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Protein-ligand Interaction Energy-based Entropy

Calculations: Fundamental Challenges For Flexible

**Systems** 

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ABSTRACT: Entropy calculations represent one of the most challenging steps in obtaining the binding free energy in biomolecular systems. A novel computationally effective approach (IE) was recently proposed to calculate the entropy based on the computation of protein-ligand interaction energy directly from molecular dynamics (MD) simulations. We present a study focused on the application of this method to flexible molecular systems and compare its performance with well-established normal mode (NM) and quasiharmonic (QH) entropy calculation approaches. Our results raise substantial concerns on the general applicability of IE in terms of reproducibility, reasonable absolute values of the entropy and agreement with NM and QM approaches. IE shows significant variation in the computed entropy values depending on the MD frames chosen for calculations. These deviations render reproducibility of IE calculations to be far from sufficient. We conclude that IE is recommended to be used after substantial modifications with respect to its sampling methodology.

#### Introduction

Understanding the principles of intermolecular recognition between diverse biomacromolecules is a key in establishing molecular basis of all biochemically relevant processes in the cell. Both experimental and theoretical approaches are being constantly developed to 40be able to answer more effectively the fundamental question of how several biomolecules form a complex and so carry out their biological function<sup>1</sup>. Proper characterization of the structure-function relationship in biomacromolecular systems is a critical step for novel drug discovery and therefore represents one of the central topics in the field of biomedical research<sup>2,3</sup>. Although experimental data on structures of molecular complexes can provide an essential insight into the function of

corresponding systems, computational approaches are of high value. They not only support, complement and explain experimental findings but could also provide the detailed analysis of the data, understanding the contribution of separate energy terms and their connection to the participating molecules, which is mostly unapproachable from experiments. Moreover, they assist the rational guiding of further experiments leading to the minimization of the time and resources invested<sup>4</sup>. However, a major sensitive feature of computational methods that always needs careful evaluation is the reliability of the produced results, especially in terms of their comparison with the corresponding experimental data. The latter is key for increasing trust and widespread use of the available methods, highly aiding understanding of a plethora of molecular processes.

Among others, calculating binding affinity, a crucial property of complex systems, represents a subject of many diverse computational approaches<sup>5</sup>. However, despite successes in free energy calculations<sup>6</sup> by using molecular dynamics (MD)-based methods as rigorous free energy perturbation and thermodynamic integration methods<sup>7</sup>, Molecular Mechanics Poisson Boltzmann Surface Area (MM-PBSA)<sup>8,9</sup> or linear interaction energy<sup>10</sup> and Monte Carlo simulation-based methodologies<sup>11</sup>, accurate estimation of the binding affinity of a molecular system still represents a serious challenge<sup>12</sup>. One of the most commonly applied MD-based free energy calculation method, MM-PBSA and its modified version MM-GBSA (Generalized Born Approximation for MM-PBSA)<sup>13</sup>, which implements an implicit solvent model, yields reliable enthalpy values<sup>14,15,16</sup>. Nevertheless, its bottleneck is the calculation of the entropic component, which can potentially increase the overall uncertainty in the calculated binding energies dramatically<sup>17,18</sup>. A popular

method to calculate entropy, which is widely used with the MM-PBSA approach, is the normalmode (NM) analysis of harmonic frequencies obtained from minimized snapshots of MD simulations<sup>9</sup> and estimates the conformational entropy of a molecular system in vacuo. This method, being computationally expensive, has a solid capability to yield potentially useful results in reproducing experimental binding affinities<sup>19</sup>. An alternative method is quasiharmonic approximation (QH), which is essentially faster but in general less accurate and has a number of limitations<sup>20</sup>. QH involves calculations of the covariance matrix of atomic coordinates from MD simulations, interprets the covariances resulting from a harmonic energy surface, while neglecting anharmonic energy surface, and computes QH force constants, which are then used to compute the entropy<sup>21</sup>. A QH approach yields the conformational entropy of a molecule in an "effective" quadratic potential, implicitly including effects of solvent, which renders its results to be different from results from the NM approach. Moreover, the conformational entropy calculated by both above-mentioned methods represent only a part of the total entropy change in a binding process. In comparison to the above-mentioned methods, a relatively novel approach to calculate the entropy based on the computation of protein-ligand interaction energy reported by Duan et al. can use the already generated data from MD simulations without any additional computational cost<sup>22</sup>. The method, which will be related to as IE in this paper, calculates the entropy as following. The gas-phase component of the binding free energy, which is obtained directly from the MD simulation is derived:

$$\begin{split} &\Delta G_{gas} \!=\! -KT \ln \frac{\int dq_{w} dq_{p} dq_{l} e^{-\beta (E_{p} + E_{l} + E_{p}^{\text{int}} + E_{w}^{\text{int}} + E_{hw}^{\text{int}})}}{\int dq_{w} dq_{p} dq_{l} e^{-\beta (E_{p} + E_{l} + E_{w} + E_{pw}^{\text{int}} + E_{hw}^{\text{int}})}} \\ &= -KT \ln \left[ \frac{1}{\langle e^{\beta E_{pl}^{\text{int}}} \rangle} \right] = KT \ln \left[ \langle e^{\beta E_{pl}^{\text{int}}} \rangle \right] \\ &= KT \ln \left[ e^{\beta \langle E_{pl}^{\text{int}} \rangle} \langle e^{\beta \langle E_{pl}^{\text{int}} - \langle E_{pl}^{\text{int}} \rangle} \rangle \right] \\ &= \langle E_{pl}^{\text{int}} \rangle + KT \ln \langle e^{\beta \Delta E_{pl}^{\text{int}}} \rangle \\ &= \langle E_{pl}^{\text{int}} \rangle - T \Delta S \end{split} \tag{1}$$

 $E_p, E_l, E_w$  are protein, ligand and solvent internal energies, respectively;  $E_{pl}^{\text{int}}$ ,  $E_{pl}^{\text{int}}$ ,  $W_{lw}^{\text{int}}$  are protein-ligand, protein-solvent and ligand-solvent interaction energies, respectively;  $\langle E_{pl}^{\text{int}} \rangle$  is the protein-ligand interaction energy averaged in the ensemble, and

 $\Delta E_{pl}^{int} = E_{pl}^{int} - \langle E_{pl}^{int} \rangle$  (2) is the difference between protein-ligand interaction energy and its average value. Hence, final equation of interaction entropy is defined by the protein-ligand interaction energy and is expressed by the following equation:

$$-T \Delta S = KT \ln \langle e^{\beta E_{pl}^{int}} \rangle$$
 (3)

The relevant ensemble averages can be calculated by averaging over MD trajectory:

$$\langle E_{pl}^{\text{int}} \rangle = \frac{1}{T} \int_{0}^{T} E_{pl}^{\text{int}}(t) dt = \frac{1}{N} \sum_{i=1}^{N} E_{pl}^{\text{int}}(t_{i})$$
 (4)

The averaging procedure can be further formulated as:  $\left\langle e^{\beta \Delta E_{pl}^{int}} \right\rangle_{arithmetic} = \frac{1}{N} \sum_{i=1}^{N} e^{\beta E_{pl}^{int}(t_i)}$  (5)

Authors claim that the method is essentially superior to the NM approach demonstrating it to be computationally efficient and numerically reliable<sup>22</sup>. Since the method was described in 2016, it has been widely applied and cited in the literature as original papers<sup>23-31</sup> and reviews<sup>32, 33</sup>.

The conceptual break-through arising from the theory behind IE method is that the entropy could be defined in a highly efficient and reliable way directly from the analysis of MD simulations without any entropy-specific post-processing calculations. The motivation to carry out an evaluative case study focusing on IE applicability<sup>22</sup> was grounded on the following observation that has drawn our attention: in practice, according to the described original protocols, the calculated entropic component of the free energy of binding is determined by a single spike (the highest positive value) of the protein-ligand interaction energies obtained in a course of the MD simulation in comparison to its average value. This could potentially suggest that if such a spike does not occur during the analyzed MD simulation or, on contrary, if a bigger spike occurs, the calculated entropy value obtained should be essentially different. Taking into account the length of the analyzed MD simulation from the original procedure, such calculations could be insufficient to compensate the effect of all the possible spikes in values of the interaction energy between the receptor and the ligand and so lead to low reproducibility of the calculated entropy values. In order to deal with the mentioned challenge in the absence of sufficient thermal equilibration of the system, it is possible to impose an energy cut-off for such an energy spike in order to smooth its impact on the calculated entropy values<sup>34</sup>. In our work, we rigorously analyzed and pointed out the limitations of IE in terms of its reliability and reproducibility for several case study systems with the focus on flexible molecular systems.

## Methods

Structures. We used the experimental structure of urokinase-type plasminogen activator in complex with 4-(aminomethyl) benzoic acid (PDB ID: 3KGP, 2.35 Å) that was previously used in the dataset in the original study on IE method of Duan et al.<sup>22</sup>. The structures for the complex between CM15 peptide and suramin were obtained by the molecular docking procedure as described further.

Ligands parameterization for docking and MD simulations. Both 4-(aminomethyl) benzoic acid and suramin (Scheme 1) were parameterized in antechamber package <sup>35</sup> of the AMBER 16 MD programs suite<sup>36</sup> using GAFF force field parameters<sup>37</sup> with AM1-BCC charge module<sup>38</sup>.

Docking with Autodock 3. In the docking calculations, CM15 was used as a receptor and suramin as a ligand. Since CM15 is a disordered peptide, its 21 conformations were obtained by several diverse MD approaches carried out in AMBER 16<sup>36</sup>, GROMACS<sup>39</sup>, ECEPP<sup>40</sup> and coarse-grained package UNRES<sup>41</sup>. The RMSD of the backbone atoms for obtained CM15 structures was 5.6±1.4 Å, which provided a high structural variety of the peptide conformations. For blind docking simulations for each of the peptide conformations, Autodock 3 (AD3)<sup>42</sup> was used with the protocol optimized for the glycosaminoglycan ligands<sup>43</sup>, which similarly to suramin are linear and contain numerous sulfate and sulfonate groups. Flexible ligand docking was carried out within a box with a grid step of 0.375 Å. The Lamarckian genetic algorithm was applied with the following parameters: initial population size of 300, 10<sup>5</sup> generations, 9995·10<sup>5</sup> energy evaluations, 10<sup>3</sup> independent runs. All torsional angles besides C-S for sulfonates were flexible during docking. In overall, this corresponded to 10 torsional angles degrees of freedom for the

suramin molecule. Fifty top docking solutions were clustered by the DBSCAN algorithm<sup>44</sup> with a neighbourhood search radius of 2 Å and minimal number of points of 2. One representative from each cluster corresponding to the highest AD3 score was further chosen for MD-based analysis.

Molecular dynamics and MM-GBSA calculations. MD simulations were run for protein-ligand complexes using the following protocol. The complexes were solvated in a periodic TIP3P truncated octahedron water box with a minimal distance of 15 Å between complex atoms and the box, and charge neutralizing counter-ions were added. The ff14SB and GAFF force field parameters were used for protein and a small molecule, respectively. Prior to MD simulations, two energy-minimization steps were carried out: first, steepest descent (1500 cycles) and conjugate gradient (100 cycles) with harmonic force restraints of 10 kcal/(mol·Å) applied on protein and small molecule atoms; then, steepest descent (3000 cycles) and conjugate gradient (3000 cycles) without any restraints. This was followed by the process of system heating from 0 to 300 K for 10 ps, and a MD equilibration at 300 K and 10<sup>6</sup> Pa in NTP for 100 ps. Following the equilibration procedure, productive MD runs of 5 ns + 1 ns (see Entropy calculations section) and 100 ns were carried out for 3KGP and CM15-suramin complexes, respectively. An 8 Å cutoff for non-bonded interactions and Particle Mesh Ewald method for long-range electrostatic interactions were applied. The simulations were carried out in NTP with Langevin temperature coupling with collision frequency parameter  $y = 1 \text{ ps}^{-1}$  and integration step of 2 fs with the SHAKE algorithm applied to the bonds containing hydrogen atoms. MD trajectories were recorded every 10 ps. Energetic post-processing of the trajectories was done using MM-GBSA with igb = 2. MD simulations for each complex structure was repeated three times.

Entropy calculations. Entropy calculations were carried out using three different methods: interaction entropy (IE), normal mode (NM) and quasiharmonic approximation (QH). In case of the 3KGP complex, it was initially simulated for 5 ns and then the last 1 ns was repeated three times with different random seeds to be consistent with Duan's protocol<sup>22</sup>. During the last 1 ns the coordinates were saved every 0.1 ps resulting in 10000 frames. IE calculation and NM analysis were performed on each 10th frame covering altogether 1000 frames, while QH was carried out for 10000 frames due to its higher pretension of more frequent sampling<sup>45</sup>. For the CM15-suramin complex the entropic contribution was determined from the last 10 ns of the corresponding trajectory using 1000 frames for each method.

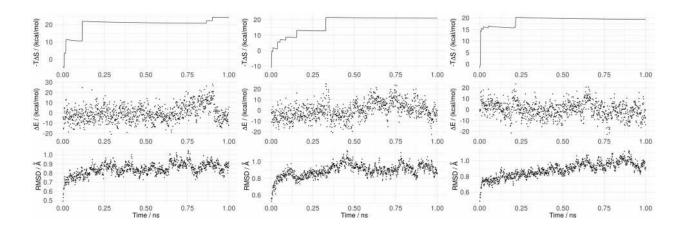
We compared the data obtained by IE to the results obtained by NM and QH for one of the systems reported in the study of Duan et al.<sup>23</sup>, containing a protein and a ligand 4-(aminomethyl) benzoic acid (Scheme 1). The second system we used contained a potential drug-carrier CM15, which is a disordered 15 aa peptide (KWKLFKKIGAVLKVL-amide), and a ligand suramin, which can induce the folding of CM15<sup>46</sup>.

**Scheme 1**. Chemical structures of 4-(aminomethyl) benzoic acid (left) and suramin (right).

Taking into account the disordered structure of CM15, entropic contribution to its binding affinity could be particularly decisive, and the system is very suitable to address potential discrepancies in entropy calculations arising from its flexible nature.

# **Results and Discussion**

First, we carried out three repetitive MD simulations of the urokinase-type plasminogen activator in complex with 4-(aminomethyl) benzoic acid taken from the dataset of Duan. The obtained results clearly show that the values of the entropy are strongly dependent on the frame corresponding to the highest spike of the interaction energy and are not significantly changed after this event due to the exponential function in the equation 3 (Figure 1).



**Figure 1**. Performance of IE method in three independent MD simulations starting from the same initial configuration. Top panel: calculated entropy. Middle panel: protein-ligand interaction energy. Bottom panel: RMSD (Å) of the ligand with the reference to itself at the start of the MD simulation.

This means that in case the simulation is prolonged, and there is a potentially higher change of the interaction energy, the value of the calculated entropy would be essentially increased. Therefore, such calculations could yield consistent results only in case a very long and probably practically unfeasible MD simulation is carried out so that the system is well converged in terms of the changes of interaction energy. Taking into account that the analyzed 1 ns is a rather short MD simulation, such expectations are not likely to be realized. The averaging procedure (5) was used to calculate the entropy in the system yielding 24.2, 20.9, 19.4 kcal/mol, respectively. These values are very close to the ones values obtained in the study of Duan et al.<sup>22</sup>, while the observed differences are expected due be to the use of a different parameterization procedure of the ligand in our study (see Supporting Materials: Methods section). Comparison of these results with the values of the entropy obtained by NM and QH reveal essential differences in the performance of the methods in terms of the absolute values obtained. Apart from this, the values obtained in three different simulations are essentially less deviating from each other in case of NM and QH than of IE (Table 1). Based on these findings, we conclude the reproducibility of IE is less robust.

When considering the more flexible CM15-suramin complex, the entropic contribution to the binding could be more representative and have a stronger impact on the full free energy of the system due to the disordered nature of the peptide. Namely, for the estimation of the binding affinities in similar complexes, entropy calculations are of special significance because MD simulations are often used as a post-processing step after docking to score the obtained docking solutions and to choose a more plausible prediction<sup>47</sup>. Comparison of the absolute values

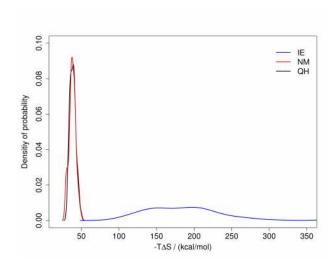
obtained from MD simulations shows significant differences in the values calculated by the three methods for this system (Table 1). The discrepancies between NM, QH and IE are essentially higher than in the urokinase system described above. This is reflected by the higher spikes of the interaction energies obtained here (Supporting Information: Figures S1, S2, S3), where the contribution of the electrostatic interactions is decisive. However, the distributions of the entropy values arising from altogether 69 MD simulations are much narrower for NM and QH, than for the IE method (Figure 2). In addition, the correlation between the values obtained by NM and QH is higher than their correlations with IE (Table 2). These observations underline the fact that according to the principle of the IE approach, the entropic term is purely determined by the interactions between the complex counterparts averaged over the equilibrated MD trajectory of their complex and neglects potential changes of their configurations occurring upon complex formation. Therefore, the application of IE for a case system from the dataset of Duan et al.<sup>22</sup> yielded far more meaningful results than in case of CM15-suramin complexes, where significant conformational changes in both complex counterparts are observed upon binding. In addition, a very high interaction energy variances observed for different CM15-suramin docked structures in the MD analysis lead to the high differences in the IE-derived entropy values for this system. This suggests that IE methodology is not intended to yield meaningful data when applied for ranking of the diverse and *a priori* not sufficiently equilibrated structures obtained by a molecular docking approach for highly flexible systems.

**Table 1**. The entropy calculations obtained by three methods: IE, NM and QH. The data for urokinase – benzoic acid complex (PDB ID: 3KGP) and several representatives of CM15-suramin complexes are shown, which were obtained from three repeated simulations<sup>a</sup>

					-ΤΔS			$\Delta G_{\rm bind}$	
name	$< E^{int} >$	$\Delta G_{\text{sol}}$	ΔΗ	IE	NM	QH	IE	NM	QH
3kgp_1	-106.1	79.1	-27.0	24.2	18.2	17.9	-2.8	-8.8	-9.1
3kgp_2	-105.0	79.0	-26.1	20.9	17.7	17.1	-5.2	-8.3	-9.0
3kgp_3	-112.9	84.8	-28.1	19.4	17.3	18.7	-8.7	-10.8	-9.4
CM-S_22_1	-971.3	936.9	-34.4	241.4	34.9	33.3	207.0	0.5	-1.1
CM-S_22_2	<b>-</b> 1191.0	1131.5	<del>-</del> 59.5	236.6	43.3	38.5	177.1	-16.2	-21.1
CM-S_22_3	-1237.8	1174.1	-63.7	156.8	40.2	39.8	93.1	-23.5	-23.9
CM-S_29_1	<b>-</b> 1337.7	1283.3	-54.4	150.2	40.9	37.9	95.8	-13.5	-16.5
CM-S_29_2	-1132.5	1085.7	-46.8	134.9	39.8	37.1	88.1	-7.0	-9.6
CM-S_29_3	-1116.7	1055.6	-61.1	104.5	38.4	40.4	43.4	-22.7	-20.7
CM-S_41_1	<b>-</b> 1275.3	1219.7	-55.6	143.4	45.3	38.1	87.8	-10.3	-17.5
CM-S_41_2	-1387.7	1326.0	-61.6	186.0	43.7	40.3	124.4	-18.0	-21.3
CM-S_41_3	-1219.6	1164.6	-55.0	107.9	42.5	39.1	52.9	-12.4	-15.9
CM-S_22_RES1	-1083.6	1039.2	-44.4	62.9	42.9	23.1	18.5	-1.5	-21.3
CM-S_22_RES2	-1049.7	1004.8	-44.9	76.1	46.2	22.6	31.2	1.3	-22.3
CM-S_22_RES3	-1084.6	1038.5	-46.1	90.5	40.2	23.1	44.4	-5.9	-23.0
CM-S_29_ES1	-1044.9	998.0	-46.9	185.3	42.0	23.2	138.4	-4.9	-23.7
CM-S_29_RES2	-1051.4	1004.6	-46.8	165.7	41.8	23.4	118.9	-5.0	-23.4
CM-S_29_RES3	-1002.0	956.0	-46.0	130.2	40.1	25.8	84.2	-5.9	-20.2
CM-S_41_RES1	-1232.1	1193.2	-38.9	70.6	42.2	22.8	31.7	3.3	-16.1
CM-S_41_RES2	-1226.1	1187.5	-38.6	54.2	42.3	22.9	15.6	3.7	-15.7
CM-S_41_RES3	-1228.6	1190.0	-38.6	56.8	42.2	23.0	18.2	3.6	-15.6
ASP_ARG_1	-22.3	21.8	-0.5	11.1	0.2	-28.6	10.6	-0.3	-29.1
ASP_ARG_2	-18.3	17.9	-0.4	7.3	0.01	-28.6	6.9	-0.4	-29.0
ASP_ARG_3	-15.2	15.0	-0.2	4.4	-0.7	-28.9	4.2	-0.9	-29.1
ASP_ARG_RES1	-32.0	31.3	-0.7	14.4	2.2	-24.6	13.7	1.5	-25.3
ASP_ARG_RES2	-45.8	44.0	-1.8	27.3	2.9	-25.1	25.5	1.1	-26.9
ASP_ARG_RES3	-32.5	32.1	-0.4	13.2	3.9	-25.0	12.8	3.5	-25.4

<sup>a</sup>The subscript after CM-S (staying for the complex CM15-suramin) corresponds to the number of the conformation of CM15 and the second subscript corresponds to the number of the simulation in three repeats. All values are in kcal/mol.

Figure 2. Density of probability for the entropy calculated by three methods: IE, NM and QH



**Table 2**. Correlations between the entropy values obtained by three methods: IE, NM and QH. The data from altogether 69 MD simulations are shown.

	IE	NM	QH
IE	1	0.16	0.11
NM	0.16	1	0.34
QH	0.11	0.34	1

In case of the flexible system IE has always at least 60 kcal/mol deviation for all obtained entropy terms when compared to NM and QH (Table 1). This consequently leads to major deviations in the obtained free energy values. Unfortunately, the obtained results by IE suggest that complex formation is highly unfavorable. Therefore the simulation results provide an unfortunate misinterpretation of the binding preference even at a qualitative level.

# **Conclusions**

To summarize our findings, we observe that: i) the occurrence of a frame with the highest spike of the interaction energy is indeed crucial and determining for the final value of the entropy in IE calculations; ii) therefore, it is dependent on the length of the MD simulation affecting the reproducibility of the method; iii) the events in the simulation occurring after this spike do not significantly affect the obtained entropy value; iv) IE method yields both absolute values as well as the distribution of the entropy essentially different from the ones in NM and QH. The latter two approaches provide rather similar entropy values even for a very flexible complex system. Despite the high potential of IE in obtaining entropy values computationally very effectively, the presented data suggest that there are strong limitations for its general applicability. It could be potentially reliable only for very well equilibrated complexes, which are fundamentally not accessible within the feasible range of MD simulations, or for those systems which according to the experiment have substantially low structural and thus intermolecular interaction energy changes. We suggest that IE should be applied after the rigorous calibration for a particular molecular system and used together with the NM approach as a general reference.

#### ASSOCIATED CONTENT

# **Supporting Information.**

**Figure S1, S2, S3.** Performance of IE method in three independent MD simulations starting from the same initial configurations for CM15-suramin system: CM\_S\_22, CM\_S\_29, CM\_S\_41.

## **AUTHOR INFORMATION**

## **Notes**

The authors declare no competing financial interests.

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