark Matter and Various Quarks”:



- Kamenik and Zupan 2011 - “monotop” from MFV and dark matter