

Hydrodynamics

## Section



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## Propeller

Lifting-Surface Corrections

## BY

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# Propeller Lifting-Surface Corrections 

by

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## TABLE OF CONTENTS

Page
Introduction ..... 1
Theoretical Background ..... 3
Presentation of the Serics ..... 8
Experimental Checks on the Theory ..... 20.
Conclusions ..... 27.
Acknowleagments ..... 28
References ..... 28

## PREFACE

During the summer of 1967 Dr, Wm. B. Morgan: Naval Architect, Hydromechanic Laboratory, Naval Ship Research and Development Center, Washington, D. C., spent three months at the Hydro- and Aeredynamics Laboratory. Although his stay was short it become very fruictul. Besides giving lectures on Propeller Design at the Technical University of Demmark he found time to prepare in co-operation with Dr. V. Silovic, Head of Propeller Section at Hy $A$, the present paper on Propeller Lifting Surface Corrections, which was finalished after Dr. Morgan's return to the N.S.R.D.C. with the further assistance of Mr, Stephen B. Denny, mathematician, N.S.R.D.C.
The Hydro- and Aerodynamics Laboratory is grateful to the Naval Ship Research and Development Center for having made this co-operation possible, also to the Applied Mathematics Laboratory of N.S.R.D.C., and to the Northern Europe University Computing Centre, Lyngby, Denmark, for assistance in the numerical calculations.

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# Propeller Lifting-Surface Corrections 

By Wm. B. Margan, ${ }^{1}$ Member, Vladimir Silovic, ${ }^{2}$ Member, and Stephen B. Denny, ${ }^{3}$ Visifor


#### Abstract

'Correction factors for camber, ideal angle due to toading, and ideal angle due to thickness, whicn are based on propeller lifting surface theory, are presanted for a series of propellers. This series consists of optimum free-running propellers with chordwise loadings the same os an NACA a $=0.8$ mean line and with NACA- 66 chordiwise thickness distributions. The results of the calculations show that the threedimensional camber and ideal angle are generally greater than the two-dimensional camber and ideal angle af the same lift coefficient. The correction factors increase with increasing expanded arsa ratio, and those for camber and ideal angle due to loading decrease with increasing number of blades. Thickness, in general, induces a positiva angle to the flow, which necessitates a correction to the blade pitch. This ideal angle is largest near the blade root and decreases to negligible values toward the blade fip and increases with increasing number of blades. Skew induces on inflow angle, necessitating a pitch change which is positive toward the blade root and negative toward the blade tip.


## Introduction

The trend in ships has been toward increasing speed, size, and horsepower. As a result, there is an increasing demand for the design of propellers which are efficient and yet produce minimum cavitation and induced vibration. Such a task requires a knowledge of the flow field in which a propeller operates and an accurate determination of the flow over the propeller blade surfaces.
Much effort has been devoted recently to the development of a more sophisticated propeller theory. This theory has proceeded from the simple momentum concepts of Rankine [1] ${ }^{\text { }}$ (1S65) to the lifting-line model of Goldstein [2] (1929) and, finally, to the lifting-surface model of

[^0]Ludwieg and Ginzel [3] (1944). More recently, because of the numerical evaluations possible with high-speed computers, emphasis has been placed on more sophisticated mathematical models than hitherto possible; e.g., Sparenberg [4]. Cox [5], Pien [6], and Kerwin [7].
A complete review of the various liftingsurface theories will not be given here since adequate reviews for work through 1904 are given by Wu [ $\$$ ], Isay [ 9 ], and Lerbs et al. [10]. Recent work not mentioned in these three references are studies by Cheng [11, 12], Kerwin and Leopold $[13]$, Nelson [14, 15], Sulmont $[16,17]$, Nishiyama and Sasajima [18], Malavard and Sulmont [19], and Murray [20]. Murray develops the lifting-surface theory of both contrarotating and conventional propellers.
In general, the investigators have dealt with the same basic mathematical model but have used different assumptions to obtain numerical solutions. The propeller-blade boundary conditions are linearized and the blade and its helical wake are replaced by vortex systems. Assumptions are made for the chordwise and spanwise loading on the blade and a complicated singular integral is derived from which the necessary distortion
to the blade can be calculated to obtain the prescribed loading. In addition to determining loading effects, Ketwin and Leopold [21], Nelson [14], and Murray [20] also considered the effect of blade thickness, in the linearized sense, and found that is contributed significantly to the required blade distortion.

For an unskewed propeller blade, the effect of loading on the blade shape is to require both a larger camber and a larger ideal angle than would be required is two-dimensional flow in order to produce the same lift. The principal effect of thickness is to distort the flow such that an increase in angle of the blade is required to maintain the desired loading. Likewise, the principal effect of skew is to necessitate a blade angle change but to require little cinange in camber.

Because of the lincarized boundiary conditions, camber and ideai angle corrections, which are independent of the magnitude of the propeller loading [10], can be derived by taking a ratio of the three-dimensional camber and ideal angle to the two-dimensional camber and ideal angle, respectively. These correction factors are, of course, dependent upon the chordwise and spanwise load (or pitch) distributions, number of blades, and blade area and shape. Similarly, a thickness correction factor can be derived which is independent of the magnitude of the thickness but is, of course, dependent apon the chordwise and spanwise thickness distributions.

Since "iting-surface correction factors can be derived which are independent of the magnitude of the propeller loading and thickness, a system-
$A_{5}=$ blate expanded wea
d, $\approx$ propeller disk area
$B T F=$ bade-thickness fraction
$C(r)=$ ceeflicient of blade outline

$$
C_{L} \approx \text { lift coeficient, } \frac{2 x G(r) V D}{c V}
$$

$C_{\text {max }}=$ two-dimensional maxinum mean-line ordinate fo: $C_{L}=1.0$
$C_{r}=$ power coeflicient, $\frac{P_{D}}{\frac{P}{2} A_{0} V_{A}{ }^{3}}$
$C_{\mathrm{Th}}=$ thrust conficient, $\frac{T}{\frac{\rho}{2} A_{0} V_{A}{ }^{2}}$
$r=$ radial cosrdisate nomdimensionaized by propeller radius
$r_{n}=$ dimensionless hub radius
$(r, \theta, s)=$ dimensionless cylindirical coordinates
$T=$ propeller thrust
$f\left({ }_{f}, x_{c}\right)=$ half-thickness ordinate
$t_{\text {maxa }}(r)=$ maxinum $\cdot$ thickness ordinate
$U(r)=$ resultant induced velocity from lifting-line theory
$U_{n}(r)=$ velocity induced by loading normal to bladesection chord
$U_{n}(r)=$ velocity induced by thickness normal to blade-section chord
$U_{\tau}(r)=$ tangential induced velocity
$V=$ ship speed
$V_{A}(r)=$ speed of advance at a given radius
$V_{t}(P)=$ induced velocity vector
$V^{\prime}(r)=$ resultant section infow velocity
$w(r)=$ circumferential mean-wake coeflicient at a given radius
$(x, y, z)=$ dimensionless cartesian coordinates
$x_{f}=$ chordwise section abscissa nondimensionalized by chord
$x_{\text {max }} \neg$ chordwise position of maximum camber
$Z=$ number of blades
$\alpha_{1}(r)=$ section ideal angle of attack
$a_{4.0}=$ twodimensional ideal angle of attack for $C_{L}=1.0$
$\alpha_{1}(r)=$ angle-of-attack correction from thekness
$\beta_{1}(r)=$ hydrodynamic pitch angle
$\lambda_{1}=$ induced advance coeflicient $=r$ tan $\beta_{1}$
$\lambda_{0}=$ apparent advance coefficient, $\frac{V}{\pi n I)}$
$\sigma(r)=$, source strength
$\theta_{1}(r)=$ angular position of blade-section leatins edge
$\theta_{1}=$ skew at blade tip, deg
$\theta_{1}(r)=$ angular position of blade-section trailing edge
atic series of these correction factors can be calculated for design use. By varying such parameters as blade number, blade area, blade skew. and hydrodynamic pitch angle, the effect of these parameters for a practical blede shape can be ascertained and a better understanding of the flow over the blade will result. These correction factors can be utilized in propeller design to obtain the blade pitch and camber without empirical adjustments. Design methods-icveloped over the years-based on empirical correction factors and on correction factors from simplified mathematical models should be regarded as outmoded. ${ }^{\circ}$ In addition to providing the designer with coefficients with which he can design propellers by modern theoretical methods, these correction factors can be used for making qualitative checks of lifting-surface computer calculations, provided that the prepeller geometry and loading do not differ radically from the series.

Lifting-surface correction factors are presented for four, five, and six blades, for blade area ratios from 0.35 to 1.15 , for hydrodynamic pitch ratios from 0.4 to 9.0 , and for three blade skews. The calculations are for optimum ${ }^{6}$ free-running propellers and the results are similar to those presented by Lerbs [10] but with two essential differences. One is that the assumed chordwise load distribution is that of the NACA $a=0.8$ mean iine: i.e., constant loading from the leading edge to 0.8 chord and then a constant slope to zero at the trailing edge. The second difference is the inclusion of a correction for thickness. Experimental results have shown that these differences are significant in applying liftingsurface theory to practical propeller designs. The thickness correction must be included if the propeller is to have sufficient pitch, and the NACA $0=0 . S$ mean line is a mean line which achieves, approximately, its theoretical lift in viscous flow. For example, the NACA $a=1.0$ mean line achieves only $7 t$ percent of its theoretical lift in a real fluid, and it is not feasible in the present lifting-surface theory to account for viscous effects on lift. In analogy with two-dimensional results, the $\mathcal{N A C A} a=0 . S$ load distribution requires an ideal angle for the mean line to operate at shock-free entry.

[^1]The following sections of the paper will review briefly the theory used for the calculations and outline the assumptions and limitations. Results of the calculations will be discussed. Some comparisons between experiment and theory will be presented and conclusions will be drawn from the apparent trends as to the applicability of the theory.

## Theoretical Background

## Statement of the Problem

The problem can be stated as follows: "For a given number, loading, area, skew, chord distribution, and thickness distribution, determine the required blade-section camber and ideal angle." This problem is similar to the inverse problem of airfoil theory, except that the thickness is specified.

## Assumptions

The brief theoretical background presented here will serve only to bring out the salient features of the procedures used. Details of the theories involved are available in references [12] and [21].
In the mathematical model for loading, a distribution of bound vortices is assumed to cover the blades, and free vortices are shed from these bound vortices downstream along helical paths. For thickness, a network of sources and sinks is assumed to be distributed over the blades. The following assumptions based on this mathematical model are generally made:

1 The fluid is inviscid and incompressible.
2 The free-stream velocity is axisymmetric and steady, allowing the propeller to be wakeadapted.
3 Each propeller blade is replaced by a distribution of bound vortices for loading effects, and sources and sinks for thickness effects. The circulation is distributed in both the chordwise and spanwise directions. It follows from vortex theory that free vortices are shed from the bound vortices and, in a coordinate system which rorates with the propeller, these free vortices form a general helical surface behind the propeller.
4 Each of the free vortices has a constant diameter and a constant pitch in the downstream direction, but the pitch may vary in the radial direction. This means that effects of slipstream contraction and centrifugal force on the shape of the free-vortex sheets are ignored.
5 The boundary conditions on the blade are linearized, which implies that the lifting surface


Fig. 1 Propeller velocity diagram
has only a small deviation from the hydrodynamic pitch. This assumption is similar to the linearized theory of two-dimensional airfuils where the boundary condition is not satisfied on the profile but on the profile chord. Also, the linearization enables separation of the loading and thickness effects.
6 The pitch of the blade and of the trailing vortex sheets is the hydrodynamic pitch obtained from lifting-line theory. Within the context of the linearized boundary conditions, this is not an assumption-except for skewed propellerssince the lifting-line theory does not account for skew.

7 The hub is assumed to be small enough that it is not necessary to satisfy the hub boundary condition.

8 Blade rake is not considered.
The assumptions listed apply to moderately loaded propeller theory. Assumption 4 would have to be removed for heavily loaded propellers, and an additional assumption would have to be made for lightly loaded propellers; namely, that the effect of the induced velocities on the pitch of the vortex sheets is negligible.

## Locding

Fig. 1 shows the velocity diagram for a propeller blade section in the absence of thickness. From this figure and within the concepts of linearized theory, the boundary condition at each section is

$$
\begin{equation*}
\alpha_{1}(r)+\frac{\partial f_{p}\left(r, x_{c}\right)}{\partial x_{c}}=\frac{U_{n}}{V_{r}}\left(r, x_{c}\right)-\frac{U}{V_{r}}(r) \tag{1}
\end{equation*}
$$

where $f_{p}$ is the camber along the chord $x_{c}, \alpha_{1}$ is the ideal angle of attack $\left(\alpha_{1} \simeq \tan \alpha_{1}\right), V_{r}$ is the re-


Fig. 2 Coordinate system
sultant inflow velocity to the blade section, $U$ is the resultant induced velocity from lifting-line theory, and $U_{n}$ is the induced velocity normal to the blade-section chord. The angle $\beta_{i}$ shown in Fig. 1 is the hydrodynamic pitch angle.

Both a carce: ian coordinate system ( $x, y, z$ ) and a cylindrical system ( $r, \theta, z$ ) will be used (see Fig. 2). The bound circulation will be assumed to have the strength $G_{r}(r)$ such that

$$
\begin{equation*}
\int_{\theta_{1}(r)}^{\theta_{1}(r)} G_{r}(r, \theta) d \theta=G(r) \tag{2}
\end{equation*}
$$

The coordinates $\theta_{1}(r)$ and $\theta_{1}(r)$ define the angular position of the blade-section leading and trailing edges, respectively, and $G(r)$ is the nondimensionalized bound circulation as determined from lifting-line theory [23]. Since $G_{r}(r, \theta)$ may vary in both the radial and chordwise directions, a free vortex may be shed at each point on the blade surface and will trail bethind along a helix of pitch $\pi r \tan \beta_{t}\left(\tan \beta_{f}\right.$ from liftingline theory).
It follows from vortex theory that the strengths of the helical free vortices are equal to the negative of the radial rate of change of the bound circulation at the point the vortex is shed. If the strength of each helical free vortex is assiumed to be $G_{f}(r)$ behind the trailing edge, then

$$
\begin{equation*}
G_{f}(r)=-\frac{d G(r)}{d r} d r \tag{3}
\end{equation*}
$$

With this equation and equation (2), the strengths of the free vortices shed from the blade are found to be

$$
\begin{align*}
G_{r}(r)= & -\frac{d}{d r} \int_{\theta_{1}(r)}^{\theta_{1}(r)} G_{i}(r, \theta) d \theta d r \\
= & -\int_{\theta_{1}(r)}^{s_{1}(r)} \partial G_{r}(r, \theta) \\
& \quad \times \frac{d \theta}{\partial \theta_{2}} d \theta d r-G_{,}\left(r, e_{1}\right)  \tag{4}\\
& =G_{r}\left(r, \theta_{t}\right) \frac{d \theta_{1}}{d r} d r
\end{align*}
$$

The first term on the right corresponds to the integral of the wee vortex strengths from the leading edge to the trailing edge, which resulits from a radial change in the bound vortices. The two remaining terms on the right are strengths oi the free vortices shed along the blade outline. It follows from this equation that, within the lifting surface, the free vortices have strength

$$
\begin{align*}
G_{1}(r, \theta) & =-G_{1}(r, \theta) \frac{d \theta_{1}}{d r} d r \\
& -\int_{0}^{g_{1}(r)} \frac{\partial G_{r}}{\partial r}\left(r, \theta_{0}\right) d \theta_{0} d r, \partial_{1} \geq \theta>\theta_{1} \tag{5}
\end{align*}
$$

From the Biot-Savart law, the induced velocity $V_{i}$ at any point $I^{\prime}$ on the lifting surface is found to be

$$
\begin{align*}
& \frac{V_{1}(P)}{V}=-\frac{1}{2} \iint_{A_{1}} G_{1}(r, \theta)\left(\frac{S \times \mathrm{d} r}{S^{2}}\right) d \theta \\
& -\frac{1}{2} \iint_{A_{1}+A_{2}} G_{f}(r, \theta)\left(\frac{X \mathrm{~d} l}{S^{3}}\right) \tag{i}
\end{align*}
$$

where $A_{1}$ is the area of the lifting surface, $A_{3}$ is the area of the helical surface behind the trailing edge of the blade, $V^{\prime}$ is the freestream velocity or ship speed, $S$ is the vector distance frove a point $Q$ on the helical surface to the point $P$ also on the heliaal surface, $\mathrm{d} /$ is the elementary vector tangent to the vortex line, and $S=|\mathbf{S}|$. The second term in this equation can be expressed in two parts by equations (3) and (5), and the quantity

$$
\begin{aligned}
& 1 \\
& 2
\end{aligned} \iint_{A=} \frac{d G(r)}{d r}\left(\frac{\mathrm{~S} \times \mathrm{d} l}{S}\right) d r
$$

is now added to and subtracted from this equation: $A_{3}$ is the area between the trailing edge and a generating line along which a lifting line would be placed. A new coordinate system is introduced such that the point $P$ always lies on the lifting surface ${ }^{7}[6]$. This means that the integral

[^2]$$
\frac{1}{2} \iint_{A_{z+A}} \frac{d G(r)}{d r}\left(\frac{S \times \mathrm{d} l}{S^{5}}\right) d r \equiv\left(\frac{V_{t}}{V}\right)_{\mathrm{lilting} \text { line }}
$$
is identical to the velocity induced at a lifting line by the trailing vortex system. The velocity induced by the lifting surface itself is
\[

$$
\begin{align*}
\frac{V_{1}(P)}{V^{\prime}}= & -\frac{1}{2} \iint_{A_{t}} G_{r}(r, \theta)\left(\frac{S \times \mathrm{d} r}{S^{2}}\right) d \theta \\
& +\underset{V_{1}}{V_{i}}+\iint_{A_{1}}\left\{G_{r}(r, \theta) \frac{d \theta_{1}}{d r}\right. \\
+ & \left.\int_{0}^{\theta_{t}(r)} \frac{\partial G_{r}}{\partial r}\left(r, \theta_{G}\right) d \theta_{0}\right\}\left(\frac{\mathrm{S} \times \mathrm{d} l}{S}\right) d r \\
& -\iint_{A_{2}} \frac{d G}{d r}(r)\left(\frac{\mathrm{S} \times \mathrm{d} /}{S^{\mathrm{a}}}\right) d r \tag{7}
\end{align*}
$$
\]

This equation gives the velocity induced by a single blade at the point $P$. The total velocity induced by all the blades at the point $P$ on one of the blades is obtained by summing this equation with respect to the blade position $m$, of which S is a function. The velocity $\mathrm{V}_{l}\left(I^{\prime}\right)$ includes the radial induced velocity as well as the axial and tangential velocities. The radial velocity does not appear in the boundary condition, equation (1), and can be neglected. Then, the normal velocity is

$$
\mathrm{n} \cdot \frac{V_{i}}{V_{\text {Hitsing line }}} \equiv \frac{V}{V}-\frac{U}{V}
$$

and the equation which must be evaluated in determining $\partial f_{p} / \partial x_{r}$ is

$$
\begin{align*}
& V_{i}^{\prime}\left(\alpha_{1}+\frac{\partial f_{p}}{\partial r_{r}}\right)=\frac{l_{n}\left(r_{r}, x_{d}\right)}{V}-\frac{U(r)}{V^{r}}= \\
& -\frac{1}{2} n \cdot\left\{\sum_{m=1}^{Z} \iint_{A_{1}} C_{r}(r, \theta)\binom{S \times \mathrm{d} r}{S^{3}} d \theta\right. \\
& +\sum_{2}^{1} \sum_{m=1}^{2}\left[\iint_{A_{1}}\left\{G_{i}\left(r, \theta_{1}\right) \frac{d \theta_{3}}{d r}+\int_{0}^{\theta_{1}(r)} \partial G_{r}\left(r, \theta_{0}\right)\right\}\right. \\
& \left.\left.\times\binom{\mathrm{S} \times \mathrm{d} l}{\frac{S^{3}}{}} d r-\iint_{A_{3}} \frac{d G(r)}{d r}\binom{\mathrm{~S} \times \mathrm{d} l}{S^{3}} d r\right]\right\} \tag{B}
\end{align*}
$$

where $Z$ is the number of blades. The blade camber is given by

To determine $\alpha_{t}$, this integral is evaluated from the leading edge to the trailing edge and $f_{p}$ is taken to be zero.

The equation for the slope of the camber line, equation (9), avoids integration limits of infinity, but the double integrals are dificult to evaluate because the integrands are singular when $S$ becomes zero. 'To facilitate the numerical evaluation of these singular integrals, Pien [6] and Cheng [12] made a number of simplifications:
1 The radial bound circulation distribution $G(r)$ and the hydrodynamic pitch ratio distribution are expanded in a half-range cosine series with a finite number of terms.
2 The singular points of the integrands are isolated and the integrals are evaluated numerically. In the region of the singular point, the radial integration is done analytically by assuming that pitch and circulation are constant.
3 For calculating the blade mean line, the induced normal velocity is expressed as a power series.

The simplifications are numerical approximations and not assumptions made in the theory The precise accuracy of the evaluations is difficult to ascertain.

Correction factors for camber and ideal angle due to loading which are independent of the magnitude of propeller loading can be derived by normalizing the three-dimensional camber and ideal angle by the two-dimensional section values. With regard to camber, only the correction factors for the maximum chordwise camber will be formed in this way. From equation (9), then

$$
\begin{align*}
& k_{r}(r)=\frac{f_{D}\left(r, x_{\max }\right)}{f_{\max }}=\frac{1}{C_{\max } C_{L}} \\
& \quad \times\left\{\int_{0}^{x_{\max }}\left[\frac{U_{n}(r, x)}{V_{r}}-\frac{U(r)}{V_{r}}\right] d x-x_{\max } \alpha_{l}(r)\right\} \tag{10}
\end{align*}
$$

where $x_{\text {wax }}$ is the chordwise position of maximum camber, $C_{L}$ is the blade-section lift coefficient, and $f_{\text {max }}$ is the two-dimensional maximum camber and equal to

$$
\begin{equation*}
f_{\text {max }}=C_{\text {max }} C_{L} \tag{11}
\end{equation*}
$$

The function $C_{\text {max }}$ is the maximum mean-line ordinate for $C_{L}=1.0$ in two-dimensional flow. For other chordwise positions, the three-dimensional cambers will be normalized by the maximum camber; i.e.,

$$
\begin{equation*}
f\left(r, x_{c}\right)=\frac{f_{p}\left(r, x_{c}\right)}{f_{p}\left(r, x_{\max }\right)} \tag{12}
\end{equation*}
$$

Correction factors for the ideal angle due to loading can be formed in the same way as the camber correction factor.

$$
\begin{align*}
k_{\mathrm{a}}(r)= & \frac{\alpha_{1}(r)}{\alpha_{1, .} C_{L}}=\frac{1}{\alpha_{L_{1} .0} C_{L}} \\
& \times \int_{0}^{1}\left[\frac{U_{n}}{V_{r}}(r, x)-\frac{U(r)}{V_{r}}\right] d x \tag{13}
\end{align*}
$$

where $\alpha_{t .0}$ is the two-dimensional ideal angle of attack for $C_{L}=1.0$ for the NACA $a=0.8$ mean line.

## Thickness

Blade thickness effects are determined by introducing a source-sink system distributed over the liiting surface [21]. The induced velocity field is given by

$$
\begin{align*}
\left(\frac{\mathbf{V}_{1}(P)}{V}\right)_{\text {thickness }} & =\int_{r_{h}}^{1} \int_{\theta_{1}(r)}^{\theta_{1}(r)} \sigma\left(r_{0}, \theta_{0}\right) \\
& \times \mathrm{H}_{s}\left(P, r_{0}, \theta_{0}\right)\left|\frac{\partial x_{c}}{\partial \theta_{0}}\right| d \theta_{0} d r_{0} \tag{14}
\end{align*}
$$

where H , is the velocity induced at the point $P$ by a unit source located at point $r_{1} \theta$; i.e.

$$
\mathbf{H}_{s}=\sum_{m=1}^{Z} \operatorname{grad}\left(\frac{-1}{4 \pi S}\right)
$$

and where $d_{i}{ }_{c}$ is the element of blade chord, and $\sigma(r, \theta)$ is the source strength. It is assumed that the source strength distribution is known and is that derived by the usual linear approximations from airfoil theory; i.e.

$$
\begin{align*}
\sigma\left(r_{0}, \theta_{0}\right)=\frac{V_{r}}{V} & \frac{\partial l\left(r_{0}, x_{c}\right)}{\partial x_{c}} \\
& =\left(\frac{r_{0}}{\lambda_{s}}-\frac{U_{r}}{V}\right) \frac{1}{\cos \beta_{i}} \frac{\partial\left(\left(r_{0}, x_{c}\right)\right.}{\partial x_{c}} \tag{15}
\end{align*}
$$

where $t\left(r_{0}, \lambda_{c}\right)$ is the thickness distributed along the chord $x_{c}, U_{T}$ is the tangential induced velocity, and $\lambda_{s}$ is the apparent advance ratio of the propeller; i.e., $V / \pi n D$.
As for the effect of loading, only the axial and tangential induced velocity components are required from equation (14). If the change in camber along the chord $x_{c}$ due to thickness is $f_{p_{s}}\left(r, x_{c}\right)$, and the ideal angle induced by thickness is $\alpha_{t}(r)$, the linearized boundary condition on the blade is

$$
\begin{equation*}
\alpha_{l}(r)+\frac{\partial f_{p_{i}}}{\partial x_{c}}\left(r, x_{c}\right)=\frac{U_{n_{t}}}{V_{r}}\left(r, x_{c}\right) \tag{i6}
\end{equation*}
$$

where $U_{n_{t}}$ is the induced velocity normal to the section chord from equation (14) and where $\alpha_{t}(r)=\tan \alpha_{t}(r)$. Equation (16) is similar to equation (1), and the change in curvature due to thickness is found by integrating equation (16) in the manner of equation (9).

To facilitate evaluating the integrals of equation (14), Kerwin and Leopold \{211 used a distribution of quadrilateral source elements over the blades. This approach is similar to that of Hess aml Smith [ 24 ] for potential flow about arbitrary borlies, except that the elements are placed along the chord rather than over the surface. Therefore, the source-sink distribution is not continuous but is made up of discrete elements. If a sufficient number of elements are used to describe the sarface, the accuracy should be good.

Calculations have shown that thickness induces an infow angle and a camber change but that camber alteration is small [ 21 ] except for very small pitch ratios. This flow distortion requires an increase in blade angle (ideal angle) to maintain the desired loading. Correction factors for this ideal angle can be made independent of the magnitude of thickness by dividing the angle (in radians) by the blade thickness fraction for which the calculations were made; e.g.

$$
\begin{equation*}
k_{1}(r)=\frac{1}{B T V} \int_{0}^{1} \frac{U_{n_{1}}}{V_{r}}\left(r, x_{c}\right) d x_{c} \tag{17}
\end{equation*}
$$

where $B T F$ is the blade thickness fraction. A camber correction due to thickness can by formed in a similar manner, but since the correction is small it will be ignored. It should be noted from equation (17) that $k_{t}(r)$ is also a function of the propeller loading since $V^{\prime}$, is a function of the loading. This effect is small, however, and no correction will be made in the results which are presented.

## Problems in Numerical Evaluation

Because of their complexity, the expressions arising in the lifting surface calculation procedure described in this paper were necessarily evaluated by digital computation. This resulted in repeated numerical operations to obtain discrete values of the functions involved and the occurrence of inaccuracies, which were dependent upon the characteristics of the functions themselves and were not always predictable by the program user.

The Cheng computer program, described in reference [12], offers many input and output data options which are helpful in describing propellers with special geometries and loadings. This flexibility, however, can lead to inaccuracies in the calculations. Points which merit special attention in both the operation of the program and the analysis of the results are:

1 The computer program, as listed in Appendix $B$ of reference [12], must be modified slightly to give correct results. The convergence criteria
cited in Fortran statements 201 to 204 of Subroutine SUB3 should rad 0.9909 and 0.0001 respectively, rather than 0.99 and 0.01 . This oversight results in considerable differences in calcuiated camber values when compared io results of the former program [11] for constant chordwise loading cases.

2 Initial input to the Cheng program is the designation of a fan angle and grid spacing which govern the location and number of control points. The fan angle input must contain the projected view of the propeller blade, and the grid spacing in that fan should be as fine as possible for the greatest accuracy in flow calculation at the control points. The total number of chordwise stations on either side of the control point, however, cannot exceed 90. La order to meet these restrictions, the angle and grid spacing should be carefully checked before each propeller calculation. and particular care should be taken for large blade areas and highly skewed propellers. No automatic checks exist in the program to assume this task.

3 A Lost critical of the input parameters is the specification of chord-wise coordinates at which the camber and induced velocities will be calculated. Considerable differences in calculated induced velocities appear for relatively small abscissa changes along the chord near the leading and trailing edges. The inclusion of these changes in the resulting integrals yield rather small percent changes in camber but relatively large percent changes in ideal angle. The series discussed in this paper was calculated with chordwise coordinates specified from 5.0 to 95.0 -percent chord. Questionable ideal angle and camber corrections arose occasionally for small blade areas and/or low pitch ratios. A change of chordiwise coordinates to a range from 2.0 to 980 percent chord improved upon the erratic results in these regions, but, for consistency in the series data, these calculated values are not presented. It should be noted also that the questionable data were obvious only in comparisons to other data in the series, alone they perhaps would have gone umoticed. Independent propeller calculations should be checked against the series presented in this paper or other similar designs if possible.

4 Lifting-surface correction data are given herein only at radial positions $r / R=0.3$ through $r / R=0.9$ for $0.2 R$ hub propellers. The liftingsurface calcu'? ions were made also at $r / R=025$ and $r / R=0.95$. However, the result were questionable which is probably due to a combination of numerical difficulties involving the singularities at the blade hub and tip.


Fig. 3 Comparison of camber correction facrors for three blades $\pi \lambda_{1}=1.0472$, and $A_{5} / A_{0}=0.75$

## Prasentation of the Series

## Procechure

The propellers for the series were free-running propellers of constant hydrodynamic pitch with : F hub diameter of 0.2 of the propeller diameter. Lifting-line calculations were from a computer program based on the induction-factor method of Lerbs [25]. With loading and pitch distributions available from the lifting-line calculations, lifting-surface calculations for loading were made using the program developed by Cheng [12] and for thickness using the program developed by Kerwin and Leopold [21]. These have been combined into one program at NSRDC for running on the IBM-7090 computer.

Because of the complicated nature of the numerical procedures, it is difficult to state the accuracy of the results. To this end, calculations were made and compared in Fig. 3 to those of Cox [5], Kerwin [13], and Lerbs et al. [10] for a three-bladed propeller with an induced advance coefficient ${ }^{8} \lambda_{1}$ of 0.3333 , expanded area ratio of 0.75 , and a constant chordwise load distribution. The camber correction factors, $k_{c}(r)$, calculated using results of the Cheng method are within 1 percent of those derived from the Lerbs method for radii between 0.4 and 0.8 but deviate considerabiy at the 0.3 and 0.9 radii. Both the Cox and the Kerwin methods give correction factors which are, in general, higher than those obtained by the Lerbs and the Cheng methods. These comparisons are consistent with the way Cox and Kerwin arrive at their camber correction factors [10].

[^3]Since the calculations are based on similar procedures, the reason for the difference ber ween the results of Lerbs and Cheng is nat known. The difference toward the blade tip is probably caused by different numerical techniques, and the difference toward the root is mainly due to the liftingline calculations; i.e., induction factors versus Goldstein factors.

For a further comparison of lifting-surface calculations, the camber correction factor for a chordwise load distribution corresponding to a NACA $a=0.8$ mean line is also plotted in Fig. 3. It is interesting to now that the camber correction factor for this mean line is everywhere smalier than the corresponding factor for the constant load distribution. The three-dimensional camber is higher, however, since $C_{\text {max }}$ is approximately 23 percent larger for the $a=0.8$ than for the $a=$ 1.0 mean line.

Lifting-surface corrections were calculated for propellers with four, five, and six blades, with expanded area ratios from 0.35 to 1.15 , with hydrodynamic pitch ratios of 0.4 to 2.0 . and with a symmetrical outline and skew angles of 7 , $1 \cdot 4$, and 21 deg. The skew angle $\theta_{s}$ is defined as the angle between two straight lines in the projected plane-one from the shaft centerline through the midchord at the root section and the other

Table 1 Ordinates for NACA 66 (Mod) Thickness Distribution and NACA a $=0.8$ Camber Distribution

| Station, $x_{r}$ | Thickness Ordinate, $t / t_{\text {max }}$ | Camber Ordinate, $f / f_{\text {max }}$ |
| :---: | :---: | :---: |
| $\stackrel{0}{0.005}$ | $\begin{gathered} 0 \\ 0.066 \overline{5} \end{gathered}$ | $\begin{array}{ll} 0 \\ 0 & 0423 \end{array}$ |
| 0.005 0.0075 | 0.00812 | 0.0595 |
| 0.0125 | 0.1044 | 0.0907 |
| 0.025 | 0.1466 | 0.1586 |
| 0.05 | 0.2068 | 0.2712 |
| 0.075 | 0.2525 | 0.3657 |
| 0.1 | 0.2907 | 0.4482 |
| 0.15 | 0.3521 | 0.5869 |
| 0.2 | 0.4000 | 0.6993 |
| 0.25 | 0.4363 | 0.7905 |
| 0.3 | 0.4637 | 0.8635 |
| 0.35 | 0.4832 | 0.9202 |
| 0.4 | 0.4952 | 0.9815 |
| 0.45 | 0.5 | 00881 |
| 0.5 | 0.4962 | 10 |
| 0.55 | 04846 | 0.9971 |
| 0.6 | 0.4653 | 0.9786 |
| 0.65 | 0.4383 | 0.9434 |
| 0.7 | 04035 | 0.8892 |
| 0.75 | 0.3612 | 0.8121 |
| 0.8 | 0.3110 | 0.7027 |
| 0.85 | 0.2539 | 0) 5425 |
| 0.9 | 0.1877 | 0.3586 |
| 0.95 | 0.1143 | 0.1713 |
| 0.975 | 0.748 | 0.0823 |
| 1.0 | 0.0333 | 0 |

Table 2 Distribution of Elade Chord

| $t$ | $C(r)$ |
| :---: | :---: |
| 0.3 | 10338 |
| 0.3 | 18082 |
| 04 | 19848 |
| 05 | 2.0967 |
| 0.0 | -1020 |
| $0 \%$ | 22350 |
| 1) 8 | $\bigcirc 1719$ |
| 0.9 | ( 8131 |
| ${ }^{6} 95$ | 153 L 2 |
| 10 | 0 |

Yable 3 Skew Distribution for 14 Kag Skew and $\pi \lambda_{8}$ of 1.2

Skew R
0 0
0
0

| 0.3 | 00037 |
| :---: | :---: |
| $0{ }_{0}$ | $0^{0} 0148$ |
| 0.5 | 00336 |
| 0.6 | 00604 |
| 07 | $0.095 \%$ |
| 0.8 | () 1402 |
| 0.9 | 0189 |
| 10 | 0.2016 |

Table 4 Calculated Propellers with Symmetrical Blade Outline for Four, Five, and Six Blades

| $\pi \lambda_{i}$ | 0.4 | 0.8 | 1.2 | 1.6 | 2.0 |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $\frac{A S}{}$ | 0.35 | 0.35 | 0.35 | 0.35 | 0.35 |
| $\lambda_{0}$ | 0.55 | 0.55 | 0.55 | 0.55 | 0.55 |
|  | 0.75 | 0.75 | 0.75 | 0.75 | 0.75 |
|  | 0.95 | 0.95 | 0.95 | 0.95 | 0.95 |
|  | 1.15 | 1.15 | 1.15 | 1.15 | 1.15 |

from the shait centerline through the midchord at the blade tip.

## Choice of Porameters

F Geometric properties of the propeller series were chosen to be consistent with present-day practice. All the propellers were designed to have the same chordwise load distribution as that of the NACA $a=0.8$ mean line, i.e., a constant chordwise load from the blade leading edge to 0.8 of the chord, and a constant slope to zero at the trailing edge. This mean line was chosen because viscous effects on its lift are small [26], and hence the potential solution closely approximates the true physical flow. The maximum ordinate of this mean line for a $C_{L}=1.0, C_{\text {max }}$, is 0.06790 , and the twodimensional ideal angle of attack for $C_{L}=1.0$, $\alpha_{t_{1},}$ is 1.54 deg [26]. The chordwise thickness distribution chosen was the NACA-66 section, TMB modification [27], which has a desirable pressure distribution. Table 1 indicates the


Fig. 4 Expanded blade outines, $A_{z} / A_{b}=0.75, Z=5$, with symmetrical blade outline and 14 -deg skew
half-thickness ordinates as normalized by the maximum thickness. How:ver, the lifting-surface correction for thickness would be expected to be essentially independent of the chordwise thickness distribution, assuming, of course, that the shapes are reasonable.

The spanwise distribution of thickness was assumed to be linear with respect to radius and was given by the following equation:

$$
\begin{equation*}
\frac{I_{\max }}{D}=(B T F-0.0003)(1-r)+0.0003 \tag{18}
\end{equation*}
$$

where $B T F$ is the blade thickness fraction. I nis choice was based on an examination of a number of propellers of current design.

The blade outlines chosen were slightly wider toward the tip than the Troost outline and were given by

$$
\begin{equation*}
c / D=\frac{C(r)}{Z}\left(\frac{A_{B}}{A_{0}}\right) \tag{19}
\end{equation*}
$$

where $C(r)$ is given in Table 2 and $\left(1_{8} /:_{0}\right)$ is the expanded area ratio. This is a mathematical outline given by Cox [5] with his constants $\sigma=0.732$ and $\bar{x}=0.7$, and is essentially the same outline as given by Schoenherr [2S] with his constants $c=0.4$ and $n=0.3$. An expanded blade for an $A_{g} / A_{0}$ of 0.75 and five blades is shown in Fig. 4.

The radial distribution of ske:y was chosen so that the blade-section mid-chord line followed a circular arc in the expanded plane, as shown in Figure 4 , and was calculated from the following equation:


$$
\begin{gather*}
\frac{\text { skew }}{R}=R_{s}-\sqrt{R_{s}^{2}-(r-0.2)^{2}} \\
R_{t}=\frac{0.32}{\text { skew }_{t_{i},}}+\frac{\text { skew }_{\text {up }}}{2} \tag{20}
\end{gather*}
$$

where

$$
\text { skew }_{\text {tip }}=\frac{\theta_{1}}{\cos \beta_{i}}
$$

and $\theta_{3}$ is the skew angle. As an example, skew in the expanded plane as calculated from this equation is given in Table 3 for a skew angle of 14 deg and a $\pi \lambda_{i}=1.2$.

## Series of Propellers With Symmetrical Blade Oulline

The propellers which were calculated for this systematic series are listed in Table 4. Correction factors for camber, ideal angle due to loading,

and ideal angle due to thickness as determined by equations (10), (13), and (17), respectively, are listed in Table 5 for these propellers. These data, in general, are direct computer outputs and are not cross-faired. The data should not be regarded as more accurate than to the second decimal place even though three places are given. Plots of the correction factors versus the propeller
radius are shown in Figs. 5 through 13 for a representative part of the data. Figs. 5, 6, and 7 show the correction factors for a five-bladed propeller with a $\pi \lambda_{i}$ of 1.2 for the range of expanded area ratios investigated. Figs. S, 9, and 10 show the correction factors for a five-bladed propeller with $A_{E} / I_{0}=0.75$ for the range of induced advance coeflicients investigated. Fi-

nally, Figs. 11, 12, and 13 show the correction factors for $\pi \lambda_{1}=1.2$ and $A_{k} / A_{0}$ for four, five, and six blades.

In addition to these results, the chordwise distributions of camber, as normalized by the maximum camber, equation (12), are listed in Tabie 6. Only the normalized coordinates for part of the series are presented since the variation
with the various parameters is small. The data given in this table are direct computer output and, obviously, there are discrepancies at some radii. These radii are marked in the table. The questionable data have been carefully checked and no apparent error could be found. The discrepancies are probably linked to the difficulty encountered in the calculations for low pitch
Table 6 Normalized Chordwise Ordinates, $f\left(r, x_{c}\right)$, for Symmetrical Outline, $Z=4$




Propeller Lifting-Surface Corrections


Fig. 5 Camber correction factor for five blades and $\pi \lambda_{i}=1.2$


Fig. 6 Correction factor for ideal angle due to loading for five blades and $\pi \lambda_{i}=1.2$


Fig. 7 Correction factor for ideal angle due to thickness for five blades, $\pi \lambda_{1}=1.2$ and $B T F=0.4$


Fig. 8 Camber correction factor, $h_{c}(r)$, for five blades and $A_{E} / A_{\rho}=0.75$


Fig. 9 Correction factor for ideal angle due to loading for five blades and $A_{2} / \Lambda_{0}=0.75$
ratios and/or small blade areas. Also, all chordwise data should not be regarded as more accurate than to the second decimal place.

A comparison of the chordwise distribution of camber with the two-dimensional value from reference [26] is shown in Table 7 for a typical propeller. The three-dimensional values are flatter near the leading edge and fuller toward the trailing edge of the blade, and the differences deczease slightly toward the blade tip for this propeller. As can be seen from Table 6 , there is a tendency for the coordinates to become flatter toward the leading edge with increasing induced advance coefficient and with increasing expanded area ratio. Also, the ordinates become slightly fuller toward the trailing edge for increasing induced advance coefficient, but show little change with area ratio. These data show only slight variation in the spanwise direction, however, and these differences are essentially within the accuracy of the calculations. This leads to the interesting supposition that the normalized camber ordinates are almost independent of their spanwise position.


Fig. 10 Correction factor for ideal angle due to thickness for five blades and $A_{E} / A_{t}=0.75$


Fig. 11 Camber correction factor for $\pi \lambda_{1}=1.2$ and $A_{k} / A_{0}=0.75$

It is quite apparent from Table 5 and Figs. 5 , 6 , and 7 that the effect of expanded area ratio dominates the correction factors. The correction factors all increase with expanded area ratio for a given number of blades and pitch ratio. In general, the correction factors for camber and


Fig. 12 Correction factor for ideal angle due to loading for $\pi \lambda_{i}=1.2$ and $\lambda_{E} / A_{3}=0.75$


Fig. 13 Correction factor for ideal angle due to thickness for $\pi \lambda_{:}=1.2$ and $A_{E} / A_{0}=0.75$
ideal angle due to loading increase with increasing induced advance coefficient, except possibly near the hub and tip of the blades, but decrease for increasing number of blades, i.e., for the range investigated. The factors for camber and ideal

Table 7 Comparison of Chordwise Distributions of Camber for Five Blades, $\pi \lambda_{i}=1.2$, and $A_{E} / A_{0}=0.75$

| Chord Position, $x_{e}$ | 2-Dimensional Distribution | $\begin{gathered} \text { Distribution } \\ \quad \text { for } \\ r=0.3 \end{gathered}$ | $\begin{gathered} \text { Distribution } \\ \quad \text { for } \\ r=0.6 \end{gathered}$ | $\begin{gathered} \text { Distribution } \\ \quad \text { for } \\ r=0.9 \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: |
| 0 | 0 | 0 | 0 | 0 |
| 0.025 | 0159 | (0) 121 | 0.132 | 0.133 |
| 0.05 | 0271 | 0230 | 0.243 | 0243 |
| 0.1 | 0.448 | 0410 | 0.424 | 0.431 |
| 0.2 | 0699 | 0.675 | 0.687 | 0.698 |
| 03 | 0864 | () 850 | 0.860 | () 871 |
| 0.4 | 0962 | 0.961 | 0.963 | 0) 965 |
| 05 | 1000 | 1000 | 1000 | 1.000 |
| 06 | 0 979 | 0) 977 | 0) 979 | 9) 973 |
| 0.7 | 0889 | () 882 | 0.889 | 1.878 |
| 0.8 | 0) 703 | () 682 | 0.691 | 0.682 |
| () 9 | 0.359 | 0.364 | 0365 | 0.365 |
| 0.95 | 0171 | () 177 | 0177 | 0.176 |
| 0.975 | 0082 | 0085 | 0.086 | 0082 |
| 10 | 0 | 0 | 0 | 0 |

Propeller Lifting-Surface Corrections


Fig. 14 Comparison of three blade outlines for fivebladed propelier with expanded area ratio of 0.75
angle are greater than unity, except for the higher numbers of blades at low induced advance coefficients. This means that the three-dimensional values are, in general, greater than the twodimensional values. In most cases, the thickness induces an ideal angle which is largest near the blade root and decreases to negligible values toward the blade tip. In general, this angle decreases with an increasing induced advance coefficient but increases with increasing number of blades.

To show the effect of blade outline, spanwise loading, and spanwise thickness distribution on the correction factors, propellers were calculated in addition to those shown in Table 4. All calculations were for a five-bladed propeller with an induced advance coefficient $\pi \lambda_{1}$, of 1.0 and a blade area ratio of (0.75. Two blade outlines, in addition to the one of the series, were investigated: one was the Cox Type 1 and the other was a Troost outline. Fig. It compares these outlines, and Figs. 15-17 show the correction factors for the two additional outlines and the series outline. The correction factors generally follow the blade chord distribution. i.e., the blade outline which is narrowest near the root and widest toward the tip (Cox Type 1) has the smallest correction factors near the root and the largest near the tip. Correction factors for the Troost outline, which is the widest near the root and the narrowest toward the tip, show opposite trends. This is true for all three correction factors, except that the ideal angles due to thickness are too small near the


Fig. 15 Comparison of camber correction factors for three blade outlines of Fig. 14


Fig. 16 Comparison of correction factors for ideal angle due to loading for three blade outlines of Fig. 14


Fig. 17 Comparison of correction factors for ideal angle due to thickness for three blade outlines of Fig. 14
tip to show any such effect. The trends of the correction factors with the blade outlines suggest that correction factors could be approximated from the series for an arbitrary outline by using an "equivalent" area ratio at each radius. The "equivalent" area ratio is defined as the expanded blade area ratio for a propeller of the series which has the same chord at the particular radius in question as the arbitrary outline. Correction factors derived from the series with "equivalent" outlines for the Cox Type i and Troost blade shapes were found to be reasonably close to the calculated values, but not in every case were


Fig. 18 Comparison of ideal angle correction factors for radial variation in thickness


Fig. 19 Comparison of three pitch distributions for five-bladed propeller with expanded area ratio of 0.75


Fig. 20 Comparison of camber correction factors for three pitch distributions of Fig. 19


Fig. 21 Comparison of correction factors for ideal angle due to loading for three pitch distributions of Fig. 19


Fig. 22 Comparison of correction factors for ideal angle due to thickness for three pitch distributions of Fig. 19

Table 8 Comparison of Radial Distributions of Thickness

| $r$ | $\begin{gathered} t_{\text {max }} I W \\ \text { series } \end{gathered}$ | (tmas ${ }_{\text {a }}$ ) |
| :---: | :---: | :---: |
| 012 | $0003 \%$ | (1) 03216 |
| 103 | () 0289 | 1) 02761 |
| 04 | 0025 | c) 022393 |
| 05 | 0015 | (1)01939 |
| 118 | (0)0178 | () 0157\% |
| 07 | 0 0141 | () 01229 |
| 118 | 0) 014 | 0108899 |
| 0) 9 | O 006\% 7 | () (x)502 |
| 10 | (3)0230 | 0 (1)3310 |

Tabie 9 Calculated Propellers with Skew

| $\begin{aligned} & \pi \lambda_{1} \\ & \frac{\lambda_{\varepsilon}}{\lambda_{0}} \end{aligned}$ |  |  |  | -0.8- |  |  | -1.2-1.2- |  |  | -2.0- |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | 0.35 | 0.75 | 115 | 0.35 | 0.78 | 1.15 | 0.35 | 0.7 | 1.15 | 0.35 | . 75 | 1.15 |
|  | 7 | 7 |  |  | 1 |  | 7 | 7 | 7 | 7 | 7 | 7 |
| $\theta_{\text {c }}$ deg der | 14 | 14 | 14 | 14 | 14 | 14 | 14 | 14 | 14 | 14 | 14 | 14 |
|  | 21 | 91 | 21 |  | 21 |  | 21 | 21 | 21 | 21 | 21 | 21 |

at the tip. The resulting pitch distributions are shown in Fig. 19 along with the wake distribution for the wake-adapted propeller. The relation between the correction factors for these two load distributions, as compared to the series results, is shown in Figs. 20, 21, and 22. It is quite apparent that the radial load distribution has a significant effect on the camber and ideal angle due to loading correction factors. These correction factors for the propeller with the reduced loading at the tip are largest toward the blade root and smallest toward the blade tip. The correction factors for the wake-adapted propeller are close to the series results, probably because the radial load distribution is much the same. An attempt was made to estimate the correction factors from Figs. S, 9, and 10 by enteris.g the figures at the $\lambda_{i}$ of each radius for the wakeadapted propeller and that with reduced loading at the tip. Although the values of the correction factors estimated in this manner were, in general, closer to the calculated values, there was not a significant improvement over using the series correction factors directly. The result of the analysis shows clearly that the effect of the radial load distribution is important.

## Series of Propellers With Skew

The propellers which were calculated with skew are listed in Table 9 . Correction factors for camber, ideal angle due to loading, and ideal angle due to thickness, as determined by equations (10), (13), and (17), respectively, are listed in Table 10 for these propellers. As stated for the symmetrical blade outline, the data are not crossfaired and should not be regarded as more accurate than to the second decimal nlace.
Tables 5 and 10 show that the camber correction factor $k_{c}$ is almost independent of skew but that it does iend to increase slightly toward the blade tip above the vaiue for the symmetrical outline for increasing skew. This effect increases slightly with increasing blade area and with increasing induced advance coefficient. The ideal angle correction due to thickness tends to be slightly smaller than the symmetrical outline value. Otherwise, the relationship botween this correction factor and skew shows no distinct trend.


Fig. 23 Ideal angle correction factor induced by skew, $Z=5, \pi \lambda_{5}=1.2$, and $A_{E} / A_{9}=0.75$

The dominant effect of skew is on the ideal angle correction due to loading ka. Fig. 23 shows the ideal angle correction factor induced by skew for this series versus the propeller radius for $Z=5$, $\pi \lambda_{1}=1.2$, and $A_{s} / A_{0}=0.75$. To obtain this factor, the ideal angle correction factor for the symmetrical outline was subtracted from the correction factor for the skewed outlines. This figure shows that skew effect is significant and induces a positive angle toward the blade root and a negative angle toward the blade tip. This effect has been found by many others [10].

The chordiwise distributions of camber, as normalized by the maximum camber, equation (12), are listed in Table 11. Only data for part of the series are presented, as the results are very close to the results obtained for the symmetrical outline. The trends shown by the data for the skewed outline are the same as for the symmetrical outline.

## Experimenial Checks on the Theory

Using lifting-suriace theory, several propellers have been designed and tested at NSRDC. None of these has been taken specifically from the series, but the computer programs used in making the lifting-surface calculations were the same as used for the series or were older versions

Table 10 Correction "atiors for Skewed Propellers, $Z=4$

questionable oats

Table 10 (sont) Correction Factiors for Skewed Propellers, $Z=5$


Table 10 （cont）Correction Factors for Skewed Propeliers，$Z=6$

|  |  |  |  |  |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | 18 | ${ }^{\text {a }}$ |  |  |  | k | ［4］ | ${ }_{c}$ | Ctis |  |
|  |  |  |  |  |  |  |  |  |  |  |
|  |  |  |  |  |  |  |  |  |  |  |
| ＋$\quad 0.4$ | H 0 | ：125 |  |  | 6， | 10 | －223： | 16088 | Cosjat | 0186 |
| t |  | －073 |  |  | \％opit | 12109 | 152 | H032 | （ass， |  |
| 1 0．8 | 1．02\％ 1.301 | －2325 |  |  | \％，157． | $1200$ | － $0_{01}$ | Rit13． | 1．342， |  |
| $0.8$ | 1.277 ， 363 | －037 | ＋ |  | 1，256 | －635 | － | 1． 212 | －685 | ． |
| ，\％ |  |  |  |  |  |  |  |  |  |  |
|  |  |  |  |  |  |  |  |  |  |  |
| 0.4 0.5 |  | ． 773 d | 2．039 | －725 | ${ }_{i}$ | 2．lay | C37 | 1， 1.489 | 300 ${ }^{2}$ | 483， |
| 0.6 | － $2 \times 5$ 3 | －1s9 110 | ［：647 | C29 | 1.195 | 1：725． | －297． | 14323 | T：974 | $: 377$ |
| 0.7 0.8 | $\begin{array}{ll}1.252 \\ 1.237 & 1.337 \\ 1.259\end{array}$ |  | ：1092 | ．171 | （1．353 |  | 1818 | （1365 |  | －189 |
| 0.9 | 1．354 1．354 |  | 1，1 |  | ${ }_{305}$ | 10956 | C092， | 128s0． | \％ 850 | C11 |
| $1.13$ |  |  |  | ， |  |  |  |  | K | Sx |
| 0.3 | 2．331 30.309 | 1.6388 | 8 F | \％ | 2．137 | 3.557 | －890 | 20268） |  | － 327 |
| 0.4 | $\begin{array}{ll}1.447 \\ 1.311 & \text { 2，} 325 \\ \text { 2，303 }\end{array}$ | 1．280 $859 \times 10$ | ＋ | 689 | le 1.654 | ${ }_{\substack{2.732 \\ 2.565}}$ | － 5175 | 11：564 | 2．575 | ． 2123 |
| 0 | ［1．113 | －9\％ | 3 | － | 1， 6.45 | 2． 297 | －112 | 1．585 | 2．452： | ． 35 |
| 0. |  | －278） | ， | － | （1．627 | 2．985 | ， 127 | 1.755 | 2．23s： | －257 |
| 0.9 | 12，432 1．804 | －083 | 6） | 4，${ }^{\text {S }}$ | 1：482 | 1：98s |  | 2，473 | 1i＋492 | 0 |
|  |  |  |  |  |  |  |  |  |  |  |
|  |  |  |  |  |  |  |  |  |  |  |
| ；$\quad \mathrm{r}=0.2$ | 1.184 | ． 372 | 1．572 | 0.419 | 11．125 | ${ }^{16} 1095$ | \％ 335 | 1＊152 | $1.927^{*}$ 2.055 | ：235 |
| 0.4 | $1 \begin{array}{ll}1.852 \\ 0.73 & 1.740 \\ 1.5 \%\end{array}$ | － 1941 | 1．0．733 | .197 | 1.64 |  | －225 | 10094： | － 20.614 | －203 |
| 0.5 | 16\％2 liss |  |  | 123 | 1－c¢13 | －i：32 | \％ | 1，13 | fecti： | －158 |
| 0．7 |  | －tas | 1．433 | 678 | ¢ | ${ }_{1}^{1.405}$ | $\cdots$ | （ 1.1813 | 51．632 | －105 |
| 0.3 | 11,312 ， 551 | 1：322 | 35 | 22 | ． 33. | 142 | 044. | 1．0985 | － 364 | 38 |
| $A_{E} / A_{0}-0.75$ $\qquad$ <br>  <br>  |  |  |  |  |  |  |  |  |  |  |
| 0.3 | 1.474 2．645 | 1.42201043 | 2.534 | 1.0301 | 1.927 | 2，330． | ． 764 | 16050 | 2.557 | 459 |
| 0.4 | －924 |  | 3.388 2.257 | ． 114 | 1 1．233 | － 2.683 | C35\％ |  | 2.223 2.704 | － 335 |
| 0.6 |  | －212 |  | ．234 | 10.205 | 2.211 | －291． | $4 \times 319$ | 2.507. | ：212 |
| 0.7 |  | －11s $0_{0}$ | 1，731 | ${ }^{172}$ | 1.35 | 1.1 .859 | －192 | 10330． |  | ：133 |
| $\begin{aligned} & 0.6 \\ & 0.9 \end{aligned}$ | （1．247 1.50 |  | 1.854 <br> .653 | ．072 | 1.1 .577 | $\begin{array}{r}1.35 \\ \hline 245\end{array}$ |  |  | － | ${ }_{10}$ |
| ，\％ $4 / 10_{0}=1.15$ |  |  |  |  |  |  |  |  |  |  |
|  | 12．351 3.371 | 1．687 12.15 | 3.172 | 1.225 | 2，149 | 3.059 | － 8721 | 1．916： | ．2．323． | 13 |
|  | 1.447 3.470 <br> 1.211  | ［．319 0151.555 | SesEs | －959 | 1． 1.553 <br> 1.537 <br> 53 | 3.355 | －237． | 1.226 | 2.934 | 145 |
| ${ }_{0}^{0.5}$ | ［1．211 $\begin{array}{ll}1.820 \\ 1.113 & 2.332\end{array}$ |  | 2.911 2.575 | ． 712 | 1．537 | 3.912 2.741 | － 5 ¢ 42 | ${ }_{1}^{1.4750}$ | － | － 320 |
| 0.7 | 11.140 1．815 | 30951 | 2.175 | ． 315 | 1.635 | 2． 350 | －295 | 1 l 5 So | $\bigcirc$ | ．259 |
| 0.8 0.9 | $\begin{array}{ll}\text { 1．43s } \\ 13.432 & 1.397 \\ 1.260\end{array}$ | ．119， 115 |  | －1132 | 1.071 2.996 | 1.752 $: 714$ | －1188 | li．791 |  | －179 |
|  |  |  |  |  |  |  |  |  |  |  |
| $s_{5}=22^{\circ}+1 / A_{0}=0.35$ <br>  |  |  |  |  |  |  |  |  |  |  |
|  |  |  |  |  |  |  |  |  |  |  |
| 0．4 |  | ：13 | \％ | 新 | 1．033 | 2．332 2.253 | ． 324 | ［1．103． | 2．425 | ． 262 |
| 0．6 | 1．061i 1.177 | ：71 | \％$\%$ | \％ | 1．233 | 2．045 | －135 | 1．161） | 2．245 | \％152 |
| 0.7 | －1037 1．6\％ | $\bigcirc$ | ， | \％${ }^{5}$ \％ | 1．165 | 1.655 | 0.005 | 1.198 | 1：934 | ：103 |
| 0.8 | $\begin{array}{ll}1.164 & 1.351 \\ 1.35\end{array}$ | ：5in $0^{2}$ | 5 | 柂衰 | 1.348 | －． 28 El | ． 044 |  | － | ．039 |
| $\lambda_{E} / A_{0}=0.75$ |  |  |  |  |  |  |  |  |  |  |
|  | 1．47： 2.53 | 1．575 1.463 | 2.779 | 1．019 | 1.527 | 2.861 | ． 753 | 1.050 | $2.504:$ | $\bullet$ |
| 0.4 | －964 3.581 | 021 1－139 | 2.887 | ． 710 | 1．232 | 5 | －585 | （1．0354： | 3．0388 | ． 432 |
| 0.5 0.6 | $\bigcirc 313$ | －215 415 | \％ | 42 | ： 1.195 | 2．693 | － 4294 | ．1．013 | 2．969． | －271 |
| 0.7 | $\therefore 062 \quad 1041$ | 0131 ［1．331 | 2.045 | －169 | ：．328 | 2.157 | 188 | 1.438 | 2.458 | ：188 |
| 0．8．8 |  |  |  |  | 1．513 | 1．350 | ． 6119. |  | （．541 | ． 191 |
|  |  |  |  |  |  |  |  |  |  |  |
|  | 2.328 | 1.6989 | $x_{6}$ |  | $2.149$ |  | － 827 | $2: 261$ |  | 5554 |
|  | $\begin{array}{ll}1.447 \\ 1.217 & 4.017 \\ 1.218 \\ 3.004\end{array}$ | $\begin{array}{r}1.275 \\ .194 \\ \hline 1\end{array}$ | － |  | ＋ $\begin{aligned} & \text { 2．552 } \\ & 1.572\end{aligned}$ | 3.873 3.491 | －．702 | $\begin{array}{r} 1 ; 795 \\ i, 710 \end{array}$ | 36966， | ． 313 |
| ，${ }^{2}$ ，${ }_{0.5}^{0.5}$ | 1.113 | ：968 |  |  | 1．565 | 4.165 | ．414 | 1，712 | 3．535． | $4{ }^{1}$ |
| 5 0.6 <br> $\mathrm{~B}_{2}$ 0.7 <br> 0 0.8 | 1.140 | －316， | \％ | 442 | 1.653 | 2.697 | －294 | ${ }^{1} 1.803$ | 3，049：－ | －254 |
|  | $\begin{array}{ll}3.433 \\ 2.433 & 1.416 \\ .617\end{array}$ | ${ }_{.082}^{147}$ | \％ 4 |  | 1.310 2.53 | 1.814 <br> .618 | －128 | 2．010， | 169938 |  |
| ckcos | 2.433 ． 617 | ．082／ 2 | ， | \％ | 2.565 | －． 018 | ．123 | 2.399 | 4393 | ［134 |


Table 11 （cont）Normalized Chordwise Ordinates，$f\left(r, x_{c}\right)$ ，for Skewed Propellers；$Z=5, A_{k} / A_{11}=0.75$

|  |  | 87\％ <br> 8ロロ <br>  <br>  <br>  <br>  <br>  |  <br>  <br>  <br>  <br> ยทำ <br>  |
| :---: | :---: | :---: | :---: |
|  |  <br>  <br>  <br>  <br>  <br> 4 <br>  <br>  |  <br>  <br>  <br>  <br> ตดต <br>  <br> ？ |  <br>  <br>  <br>  <br>  <br>  <br>  |
|  |  <br>  <br>  <br>  <br>  <br>  <br>  |  <br>  <br>  <br>  <br>  <br>  <br> Eax |  <br>  <br> 号： <br>  <br>  <br>  <br> 气anconan |
|  |  |  |  |



Table 12 Comparion Between Theory and Experiment for a Constont Thrust Loading

| Propeller | $\%$ | $A E / A$ | $\begin{gathered} P / D \text { at } \\ 0.7 R \end{gathered}$ | $B T F$ | $C_{r}$ | $\begin{aligned} & \text { Experi- } \\ & \text { mental } \lambda_{0} \\ & \text { Design } \lambda_{1} \end{aligned}$ | Enperimental $C_{p}$ Design $C_{r}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| A | 5 | 1.318 | 1.473 | 0.054 | 0.5663 | 1.013 | 0.976 |
| B | 3 | 0.603 | 1.077 | 0.040 | ง. 563 | 0.904 | 1.025 |
| C | 3 | 0.806 | 1.081 | 0.080 | 0.503 | 0.087 | 0.089 |
| D | 3 | 0.303 | 1.086 | 0.057 | 0.570 | 0.995 | 1.016 |
| F | 3 | 1.212 | 1.073 | 0.028 | 0.549 | 1.014 | 0.965 |
| G | 5 | 1.48) | 1.503 | 0.036 | 0.361 | 1.008 | 0.971 |

of the programs. Table 12 compares theoretical and experimental values for six of these propellers, designated A, B, C, D, F, and G; all used the NACA $a=0.8$ mean line and had no skew. In general, the theoretical values are within the accuracy of the experiments and within the accuracy with which the blade-section viscous drag can be chosen.
There appears to be a trend for wide-bladed propellers A, F, and G to be slightly over-pitched. Also, the thick propeller $C$ tended to be slightly under-pitched as compared to the propeller of standard thickness, B. It is possible that both of these differences could arise from determination of the section viscous drag from airfoil data. Propellers A, F, and G have blades with thicknesschord ratios lower than those for which experimental data are available, and propeller B has blades with thickness-chord ratios which are higher than common. A further discussion of most of these propellers will be found in reference [29].

## Conclusiens

From the numerous calculations made and the series data presented, a number of conclusions can be drawn with reference to lifting-surface corrections for propellers. Many of these conclusions, of course, have been made previously by other investigators.

1 Use of a realistic chordwise load similar to that for an NACA $a=0.8$ mean line results in both an induced camber and ideal angle. The camber correction factors for this load distribution are somewhat less than for the constant-load mean line, but, in general, they are greater than unity, i.e., greater than the two-dimensional camber for the same lift. The ideal angle correction factor due to loading is, of course, zero for the constant load.

2 Correction factors can be formed for the camber and ideal angle due to loading which are independent of the magnitude of propeller loading.

3 Thickness induces an angle and a camber,
but the camber correction is negligible except for low pitches.
$\pm$ A correction factor can be formed for the ideal angle induced by thickness which is independent of the magnitude of thickness but dependent on the thickness distribution.

5 Correction factors for camber and ideal angle due to loading are largest near the blade tip and smallest at 0.4 and 0.5 radii. They increase with increasing expanded area ratio and, generally, with increasing induced advance coefficient, but they decrease with increasing number of blades.

6 The ideal angle due to thickness is largest near the blade root and decreases to negligible values toward the blade tip. In general, this angle decreases with increasing induced advance coefficient but increases with increasing blade number.
7 The chordwise distribution of camber is somewhat flatter toward the blade leading edge and, in general, fuller toward the trailing edge as compared with the two-dimensional values. The spanwise change in the camber chordwise distribution is small. In practice, the use of the two-dimensional distribution is probably reasonable.

8 The shape of the blade outline has a significant effect on the correction factors for camber and ideal angle due to loading.

9 The spanwise load distribution has a significant effect on the correction factors for camber and ideal angle due to loading.

10 Skew has little effect on the camber and ideal angle due to thickness but has a large effect on the ideal angle due to loading. Skew induces a positive angle near the blade root and a negative angle toward the blade tip.

11 Experimental results indicate that, for a chordwise load distribution corresponding to that of the NACA $a=0.8$ mean line, the use of liftingsurface corrections gives propellers which, in general, meet their predicted performance within the accuracy of the experiments and the accuracy
with which the blade-section viscous cirag can be chosen.

12 Correction factors derived herein should replace those calculated by less sophisticated methods; for example, those in reference [22].

13 The numerical evaluation of the complicated theoretical equations may lead to solutions which are erroneous, and calculations must be carefully checked.

## Acknowledgments

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    - Numbers in brackets designate References at end of paper.
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[^1]:    ${ }^{5}$ Reference [2:] is an example of an outmoded method. In this particular reference the camber correction factors $k_{1}$ and $k_{2}$ are replaced by the camber correction factor $k_{c}$. and the pitch correction procedure is replaced by the ideal angle due to loading and thickness.
    -The word optimum used here means the propeller las a constant hydrodynamic pitch. Strictly speaking, such a propeller would be optimum only when lightly loaded, when operating in an inviscid fluid, and if it had an infinite number of blades.

[^2]:    ${ }^{7}$ For this clange in coordinate system, the lifting line may deviate from the blade surface at radial points other than $P$ when the propeller is not of constant pitch. The effect on the results is probably small if the pitch does not deviate too much from a constant value.

[^3]:    The induced advance coeflicient $\lambda_{1}$ is related to the t:ydrodynamic pitch ratio by $\pi \lambda_{1}=r \tan \beta_{1}=(P / D)_{1}$.

