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Pulsed Magnetron Sputtering of Strongly Thermochromic VO₂-Based Coatings with a Transition Temperature of 22 °C onto Ultrathin Flexible Glass



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Abstract: The reversible semiconductor-to-metal transition of vanadium dioxide (VO₂) makes VO_2 -based coatings a promising candidate for thermochromic smart windows, reducing the energy consumption of buildings. This paper deals with maximizing the application potential of these coatings in terms of their performance, an industry-friendly preparation technique, and an industrially relevant substrate. We present a scalable sputter deposition technique for the preparation of strongly thermochromic ZrO₂/V_{0.984}W_{0.016}O₂/ZrO₂ coatings on ultrathin flexible glass and standard glass at a relatively low substrate surface temperature (330 °C) and without any substrate bias voltage. The $V_{0.984}W_{0.016}O_2$ layers were deposited by a controlled high-power impulse magnetron sputtering of a V target, combined with a simultaneous pulsed dc magnetron sputtering of a W target. We explain the fundamental principles of this technique using the discharge characteristics measured for both discharges. We characterize the coating structure (X-ray diffraction) and a wide range of optical properties (spectrophotometry and spectroscopic ellipsometry). We find that the coatings combine a transition temperature of 22 °C, a luminous transmittance approaching 50%, a modulation of the solar energy transmittance over 10% and a temperature-independent color. The results in general, and the successful transfer from a standard glass to the ultrathin flexible glass in particular, are crucial for future applications of the coatings on smart windows.

Keywords: vanadium dioxide; strongly thermochromic coatings; low transition temperature; pulsed magnetron sputtering; low deposition temperature; ultrathin flexible glass; smart windows

1. Introduction

Global warming causes numerous problems in human society and drives a focus on energy-saving materials. Owing to the excessive use of heating, cooling, lighting and ventilation, buildings have been estimated to produce about 30% of all anthropogenic greenhouse gas emissions [1] and are responsible for more than 30% of the primary energy consumption in the world [2]. It is evident that energy-saving smart windows with adjustable throughput of solar energy and visible light can lower the energy expenditure. Such windows could incorporate chromogenic materials, in particular those with thermochromic and electrochromic properties [3].

Vanadium dioxide (VO₂) exhibits a reversible phase transition from a low-temperature monoclinic VO₂(M1) semiconductive phase to a high-temperature tetragonal VO₂(R) metallic phase at a transition temperature (T_{tr}) of approximately 68 °C for the bulk material [4]. The abrupt decrease of infrared transmittance without attenuation of luminous transmittance in the metallic state makes VO₂-based coatings a promising candidate for thermochromic smart windows reducing the energy consumption of buildings.

There have been many investigations and much progress in thermochromic VO₂-based materials in recent years (see, for example, reviews [3,5–9] and the works cited therein). However, there are still obstacles impeding their application in smart windows. These are: (1) a high temperature needed for fabrication, (2) a high transition temperature, (3) a low luminous transmittance (T_{lum}), (4) a low modulation of the solar energy transmittance (ΔT_{sol}), (5) low environmental stability and (6) an unfavorable yellowish color. To meet the requirement for large-scale implementation on building glass, VO₂-based coatings should satisfy the following criteria simultaneously: a deposition temperature close to 300 °C or lower [10–13], T_{tr} close to 20 °C [14], $T_{lum} > 60\%$ [5,15,16], $\Delta T_{sol} > 10\%$ [17–19], long-time environmental stability [9,20–22], and a more appealing color (as investigated in [23,24]).

Decrease of the deposition temperature of thermochromic VO₂-based coatings to 300 °C is of key importance: (1) to facilitate their large-scale production by reducing the energy consumption, simplifying substrate heating and cooling procedures and minimizing problems with a temperature non-uniformity over large substrate surfaces, and (2) to allow deposition of these coatings onto temperature-sensitive flexible substrates.

Usual approaches in the literature [3,5–9,25] are: (1) to optimize only T_{lum} and ΔT_{sol} regardless of T_{tr} and the substrate surface temperature (T_{s}), (2) to decrease T_{s} at the cost of making the process significantly less industry-friendly, e.g., by introducing substrate bias voltage or post-deposition annealing, (3) to decrease T_{tr} using doping of VO₂ by other elements (such as W) at the cost of much lower ΔT_{sol} , or (4) to choose antireflection (AR) layers (both below and above the thermochromic VO₂-based layer) thickness in order to optimize only T_{lum} (typically utilizing the first-order interference maximum) regardless of ΔT_{sol} .

Magnetron sputter deposition, with its versatility and ease of scaling up to large substrate sizes, is most likely the most important preparation technique of thermochromic VO₂-based coatings [9,12,13].

In our recent paper [25], we reported on high-performance three-layer thermochromic $ZrO_2/V_{0.982}W_{0.018}O_2/ZrO_2$ coatings prepared on soda-lime glass (SLG) using a pulsed magnetron sputtering at a relatively low substrate surface temperature $T_s = 330$ °C and without any substrate bias voltage. A coating design utilizing a second-order interference in the ZrO_2 AR layers was applied to increase both the luminous transmittance and the modulation of the solar energy transmittance [26]. The active $V_{0.982}W_{0.018}O_2$ layers were deposited by a controlled high-power impulse magnetron sputtering (HiPIMS) of a V target, combined with a simultaneous pulsed dc magnetron sputtering of a W target to reduce the transition temperature to $T_{tr} = 20-21$ °C, at $T_s = 330$ °C in an argon-oxygen gas mixture. The $ZrO_2/V_{0.982}W_{0.018}O_2/ZrO_2$ coatings exhibited T_{lum} up to 50% at ΔT_{sol} above 10% for a $V_{0.982}W_{0.018}O_2$ thickness of 69 nm. Here, it should be mentioned that HiPIMS techniques have the advantage of using essentially conventional magnetron sputtering equipment, while only replacing the power supply. Thus, these techniques can be implemented into industrial-size deposition systems with various target geometries [27].

In this paper, we present in more detail the fundamental principles of the deposition technique, which were reported only briefly in [25]. Here, the deposition technique was optimized for preparation of three-layer thermochromic $ZrO_2/V_{1-x}W_xO_2/ZrO_2$ coatings on ultrathin flexible glass (FG) and standard glass in large-scale deposition systems. A high maximum target power density in a pulse during the HiPIMS of the V target [25] was decreased three times to 1.4 kW cm⁻² at almost the same deposition-averaged target power density of 13.5 W cm⁻². To avoid atmospheric contamination and to reduce the total deposition time, the preparation of the three-layer $ZrO_2/V_{0.984}W_{0.016}O_2/ZrO_2$ coatings was performed in the same vacuum chamber without venting it to the atmosphere between the

depositions of individual layers. In addition, the structure, optical properties (not only transmittance as in [25], but also reflectance and absorption) and colors of these thermochromic coatings deposited onto 0.1 mm thick FG and 1 mm thick SLG substrates are presented.

The main aim of this work was to present a scalable technique for the preparation of strongly thermochromic VO₂-based coatings with a transition temperature of 22 °C on ultrathin FG at a relatively low substrate surface temperature $T_s = 330$ °C and without any substrate bias voltage.

Ultrathin FG was introduced to the market recently [28] as an alternative both to thick (rigid) glass sheets and to flexible polymer films. Its material properties are similar to those of rigid glass, but the thickness is much lower, ranging from 0.2 mm down to 0.02 mm [29]. In contrast to most polymers, the ultrathin FG is stable at much higher temperatures (even above 500 °C). Therefore, it can be used as a substrate for deposition of thermochromic VO₂-based coatings. Bendability of the ultrathin FG enables its high-volume processing using a roll-to-roll deposition device with magnetron sputter sources [28–30]. This opens up the possibility for new smart-window applications of thermochromic VO₂-based coatings, such as retrofitting of existing low-efficiency glass windows [24] and manufacturing of new high-efficiency insulated glass units with multiple functionalities.

2. Materials and Methods

2.1. Coating Preparation

In this study, the $ZrO_2/V_{0.984}W_{0.016}O_2/ZrO_2$ coatings were deposited onto 0.1 mm thick FG (Nippon Electric Glass Co., Ltd., Otsu, Japan) and 1 mm thick SLG substrates in argon-oxygen gas mixtures at the argon partial pressure $p_{ar} = 1$ Pa in an ultra-high vacuum sputter device with four unbalanced magnetrons, located symmetrically around the vacuum chamber axis (ATC 2200-V AJA International Inc., Scituate, MA, USA). The stainless-steel vacuum chamber (diameter of 560 mm and length of 430 mm) was evacuated by a turbomolecular pump (1200 L·s⁻¹) backed up with a double stage Roots pump (27 m³·h⁻¹). The base pressure before deposition was below 10⁻⁴ Pa. The rotating (20 rpm) substrates at the distance of 145 mm from the targets were at a floating potential. The values of p_{ar} and of $p_{ar} + p_{ox}$, where p_{ox} is the oxygen partial pressure in the chamber, were measured at the chamber wall using a high-stability capacitance manometer (Baratron, Type 127, MKS, Andover, MA, USA) with an accuracy much better than 1%.

The V_{0.984}W_{0.016}O₂ layers were deposited by controlled HiPIMS of a planar V target (99.9% purity, diameter of 50 mm and thickness of 6 mm), combined with a simultaneous pulsed dc magnetron sputtering of a planar W target (99.9% purity, diameter of 50 mm and thickness of 6 mm) at the substrate surface temperature $T_s = 330$ °C (see Figure 1). This T_s value was maintained during the deposition by a built-in infrared heating system calibrated using a thermocouple directly attached to the substrate surface (additional heating effect by the plasma included). The magnetron with the V target was driven by a unipolar high-power pulsed dc power supply (TruPlasma Highpulse 4002 TRUMPF Huettinger, Zielonka, Poland). In this work, the voltage pulse duration was 80 µs at a repetition frequency of 625 Hz (duty cycle of 5%) and the deposition-averaged target power density (spatially averaged over the total target area) was 13.5 W·cm⁻². The magnetron with the W target was driven by a unipolar bigh power supply (IAP-1010 EN Technologies Inc., Gunpo-si, Korea). The voltage pulse duration was 16 µs at a repetition frequency of 5 kHz (duty cycle of 8%) and the deposition-averaged target power density was 35 mW·cm⁻².

Oxygen was admitted into the vacuum chamber via mass flow controller and two corundum conduits (Figure 1). Two O_2 inlets with a diameter of 1 mm were placed symmetrically above the V target racetrack at the same distance of 20 mm from the V target surface and oriented to the substrate. The to-substrate O_2 injection into the dense plasma in front of the sputtered target is very important for reactive HiPIMS depositions of oxide films [31,32]. It leads to a substantially decreased compound fraction in the target surface layer [32,33], resulting in reduced arcing, increased sputtering of metal atoms, and low production of O^- ions at the target [31,34,35], and to a substantially increased

compound fraction in the substrate layer due to a significantly increased chemisorption flux of O atoms onto the substrate [32,33]. This is caused by a substantially (two to three times in [32]) increased local oxygen partial pressure (p'_{ox}) in front of the O₂ inlets and by a very high degree of dissociation of O₂ molecules in the high-density plasma in front of the target.



Figure 1. Schematic diagram of the magnetron sputter device showing two magnetrons with V and W targets, located in opposite positions, which were used for the deposition of the active $V_{0.984}W_{0.016}O_2$ layers. Two O_2 inlets were placed 20 mm from the V target surface and oriented to the substrate. Positions of the pressure sensor and the Ar inlet in the vacuum chamber wall are also shown. An increased local value of the oxygen partial pressure due to the O_2 injection is denoted as p'_{ox} .

The argon flow rate was 60 sccm corresponding to $p_{ar} = 1$ Pa, while the total oxygen flow rate (Φ_{ox}) in both conduits was not fixed but alternating between 1.45 and 1.85 sccm (see Figure 2). The moments of switching of the Φ_{ox} pulses were determined during the deposition by a programmable logic controller using a pre-selected critical value of the average discharge current on the V target in a period of the power supply (\bar{I}_d)_{cr} = 0.485 A: when $\bar{I}_d < (\bar{I}_d)_{cr}$, $\Phi_{ox} = 1.85$ sccm and when $\bar{I}_d \ge (\bar{I}_d)_{cr}$, $\Phi_{ox} = 1.45$ sccm.

The basic principle of the pulsed O₂ flow control is illustrated in Figure 2, which shows the time evolution of the magnetron voltage ($U_d(t)$) and the target current density ($J_t(t)$), averaged over the total target area, for both targets at the minimum ($p_{ox} \approx 0.02$ Pa) and maximum ($p_{ox} \approx 0.05$ Pa) value of the oxygen partial pressure in the vacuum chamber corresponding to the minimum and maximum \bar{I}_d , respectively, during the deposition. The much higher values of $J_t(t)$ for the V target at the highest p_{ox} , leading to a maximum V target power density of 1.4 kW cm⁻² in a pulse, can be explained [35]: (1) by a significantly enlarged flux of the O₂⁺ and O⁺ ions, arising in front of the partly oxidized V target (due to a higher p'_{ox}), and (2) by an increased secondary-electron emission yield of the V target with a larger compound coverage, particularly for the impacting O₂⁺ and O⁺ ions [36]. A slow decrease in $U_d(t)$ after the negative voltage pulse measured for the V target at $p_{ox} \approx 0.02$ Pa (see Figure 2) is a result of a slow discharging of the power supply capacitor owing to a very low $J_t(t)$ at the moment of the U_d pulse termination.



Figure 2. Waveforms of the magnetron voltage (U_d) and the target current density (J_t) for preset deposition-averaged target power densities of 13.5 Wcm⁻² and 35 mWcm⁻² for V and W target, respectively, during a deposition of the V_{0.984}W_{0.016}O₂ layers (the J_t values for W target are magnified 500 times). Time evolution of the average discharge current on the V target in a period of its power supply (\bar{I}_d) during the deposition is shown in the inset. A pre-selected critical value (\bar{I}_d)_{cr} = 0.485 A determining the switching between the oxygen flow rates $\Phi_{ox} = 1.45$ sccm and $\Phi_{ox} = 1.85$ sccm is marked by dots.

This feedback process control makes it possible to deliver a high power into discharge pulses without arcing on the V target surface and thus, to utilize two exclusive benefits of the HiPIMS discharges in the preparation of crystalline thermochromic VO₂-based layers on unbiased non-conductive substrates at relatively low T_s [34,35]. The first benefit is highly ionized fluxes of particles with many vanadium ions onto the substrate and enhanced energies (up to 50 eV relative to grounded electrode in [34,35]) of the ions bombarding the growing films, allowing us to achieve the VO₂ crystallinity at a relatively low T_s and without any substrate bias voltage. The second benefit is a very high degree of dissociation of the O₂ molecules injected into the high-density plasma in front of the V target, allowing us to achieve the required VO₂ stoichiometry at a low compound fraction in the V target surface layer.

The advantages of the used pulsed O_2 flow control are: (1) very high process stability (no problems with inertia of the inlet system, delay of valves and sensors, and hysteresis effects) as the controller does not try to keep one value but allows for a pre-selected interval of an output variable in accordance with the control-theory literature dealing with the control of non-linear systems [37], (2) simplicity, as no additional measurement devices (such as a plasma emission monitoring system, mass spectrometer or Lambda sensor) are needed, and (3) applicability to large-area coaters as a multi-segment O_2 injection control can be used.

Both the bottom and the top ZrO_2 AR layers (see Figure 3) were deposited by reactive pulsed dc magnetron sputtering without ohmic heating ($T_s < 60 \text{ °C}$). The p_{ox} was 0.12 Pa (oxide mode) at $\Phi_{ox} = 9$ sccm. The depositions were performed using an unbalanced magnetron with a planar Zr target (99.9% purity, diameter of 50 mm and thickness of 6 mm) driven by an asymmetric bipolar pulsed dc power

supply (TruPlasma Bipolar 4010, TRUMPF Huettinger, Freiburg, Germany) with a reverse (positive) voltage being close to 10% of the operational (negative) voltage of 530 to 560 V during the pulses. The negative voltage pulse duration was 6.25 μ s at a repetition frequency of 80 kHz (duty cycle of 50%) and the average target power density was 12.6 Wcm⁻² in these negative pulses.



Figure 3. The three-layer thermochromic VO₂-based coatings on ultrathin flexible glass (**a**) and soda-lime glass (**b**) substrates investigated in this paper. Here, h_b , h, and h_t represent the thickness of the bottom ZrO₂ layer, the thickness of the active V_{0.984}W_{0.016}O₂ layer, and the thickness of the top ZrO₂ layer, respectively. The refractive index (n_{550}) and extinction coefficient (k_{550}) at the wavelength of 550 nm of all layers were measured at $T_m = 23$ °C.

2.2. Coating Characterization

The W content in the metal sublattice of $V_{0.984}W_{0.016}O_2$, i.e., 1.6 ± 0.6 at.%, was measured on a dedicated 290 nm thick layer (using the same substrate) in a scanning electron microscope (SU-70, Hitachi, Kyoto, Japan) using wave-dispersive spectroscopy (Magnaray, Thermo Scientific, Waltham, MA, USA) at a low primary electron energy of 7 kV. The higher thickness of the dedicated film and the low primary energy collectively facilitate using the bulk procedure for processing the data. The room-temperature (25 °C) crystal structure of both $V_{0.984}W_{0.016}O_2$ and ZrO_2 was characterized by X-ray diffraction (XRD) using a PANalytical X'Pert PRO diffractometer (Malvern, UK) working with a CuK α (40 kV, 40 mA) radiation at a glancing incidence of 1°.

The thickness and room-temperature (23 °C) optical constants (refractive index, *n*, and extinction coefficient, *k*) of individual layers were measured by spectroscopic ellipsometry using the J.A. Woollam Co. Inc. (Lincoln, NE, USA) VASE instrument. The measurements were performed in the wavelength (λ) range 300–2000 nm (below we present properties at $\lambda = 550$ nm, n_{550} and k_{550}) at the angles of incidence of 55°, 60° and 65° in reflection. The optical model [26] consisted of the FG or SLG substrate, a bottom ZrO₂ layer, a bottom intermix layer, a V_{0.984}W_{0.016}O₂ layer, a top intermix layer, a top ZrO₂ layer and a surface roughness layer. The intermix layers and half of the surface roughness layer are included in the presented thickness values. The two ZrO₂ layers were expected to have the same properties and were described by the Cauchy dispersion formula. The V_{0.984}W_{0.016}O₂ layer was described by a combination of a Cody-Lorentz oscillator with Lorentz oscillators.

The coating transmittance (T) and reflectance (R) were measured by spectrophotometry using the Agilent CARY 7000 instrument (Mulgrave, Australia) and an in-house made heat/cool cell. In the case of the thick SLG substrate, the samples in the cell were freestanding. In the case of the thin FG substrate, the samples in the cell were supported by an underlying 1 mm thick glass in order to guarantee the homogeneity of the temperature field. There was most likely a very thin air layer at some parts of the contact area between the FG substrate and the supporting glass, not only leading to an interference

on FG, which is averaged out in Figures 4 and 6, but even affecting the averaged properties in one direction, as discussed in Section 3.4. The measurements were performed in the wavelength range 300–2500 nm (below we present not only the spectral dependences but also *T* at $\lambda = 2500$ nm, T_{2500}) at the angle of incidence of 7° (*R* could not be measured at 0°; however, it has been cross-checked that the *T* differences resulting from this angle difference can be safely neglected) and the results were averaged over the p- and s-polarization. The coating absorption (*A*) was calculated as 1-T-R. The measurement temperature (T_m) was varied between $T_{ms} = -10$ °C and $T_{mm} = 70$ °C (between the semiconducting and the metallic state of the $V_{0.984}W_{0.016}O_2$ layer, respectively). Five measurements of the spectral dependences $T(\lambda)$ and $R(\lambda)$, and the temperature dependence $T_{2500}(T_m)$ for each coating on the same substrate proved their full reproducibility.

The coating performance is quantified by means of integral luminous transmittance, reflectance and absorption (T_{lum} , R_{lum} and A_{lum} , respectively), their modulations (ΔT_{lum} , ΔR_{lum} and ΔA_{lum}), integral solar energy transmittance, reflectance and absorption (T_{sol} , R_{sol} and A_{sol} , respectively) and again their modulations (ΔT_{sol} , ΔR_{sol} and ΔA_{sol}). The quantities are defined as

$$T_{\rm lum}(T_{\rm m}) = \frac{\int_{380}^{780} \varphi_{\rm lum}(\lambda)\varphi_{\rm sol}(\lambda)T(T_{\rm m},\lambda)d\lambda}{\int_{380}^{780} \varphi_{\rm lum}(\lambda)\varphi_{\rm sol}(\lambda)d\lambda},$$
(1)

$$\Delta T_{\text{lum}} = T_{\text{lum}}(T_{\text{ms}}) - T_{\text{lum}}(T_{\text{mm}}), \qquad (2)$$

$$T_{\rm sol}(T_{\rm m}) = \frac{\int_{300}^{2500} \varphi_{\rm sol}(\lambda) T(T_{\rm m}, \lambda) d\lambda}{\int_{300}^{2500} \varphi_{\rm sol}(\lambda) d\lambda},$$
(3)

$$\Delta T_{\rm sol} = T_{\rm sol}(T_{\rm ms}) - T_{\rm sol}(T_{\rm mm}), \qquad (4)$$

and similarly, for the reflectance and for the absorption, where φ_{lum} is the luminous sensitivity of the human eye and φ_{sol} is the sea-level solar irradiance spectrum [38] at an air mass of 1.5.

The coating color and its temperature dependence was evaluated in the CIE XYZ color space (used to visualize the chromaticity in terms of x = X/[X + Y + Z] and y = Y/[X + Y + Z]) as well as in the related CIE L* a* b* color space (being more perceptually uniform). The evaluation was done for the CIE 1931 2° standard observer and the CIE standard daylight illuminant D65, both in transmission and in reflection (using the aforementioned $T(T_m)$ and $R(T_m)$, respectively, measured by spectrophotometry).

3. Results and Discussion

This section is divided into four subsections. In the first one, the design of the three-layer thermochromic $ZrO_2/V_{0.984}W_{0.016}O_2/ZrO_2$ coatings is presented (Figure 3) and explained. In the second one, their transition temperature, lowered to 22 °C, is presented (Figure 4). The third subsection deals with the crystal structure of these coatings (Figure 5). Finally, their optical properties are presented and discussed. The spectral transmittance, reflectance, and absorption of the coatings are shown in Figure 6, and the corresponding integral quantities are given in Table 1. The colors of these thermochromic coatings are shown in Figure 7.



Figure 4. Temperature ($T_{\rm m}$) dependence of the transmittance (T_{2500}) at 2500 nm for the $ZrO_2/V_{0.984}W_{0.016}O_2/ZrO_2$ coatings deposited on 0.1 mm thick FG and 1 mm thick SLG substrates (Figure 3). The transition temperature $T_{\rm tr} = 22$ °C is presented.



Figure 5. X-ray diffraction patterns taken at $T_{\rm m} = 25$ °C from the ZrO₂ (172 nm)/V_{0.984}W_{0.016}O₂ (71 nm)/ZrO₂ (178 nm) coating on 0.1 mm thick FG substrate and the ZrO₂ (175 nm)/V_{0.984}W_{0.016}O₂ (79 nm)/ZrO₂ (190 nm) coating on 1 mm thick SLG substrate. For comparison, the X-ray diffraction patterns from the corresponding ZrO₂ layers with the thickness $h_{\rm b} = 172$ nm and 175 nm are given. The main diffraction peaks of VO₂(R), VO₂(M1), ZrO₂(m) and ZrO₂(t) are marked.



Figure 6. Spectral transmittance (*T*), reflectance (*R*) and absorption (*A*) at an angle of incidence of 7° for the $ZrO_2/V_{0.984}W_{0.016}O_2/ZrO_2$ coatings on 0.1 mm thick FG (solid and dashed lines) and 1 mm thick SLG (dash-dotted and dotted lines) substrates at $T_{ms} = -10$ °C (significantly below $T_{tr} = 22$ °C) and $T_{mm} = 70$ °C (significantly above $T_{tr} = 22$ °C). The boundaries of the shaded areas represent the luminous sensitivity of the human eye (φ_{lum}) in panel (**a**), and the sea-level solar irradiance spectrum (φ_{sol}) at an air mass of 1.5 in panels (**b**,**c**), normalized to maxima of 100%.

Table 1. Thermochromic properties of the $ZrO_2/V_{0.984}W_{0.016}O_2/ZrO_2$ coatings deposited on 0.1 mm thick flexible glass (FG) and 1 mm thick soda-lime glass (SLG) substrates are presented. For the coatings on the FG substrate, the measured *T* values may be underestimated by up to 4% and the measured *R* values may be overestimated by approximately 1%, owing to the effect of the very thin air layer arising between the FG substrate and an underlying supporting 1 mm thick SLG used during the measurements (see Section 2.2).

Substrate	T _{lum} (T _{ms}) (%)	T _{lum} (T _{mm}) (%)	ΔT_{lum} (%)	T _{sol} (T _{ms}) (%)	T _{sol} (T _{mm}) (%)	ΔT_{sol} (%)
FG	45.7	42.2	3.5	39.6	30.0	9.6
SLG	48.6	45.5	3.1	43.1	32.4	10.7
Substrate	$R_{lum} (T_{ms})$ (%)	R _{lum} (T _{mm}) (%)	ΔR_{lum} (%)	R _{sol} (T _{ms}) (%)	R _{sol} (T _{mm}) (%)	ΔR_{sol} (%)
FG	8.2	8.3	-0.1	16.6	18.9	-2.3
SLG	6.4	7.1	-0.7	15.1	16.9	-1.8
Substrate	A _{lum} (T _{ms}) (%)	A _{lum} (T _{mm}) (%)	ΔA_{lum} (%)	A _{sol} (T _{ms}) (%)	A _{sol} (T _{mm}) (%)	ΔA_{sol} (%)
FG	46.1	49.5	-3.4	43.8	51.1	-7.3
SLG	45.0	47.4	-2.4	41.8	50.7	-8.9



Figure 7. Chromaticity of the coatings shown in the CIE xy chromaticity diagram. There are data below T_{tr} (full symbols) and above T_{tr} (empty symbols), obtained on the ultrathin FG substrate (triangles up) and SLG substrate (triangles down) using the transmittance (*T*) and the reflectance (*R*) presented in Figure 6.

3.1. Coating Design

As indicated by the coating formula description, the coating design shown in Figure 3 consists of ZrO_2 AR layers both below and above the thermochromic layer. Note that ZrO_2 layers are being increasingly applied to architectural glass as a protective overcoat for advanced low-emissivity stacks [39]. The specific advantages of ZrO_2 for the present purpose are: (1) a proper value of *n*, which is sufficiently close to the geometric mean of those of $V_{0.984}W_{0.016}O_2$ and air (top AR layer) as well as those of $V_{0.984}W_{0.016}O_2$ and FG or SLG (bottom AR layer) given in Figure 3, (2) a very small value of *k* in the λ range of interest (measured k_{550} —which also confirms the ZrO₂ stoichiometry—was on the order of 10^{-3} or lower), (3) achievable crystallinity at a low deposition temperature, which gives the bottom AR layer a second role of a template which improves the crystallinity of the thermochromic layer, and (4) high hardness (for an oxide) which gives the top AR layer the aforementioned second role of a mechanical protection of the thermochromic layer. Furthermore, the design of the coating and the pathway for its preparation is prone to maximize its environmental stability due to the densification of $V_{0.984}W_{0.016}O_2$ resulting from using HiPIMS [20] and due to the dual protection [21] of $V_{0.984}W_{0.016}O_2$

by both stable AR layers.

The coating performance strongly depends on the thickness of the bottom and the top AR layer (h_b and h_t , respectively). It has been detailed [26] and experimentally justified [25,26] that while the usual thin quarter-wavelength AR layers lead to a maximum T_{lum} at a very low ΔT_{sol} , thicker three quarter-wavelength AR layers lead to maxima of both T_{lum} (due to the second-order interference in the visible) and ΔT_{sol} (due to the first-order interference in the near infrared where most of the energy saving takes place). This leads to an optimum optical path length of ZrO₂ corresponding to, e.g., $h_b = h_t = 180$ nm at $n_{550} = 2.15$ (value considered in [26]) and similarly for other n_{550} values. Figure 3 shows that first, the optical path length of ZrO₂ is about the same on both FG and SLG (the coating on FG exhibits slightly lower h_b and h_t at slightly higher n_{550} and vice versa). Second, the optical path length is close to the optimum value (e.g., for the top AR layer on FG [180 × 2.15]/[178 × 2.09] = 104%) and consequently allows us to once again demonstrate the advantages of the aforementioned design (following the design, even more exactly, is desired for the industrialization, but this is beyond the aim of this paper).

3.2. Transition Temperature

Figure 4 shows the transition temperature (T_{tr}) of the coatings prepared on FG and on SLG, defined as the temperature leading to a half of the total T_{2500} change (averaged over the heating and the cooling part of the hysteresis curve). Let us recall that the hysteresis curves, and thus also the T_{tr} values, are fully reproducible for these stable thermochromic coatings (see Section 2.2). It can be seen that in agreement with the same elemental composition of the active layers in both coatings (guaranteed by a rotating substrate holder; see also Figure 3 for almost substrate-independent n_{550} and k_{550} of these layers), both coatings exhibit the same $T_{tr} = 22$ °C. Doping the metal sublattice by 1.6 at.% of W decreased T_{tr} from 57 °C, which was exhibited by thin films of pure VO₂ prepared by us using the same technique [25], i.e., at a gradient rate of -22 °C/at.% (consistent with the overview in [26]). Let us emphasize that the presented deposition technique allowed us not only to achieve a smooth and reproducible control of T_{tr} in itself, but it allowed us to achieve it at a preserved strength of the thermochromic effect (Section 3.4). The presented combination of a lowered T_{tr} with a high T_{lum} and ΔT_{sol} is unlike the previously reported lowering of T_{tr} at the cost of a drop of ΔT_{sol} [15,23]. Let alone those efforts to prepare VO₂-based coatings onto PET foils using hydrothermal synthesis, which did not lead to similarly low T_{tr} at all [24,40].

3.3. Coating Structure

The crystal structure of the bottom ZrO₂ layers and of the whole three-layer coatings is shown in the bottom and in the top part of Figure 5, respectively. Regarding bottom ZrO₂ layers by themselves, the most important piece of information is their crystallinity in itself, confirming its role of a template layer for the growth of $V_{0.984}W_{0.016}O_2$. There are mostly monoclinic ZrO₂ crystals (PDF#04-013-6875 [41]), with a contribution of tetragonal ZrO₂ crystals (PDF#01-081-1544 valid for ZrO_{1.95}, particularly visible in that XRD spectrum on FG which includes also the top ZrO₂ layer). Regarding $V_{0.984}W_{0.016}O_2$, the most important piece of information is that there is no trace of any undesired VO_x phase. All peaks are identified as diffraction peaks of either the low-temperature thermochromic phase VO₂(M1) (PDF#04-003-2035) or the high-temperature thermochromic phase VO₂(R) (PDF#01-073-2362). These two phases are difficult to distinguish and they are actually expected to be present simultaneously because the room-temperature XRD measurements took place almost exactly at the thermochromic transition temperature. See, e.g., the strongest peak around $2\Theta = 27.8^{\circ}$, corresponding well to theoretical positions of the VO₂(M1) (011) planes and the VO₂(R) (110) planes ($2\Theta = 27.80^{\circ}$ and 27.91° , respectively). In agreement with Figures 3 and 4, Figure 5 also does not show any convincing difference between the two coatings and confirms the successful transfer from the SLG to the ultrathin FG substrate.

3.4. Optical Properties

The spectral *T*, *R*, and *A* of the coatings are shown in Figure 6a–c, respectively, and the corresponding integral quantities are given in Table 1. Figure 6a clearly shows the desired thermochromic effect by itself: a strong modulation of *T* in the near infrared, where it is multiplied by still a relatively strong solar irradiance in the calculation of ΔT_{sol} according to Equations (3) and (4). This effect is combined with an almost preserved transmittance in the visible (ΔT_{lum} as low as $\approx 3\%$), which constitutes a comparative advantage of the presented solution over other types of energy-saving coatings. Figure 6b shows the effectiveness of the antireflection layers: *R* drops almost to zero around 550 nm (*R*_{lum} as low as $\approx 7\%$, second-order interference which guarantees high T_{lum}), and again around 1650 nm (first-order interference which guarantees high ΔT_{sol}). The positions of interference maxima of *T* well correspond to the positions of interference minima of *R* and vice versa. Furthermore, these positions are the same for both coatings, confirming the same optical path length of ZrO₂ layers.

Values of the most important integral quantities measured on SLG are: T_{lum} approaching 50% and ΔT_{sol} over 10%. Let us recall the application potential resulting from obtaining this result at a lowered $T_{\text{tr}} = 22 \,^{\circ}\text{C}$. Moreover, the recent literature overview in [25] even shows that the presented T_{lum} and ΔT_{sol} values are competitive with all reported values achieved at any T_{tr} up to 40 °C. Note Figure 6a, in which T measured on FG is slightly lower than that measured on SLG (with the complement to 100% observable as a slightly enhanced R and A). While this cannot be explained by the measured $V_{0.984}W_{0.016}O_2$ thickness and properties, it can be explained as an artifact of the aforementioned (Section 2.2) very thin air layer between thin FG and the supporting thick glass during the spectrophotometry. This air layer is not included in the optimized coating design, which maximizes T and consequently may decrease the measured T below that, which is very achievable. Optical calculations which include the air layer (not shown) indicate that T_{lum} may in that case be underestimated by $\approx 4\%$ (consistently with its measured difference shown in Table 1), at a comparatively weaker effect on ΔT_{sol} (where the underestimations of T below and above T_{tr} partially cancel out).

Furthermore, Figure 6 and Table 1 collectively demonstrate how the drop in *T*, resulting from the transition of both coatings from the semiconducting (-10 °C) to the metallic (70 °C) state, is divided between *R* (actually lowered around 800 nm but strongly enhanced around 1100 nm) and *A* (strongly enhanced above 1100 nm). Quantitatively, ΔT_{sol} of $\approx 10\%$ corresponds to ΔR_{sol} of $\approx 2\%$ (in agreement with the effectiveness of the antireflection layers) and ΔA_{sol} of $\approx 8\%$. This phenomenon, in particular the enhanced *A* leading to enhanced radiation, is important for future calculations of energy fluxes in the thermochromic coatings.

The chromaticity of the presented coatings, assuming daylight illumination from the coated side and no illumination from the backside, is shown in Figure 7. It can be seen that in transmission (smart window as seen from the inside) the coatings are yellowish, while in reflection (smart window as seen from the outside) the coatings are somewhat violet in color. From the application point of view, it is very important that the corresponding full and empty symbols in Figure 7 are very close to each other, i.e., that there is a negligible chromaticity change resulting from the thermochromic transition. This is consistent with negligible changes of optical properties in the visible as shown in Figure 6. Furthermore, the corresponding triangles up and triangles down in Figure 7 are almost overlapping (in transmission) or very close to each other (in reflection), which constitutes another confirmation of the successful transfer from SLG to ultrathin FG. The color coordinates in the perceptually uniform CIE L* a* b* color space (chromaticity given by *a** and *b** and lightness given by *L** between 0 and 100), averaged over the corresponding four datapoints in each case, include *L** = 73, *a** = -4, *b** = 39 in transmission and *L** = 33, *a** = 16, *b** = -30 in reflection. Note that in agreement with the non-linear nature of *L**, both its values are well above the values of *T*_{lum} and *R*_{lum}, respectively. All the details presented in Figures 6 and 7 complement the basic information given about this type of coatings in [25] and are relevant for the optimization and usage of smart windows.

4. Conclusions

We have presented and explained a scalable sputter deposition technique for the preparation of strongly thermochromic $ZrO_2/V_{0.984}W_{0.016}O_2/ZrO_2$ coatings on ultrathin FG and standard SLG at a relatively low substrate surface temperature (330 °C) and without any substrate bias voltage. The $V_{0.984}W_{0.016}O_2$ layers were deposited by a controlled HiPIMS of a V target, combined with a simultaneous pulsed dc magnetron sputtering of a W target (doping of VO₂ by W to reduce the transition temperature without any degradation of thermochromic properties). The coatings exhibit a transition temperature of 22 °C, a luminous transmittance approaching 50%, a modulation of the solar energy transmittance over 10% and a temperature, has not yet been published by other teams for thermochromic VO₂-based coatings in the literature. This deposition technique is compatible with the existing magnetron sputtering systems in glass production lines and in roll-to-roll deposition devices. Its successful transfer to a large-scale roll-to-roll device producing coatings on ultrathin FG would open up the possibility for new smart-window applications of thermochromic VO₂-based coatings, such as the retrofitting of existing low-efficiency glass windows and the manufacturing of new high-efficiency insulated glass units with multiple functionalities.

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