

Progress during TOGA in understanding and modeling global teleconnections associated with tropical sea surface temperatures

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Abstract. The primary focus of this review is tropical-extratropical interactions and especially the issues involved in determining the response of the extratropical atmosphere to tropical forcing associated with sea surface temperature (SST) anomalies. The review encompasses observations, empirical studies, theory and modeling of the extratropical teleconnections with a focus on developments over the Tropical Oceans-Global Atmosphere (TOGA) decade and the current state of understanding. In the tropical atmosphere, anomalous SSTs force anomalies in convection and large-scale overturning with subsidence in the descending branch of the local Hadley circulation. The resulting strong upper tropospheric divergence in the tropics and convergence in the subtropics act as a Rossby wave source. The climatological stationary planetary waves and associated jet streams, especially in the northern hemisphere, can make the total Rossby wave sources somewhat insensitive to the position of the tropical heating that induces them and thus can create preferred teleconnection response patterns, such as the Pacific-North American (PNA) pattern. However, a number of factors influence the dispersion and propagation of Rossby waves through the atmosphere, including zonal asymmetries in the climatological state, transients, and baroclinic and nonlinear effects. Internal midlatitude sources can amplify perturbations. Observations, modeling, and theory have clearly shown how storm tracks change in response to changes in quasi-stationary waves and how these changes generally feedback to maintain or strengthen the dominant perturbations through vorticity and momentum transports. The response of the extratropical atmosphere naturally induces changes in the underlying surface, so that there are changes in extratropical SSTs and changes in land surface hydrology and moisture availability that can feedback and influence the total response. Land surface processes are believed to be especially important in spring and summer. Anomalous SSTs and tropical forcing have tended to be strongest in the northern winter, and teleconnections in the southern hemisphere are weaker and more variable and thus more inclined to be masked by natural variability. Occasional strong forcing in seasons other than winter can produce strong and identifiable signals in the northern hemisphere and, because the noise of natural variability is less, the signal-to-noise ratio can be large. The relative importance of tropical versus extratropical SST forcings has been established through numerical experiments with atmospheric general circulation models (AGCMs). Predictability of anomalous circulation and associated surface temperature and precipitation in the extratropics is somewhat limited by the difficulty of finding a modest signal embedded in the high level of noise from natural variability in the extratropics, and the complexity and variety of the possible feedbacks. Accordingly, ensembles of AGCM runs and time averaging are needed to identify signals and make predictions. Strong anomalous tropical forcing provides opportunities for skillful forecasts, and the accuracy and usefulness of forecasts is expected to improve as the ability to forecast the anomalous SSTs improves, as models improve, and as the information available from the mean and the spread of ensemble forecasts is better utilized.

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1. Introduction

The Tropical Oceans-Global Atmosphere (TOGA) program focused on the coupling of the tropical oceans with the atmosphere; in particular, it focused regionally on the Pacific and the El Niño-Southern Oscillation (ENSO) phenomenon. The manifestations of the changes in atmospheric circulation in the tropics are felt throughout the global atmosphere via teleconnections and large-scale monsoonal overturning. In turn, the changes in extratropical atmospheric circulation can influence the oceans and force an identifiable signal there which is much stronger than a connection through the ocean itself. The other review papers in this issue deal mainly with the tropical aspects of TOGA; this review will focus on the global atmosphere, especially the extratropical circulation and the linkages with the tropics.

In the buildup to the TOGA program a number of studies had documented empirical connections between the tropics and the atmospheric circulation at higher latitudes (as reviewed in section 2), and theoretical breakthroughs (section 3) had provided a basis for beginning to understand these relationships. The result was that perturbations in the pattern of atmospheric heating over the equatorial Pacific were thought to clearly affect the planetary wave structure over much of the North Pacific Ocean, North America, and probably other parts of the globe [*National Academy of Sciences*, 1983]. Moreover, associated with the atmospheric circulation changes were seasonal changes in surface temperatures and precipitation, so there were prospects for seasonal forecasts that might have great utility.

Prior to TOGA, there had been a considerable debate about the relative role of extra-tropical sea surface temperatures (SSTs) and tropical SSTs in forcing teleconnections in the atmosphere [*Namias*, 1963, 1969; *Bjerknes*, 1969]. However, the TOGA program focused on the tropics, and it was recognized that the way in which the tropical SST changes influenced the atmosphere was through surface fluxes of moisture, heat, and momentum and a readjustment of the tropical circulation in a thermally direct sense. Thus the distribution of deep convection is altered along with associated changes in heating, low level convergence, and upper level divergence, thereby altering the generation of the horizontal component of atmospheric vorticity and the forcing of large-scale atmospheric Rossby waves which could propagate into higher latitudes. In particular, the deepening of the Aleutian low-pressure system in the North Pacific in winter in association with El Niño events was noted by *Bjerknes* [1969]. This feature is now recognized to be one lobe of the Pacific-North American (PNA) teleconnection pattern. Moreover, it was also recognized that such a pattern was not always present during the warm phase of ENSO, and it was a topic for TOGA to understand why. In addition, the evolution of these pat-

terns and their interactions with transient phenomena such as storm tracks were identified as vital pieces of the puzzle not yet adequately described or understood.

Global atmospheric general circulation models (AGCMs) had been used to attempt to simulate the extratropical response to tropical SST anomalies (section 4) with some success but also with some puzzling results, which indicated that the extratropical response was not explainable by simple wave propagation but involved interactions with the extratropical stationary planetary waves and perhaps "modal" behavior. This excitation of possible preferred modes of the extratropical circulation is one aspect that has been studied during TOGA.

The early promise of major breakthroughs in understanding, modeling, and predicting extratropical responses to tropical forcing has been slow in coming, and recent results have tempered future expectations with the realization that potentially predictable tropical influences must compete with extratropical chaotic weather that is essentially unpredictable beyond 2 weeks or so.

Section 2 describes observed teleconnections and especially the linkages between the tropics and extratropics and how these are manifested in terms of changes in storm tracks, temperatures, and precipitation. The 1986-1987 El Niño case is used as a typical example of many aspects. The interactions between the tropics and the extratropics are discussed in detail in section 3 with a focus on theoretical understanding and planetary wave modeling, including recent advances in storm track modeling. Extratropical influences on the tropical atmosphere are also mentioned.

Section 4 describes the state of the art in AGCM modeling as it pertains to the TOGA problem, with special attention given to model runs forced with observed SSTs, and the relative importance of SSTs in different parts of the world. It also addresses the use of models for predictions and how to deal with the extratropical chaotic dynamics by making extensive use of ensemble averaging to distinguish predictable signals from the noise.

The extratropical oceans also have identifiable signals associated with both phases of ENSO, in part through direct propagation of Kelvin waves along eastern boundaries of oceans such as the Americas in the Pacific but also as a result of changes in the extratropical atmosphere and surface fluxes, and these aspects are discussed in section 5 along with possible feedback effects over land. Unsolved problems and issues for the future are addressed in section 6.

2. Observed Relationships

2.1. Teleconnections

The term "teleconnections" does not appear in the glossary [*Huschke*, 1959] although it dates back to *Angstroem* [1935]. It has been in common use by operational long-range forecasters, at least those in the United States, since the 1950s. It has become common

for contemporary climate scientists to use correlation or teleconnection patterns to describe relationships in the variability of large-scale features of the atmospheric circulation as well as tropical and extratropical precipitation and temperature relationships, especially those related to the ENSO. Implied in the term teleconnection is that there is a physical reason for the simultaneous variations, often of opposite sign, over distant parts of the globe, and it now appears that in the extratropics the primary reason is the presence of Rossby waves, whose theory is discussed in section 3. The TOGA era has seen an increase in the widespread use of the term teleconnections as well as the wide acceptance of correlation and composite analyses in attempts to understand and document modes of interannual climate variability.

Interest in prediction of seasonal surface temperature and precipitation motivated the earliest teleconnection studies [e.g., Walker, 1923, 1924], although not called by that name. Walker's studies were primarily focused on the Indian summer monsoon, and they provided the first evidence of the significant global-scale ENSO-related relationships with temperature and precipitation. After a long hiatus in teleconnection research these ENSO teleconnections patterns were later confirmed.

2.1.1. Circulation. Several studies in the early 1980s set the pre-TOGA stage for the use of teleconnection patterns to describe northern hemisphere circulation features and their relationships to tropical anomalies. The North Pacific (NP) oscillation and North Atlantic Oscillation (NAO), which had a long history in the literature [e.g., Walker and Bliss, 1932], were examined closely by van Loon and Rogers [1978] and Rogers and van Loon [1979]. Wallace and Gutzler [1981] reviewed the teleconnection studies up to that time and extended them through analysis of 500 mbar and sea level pressure data for the northern hemisphere winter for timescales of a month and longer. They identified a number of major teleconnection patterns including a pattern they recognized as the now well-known PNA (see Figure 1). They noted for the sea level pressure teleconnections that the anomalies of opposite sign were predominantly between temperate and higher latitudes (e.g., the NP and NAO), while at 500 mbar the patterns were more regional in scale and more wavelike in appearance, and had equivalent barotropic vertical structure.

All these patterns were recognized as preferred modes of variability of the atmosphere, but only some have been clearly identified as also being associated with forcing from SSTs. Horel and Wallace [1981] documented the link between the equatorial SST anomalies and the PNA. Van Loon and Madden [1981] presented global relationships of the Southern Oscillation (SO) in the northern winter with sea level pressures and temperatures, while Trenberth and Paolino [1981] explored northern hemisphere modes of variability in sea level pressure and relationships with the SO in all four seasons using data beginning in 1899. Thus, by the be-

ginning of the TOGA era, much of the foundation had been laid for more detailed teleconnection studies.

TOGA diagnostic research has fostered the introduction and growth of a number of analysis tools previously not widely applied in the meteorological literature. These include the large array of techniques that fall under the general heading of eigenvalue analysis and are generally identified in climate studies as principal component (PC), empirical orthogonal function (EOF), and singular value decomposition (SVD) analyses. Horel [1981] expanded the concept of teleconnection from that given by a simple contemporaneous correlation between a base point and values at every data point in the analysis domain to the use of rotated principal component analysis (RPCA) to identify principal height patterns. During the TOGA era, Barnston and Livezey [1987] exploited RPCA analysis to document the annual cycle of monthly 700 mbar height field patterns as well as longer-term secular variability with a 35 year data set and provided the impetus for operational monitoring (e.g., the National Oceanic and Atmospheric Administration's (NOAA) *Climate Diagnostics Bulletin*) of principal circulation patterns [Bell and Halpert, 1995].

At least 13 distinct teleconnection patterns can be identified in the northern hemisphere extratropics throughout the year, and many of these patterns have appeared previously in the meteorological literature [Bell and Halpert, 1995]. Only some of these patterns and a portion of their variability arises from SST variability. The nomenclature of both Wallace and Gutzler [1981] and Barnston and Livezey [1987] is used. Six prominent patterns are found over the North Pacific-North American sector: the West Pacific pattern (WP), which exists in all months; the East Pacific pattern (EP), which exists in all months except August and September; the NP pattern, which exists from March to July; the PNA, which exists in all months except June and July; the Tropical/Northern Hemisphere pattern (TNH), which exists from November to January; and the Pacific Transition pattern, which exists from May to August. This latter pattern is somewhat similar to another teleconnection pattern that is important in the northern summer and variously known as the northern hemisphere summer pattern, Pacific-Japan (PJ) pattern (or earlier, the south Japan pattern) [Gambo and Kudo, 1983; Nitta, 1986, 1987], and Asia-North America (ANA) pattern [Lau, 1992] that features a wave train sequence of at least four centers stretching from Japan along about 45°N to North America. The PJ pattern is especially important in the far-western Pacific primarily as a north-south dipole pattern.

Most teleconnections patterns associated with changes in SSTs in the tropical Pacific over the northern hemisphere and associated times series for 1963 to 1995 are shown in Figure 1. On the basis of RPCA, shown are the PNA, WP, NP, and the TNH patterns as seen at 700 mbar. For PNA, WP, and TNH, the patterns and time series are for December-January, while NP applies to March-April-May [after Bell and Halpert, 1995]. The

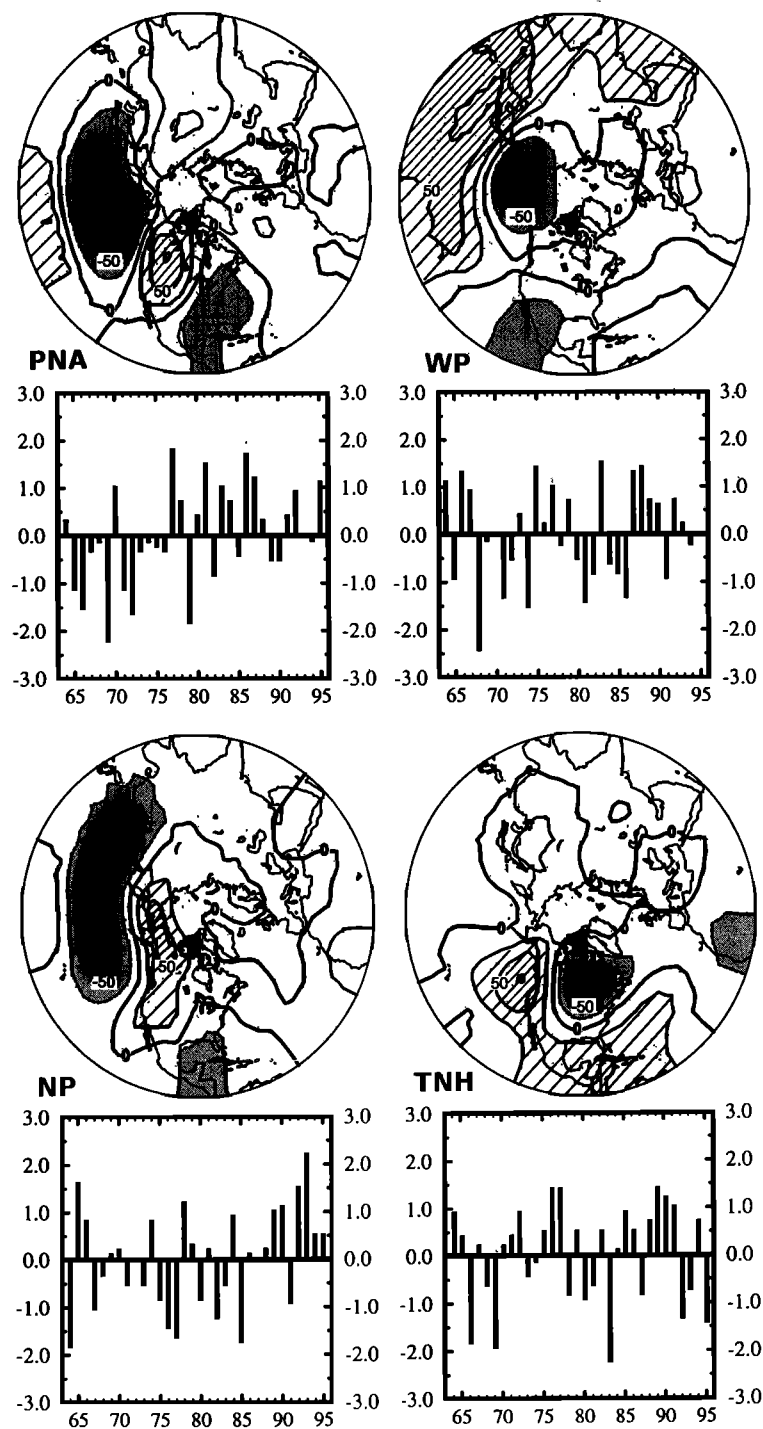


Figure 1. Teleconnections patterns associated with changes in sea surface temperatures (SST) in the tropical Pacific over the northern hemisphere and associated times series for 1963 to 1995. On the basis of rotated principal component analysis (RPCA), shown are the Pacific-North American (PNA), Western Pacific (WP), North Pacific (NP), and the Tropical-Northern Hemisphere (TNH) patterns as seen at 700 mbar. For PNA, WP and TNH, the patterns and time series are for December-January, while NP applies to March-April-May. Positive values are hatched, and negative values are stippled, with the contour interval of 25 m. All time series are standardized to have mean zero and standard deviation of 1. After *Bell and Halpert [1995]*.

PNA pattern stands out because of the prominence of its four centers of alternating sign. It arches from the tropical Pacific across North America and is a good example of a wave train. It is one of the most prominent of all teleconnection patterns in the winter half year. The PNA tends to be positive during El Niño events, and the time series also reflects the trend for more positive values in recent years, following the SST trends in the central equatorial Pacific [Trenberth and Hoar, 1996]. Both WP and NP patterns consist primarily of two centers as a dipole oriented north-south in the western North Pacific but with the centers in each pattern at differing latitudes. The WP pattern reflects latitudinal variations in the subtropical jet stream. Its most persistent positive phase was from September 1986 to June 1987 when there was an El Niño event, and the most persistent negative phase was from March to December 1989 during the La Niña cold event. The NP pattern has some evidence of a third center over the Gulf of Mexico and resembles some aspects of the PNA pattern. It seems to be a preferred response to El Niño during the northern spring, most notably during 1992 and 1993 when it was strongly positive. The TNH pattern really consists of three centers of action: one off the west coast of North America, one near the Great Lakes, and one to the southeast of the United States. Pronounced negative phases of the TNH accompany El Niño events, as seen, for example, in 1982-1983, especially in December-January, and an example of the reverse situation of positive values began in the cold event of 1988-1989. The patterns in Figure 1 are found to project well on the observed atmospheric circulation associated with both the cold and warm phases of ENSO, and they account for substantial percentages of the variance during the seasons indicated. Nevertheless, variations in the teleconnections can also arise from sources other than SSTs. Further details and background are given by Bell and Halpert [1995].

In addition, the PJ pattern and associated NH summer pattern (not shown) appear to be generated by strong convective activity over the Philippine Sea area associated with warmer SSTs there, which often occur during La Niña conditions [Gambo and Kudo, 1983; Nitta, 1987]. Grimm and Silva Dias [1995b] and Lau [1992] suggest that this wave train pattern can be excited by tropical heating almost anywhere in the Pacific, so that it may be a preferred modal response in summer.

During the TOGA era the frequency dependence of the structural and propagation characteristics of various extratropical teleconnection patterns has been documented in considerable detail through analyses of time-filtered data sets. Esbensen [1984] reported that a majority of the teleconnection patterns are prominent in either the intermonthly or the interannual band, but not both. The PNA pattern is among the few exceptions that are well defined within the broadband ranging from a month to several years. Kushnir and Wallace [1989] further noted that the PNA pattern stands out most clearly above the myriad anomaly structures on interannual timescales. Blackmon et al. [1984a, 1984b] and

Schubert [1986] reported that teleconnection phenomena could be grouped into three broad categories. The behavior of fluctuations with timescales much longer than a month is similar to that described by Wallace and Gutzler [1981]. The fluctuations exhibit a notable geographical dependence, with zonally elongated geopotential height anomalies organized about north-south oriented dipoles and with variations in the opposing poles being almost exactly 180° out of phase with each other. Such dipolar structures are prevalent near the jet exit regions over the oceans (e.g., the WP and NP patterns in Figure 1). Perturbations with intermediate timescales (10-30 days) are less anchored to the underlying geography. These patterns appear as eastward and equatorward oriented wave trains, and they are particularly active near the jet entrance regions over the continents. In the southern hemisphere, wave trains on this timescale in both summer and winter tend to be more zonally oriented and exhibit more phase propagation of individual nodes and antinodes than in the northern hemisphere [Kidson, 1991; Hoskins and Ambrizzi, 1993; Ambrizzi et al., 1995]. Variations with still shorter timescales (several days) are characterized by a distinct baroclinic structure, with successive pressure ridges and troughs migrating eastward through continuous phase propagation. These high-frequency features are prominent over the midlatitude oceans and exhibit strong relationships with the more slowly varying teleconnection patterns (see section 2.3 for details). Lau and Nath [1987] demonstrated that the observed properties of the above three classes of teleconnection patterns are reproducible in an AGCM.

In addition to the local forcing of teleconnections, the enhanced tropical heating during an El Niño event also affects the zonal mean circulation in the tropics and subtropics [van Loon and Rogers, 1981]. There is a zonal mean temperature increase in the middle to upper tropical troposphere during an El Niño event associated with latent heat release in the region of enhanced convection and adiabatic heating in the descending branches of the anomalous Walker and Hadley circulations [Horel and Wallace, 1981; Pan and Oort, 1983], although anomalous water vapor amounts, cloudiness, and radiative and sensible heating presumably also play a role. Consequently, the zonal mean flow exhibits an enhanced Hadley circulation together with increased equatorial easterlies in the tropical upper troposphere and enhanced subtropical westerlies.

Anomalous tropical forcing has generally been strongest in the northern winter, coinciding with the mature stage of the El Niño events, while there has been considerably weaker forcing in the southern winter (see Figure 11 and section 3.7). In the southern hemisphere, observational studies have shown that a meridional teleconnection pattern exists across the South Pacific Ocean and South America (sometimes called the Pacific-South American (PSA) pattern, analogous to the PNA pattern in the northern hemisphere) during the warm phase of ENSO in the southern winter, though it appears to be weaker and more variable than the

northern hemisphere (NH) teleconnections [van Loon and Shea, 1985, 1987; Karoly, 1989]. The weaker and more variable response in southern winter is consistent with the relatively weaker tropical forcing, as well as the weaker zonal gradients in the mean flow providing less of a geographical focus for the midlatitude teleconnections (see section 3.3).

2.1.2. Temperature and precipitation. The TOGA era witnessed a significant increase in documentation and understanding of ENSO-precipitation relationships. Consistent warm phase ENSO-precipitation teleconnections were documented by Stoeckenius [1981] Ropelewski and Halpert [1986, 1987], and Lau and Sheu [1988], and later expanded to include precipitation re-

lated to the cold phase of ENSO [Ropelewski and Halpert, 1989] as well as temperature during both phases [Halpert and Ropelewski, 1992]. Correlation, composite, and EOF analyses of global precipitation [e.g., Lau and Sheu, 1988; Kiladis and Diaz, 1989] also helped to define the global ENSO-precipitation patterns. While these studies established the importance of the ENSO-precipitation relationships on a global scale, they also spawned many efforts to document ENSO teleconnections on the local scale. Our current understanding of consistent ENSO relations with precipitation and temperature during the northern winter and summer seasons is summarized in schematic form (Figure 2). To a good approximation, the anomalies during the cold

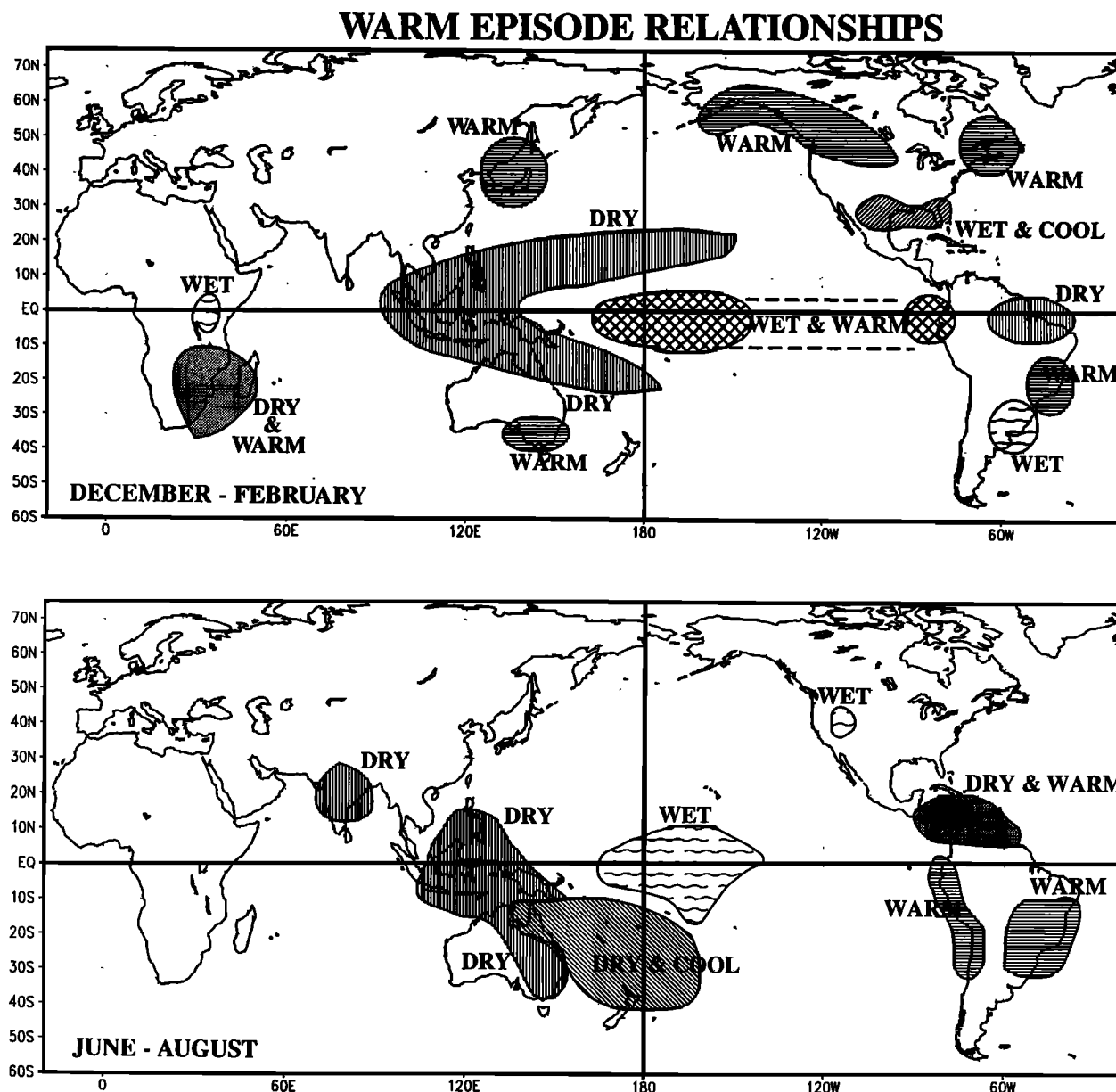


Figure 2. Schematic of temperature and precipitation anomalies generally associated with the warm phase of El Niño-Southern Oscillation (ENSO) during the northern winter and summer seasons. To a good approximation, relationships with the cold phase of ENSO are simply reversed in sign. (After Ropelewski and Halpert [1986, 1987, 1989] and Halpert and Ropelewski [1992] and supplemented by Aceituno [1988].)

phase of ENSO are the reverse of those in Figure 2. This figure is based upon surface measurements and is incomplete over the oceans.

Many of the ENSO-precipitation studies have been limited by the irregular spatial coverage of gauge precipitation data. This limitation motivated the use of satellite observations for the estimation of precipitation especially over the oceans. During the TOGA era, outgoing longwave radiation (OLR), measured by operational NOAA satellites since 1974 [Chelliah and Arkin, 1992], has become the most widely used satellite-based indicator of convective precipitation. OLR measures the outgoing radiation from the Earth and tends to be dominated either by the radiation from the surface, where there are no clouds, or from cloud tops, so that it serves as a proxy index of the amount of convection in the tropics. The relatively long period of record, at least for satellite measurements, has made OLR a mainstay for a multitude of research studies. The continued operation of the weather satellites has also made them a major observational tool for monitoring current climate (e.g., the *Climate Diagnostics Bulletins* issued by NOAA's Climate Prediction Center). While OLR has proven invaluable for several TOGA-related studies, its limitations have also spurred efforts to provide precipitation estimates with higher spatial and temporal coverage as well as to provide better estimates for nonconvective precipitation. Many of these activities are taking place under the World Climate Research Program's Global Precipitation Climatology Project [Arkin and Xie, 1994].

The above satellite observations reveal that in the Pacific the Intertropical Convergence Zone (ITCZ) is typically stronger and located closer to the equator in El Niño years, while the South Pacific Convergence Zone (SPCZ) shifts north and east of normal. During the ENSO warm phase in the northern winter for December-January-February (DJF), Figure 2 reveals the warm and wet conditions that typically occur throughout the tropical central and eastern Pacific and along the Pacific coast of South America. Dry conditions, corresponding to higher sea level pressures, exist in a characteristic boomerang-shaped pattern over the western Pacific and extend to southern Africa and across the Atlantic to northeast Brazil and Columbia. Cool and wet conditions in the southeastern United States are linked via a teleconnection to the warm conditions over western Canada and Alaska. Warmth in Uruguay expands in area in southern winter in June-July-August (JJA). Dryness in parts of Southeast Asia, Indonesia, Central America, Australia, and New Zealand in JJA contrasts with wet conditions in the tropical central Pacific.

2.1.3. An example: The El Niño of 1986-1987. To illustrate the changes in circulation in the extratropics and the links with the tropics, the example of the 1986-1987 El Niño event is used. This was a modest event by most standards, and the warmest water near the equator was displaced eastward to near 170°W. Even though the SST anomalies there were only

1°C or so, it was sufficient to produce ensuing strong convection in that area [Trenberth, 1996] (see Figure 3). We use the National Centers for Environmental Prediction (NCEP, formerly National Meteorological Center (NMC)) NCAR reanalyses for 1985 to 1993 (9 years) to define a base period to determine circulation anomalies for this event, and the same base period is used for OLR and SST.

In Figure 3, for DJF 1986-1987 the anomalous SST, OLR, and sea level pressure fields are shown, along with corresponding anomalies at 200 mbar in the divergent wind component, velocity potential field, and the streamfunction field, and the geopotential height field at 300 mbar. In the SST field, besides the warming in the equatorial Pacific, features of note are the cooling in the North Pacific and around New Zealand, with the latter especially pronounced (see section 5.1). The OLR anomalies are as large as -50 W m^{-2} , and experience indicates that magnitudes exceeding about 10 W m^{-2} for seasonal means (e.g., see Figure 11) are nonrandom and linked to surface forcing. Figure 3 indicates the above normal convective activity centered on the equator at 170°W over the warmest water but with evidence of reduced activity to the north in the vicinity of the Hawaiian Islands, to the west over Indonesia, and to the southwest, off Australia. Note also the eastward shifted SPCZ. Regions of low level convergence over the warmest water are mirrored in the upper troposphere by regions of divergence, as seen by the anomalous outflow at 200 mbar in Figure 3. The divergent component of the flow provides a forcing for Rossby waves in the atmosphere through the advection of the Earth's vorticity by the anomalous divergent flow, giving an anticyclonic forcing in the upper troposphere (compare Figure 7, presented later, and Rasmusson and Mo [1993]). Note also the upper tropospheric convergence and thus suppressed rainfall over northeast Brazil. All the main tropical and subtropical features in the OLR field have corresponding upper tropospheric divergence anomalies.

The nondivergent component of the flow is revealed by the streamfunction (ψ) anomalies in Figure 3. Most studies of the atmospheric circulation, especially before TOGA, used geopotential height (z) as the primary variable. However, because small gradients in height can be associated with strong winds in low latitudes, it has been recognized that the streamfunction is more appropriate for revealing global aspects of the flow. Figure 3 presents the 300 mbar z as well as the ψ field (although at 200 mbar) to illustrate the difference in perception that results. The z variations fail to reveal, the tropical circulation changes that often lead to the extratropical wave trains. Together with the sea level pressure field, the dominant equivalent barotropic nature of the extratropical anomalies is revealed, and the anomalies in the Pacific-North American region stand out.

The streamfunction anomalies in Figure 3 show the strong anticyclonic couplet in both hemispheres straddling the equator in the region of anomalous convection,

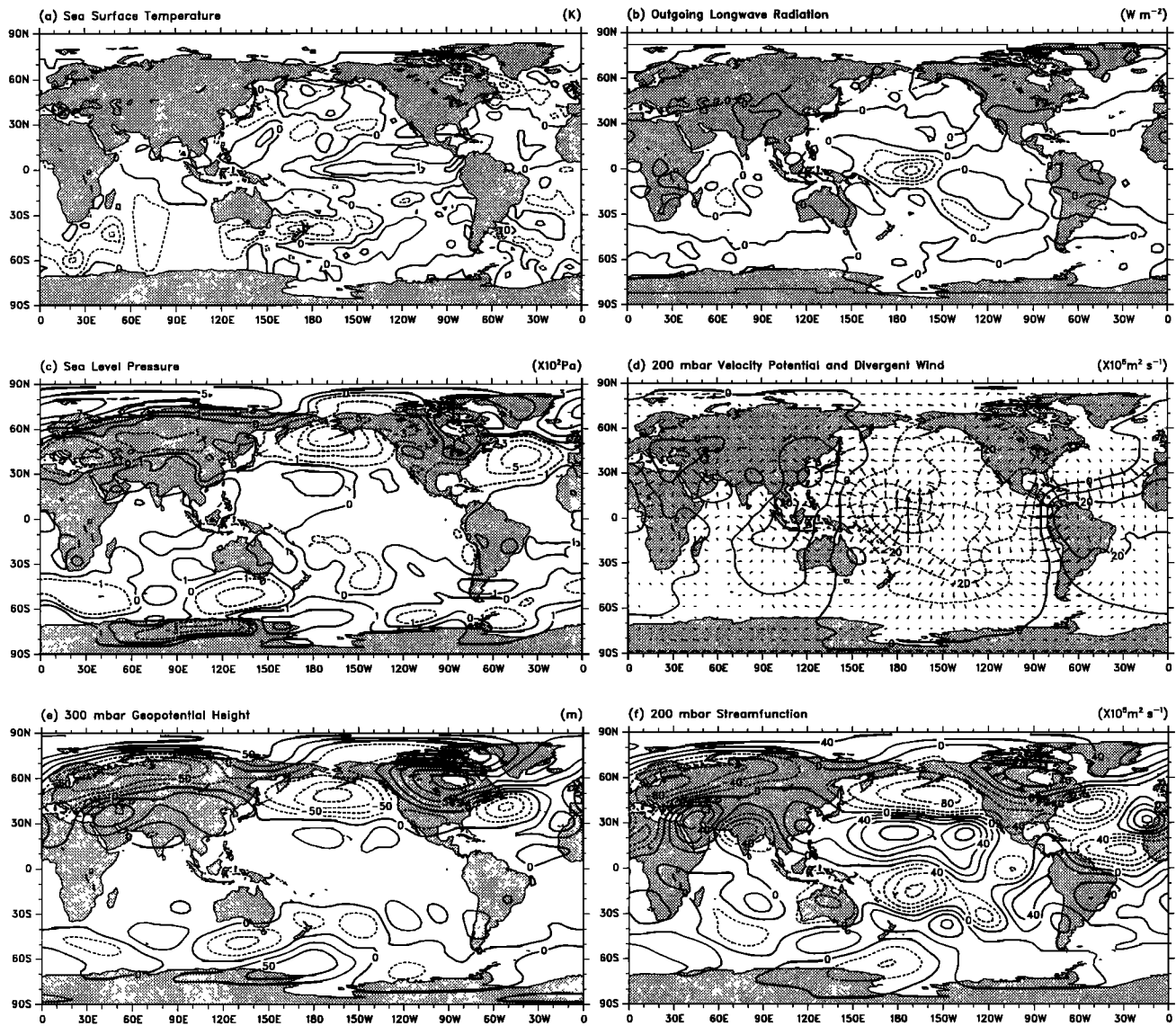


Figure 3. Anomalies for the ENSO event during December-January-February (DJF) 1986-1987, all relative to a base period from 1985 to 1993 as (a) SST, contour interval 0.5°C and (b) outgoing longwave radiation (OLR), contour interval 10 W m^{-2} , and based on NCEP/NCAR reanalyses of (c) sea level pressure, contours are $0, \pm 1, \pm 3\text{ mbar}$, etc.; (d) velocity potential and divergent component of the wind vectors at 200 mbar, with the longest wind vector shown equal to 2.1 m s^{-1} ; (e) 300 mbar geopotential height, contour interval 25 m ; and (f) streamfunction at 200 mbar. The contour intervals for velocity potential and streamfunction are 10 and $20 \times 10^5\text{ m}^2\text{ s}^{-1}$, respectively.

and teleconnections into higher latitudes. The PNA teleconnection pattern is strongly evident across North America, and anticyclonic conditions prevail over Australia, but cyclonic flow prevails over the New Zealand region. Breaking down the pattern in Figure 3 into the teleconnection patterns of Figure 1 reveals standardized index values of the PNA of $+1.2$, TNH -0.9 and WP $+1.3$, so that all contribute substantially. Associated with the changes in quasi-stationary waves are changes in storm tracks, as detailed in section 2.3 (see Figure 5, presented later). All these features are fairly typical of

most El Niño events and give rise to the precipitation and temperature anomalies in Figure 2; many processes play a role in their formation (section 3).

A schematic of the tropical forcing and dominant northern hemisphere atmospheric response in the form of a wave train forced from the tropical upper tropospheric divergence and subtropical convergence is shown in Figure 4. Aspects of this simplified picture can be seen in Figure 3; however, it is further discussed, the physical relationships are revealed, and the features are depicted in subsequent sections.

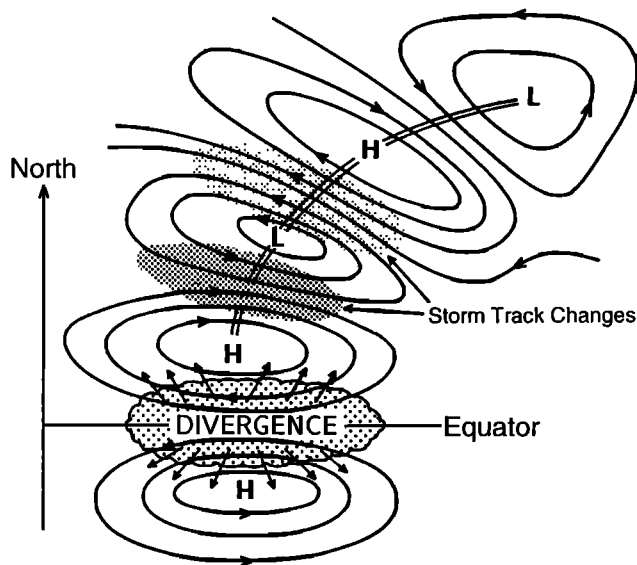


Figure 4. Schematic view of the dominant changes in the upper troposphere, mainly in the northern hemisphere, in response to increases in SSTs, enhanced convection, and anomalous upper tropospheric divergence in the vicinity of the equator (scalloped region). Anomalous outflow into each hemisphere results in subtropical convergence and an anomalous anticyclone pair straddling the equator, as indicated by the streamlines. A wave train of alternating high and low geopotential and streamfunction anomalies results from the quasi-stationary Rossby wave response (linked by the double line). In turn, this typically produces a southward shift in the storm track associated with the subtropical jet stream, leading to enhanced storm track activity to the south (dark stipple) and diminished activity to the north (light stipple) of the first cyclonic center. Corresponding changes may occur in the southern hemisphere.

2.2. Decadal Variability Throughout the Pacific

An example of an important teleconnection emerging from the tropical Pacific on longer timescales has been documented in the North Pacific in winter, with the period of the fluctuations exceeding 20 years. In particular, a decade-long change in the North Pacific atmosphere and ocean beginning around 1976 and lasting until at least 1988 has been noted by Trenberth [1990]. A comprehensive review of many related aspects, including linkages to changes in the tropical Pacific SSTs which generally increased after 1976, has been given by Trenberth and Hurrell [1994].

Observed significant changes in the atmospheric circulation for the 1976-1988 period involve the PNA teleconnection pattern (Figure 1), so that there was a deeper and eastward shifted Aleutian low pressure system (positive PNA) in the winter half year which advected warmer and moister air along the west coast of North America and into Alaska and colder air over the North Pacific. Consequently, there were increases in temperatures and SSTs along the west coast of North America and Alaska but decreases in SSTs over the central North Pacific [Tanimoto *et al.*, 1993; Kawamura, 1994], as well as changes in coastal rainfall and streamflow and decreases in sea ice in the Bering Sea. Associated changes occurred in the surface wind stress and, by inference, in the Sverdrup transport in the North Pacific Ocean and had a distinctive signature throughout the upper ocean to greater than 400 m depth [Deser *et al.*, 1996]. Changes in the mean flow were accompanied by a southward shift in the storm tracks and associated synoptic eddy activity [Trenberth and Hurrell, 1994] and in the surface ocean sensible and latent heat fluxes. The deeper Aleutian low and associated changes

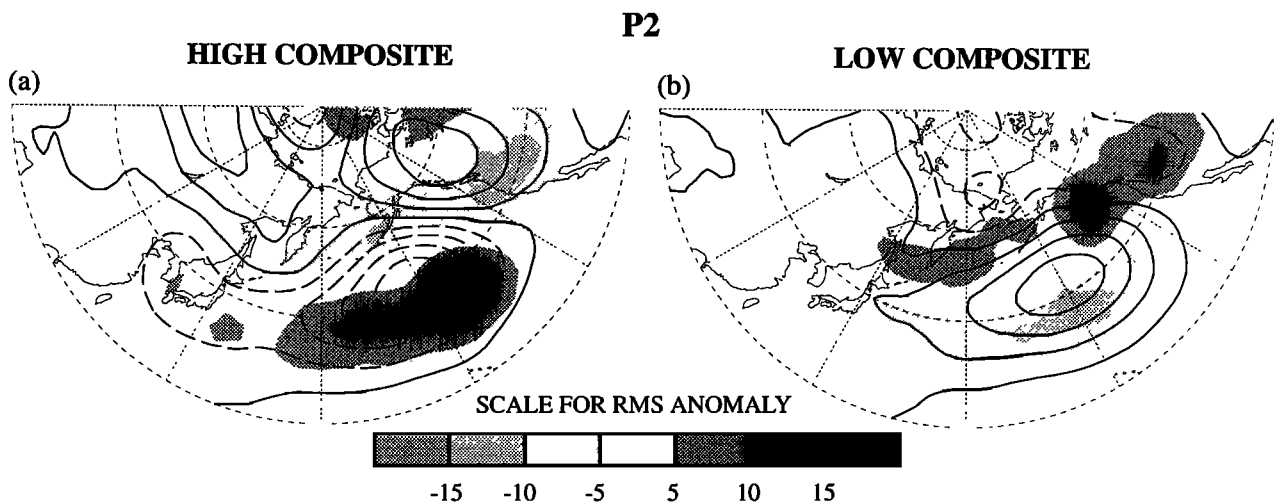


Figure 5. Composites of observed anomalies in monthly averaged 500 mbar height (solid contours are positive, dashed contours negative, contour interval 25 m) and in root-mean-squares of 500 mbar height fluctuations with timescales of 2.5-6 days (stippling, see scale bar at bottom). These composites are based on individual groups of 10 winter months corresponding to opposite extremes of selected principal modes of storm track variability in the North Pacific. The specific months in each group are listed in Table 2 of Lau (1988) under the headings of (a) P2 High Composite, and (b) P2 Low Composite.

in upper ocean mixing and Ekman transport (section 5.1) increased the nutrient supply, phytoplankton, and zooplankton in the central North Pacific Ocean. These changes, along with the altered ocean currents and temperatures, shifted in the migration patterns and increased the stock of many fish species (e.g., Alaskan and Canadian salmon) but decreased some other species (e.g., West Coast U.S. salmon, coastal species) [cf. *McFarlane and Beamish*, 1992; *Beamish and Bowillon*, 1993; *Polovina et al.*, 1994].

The dominant atmosphere-ocean relation in the North Pacific is one where atmospheric changes lead SSTs by 1 to 2 months, apparently because of the changes in surface sensible and latent heat fluxes combined with mixing in the ocean and entrainment. A southward shift in the storm tracks in the North Pacific helps to reinforce and maintain the anomalous circulation in the upper troposphere. Observational evidence suggests a link with the tropical Pacific SSTs and, after 1976, there is also evidence for increased tropospheric temperatures and water vapor in the western tropical Pacific. Several aspects of the decadal scale change beginning around 1976 have been simulated with AGCMs using specified SSTs (see section 4.2), which confirm that the atmospheric changes are tied to the changes in SSTs and, further, that the changes over the North Pacific and surrounding areas are substantially controlled by the anomalous SST forcing from the tropical Pacific and thus linked to the changes in ENSO. However, whereas changes in the tropical Pacific have continued in the 1990s, the positive PNA has not been as strongly present since 1989 (e.g., Figure 1), indicating either that there was a change in the nature of the teleconnections or that there are other influences of importance, such as those given in section 3, yet to be clarified.

2.3. Variability of Storm Tracks Associated with Teleconnection Patterns

In the wintertime extratropics, transient phenomena with timescales of several days consist mainly of migratory baroclinic cyclone waves. The relationships between these disturbances and the stationary flow field have been documented by *Blackmon et al.* [1977] and *Trenberth* [1991], among others, using climatological circulation statistics for time-filtered data. It is apparent that the cyclone waves are particularly active in maritime sites located downstream and poleward of the time-averaged jet streams. The preferred paths of the weather-scale fluctuations in these regions of activity are often referred to as the "storm tracks." The studies support the notion that the initiation and subsequent propagation of the cyclone eddies are largely governed by the preferred sites of cyclogenesis and steering effects associated with the stationary time-averaged flow field.

By virtue of the intense transports of heat and vorticity by the transient eddies along the storm tracks, strong dynamical interactions between the time-varying and stationary components of the atmospheric circulation take place in these regions. As pointed out by *Lau and Holopainen* [1984], the long-term-averaged geopo-

tential height tendencies induced by the convergence of eddy vorticity fluxes reinforce the upper tropospheric climatological troughs and jet streams over eastern Asia and North America, whereas the corresponding eddy heat transports dissipate the same stationary features. Because the magnitude of the forcing in the upper troposphere due to vorticity fluxes is larger than that due to heat fluxes, the net effect of the disturbances is to enhance the zonally asymmetric component of the time mean flow. Thus the stationary waves generated by orographic and thermal forcings are reinforced by the transient eddies in the upper troposphere. In the lower troposphere the stationary wave temperature perturbations are maintained by advection by the mean flow but dissipated by the transient eddies. The intense eddy-mean flow interactions occurring in the storm tracks dictate that a full understanding of the steady state atmospheric response to perturbed boundary conditions (such as SST anomalies) must take into account the aggregate effects of the transient waves, including their influence on the surface fluxes and diabatic heating [*Hoskins and Valdes*, 1990].

It has been shown that the above relationships between the storm tracks and the climatologically averaged flow are also applicable to understanding the interactions between the high-frequency transients and more slowly varying circulation anomalies with timescales ranging from subseasonal periods [e.g., *Hoskins et al.*, 1983; *Mullen*, 1987; *Metz*, 1989; *Lau and Nath*, 1991; *Cai and van den Dool*, 1991, 1992], to ENSO cycles [e.g., *Kok and Opsteegh*, 1985; *Held et al.*, 1989; *Hoerling and Ting*, 1994], and to interdecadal changes (see *Trenberth and Hurrell*, [1994], and section 2.2). In particular, some of these studies indicate that the storm track variations accompanying the changes in the background flow exert a positive feedback on the low-frequency component of the circulation. It is also established that changes in the mean flow and the associated jet stream cause corresponding changes in the atmospheric storm tracks (see *Branstator* [1995] and section 3.4.2).

The strong association between storm track variability and Pacific teleconnection patterns (section 2.1) is illustrated in Figure 5, which shows composites of observed anomalies in monthly averaged 500 mbar height and root-mean-squares (rms) of transient (2.5–6 day) fluctuations of 500 mbar height. The rms field presented here is a good indicator of the intensity and location of the storm tracks. The patterns in Figure 5 were constructed by averaging over those 10 winter months corresponding to the two extremes of a selected mode of variability in this rms field in the Pacific. This mode has been identified by *Lau* [1988] using an EOF analysis and is labeled "P2." It is evident (Figure 5) that the monthly mean patterns for P2 correspond to the two opposing polarities of the PNA pattern. The positive height anomaly center over western North America and the enhanced eastward geostrophic flow over the subtropical central Pacific (Figure 5a) are accompanied by an equatorward displacement of the Pacific storm track and diminished cyclone activity along the west-

ern seaboard of North America. Reversal of the polarity in the monthly mean pattern (Figure 5b) occurs in conjunction with a poleward storm track displacement and enhanced activity off the west coast of North America. These relationships are depicted schematically in Figure 4.

In the southern hemisphere, larger interannual variations are found in the zonal mean flow than in the wavelike teleconnections, but there also exists a profound influence of the low-frequency circulation on the storm tracks with similar feedbacks coming into play. The dominant mode is one of switching from a single to a double jet structure [Trenberth and Christy, 1985]. Kidson [1988b] detected variations in the westerlies several days after the occurrence of barotropic forcing due to the convergence of eddy westerly momentum transports. Karoly [1990] noted that the storm track displacements are accompanied by changes in the transient eddy poleward heat flux and thus in the Eliassen-Palm flux. Randel [1989] traced the evolution of the zonal mean flow in response to life cycle dynamics of baroclinic waves. However, Trenberth [1984], Kidson [1988a, 1988b], Karoly [1990], Shiotani [1990], and Cuff and Cai [1995] emphasize that the low-frequency variations in zonal mean flow are primarily equivalent barotropic and dominated by changes in the transient momentum (or equivalently vorticity) fluxes, with the heat flux playing a much smaller role. Both the southern hemisphere jet stream and storm tracks vary considerably from year to year in models without external forcing [Zwiers, 1987; Yu and Hartmann, 1993; Hartmann, 1995], so the effects of eddy-mean flow interactions contribute significantly to internal atmospheric variability, which could mask any influences of tropical SST forcing.

3. Theory and Diagnostics of Interactions of the Tropics with the Extratropics

3.1. Introduction

We now focus on the theoretical basis for the linkages between the anomalous tropical circulation associated with ENSO and changes in the extratropical circulation. Prior to TOGA, the modeling studies of Hoskins *et al.* [1977], Hoskins and Karoly [1981], and Webster [1981] had shown that the extratropical response to large-scale tropical forcing could be understood in terms of anomalous planetary wave propagation from regions of tropical upper tropospheric divergence (Figure 4). The modeled teleconnections from tropical heating were shown by Hoskins and Karoly [1981] and Webster [1981] to be remarkably similar to some of the observed teleconnections described in section 2. We shall refer to this pre-TOGA conceptual model as the "protomodel" of the extratropical response to tropical forcing.

There have been many refinements to the theoretical understanding of these linkages during the TOGA decade, and these are outlined in the remainder of this section. A number of factors are important in deter-

mining the extratropical circulation during an El Niño event, including the location and intensity of the tropical circulation anomalies, the effects of the mean flow on planetary wave propagation and forcing, interactions with midlatitude storm tracks, and interference from the internal chaotic variability of the midlatitude circulation. However, Rossby wave propagation still provides the underpinning for all theories of how the tropics influence midlatitudes. Hence we start with a description of the protomodel, including a brief review of the theory of Rossby wave forcing and propagation as it stood at the start of the TOGA period, then continue with the more recent refinements.

The vorticity equation in pressure coordinates is given by

$$\frac{\partial \zeta}{\partial t} + \mathbf{v} \cdot \nabla(\zeta + f) + \omega \frac{\partial \zeta}{\partial p} = -(\zeta + f) \nabla \cdot \mathbf{v} + \mathbf{k} \cdot \left(\frac{\partial \mathbf{v}}{\partial p} \times \nabla \omega \right) - \mathbf{F} \quad (1)$$

where ζ is the vertical component of relative vorticity, $\mathbf{v} = (u, v)$ is horizontal velocity, ω is vertical velocity, $f = 2\Omega \sin \phi$ is the planetary vorticity, and \mathbf{F} is due to friction. When applied to the time-averaged flow in the upper troposphere, the vorticity equation can be written (using an overbar to represent time-averaged quantities)

$$\frac{\partial \bar{\zeta}}{\partial t} + \bar{\mathbf{v}} \cdot \nabla(\bar{\zeta} + f) = -(\bar{\zeta} + f) \nabla \cdot \bar{\mathbf{v}} - \nabla \cdot (\bar{\mathbf{v}}' \bar{\zeta}') - \bar{\mathbf{F}} \quad (2)$$

neglecting terms involving ω since the vertical velocity is small near the tropopause. This is the equivalent barotropic vorticity equation for the mean vorticity $\bar{\zeta}$ with forcing by the mean stretching of absolute vorticity $(\bar{\zeta} + f) \nabla \cdot \bar{\mathbf{v}}$ or by transient eddy convergence of vorticity $-\nabla \cdot (\bar{\mathbf{v}}' \bar{\zeta}')$.

Hoskins and Karoly [1981] showed that tropical heating is balanced by vertical motion, with large divergence in the tropical upper troposphere, leading to substantial forcing in the barotropic vorticity equation. In the protomodel the effect of this divergence is approximated by solving (2) with specified divergence, with the transient fluxes ignored, and with all other terms linearized about a basic state, independent of time and longitude. The response of this model can be considered in terms of the dispersion of Rossby waves, which are governed by the first two terms of (1). After linearizing (1) about a basic state consisting of a zonal mean zonal flow $[\bar{u}]$, we consider small amplitude perturbations ζ^{\dagger} to the basic state vorticity, giving a wave equation for the perturbation streamfunction ψ^{\dagger} on a β plane

$$\left(\frac{\partial}{\partial t} + [\bar{u}] \frac{\partial}{\partial x} \right) \nabla^2 \psi^{\dagger} + \frac{\partial \psi^{\dagger}}{\partial x} \left(\beta - \frac{\partial^2 [\bar{u}]}{\partial y^2} \right) = \text{forcing} \quad (3)$$

Assuming a separation in scale between the variations in the basic state and the perturbations, a WKB solution to (3) exists of the form $\psi^{\dagger} = A \exp[i(kx + ly - \sigma t)]$, with horizontal wavenumber $\mathbf{k} = (k, l)$ and frequency σ

which satisfies the dispersion equation $\sigma = [\bar{u}]k - (\beta - [\bar{u}]_{yy})k/K^2$, with $K = |k|$. The dispersion of wave energy from the forcing is in the direction of the group velocity

$$c_g = (\sigma/k + 2(\beta - [\bar{u}]_{yy})k^2/K^4, 2(\beta - [\bar{u}]_{yy})kl/K^4). \quad (4)$$

For stationary waves on the sphere in a climatological mean zonal flow, as considered by *Hoskins and Karoly* [1981], wavelike solutions are possible for low-latitude forcing in a westerly mean flow, and wave energy disperses poleward and eastward from the forcing region before arcing back toward low-latitudes and being absorbed (according to linear theory) at the low latitude "critical line" where the mean flow is equal to the phase speed of the waves (zero) (compare Figure 4).

The dispersion of Rossby waves from tropical forcing described by *Hoskins et al.* [1977] and *Hoskins and Karoly* [1981] agreed well with the observed and modeled responses to anomalous tropical heating during El Niño events in a number of ways, particularly for the structure and direction of the teleconnection patterns emanating from the tropical forcing. It helped to explain the stronger teleconnections in the winter rather than the summer hemisphere (as wave propagation is stronger in stronger westerly flow and not possible through equatorial easterlies), the arcing teleconnection patterns, the spacing between successive anomalous highs and lows, and the timescale for the establishment of the teleconnection patterns. Figure 6 shows the steady response in the upper troposphere to tropical heating in NH winter from a linearized primitive equation model, the horizontal flux of stationary wave activity diagnosed from this model [*Karoly et al.*, 1989], and the Rossby wave rays for a similar case from *Hoskins and Karoly* [1981]. This shows clearly all the typical characteristics of the protomodel described above.

There were a number of limitations of this protomodel. Principal drawbacks were (1) the need to have the Rossby wave forcing in mean westerly flow, whereas the observed tropical forcing was in mean easterlies; (2) no explanation for the observed and AGCM-modeled preference for geographically fixed midlatitude anomalous wave trains; and (3) reduced amplitudes in the linear model solutions when forced with realistic tropical divergence anomalies.

Each of these limitations has been addressed in refinements to the theoretical understanding developed over the period of the TOGA decade and is discussed below, together with new processes now recognized to be important. We discuss these processes in the order

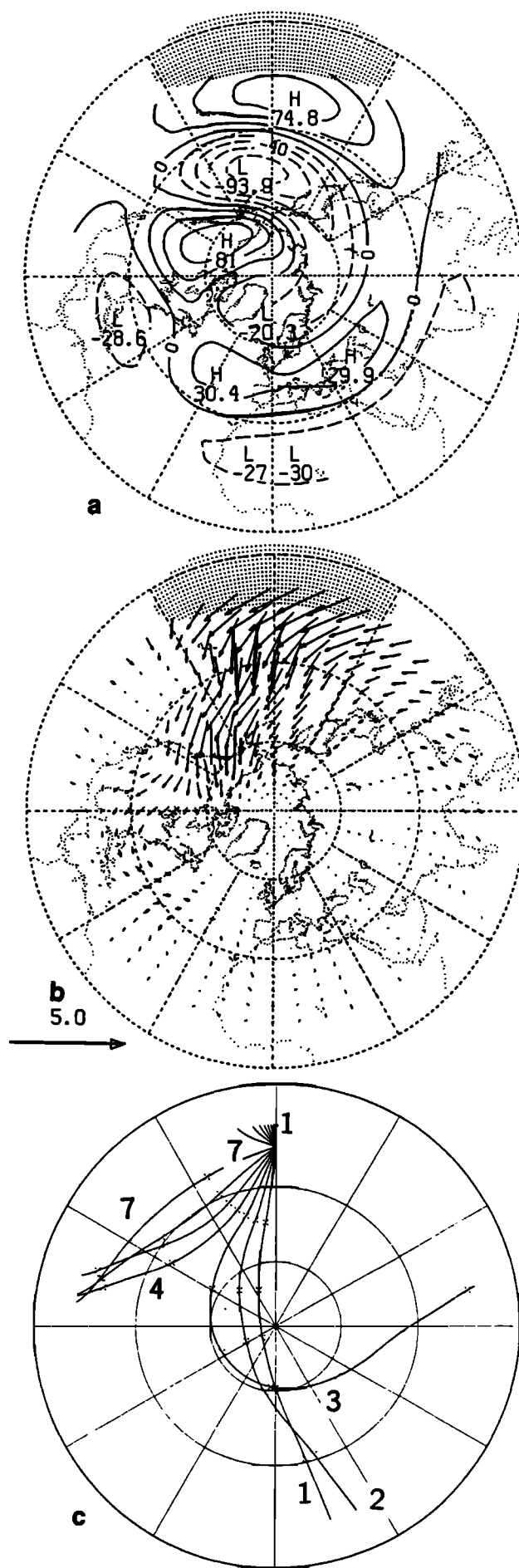


Figure 6. Linearized, steady state primitive equation model solution to thermal forcing in the tropics in northern hemisphere (NH) winter, showing (a) the height field anomalies at 200 mbar, contour interval 20 m, (b) the associated horizontal flux of stationary wave activity [from *Karoly et al.*, 1989], and (c) the barotropic Rossby wave rays for tropical forcing at 15°N in a 300 mbar NH winter basic state; crosses indicate phases every 180° [from *Hoskins and Karoly*, 1981].

in which they occur in nature, starting with the tropical forcing, then continuing with propagation into midlatitudes, and internal midlatitude forcing and other issues.

3.2. Tropical Forcing

The focus of the ENSO phenomenon is the tropical Pacific, and the observed global influence arises from atmospheric teleconnections from the regions of anomalous tropical heating. Small changes in SST and SST gradients can lead to shifts in the location of the large-scale organized convection in the tropics and also to changes in the intensity of the convection. These result in large anomalies in atmospheric heating, mainly through latent heat release in precipitation, and upper tropospheric divergence (Figure 3). The tropical response to this anomalous heating can be explained in terms of equatorially trapped Kelvin and mixed Rossby-gravity waves with a first internal mode vertical structure in the troposphere [Gill, 1980]. In the vicinity of the tropical heating the response in the upper troposphere takes the form of a pair of anomalous anticyclones which straddle the equator (compare Figures 3 and 4). In addition to the anomalous divergence over the tropical heating, there is usually an anomalous circulation involving an east-west Walker circulation linking a region of anomalous convergence over the Indonesian region to the anomalous divergence (see also section 2.1.3).

Hence one factor which directly influences the tropical forcing during ENSO events is the differences of SST anomaly patterns between events. These lead to differences in the location of the tropical heating anomalies, apparent in the anomalous tropical rainfall and associated OLR fields observed during different events. It is not clear how much the changes in the location and magnitude of the tropical heating anomalies account for some of the variations of the extratropical circulation between ENSO events.

The forcing of extratropical Rossby waves by the anomalous tropical heating can be best understood by partitioning the horizontal flow \mathbf{v} into nondivergent \mathbf{v}_ψ and irrotational \mathbf{v}_χ components, such that $\mathbf{v} = \mathbf{v}_\psi + \mathbf{v}_\chi = \mathbf{k} \times \nabla\psi + \nabla\chi$, with streamfunction ψ and velocity potential χ . Then (2) can be rewritten as

$$\begin{aligned} \frac{\partial \bar{\zeta}}{\partial t} + \bar{\mathbf{v}}_\psi \cdot \nabla(\bar{\zeta} + f) = & -(\bar{\zeta} + f)\nabla \cdot \bar{\mathbf{v}}_\chi \\ & - \bar{\mathbf{v}}_\chi \cdot \nabla(\bar{\zeta} + f) - \nabla \cdot (\bar{\mathbf{v}}'_\chi \bar{\zeta}') \\ & - \nabla \cdot (\bar{\mathbf{v}}'_\psi \bar{\zeta}') - \bar{\mathbf{F}}. \end{aligned} \quad (5)$$

The left-hand side retains all terms necessary to support Rossby wave propagation, while the right-hand side involves an additional term compared with (2) involving the advection of vorticity by the divergent flow. This form of the barotropic vorticity equation was discussed by Sardeshmukh and Hoskins [1988], who described the first three terms on the right-hand side as the Rossby wave source, $\bar{S} = -\nabla \cdot \bar{\mathbf{v}}_\chi(\bar{\zeta} + f)$, which is the forcing associated with the divergent flow. The anomalous

Rossby wave source term, which contributes to the forcing of seasonal anomalies, can be written as

$$\begin{aligned} -S_a = fD_a + \beta \mathbf{v}_{\chi a} \cdot \zeta_c D_a + \zeta_c D_a + \zeta_a D_c + \zeta_a D_a \\ + \mathbf{v}_{\chi c} \cdot \nabla \zeta_a + \mathbf{v}_{\chi a} \cdot \nabla \zeta_c + \mathbf{v}_{\chi a} \cdot \nabla \zeta_a \\ + \overline{(\nabla \cdot (\mathbf{v}'_\chi \zeta'))}_a \end{aligned} \quad (6)$$

where the overbar represents a single season mean, the primes indicate the daily departures from the single season mean, subscript represents the climatological mean value, and subscript represents the departure of the single season mean from the climatological value, and D is the divergence.

Rasmusson and Mo [1993] used a diagnostic analysis to identify the forcing of midlatitude Rossby waves in the upper troposphere during the 1986–1989 ENSO warm event and cold event sequence. Fields from their analysis are similar to those for the 1986–1987 event alone (Figure 3). They found that the anomalous Rossby wave source (see Figure 7) is small within 15° latitude of the equator, despite large upper tropospheric divergence, as the absolute vorticity is small in this region. Poleward of 15°, the strongest Rossby wave source anomaly results from the anomalous subtropical convergence associated with the descending branch of the local Hadley circulation from the tropical heating, $-fD_a$. This provides an effective Rossby wave forcing, as it occurs in mean westerly flow and, as the experiments of Held and Kang [1987] show, makes a major contribution to the midlatitude wave train. In addition, vorticity advection by the anomalous divergent outflow $\mathbf{v}_{\chi a} \cdot \nabla(\zeta_c + f)$ can be an important contributor to the Rossby wave forcing [Sardeshmukh and Hoskins, 1988]. Divergent flow anomalies in regions of strong mean vorticity gradients, such as the winter subtropical jets, can lead to large Rossby wave forcing anomalies, and the location of these sources can be somewhat insensitive to the position of the heating that induces them.

Although the upper level tropical divergence is relatively ineffective in directly forcing midlatitude Rossby waves, the divergent outflow and the subtropical convergence associated with the meridional overturning in local Hadley circulations away from the anomalous tropical heating are both effective in forcing Rossby waves in the subtropics in westerly mean flow, which can then propagate to higher latitudes (Figure 4). Hence factors which determine the locations of the upper tropospheric subtropical convergence and the subtropical vorticity gradients will greatly influence the location of the subtropical Rossby wave forcing.

3.3. Energy Propagation

From the discussion in the previous subsection it is apparent that wave propagation is not the only means by which the influence of tropical heating reaches beyond the source region, but wave propagation is still recognized as the primary mechanism by which tropical heating influences regions thousands of kilometers away. As seen in section 3.1, the basic theory of large-scale energy propagation can be understood from a very simple

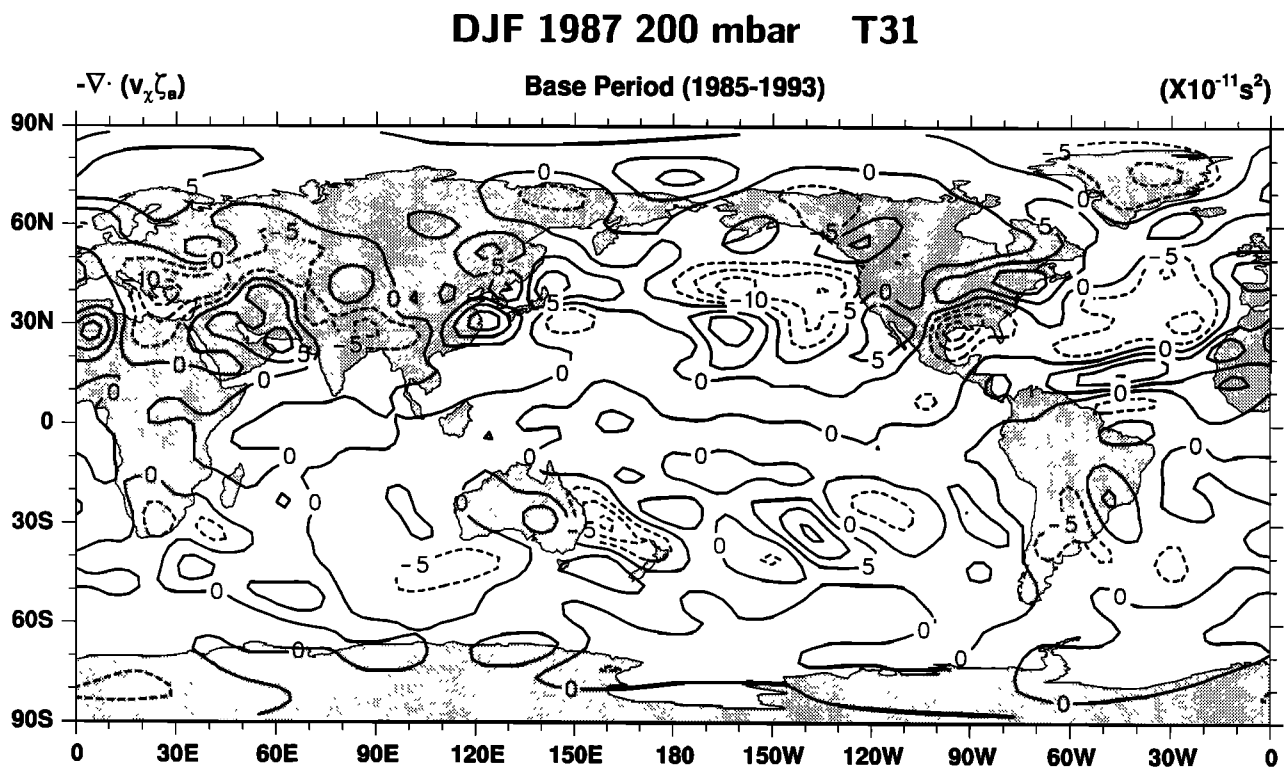


Figure 7. Rossby wave source anomalies for the ENSO event during DJF 1986-1987, relative to a base period from 1985 to 1993 and based on NCEP reanalyses. Contour interval is $5 \times 10^{-11} \text{ s}^{-1}$, and negative values are dashed.

dynamical model. By relaxing some of the simplifying assumptions of that protomodel, a more complete picture of how and where energy will propagate in a given situation can be constructed.

One example of this is consideration of the influence of the time average divergent circulation on the transmission properties of the atmosphere. As pointed out by *Schneider and Watterson* [1984], steady heat sources, which the simple theory suggests as being too deep in the tropics to affect midlatitudes in the winter hemisphere because of intervening upper tropospheric easterlies, will influence midlatitudes if the effects of the climatological zonal mean poleward tropical flow are taken into account. Furthermore, the mean convergence associated with these poleward winds tends to produce a subtropical source of vorticity in the winter subtropical upper troposphere through its stretching effect on perturbation vorticity, thus making the influence of the heating in the winter hemisphere even stronger [*Sardeshmukh and Hoskins*, 1988]. *Grimm and Silva Dias* [1995a, b] confirm the importance of including the effects of the divergent part of the mean flow in computing the response to forcing in a barotropic model.

Another example of added realism comes from accounting for the timescale of the tropical heat sources. Though often approximated by steady sources, the heating associated with El Niño consists of episodes of varying duration, and this variability affects the path taken by energy propagating away from the source regions.

The analyses of *Karoly* [1983] and *Li and Nathan* [1994] and the dependence of the group velocity on wave frequency in (4) show that the shorter the timescale, the more equatorially trapped is the response. So even if the heating remains in exactly the same location throughout an El Niño event, the midlatitude regions affected may vary.

A third complexity that can be included is that of a baroclinic atmosphere. Since the work of *Bretherton* [1964] it has been recognized that barotropic disturbances more readily propagate to midlatitudes than baroclinic disturbances do. Thus the simple barotropic protomodel would seem to be overstating the midlatitude influence of ENSO-related heating because a heat source concentrated in the middle troposphere produces a local response that is baroclinic. However, numerous investigations have bolstered confidence in the validity of the barotropic theories by showing that perturbations forced by tropical heating are primarily equivalent barotropic by the time they reach midlatitudes and that much of the behavior found in barotropic models continues to hold in tropically forced baroclinic models. Some view this as the result of a natural filtering caused by the trapping of the baroclinic component of disturbances in the tropics. Others [e.g., *Kasahara and Silva Dias*, 1986; *Lim and Chang*, 1986] argue that baroclinic disturbances are transformed into barotropic disturbances by the action of the time-averaged shear. Interestingly, *Tomas and Webster* [1994] have presented striking ob-

servational evidence of analogous behavior for perturbations entering the tropics; the vertical structure of such disturbances changes from equivalent barotropic to baroclinic as they encounter the region of low-level easterlies.

Given the prominent equivalent barotropic structure of midlatitude anomalies, it is quite common to examine extratropical responses to tropical forcings by using a barotropic model forced by tropical divergence. However, results are quite sensitive to the level chosen to represent the atmosphere owing to considerable sensitivity to the strength of the zonal mean wind and its effect on Rossby wave dispersion. Best results (in the sense that the midlatitude response to tropical heating agree with a full baroclinic model) occur for a barotropic model applied at the 350 mbar level [Ting, 1996]. This roughly agrees with the results of the Held *et al.* [1985] study of the appropriate level for use of the barotropic vorticity equation in the study of the forced response of the external mode. Further, as Ting and Sardeshmukh [1993] have demonstrated, the vertical distribution of tropical heating can influence the amplitude of the resulting midlatitude anomalies because meridional propagation is influenced by the vertical structure of a disturbance, as explained above. Because barotropic models cannot represent the spatially varying vertical structure of tropically forced perturbations, baroclinic models are required for accurate representation of their propagation. Therefore, although barotropic dynamics are able to capture the essence of the tropical-extratropical character during ENSO, baroclinic models are needed for quantitative understanding.

A complementary view of the links between tropical forcing and the extratropical response can be obtained using "potential vorticity (PV) thinking" [Hoskins *et al.*, 1985]. Rossby waves propagate along isentropic surfaces on potential vorticity gradients, so the large PV gradients at the tropopause provide preferred channels for wave propagation. Anomalous tropical heating and the associated vertical motion lead to disturbances of the isentropic surfaces in the upper troposphere in the tropics and subtropics. These disturbances can radiate Rossby waves along the PV gradients into midlatitudes, leading to the teleconnections described earlier.

A fourth elaboration arises from recognizing that perturbations stimulated from the tropics are influenced by a time-mean state that is a function of longitude as well as latitude. This means that the midlatitude response to tropical heating is very sensitive to the longitude of the heating as, for example, Simmons [1982] demonstrated. Investigations have found that because of the large scale of longitudinal variations in the mean state, WKB theory holds. This means that the general properties of wave propagation developed in section 3.1 for a zonally symmetric basic state can be applied to the local conditions in a zonally varying mean state. This is true for propagation in both the tropics and in midlatitudes. Applying this principle to the tropics, Webster and Holton [1982] showed that meridional propagation across the tropical belt is possible if it occurs through

a patch of local westerlies, while Webster and Chang [1988] and Hoskins and Jin [1991] suggested that longitudinal variations in the tropical background can affect the zonal propagation of tropical perturbations leading to an accumulation of perturbation energy to the east of a westerly maximum.

In midlatitudes, nonuniformities in the background state can act to refract perturbations to certain preferred locations, as calculations by Branstator [1983], Karoly [1983], and Hoskins and Ambrizzi [1993] have demonstrated. In particular, there is a tendency for the response to tropical forcing to be guided along the Asian and Atlantic subtropical jets. Thus heating anomalies in the tropics well to the west of the dateline tend to influence the North Pacific, while wave trains stimulated by El Niño extend well into the Atlantic. An example of this is seen in Figure 8, which contrasts the simple arching wave train produced by a tropical source in a solid body rotation basic state with the undulating wave train produced by the refractive properties of the local jets in a zonally varying basic state.

3.4. Internal Midlatitude Sources

In addition to the two fundamental elements that were recognized before TOGA, namely heating-induced sources and energy propagation, other basic mechanisms are now recognized to influence the midlatitude impact of tropical heating. In particular, the tropical heat source is not the only significant contributor of low-frequency energy to the anomalies that accompany that heating; other internal dynamical sources also make substantial contributions.

3.4.1. Sources from mean zonal variations.

An important source was discovered in studies of the influence of longitudinal inhomogeneities in the background state. Not only do such nonuniformities affect the local propagation properties of the medium, they also can act as sources of perturbation energy, much as gradients serve as energy sources for various types of hydrodynamical instability. It has been found that perturbations that can gain energy from this source tend to grow from longitudinal gradients in the jet exit regions [Simmons *et al.*, 1983]. Further, they have natural low frequencies and structures similar to wave trains stimulated from the tropics (see Figure 9). This means that wave trains which are actually forced from the tropics and that encounter the jet exit regions will tend to benefit from this internal source. As an example of this, Branstator [1985a] presented a linear barotropic counterpart to an ENSO case (Figure 8b) in which about 40% of the global perturbation kinetic energy was derived from this internal source rather than from the tropical heating anomaly.

Not only can this mechanism make midlatitude anomalies stronger than they would be if tropical heating were their only energy source, it also changes their structure since some components of a pattern stimulated from the tropics will be amplified by this mechanism while others will not. The natural advantage that this mechanism gives some structures is thought (together with the ten-

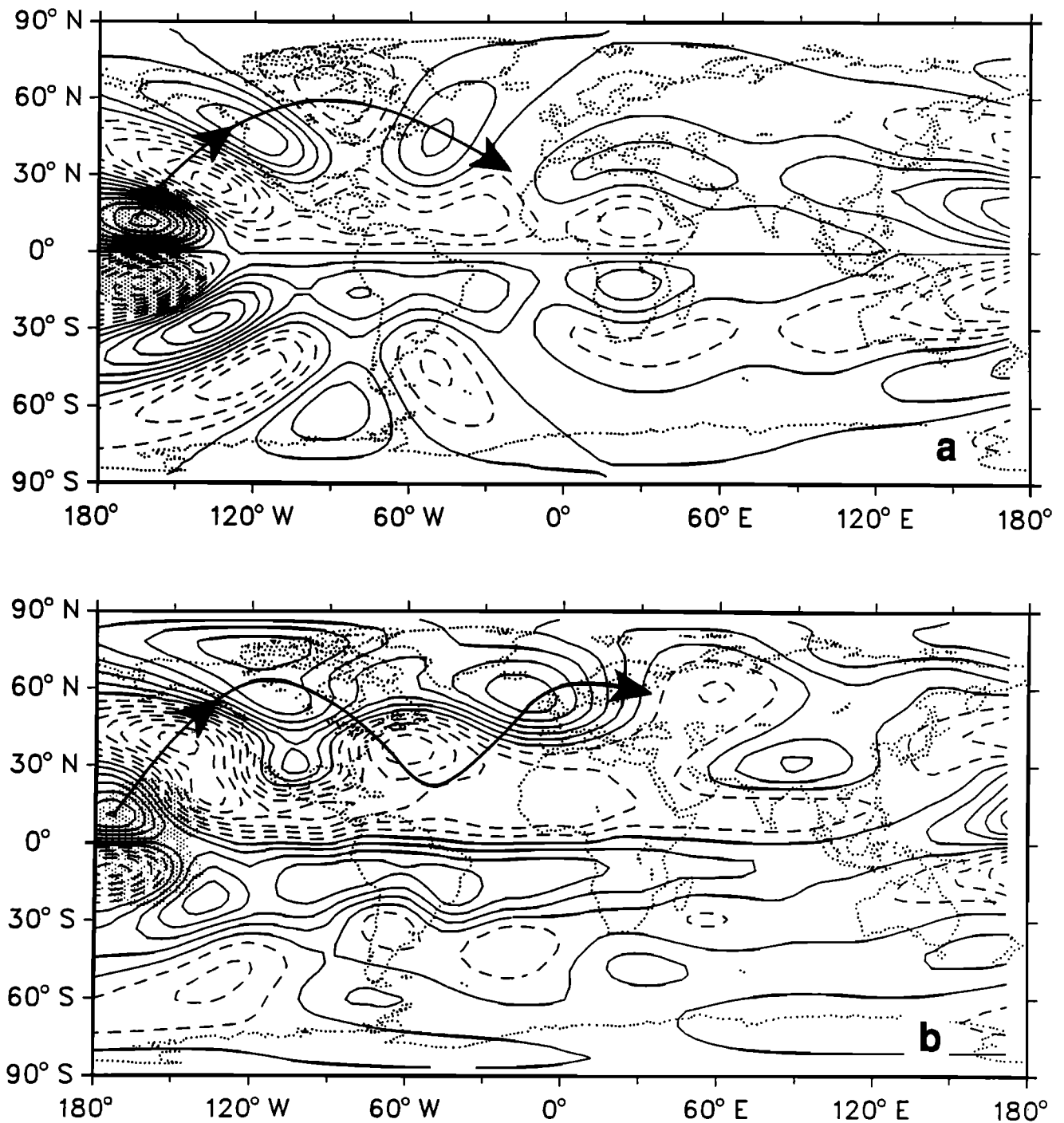


Figure 8. Anomalous streamfunction response to a tropical source in the shaded region for the barotropic vorticity equation linearized about (a) solid body rotation and (b) 300 mbar time mean January conditions. The arrowed thick line indicates the direction of the local group velocity. From *Branstator* [1985a].

dency for the location of the Rossby wave source to be rather insensitive to heating position) to make the response of the atmosphere less sensitive to details of the distribution of tropical heating than one might expect from the protomodel. As extreme examples, *Geisler et al.*, [1985] found that the structure of an AGCM's response to tropical Pacific SST anomalies was nearly constant, even for anomalies placed about 60° of lon-

gitude apart, while *Branstator* [1990] found in a linear baroclinic model that the dominant response to steady heating anomalies with random spatial structure was an arching North Pacific pattern much like that expected from a tropical dateline source.

One consequence of the presence of the internal source and its tendency to produce structures similar to those known to be initiated by tropical sources is that it is

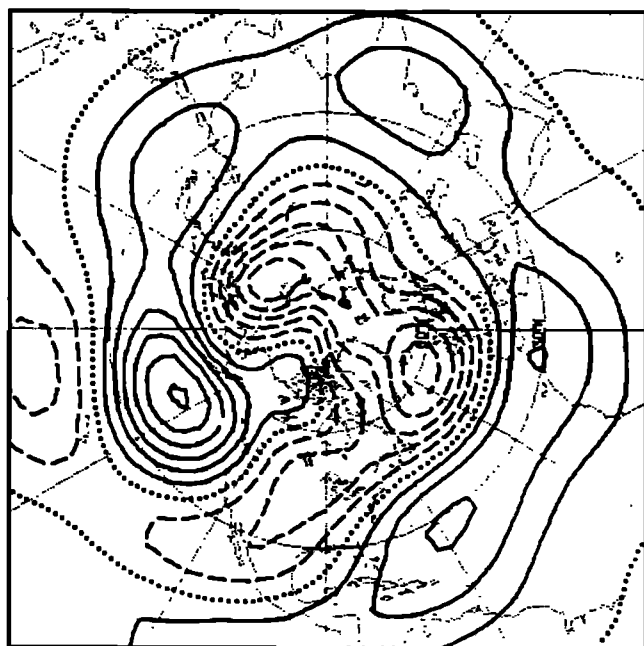


Figure 9. One phase of the fastest growing mode of the barotropic vorticity equation linearized about 300 mbar time mean in January. From *Simmons et al.* [1983].

often difficult to distinguish externally and internally initiated low-frequency perturbations. Even their energy sources can be similar.

3.4.2. Effects of midlatitude storm tracks. A second source of low-frequency energy that was ignored in early models but which is now thought to play a significant role in shaping midlatitude anomalies forced from the tropics are momentum fluxes from the high-frequency transients (terms on the far right-hand side of (5)) when organized into storm tracks, as described from observations in section 2.3. Although fluxes from high-frequency transients have long been recognized as a major contributor to the maintenance of the climatological state, it was the analysis of *Kok and Opsteegh* [1985] that demonstrated the importance of these fluxes in producing the midlatitude anomalies associated with the 1982–1983 El Niño. There were strong anomalous midlatitude eddy flux anomalies during that event, and the Kok and Opsteegh analysis indicated that without them the observed midlatitude time-mean anomalies would have been much weaker. *Held et al.* [1989] found a similar result for an AGCM-simulated El Niño event (see Figure 10) and, furthermore, showed that the subtropical convergence anomaly typically found poleward of a tropical heat source is probably a dynamical response to momentum fluxes from anomalous transients in that region and not just a manifestation of a local Hadley circulation anomaly.

Recent investigations have suggested that anomalous momentum fluxes by high-frequency eddies may do more than just amplify tropically forced wave trains; they may actually control the structure of these wave trains and produce a certain commonality in the mid-

latitude response to various tropical heating anomalies. *Branstator's* [1995] results indicate that low-frequency structures are likely to affect storm tracks and thus alter distributions of high-frequency fluxes, but it may be that only some will do it in such a way as to produce a positive feedback, and these structures will then have a dynamical advantage. In particular, there should be a tendency for perturbations stimulated from the tropics to be composed of those structures that induce a positive feedback.

Hoerling and Ting [1994], from case studies of the circulation associated with a few El Niño events, have concluded that midlatitude transients can influence the character of the extratropical response by causing it to occur at preferred geographic locations. Their work suggests that in the northern hemisphere the enhanced westerlies on the northward flank of the subtropical high induced by El Niño will tend to elongate the jet stream in about the same place for SST anomalies in a variety of equatorial locations. Reasoning that this common jet stream shift is responsible for transient momentum

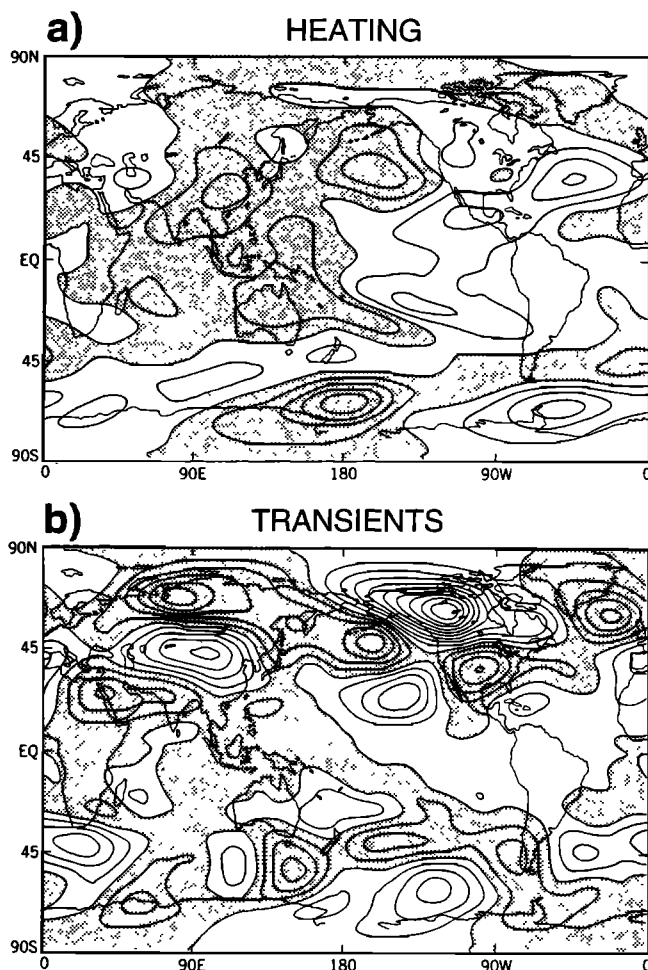


Figure 10. 300 mbar eddy geopotential height response of a linear model to (a) anomalous diabatic heating and to (b) forcing from anomalous transients taken from a GCM simulation of ENSO. The contour interval is 10 m and negative values are shaded. From *Held et al.* [1989].

fluxes having similar structure during each of the events they studied, they argue that it also leads to a commonality in the location of the extratropical response because the response is significantly forced by anomalous momentum fluxes.

3.5. Optimal Source Regions

In the protomodel of the impact of tropical heating on midlatitudes the position, strength, and structure of the midlatitude reaction was largely controlled by the position and structure of the tropical heat source. To a degree, control of these attributes is not as strongly dominated by the external forcing when some of the above discussed modifications to the basic model are taken into account. Local background gradients can tend to fix the position of the Rossby wave source; steering effects by waveguides can make the regions affected by tropical heating somewhat insensitive to the position of the heating; the action of internal sources can effectively limit the degrees of freedom the system uses when responding to tropical sources. However, the midlatitude response is still very sensitive to the specification of the heating. For example, depending on whether or not the forcing is embedded in local westerlies, the midlatitude response may be weak or strong. Similarly, whether or not the source stimulates a wave train that crosses midlatitude regions with strong gradients will affect the strength and structure of that wave train. Moreover, these influences vary in importance with the annual cycle and in the southern hemisphere (see sections 2.1 and 3.7).

Investigations have been undertaken to determine which regions are the most effective sites for the stimulation of midlatitude northern hemisphere anomalies. Though a few observational [e.g., Kiladis and Weickmann, 1992] and AGCM studies [e.g., Geisler et al., 1985] have considered the sensitivity of the midlatitude response to the position and structure of forcing, most information about this question comes from linear investigations, and several tools are available for addressing this issue.

One useful tool for producing information about optimal forcing is the Green's function (response function). For example, Branstator [1990] has calculated the response of a linear baroclinic model to point sources of heating and found that the east coast of Asia and the western Pacific are especially effective at simulating responses that arc across the Pacific and North America. Grimm and Silva Dias [1995a, b] note the dependence of these results on whether or not the basic state divergence is included in computing the Green's function and that the western tropical Pacific becomes less important if it is included. Also, they formulate the problem in terms of "influence functions," which relate directly to anomalous upper tropospheric divergence (rather than the Rossby wave forcing, which can be difficult to interpret). The upper tropospheric divergence can be more directly linked to anomalous convection and tropical heating.

Another approach that has proven to give insight is eigenanalysis. Assuming the atmosphere can be approximated by a discrete linear model \mathbf{L} with N degrees of freedom, the problem is simplified by using the eigenfunctions E_j , $j = 1, \dots, N$ of that model as a basis. If \mathbf{L} is stable, has eigenvalues σ_j , and is forced by an oscillating source with frequency σ and structure r ,

the response will asymptote to $R = \sum_{j=1}^N \left(\frac{r_j}{\sigma - \sigma_j} \right) E_j$.

Here r_j is the dot product of r with F_j , the eigenvector of the adjoint of \mathbf{L} , whose eigenvalue is equal to the conjugate of σ_j [see Branstator, 1985b]. It is apparent that to learn about which forcing distributions can give especially strong responses to, say, steady forcing, one should examine the adjoint eigenfunctions with low-frequency, nearly neutral eigenvalues. Simmons et al. [1983] discovered that these dominant low-frequency modes tend to be concentrated over the North Pacific and North Atlantic when climatological January conditions are used for background states, as the example in Figure 9 demonstrates. Branstator [1985b] and Ferranti et al. [1990] found that the eigenfunctions which influence low-frequency structures over the North Pacific tend to have adjoint structures that are concentrated in the south Asia region.

3.6. Nonlinearities

Much of the above analysis is based on linear dynamics, and studies ranging from the idealized barotropic experiments of Haarsma and Opsteegh [1989] to the AGCM simulations of Ting and Held [1990] support the reliance on linear theory for tropical heat sources with amplitudes near those observed. However, there is still potential for nonlinearity to be important, especially if more than a qualitative understanding is desired. For example, according to work by Hendon [1986] and Sardeshmukh and Hoskins [1988], something as basic as the position of the subtropical high forced by equatorial heating is affected by nonlinearities. Midlatitude features forced from the tropics can also be influenced, as nonlinearity alters both the propagation properties of the flow and the Rossby wave source, as examples by Sardeshmukh and Hoskins [1988] demonstrate. Nonlinearity may also come into play in more complicated, two-way interactions between the tropics and midlatitudes. Observational studies [e.g., Lau and Phillips, 1986; Kiladis and Weickmann, 1992] show that tropical convection can be stimulated by wave trains originating in midlatitudes. Combining this behavior with the conventional forcing of midlatitude anomalies from the tropics, Molteni et al. [1993] concluded from a diagnosis of AGCM experiments that tropical forcing anomalies can be affected by changes in midlatitude anomalies which have their origins in the tropics.

One commonly considered source of nonlinearity, which does not appear to be very important, is the interaction between equatorial heating-induced changes in the zonal mean state and climatological departures from the zonal mean. Numerous studies have shown

how the zonal mean state influences the climatological waves (see section 2.1.1), so it is natural to wonder if the zonal mean anomalies induced by tropical heating significantly change the climatological waves. *Hoerling et al.* [1995] have found, at least in the northern hemisphere winter, that these zonal mean anomalies are too far south of the large amplitude climatological waves to make large changes in their structure and strength.

Although some investigations suggest that the role of nonlinearity is to introduce fairly minor modifications to what is essentially a linearly controlled system, others suggest that the midlatitude response to tropical heating is fundamentally nonlinear. One piece of evidence in support of this latter idea comes from a study by *Ferranti et al.* [1994] of the reaction of an AGCM to an idealized Indonesian SST anomaly. Associated heating leads to enhanced blocking activity in both the North Pacific and North Atlantic. Since the enhancement cannot be explained as a simple additive effect of the time-averaged anomaly produced by the heating, it is reasoned that this is evidence of the nonlinear character of the heating's influence. Furthermore, it appears that the way the heating enhances blocking frequency is to increase the population of one of the model's nonlinear regimes, one in which blocking is likely to occur because the regime has stronger than normal ridges in the northeast Pacific and northeast Atlantic. From this viewpoint, tropical heating is seen to influence midlatitudes by making some transitions between regimes more likely than others, just as forcing is seen to bias transitions in very simple dynamical systems [*Palmer*, 1993].

3.7. Seasons Other Than Winter

Although most of the observational and modeling studies of ENSO have focused on the teleconnections from the tropical Pacific into the northern hemisphere in December-February, the few studies that have considered the other seasons or the southern hemisphere suggest that the influence of the tropics on midlatitudes may be greater than the theories prevalent at the beginning of the TOGA decade indicated. An example of a strong signal during the northern spring and early summer occurred in 1988. Studies of the 1988 North American drought have shown that subtropical SST anomalies can have a substantial influence on midlatitude standing wave anomalies at that time of year [*Trenberth et al.*, 1988; *Mo et al.*, 1991; *Trenberth and Branstator*, 1992]. Another example is the mature El Niño of 1993, which has been shown to have an influence on setting up conditions leading to the North American spring-summer floods [*Trenberth and Guillemot*, 1996]. The midlatitude summer influence may be stronger than once thought possible because the zonal mean tropical easterlies are not as strong a barrier to influences from the tropics as once thought, owing to the effects of the local Hadley circulation as a Rossby wave source and the wavy basic state enhancing wave propagation [*Lau and Peng*, 1992; *Grimm and Silva Dias*, 1995b]. The NH summer pattern (related to the PJ or ANA pattern, see section 2.1.1) may have played a

role in both cases [e.g., *Grimm and Silva Dias*, 1995b]. Nevertheless, in the summer hemisphere the influence of the tropics is usually more regionally confined than in the winter hemisphere [*Chen and Yen*, 1993].

An alternative explanation for larger-than-expected boreal summer influence might involve enhanced tropical forcing, but all indications are that tropical heating anomalies in the ENSO region are at a minimum then. This is demonstrated by Figure 11, which shows the annual cycle from 1979 to 1995 of the interannual standard deviation of the SST and OLR fields from 5°N to 5°S. The maximum SST variance throughout this period occurred for December coinciding with the mature stage of El Niño events, although the maximum variance in SST along the coast of South America occurs several months later. Standard deviations of SST of less than 0.7°C prevail from May through September, compared with values over 1°C from November through February. The manifestations of this in terms of changes in tropical convection are altered by the mean SSTs on which the anomalies are imposed. Because the cold tongue is most pronounced in September and October, SST gradients along the equator remain strong regardless of SST anomalies, and the near-equatorial convective response in the northern fall (Figure 11) is weak. The weakest SST gradients occur in the northern spring, and it is at that time of year that convection occurs in the vicinity of the equator. Thus the annual cycle of OLR variance is shifted westward, toward the warmer waters, and shifted in time by a month or so relative to SST to have its maximum in January, with standard deviations exceeding 20 W m⁻² only from December to April. Values are lowest from June through October. During these months the maximum variance occurs on the equator slightly west of the dateline, but there is a weaker maximum of about 10–12 W m⁻² along about 10°N. Thus the values in Figure 11 accurately depict the annual cycle in the tropical forcing signal.

4. Modeling and Prediction

4.1. Sensitivity to Idealized and Realistic Forcings

The first attempts to determine the full dynamical response to tropical SST anomalies using AGCMs employed simplified scenarios, such as a perpetual January and a prescribed SST anomaly with a constant and enhanced amplitude [e.g., *Julian and Chervin*, 1978; *Blackmon et al.*, 1983; *Shukla and Wallace*, 1983]. The perpetual January runs were necessitated in part by computer limitations and the need for long enough runs so that a signal would emerge from the anomalous forcing. The results of these experiments were encouraging in that the models were able to capture the gross features of response observed in nature. However, they also served to point out inadequacies in the theory by suggesting that the amplitude of the response was not a linear function of the heating anomaly and by indicating that the response pattern did not simply shift posi-

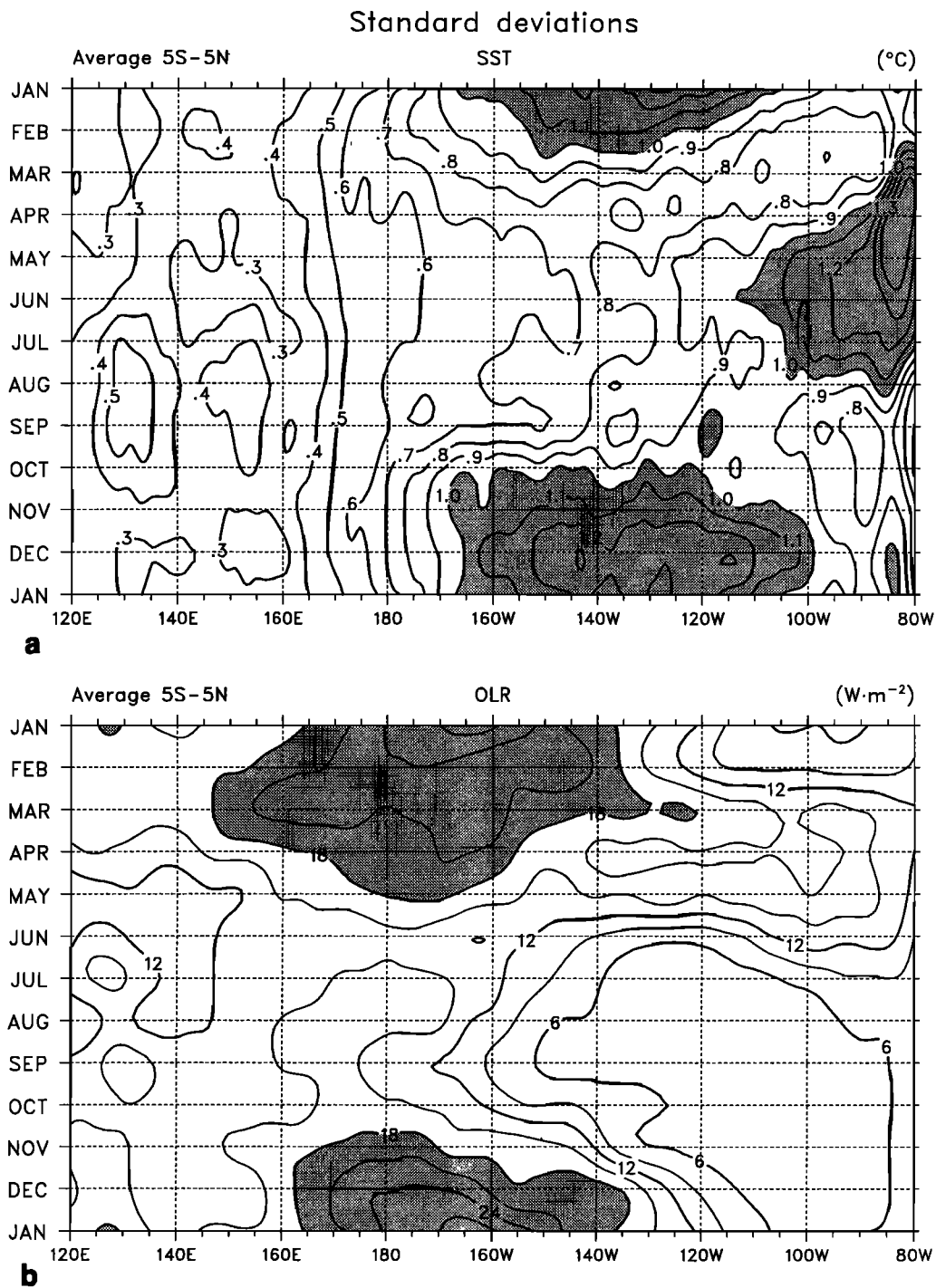


Figure 11. The mean annual cycle of the interannual standard deviation of monthly mean (a) SST and (b) OLR for 5°N to 5°S presented as a function of longitude and month based on 1979 to 1995. For SST the contour interval is 0.1°C , and values greater than 1.0°C are stippled. For OLR the contour interval is 3 W m^{-2} , and values greater than 18 W m^{-2} are stippled.

tions if the location of the heating was shifted [Geisler *et al.*, 1985].

The use of perpetual season integrations and idealized SST anomalies led to uncertainties in how to compare the results of these early experiments with observations. Nevertheless, there is still considerable scope for advancing understanding by making numerical experiments with some aspects of the system changed or ide-

alized. In the TOGA decade these included exploring the differences in the extratropical responses with differing climatological features present in the mean flow, such as those created by removing mountains [Blackmon *et al.*, 1987], considering different times of year [e.g., Peng *et al.*, 1995; Ting and Peng, 1995], and specifying a zonal mean climatology [Ting and Held, 1990], for example. Other examples are given in section 4.2.

Palmer and Mansfield [1986a] also noted great sensitivity in their response experiments to the basic climatology of their model, thus helping to reinforce the idea suggested by linear calculations that the atmospheric response to El Niño heating cannot be well simulated unless a model has an accurate climatology under undisturbed conditions. The importance of the SST field itself was highlighted by *Palmer and Mansfield* [1986b], who noted that details in the SST anomalies in different El Niño events are reflected in very different atmospheric teleconnection responses over the North Pacific and North America.

While idealized experiments with AGCMs proved to be very useful during the early TOGA period, the spatial patterns of the SST perturbations were often prescribed using certain analytic functions, and the forcing was assumed to be temporally fixed. This model scenario differs considerably from the observed situation. The constraints inherent in the idealized experiments can be removed by incorporating the observed month-to-month SST variability [*Lau, 1985*]. The added degree of realism allows for more incisive diagnoses of the transient atmospheric responses through complete ENSO cycles and facilitates comparison with observations. The seasonal dependence of such responses, as well as the commonalities and distinctions among the atmospheric behavior in different El Niño events, is also highlighted.

This experimental design is adopted in the Atmospheric Model Intercomparison Project (AMIP) [see *Gates, 1992*], which is an extensive program aimed at assessing the capability of various AGCMs during the 1979–1988 decade. All models being evaluated are generally restricted to a single integration subjected to the same temporal sequence of SST and sea ice forcing observed during this decade. AMIP provides a valuable data set for evaluating model simulation of certain atmospheric features, however, outside of the tropics, the basic design of AMIP makes the comparison of the model's sensitivity to the SST forcing difficult because single realizations are not sufficient to suppress that noise unrelated to SST. In view of the high level of "background" variability in both the real and simulated atmospheres, conclusions drawn from single realizations of observed or model responses to a certain SST anomaly must be viewed with caution, and ensembles of AGCM integrations are needed (section 4.3).

4.2. Relative Importance of Tropical and Extratropical SST Anomalies

An ongoing debate in the research and operational communities is concerned with the relative roles of SST anomalies in different parts of the world ocean in modulating midlatitude atmospheric variability. Extratropical SST anomalies have long been used as a prediction tool in the extended range by forecasters in North America and Europe [e.g., *Namias, 1969; Ratcliffe and Murray, 1970*]. The observational and theoretical studies reviewed in sections 2 and 3 have shifted some of the focus to the remote link between SST anomalies in the

tropics and the atmospheric circulation in the extratropics. Because the observed SST perturbations in various maritime sites seldom occur in isolation and the limited duration of historical data sets does not provide for enough examples of the anomalous events, the question of whether tropical or extratropical SST anomalies are more effective contributors to midlatitude atmospheric variability cannot be fully answered by analyzing observational data alone. This problem is further complicated by the notable temporal correlation between the tropical and extratropical SST anomalies (see section 5.1). As an alternative, the controllable environment of a general circulation model (GCM) can be used for addressing issues of this kind. Assessment of the relative impacts of SST anomalies at different sites can be made by comparing model responses in an array of experiments, with each member being subjected to forcing confined to one of these sites only. Any limitations on data sampling inherent in observational studies can readily be removed in the model studies by lengthening the integration (for temporally fixed SST anomalies) or by conducting multiple integrations (for time-varying SST anomalies).

Attempts to evaluate the influences of tropical and extratropical SST changes on the midlatitude flow pattern using AGCMs have been made by *Pitcher et al. [1988]*, *Lau and Nath [1990, 1994]*, *Graham et al. [1994]*, *Ferranti et al. [1994]*, and *Kharin [1995]*. Notwithstanding the differences in model numerics and physics, experimental design, and analysis procedures, a majority of these investigations conclude that SST anomalies introduced in the tropical Pacific produce stronger and more reproducible atmospheric responses than those in the extratropics.

To illustrate the model responses to SST anomalies prescribed in various sites, the composites of the midtropospheric height anomalies for selected winters are shown in Figure 12. Figures 12a, 12c, 12e, and 12g are based on those six winters in which the tropical Pacific is much warmer than normal and the central extratropical North Pacific is much colder than normal (referred to hereafter as WTP/CNP). Figures 12b, 12d, 12f, and 12h correspond to those six winters with much colder tropical Pacific and warmer central North Pacific SSTs (referred to hereafter as CTP/WNP). These composites have been constructed by *Lau and Nath [1994]* using observational data (Figures 12a and 12b) and output from model runs in which observed intermonthly SST variations within the 1946–1988 period are prescribed at grid points throughout much of the global ocean (the Global Ocean-Global Atmosphere (GOGA) runs, Figures 12c and 12d), within the tropical Pacific only (the TOGA runs, Figures 12e and 12f), and within the extratropical North Pacific only (the Midlatitude Ocean-Global Atmosphere (MOGA) runs, Figures 12g and 12h). Each of the GOGA, TOGA, and MOGA experiments has been repeated four times using independent atmospheric initial conditions, so as to enhance the sample size.

A considerable qualitative resemblance exists (Fig-

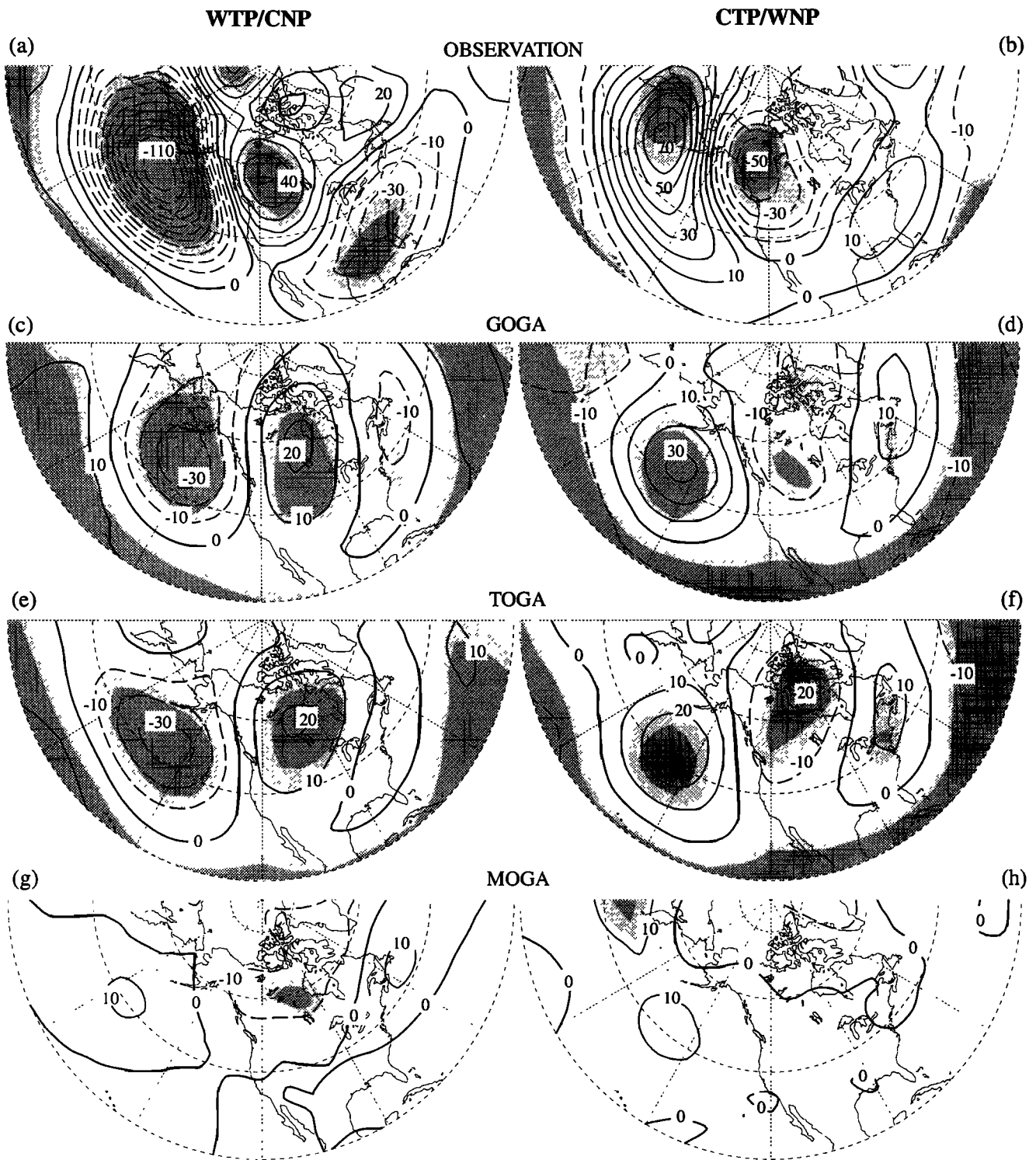


Figure 12. Composites of the anomalies in (a,b) observed 500 mbar height, and in 515 mbar height averaged over four parallel runs in the (c,d) GOGA, (e,f) TOGA, and (g,h) MOGA experiments. These composites are performed over six winter seasons with warm SST anomalies in the tropical Pacific and cold anomalies in the North Pacific (WTP/CNP, Figures 12a, 12c, 12e, and 12g) and over six seasons with cold SST anomalies in the tropical Pacific and warm anomalies in the North Pacific (CTP/WNP, Figures 12b, 12d, 12f, and 12h). Contour interval is 10 m. Anomalies exceeding the thresholds for significance at the 90% and 95% levels are indicated by light and dark stippling, respectively. From *Lau and Nath* [1994].

ure 12) between the patterns based on observed, GOGA, and TOGA data, thus indicating that the model atmosphere is capable of producing a realistic and statistically significant response when subjected to SST changes in the global ocean and that much of this response may be attributed to forcing in the tropical Pacific only. The much weaker and less significant patterns for the MOGA experiment illustrate the relative inefficiency of midlatitude SST anomalies in forcing the model atmosphere.

The amplitude of the simulated height anomalies in all three model experiments shown in Figure 12 is notably lower than the observed data. This discrepancy is likely related to the following factors. The low resolution (rhomboidal truncation at 15 wavenumbers) used in this suite of integrations is not sufficient to fully reproduce the constructive effects of the storm tracks on the slowly varying background flow (see sections 2.3 and 3.4.2). The precipitation response (which in turn governs the magnitude of the atmospheric forcing due to diabatic heat sources) to SST changes in the particular AGCM used in this study is also known to be weaker than observed. The four separate realizations of the model response for a given SST forcing exhibit considerable sampling fluctuations, so that the magnitude of the ensemble average presented in Figures 12c-12h is typically less than that of the individual samples.

Whereas most of the modeling studies have considerable success in simulating the extratropical atmospheric response to tropical SST anomalies, the corresponding results for the in situ response to prescribed midlatitude SST forcing are much more tenuous. *Pitcher et al.* [1988] and *Kushnir and Lau* [1992] have pointed out the nonlinear nature of the model response with respect to the polarity of the midlatitude SST anomaly (i.e., the response pattern does not undergo a straightforward sign reversal when a warm anomaly is replaced by a cold anomaly). This behavior is in sharp contrast to the quasi-linear response to tropical SST forcing (e.g., note the change in polarity of the spatial patterns simulated in the TOGA runs for warm and cold SST anomalies in the tropical Pacific in Figures 12e-12f). There also exist notable discrepancies among various AGCM experiments concerning the amplitude of atmospheric anomalies associated with extratropical SST changes (e.g., contrast the strong signals reported by *Latif and Barnett* [1994] with the feeble response illustrated in Figures 12g-12h).

A partial explanation for the above puzzling findings has been offered by *Trenberth and Hurrell* [1994]. While the changes in eddy transports due to the altered synoptic systems (section 2.3) constitute a major complication, another potentially important factor is that the atmospheric heating effects may not be local in the extratropics. The sensible heat exchanged between the ocean and atmosphere is realized locally, but the latent heat lost by the ocean through evaporation is realized only as an increase in moisture, and the actual atmospheric heating is not realized until precipitation occurs, often far downstream. The occurrence of conden-

sational heating hence depends on the prevailing synoptic situation and varies geographically according to the ambient flow field. These nonlocal effects therefore exhibit a strong regional dependence, and they add a large variable component to any forcing, thus complicating attempts to detect any systematic effects in both the real atmosphere and models. This issue may also contribute to discrepancies among the findings based on different model experiments, in which the prescribed SST anomalies vary in location and intensity, and the model climatologies vary. Placing "super SST anomalies" into a model will enhance the local effects so that results are more likely to appear as significant, but they are also much more likely to be unrealistic and inappropriate for the real atmosphere.

The extended integrations examined by *Kitoh* [1991], *Lau and Nath* [1994], *Graham et al.* [1994], *Kumar et al.* [1994], and *Kawamura et al.* [1995] demonstrate that the current AGCMs are not only capable of reproducing the essential features accompanying individual ENSO cycles on interannual timescales, but they also yield longer-term atmospheric signals in response to decadal changes in the SST conditions, such as those described in section 2.2. In particular, aspects of the decadal scale change beginning around 1976 have been simulated with atmospheric models using specified SSTs [see also *Miller et al.*, 1994]. Further model experiments of this kind should shed new light on the mechanisms contributing to air-sea interaction on interdecadal timescales, as well as on the relationships between such slowly varying phenomena and the individual El Niño events.

4.3. Predictability and Prediction of the Extratropical Atmospheric Circulation

Previous sections have shown that SST variations in the tropical Pacific can have considerable impact on the extratropical mean state. It is desirable to determine a "signal" of tropical SST anomalies in the extratropical atmosphere. However, any such signal will be embedded within the "noise" of natural variability associated with midlatitude weather systems. The noise must be regarded as unpredictable while the signal might be potentially predictable given a very good model and a good forecast of the tropical SSTs. The signal-to-noise ratio then provides a measure of the predictability.

One estimate of the signal comes from observed fields composited with respect to the tropical SSTs. This method brings out the common features of each case but ignores differences in forcings coming from other parts of the tropics and elsewhere. It also ignores differences in details such as magnitudes and spatial patterns of SST anomalies in the region used for compositing [e.g., *Tomita and Yasunari*, 1993; *Trenberth*, 1996]. In principle, this could be overcome with a long enough observational record but in practice, the observational record is very short, and no two El Niño events, for instance, are alike. Consequently, any such estimates of potential predictability from this technique will be very conservative.

Another method is to use AGCMs forced with tropical SST anomalies. If the same SST anomaly is used for many cases with different starting conditions, an ensemble is created that can be used to generate statistics of the results. In this case, the signal is the reproducible component, while the scatter among the members of the ensemble is a measure of the noise. The signal can be regarded as “externally” forced, while the noise depicts the “internal” variability. Such results should provide estimates for the predictability but will depend on the veracity of the model and thus will be model dependent. Thus the interannual variability of the atmosphere may be separated into internal and external components [Harzallah and Sadourny, 1995; Rowell *et al.*, 1995; Kumar and Hoerling, 1995].

An example of the spatial distribution of internal and external variance for January mean 200 mbar height based on a nine member ensemble of AGCM integrations forced with observed SSTs is shown in Figure 13 (for more details see Kumar and Hoerling [1995]). Over

almost the entire extratropical region the total variance is dominated by the internal part with the external variance playing a minor role. Note that the spatial pattern of maximum external variability in northern extratropics is similar to a PNA-like pattern, confirming results from earlier empirical, theoretical, and modeling studies. Similar results using different GCMs have been reported by Barnett [1995], Smith [1995], and Dix and Hunt [1995].

To assess the seasonality of the atmospheric response, area-averaged internal and external variances of monthly mean 200 mbar height together with the ratio of external to internal variance over the PNA region (defined as 180°–60°W and 20°–80°N) for different months during northern winter are given in Table 1. For contrast, similar variances for a summer month are also shown. During northern winter the SST-forced variability maximizes during February with a sharp decrease in the following spring. Atmospheric internal variability also shows qualitatively similar behavior. The ratio of total-

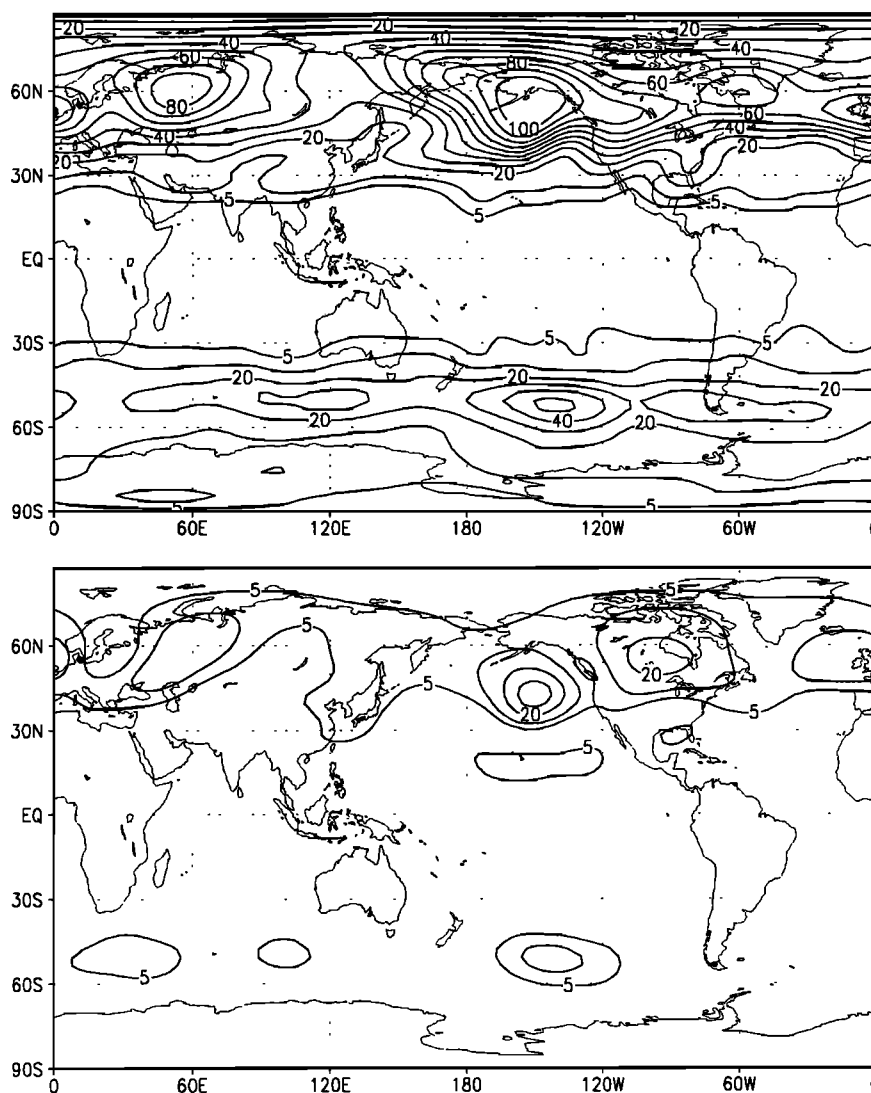


Figure 13. An example of the spatial distribution of internal and external variance for January mean 200 mbar height based on a nine member ensemble of AGCM integrations forced with observed SSTs. The contour interval is 10 square decameters (dam^2) except a 5 dam^2 contour is also plotted. (For details, see Kumar and Hoerling [1995]).

Table 1. Area Average Variances of Monthly Mean 200 mbar Height Over the NA Region for different calendar months

	December	January	February	March	April	July
Internal variance (I)	4400	4925	4001	3285	3410	1736
External variance (E)	759	867	1117	711	784	294
Ratio (E/I)	0.17	0.18	0.28	0.22	0.23	0.17

The PNA region is defined as an area bounded by 180°–60°W and 20°–80°N. Variances are calculated from an ensemble of 9 AGCM integrations for 1982–1993 forced by observed SSTs. Units of I and E are square meters.

to-internal variance indicates potential predictability to be highest during February. Decreases in the external variance during spring are offset by a similar decrease in the internal variability with the consequence that ratio of variances does not show a marked decrease from winter to spring.

Because of reduced baroclinicity, there are reasons to expect reductions in atmospheric internal variability during spring and summer. However, there is also a decrease in the external variance which could be partly due to changes in the atmospheric mean state and its impact on the tropical-extratropical interactions [Webster, 1982]. It is also due to the size of tropical SST anomalies which show a sharp decline in spring (Figure 11). As noted in section 3.7, anomalous forcing has been historically weakest during the northern summer. In the few cases where large tropical SST anomalies have occurred, predictability may be enhanced. For the summer month of July both internal and external variances over the PNA region in Table 1 are much smaller than during winter. However, the potential predictability for the July monthly mean anomalies over the PNA region is comparable to the winter months. Whether this is an artifact of this AGCM or also holds for nature remains to be seen. The results from the AGCM, however, should provide further impetus for observational predictability studies during the northern summer season.

Another prospect for increasing signal-to-noise ratios from those for monthly means is to take longer time averages. Indeed, if increased time averaging, without unduly affecting the atmospheric signal, leads to a reduction in the unpredictable noise, it will also lead to enhanced predictability. To demonstrate, area-averaged variances for 1, 3, and 5 month time means are given in Table 2. Also shown are the relative changes in the variances from their respective values for the monthly time means. As can be seen, time averaging reduces internal variability, and for 5 month means the variance is 4.5 times smaller than that corresponding to the monthly means. Increased time averaging, however, leads to some reduction in the external variance as a consequence of nonstationarity of the atmospheric signal. Nonetheless, time averaging indeed improves the signal-to-noise ratio, and this has practical forecasting implications because only one possible outcome is realized. For 5 month time means, SST forcing results in a 55% signal-to-noise ratio as compared to 21% for

the monthly means for this case. This increase in the fraction of variance due to external forcing makes the 5 month time averages better suited for identifying the low-frequency atmospheric signal [Trenberth and Hurrell, 1994].

A complementary method of analysis of predictability is based on the probability density function (PDF) of atmospheric states. The PDFs could be for some variable at a geographical location or some spatial pattern, for example, EOFs of seasonal variability [Barnett, 1995]. Any point on the PDF is a possible realization for that variable (or the spatial pattern). For a specified SST anomaly an ensemble of AGCM integrations is an attempt to sample the PDF. Because different points on the PDF have a certain probability of occurrence, ensemble prediction is inherently probabilistic in nature. Thus second moment information from the PDF is also useful, and knowledge of the PDF allows probabilities to be assigned. Because the observed value may be on the wings of a PDF, it follows that even with a perfect model the skill of the ensemble prediction need not always be high [Kumar and Hoerling, 1995; Smith, 1995]. Further, the average of the correlation of the ensemble mean AGCM prediction with the individual events, depends on the signal-to-noise ratio. Although the ensemble mean is but one candidate among many AGCM integrations for possible predictions, it offers the particular advantage that it is also the region of maximum likelihood for individual events and thus it is the best possible prediction. Other states on the PDF different from the ensemble mean will fare poorly as forecasts

Table 2. Area Average Variances of Monthly Mean 200 mbar Height Over the PNA Region for different time means

	1 Month	3 Month	5 Month
Internal variance (I)	4004	1480	888
External variance (E)	847	577	489
Ratio (E/I)	0.21	0.39	0.55
$I(1\text{Month})/I$	1.00	2.70	4.50
$E(1\text{Month})/E$	1.00	1.47	1.73

Variances for 1 month are averages of those from December to April. Variances for 3 month time means correspond to the average of variance for DJF, JFM, FMA, and variances for 5 month time means for the DJFMA. Units are in square meters.

on average. Further, compared to linear statistical prediction methods, the ensemble AGCM predictions have a better chance of incorporating the possible nonlinear dependence between the SST and atmospheric anomalies. Of course, AGCM errors, small ensembles (which lead to unreliable estimates of the PDF), and errors in estimation of the first moment can further compromise estimates of the upper bounds on the predictability.

Practical implementation of an ensemble-based prediction system requires an estimate for the needed ensemble size and this can be done a priori. For 200 mbar extratropical heights, ensembles of size greater than six appear to be sufficient to obtain a reliable estimate for the monthly anomalies [Kumar and Hoerling, 1995]. However, the ensemble size required to make the tropical rainfall anomalies pass a statistical test of significance is different [Dix and Hunt, 1995]. For example, over the equatorial tropical Pacific where the interannual variation in the SST maximizes, a single AGCM realization is sufficient to bring out the signal in the seasonal mean precipitation anomaly.

To summarize, although the accuracy of tropical Pacific SST predictions is expected to improve, analysis of the AGCM ensemble integrations forced with observed SSTs indicates that for the extratropical atmosphere the usefulness of the tropical SSTs as predictors may be limited. The dominance of the extratropical natural variability over modest SST-related signals implies that the extratropical atmosphere is only weakly constrained to select the SST-forced signal. Alternatively, it is only when the expected anomalies of SST are large in the tropics that a reliable forecast can be expected for the extratropics. An example is given for 1982-1983 in Bengtsson *et al.* [1993]. Ensemble averages have clearly demonstrated their usefulness as analysis tools and have also established a basis for dynamical seasonal prediction, the usefulness of which can be exploited under special circumstances.

Much of the current analysis is related to the AGCMs forced with observed SSTs. Predictability estimates from these can only be realized if the SSTs are known perfectly. For coupled prediction systems, errors in the SST forecast also have an adverse impact on atmospheric predictions. To a certain extent, the prediction errors in the SSTs can be corrected by statistical intervention, and the AGCM can be reintegrated with corrected SSTs. Such "two-tiered" schemes for SST-based seasonal prediction are the first logical step toward a fully coupled seasonal prediction system [Bengtsson *et al.*, 1993].

5. Extratropical Feedbacks

5.1. Extratropical Atmosphere-Ocean Coupling

Much of the discussion in the preceding sections has been concerned with the processes through which the ocean influences the atmosphere. However, changes in the atmospheric circulation could also influence the ocean, and the resulting SST anomalies could in turn feedback on the atmosphere. The empirical evidence

presented by Davis [1976], Wallace *et al.* [1990], Cayan [1992], and Trenberth and Hurrell [1994] indicates that observed SST perturbations in the midlatitude oceans exhibit a temporal lag behind the local atmospheric anomalies. This relationship is consistent with the notion that the observed SST changes in the extratropics are primarily responses to, rather than causes of, atmospheric changes.

As illustrated schematically in Figure 14, the atmospheric circulation (in particular, the quasi-stationary flow and the oceanic storm tracks) essentially serves as a "bridge" linking tropical SST changes with those in the midlatitude oceans. This conceptual framework has proven to be helpful in understanding the prominent mode of SST variability in the entire Pacific Basin, with temperature of the near-equatorial waters in the ENSO region being positively correlated with that off the western seaboard of North America and negatively correlated with that in the central North Pacific [e.g., see Weare *et al.* [1976], and Figure 15a in the present paper). This framework helps to clarify the nature of some statistical relationships between atmospheric changes and extratropical SSTs which, rather than being causal, may simply reflect a common relationship to other changes.

During an El Niño event the midlatitude atmospheric response to tropical SST forcing often takes the form of

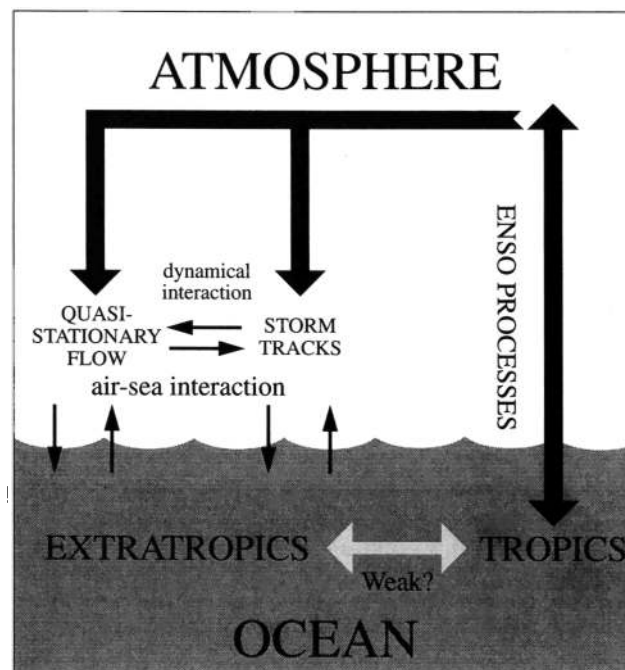


Figure 14. Schematic diagram illustrating the processes associated with the atmospheric bridge linking tropical SST anomalies to changes in the extratropical oceans. Bold arrows depict the remote influences of ENSO-related SST variations on the midlatitude quasi-stationary flow and storm tracks, which in turn affect the underlying ocean through in situ air-sea interaction. Also highlighted is the dynamical interaction between the low-frequency atmospheric fluctuations and the storm tracks.

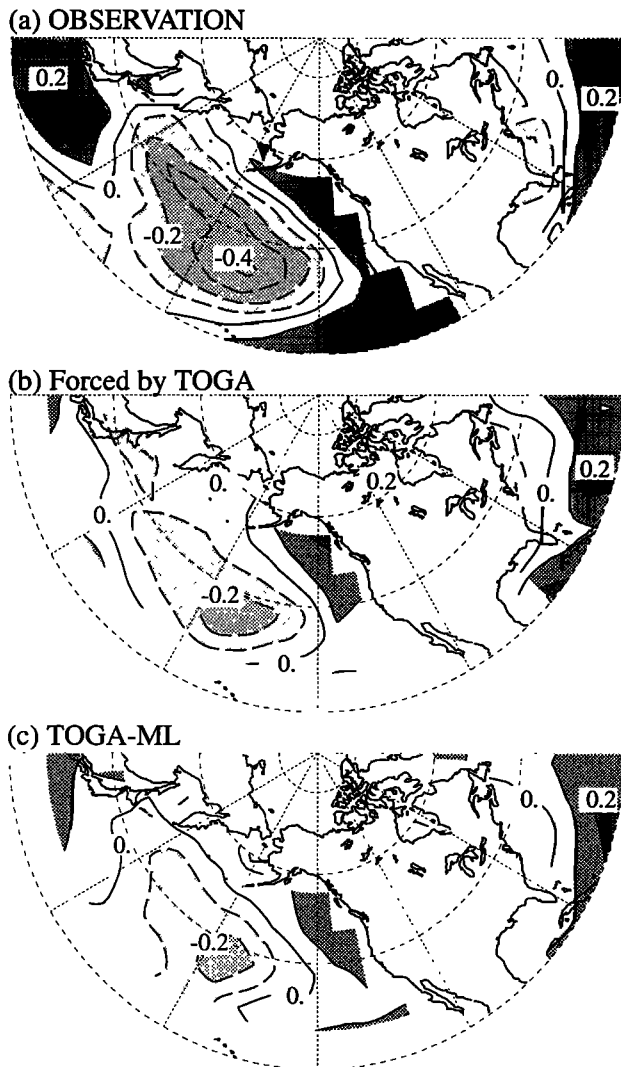


Figure 15. Distributions of temporal regression coefficients of (a) observed SST, (b) mixed-layer (ML) temperature forced by surface fluxes generated in the TOGA experiment and modulated by linear damping with a 5 month timescale, and (c) mixed-layer temperature predicted in the fully coupled TOGA-ML experiment, versus an index of the tropical Pacific SST field. See *Lau and Nath [1996]* for further details. Contour interval is 0.1°C . The regression values shown here correspond to typical extratropical SST changes per 1°C of warming in the central equatorial Pacific.

a PNA pattern (see section 2.1), with lowered sea level pressure over the central North Pacific. The enhanced cyclonic circulation around the deepened Aleutian low is accompanied by cold and dry advection to the west of the pressure anomaly, warm and moist advection to the east, and increased wind speeds and storminess to the south. The resulting changes in the sensible and latent fluxes across the air-sea interface are consistent with the cold SST anomaly observed in the central North Pacific and the warm anomaly near the North American coast during an El Niño episode. These patterns extend throughout the upper ocean [*Deser et al., 1996*]. In the latter region, especially from California southward,

upper ocean warming is also produced by downwelling Kelvin waves that propagate along the coast from the tropical Pacific in association with El Niño events, and subsequently, offshore propagation of Rossby waves can spread the signal into the central Pacific, with modifications by local wind forcing, [e.g., *Johnson and O'Brien, 1990; Jacobs et al., 1994*].

The reality of the “atmospheric bridge” mechanism has been critically tested by *Alexander [1990, 1992a, 1992b]*, *Luksch and von Storch [1992]*, *Miller et al. [1994]*, *Gallimore [1995]*, and *Lau and Nath [1996]* using ocean-atmosphere models and experimental designs of various degrees of complexity. These investigations entail the forcing of a model of the midlatitude ocean by atmospheric variations accompanying El Niño events. The atmospheric driving is obtained either directly from observations during prominent El Niño events or from AGCM simulations subjected to SST perturbations in the tropics. The model evidence gathered thus far indicates that the atmosphere indeed constitutes an important link between SST variability in the tropics and extratropics.

Lau and Nath [1996] demonstrated that the essence of midlatitude air-sea interaction pertinent to the atmospheric bridge mechanism can be captured by forcing the temperature of a motionless, constant-depth ocean mixed layer with the surface heat flux anomalies generated remotely by tropical SST changes. The negative feedback of the mixed layer temperature response on the imposed fluxes can be represented by a linear damping term. The resulting tendency equation for mixed-layer temperature bears a close analogy with that of a first-order Markov process.

In Figure 15 are shown the distributions of the temporal regression coefficients between an index of the central equatorial Pacific SST and (Figure 15a) the observed SST, (Figure 15b) the simulated temperature response of a 50 m mixed layer driven by surface fluxes produced in the TOGA experiment described in section 4.2 and modulated by a linear damping term with a dissipative timescale of 5 months, and (Figure 15c) the temperature variations of the same 50 m mixed layer which is fully coupled to an AGCM with time-varying SST forcing in the tropical Pacific. The mixed-layer (ML) temperature used for constructing Figure 15b is computed in an “off-line” fashion using predetermined surface flux data from the TOGA runs examined by *Lau and Nath [1994]*, with no consideration of the effects of such temperature changes on the overlying atmosphere. However, Figure 15c is based on the output from a new suite of four coupled GCM integrations (labeled by *Lau and Nath [1996]* as the “TOGA-ML” experiment), which incorporates a full dynamic exchange of information between the mixed layer and the atmosphere.

The patterns in Figure 15 indicate that both modeling approaches have considerable success in emulating the qualitative SST anomaly pattern in the extratropical northern oceans observed during El Niño events (Figure 15a). The realistic simulation depicted in Fig-

ure 15b confirms that this SST pattern can largely be attributed to the oceanic response to atmospheric driving. Comparison between Figure 15b and 15c reveals that the incorporation of feedbacks from the ocean to the atmosphere in the TOGA-ML runs does not significantly alter the local SST response. The notably lower amplitudes of the simulated SST anomalies (Figures 15b-15c) as compared to the observed values (Figure 15a) are likely a consequence of the relatively weaker extratropical responses to tropical SST forcing in the model atmosphere, as pointed out in section 4.2 (see Figure 12).

Additional diagnosis [Lau and Nath, 1996] indicates that the presence of two-way air-sea coupling in the TOGA-ML experiment leads to a noticeable enhancement of the atmospheric signal. This finding suggests that the SST response to atmospheric driving in mid-latitudes could in turn affect the overlying atmosphere through positive feedback processes, as hypothesized by Trenberth and Hurrell [1994].

It is anticipated that further progress in understanding extratropical air-sea coupling can be made by improving the highly idealized mixed-layer treatment of the surface ocean. For instance, the incorporation of seasonal variations in the thermocline processes at the bottom of the mixed layer would link the near-surface changes to heat storage in the deeper ocean. For atmosphere-ocean changes on timescales of a decade and longer, other low-frequency processes involving transports by anomalous ocean currents might also need to be taken into consideration [Trenberth and Hurrell, 1994; Deser et al., 1996].

5.2. Interactions With Land

It is well established that ENSO events lead to changes in rainfall all over the globe (Figure 2), often producing droughts in some regions and floods in other regions. The regions most affected are in the tropics and subtropics. These changes in rainfall clearly influence the availability of surface water over land, and in particular, the soil moisture and transpiration in vegetation is changed, providing a potentially powerful feedback on the atmospheric circulation.

The importance of surface hydrology in the low-frequency variability in the Asian monsoon has been postulated by Webster [1983]. The influence of snow cover and soil moisture on the Asian monsoon has been subject to considerable study and is one domain where effects seem to be important. Barnett et al. [1989] and Yasunari et al. [1991] showed with climate model experiments how anomalous snowfall over Asia could influence the subsequent summer monsoon, with heavier snow causing a delay in continental heating. Moreover, there are indications that the subsequent changes in the surface winds over the Pacific could possibly trigger an El Niño event.

Over North America the 1988 spring-summer drought has been linked to the strong La Niña and related tropical SST anomalies at that time, but with dry soil conditions and low-moisture availability contributing to

heat waves in ways that are likely to provide important feedbacks [Trenberth et al., 1988; Trenberth and Branstator, 1992; Atlas et al., 1993]. The North American spring-summer floods in 1993 during mature El Niño conditions is another case, this time of positive SST anomalies in the tropical Pacific. Trenberth and Guillemot [1996] show the influence of the tropics on changes in storm tracks, moisture sources, and transient eddy moisture transports that contributed to the North American floods which in turn altered the Bowen ratio to be close to its potential value as the surface became saturated in the upper Mississippi Basin. The recycling of moisture diminishes the sensible heating at the expense of increased evapotranspiration, and these effects combine with changes in atmospheric heating through altered latent heating and increased cloudiness to affect the atmospheric circulation and subsequent moisture transports. Preconditioning of the early spring soil moisture to dry conditions, as in 1988, or moist conditions, as in 1993, are important in the beginning stages [Dirmeyer and Shukla, 1993]. These considerations suggest the potential for significant feedback effects, most likely in the spring and summer, but their importance remains to be quantified.

6. The Future

The main thrust of this review has been on the improved understanding gained throughout TOGA of the extratropical response to tropical atmospheric forcing. Considerable progress has been made, but much of that progress consisted of unearthing the full complexity of the problem and the difficulties in capitalizing on the predictability that exists.

From realization of the complexity of the problem we now know that several issues must be addressed if the midlatitude response to tropical forcing is to be successfully represented and predicted. The first issue is one of obtaining the tropical atmospheric forcings, driven by anomalous SSTs. The second is one of determining reliably the response of the atmosphere. The processes involved in transforming the thermally forced internal mode in the tropics into an external, equivalent barotropic mode in the extratropics are not fully explained; neither are the mechanisms that produce extratropical convergence anomalies associated with the Rossby wave source. The response also involves sorting out the signal from the high level of natural variability that gives rise to noise and properly accounting for many and varied feedbacks. The climatological mean planetary waves create an environment favorable for producing internal sources of the energy that can compete in magnitude with the original perturbation. Further research into the role of nonlinearity is needed. Changes in the planetary waves produce changes in jet streams and storm tracks, and the latter feedback to alter and often reinforce the initial response. But such a feedback only applies on average, and limited periods can behave quite differently according to the synoptic situation. A second major feedback arises from

subsequent changes in extratropical oceans and SSTs (section 5.1), which in turn also influence the evolution of the atmosphere. A third major feedback occurs over land, especially in spring and summer, where changes in storm tracks alter the hydrological cycle, and the distribution of floods and droughts, and thus feedbacks from land where processes such as those embodied in the Bowen ratio (of sensible to latent heating) are altered. Changes in heating over land can also further influence moisture convergence and circulation.

All of these feedbacks and processes can in principle be modeled. As noted above, AGCMs can be very useful, but it is difficult to separate out the signal from the weather noise, and ensembles must be generated. Moreover, the models are not perfect so that results concerning predictability are model dependent. *Trenberth and Branstator* [1992] discuss the criteria for performance of these models if the simulations and/or predictions are to be useful. To be able to properly determine the total extratropical response and thus make forecasts, a number of modeling challenges exist. On the basis of our current understanding of the problem, as outlined in this paper, a modified set of criteria include the need to be able to simulate the following with fidelity: (1) the mean flow, which includes tropical rainfall, the ITCZ, SPCZ, and major monsoonal regions whose movements determine the tropical heating anomalies; vertical heating profiles, upper tropospheric divergences, and thus the Rossby wave sources; climatological winds, jet streams, and thus the environment through which Rossby waves must propagate; (2) a realistic annual cycle; (3) the anomalous SSTs and associated tropical forcing to determine the anomalous forcing of the atmosphere; (4) the temporal character of anomalous tropical forcing; (5) the changes in midlatitude storm tracks and associated rainfall (which requires adequate model resolution to replicate the processes involved), which provide important feedbacks; (6) the changes in extratropical SST and mixed-layer ocean conditions and their interactions with the atmosphere, all of which also provide important feedbacks; and (7) the changes in land surface conditions, soil moisture, vegetation (albedo, moisture stress, etc.), moisture transports, the hydrological cycle (which requires good resolution of mountains and orographic influences), and a realistic diurnal cycle, which also influence important feedbacks through land surface processes.

The above features highlight factors of importance in the extratropical teleconnections. Consequently, they provide a list of quantities and/or processes that are often worthwhile to examine diagnostically to determine what is going on, whether in the global observed data sets or model results. All models are deficient in various ways. Improved models will allow the predictability that does exist to be better exploited. In addition, better theoretical understanding of such issues is needed to enhance our confidence that the level of sensitivity of models also occurs in nature.

Because predictability depends on both signal and

the noise, it is important to analyze the dependence of the internal variability on SST forcing. On 20–60 day timescales, monsoon break and active periods occur and large Madden-Julian Oscillations perturb the tropical atmosphere but are not closely tied to SSTs, thereby creating one form of noise. It is not yet clear how much changes in the location and magnitude of the tropical heating anomalies account for the variations of the extratropical circulation during the ENSO cycle. On one hand, there are different “flavors” of El Niño that alter the location and magnitude of tropical forcing of the atmosphere, leading to expectations from linear theory and the protomodel for the extratropical response to be controlled by the position and structure of the tropical heat source. On the other hand, relative to the protomodel, changes in the actual anomalous Rossby wave source and Rossby wave propagation are diminished in the northern hemisphere owing to the presence of climatological planetary waves and excitation of preferred modes. Nevertheless, it does seem to be established that it is only when the expected anomalies of SST are large in the tropics that a reliable forecast can be expected for the extratropics. Even then, the influence of the tropical forcing depends on the time of year and is much less in the boreal summer (section 3.7).

All of these aspects highlight the potential for improved short-term climate forecasts in the extratropics, although with the realization that predictability is somewhat limited, so that a major challenge will be to utilize the uncertain forecast information in the best way possible throughout different sectors of society. The need for regional specificity presents a particular challenge. Both the applications communities and the social science communities require regional information that is as detailed as possible. Yet predictability inherently becomes less on smaller space and timescales because the natural variability is larger.

Some of the research focus in the post-TOGA era has turned to investigation of ways for the more useful application of these teleconnection relationships to practical seasonal predictions. One approach [*Ropelewski and Halpert*, 1996] is to document conditional climatologies of seasonal rainfall amounts as a function of SO phase. There is also interest in the examination of the modulation in daily rainfall frequency and intensity as a function of SO phase. The firm documentation of the ENSO-precipitation relationships has also inspired interest in the more general problem of the relationships between the large-scale boundary conditions, both SST and land surface, and the atmospheric circulation. Among others it motivated the Pan-American Climate Studies (PACS) program to examine “warm season” precipitation in the Americas and its relationships to SST as well as to the land surface. Feedbacks from land surface processes are an important process not adequately addressed in TOGA, but these are high on the agenda for the follow-on to TOGA, namely GOALS (the Global Ocean-Atmosphere-Land System program), and GEWEX (Global Energy and Water

Cycle Experiment). Links to surface hydrology will need to be explored. The international framework for GOALS under the World Climate Research Program is the Climate Variability and Predictability program (CLIVAR). Also, under the CLIVAR umbrella is a focus on longer timescales, and an emerging high profile topic in need of further investigation is the decadal and longer variations in ENSO and its associated teleconnections.

Acknowledgments. This review is partly sponsored by the NOAA Climate and Global Change Program. Special thanks to Gerald Bell and Michael Halpert (Figure 1), Doug Lindholm (Figure 3), Jim Hurrell (Figure 7), and Chris Guillemot for Figure 11. The National Center for Atmospheric Research is sponsored by the National Science Foundation.

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(Received April 4, 1996; revised February 17, 1997; accepted May 15, 1997.)