# QUADRATIC OPTIMAL CONTROL OF STABLE WELL-POSED LINEAR SYSTEMS 

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#### Abstract

We consider the infinite horizon quadratic cost minimization problem for a stable time-invariant well-posed linear system in the sense of Salamon and Weiss, and show that it can be reduced to a spectral factorization problem in the control space. More precisely, we show that the optimal solution of the quadratic cost minimization problem is of static state feedback type if and only if a certain spectral factorization problem has a solution. If both the system and the spectral factor are regular, then the feedback operator can be expressed in terms of the Riccati operator, and the Riccati operator is a positive self-adjoint solution of an algebraic Riccati equation. This Riccati equation is similar to the usual algebraic Riccati equation, but one of its coefficients varies depending on the subspace in which the equation is posed. Similar results are true for unstable systems, as we have proved elsewhere.


## 1. Introduction

This work treats the infinite horizon quadratic cost minimization problem for a time-invariant well-posed linear system in the sense of Salamon and Weiss, and extends the results presented in [37] to externally stable Salamon-Weiss systems (the unstable case is discussed in [40]). The main idea behind this work is the same as in [37], and it can be described briefly as follows.

It is well-known that it is possible to solve a certain canonical Wiener-Hopf factorization problem by a state space method. The problem is to factor $\mathcal{D}^{*} \mathcal{D}$, where $\mathcal{D}$ is a stable input/output map and $\mathcal{D}^{*}$ is its adjoint, in the form

$$
\mathcal{X}^{*} \mathcal{X}=\mathcal{D}^{*} \mathcal{D}
$$

where $\mathcal{X}$ is a stable input/output map from the control space $U$ into itself with a stable inverse $\mathcal{X}^{-1}$. The solution $\mathcal{X}$ is called a spectral factor of $\mathcal{D}^{*} \mathcal{D}$. In the state space approach we first construct a realization $\Psi$ of $\mathcal{D}$ (or alternatively, we assume that a system $\Psi$ is given whose input/output map is $\mathcal{D}$ ). Then we pose an infinite horizon quadratic cost minimization problem in the state space $H$ of $\Psi$, and solve the algebraic Riccati equation to get the Riccati operator $\Pi$ describing the optimal cost. From the Riccati operator we get the feedback operator $K$, and this feedback operator is then used to construct the spectral factor $\mathcal{X}$.

Although the state space solution of the factorization problem is the most common one today, it is not the only one available, and it is not universally applicable.

[^0]In particular, it requires a certain smoothness of the symbol of the input/output map $\mathcal{D}$. Without this smoothness every possible realization of $\mathcal{D}$ must have both an unbounded control operator $B$ and an unbounded observation operator $C$ (with respect to the same space), and there has been no general Riccati equation theory available for that case.

The spectral factorization problem always has a solution (see Lemma 18 below), even though it has not been known how to compute this factorization in general by a state space method. Furthermore, there exist some independent ways of computing the factorization; see, e.g., [16] and [51]. We take advantage of this fact and study the converse of the problem presented above, i.e., we use the spectral factorization to solve the quadratic cost minimization problem in feedback form. More specifically, in Theorem 27 we show that these two problems are equivalent in the sense that the factorization problem has a solution if and only if the solution of the quadratic cost minimization problem is of static state feedback type. Moreover, there is a simple relation between the spectral factor and the feedback solution, which means that they can easily be computed from each other.

In order to derive the Riccati equation solution to the optimal feedback problem we have to make one additional assumption, namely that both the input/output map $\mathcal{D}$ and the spectral factor $\mathcal{X}$ are regular together with their adjoints in the sense of [45] (a much stronger version of this assumption holds in the classical case). We show that in this case the Riccati operator of an arbitrary realization $\Psi$ of $\mathcal{D}$ satisfies a certain algebraic Riccati equation, and that the feedback operator $K$ can be computed from $\Pi$. However, this Riccati equation contains one unexpected feature (first discovered in [37]): the positive self-adjoint weighting operator in the quadratic term in the Riccati equation is $\left(X^{*} X\right)^{-1}$, where $X$ is the feed-through operator of the spectral factor $\mathcal{X}$, instead of the expected $\left(D^{*} D\right)^{-1}$, where $D$ is the feed-through operator of the input/output map $\mathcal{D}$. If the control and observation operators are bounded then $X^{*} X=D^{*} D$, so the new theory agrees with the old one in this respect. However, it is not true in general that $X^{*} X=D^{*} D$, and the computation of $X$ is a nontrivial task. In some cases $X$ can be computed from a discrete time Riccati equation [41], and there is also a fairly general method available for the case of a scalar control (see [18, Chapter VII]). An alternative formula for $X^{*} X$ is given in [40].

As explained above, under certain regularity assumptions we prove that the Riccati operator (the optimal cost operator) necessarily satisfies a certain Riccati equation. In the finite dimensional case this is the "difficult" direction to prove, the converse direction being almost trivial. In our case the converse direction is still open, partly due to the fact that the Riccati equation that we get is much more complicated than the classical one. In some sense it resembles the discrete time Riccati equation; cf. [41]. Moreover, because of the required regularity hypothesis, it is still an unsettled question what the appropriate conditions are for our results to apply to boundary control problems for partial differential equations. For time delay systems the regularity problem is less severe; see the examples in [38].

We work formally in the time domain the whole time, but some of our proofs are adapted frequency domain proofs, with some added state space features. The key addition is the facorization of the Hankel operator induced by the input/output map as the product of the controllability and observability maps. This makes it possible to connect the state space and the frequency domain theories to each other. See Definition 1 and the paragraph following Definition 4.

The first preprints of this paper and [40] were circulated in the spring of 1995, and a preliminary version was presented in the summer of 1995 at the conference [36]. In the spring of 1996 we received the preprint [50] where some of the results presented here and in [40] are also found. That preprint uses quite different notations, it deals with weakly regular systems, it contains some additional examples (one of which is essentially the same as the one in [41]), and it carries the analysis further than we do before resorting to spectral factorization. In particular, [50, Proposition 10.5] contains one additional Riccati equation valid on the domain of the closed loop generator in the case where the observation operator is bounded.

There is a significant overlap between some of the results presented here (those that are not related to spectral factorization) and those in [17]. In that paper the output operator $C$ is bounded, but the system need not be stable, and no regularity assumptions are imposed on the spectral factor (the spectral factor is not even mentioned). The well-posedness of the closed loop system (our first main result, presented in Theorem 27) is not studied in [17]. The Riccati equation given in [17, Theorem 2.2] looks different from our Riccati equation (it does not contain the nonstandard term), due to the fact that [17, Corollary 4.9] extends the operator $B^{*} \Pi$ in a different way than our Proposition 36. An example highlighting this difference is given in [47].

For a further discussion of the existing literature we refer the reader to [38], [40]. The recent literature includes (but is certainly not restricted to) [1], [2], [4], [5], [6], [7], [9], [12] [13], [14], [20], [21], [23], [24], [25], [26], [27], [28], [30], [31], [42], [48], [49], [51], and the other papers in our list of references.

We use the following notation:
$\mathcal{L}(U ; Y), \mathcal{L}(U): \quad$ The set of bounded linear operators from $U$ into $Y$ or from $U$ into itself, respectively.
$I: \quad$ The identity operator.
$A^{*}: \quad$ The (Hilbert space) adjoint of the operator $A$.
$\operatorname{dom}(A)$ : The domain of the (unbounded) operator $A$.
range $(A)$ : $\quad$ The range of the operator $A$.
$\mathbf{R}, \mathbf{R}^{+}, \mathbf{R}^{-}: \quad \mathbf{R}=(-\infty, \infty), \mathbf{R}^{+}=[0, \infty)$, and $\mathbf{R}^{-}=(-\infty, 0]$.
$L^{2}(J ; U)$ : $\quad$ The set of $U$-valued $L^{2}$-functions on $J$.
$W^{1,2}(J ; U)$ : The set of functions in $L^{2}(J ; U)$ with a (distribution) derivative in $L^{2}(J ; U)$.
$C(J ; U), C^{1}(J ; U): \quad C(J ; U)$ is the set of $U$-valued continuous functions on $J$, and $C^{1}(J ; U)$ is the set of $U$-valued continuously differentiable functions on $J$.
$B C(J ; U), B C^{1}(J ; U): \quad B C(J ; U)$ is the set of $U$-valued bounded continuous functions on $J$, and $B C^{1}(J ; U)$ is the set of $U$-valued bounded continuously differentiable functions on $J$ with a bounded derivative.
$\langle\cdot, \cdot\rangle_{H}: \quad$ The inner product in the Hilbert space $H$.
$\tau(t): \quad$ The bilateral time shift operator $\tau(t) u(s)=u(t+s)$ (this is a left-shift when $t>0$ and a right-shift when $t<0)$.
$\pi_{J}: \quad\left(\pi_{J} u\right)(s)=u(s)$ if $s \in J$ and $\left(\pi_{J} u\right)(s)=0$ if $s \notin J$. Here $J$ is a subset of $\mathbf{R}$. This operator is used both as a projection operator $L^{2}(\mathbf{R}) \rightarrow L^{2}(\mathbf{R})$ and as an embedding operator $L^{2}(J) \rightarrow L^{2}(\mathbf{R})$.
$\pi_{+}, \pi_{-}: \quad \pi_{+}=\pi_{\mathbf{R}^{+}}$and $\pi_{-}=\pi_{\mathbf{R}^{-}}$.

We extend an $L^{2}$-function $u$ defined on a subinterval $J$ of $\mathbf{R}$ to the whole real line by requiring $u$ to be zero outside of $J$, and we denote the extended function by $\pi_{J} u$. Thus, we use the same symbol $\pi_{J}$ both for the embedding operator $L^{2}(J) \rightarrow L^{2}(\mathbf{R})$ and for the corresponding orthogonal projection operator $L^{2}(\mathbf{R}) \rightarrow \operatorname{range}\left(\pi_{J}\right)$. With this interpretation, $\pi_{+} L^{2}(\mathbf{R} ; U)=L^{2}\left(\mathbf{R}^{+} ; U\right) \subset L^{2}(\mathbf{R} ; U)$ and $\pi_{-} L^{2}(\mathbf{R} ; U)=$ $L^{2}\left(\mathbf{R}^{-} ; U\right) \subset L^{2}(\mathbf{R} ; U)$.

## 2. Well-Posed Linear Systems and Time-Invariant Operators

In order to fix the notation and describe the basic setting we first give a brief presentation of the Salamon-Weiss class of time-invariant well-posed linear systems, including only those parts of the theory that we need. This theory has been developed in [33], [34], [35], [8], [11], and [43], [44], [45], [46] (and many other papers), and we refer the reader to these sources for additional reading. (Salamon calls these systems "well-posed semigroup control systems" and Weiss calls them "abstract linear systems".) Since we need only externally stable systems here, we restrict our discussion to this class of systems.

In the externally stable case it is possible to define a well-posed linear system in a slightly nonstandard way that turns out to be very convenient for our purposes. Usually the axioms of the system are formulated in terms of operators defined on local $L^{2}$-spaces, but in the stable case it is much more convenient to work with global $L^{2}$-spaces. We shall do so throughout, since it simplifies the formulation significantly; in particular, it leads to a much simpler dual theory. This setting is also the best one for $H^{\infty}$-related work. Parts of this setting are found in [35].

In order to formulate the axioms satisfied by a well-posed linear system we need the "past time" projection operator $\pi_{-}$, the "future time" projection operator $\pi_{+}$, and the bilateral "time shift" group $\tau(t)$. These operate on functions $u$ defined on $\mathbf{R}=(-\infty, \infty)$ in the following way:

$$
\begin{aligned}
\left(\pi_{-} u\right)(s) & = \begin{cases}u(s), & s \in \mathbf{R}^{-} \\
0, & s \in \mathbf{R}^{+}\end{cases} \\
\left(\pi_{+} u\right)(s) & = \begin{cases}u(s), & s \in \mathbf{R}^{+} \\
0, & s \in \mathbf{R}^{-}\end{cases} \\
(\tau(t) u)(s)=u(t+s), & t, s \in \mathbf{R}
\end{aligned}
$$

Definition 1. Let $U, H$, and $Y$ be Hilbert spaces. A (causal) externally stable (time-invariant) well-posed linear system on $(U, H, Y)$ is a quadruple $\Psi=\left[\begin{array}{cc}\mathcal{A} & \mathcal{B} \\ \mathcal{C} & \mathcal{D}\end{array}\right]$, where $\mathcal{A}, \mathcal{B}, \mathcal{C}$, and $\mathcal{D}$ are bounded linear operators of the following type:
(i) $\mathcal{A}$ is a strongly continuous semigroup on $H$;
(ii) $\mathcal{B}: L^{2}(\mathbf{R} ; U) \rightarrow H$ satisfies $\mathcal{A}(t) \mathcal{B} u=\mathcal{B} \tau(t) \pi_{-} u$ for all $u \in L^{2}(\mathbf{R} ; U)$ and $t \in \mathbf{R}^{+} ;$
(iii) $\mathcal{C}: H \rightarrow L^{2}(\mathbf{R} ; Y)$ satisfies $\mathcal{C} \mathcal{A}(t) x=\pi_{+} \tau(t) \mathcal{C} x$ for all $x \in H$ and $t \in \mathbf{R}^{+}$;
(iv) $\mathcal{D}: L^{2}(\mathbf{R} ; U) \rightarrow L^{2}(\mathbf{R} ; Y)$ satisfies $\tau(t) \mathcal{D} u=\mathcal{D} \tau(t) u, \pi_{-} \mathcal{D} \pi_{+} u=0$, and $\pi_{+} \mathcal{D} \pi_{-} u=\mathcal{C B} u$ for all $u \in L^{2}(\mathbf{R} ; U)$ and $t \in \mathbf{R}$.
If, furthermore, $\sup _{t \in \mathbf{R}^{+}}\|\mathcal{A}(t)\|<\infty$, then $\Psi$ is called stable rather than externally stable, and if $\mathcal{A}(t) x \rightarrow 0$ as $t \rightarrow \infty$ for all $x \in H$, then $\Psi$ is called strongly stable.

The different components of $\Psi$ are named as follows: $U$ is the input space, $H$ the state space, $Y$ the output space, $\mathcal{A}$ the semigroup, $\mathcal{B}$ the controllability (or reachability) map, $\mathcal{C}$ the observability map, $\mathcal{D}$ the input/output map, $\mathcal{B B}^{*}$ the controllability gramian, and $\mathcal{C}^{*} \mathcal{C}$ the observability gramian of $\Psi$.

Remark 2. The same axioms can be used to define an arbitrary (unstable) wellposed linear system. We simply replace all the $L^{2}$-spaces by the corresponding weighted $L^{2}$-spaces $L_{\omega}^{2}=\left\{f \in L_{\text {loc }}^{2} \mid\left(t \mapsto \mathrm{e}^{-\omega t} f(t)\right) \in L^{2}\right\}$. Here $\omega$ is an arbitrary number larger than the exponential growth rate of $\mathcal{A}$. See [39, Section 2] for details.

The axioms listed above describe standard properties of the corresponding maps induced by exponentially stable systems with bounded control and observation operators. Whenever we refer to a stable "classical" system, we mean a system of the following type: we let $A$ be the generator of an exponentially stable semigroup $\mathcal{A}$ on a Hilbert space $H$, let $U$ and $Y$ be Hilbert spaces, let $B \in \mathcal{L}(U ; H), C \in \mathcal{L}(H ; Y)$, and $D \in \mathcal{L}(U ; Y)$, and consider the system

$$
\begin{align*}
x^{\prime}(t) & =A x(t)+B u(t), \\
y(t) & =C x(t)+D u(t), \quad t \geq T,  \tag{1}\\
x(T) & =x_{T},
\end{align*}
$$

where $T$ is a given initial time and $x_{T}$ a given initial value. We call $u$ the control, $x$ the state, $y$ the output (or observation), $A$ the generator, $B$ the control operator, $C$ the observation operator, and $D$ the feed-through operator of this classical system. The state $x$ is required to be a strong solution of (1), i.e., the state $x$ and output $y$ are given by

$$
\begin{align*}
& x(t)=\mathcal{A}(t-T) x_{T}+\int_{T}^{t} \mathcal{A}(t-s) B u(s) d s, \quad t \geq T \\
& y(t)=C \mathcal{A}(t-T) x_{T}+\int_{T}^{t} C \mathcal{A}(t-s) B u(s) d s+D u(t), \quad t \geq T \tag{2}
\end{align*}
$$

In this case we define $\mathcal{B}, \mathcal{C}$, and $\mathcal{D}$ by

$$
\begin{align*}
\mathcal{B} u & =\int_{-\infty}^{0} \mathcal{A}(-s) B u(s) d s  \tag{3}\\
\mathcal{C} x & =\left(t \mapsto C \mathcal{A}(t) x, \quad t \in \mathbf{R}^{+}\right)  \tag{4}\\
\mathcal{D} u & =\left(t \mapsto \int_{-\infty}^{t} C \mathcal{A}(t-s) B u(s) d s+D u(t), \quad t \in \mathbf{R}\right) \tag{5}
\end{align*}
$$

Thus, $\mathcal{B}$ is the mapping from the control $u \in L^{2}\left(\mathbf{R}^{-} ; U\right)$ to the final state $x(0) \in H$ (take $T=-\infty, x_{T}=0$, and $t=0$ ) $\mathcal{C}$ is the mapping from the initial state $x_{0} \in H$ to the output $y \in L^{2}\left(\mathbf{R}^{+} ; Y\right)$ (take $T=0$ and $u=0$ ), and $\mathcal{D}$ is the mapping from the control $u \in L^{2}(\mathbf{R} ; U)$ to the output $y \in L^{2}(\mathbf{R} ; Y)$ (take $T=-\infty$ and $x_{T}=0$ ).

We shall also need the concept of an input/output map $\mathcal{D}$ without knowing that it is a part of a well-posed linear system. In this case we apply the following definition:

Definition 3. Let $U$ and $Y$ be two Hilbert spaces. A bounded linear operator $\mathcal{D}: L^{2}(\mathbf{R} ; U) \rightarrow L^{2}(\mathbf{R} ; Y)$ is called time-invariant iff it commutes with bilateral time shifts, i.e., $\tau(t) \mathcal{D} u=\mathcal{D} \tau(t) u$ for all $u \in L^{2}(\mathbf{R} ; U)$ and all $t \in \mathbf{R}$. The Hankel operator induced by $\mathcal{D}$ is the operator $\pi_{+} \mathcal{D} \pi_{-}$, and the anti-Hankel operator induced
by $\mathcal{D}$ is the operator $\pi_{-} \mathcal{D} \pi_{+}$. The Toeplitz operator induced by $\mathcal{D}$ is the operator $\pi_{+} \mathcal{D} \pi_{+}$, and the anti-Toeplitz operator induced by $\mathcal{D}$ is the operator $\pi_{-} \mathcal{D} \pi_{-}$.

The word "causal" that we have included in the definition of a well-posed linear system relates to the fact that all the components of $\Psi$ in Definition 1 are causal:

Definition 4. An operator $\mathcal{B}: L^{2}(\mathbf{R} ; U) \rightarrow H$ is causal [anti-causal] if $\mathcal{B} \pi_{+}=0$ $\left[\mathcal{B} \pi_{-}=0\right]$. An operator $\mathcal{C}: H \rightarrow L^{2}(\mathbf{R} ; Y)$ is causal [anti-causal] if $\pi_{-} \mathcal{C}=0$ $\left[\pi_{+} \mathcal{C}=0\right]$. An time-invariant operator $\mathcal{D}: L^{2}(\mathbf{R} ; U) \rightarrow L^{2}(\mathbf{R} ; Y)$ is causal [anticausal $]$ if $\pi_{-} \mathcal{D} \pi_{+}=0\left[\pi_{+} \mathcal{D} \pi_{-}=0\right]$, and it is static (or memoryless) if it is both causal and anti-causal.

Thus, a time-invariant operator is causal iff its anti-Hankel operator vanishes, it is anti-causal iff its Hankel operator vanishes, and it is static if both its Hankel operator and its anti-Hankel operator vanish. The condition imposed on the input/output map $\mathcal{D}$ in Definition 1 requires $\mathcal{D}$ to be a causal time-invariant operator whose Hankel operator is equal to $\mathcal{C B}$. Intuitively, a causal controllability map $\mathcal{B}$ maps past inputs into the present state, a causal observability map $\mathcal{C}$ maps the present state into future outputs, and future inputs to a causal input/output map $\mathcal{D}$ have no effect on past outputs.

The following result is immediate:
Lemma 5. Let $\mathcal{D}: L^{2}(\mathbf{R} ; U) \rightarrow L^{2}(\mathbf{R} ; Y)$ be time-invariant.
(i) The adjoint $\mathcal{D}^{*}: L^{2}(\mathbf{R} ; Y) \rightarrow L^{2}(\mathbf{R} ; U)$ of $\mathcal{D}$ is time-invariant.
(ii) If $\mathcal{D}$ is invertible, then the inverse $\mathcal{D}^{-1}$ is time-invariant (but the causality of $\mathcal{D}$ does not imply the causality of $\mathcal{D}^{-1}$ ).
(iii) If $\mathcal{D}$ is causal then $\mathcal{D} \pi_{+}=\pi_{+} \mathcal{D} \pi_{+}$and $\pi_{-} \mathcal{D}=\pi_{-} \mathcal{D} \pi_{-}$, and if $\mathcal{D}$ is anti-causal then $\mathcal{D} \pi_{-}=\pi_{-} \mathcal{D} \pi_{-}$and $\pi_{+} \mathcal{D}=\pi_{+} \mathcal{D} \pi_{+}$.

Below we use this lemma repeatedly without explicit reference.
(Different authors use different names for the same concepts. For example, [32] use the names " $S$-analytic" for a causal time-invariant operator and " $S$-constant" for a static time-invariant operator, and they identify the class of anti-causal timeinvariant operators with the class of operators whose adjoints are $S$-analytic. Here $S$ stands for the (unilateral) Laguerre shift on $L^{2}\left(\mathbf{R}^{+} ; U\right)$; see [32, p. 16].)

Static time-invariant operators have a very simple characterization:
Lemma 6. The set of bounded linear static time-invariant operators $\mathcal{D}: L^{2}(\mathbf{R} ; U) \rightarrow L^{2}(\mathbf{R} ; Y)$ can be identified with the set of bounded linear operators $M \in \mathcal{L}(U ; Y)$ in the following sense: For every $M \in \mathcal{L}(U ; Y)$ the multiplication operator $u \mapsto M u$ induced by $M$ is a static time-invariant operator mapping $L^{2}(\mathbf{R} ; U)$ into $L^{2}(\mathbf{R} ; Y)$, and conversely, every static time-invariant operator $\mathcal{D}$ has a representation of this type.

This well-known result can be derived from, e.g., [32, Theorem 5.2C, p. 96] or [15, Chapter 1].

Corollary 7. The Hankel operator of a causal [the anti-Hankel operator of an anticausal] time-invariant operator $\mathcal{D}: L^{2}(\mathbf{R} ; U) \rightarrow L^{2}(\mathbf{R} ; Y)$ determines $\mathcal{D}$ uniquely, modulo a multiplication operator $u \mapsto M u$, where $M \in \mathcal{L}(U ; Y)$.

Proof. If two different causal time-invariant operators have the same Hankel operator, then their difference is static.

Thus, the control and observation operators $\mathcal{B}$ and $\mathcal{C}$ of a well-posed linear system determine the input/output map $\mathcal{D}$ uniquely, except for an undetermined multiplication operator (which in the classical case is represented by the feed-through operator $D$ in (1)). We remark that every bounded causal time-invariant operator $\mathcal{D}$ can be represented as the input/output map of a strongly stable well-posed linear system $\Psi$; see [35, Theorem 4.3] or [39, Definition 2.10]. Such a system $\Psi$ is called a realization of $\mathcal{D}$.

The adjoint of the system $\Psi$ can be defined in two different ways, that differ from each other through a reflection of the time axis. We shall follow the classical tradition and let the adjoint system evolve in the backward time direction. To get this system we use the standard inner products in $L^{2}(\mathbf{R} ; U)$ and $L^{2}(\mathbf{R} ; Y)$ (instead of inner products that reflect the time axis), replace all the operators in Definition 1 by their adjoints, and finally interchange $\mathcal{B}^{*}$ and $\mathcal{C}^{*}$. Doing so we get a system of the following type:

Definition 8. Let $Y, H$, and $U$ be Hilbert spaces. An anti-causal externally stable well-posed linear system on $(Y, H, U)$ is a quadruple $\Psi^{*}=\left[\begin{array}{cc}\mathcal{A}^{*} & \mathcal{C}^{*} \\ \mathcal{B}^{*} & \mathcal{D}^{*}\end{array}\right]$, where $\mathcal{A}^{*}, \mathcal{C}^{*}$, $\mathcal{B}^{*}$, and $\mathcal{D}^{*}$ are bounded linear operators of the following type:
(i) $\mathcal{A}^{*}$ is a strongly continuous semigroup on $H$;
(ii) $\mathcal{C}^{*}: L^{2}(\mathbf{R} ; Y) \rightarrow H$ satisfies $\mathcal{A}^{*}(-s) \mathcal{C}^{*} y^{*}=\mathcal{C}^{*} \tau(s) \pi_{+} y^{*}$ for all $y^{*} \in L^{2}(\mathbf{R} ; Y)$ and $s \in \mathbf{R}^{-}$;
(iii) $\mathcal{B}^{*}: H \rightarrow L^{2}(\mathbf{R} ; U)$ satisfies $\mathcal{B}^{*} \mathcal{A}^{*}(-s) x^{*}=\pi_{-} \tau(s) \mathcal{B}^{*} x^{*}$ for all $x^{*} \in H$ and $s \in \mathbf{R}^{-}$;
(iv) $\mathcal{D}^{*}: L^{2}(\mathbf{R} ; Y) \rightarrow L^{2}(\mathbf{R} ; U)$ satisfies $\tau(s) \mathcal{D}^{*} y^{*}=\mathcal{D}^{*} \tau(s) y^{*}, \pi_{+} \mathcal{D}^{*} \pi_{-} y^{*}=0$, and $\pi_{-} \mathcal{D}^{*} \pi_{+} y^{*}=\mathcal{B}^{*} \mathcal{C}^{*} y^{*}$ for all $y^{*} \in L^{2}(\mathbf{R} ; Y)$ and $s \in \mathbf{R}$.
If, furthermore, $\sup _{t \in \mathbf{R}^{+}}\left\|\mathcal{A}^{*}(t)\right\|<\infty$, then $\Psi^{*}$ is called stable rather than externally stable, and if $\mathcal{A}^{*}(t) x^{*} \rightarrow 0$ as $t \rightarrow \infty$ for all $x^{*} \in H$, then $\Psi^{*}$ is called strongly stable.

The different components of $\Psi^{*}$ are named as follows: $Y$ is the input space, $H$ the state space, $U$ the output space, $\mathcal{A}^{*}$ the semigroup, $\mathcal{C}^{*}$ the controllability map, $\mathcal{B}^{*}$ the observability map, $\mathcal{D}^{*}$ the input/output map, $\mathcal{C}^{*} \mathcal{C}$ the controllability gramian, and $\mathcal{B B}^{*}$ the observability gramian of $\Psi^{*}$.

It is easy to see that the adjoint of the causal system $\Psi$ in Definition 1 is an anticausal system $\Psi^{*}$ of the type described in Definition 8. The two key observations are that $\pi_{+}$and $\pi_{-}$are self-adjoint, and that the adjoint of the left-shift $\tau(t)$ is the right-shift $\tau(-t)$. The reason for the name "anti-causal" is obvious: all the parts of the anti-causal system are anti-causal according to Definition 4. By reversing the time direction in which $\Psi^{*}$ evolves we get a causal system of the type described in Definition 1 (i.e., let $R$ be the reflection operator $R u(t)=u(-t)$, and replace $\mathcal{C}^{*}, \mathcal{B}^{*}$, and $\mathcal{D}^{*}$ by $\mathcal{C}^{*} R, R \mathcal{B}^{*}$, and $R \mathcal{D}^{*} R$, respectively). Thus, all results proved for causal systems can be applied to anti-causal systems, too (with some trivial modifications due to the reflection of the time direction).

In the classical stable case (1) the adjoint system $\Psi^{*}$ is induced by the (antistable) system

$$
\begin{align*}
\left(x^{*}\right)^{\prime}(s) & =-A^{*} x^{*}(s)-C^{*} y^{*}(s) \\
u^{*}(s) & =B^{*} x^{*}(s)+D^{*} y^{*}(s), \quad s \leq S  \tag{6}\\
x^{*}(S) & =x_{S}^{*}
\end{align*}
$$

whose strong solution $x$ and output $y$ are given by

$$
\begin{align*}
& x^{*}(s)=\mathcal{A}^{*}(S-s) x_{S}^{*}+\int_{s}^{S} \mathcal{A}^{*}(t-s) C^{*} y^{*}(t) d t, \quad s \leq S  \tag{7}\\
& u^{*}(s)=B^{*} \mathcal{A}^{*}(S-s) x_{S}^{*}+\int_{s}^{S} B^{*} \mathcal{A}^{*}(t-s) C^{*} y^{*}(t) d t+D^{*} y^{*}(s), \quad s \leq S
\end{align*}
$$

The different maps of $\Psi^{*}$ are now given by

$$
\begin{align*}
\mathcal{C}^{*} y^{*} & =\int_{0}^{\infty} \mathcal{A}^{*}(t) C^{*} y^{*}(t) d t  \tag{8}\\
\mathcal{B}^{*} x^{*} & =\left(s \mapsto B^{*} \mathcal{A}^{*}(-s) x^{*}, \quad s \in \mathbf{R}^{-}\right)  \tag{9}\\
\mathcal{D}^{*} y^{*} & =\left(s \mapsto \int_{s}^{\infty} B^{*} \mathcal{A}^{*}(t-s) C^{*} y^{*}(t) d t+D^{*} y^{*}(s), \quad s \in \mathbf{R}\right) \tag{10}
\end{align*}
$$

Two of the central notions of a well-posed linear system are still missing, namely the notions of the controlled state and the output:

Definition 9. Let $\Psi=\left[\begin{array}{cc}\mathcal{A} & \mathcal{B} \\ \mathcal{C}\end{array}\right]$ be a causal externally stable well-posed linear system on $(U, H, Y)$, and let $\Psi^{*}=\left[\begin{array}{cc}\mathcal{A}^{*} & \mathcal{C}^{*} \\ \mathcal{D}^{*}\end{array}\right]$ be its anti-causal adjoint system on $(Y, H, U)$. Let $x_{0} \in H, x_{0}^{*} \in H, u \in L^{2}(\mathbf{R} ; U)$, and $y^{*} \in L^{2}(\mathbf{R} ; Y)$. In the time-invariant setting the controlled state $x(t)$ at time $t \in \mathbf{R}$ and the output $y$ of $\Psi$ with control $u$ are given by

$$
\binom{x(t)}{y}=\binom{\mathcal{B} \tau(t) u}{\mathcal{D} u}
$$

and the controlled state $x^{*}(s)$ at time $s \in \mathbf{R}$ and the output $u^{*}$ of $\Psi^{*}$ with control $y^{*}$ are given by

$$
\binom{x^{*}(s)}{u^{*}}=\binom{\mathcal{C}^{*} \tau(s) y^{*}}{\mathcal{D}^{*} y^{*}}
$$

In the initial value setting with initial time zero, the controlled state $x(t)$ at time $t \in \mathbf{R}^{+}$and the output $y$ of $\Psi$ with initial value $x_{0}$ and control $u$ are given by

$$
\binom{x(t)}{y}=\left(\begin{array}{cc}
\mathcal{A}(t) & \mathcal{B} \tau(t) \\
\mathcal{C} & \mathcal{D}
\end{array}\right)\binom{x_{0}}{\pi_{+} u}=\binom{\mathcal{A}(t) x_{0}+\mathcal{B} \tau(t) \pi_{+} u}{\mathcal{C} x_{0}+\mathcal{D} \pi_{+} u}
$$

and the controlled state $x^{*}(s)$ at time $s \in \mathbf{R}^{-}$and the output $u^{*}$ of $\Psi^{*}$ with initial value $x_{0}^{*}$ and control $y^{*}$ are given by

$$
\binom{x^{*}(s)}{u^{*}}=\left(\begin{array}{cc}
\mathcal{A}^{*}(-s) & \mathcal{C}^{*} \tau(s) \\
\mathcal{B}^{*} & \mathcal{D}^{*}
\end{array}\right)\binom{x_{0}^{*}}{\pi_{-} y^{*}}=\binom{\mathcal{A}^{*}(-s) x_{0}^{*}+\mathcal{C}^{*} \tau(s) \pi_{-} y^{*}}{\mathcal{B}^{*} x_{0}^{*}+\mathcal{D}^{*} \pi_{-} y^{*}}
$$

Let us remark that the most common problem is the initial value problem for $\Psi$ with initial time zero, and in many papers this is the only one that is treated. Frequently the adjoint system $\Psi^{*}$ is also studied in an initial value setting with a positive initial time $S$; see, e.g., [34, Section 5]. This is in particular true in studies of the finite horizon quadratic cost minimization problem. See [39, Section 2] for the appropriate definition.


Figure 1. Input/State/Output Diagram of $\Psi$
In the case of the stable classical system (1) and its anti-stable adjoint (6), in the time-invariant setting the states of $\Psi$ and $\Psi^{*}$ are given by

$$
\begin{align*}
x(t) & =\int_{-\infty}^{t} \mathcal{A}(t-s) B u(s) d s, \quad t \in \mathbf{R}  \tag{11}\\
x^{*}(s) & =\int_{s}^{\infty} \mathcal{A}^{*}(t-s) C^{*} y^{*}(t) d t, \quad s \in \mathbf{R} \tag{12}
\end{align*}
$$

and the outputs are given by (5) and (10). In the initial value setting with initial time zero the states and outputs of the same systems are given by

$$
\begin{align*}
x(t) & =\mathcal{A}(t) x_{0}+\int_{0}^{t} \mathcal{A}(t-s) B u(s) d s, \quad t \in \mathbf{R}^{+}  \tag{13}\\
y(t) & =C \mathcal{A}(t) x_{0}+\int_{0}^{t} C \mathcal{A}(t-s) B u(s) d s+D u(t), \quad t \in \mathbf{R}^{+}  \tag{14}\\
x^{*}(s) & =\mathcal{A}^{*}(-s) x_{0}^{*}+\int_{s}^{0} \mathcal{A}^{*}(t-s) C^{*} y^{*}(t) d t, \quad s \in \mathbf{R}^{-}  \tag{15}\\
u^{*}(s) & =B^{*} \mathcal{A}^{*}(-s) x_{0}^{*}+\int_{s}^{0} B^{*} \mathcal{A}^{*}(t-s) C^{*} y^{*}(t) d t+D^{*} y^{*}(s), \quad s \in \mathbf{R}^{-} \tag{16}
\end{align*}
$$

Compare these formulas to (2) and (7).
An important property of the two different settings is that if we "restrict the causal time-invariant problem to $\mathbf{R}^{+}, "$ then we get an initial value problem, or more precisely, on $\mathbf{R}^{+}$the state and the output of the causal time-invariant problem are identical to the state and the output of the initial value problem with $x_{0}$ replaced by $x(0)=\mathcal{B} u$. A similar statement is true in the anti-causal case (replace $\mathbf{R}^{+}$by $\mathbf{R}^{-}$and $x_{0}^{*}$ by $\left.\mathcal{C}^{*} y^{*}\right)$.
Remark 10. Because of Definition 9, we also use the alternative notation $\left[\begin{array}{c}\mathcal{A} \\ \mathcal{C} \\ \mathcal{D} \boldsymbol{\mathcal { D }}\end{array}\right]$ for the system $\left[\begin{array}{cc}\mathcal{A} & \mathcal{B} \\ \mathcal{C} & \mathcal{D}\end{array}\right]$.

We use diagrams of the type drawn in Figure 1 to represent the relation between the state $x$, the output $y$, the initial value $x_{0}$, and the control $u$ of $\Psi$ in the initial value setting. In our diagrams we use the following conventions throughout:
(i) Initial states and controls enter at the top or bottom, and they are acted on by all the operators located in the column to which they are attached. In particular, note that $x_{0}$ is attached to the first column and $u$ to the second.
(ii) Final states and outputs leave to the left or right, and they are the sums of all the elements in the row to which they are attached. In particular, note that $x$ is attached to the top row, and $y$ to the bottom row.
A similar diagram is used to describe the adjoint system $\Psi^{*}$. In the time-invariant setting we use the same diagram without the initial vector $x_{0}$ and the projection $\pi_{+}$.

## 3. Quadratic Cost Minimization

This study centers around four problems that are closely connected to each other, namely quadratic cost minimization, spectral factorization, inner-outer factorization, and various types of feedback. The first of these problems is presented here, the next two are presented in Section 4, and the feedback configurations in Section 5.

We formulate the quadratic cost minimization problem for a causal well-posed linear system $\Psi$ in the initial value setting with initial time zero:
Definition 11. Let $\Psi=\left[\begin{array}{ll}\mathcal{A} & \mathcal{B} \\ \mathcal{C} & \mathcal{D}\end{array}\right]$ be an externally stable causal well-posed linear system on $(U, H, Y)$. The quadratic cost minimization problem for $\Psi$ consists of finding, for each $x_{0} \in H$, the infimum over all $u \in L^{2}\left(\mathbf{R}^{+} ; U\right)$ of the cost

$$
Q\left(x_{0}, u\right)=\|y\|_{L^{2}\left(\mathbf{R}^{+} ; Y\right)}^{2}
$$

where $y=\mathcal{C} x_{0}+\mathcal{D} \pi_{+} u$ is the output of $\Psi$ with initial value $x_{0}$ and control $u$.
As is well-known, the existence of a unique minimizing solution $u_{\text {opt }}$ of the quadratic cost minimization problem is guaranteed if the input/output map $\mathcal{D}$ of $\Psi$ is coercive:
Definition 12. A bounded linear operator $\mathcal{D}: L^{2}(\mathbf{R} ; U) \rightarrow L^{2}(\mathbf{R} ; Y)$ is coercive (or bounded from below) iff there exists some $\epsilon>0$ such that

$$
\|\mathcal{D} u\|_{L^{2}(\mathbf{R} ; Y)}^{2} \geq \epsilon\|u\|_{L^{2}(\mathbf{R} ; U)}^{2}
$$

for all $u \in L^{2}(\mathbf{R} ; U)$. Equivalently, $\mathcal{D}$ is coercive iff

$$
\mathcal{D}^{*} \mathcal{D} \geq \epsilon I
$$

for some $\epsilon>0$. Equivalently (by the closed graph theorem), $\mathcal{D}$ is coercive iff $\mathcal{D}$ is one-to-one and has a closed range.
Lemma 13. Let $\Psi=\left[\begin{array}{cc}\mathcal{A} & \mathcal{B} \\ \mathcal{C} & \mathcal{D}\end{array}\right]$ be an externally stable well-posed linear system on $(U, H, Y)$.
(i) $\mathcal{D}$ is coercive if and only if the Toeplitz operator $\pi_{+} \mathcal{D}^{*} \mathcal{D} \pi_{+}$, regarded as an operator mapping $L^{2}\left(\mathbf{R}^{+} ; U\right)$ into itself, has a bounded inverse. We denote this inverse by $\left(\pi_{+} \mathcal{D}^{*} \mathcal{D} \pi_{+}\right)^{-1}$.
(ii) Suppose that $\mathcal{D}$ is coercive. Then, for each $x_{0} \in H$, there is a unique control $u_{\mathrm{opt}} \in L^{2}\left(\mathbf{R}^{+} ; U\right)$ that minimizes the cost function $Q\left(x_{0}, u\right)$ in Definition 11, namely

$$
\begin{equation*}
u_{\mathrm{opt}}=-\left(\pi_{+} \mathcal{D}^{*} \mathcal{D} \pi_{+}\right)^{-1} \pi_{+} \mathcal{D}^{*} \mathcal{C} x_{0} \tag{17}
\end{equation*}
$$

The corresponding state $x_{\mathrm{opt}}$, the output $y_{\mathrm{opt}}$, and the minimal value of the cost function $Q$ are given by

$$
\begin{aligned}
x_{\mathrm{opt}} & =\mathcal{A} x_{0}-\mathcal{B} \tau \pi_{+}\left(\pi_{+} \mathcal{D}^{*} \mathcal{D} \pi_{+}\right)^{-1} \pi_{+} \mathcal{D}^{*} \mathcal{C} x_{0}, \\
y_{\mathrm{opt}} & =(I-P) \mathcal{C} x_{0}, \\
Q\left(x_{0}, u_{\mathrm{opt}}\right) & =\left\langle x_{0}, \mathcal{C}^{*}(I-P) \mathcal{C} x_{0}\right\rangle_{H},
\end{aligned}
$$

where $P=\mathcal{D} \pi_{+}\left(\pi_{+} \mathcal{D}^{*} \mathcal{D} \pi_{+}\right)^{-1} \pi_{+} \mathcal{D}^{*}$ is the orthogonal projection in $L^{2}(\mathbf{R} ; Y)$ onto the range of $\mathcal{D} \pi_{+}$. Moreover,

$$
\pi_{+} \mathcal{D}^{*} y_{\mathrm{opt}}=\pi_{+} \mathcal{D}^{*}\left(\mathcal{C} x_{0}+\mathcal{D} \pi_{+} u_{\mathrm{opt}}\right)=0
$$

that is, $y_{\mathrm{opt}}$ is orthogonal to the range of $\mathcal{D} \pi_{+}$.

Proof. (i) By definition, if $\mathcal{D}$ is coercive, then $\|\mathcal{D} u\|_{L^{2}(\mathbf{R} ; Y)}^{2} \geq \epsilon\|u\|_{L^{2}(\mathbf{R} ; U)}^{2}$ for some $\epsilon>0$ and all $u \in L^{2}(\mathbf{R} ; U)$. In particular, this implies that $\left\|\mathcal{D} \pi_{+} u\right\|_{L^{2}(\mathbf{R} ; Y)}^{2} \geq$ $\epsilon\left\|\pi_{+} u\right\|_{L^{2}(\mathbf{R} ; U)}^{2}$, and thus $\pi_{+} \mathcal{D}^{*} \mathcal{D} \pi_{+}$is positive and bounded away from zero on $L^{2}\left(\mathbf{R}^{+} ; U\right)$, hence invertible.

Conversely, if $\pi_{+} \mathcal{D}^{*} \mathcal{D} \pi_{+}$is invertible on $L^{2}\left(\mathbf{R}^{+} ; U\right)$, then it is strictly positive, and by the time-invariance of $\mathcal{D}$, there is some $\epsilon>0$ such that for for all $u \in$ $L^{2}(\mathbf{R} ; U)$,

$$
\begin{aligned}
\left\|\mathcal{D} \pi_{[t, \infty)} u\right\|_{L^{2}(\mathbf{R} ; Y)}^{2} & =\left\|\mathcal{D} \tau(-t) \pi_{+} \tau(t) u\right\|_{L^{2}(\mathbf{R} ; Y)}^{2} \\
& =\left\|\tau(-t) \mathcal{D} \pi_{+} \tau(t) u\right\|_{L^{2}(\mathbf{R} ; Y)}^{2} \\
& =\left\|\mathcal{D} \pi_{+} \tau(t) u\right\|_{L^{2}(\mathbf{R} ; Y)}^{2} \\
& \geq \epsilon\left\|\pi_{+} \tau(t) u\right\|_{L^{2}(\mathbf{R} ; U)}^{2} \\
& =\epsilon\|u\|_{L^{2}([t, \infty) ; U)}^{2}
\end{aligned}
$$

Let $t \rightarrow-\infty$ to conclude that $\mathcal{D}^{*} \mathcal{D} \geq \epsilon I$.
(ii) If $\mathcal{D}$ is coercive, then the cost function $Q\left(x_{0}, u\right)$ is convex and coercive with respect to $u \in L^{2}\left(\mathbf{R}^{+} ; U\right)$; hence there is a unique minimizing control $u_{\text {opt }} \in$ $L^{2}\left(\mathbf{R}^{+} ; U\right)$. To show that the corresponding output $y_{\text {opt }}$ satisfies (21) we argue as follows. Without loss of generality, let us suppose that $U$ is a real Hilbert space (if not, then we replace the inner product in $U$ by the real inner product $\Re\langle\cdot, \cdot\rangle$ ), and let us compute the Fréchet derivative of the cost function $Q\left(x_{0}, u\right)$ with respect to $u$ at the optimal $u_{\text {opt }}$. For each variation $\eta \in L^{2}\left(\mathbf{R}^{+} ; U\right)$, we have

$$
\begin{aligned}
d Q\left(x_{0}, u_{\mathrm{opt}}\right) \eta & =2\left\langle\mathcal{C} x_{0}+\mathcal{D} \pi_{+} u_{\mathrm{opt}}, \mathcal{D} \pi_{+} \eta\right\rangle_{L^{2}\left(\mathbf{R}^{+} ; U\right)} \\
& =2\left\langle y_{\mathrm{opt}}, \mathcal{D} \pi_{+} \eta\right\rangle_{L^{2}\left(\mathbf{R}^{+} ; U\right)} \\
& =2\left\langle\mathcal{D}^{*} y_{\mathrm{opt}}, \eta\right\rangle_{L^{2}\left(\mathbf{R}^{+} ; U\right)}
\end{aligned}
$$

This is zero for all $\eta \in L^{2}\left(\mathbf{R}^{+} ; U\right)$ iff (21) holds. Clearly, (17) follows from part (i) and (21). By substituting this value for $u_{\mathrm{opt}}$ into $x_{\mathrm{opt}}=\mathcal{A} x_{0}+\mathcal{B} \tau \pi_{+} u_{\mathrm{opt}}$, $y_{\mathrm{opt}}=\mathcal{C} x_{0}+\mathcal{D} \pi_{+} u_{\mathrm{opt}}$, and $Q\left(x_{0}, u_{\mathrm{opt}}\right)$ (and making a straightforward computation) we get (18), (19), and (20).

Definition 14. Under the hypotheses of part (ii) of Lemma 13, define

$$
\begin{aligned}
\mathcal{A}_{\circlearrowleft} & =\mathcal{A}-\mathcal{B} \tau \pi_{+}\left(\pi_{+} \mathcal{D}^{*} \mathcal{D} \pi_{+}\right)^{-1} \pi_{+} \mathcal{D}^{*} \mathcal{C} \\
\mathcal{C}_{\circlearrowleft} & =\left(I-\mathcal{D} \pi_{+}\left(\pi_{+} \mathcal{D}^{*} \mathcal{D} \pi_{+}\right)^{-1} \pi_{+} \mathcal{D}^{*}\right) \mathcal{C} \\
\mathcal{K}_{\circlearrowleft} & =-\left(\pi_{+} \mathcal{D}^{*} \mathcal{D} \pi_{+}\right)^{-1} \pi_{+} \mathcal{D}^{*} \mathcal{C} \\
\Pi & =\mathcal{C}^{*}\left(I-\mathcal{D} \pi_{+}\left(\pi_{+} \mathcal{D}^{*} \mathcal{D} \pi_{+}\right)^{-1} \pi_{+} \mathcal{D}^{*}\right) \mathcal{C}
\end{aligned}
$$

We call $\Pi$ the Riccati operator of $\Psi$.
Thus, according to Lemma $13, x_{\mathrm{opt}}=\mathcal{A}_{\circlearrowleft} x_{0}, y_{\mathrm{opt}}=\mathcal{C}_{\circlearrowleft} x_{0}, u_{\mathrm{opt}}=\mathcal{K}_{\circlearrowleft} x_{0}$, and $Q\left(x_{0}, u_{\mathrm{opt}}\right)=\left\langle x_{0}, \Pi x_{0}\right\rangle_{H}$.
Lemma 15. Make the same assumption as in part (ii) of Lemma 13, and introduce the same notation as in Definition 14. Then the following claims are true:
(i) The operators $\mathcal{A}_{\circlearrowleft}, \mathcal{C}_{\circlearrowleft}, \mathcal{K}_{\circlearrowleft}$, and $\Pi$ satisfy

$$
\begin{aligned}
\mathcal{A}_{\circlearrowleft} & =\mathcal{A}+\mathcal{B} \tau \mathcal{K}_{\circlearrowleft}, \\
\mathcal{C}_{\circlearrowleft} & =\mathcal{C}+\mathcal{D} \mathcal{K}_{\circlearrowleft}, \\
\Pi & =\mathcal{C}_{\circlearrowleft}^{*} \mathcal{C}_{\circlearrowleft}=\mathcal{C}^{*} \mathcal{C}_{\circlearrowleft}=\mathcal{C}_{\circlearrowleft}^{*} \mathcal{C} .
\end{aligned}
$$

(ii) $\mathcal{A}_{\circlearrowleft}$ is a strongly continuous semigroup on $H$, and $\mathcal{C}_{\circlearrowleft}$ and $\mathcal{K}_{\circlearrowleft}$ are admissible stable observability maps for $\mathcal{A}_{\circlearrowleft}$ in the sense that $\mathcal{C}_{\circlearrowleft} \in \mathcal{L}\left(H ; L^{2}(\mathbf{R} ; Y)\right)$, $\mathcal{K}_{\circlearrowleft} \in \mathcal{L}\left(H ; L^{2}(\mathbf{R} ; U)\right)$, and

$$
\begin{aligned}
\mathcal{C}_{\circlearrowleft} \mathcal{A}_{\circlearrowleft}(t) & =\pi_{+} \tau(t) \mathcal{C}_{\circlearrowleft}, \\
\mathcal{K}_{\circlearrowleft} \mathcal{A}_{\circlearrowleft}(t) & =\pi_{+} \tau(t) \mathcal{K}_{\circlearrowleft},
\end{aligned}
$$

for all $t \in \mathbf{R}^{+}$.
Proof. (i) This follows from Lemma 13 and Definition 14.
(ii) We fix some $t>0$, and consider controls $u$ of the type $u=u_{\mathrm{opt}}+\eta$, where $\eta \in L^{2}\left(\mathbf{R}^{+} ; U\right)$ satisfies $\pi_{[0, t]} \eta=0$ but is otherwise arbitrary. Then the state $x(t)$ at time $t$ and the output $y$ of $\Psi$ are given by

$$
\begin{aligned}
x(t) & =x_{\mathrm{opt}}(t), \\
y(s) & =y_{\mathrm{opt}}(s), \quad 0 \leq s \leq t \\
\pi_{+} \tau(t) y & =\pi_{+} \tau(t) y_{\mathrm{opt}}+\mathcal{D} \tau(t) \eta \\
& =\mathcal{C} x_{\mathrm{opt}}(t)+\mathcal{D} \pi_{+} \tau(t)\left(u_{\mathrm{opt}}+\eta\right)
\end{aligned}
$$

We write the cost $Q\left(x_{0}, u\right)$ in the form

$$
\begin{aligned}
Q\left(x_{0}, u\right) & =\int_{0}^{t}\left\langle y_{\mathrm{opt}}(s), y_{\mathrm{opt}}(s)\right\rangle_{Y} d s+\int_{t}^{\infty}\langle y(s), y(s)\rangle_{Y} d s \\
& =\int_{0}^{t}\left\langle y_{\mathrm{opt}}(s), y_{\mathrm{opt}}(s)\right\rangle_{Y} d s+\int_{0}^{\infty}\langle(\tau(t) y)(s),(\tau(t) y)(s)\rangle_{Y}
\end{aligned}
$$

Since $u_{\text {opt }}$ is the minimizing control, the derivative of $Q\left(x_{0}, u\right)$ with respect to $\eta$ must vanish at the point $\eta=0$, and this implies that (cf. the proof of Lemma 13)

$$
\pi_{+} \mathcal{D}^{*} \pi_{+} \tau(t) y_{\mathrm{opt}}=\pi_{+} \mathcal{D}^{*}\left(\mathcal{C} x_{\mathrm{opt}}(t)+\mathcal{D} \pi_{+} \tau(t) u_{\mathrm{opt}}\right)=0
$$

i.e., (21) holds with $x_{0}, y_{\mathrm{opt}}$, and $u_{\mathrm{opt}}$ replaced by $x_{\mathrm{opt}}(t), \pi_{+} \tau(t) y_{\mathrm{opt}}$, and $\pi_{+} \tau(t) u_{\mathrm{opt}}$, respectively. From this equation we can solve $\pi_{+} \tau(t) u_{\mathrm{opt}}$, and get (cf. Lemma 13 and Definition 14)

$$
\begin{aligned}
& \pi_{+} \tau(t) \mathcal{K}_{\circlearrowleft} x_{0}=\pi_{+} \tau(t) u_{\mathrm{opt}}=\mathcal{K}_{\circlearrowleft} x_{\mathrm{opt}}(t)=\mathcal{K}_{\circlearrowleft} \mathcal{A}_{\circlearrowleft}(t) x_{0}, \\
& \pi_{+} \tau(t) \mathcal{C}_{\circlearrowleft} x_{0}=\pi_{+} \tau(t) y_{\mathrm{opt}}=\mathcal{C}_{\circlearrowleft} x_{\mathrm{opt}}(t)=\mathcal{C}_{\circlearrowleft} \mathcal{A}_{\circlearrowleft}(t) x_{0}, \\
& \mathcal{A}_{\circlearrowleft}(s+t) x_{0}=x_{\mathrm{opt}}(s+t)=\mathcal{A}_{\circlearrowleft}(s) x_{\mathrm{opt}}(t)=\mathcal{A}_{\circlearrowleft}(s) \mathcal{A}_{\circlearrowleft}(t) x_{0} .
\end{aligned}
$$

Thus, $\mathcal{A}_{\circlearrowleft}$ is a semigroup, and $\mathcal{K}_{\circlearrowleft}$ and $\mathcal{C}_{\circlearrowleft}$ are admissible observability maps for $\mathcal{A}_{\circlearrowleft}$. The strong continuity of $\mathcal{A}_{\circlearrowleft}$ is immediate.

A stronger version of this lemma was proved independently by Zwart [52].

## 4. Spectral Factorization and Inner-Outer Factorization

As we shall prove below, the quadratic cost minimization problem is closely related to the spectral factorization problem, where we factor $\mathcal{D}^{*} \mathcal{D}$ into two invertible factors, out of which one is causal, and the other is anti-causal:

Definition 16. Let $\mathcal{D}: L^{2}(\mathbf{R} ; U) \rightarrow L^{2}(\mathbf{R} ; Y)$ be bounded, linear, and time-invariant. A bounded linear time-invariant operator $\mathcal{X}: L^{2}(\mathbf{R} ; U) \rightarrow L^{2}(\mathbf{R} ; U)$ is called a (canonical and invertible) spectral factor of $\mathcal{D}^{*} \mathcal{D}$ iff $\mathcal{X}$ is causal and has a bounded causal inverse (this inverse is necessarily time-invariant), and

$$
\mathcal{X}^{*} \mathcal{X}=\mathcal{D}^{*} \mathcal{D}
$$

As is well-known, the spectral factorization problem can be reduced to an innerouter factorization problem.
Definition 17. Let $\mathcal{D}: L^{2}(\mathbf{R} ; U) \rightarrow L^{2}(\mathbf{R} ; Y)$ be bounded, linear, causal, and time-invariant. An inner-outer factorization of $\mathcal{D}$ is a factorization of the form

$$
\mathcal{D}=\mathcal{Y} \mathcal{X}
$$

where both $\mathcal{Y}: L^{2}(\mathbf{R} ; U) \rightarrow L^{2}(\mathbf{R} ; Y)$ and $\mathcal{X}: L^{2}(\mathbf{R} ; U) \rightarrow L^{2}(\mathbf{R} ; U)$ are bounded linear causal time-invariant operators, $\mathcal{Y}$ is an isometry, i.e., $\mathcal{Y}^{*} \mathcal{Y}=I$, and the range of $\mathcal{X} \pi_{+}$is dense in $L^{2}\left(\mathbf{R}^{+} ; U\right)$. We call $\mathcal{Y}$ an inner factor and $\mathcal{X}$ an outer factor of $\mathcal{D}$.

Note that $\mathcal{X} \pi_{+}$is the Toeplitz operator induced by $\mathcal{X}$.
Lemma 18. Let $\mathcal{D}: L^{2}(\mathbf{R} ; U) \rightarrow L^{2}(\mathbf{R} ; Y)$ be bounded, linear, causal, and timeinvariant.
(i) If $\mathcal{X}$ is a spectral factor of $\mathcal{D}^{*} \mathcal{D}$, then $\mathcal{D}$ is coercive and $\mathcal{Y} \mathcal{X}=\left(\mathcal{D} \mathcal{X}^{-1}\right) \mathcal{X}$ is an inner-outer factorization of $\mathcal{D}$. Conversely, if $\mathcal{D}$ is coercive and $\mathcal{Y} \mathcal{X}$ is an inner-outer factorization of $\mathcal{D}$, then $\mathcal{X}$ is a spectral factor of $\mathcal{D}^{*} \mathcal{D}$.
(ii) $\mathcal{D}^{*} \mathcal{D}$ has a spectral factor $\mathcal{X}$ if and only if $\mathcal{D}$ is coercive.
(iii) The set of all spectral factors of $\mathcal{D}^{*} \mathcal{D}$ can be parametrized as JX , where $\mathcal{X}$ is a fixed spectral factor and $J \in \mathcal{L}(U)$ is an arbitrary unitary operator.
(iv) The inverse of the Toeplitz operator $\pi_{+} \mathcal{D}^{*} \mathcal{D} \pi_{+}$in part (i) of Lemma 13 can be written in the form $\left(\pi_{+} \mathcal{D}^{*} \mathcal{D} \pi_{+}\right)^{-1}=\mathcal{X}^{-1} \pi_{+}\left(\mathcal{X}^{*}\right)^{-1}$, where $\mathcal{X}$ is an arbitrary spectral factor of $\mathcal{D}^{*} \mathcal{D}$.
Inner-outer factorizations (without any coercivity assumptions) can be parametrized in the same way; see [32, Theorem B, p. 101].
Proof of Lemma 18. (i) The first claim is obvious. Conversely, suppose that $\mathcal{D}$ is coercive, and that $\mathcal{Y} \mathcal{X}$ is an inner-outer factorization of $\mathcal{D}$. Then $\mathcal{D}^{*} \mathcal{D}=\mathcal{X}^{*} \mathcal{Y}^{*} \mathcal{Y} \mathcal{X}=$ $\mathcal{X}^{*} \mathcal{X}$. Thus, in order to prove that $\mathcal{X}$ is a spectral factor of $\mathcal{D}^{*} \mathcal{D}$, it suffices to show that $\mathcal{X}$ has a bounded causal inverse.

Since $\mathcal{D}$ is coercive and $\mathcal{Y}$ is an isometry, also $\mathcal{X}$ is coercive; hence $\mathcal{X}$ is one-to-one and has a closed range. By the time-invariance of $\mathcal{X}$, we have $\mathcal{X} \tau(t) \pi_{+}=\tau(t) \mathcal{X} \pi_{+}$ for all $t \in \mathbf{R}^{+}$, and this combined with the fact that the range of $\mathcal{X} \pi_{+}$is dense in $L^{2}\left(\mathbf{R}^{+} ; U\right)$ implies that the range of $\mathcal{X}$ is dense in $L^{2}(\mathbf{R} ; U)$. Thus, $\mathcal{X}$ is a bijection and, by the closed graph theorem, it has a bounded inverse $\mathcal{X}^{-1}$. To prove that this inverse is causal, it suffices to show that the range of $\mathcal{X} \pi_{+}$is equal to $L^{2}\left(\mathbf{R}^{+} ; U\right)$. But this follows from the fact that the range of $\mathcal{X} \pi_{+}$is both dense in $L^{2}\left(\mathbf{R}^{+} ; U\right)$ and (by the coercivity of $\mathcal{X} \pi_{+}$) closed in $L^{2}\left(\mathbf{R}^{+} ; U\right)$.
(ii) If $\mathcal{D}^{*} \mathcal{D}$ has a spectral factor, then $\mathcal{D}^{*} \mathcal{D}$ is invertible, and $\mathcal{D}$ must be coercive. Conversely, if $\mathcal{D}$ is coercive, then it follows from [32, Theorem 3.4, p. 50 and Theorem 3.7, p. 54] that $\mathcal{D}^{*} \mathcal{D}$ has a factorization $\mathcal{X}^{*} \mathcal{X}$ where $\mathcal{X}$ is outer. By the argument given above $\mathcal{X}$ is invertible; hence $\mathcal{X}$ is a spectral factor of $\mathcal{D}^{*} \mathcal{D}$.
(iii) To prove the claim about the uniqueness of a spectral factor $\mathcal{X}$, suppose that both $\mathcal{X}$ and $\mathcal{Z}$ are spectral factors, i.e.,

$$
\mathcal{X}^{*} \mathcal{X}=\mathcal{D}^{*} \mathcal{D}=\mathcal{Z}^{*} \mathcal{Z}
$$

and that both $\mathcal{X}$ and $\mathcal{Z}$ are causal and have causal inverses. Then

$$
\mathcal{X} \mathcal{Z}^{-1}=\left(\mathcal{X}^{*}\right)^{-1} \mathcal{Z}^{*}
$$

The left-hand side is causal and the right-hand side is anti-causal, so these operators are static and, by Lemma 6,

$$
\mathcal{X} \mathcal{Z}^{-1}=\left(\mathcal{X}^{*}\right)^{-1} \mathcal{Z}^{*}=J
$$

for some operator $J \in \mathcal{L}(U)$. Thus, $\mathcal{X}=J \mathcal{Z}$ and $\mathcal{Z}^{*}=\mathcal{X}^{*} J$; hence $\mathcal{Z}=J^{*} \mathcal{X}=$ $J^{*} J \mathcal{Z}$. Since both $\mathcal{X}$ and $J$ are invertible, this implies that $J$ is unitary, i.e., $J^{*} J=J J^{*}=I$.
(iv) Use the causality of $\mathcal{X}$ and anti-causality of $\left(\mathcal{X}^{*}\right)^{-1}$ to get

$$
\begin{aligned}
\mathcal{X}^{-1} \pi_{+}\left(\mathcal{X}^{*}\right)^{-1} \pi_{+} \mathcal{D}^{*} \mathcal{D} \pi_{+} & =\mathcal{X}^{-1} \pi_{+}\left(\mathcal{X}^{*}\right)^{-1} \pi_{+} \mathcal{X}^{*} \mathcal{X} \pi_{+} \\
& =\mathcal{X}^{-1} \pi_{+}\left(\mathcal{X}^{*}\right)^{-1} \mathcal{X}^{*} \mathcal{X} \pi_{+} \\
& =\mathcal{X}^{-1} \pi_{+} \mathcal{X} \pi_{+} \\
& =\mathcal{X}^{-1} \mathcal{X} \pi_{+} \\
& =\pi_{+}
\end{aligned}
$$

Thus, $\mathcal{X}^{-1} \pi_{+}\left(\mathcal{X}^{*}\right)^{-1}$ is a left inverse of $\pi_{+} \mathcal{D}^{*} \mathcal{D} \pi_{+}$on $L^{2}\left(\mathbf{R}^{+} ; U\right)$. Since we know $\pi_{+} \mathcal{D}^{*} \mathcal{D} \pi_{+}$to be invertible, this operator must also be a right inverse of $\pi_{+} \mathcal{D}^{*} \mathcal{D} \pi_{+}$.

## 5. Static Output Feedback, State Feedback, and Output Injection

To connect the quadratic cost minimization problem to the spectral factorization problem (or alternatively, to the inner-outer factorization problem), we need the notion of a state feedback, which can be reduced to the notion of a (static) output feedback. The idea is to feed a part $z=L y$ of the output $y$ of a well-posed linear system $\Psi$ back into the input, as drawn in Figure 2. Here $L$ is a bounded linear operator from the output space into the input space. Then, in the initial value setting with initial value $x_{0}$ and input $v$, we find that the effective input $u$, the state $x(t)$ at time $t \geq 0$, the output $y$, and the feedback control signal $z$ satisfy the equations

$$
\begin{align*}
u & =z+\pi_{+} v \\
x(t) & =\mathcal{A}(t) x_{0}+\mathcal{B} \tau(t) u, \\
y & =\mathcal{C} x_{0}+\mathcal{D} u,  \tag{22}\\
z & =L y
\end{align*}
$$



Figure 2. Static Output Feedback
which formally can be solved as

$$
\begin{align*}
u & =(I-L \mathcal{D})^{-1}\left(L \mathcal{C} x_{0}+\pi_{+} v\right), \\
x(t) & =\left(\mathcal{A}(t)+\mathcal{B} \tau(t) L(I-\mathcal{D} L)^{-1} \mathcal{C}\right) x_{0}+\mathcal{B}(I-L \mathcal{D})^{-1} \tau(t) \pi_{+} v,  \tag{23}\\
y & =(I-\mathcal{D} L)^{-1}\left(\mathcal{C} x_{0}+\mathcal{D} \pi_{+} v\right) \\
z & =(I-L \mathcal{D})^{-1} L\left(\mathcal{C} x_{0}+\mathcal{D} \pi_{+} v\right) .
\end{align*}
$$

We say that the feedback operator $L$ is admissible whenever these equations are valid:

Definition 19. Let $\Psi=\left[\begin{array}{cc}\mathcal{A} & \mathcal{B} \\ \mathcal{C}\end{array}\right]$ be an externally stable causal [anti-causal] wellposed linear system on $(U, H, Y)$. The operator $L \in \mathcal{L}(Y ; U)$ is called an admissible stable output feedback operator for $\Psi$ iff the time-invariant operator $I$ $L \mathcal{D}: L^{2}(\mathbf{R} ; U) \rightarrow L^{2}(\mathbf{R} ; U)$ has a bounded causal [anti-causal] inverse, or equivalently, iff the time-invariant operator $I-\mathcal{D} L: L^{2}(\mathbf{R} ; Y) \rightarrow L^{2}(\mathbf{R} ; Y)$ has a bounded causal [anti-causal] inverse.

Observe that the inverses are necessarily time-invariant whenever they exist. That the two different invertibility conditions are equivalent is well-known, and so are the facts that

$$
\begin{aligned}
(I-\mathcal{D} L)^{-1} & =I+\mathcal{D}(I-L \mathcal{D})^{-1} L \\
(I-L \mathcal{D})^{-1} & =I+L(I-\mathcal{D} L)^{-1} \mathcal{D} \\
(I-\mathcal{D} L)^{-1} \mathcal{D} & =\mathcal{D}(I-L \mathcal{D})^{-1} \\
(I-L \mathcal{D})^{-1} L & =L(I-\mathcal{D} L)^{-1}
\end{aligned}
$$

As Weiss [46, Section 6] proved, $x$ and $y$ in (23) can be interpreted as the state and output of another well-posed linear system:

Proposition 20. Let $\Psi=\left[\begin{array}{ll}\mathcal{A} & \mathcal{B} \\ \mathcal{C} & \mathcal{D}\end{array}\right]$ be an externally stable causal well-posed linear system, and let $L \in \mathcal{L}(Y ; U)$ be an admissible stable output feedback operator for $\Psi$.

Then the system

$$
\begin{aligned}
\Psi_{L} & =\left[\begin{array}{cc}
\mathcal{A}_{L} & \mathcal{B}_{L} \tau \\
\mathcal{C}_{L} & \mathcal{D}_{L}
\end{array}\right]=\left[\begin{array}{cc}
\mathcal{A}+\mathcal{B} \tau L(I-\mathcal{D} L)^{-1} \mathcal{C} & \mathcal{B}(I-L \mathcal{D})^{-1} \tau \\
(I-\mathcal{D} L)^{-1} \mathcal{C} & \mathcal{D}(I-L \mathcal{D})^{-1}
\end{array}\right] \\
& =\left[\begin{array}{cc}
\mathcal{A} & \mathcal{B} \tau \\
\mathcal{C} & \mathcal{D}
\end{array}\right]+\left[\begin{array}{c}
\mathcal{B} \tau \\
\mathcal{D}
\end{array}\right] L(I-\mathcal{D} L)^{-1}\left[\begin{array}{ll}
\mathcal{C} & \mathcal{D}
\end{array}\right] \\
& =\left[\begin{array}{cc}
\mathcal{A} & \mathcal{B} \tau \\
\mathcal{C} & \mathcal{D}
\end{array}\right]+\left[\begin{array}{c}
\mathcal{B} \tau \\
\mathcal{D}
\end{array}\right] L\left[\begin{array}{ll}
\mathcal{C}_{L} & \mathcal{D}_{L}
\end{array}\right] \\
& =\left[\begin{array}{cc}
\mathcal{A} & \mathcal{B} \tau \\
\mathcal{C} & \mathcal{D}
\end{array}\right]+\left[\begin{array}{c}
\mathcal{B}_{L} \tau \\
\mathcal{D}_{L}
\end{array}\right] L\left[\begin{array}{ll}
\mathcal{C} & \mathcal{D}
\end{array}\right]
\end{aligned}
$$

is another externally stable causal well-posed linear system on $(U, H, Y)$. We call this system the closed loop system with output feedback operator L. In the initial value setting with initial time zero, initial value $x_{0}$, and control $v$, the controlled state $x(t)$ at time $t$ and the output $y$ of $\Psi_{L}$ are given by (23). (An analogous result is valid in the anti-causal case.)

To get a better feeling for the feedback formula in Proposition 20 we recommend that the reader carry out the following exercise. If in the classical system (1) we replace $u$ by $u=L y+v$, then we get a new well-defined system of the same type iff $I-D L$ is invertible, or equivalently, iff $I-L D$ is invertible. In the new system the operators $\left[\begin{array}{cc}A & B \\ C & D\end{array}\right]$ have been replaced by

$$
\begin{align*}
{\left[\begin{array}{ll}
A_{L} & B_{L} \\
C_{L} & D_{L}
\end{array}\right] } & =\left[\begin{array}{cc}
A+B L(I-D L)^{-1} C & B(I-L D)^{-1} \\
(I-D L)^{-1} C & D(I-L D)^{-1}
\end{array}\right] \\
& =\left[\begin{array}{ll}
A & B \\
C & D
\end{array}\right]+\left[\begin{array}{c}
B \\
D
\end{array}\right] L(I-D L)^{-1}\left[\begin{array}{ll}
C & D
\end{array}\right]  \tag{24}\\
& =\left[\begin{array}{ll}
A & B \\
C & D
\end{array}\right]+\left[\begin{array}{c}
B \\
D
\end{array}\right] L\left[\begin{array}{ll}
C_{L} & D_{L}
\end{array}\right] \\
& =\left[\begin{array}{ll}
A & B \\
C & D
\end{array}\right]+\left[\begin{array}{l}
B_{L} \\
D_{L}
\end{array}\right] L\left[\begin{array}{ll}
C & D
\end{array}\right]
\end{align*}
$$

This formula is probably not that familiar to most readers, since it is not common to allow a feed-through operator $D$ inside a feedback loop, and the formula simplifies significantly when $D=0$. In addition, the invertibility condition on $I-D L$ drops out when $D=0$. However, there is a striking similarity between this formula and the one given in Proposition 20. A general well-posed linear system need not have a well-defined feed-through operator $D$ (see Definition 35), so in the general case it is not possible to normalize $D$ to be zero. Similar comments apply to our subsequent feedback formulas.

As the following lemma shows, stability and strong stability are preserved under stable feedback:

Lemma 21. Let $\Psi=\left[\begin{array}{ll}\mathcal{A} & \mathcal{B} \\ \mathcal{C} & \mathcal{D}\end{array}\right]$ be an externally stable causal well-posed linear system, and let $L \in \mathcal{L}(Y ; U)$ be an admissible stable output feedback operator for $\Psi$. Then the closed loop system $\Psi_{L}$ is stable iff $\Psi$ is stable, i.e., $\sup _{t \in \mathbf{R}^{+}}\|\mathcal{A}(t)\|<\infty$ iff $\sup _{t \in \mathbf{R}^{+}}\left\|\mathcal{A}_{L}(t)\right\|<\infty$. Moreover $\Psi_{L}$ is strongly stable iff $\Psi$ is strongly stable, i.e., $\mathcal{A}(t) x \rightarrow 0$ as $t \rightarrow \infty$ for all $x \in H$ iff $\mathcal{A}_{L}(t) x \rightarrow 0$ as $t \rightarrow \infty$ for all $x \in H$.


Figure 3. State Feedback Connection

Proof. The first claim follows from the fact that the difference $\mathcal{A}_{L}-\mathcal{A}=$ $\mathcal{B} \tau L(I-\mathcal{D} L)^{-1} \mathcal{C}$ is bounded.

To prove the second claim it suffices to show that

$$
\mathcal{B} \tau(t) L(I-\mathcal{D} L)^{-1} \mathcal{C} x \rightarrow 0 \text { as } t \rightarrow \infty
$$

for every $x \in H$. Fix $x \in H$, and split the expression above into

$$
\begin{aligned}
\mathcal{B} \tau(t) L(I-\mathcal{D} L)^{-1} \mathcal{C} x= & \mathcal{B} \tau(t-T)\left(\pi_{+}+\pi_{-}\right) \tau(T) L(I-\mathcal{D} L)^{-1} \mathcal{C} x \\
= & \mathcal{B} \tau(t) \pi_{[T, \infty)} L(I-\mathcal{D} L)^{-1} \mathcal{C} x \\
& +\mathcal{A}(t-T) \mathcal{B} \tau(T) L(I-\mathcal{D} L)^{-1} \mathcal{C} x
\end{aligned}
$$

Here the first term tends to zero as $T \rightarrow \infty$, uniformly in $t \geq T$, and the second term tends to zero as $t \rightarrow \infty$ and $T$ is fixed.

As we mentioned above, a state feedback can be reduced to an output feedback as follows. The appropriate connection has been drawn in Figure 3.

Definition 22. Let $\Psi=\left[\begin{array}{cc}\mathcal{A} & \mathcal{B} \\ \mathcal{C}\end{array}\right]$ be an externally stable causal [anti-causal] wellposed linear system on $(U, H, Y)$. The pair $(\mathcal{K} \mathcal{F})$ is called an admissible stable state feedback pair for $\Psi$ iff the extended system

$$
\Psi_{\text {ext }}=\left[\begin{array}{cc}
\mathcal{A} & \mathcal{B} \\
\binom{\mathcal{C}}{\mathcal{K}} & \binom{\mathcal{D}}{\mathcal{F}}
\end{array}\right]
$$

is an externally stable causal [anti-causal] well-posed linear system, and ( $0 \quad I \quad I)$ is an admissible stable output feedback operator for $\Psi_{\text {ext }}$, i.e., $I-\mathcal{F}$ has a bounded causal [anti-causal] inverse.

In the state feedback case with initial value $x_{0}$ and input $v$, we obtain the effective input $u$, the state $x(t)$, the output $y$, and the feedback control signal $z$ of $\Psi_{\text {ext }}$ from (22) and (23) by replacing $\mathcal{C}, \mathcal{D}, L$, and $y$ by $\binom{\mathcal{C}}{\mathcal{K}},\binom{\mathcal{D}}{\mathcal{F}},\left(\begin{array}{ll}0 & I\end{array}\right)$, and $\binom{y}{z}$, respectively. This leads to the formulas

$$
\begin{align*}
u & =z+\pi_{+} v, \\
x(t) & =\mathcal{A}(t) x_{0}+\mathcal{B} \tau(t) u, \\
y & =\mathcal{C} x_{0}+\mathcal{D} u,  \tag{25}\\
z & =\mathcal{K} x_{0}+\mathcal{F} u,
\end{align*}
$$

and

$$
\begin{align*}
u & =(I-\mathcal{F})^{-1}\left(\mathcal{K} x_{0}+\pi_{+} v\right) \\
x(t) & =\left(\mathcal{A}(t)+\mathcal{B} \tau(t)(I-\mathcal{F})^{-1} \mathcal{K}\right) x_{0}+\mathcal{B}(I-\mathcal{F})^{-1} \tau(t) \pi_{+} v, \\
y & =\left(\mathcal{C}+\mathcal{D}(I-\mathcal{F})^{-1} \mathcal{K}\right) x_{0}+\mathcal{D}(I-\mathcal{F})^{-1} \pi_{+} v,  \tag{26}\\
z & =(I-\mathcal{F})^{-1}\left(\mathcal{K} x_{0}+\mathcal{F} \pi_{+} v\right)
\end{align*}
$$

We get the closed loop state feedback system $\Psi_{\circlearrowleft}$ by making the same replacements in Proposition 20:

Proposition 23. Let $\Psi=\left[\begin{array}{ll}\mathcal{A} & \mathcal{B} \\ \mathcal{C} & \mathcal{D}\end{array}\right]$ be an externally stable causal well-posed linear system, and let $\left(\begin{array}{ll}\mathcal{K} & \mathcal{F}\end{array}\right)$ be an admissible stable state feedback pair for $\Psi$. Then the system

$$
\begin{aligned}
\Psi_{\circlearrowleft} & =\left[\begin{array}{cc}
\mathcal{A}_{\circlearrowleft} & \mathcal{B}_{\circlearrowleft} \tau \\
\binom{\mathcal{C}_{\circlearrowleft}}{\mathcal{K}_{\circlearrowleft}} & \binom{\mathcal{D}_{\circlearrowleft}}{\mathcal{F}_{\circlearrowleft}}
\end{array}\right] \\
= & {\left[\begin{array}{cc}
\mathcal{A}+\mathcal{B} \tau(I-\mathcal{F})^{-1} \mathcal{K} & \mathcal{B}(I-\mathcal{F})^{-1} \tau \\
\left(\begin{array}{c}
\mathcal{C}+\mathcal{D}(I-\mathcal{F})^{-1} \\
\mathcal{K} \\
(I-\mathcal{F})^{-1} \mathcal{K}
\end{array}\right) & \binom{\mathcal{D}(I-\mathcal{F})^{-1}}{(I-\mathcal{F})^{-1}-I}
\end{array}\right] } \\
& =\left[\begin{array}{cc}
\mathcal{A} & \mathcal{B} \tau \\
\binom{\mathcal{C}}{\mathcal{K}} & \binom{\mathcal{D}}{\mathcal{F}}
\end{array}\right]+\left[\begin{array}{c}
\mathcal{B} \tau \\
\binom{\mathcal{D}}{\mathcal{F}}
\end{array}\right](I-\mathcal{F})^{-1}\left[\begin{array}{ll}
\mathcal{K} & \mathcal{F}]
\end{array}\right.
\end{aligned}
$$

is another externally stable well-posed linear system on $(U, H, Y \times U)$. We call this system the closed loop state feedback system with state feedback pair $(\mathcal{K} \quad \mathcal{F})$. In the initial value setting with initial time zero, initial value $x_{0}$, and control $v$, the controlled state $x(t)$ at time $t$ and the outputs $y$ and $z$ of $\Psi_{\circlearrowleft}$ are given by (26). (An analogous result is valid in the anti-causal case.)

Remark 24. We shall frequently regard the signal $u$ in Figure 3 as an additional output of the closed loop system. This output has the same observability map $\mathcal{K}_{\circlearrowleft}=(I-\mathcal{F})^{-1} \mathcal{K}$ as the output $z$, and its input/output map is $\mathcal{F}_{\circlearrowleft}+I=(I-\mathcal{F})^{-1}$. A similar remark applies to the signal $z^{*}$ in the adjoint output injection connection in Figure 4.

The adjoint of the state feedback system $\Psi_{\circlearrowleft}$ is the output injection system drawn in Figure 4.

Definition 25. Let $\Psi^{*}=\left[\begin{array}{cc}\mathcal{A}^{*} & \mathcal{C}^{*} \\ \mathcal{B}^{*} & \mathcal{D}^{*}\end{array}\right]$ be an externally stable anti-causal [causal] wellposed linear system on $(Y, H, U)$. Then the pair $\left(\begin{array}{c}\mathcal{F}^{*}\end{array}\right)$ is called an admissible stable output injection pair for $\Psi^{*}$ iff the extended system

$$
\Psi_{\text {ext }}^{*}=\left[\begin{array}{ccc}
\mathcal{A}^{*} & \left(\mathcal{C}^{*}\right. & \left.\mathcal{K}^{*}\right) \\
\mathcal{B}^{*} & \left(\mathcal{D}^{*}\right. & \left.\mathcal{F}^{*}\right)
\end{array}\right]
$$

is an externally stable anti-causal [causal] well-posed linear system, and ( $\left.\begin{array}{l}0 \\ I\end{array}\right)$ is an admissible stable output feedback operator for $\Psi_{\text {ext }}^{*}$, i.e., $I-\mathcal{F}^{*}$ has a bounded anti-causal [causal] inverse.


Figure 4. Adjoint Output Injection Connection

The appropriate formulas for the output injection connection are in the initial value setting

$$
\begin{align*}
z^{*} & =u^{*}+\pi_{-} v^{*}, \\
x^{*}(s) & =\mathcal{A}^{*}(-s) x_{0}^{*}+\mathcal{C}^{*} \tau(s) \pi_{-} y^{*}+\mathcal{K}^{*} \tau(s) z^{*},  \tag{27}\\
u^{*} & =\mathcal{B}^{*} x_{0}^{*}+\mathcal{D}^{*} \pi_{-} y^{*}+\mathcal{F}^{*} z^{*},
\end{align*}
$$

and

$$
\begin{align*}
z^{*}= & \left(I-\mathcal{F}^{*}\right)^{-1}\left(\mathcal{B}^{*} x_{0}^{*}+\mathcal{D}^{*} \pi_{-} y^{*}+\pi_{-} v^{*}\right) \\
x^{*}(s)= & \left(\mathcal{A}^{*}(-s)+\mathcal{K}^{*} \tau(s)\left(I-\mathcal{F}^{*}\right)^{-1} \mathcal{B}^{*}\right) x_{0}^{*}  \tag{28}\\
& +\left(\mathcal{C}^{*}+\mathcal{K}^{*}\left(I-\mathcal{F}^{*}\right)^{-1} \mathcal{D}^{*}\right) \tau(s) \pi_{-} y^{*}+\mathcal{K}^{*}\left(I-\mathcal{F}^{*}\right)^{-1} \tau(s) \pi_{-} v^{*} \\
u^{*}= & \left(I-\mathcal{F}^{*}\right)^{-1}\left(\mathcal{B}^{*} x_{0}^{*}+\mathcal{D}^{*} \pi_{-} y^{*}+\mathcal{F}^{*} \pi_{-} v^{*}\right)
\end{align*}
$$

We shall actually use this system in the time-invariant setting, for which slightly simpler formulas are valid ( $\operatorname{set} x_{0}^{*}=0$ and delete $\pi_{-}$).

We get the closed loop output injection system $\Psi_{\circlearrowleft}^{*}$ by making the appropriate replacements in Proposition 20:

Proposition 26. Let $\Psi^{*}=\left[\begin{array}{cc}\mathcal{A}^{*} & \mathcal{C}^{*} \\ \mathcal{B}^{*} & \mathcal{D}^{*}\end{array}\right]$ be an externally stable anti-causal well-posed linear system, and let $\binom{\mathcal{K}^{*}}{\mathcal{F}^{*}}$ be an admissible stable output injection pair for $\Psi$. Then the system $\Psi_{\circlearrowleft}^{*}$ given by

$$
\begin{aligned}
& \Psi_{\circlearrowleft}^{*}=\left[\begin{array}{ccc}
\mathcal{A}_{\circlearrowleft}^{*} & \left(\mathcal{C}_{\circlearrowleft}^{*} \tau\right. & \left.\mathcal{K}_{\circlearrowleft}^{*} \tau\right) \\
\mathcal{B}_{\circlearrowleft}^{*} & \left(\mathcal{D}_{\circlearrowleft}^{*}\right. & \left.\mathcal{F}_{\circlearrowleft}^{*}\right)
\end{array}\right] \\
&=\left[\begin{array}{ccc}
\mathcal{A}^{*}+\mathcal{K}^{*} \tau^{*}\left(I-\mathcal{F}^{*}\right)^{-1} \mathcal{B}^{*} & \left(\mathcal{C}^{*} \tau+\mathcal{K}^{*}\left(I-\mathcal{F}^{*}\right)^{-1} \mathcal{D}^{*} \tau\right. & \left.\mathcal{K}^{*}\left(I-\mathcal{F}^{*}\right)^{-1} \tau\right) \\
\left(I-\mathcal{F}^{*}\right)^{-1} \mathcal{B}^{*} & \left(\left(I-\mathcal{F}^{*}\right)^{-1} \mathcal{D}^{*}\right. & \left.\left(I-\mathcal{F}^{*}\right)^{-1}-I\right)
\end{array}\right]
\end{aligned}
$$

is an externally stable anti-causal well-posed linear system. We call this system the closed loop output injection system with output injection pair $\left(\mathcal{K}_{\mathcal{F}^{*}}{ }^{*}\right)$. The controlled state $x^{*}(s)$ at time $s \leq 0$ and output $u^{*}$ of $\Psi_{\circlearrowleft}^{*}$ in the initial value setting with initial time zero, initial state $x_{0}^{*}$, and controls $y^{*}$ and $v^{*}$ are given by (28). (An analogous result is valid in the causal case.)

## 6. The Connection Between Quadratic Cost Minimization and Spectral Factorization

The following theorem is our first main result. It shows that the spectral factorization problem is equivalent to the problem of finding a state feedback solution to the quadratic cost minimization problem:

Theorem 27. Let $\Psi=\left[\begin{array}{cc}\mathcal{A} & \mathcal{B} \\ \mathcal{C} & \mathcal{D}\end{array}\right]$ be an externally stable causal well-posed linear system on $(U, H, Y)$, and let $\mathcal{D}$ be coercive (in the sense of Definition 12). For each $x_{0} \in H$, let $x_{\mathrm{opt}}, u_{\mathrm{opt}}$, and $y_{\mathrm{opt}}$ be the optimal state, control, and output for the quadratic cost minimization problem (see Lemma 13). Let $\Pi$ be the Riccati operator defined in Definition 14.
(i) If $\mathcal{Y X}$ is an inner-outer factorization of $\mathcal{D}$ (in the sense of Definition 17), and if $E$ is an arbitrary invertible operator in $\mathcal{L}(U)$, then

$$
\left(\begin{array}{ll}
\mathcal{K} & \mathcal{F}
\end{array}\right)=\left(-\pi_{+} E^{-1} \mathcal{Y}^{*} \mathcal{C} \quad\left(I-E^{-1} \mathcal{X}\right)\right)
$$

is an admissible stable state feedback pair for $\Psi$, and

$$
\begin{aligned}
\left(\begin{array}{c}
x_{\mathrm{opt}} \\
y_{\mathrm{opt}} \\
u_{\mathrm{opt}}
\end{array}\right) & =\left(\begin{array}{c}
\mathcal{A}_{\circlearrowleft} \\
\mathcal{C}_{\circlearrowleft} \\
\mathcal{K}_{\circlearrowleft}
\end{array}\right) x_{0}=\left(\begin{array}{c}
\mathcal{A}+\mathcal{B} \mathcal{X}^{-1} \tau E \mathcal{K} \\
\mathcal{C}+\mathcal{Y} E \mathcal{K} \\
\mathcal{X}^{-1} E \mathcal{K}
\end{array}\right) x_{0} \\
& =\left(\begin{array}{c}
\mathcal{A} \\
\mathcal{C} \\
0
\end{array}\right) x_{0}-\left(\begin{array}{c}
\mathcal{B} \mathcal{X}^{-1} \tau \\
\mathcal{Y} \\
\mathcal{X}^{-1}
\end{array}\right) \pi_{+} \mathcal{Y}^{*} \mathcal{C} x_{0}
\end{aligned}
$$

is equal to the state and output of the closed loop system

$$
\begin{align*}
\Psi_{\circlearrowleft} & =\left[\begin{array}{cc}
\mathcal{A}_{\circlearrowleft} & \mathcal{B}_{\circlearrowleft} \\
\binom{\mathcal{C}_{\circlearrowleft}}{\mathcal{K}_{\circlearrowleft}} & \binom{\mathcal{D}_{\circlearrowleft}}{\mathcal{F}_{\circlearrowleft}}
\end{array}\right]=\left[\begin{array}{cc}
\mathcal{A}+\mathcal{B} \mathcal{X}^{-1} \tau E \mathcal{K} & \mathcal{B} \mathcal{X}^{-1} E \\
\binom{\mathcal{C}+\mathcal{Y} E \mathcal{K}}{\mathcal{X}^{-1} E \mathcal{K}} & \binom{\mathcal{Y} E}{\mathcal{X}^{-1} E-I}
\end{array}\right] \\
& =\left[\begin{array}{c}
\mathcal{A}-\mathcal{B} \mathcal{X}^{-1} \tau \pi_{+} \mathcal{Y}^{*} \mathcal{C} \\
\left(\begin{array}{c}
\mathcal{X} \mathcal{X}^{-1} E \\
\left(I-\mathcal{Y}_{+} \mathcal{Y}^{*}\right) \mathcal{C} \\
-\mathcal{X}^{-1} \pi_{+} \mathcal{Y}^{*} \mathcal{C}
\end{array}\right)
\end{array} \begin{array}{c}
\mathcal{Y} E \\
\mathcal{X}^{-1} E-I
\end{array}\right) \tag{29}
\end{align*}
$$

with this feedback pair, initial value $x_{0}$, initial time zero, and zero control $u$. The Riccati operator $\Pi$ of $\Psi$ can be written in the following alternative forms:

$$
\Pi=\mathcal{C}^{*} \mathcal{C}-\mathcal{K}^{*} E^{*} E \mathcal{K}=\mathcal{C}^{*}\left(I-\mathcal{Y} \pi_{+} \mathcal{Y}^{*}\right) \mathcal{C}=\mathcal{C}^{*} \mathcal{C}_{\circlearrowleft}=\mathcal{C}_{\circlearrowleft}^{*} \mathcal{C}_{\circlearrowleft}
$$

In particular, $\Pi$ is the observability gramian of the closed loop system $\Psi_{\circlearrowleft}$.
(ii) Conversely, if $\binom{y_{\text {opt }}}{u_{\text {opt }}}$ is equal to the output of some externally stable closed loop state feedback extension $\Psi_{\circlearrowleft}$ of $\Psi$ with initial value $x_{0}$, initial time 0 , zero control $u$, and some admissible stable state feedback pair $(\mathcal{K} \mathcal{F})$, then there exists an invertible operator $E \in \mathcal{L}(U)$ such that $\mathcal{Y} \mathcal{X}$ is an inner-outer factorization of $\mathcal{D}$, where $\mathcal{Y}=\mathcal{D}(I-\mathcal{F})^{-1} E^{-1}$ and $\mathcal{X}=E(I-\mathcal{F})$. Moreover, $\mathcal{K}$ is given by $\mathcal{K}=-\pi_{+} E^{-1} \mathcal{Y}^{*} \mathcal{C}$. Thus, every state feedback solution of the quadratic cost minimization problem is of the type described in part (i).

Remark 28. The operator $\mathcal{Y} \pi_{+} \mathcal{Y}^{*}$ in the formulas above is equal to the projection operator $P$ in Lemma 13. The operator $I-\mathcal{Y} \pi_{+} \mathcal{Y}^{*}$ can be interpreted as the Riccati operator of a particular realization of $\mathcal{D}$; see [37].

The explanation for the appearance of the extra undetermined operator $E$ in Theorem 27 is the following. In the case of the classical system (1) a standard state
feedback law has the form $u=L x$ for some bounded linear operator $L$. The righthand side of this feedback equation can be thought of as an extra output of the original system with the property that the feed-through term has been normalized to be zero. As we mentioned in connection with the formula for the output feedback connection, it is not in general possible to normalize the feed-through term of a well-posed linear system to be zero; to do this we need the extra regularity condition introduced in Definition 35. If we do not normalize the feed-through term in the feedback loop to be zero, the we must write the feedback equation in the form $u=K x+F u$ for some bounded linear operators $K$ and $F$. This equation is equivalent to the earlier equation $u=L x$ iff $I-F$ is invertible, and $L=(I-F)^{-1} K$. Define $E=(I-F)^{-1}$. Then the feedback equation $u=K x+F u$ becomes $u=E^{-1} L x+\left(I-E^{-1}\right) u$, which begins to resemble the actual formula used in Theorem 27. Notice, in particular, that the "effective" feedback is $L=E K$ instead of $K$.

For later reference, let us write out the expression for the extended open loop system $\Psi_{\text {ext }}$ in Definition 22 with the particular state feedback pair $(\mathcal{K} \mathcal{F})=$ $\left(-\pi_{+} E^{-1} \mathcal{Y}^{*} \mathcal{C} \quad\left(I-E^{-1} \mathcal{X}\right)\right)$ given by Theorem 27 , and also the expressions for the adjoints of $\Psi_{\text {ext }}$ and $\Psi_{\circlearrowleft}$ :

$$
\begin{align*}
& \Psi_{\text {ext }}=\left[\begin{array}{cc}
\mathcal{A} & \mathcal{B} \\
\binom{\mathcal{C}}{\mathcal{K}} & \binom{\mathcal{D}}{\mathcal{F}}
\end{array}\right]=\left[\begin{array}{cc}
\mathcal{A} & \mathcal{B} \\
\binom{\mathcal{C}}{-\pi_{+} E^{-1} \mathcal{Y}^{*} \mathcal{C}} & \binom{\mathcal{D}}{I-E^{-1} \mathcal{X}}
\end{array}\right],  \tag{30}\\
& \Psi_{\text {ext }}^{*}=\left[\begin{array}{lll}
\mathcal{A}^{*} & \left(\mathcal{C}^{*}\right. & \left.\mathcal{K}^{*}\right) \\
\mathcal{B}^{*} & \left(\mathcal{D}^{*}\right. & \left.\mathcal{F}^{*}\right)
\end{array}\right]=\left[\begin{array}{lll}
\mathcal{A}^{*} & \left(\mathcal{C}^{*}\right. & \left.-\mathcal{C}^{*} \mathcal{Y}\left(E^{*}\right)^{-1} \pi_{+}\right) \\
\mathcal{B}^{*} & \left(\mathcal{D}^{*}\right. & \left.I-\mathcal{X}^{*}\left(E^{*}\right)^{-1}\right)
\end{array}\right],  \tag{31}\\
& \Psi_{\circlearrowleft}^{*}=\left[\begin{array}{ccc}
\mathcal{A}_{\circlearrowleft}^{*} & \left(\mathcal{C}_{\circlearrowleft}^{*}\right. & \left.\mathcal{K}_{\circlearrowleft}^{*}\right) \\
\mathcal{B}_{\circlearrowleft}^{*} & \left(\mathcal{D}_{\circlearrowleft}^{*}\right. & \left.\mathcal{F}_{\circlearrowleft}^{*}\right)
\end{array}\right]  \tag{32}\\
& =\left[\begin{array}{ccc}
\mathcal{A}^{*}-\mathcal{C}^{*} \mathcal{Y} \pi_{+} \tau^{*}\left(\mathcal{X}^{*}\right)^{-1} \mathcal{B}^{*} & \left(\mathcal{C}^{*}\left(I-\mathcal{Y} \pi_{+} \mathcal{Y}^{*}\right)\right. & \left.-\mathcal{C}^{*} \mathcal{Y} \pi_{+}\left(\mathcal{X}^{*}\right)^{-1}\right) \\
E^{*}\left(\mathcal{X}^{*}\right)^{-1} \mathcal{B}^{*} & \left(E^{*} \mathcal{Y}^{*}\right. & \left.E^{*}\left(\mathcal{X}^{*}\right)^{-1}-I\right)
\end{array}\right] .
\end{align*}
$$

Proof of Theorem 27. (i) We suppose that $\mathcal{Y X}$ is an inner-outer factorization of $\mathcal{D}$. Let $E \in \mathcal{L}(H)$ be invertible, and define

$$
\left(\begin{array}{ll}
\mathcal{K} & \mathcal{F}
\end{array}\right)=\left(-\pi_{+} E^{-1} \mathcal{Y}^{*} \mathcal{C} \quad\left(I-E^{-1} \mathcal{X}\right)\right)
$$

We claim that the system $\Psi_{\text {ext }}$ defined in (30) is an externally stable well-posed linear system on $(U, H, Y \times U)$. It is obvious that all the operators in this system are causal, bounded linear operators on their appropriate domains. The only nonobvious parts of the claim that $\Psi_{\text {ext }}$ is well-posed are that

$$
\mathcal{K} \mathcal{A}(t)=\pi_{+} \tau(t) \mathcal{K}, \quad t \in \mathbf{R}^{+}
$$

and that the Hankel operator of $\mathcal{F}$ is equal to $\mathcal{K} \mathcal{B}$. To compute $\mathcal{K} \mathcal{A}(t)$, we recall that $E^{-1} \mathcal{Y}^{*}$ is time-invariant and anti-causal; hence

$$
\begin{aligned}
\mathcal{K} \mathcal{A}(t) & =-\pi_{+} E^{-1} \mathcal{Y}^{*} \mathcal{C} \mathcal{A}(t) \\
& =-\pi_{+} E^{-1} \mathcal{Y}^{*} \pi_{+} \tau(t) \mathcal{C} \\
& =-\pi_{+} E^{-1} \mathcal{Y}^{*} \tau(t) \mathcal{C} \\
& =-\pi_{+} \tau(t) E^{-1} \mathcal{Y}^{*} \mathcal{C} \\
& =-\pi_{+} \tau(t) \pi_{+} E^{-1} \mathcal{Y}^{*} \mathcal{C} \\
& =\pi_{+} \tau(t) \mathcal{K}
\end{aligned}
$$

In the computation of the Hankel operator of $\mathcal{F}$, we use the facts that $\mathcal{F}$ and $-E^{-1} \mathcal{X}$ have the same Hankel operator, that $\mathcal{X}=\mathcal{Y}^{*} \mathcal{Y} \mathcal{X}=\mathcal{Y}^{*} \mathcal{D}$, that $\pi_{+}+\pi_{-}=I$, that $E^{-1} \mathcal{Y}^{*}$ is anti-causal, and that the Hankel operator of $\mathcal{D}$ is $\mathcal{C B}$ :

$$
\begin{aligned}
\pi_{+} \mathcal{F} \pi_{-} & =-\pi_{+} E^{-1} \mathcal{X} \pi_{-} \\
& =-\pi_{+} E^{-1} \mathcal{Y}^{*} \mathcal{D} \pi_{-} \\
& =-\pi_{+} E^{-1} \mathcal{Y}^{*} \pi_{+} \mathcal{D} \pi_{-} \\
& =-\pi_{+} E^{-1} \mathcal{Y}^{*} \mathcal{C B} \\
& =\mathcal{K B}
\end{aligned}
$$

To complete the proof of (i) it suffices to insert the formula for $\left(\pi_{+} \mathcal{D}^{*} \mathcal{D} \pi_{+}\right)^{-1}$ given in part (iv) of Lemma 18 into the formulas of Lemma 13, and to compare the result to the output of the closed loop system $\Psi_{\circlearrowleft}$ with initial time zero, initial state $x_{0}$, and zero control $u$.
(ii) Conversely, suppose that $\binom{y_{\text {opt }}}{u_{\text {opt }}}$ is equal to the output of some externally stable closed loop state feedback extension $\Psi_{\circlearrowleft}$ of $\Psi$ with initial value $x_{0}$, initial time 0 , zero control $u$, and some admissible stable state feedback pair $(\mathcal{K} \mathcal{F})$. According to Proposition 23,

$$
\begin{align*}
& y_{\mathrm{opt}}=\mathcal{C}_{\circlearrowleft} x_{0}=\left(\mathcal{C}+\mathcal{D}(I-\mathcal{F})^{-1} \mathcal{K}\right) x_{0}  \tag{33}\\
& u_{\mathrm{opt}}=\mathcal{K}_{\circlearrowleft} x_{0}=(I-\mathcal{F})^{-1} \mathcal{K} x_{0}
\end{align*}
$$

The first step of this proof is to show that it is possible to write $\mathcal{D}^{*} \mathcal{D}$ in the form $\mathcal{Z}^{*} \mathcal{W}$, where both $\mathcal{Z}$ and $\mathcal{W}$ are causal and $\mathcal{W}$ has a bounded causal inverse. Motivated by the formula $\mathcal{F}=I-E^{-1} \mathcal{X}$ in part (i), we take $E=I$, and choose $\mathcal{W}$ and $\mathcal{Z}^{*}$ to be

$$
\mathcal{W}=I-\mathcal{F}, \quad \mathcal{Z}^{*}=\mathcal{D}^{*} \mathcal{D} \mathcal{W}^{-1}=\mathcal{D}^{*} \mathcal{D}(I-\mathcal{F})^{-1}
$$

Since $\left(\begin{array}{ll}\mathcal{K} & \mathcal{F}\end{array}\right)$ is an admissible stable state feedback pair for $\Psi, \mathcal{W}=I-\mathcal{F}$ is causal and has a bounded causal inverse (which makes it possible to define $\mathcal{Z}^{*}$ in the way that we did above).

We claim that $\mathcal{Z}^{*}$ is anti-causal. To prove this we choose an arbitrary $u \in$ $L^{2}(\mathbf{R} ; U)$, and choose $x_{0}$ in the quadratic cost minimization problem to be $x_{0}=$ $\mathcal{B}_{\circlearrowleft} u$. Recall that the Hankel operator of $\mathcal{D}_{\circlearrowleft}=\mathcal{D}(I-\mathcal{F})^{-1}$ is $\mathcal{C}_{\circlearrowleft} \mathcal{B}_{\circlearrowleft}$ since $\Psi_{\circlearrowleft}$ is a well-posed system, and compute

$$
\begin{aligned}
y_{\mathrm{opt}} & =\mathcal{C}_{\circlearrowleft} x_{0}=\mathcal{C}_{\circlearrowleft} \mathcal{B}_{\circlearrowleft} u \\
& =\pi_{+} \mathcal{D}_{\circlearrowleft} \pi_{-} u=\pi_{+} \mathcal{D}(I-\mathcal{F})^{-1} \pi_{-} u
\end{aligned}
$$

By (21), $\pi_{+} \mathcal{D}^{*} y_{\text {opt }}=0$, and this together with the computation above and the anti-causality of $\mathcal{D}^{*}$ shows that $\pi_{+} \mathcal{Z}^{*} \pi_{-}=\pi_{+} \mathcal{D}^{*} \mathcal{D}(I-\mathcal{F})^{-1} \pi_{-}=0$, i.e., $\mathcal{Z}^{*}$ is anti-causal.

By now we know that $\mathcal{Z}^{*} \mathcal{W}=\mathcal{D}^{*} \mathcal{D}$, that $\mathcal{Z}^{*}$ is anti-causal, that $\mathcal{W}$ is causal and invertible, and that the inverse of $\mathcal{W}$ is causal. Taking adjoints in the equation $\mathcal{Z}^{*} \mathcal{W}=\mathcal{D}^{*} \mathcal{D}$, we get

$$
\mathcal{W}^{*} \mathcal{Z}=\mathcal{D}^{*} \mathcal{D}=\mathcal{Z}^{*} \mathcal{W}
$$

hence

$$
\mathcal{Z W}^{-1}=\left(\mathcal{W}^{*}\right)^{-1} \mathcal{D}^{*} \mathcal{D} \mathcal{W}^{-1}=\left(\mathcal{W}^{*}\right)^{-1} \mathcal{Z}^{*}
$$

Since $\mathcal{Z \mathcal { W } ^ { - 1 }}$ is causal and $\left(\mathcal{W}^{*}\right)^{-1} \mathcal{Z}^{*}$ is anti-causal, the operator $\left(\mathcal{W}^{*}\right)^{-1} \mathcal{D}^{*} \mathcal{D} \mathcal{W}^{-1}$ is static. By Lemma 6, it is equal to a multiplication operator by an operator $M \in \mathcal{L}(U)$, i.e., $\left(\mathcal{W}^{*}\right)^{-1} \mathcal{D}^{*} \mathcal{D} \mathcal{W}^{-1}=M$ for some $M \in \mathcal{L}(U)$. The operator $M$ must be positive and invertible since $\left(\mathcal{W}^{*}\right)^{-1} \mathcal{D}^{*} \mathcal{D} \mathcal{W}^{-1}$ is positive and invertible. Let $E \in \mathcal{L}(U)$ be an arbitrary invertible operator satisfying $E^{*} E=M$; for example, we may take $E$ to be the positive square root of $M$. Then $\left(\mathcal{W}^{*}\right)^{-1} \mathcal{D}^{*} \mathcal{D} \mathcal{W}^{-1}=E^{*} E$, i.e.,

$$
\mathcal{D}^{*} \mathcal{D}=\mathcal{W}^{*} E^{*} E \mathcal{W}=\left(I-\mathcal{F}^{*}\right) E^{*} E(I-\mathcal{F})
$$

Consequently, if we define

$$
\mathcal{X}=E(I-\mathcal{F})
$$

then $\mathcal{X}$ is a spectral factor of $\mathcal{D}^{*} \mathcal{D}$, and $\mathcal{Y} \mathcal{X}=\left(\mathcal{D}(I-\mathcal{F})^{-1} E^{-1}\right)(E(I-\mathcal{F}))$ is an inner-outer factorization of $\mathcal{D}$.

To prove that $\mathcal{K}$ must be of the form $\mathcal{K}=-\pi_{+} E^{-1} \mathcal{Y}^{*} \mathcal{C}$, we observe that by (21) and (33),

$$
\pi_{+} \mathcal{D}^{*}\left(\mathcal{C}+\mathcal{D}(I-\mathcal{F})^{-1} \mathcal{K}\right)=0
$$

Replacing $\mathcal{D}^{*} \mathcal{D}$ by $\mathcal{X}^{*} \mathcal{X}=\left(I-\mathcal{F}^{*}\right) E^{*} E(I-\mathcal{F})$, we get

$$
\pi_{+} \mathcal{X}^{*} E \mathcal{K}=-\pi_{+} \mathcal{D}^{*} \mathcal{C}
$$

which we can solve for $\mathcal{K}$ (since $E^{-1}\left(\mathcal{X}^{*}\right)^{-1}$ is anti-causal) in the form

$$
\mathcal{K}=-\pi_{+} E^{-1}\left(\mathcal{X}^{*}\right)^{-1} \mathcal{D}^{*} \mathcal{C}=-\pi_{+} E^{-1} \mathcal{Y}^{*} \mathcal{C}
$$

## 7. The Generators of a Well-Posed Linear system

By a result due to Salamon [35] (and apparently discovered independently by Weiss [43, 44]), every well-posed linear system $\Psi$ has a well-defined (unbounded) control operator $B$ and a well-defined (unbounded) observation operator $C$, and formulas (3), (4), (8), (9), (11), (12), (13), and (15) hold in a weak sense (see Remark 30 below). In order to present this result we need some additional definitions.

Let $\mathcal{A}$ be the semigroup of the well-posed linear system $\Psi$ on $(U, H, Y)$, i.e., $\mathcal{A}$ is a semigroup on $H$. Denote the generator of $\mathcal{A}$ by $A$, and denote the domain $\operatorname{dom}(A)$ of $A$ by $W$. Choose an arbitrary number $\alpha$ from the resolvent set of $A$. Then $W=(\alpha I-A)^{-1} H$, and we can choose the norm in $W$ to be $\|x\|_{W}=\|(\alpha I-A) x\|_{H}$. Let $V$ be the completion of $H$ under the norm $\left\|(\alpha I-A)^{-1} x\right\|_{H}$. Then

$$
W \subset H \subset V
$$

with dense and continuous embeddings, $(\alpha I-A)$ is an (isometric) isomorphism of $W$ onto $H$, and $(\alpha I-A)$ extends to an (isometric) isomorphism of $H$ onto $V$. The semigroup $\mathcal{A}$ can be restricted to a semigroup on $W$, and extended to a semigroup on $V$. The three semigroups that we get in this way are (isometrically) isomorphic, and we denote them all with the same letter $\mathcal{A}$.

We repeat the same construction with $\mathcal{A}$ and $A$ replaced by their adjoints $\mathcal{A}^{*}$ and $A^{*}$ to get two more spaces $V^{*}=\operatorname{dom}\left(A^{*}\right)$ and $W^{*}$, with

$$
V^{*} \subset H \subset W^{*}
$$

It is possible to identify $W^{*}$ with the dual of $W$ and $V$ with the dual of $V^{*}$ if we use $H$ as the pivot space.

Proposition 29. Let $\Psi=\left[\begin{array}{cc}\mathcal{A} & \mathcal{B} \\ \mathcal{C} & \mathcal{D}\end{array}\right]$ be a causal and $\Psi^{*}=\left[\begin{array}{cc}\mathcal{A}^{*} & \mathcal{C}^{*} \\ \mathcal{B}^{*} & \mathcal{D}^{*}\end{array}\right]$ an anti-causal externally stable well-posed linear system on $(U, H, Y)$, respectively on $(Y, H, U)$. These systems have (unique) generating operators $\left.\left[\begin{array}{c}A \\ C\end{array}\right] \begin{array}{l}B \\ ?\end{array}\right]$, respectively $\left[\begin{array}{cc}A^{*} & C^{*} \\ B^{*}\end{array}\right]$, with the following properties (the question mark stands for a missing entry):
(i) $A \in \mathcal{L}(W ; H) \cap \mathcal{L}(H ; V)$ is the generator of $\mathcal{A}, A^{*} \in \mathcal{L}\left(V^{*} ; H\right) \cap \mathcal{L}\left(H ; W^{*}\right)$ is the generator of $\mathcal{A}^{*}, B \in \mathcal{L}(U ; V), C^{*} \in \mathcal{L}\left(Y ; W^{*}\right), C \in \mathcal{L}(W ; Y)$, and $B^{*} \in \mathcal{L}\left(V^{*} ; U\right)$. Moreover, for every sufficiently large real number $\alpha, \alpha I-$ $A$ has an inverse in $\mathcal{L}(H ; W) \cap \mathcal{L}(V ; H)$, and $\alpha I-A^{*}$ has an inverse in $\mathcal{L}\left(H ; V^{*}\right) \cap \mathcal{L}\left(W^{*} ; H\right)$.
(ii) $\mathcal{C} \in \mathcal{L}\left(W ; W^{1,2}\left(\mathbf{R}^{+} ; Y\right)\right), \mathcal{B}^{*} \in \mathcal{L}\left(V^{*} ; W^{1,2}\left(\mathbf{R}^{+} ; U\right)\right)$, and (4) and (9) hold for every $x \in W$ and $x^{*} \in V^{*}$. Moreover, for such $x$ and $x^{*},(\mathcal{C} x)^{\prime}=\mathcal{C} A x$ and $\left(\mathcal{B}^{*} x^{*}\right)^{\prime}=-\mathcal{B}^{*} A^{*} x^{*}$.
(iii) In the time-invariant setting, if $u \in W^{1,2}(\mathbf{R} ; U)$, then the controlled state $x=\mathcal{B} \tau u$ and the output $y=\mathcal{D} u$ of $\Psi$ satisfy $x \in B C^{1}(\mathbf{R} ; H), y \in W^{1,2}(\mathbf{R} ; Y)$, $x^{\prime}=\mathcal{B} \tau u^{\prime}=A x+B u \in B C(\mathbf{R} ; H), y^{\prime}=\mathcal{D} u^{\prime}$, and $x(t) \rightarrow 0$ and $x^{\prime}(t) \rightarrow$ 0 in $H$ as $t \rightarrow-\infty$. Analogously, in the time-invariant setting, if $y^{*} \in$ $W^{1,2}(\mathbf{R} ; Y)$, then the controlled state $x^{*}=\mathcal{C}^{*} \tau y^{*}$ and the output $u^{*}=\mathcal{D}^{*} y^{*}$ of $\Psi^{*}$ satisfy $x^{*} \in B C^{1}(\mathbf{R} ; H), u^{*} \in W^{1,2}(\mathbf{R} ; U),\left(x^{*}\right)^{\prime}=\mathcal{C}^{*} \tau\left(y^{*}\right)^{\prime}=-A^{*} x^{*}-$ $C^{*} y^{*} \in B C(\mathbf{R} ; H),\left(u^{*}\right)^{\prime}=\mathcal{D}^{*}\left(y^{*}\right)^{\prime}$, and $x^{*}(s) \rightarrow 0$ and $\left(x^{*}\right)^{\prime}(s) \rightarrow 0$ in $H$ as $s \rightarrow \infty$.
(iv) In the initial value setting, if $u \in W^{1,2}\left(\mathbf{R}^{+} ; U\right), x_{0} \in H$, and $A x_{0}+B u(0) \in$ $H$, then the controlled state $x=\mathcal{A} x_{0}+\mathcal{B} \tau \pi_{+} u$ and the output $y=\mathcal{C} x_{0}+\mathcal{D} \pi_{+} u$ of $\Psi$ satisfy

$$
\begin{gathered}
x \in C^{1}\left(\mathbf{R}^{+} ; H\right), \quad y \in W^{1,2}\left(\mathbf{R}^{+} ; Y\right), \\
x^{\prime}=\mathcal{A}\left(A x_{0}+B u(0)\right)+\mathcal{B} \tau u^{\prime}=A x+B u \in C\left(\mathbf{R}^{+} ; H\right),
\end{gathered}
$$

$x(t)$ is given by (13), and $y^{\prime}=\mathcal{C}\left(A x_{0}+B u(0)\right)+\mathcal{D} u^{\prime}$. Analogously, in the initial value setting, if $y^{*} \in W^{1,2}\left(\mathbf{R}^{-} ; Y\right), x_{0}^{*} \in H$, and $A^{*} x_{0}^{*}+C^{*} y^{*}(0) \in H$, then the controlled state $x^{*}(s)=\mathcal{A}^{*}(-s) x_{0}^{*}+\mathcal{C}^{*} \tau(s) \pi_{-} y^{*}$ and the output $u^{*}=\mathcal{B}^{*} x_{0}^{*}+\mathcal{D}^{*} \pi_{-} y^{*}$ of $\Psi^{*}$ satisfy

$$
\begin{aligned}
x^{*} \in & C^{1}\left(\mathbf{R}^{-} ; H\right), \quad u^{*} \in W^{1,2}\left(\mathbf{R}^{+} ; U\right) \\
\left(x^{*}\right)^{\prime}(s)= & -\mathcal{A}^{*}(-s)\left(A^{*} x_{0}^{*}+C^{*} y^{*}(0)\right)+\mathcal{C}^{*} \tau(s)\left(y^{*}\right)^{\prime} \\
& =-A^{*} x^{*}-C^{*} y^{*} \in C\left(\mathbf{R}^{-} ; H\right)
\end{aligned}
$$

$x^{*}(s)$ is given by (15), and $\left(u^{*}\right)^{\prime}=-\mathcal{B}^{*}\left(A^{*} x_{0}^{*}+C^{*} y^{*}(0)\right)+\mathcal{D}^{*}\left(y^{*}\right)^{\prime}$.

Outline of Proof. Most of these claims follow directly from [34, Lemmas 2.3 and 2.5] and [35, Theorem 3.1 and Lemma 3.2]. The key step is to prove that $B \in \mathcal{L}(U ; V)$ and $C \in \mathcal{L}(W ; Y)$ and that (4) and (13) hold in the appropriate spaces, because that enables us to regard $\left[\begin{array}{cc}\mathcal{A} & \mathcal{B} \\ 0 & 0\end{array}\right]$ as a classical system on $(U, V, ?)$ and $\left[\begin{array}{cc}\mathcal{A} & 0 \\ \mathcal{C} & 0\end{array}\right]$ as a classical system on $(?, W, Y)$ (the question marks represent irrelevant spaces). This makes it is possible to use standard semigroup theory as presented in, e.g., [29]. To show that $x(t) \rightarrow 0$ in $H$ as $t \rightarrow-\infty$ in (iii) is suffices to observe that

$$
\begin{aligned}
\|x(t)\|_{H} & =\|\mathcal{B} \tau(t) u\|_{H} \\
& =\left\|\mathcal{B} \pi_{-} \tau(t) u\right\|_{H} \\
& =\left\|\mathcal{B} \tau(t) \tau(-t) \pi_{-} \tau(t) u\right\|_{H} \\
& =\left\|\mathcal{B} \tau(t) \pi_{(-\infty, t]} u\right\|_{H} \\
& \leq\|\mathcal{B}\|_{\mathcal{L}\left(L^{2}(\mathbf{R} ; U) ; H\right)}\left\|\pi_{(-\infty, t]} u\right\|_{L^{2}(\mathbf{R} ; U)}
\end{aligned}
$$

which tends to zero as $t \rightarrow-\infty$. A similar proof with $u$ replaced by $u^{\prime}$ shows that $x^{\prime}(t) \rightarrow 0$ in $H$ as $t \rightarrow-\infty$. To deduce the statements about the adjoint system $\Psi^{*}$ it suffices to reflect the time axis and to apply the corresponding result for $\Psi$.

Remark 30. According to Proposition 29, (4) holds for each $x \in W$, and (13) holds as an equation in $H$, provided that the integral is interpreted as a Bochner integral in $V$. To derive a weak version of (11) we argue as follows. Fix some $u \in L^{2}(\mathbf{R} ; U)$, and define $u_{T}=\pi_{[T, \infty)} u$ and $x_{T}=\mathcal{B} \tau u_{T}$. Then $u-u_{T}=\pi_{(-\infty, T)} u \rightarrow 0$ in $L^{2}(\mathbf{R} ; U)$ as $T \rightarrow-\infty$, so $x_{T}(t) \rightarrow x(t)=\mathcal{B} \tau(t) u$ in $H$, uniformly in $t$. Moreover, $x_{T}(t)=0$ for $t \leq T$. Since we can regard $\left[\begin{array}{cc}\mathcal{A} & \mathcal{B} \\ 0 & 0\end{array}\right]$ as a classical system on $(U, V, ?)$ (the output space is irrelevant), it follows from (2) that

$$
x_{T}(t)=\int_{T}^{t} \mathcal{A}(t-s) B u(s) d s, \quad t \geq T
$$

We conclude that (11) holds in the sense that

$$
x(t)=\lim _{T \rightarrow-\infty} \int_{T}^{t} \mathcal{A}(t-s) B u(s) d s, \quad t \in \mathbf{R}
$$

where the integral is computed in $V$, but the limit is taken in $H$ and is uniform in $t$ (we interpret the integral as zero if $t<T$ ). In particular, taking $t=0$ we find that

$$
\mathcal{B} u=\lim _{T \rightarrow-\infty} \int_{T}^{0} \mathcal{A}(-s) B u(s) d s
$$

where the integral is computed in $V$ but the limit is taken in $H$. Analogous interpretations are valid for the adjoint equations (8), (9), (12), and (15).

Proposition 29 is not quite sufficient for our purposes. We need a small enhancement, which involves some additional Hilbert spaces. According to Proposition 29, the values of the states $x$ and $x^{*}$ and their derivatives lie in $H$, but we need to know that the states actually lie in some smaller spaces. These smaller spaces are not uniquely determined. We choose to follow Salamon and use the smallest possible spaces, instead of the larger spaces introduced by G. Weiss. These smallest possible spaces are "range spaces" of certain mappings, constructed as follows:

Lemma 31. Let $E \in \mathcal{L}(X ; Y)$, where $X$ and $Y$ are Banach spaces. For each $y \in \operatorname{range}(E)$, let $\|y\|_{Z}=\inf \left\{\|x\|_{X} \mid y=E x\right\}$. This defines a norm on range $(E)$ that makes range $(E)$ a Banach space $Z$, and this norm is (not necessarily strictly) stronger than the norm of $Y$. The operator $E$ induces an isometric isomorphism of the quotient $X / \operatorname{ker}(E)$ onto $Z$, and $E$ is continuous and open from $X$ onto $Z$. If $X$ is a Hilbert space, then so is $Z$, and then $E$ is an isometric isomorphism of $\operatorname{ker}(E)^{\perp}$ onto $Z$.

This result is standard but difficult to find in the literature. The idea is to factor $E$ into $E=F \pi$, where $\pi: X \rightarrow X / \operatorname{ker}(E)$ is the (continuous and open) quotient map, and $F: X / \operatorname{ker}(E)$ is injective [3, Problem F, p. 266], and to let the norm in range $(E)=\operatorname{range}(F)$ be the one induced by $F$. This makes $F$ an isometric isomorphism of $X / \operatorname{ker}(E)$ onto $Z$.

Let $\Psi$ be a well-posed linear system on $(U, H, Y)$ with generating operators $\left[\begin{array}{ccc}A & B \\ C & ?\end{array}\right]$. Choose an arbitrary number $\alpha$ from the resolvent sets of $A$. We denote the range space (constructed as in Lemma 31) of the mapping $(\alpha I-A)^{-1}\left(\begin{array}{ll}I & B\end{array}\right): H \times$ $U \rightarrow H$ by $W_{B}$, and the corresponding space that we get by replacing $\Psi$ by its adjoint by $V_{C}^{*}$. These spaces are independent of the particular value of $\alpha$.

The elements of $W_{B}$ and $V_{C}^{*}$ can be characterized in the following way:
Lemma 32. The following conditions are equivalent:
(i) $x \in W_{B}$;
(ii) $x=(\alpha I-A)^{-1}\left(x_{1}+B u\right)$ for some $x_{1} \in H$ and $u \in U$;
(iii) $x=x_{2}+(\alpha I-A)^{-1} B u$ for some $x_{2} \in W$ and $u \in U$;
(iv) $x \in H$ and $A x+B u \in H$ for some $u \in U$.

Likewise, the following conditions are also equivalent:
(v) $x^{*} \in V_{C}^{*}$;
(vi) $x^{*}=\left(\alpha I-A^{*}\right)^{-1}\left(x_{1}^{*}+C^{*} y^{*}\right)$ for some $x_{1}^{*} \in H$ and $y^{*} \in Y$;
(vii) $x^{*}=x_{2}^{*}+\left(\alpha I-A^{*}\right)^{-1} C^{*} y^{*}$ for some $x_{2}^{*} \in V^{*}$ and $y^{*} \in Y$;
(viii) $x^{*} \in H$ and $A^{*} x^{*}+C^{*} y^{*} \in H$ for some $y^{*} \in Y$.

Proof. The equivalence of (i) and (ii) follows from the definition of $W_{B}$. Take $x_{2}=(\alpha I-A)^{-1} x_{1}$ to show that (ii) and (iii) are equivalent. If $x$ is of the form (ii), then $x \in H$ and $A x+B u=-x_{1}+\alpha(\alpha I-A)^{-1}\left(x_{1}+B u\right) \in H$, so (iv) holds. Conversely, if (iv) holds, then (ii) is satisfied with $x_{1}=\alpha x-(A x+B u)$. The equivalence of (v)-(viii) is proved analogously.

It is easy to see that

$$
W \subset W_{B} \subset H \text { and } V^{*} \subset V_{C}^{*} \subset H
$$

with continuous injections.
Obviously, $W_{B}$ and $V_{C}^{*}$ are dense in $H$ since $W$ and $V^{*}$ are dense in $H$. However, $W$ need not be dense in $W_{B}$, and $V^{*}$ need not be dense in $V_{C}^{*}$. In particular, in the realizations discussed in [37] and [38] $W$ is a closed subspace of $W_{B}$ and $V^{*}$ is a closed subspace of $V_{C}^{*}$. For example, the realization that is denoted by $\Xi^{-}$in [37] is Salamon's exactly controllable realization which we get by taking $H$ to be $H=L^{2}\left(\mathbf{R}^{-} ; U\right), \mathcal{A}(t)$ to be the unilateral left-shift $\tau(t) \pi_{-}$on $L^{2}\left(\mathbf{R}^{-} ; U\right), \mathcal{B} u=$ $\pi_{-} u$, and $\mathcal{C} x=\pi_{+} \mathcal{D} x$ [39, Definition 2.10]. For this realization $W=W_{0}^{1,2}\left(\mathbf{R}^{-} ; U\right)$ and $W_{B}=W^{1,2}\left(\mathbf{R}^{-} ; U\right)$, so $W$ is a closed subspace of $W_{B}$. Salamon introduces the space $W_{B}$ in $[34]$ (denoting it by $Z$ ) in the special case where the mapping
$(\alpha I-A)^{-1}\left(\begin{array}{ll}I & B\end{array}\right)$ is one-to-one (a case that he refers to as an abstract boundary control system; the exactly controllable realization is of this type).

Proposition 29 implies the following result:
Proposition 33. The conclusion of Proposition 29 can be strengthened as follows:
(i) In part (iii), $x \in B C\left(\mathbf{R} ; W_{B}\right), x^{*} \in B C\left(\mathbf{R} ; V_{C}^{*}\right), x(t) \rightarrow 0$ in $W_{B}$ as $t \rightarrow-\infty$, and $x^{*}(s) \rightarrow 0$ in $V_{C}^{*}$ as $s \rightarrow+\infty$.
(ii) In part (iv), $x \in C\left(\mathbf{R}^{+} ; W_{B}\right)$ and $x^{*} \in C\left(\mathbf{R}^{+} ; V_{C}^{*}\right)$.

Proof. Use the fact that $x^{\prime}=A x+B u$ to write $x$ in the form

$$
x=(\alpha I-A)^{-1}\left(\alpha x-x^{\prime}+B u\right) .
$$

Since $x$ and $x^{\prime}$ are continuous in $H$ and $u$ is continuous in $U$, this implies that $x$ is continuous in $W_{B}$. The other claims are proved in a similar way.

Remark 34. By Lemma 32, the spaces $W_{B}$ and $V_{C}^{*}$ can be characterized as the spaces of permitted initial values $x_{0}$ and $x_{0}^{*}$ in part (iv) of Proposition 29. By Proposition 33, they are invariant in the sense that the states $x$ and $x^{*}$ in parts (iii) and (iv) of Proposition 29 stay in these spaces for all time. Moreover, by Proposition 37 below, they are large enough to contain the domains of all state feedback perturbations of $A$ and $A^{*}$.

In general, a well-posed linear system need not have a well-defined feed-through operator $D$, so equations (5), (10), (14), and (16) cannot possibly be true in general. This fact motivated G. Weiss to introduce the class of regular systems. In [45, Theorem 5.8] Weiss gives eight equivalent characterizations of regularity, one of which is the following:

Definition 35. A causal time-invariant operator $\mathcal{D}: L^{2}(\mathbf{R} ; U) \rightarrow L^{2}(\mathbf{R} ; Y)$ is regular iff, for every $u_{0} \in U$, the strong Cesàro mean of order one

$$
D u_{0}=\lim _{t \rightarrow 0+} \frac{1}{t} \int_{0}^{t}\left(\mathcal{D} \chi_{\mathbf{R}^{+}} u_{0}\right)(s) d s
$$

exists (here $\chi_{\mathbf{R}^{+}}$stands for the characteristic function of $\mathbf{R}^{+}$, and $\mathcal{D} \chi_{\mathbf{R}^{+}} u_{0}$ is the step response corresponding to the constant input $u_{0}$ ). An anti-causal timeinvariant operator $\mathcal{D}^{*}: L^{2}(\mathbf{R} ; Y) \rightarrow L^{2}(\mathbf{R} ; U)$ is regular iff, for every $y_{0}^{*} \in Y$, the strong Cesàro mean of order one

$$
D^{*} y_{0}^{*}=\lim _{s \rightarrow 0-} \frac{1}{s} \int_{0}^{s}\left(\mathcal{D}^{*} \chi_{\mathbf{R}^{-}} y_{0}^{*}\right)(t) d t
$$

exists (here $\chi_{\mathbf{R}^{-}}$stands for the characteristic function of $\mathbf{R}^{-}$, and $\mathcal{D}^{*} \chi_{\mathbf{R}^{+}} y_{0}^{*}$ is the step response corresponding to the constant input $y_{0}^{*}$ ). The operators $D: U \rightarrow Y$ and $D^{*}: Y \rightarrow U$ defined above are called the feed-through operators of $\mathcal{D}$ and $\mathcal{D}^{*}$. The systems $\Psi$ and $\Psi^{*}$ are regular iff their input/output maps are regular.

Thus, a system is regular iff its step response has a one-sided limit in the strong $C_{1}$-sense at the origin. It is known [46, p. 38] that the adjoint of a regular input/output map need not be regular.

Let us rewrite equations (5), (10), (14), and (16) in terms of the controlled states $x(t)$ and $x^{*}(s)$ of $\Psi$ and $\Psi^{*}$ in the common forms

$$
\begin{aligned}
y(t) & =C x(t)+D u(t) \\
u^{*}(s) & =B^{*} x^{*}(s)+D^{*} y^{*}(s)
\end{aligned}
$$

Then, as Weiss proved, for regular systems these equations remain valid in a weak sense:

Proposition 36. Let $\Psi$ and $\Psi^{*}$ be well-posed linear systems on $(U, H, Y)$, respectively, $(Y, H, U)$, with generating operators $\left[\begin{array}{cc}A & B \\ C & ?\end{array}\right]$, respectively $\left[\begin{array}{cc}A^{*} & C^{*} \\ B^{*} & ?\end{array}\right]$.
(i) $\Psi$ is regular if and only if the strong Abel limit

$$
\bar{C} x=\lim _{\lambda \rightarrow+\infty} C \lambda(\lambda I-A)^{-1} x
$$

exists for every $x \in W_{B}$, and $\Psi^{*}$ is regular if and only if the strong Abel limit

$$
\bar{B}^{*} x^{*}=\lim _{\lambda \rightarrow+\infty} B^{*} \lambda\left(\lambda I-A^{*}\right)^{-1} x^{*}
$$

exists for every $x^{*} \in V_{C}^{*}$. The operators $\bar{C}$ and $\bar{B}^{*}$ defined in this way are bounded linear operators from $W_{B}$ into $Y$ and from $V_{C}^{*}$ into $U$, respectively, and they are extensions of the operators $C \in \mathcal{L}(W ; Y)$ and $B^{*} \in \mathcal{L}\left(V^{*} ; U\right)$.
(ii) In the regular case, denote the feed-through operators of $\Psi$ and $\Psi^{*}$ by $D$ and $D^{*}$, respectively. Then the conclusion of Proposition 29 can be strengthened as follows: Both in parts (iii) and (iv), $y=\bar{C} x+D u$ and $u^{*}=\bar{B}^{*} x^{*}+D^{*} y^{*}$.
Proof. Part (i) follows easily from the equivalence of (1) and (4) in [45, Theorem 5.8], and Part (ii) from [45, Remark 6.2].

Let us remark that our operator $\bar{C}$ is a restriction of the operator that Weiss denotes by $\widetilde{C}_{L}$ in [45] and by $C_{\Lambda}$ in [46], and also a restriction of Weiss' operator $C_{L}$. He does not explicitly study $\bar{B}^{*}$ and the time-invariant setting.

Since $W$ and $V^{*}$ need not be dense in $W_{B}$, respectively $V_{C}^{*}$, the extensions $\bar{C}$ and $\bar{B}^{*}$ are not uniquely determined by Weiss $C$ and $B^{*}$. Salamon uses a different extension $K$ of $C$ in [34, Section 2.2]. As shown in [40], some of our final results can be extended to the non-regular case through the use of Salamon's extension of $C$ and the corresponding extension of $B^{*}$ instead of our extensions $\bar{C}$ and $\bar{B}^{*}$ (induced by Weiss' extensions).

In this work we make no explicit use of the notion of the transfer function of the input/output map of a well-posed linear system, and we refer the reader to [39], [45], [46], and [50] for discussions of these functions. In our setting, the transfer function of $\mathcal{D}$ is an $H^{\infty}$-function over the right half-plane, and the transfer function of $\mathcal{D}^{*}$ is an $H^{\infty}$-function over the left half-plane (in Weiss' setting, both of these are defined on right-half planes).

## 8. The Generators of the Closed Loop System

Consider the closed loop output feedback system $\Psi_{L}$ introduced in Proposition 20. We denote the generating operators of this system by $\left[\begin{array}{cc}A_{L} & B_{L} \\ C_{L} & ?\end{array}\right]$, and in the regular case by $\left[\begin{array}{cc}A_{L} & B_{L} \\ C_{L} & D_{L}\end{array}\right]$. In the classical case we know that these generators are given by (24), and Salamon and Weiss have extended this relation to general wellposed closed loop linear output feedback systems.

The generator $A_{L}$ of the closed loop semigroup $S_{L}$ was computed by Salamon in [34, Section 4]. His description of the domain $\operatorname{dom}\left(A_{L}\right)$ of $A_{L}$ was incomplete in the case where the observation operator is unbounded and the control space $U$ infinitedimensional, due to the fact that Salamon never makes the connection between the well-posedness of the closed loop system and the invertibility of the loop gain transfer function in some right-half plane that Weiss makes in [46, Proposition 3.6].

However, by combining [34, Lemma 4.4(iii)] with [46, Proposition 3.6] we get a satisfactory result (presented in [46] in the regular case). We shall not need the exact result here, but we shall need the following consequence of this result, that can also be derived from [46, Proposition 6.6].

Proposition 37. Denote the generators of the closed loop system $\Psi_{L}$ in Proposition 20 by $\left[\begin{array}{cc}A_{L} & B_{L} \\ C_{L} & ?\end{array}\right]$, and let $W_{B_{L}}$ and $V_{C_{L}}^{*}$ be the closed loop versions of $W_{B}$ and $V_{C}^{*}$, respectively. Then $W_{B_{L}}=W_{B}$ and $V_{C}^{*}=V_{C_{L}}^{*} ;$ hence $\operatorname{dom}\left(A_{L}\right) \subset W_{B}$ and $\operatorname{dom}\left(A_{L}^{*}\right) \subset V_{C}^{*}$.

Thus, in the sequel we shall throughout replace $W_{B_{L}}$ and $V_{C_{L}}^{*}$ by $W_{B}$ and $V_{C}^{*}$, respectively.

In the regular case the closed loop generating operators have been determined by Weiss [46]; see, in particular, Theorem 4.7, Proposition 4.8, and Section 7 in that paper. From there we extract the following results:
Proposition 38. Let $\Psi$ and $\Psi_{L}$ be the open and closed loop systems in Proposition 20, with admissible output feedback operator L.
(i) If both $\Psi$ and $\Psi^{*}$ are regular, then both $\Psi_{L}$ and $\Psi_{L}^{*}$ are regular, $I-D L$ is invertible in $\mathcal{L}(Y)$, and $I-L D$ is invertible in $\mathcal{L}(U)$.
(ii) Suppose that $\Psi$ and $\Psi^{*}$ are regular. Denote the generators of $\Psi$ and $\Psi_{L}$ by $\left[\begin{array}{cc}A & B \\ C & B\end{array}\right]$ and $\left[\begin{array}{cc}A_{L} & B_{L} \\ C_{L} & D_{L}\end{array}\right]$, respectively. Then the relation

$$
\left[\begin{array}{cc}
A_{L} & B_{L} \\
\bar{C}_{L} & D_{L}
\end{array}\right]=\left[\begin{array}{cc}
A+B L(I-D L)^{-1} \bar{C} & B(I-L D)^{-1} \\
(I-D L)^{-1} \bar{C} & D(I-L D)^{-1}
\end{array}\right]
$$

holds in the following sense: the equation for $A_{L}$ is valid on $\operatorname{dom}\left(A_{L}\right) \subset$ $W_{B_{L}}=W_{B}$, the equation for $\bar{C}_{L}$ is valid on $W_{B_{L}}=W_{B}$, the equation for $D_{L}$ is valid on $U$, and the equation for $B_{L}$ should be interpreted as $\bar{B}_{L}^{*}=$ $\left(I-D^{*} L^{*}\right)^{-1} \bar{B}^{*}$, which is valid on $V_{C_{L}}^{*}=V_{C}^{*}$.
Another different interpretation of the formula $B_{L}=B(I-L D)^{-1}$ is given in [46, Section 7].

Applying the preceding result with $L=(0 I)$ to the system $\Psi_{\text {ext }}$ in (30), we get the following result:
Proposition 39. In addition to the hypothesis of Theorem 27, suppose that both the extended system $\Psi_{\text {ext }}$ given by (30) and its adjoint system $\Psi_{\text {ext }}^{*}$ given by (31) are regular, i.e., $\mathcal{D}, \mathcal{D}^{*}, \mathcal{F}$, and $\mathcal{F}^{*}$ are regular. Denote the generating operators of $\Psi_{\text {ext }}$ by $\left[\begin{array}{c}A \\ \binom{C}{K}\end{array}\left(\begin{array}{c}B \\ F \\ F\end{array}\right)\right]$ and the generating operators of $\Psi_{\circlearrowleft}$ by $\left.\left[\begin{array}{c}A_{\circlearrowleft} \\ C_{\circlearrowleft} \\ K_{\circlearrowleft}\end{array}\right)\left(\begin{array}{c}B_{\circlearrowleft} \\ D_{\circlearrowleft} \\ F_{\circlearrowleft}\end{array}\right)\right]$ (the regularity of $\Psi_{\circlearrowleft}$ follows from Proposition 38.)
(i) Under these assumptions the outer factor $\mathcal{X}$ and the inner factor $\mathcal{Y}$ are regular together with their adjoints. If we denote the feed-through operator of $\mathcal{X}$ by $X$ and the feed-through operator of $\mathcal{Y}$ by $Y$, then $X=E(I-F), X$ is invertible, and $D=Y X$.
(ii) Construct the range space $V_{(C, K)}^{*}$ in the same way as the space $V_{C}^{*}$ was constructed, except that $\Psi$ is replaced by $\Psi_{\text {ext }}$. Then $V_{C}^{*} \subset V_{(C, K)}^{*}$, and the relation

$$
\left[\begin{array}{cc}
A_{\circlearrowleft} & B_{\circlearrowleft} \\
\binom{\bar{C}_{\circlearrowleft}}{\bar{K}_{\circlearrowleft}} & \binom{D_{\circlearrowleft}}{F_{\circlearrowleft}}
\end{array}\right]=\left[\begin{array}{cc}
A+B X^{-1} E \bar{K} & B X^{-1} E \\
\binom{\bar{C}+Y E \bar{K}}{X^{-1} E \bar{K}} & \binom{Y E}{X^{-1} E-I}
\end{array}\right]
$$

holds in the following sense: the equation for $A_{\circlearrowleft}$ is valid on $\operatorname{dom}\left(A_{\circlearrowleft}\right) \subset W_{B}$, the equations for $\bar{C}_{\circlearrowleft}$ and $\bar{K}_{\circlearrowleft}$ are valid on $W_{B}$, the equations for $D_{\circlearrowleft}$ and $F_{\circlearrowleft}$ are valid on $U$, the equation for $F_{\circlearrowleft}$ is valid on $Y$, and the equation for $B_{\circlearrowleft}$ should be interpreted as $\bar{B}_{\circlearrowleft}^{*}=E^{*}\left(X^{*}\right)^{-1} \bar{B}^{*}$, which is valid on $V_{(C, K)}^{*}$.
This follows directly from Theorem 27 and Proposition 38. It is also true that $\operatorname{dom}\left(A_{\circlearrowleft}^{*}\right) \subset V_{K}^{*} \subset V_{(C, K)}^{*}$, with continuous inclusions where $V_{K}^{*}$ is the range space constructed in the same way as $V_{C}^{*}$ and $V_{(C, K)}^{*}$ were constructed, but with $\Psi$ replaced by the system $\left[\begin{array}{c}\mathcal{A} \\ \mathcal{K} \\ \mathcal{F}\end{array}\right]$. That this inclusion is true follows from the fact that the closed loop semigroup $\mathcal{A}_{\circlearrowleft}$ depends only on $\mathcal{A}, \mathcal{B}, \mathcal{K}$, and $\mathcal{F}$, and not explicitly on $\mathcal{C}$.

We have the following corollary to Theorem 27 and Proposition 38:
Corollary 40. Let the assumption of Theorem 27 hold. If $\Psi$ is regular and $\mathcal{D}^{*} \mathcal{D}$ has a regular spectral factor $\mathcal{X}$, then every spectral factor of $\mathcal{D}^{*} \mathcal{D}$ is regular, and this is the case if and only if the extended system $\Psi_{\mathrm{ext}}$ in (30) is regular. If $\Psi^{*}$ is regular and $\mathcal{D}^{*} \mathcal{D}$ has a spectral factor $\mathcal{X}$ with a regular adjoint $\mathcal{X}^{*}$, then every spectral factor of $\mathcal{D}^{*} \mathcal{D}$ has a regular adjoint, and this is the case if and only if the adjoint extended system $\Psi_{\mathrm{ext}}^{*}$ in (31) is regular. If both $\Psi$ and $\Psi^{*}$ are regular and $\mathcal{D}^{*} \mathcal{D}$ has a regular spectral factor $\mathcal{X}$ with a regular adjoint $\mathcal{X}^{*}$, then both the closed loop systems $\Psi_{\circlearrowleft}$ and $\Psi_{\circlearrowleft}^{*}$ are regular, the feed-through operator $X$ of $\mathcal{X}$ is invertible, and there is a unique spectral factor of $\mathcal{D}^{*} \mathcal{D}$ with a positive self-adjoint feed-through operator. Moreover, in this case there is a unique feedback pair $(\mathcal{K}, \mathcal{F})$ in Theorem 27 for which the feed-through operator of $\mathcal{F}$ is zero, namely

$$
\left(\begin{array}{ll}
\mathcal{K} & \mathcal{F}
\end{array}\right)=\left(-\pi_{+} X^{-1} \mathcal{Y}^{*} \mathcal{C} \quad\left(I-X^{-1} \mathcal{X}\right)\right)
$$

where $\mathcal{Y X}$ is an arbitrary inner-outer factorization of $\mathcal{D}$, and $X$ is the feed-through operator of $\mathcal{X}$.

Proof. Most of these claims follow immediately from Theorem 27 and Proposition 38. To get the unique spectral factor that has a positive self-adjoint feedthrough operator we take an arbitrary spectral factor $\mathcal{X}$ with feed-through operator $X$, and multiply it to the left by the unitary operator $\left(X^{*} X\right)^{1 / 2} X^{-1}$, where $\left(X^{*} X\right)^{1 / 2}$ is the unique positive self-adjoint square root of $X^{*} X$. To get the unique feedback pair $(\mathcal{K}, \mathcal{F})$ for which the feed-through part of $\mathcal{F}$ is zero we take an arbitrary spectral factor $\mathcal{X}$ with feed-through operator $X$, and choose the operator $E$ in Theorem 27 to be $E=X$.

We remark that the case treated in [37] was regular, and there we throughout normalized $\mathcal{X}$ to have a positive self-adjoint feed-through operator, and took $E$ in Theorem 27 to be the feed-through operator of $\mathcal{X}$.

## 9. The Open and Closed Loop Lyapunov Equations

In the sequel we suppose throughout that $\mathcal{D}$ is coercive. We claim that the feedback operator can be expressed in terms of the Riccati operator, and conversely, that the Riccati operator can be computed from the feedback operator through a Lyapunov equation. By combining these two facts we get an algebraic Riccati equation for the Riccati operator. Since the proof of the fact that the Riccati operator can be computed from the feedback operator through a Lyapunov equation
is quite elementary and does not require any regularity assumptions, we start with this part.

It is possible to develop (at least) three different Lyapunov equations, one for the open loop system $\Psi_{\text {ext }}$, one for the closed loop system $\Psi_{\circlearrowleft}$, and a mixture of these two:

Theorem 41. Let $\Psi=\left[\begin{array}{cc}\mathcal{A} & \mathcal{B} \\ \mathcal{C} & \mathcal{D}\end{array}\right]$ be a causal externally stable well-posed linear system on $(U, H, Y)$ with a coercive input/output map $\mathcal{D}$. Let $E$ be the operator in Theorem 27, denote the generating operators of the system $\Psi_{\text {ext }}$ in (30)by $\left.\left[\begin{array}{cc}A \\ C \\ K\end{array}\right)\left(\begin{array}{l}B \\ ? \\ ?\end{array}\right)\right]$, and denote the generating operators of the system $\Psi_{\circlearrowleft i n}$ (29) by $\left.\left[\begin{array}{cc}A_{\circlearrowleft} & B_{\circlearrowleft} \\ C_{\circlearrowleft} \\ K_{\circlearrowleft}\end{array}\right)\binom{?}{?}\right]$. With this notation, the Riccati operator $\Pi$ of $\Psi$ satisfies the open loop Lyapunov equation

$$
\begin{aligned}
\left\langle A x_{1}, \Pi x_{0}\right\rangle_{H}+\left\langle x_{1}, \Pi A x_{0}\right\rangle_{H}=-\left\langle C x_{1}, C x_{0}\right\rangle_{Y}+\left\langle E K x_{1}, E K x_{0}\right\rangle_{U} \\
x_{0}, x_{1} \in \operatorname{dom}(A)
\end{aligned}
$$

the closed loop Lyapunov equation

$$
\left\langle A_{\circlearrowleft} x_{1}, \Pi x_{0}\right\rangle_{H}+\left\langle x_{1}, \Pi A_{\circlearrowleft} x_{0}\right\rangle_{H}=-\left\langle C_{\circlearrowleft} x_{1}, C_{\circlearrowleft} x_{0}\right\rangle_{Y}, \quad x_{0}, x_{1} \in \operatorname{dom}\left(A_{\circlearrowleft}\right),
$$

and the mixed Lyapunov equation
$\left\langle A_{\circlearrowleft} x_{1}, \Pi x_{0}\right\rangle_{H}+\left\langle x_{1}, \Pi A x_{0}\right\rangle_{H}=-\left\langle C_{\circlearrowleft} x_{1}, C x_{0}\right\rangle_{Y}, \quad x_{0} \in \operatorname{dom}(A), x_{1} \in \operatorname{dom}\left(A_{\circlearrowleft}\right)$.
If $\Psi$ is strongly stable, then each of these equations determines $\Pi$ uniquely (as a function of the other operators).

Proof. By Theorem 27 and Proposition 29, we have for all $x_{0}$ and $x_{1}$ in $W=$ $\operatorname{dom}(A)$ and for all $t \in \mathbf{R}^{+}$,

$$
\begin{aligned}
\left\langle\mathcal{A}(t) x_{1}, \Pi \mathcal{A}(t) x_{0}\right\rangle_{H}= & \left\langle\mathcal{C} \mathcal{A}(t) x_{1}, \mathcal{C A}(t) x_{0}\right\rangle_{L^{2}\left(\mathbf{R}^{+} ; Y\right)} \\
& -\left\langle E \mathcal{K} \mathcal{A}(t) x_{1}, E \mathcal{K A}(t) x_{0}\right\rangle_{L^{2}\left(\mathbf{R}^{+} ; U\right)} \\
= & \int_{0}^{\infty}\left\langle C \mathcal{A}(s+t) x_{1}, C \mathcal{A}(s+t) x_{0}\right\rangle_{Y} d s \\
& -\int_{0}^{\infty}\left\langle E K \mathcal{A}(s+t) x_{1}, E K \mathcal{A}(s+t) x_{0}\right\rangle_{U} d s \\
= & \int_{t}^{\infty}\left\langle C \mathcal{A}(s) x_{1}, C \mathcal{A}(s) x_{0}\right\rangle_{Y} d s \\
& -\int_{t}^{\infty}\left\langle E K \mathcal{A}(s) x_{1}, E K \mathcal{A}(s) x_{0}\right\rangle_{U} d s
\end{aligned}
$$

Differentiating both sides of this equation with respect to $t$ and taking $t=0$, we get the open loop Lyapunov equation. To get the closed loop Lyapunov equation we argue in the same way, except that we start from the equation $\mathcal{A}_{\circlearrowleft}^{*} \Pi \mathcal{A}_{\circlearrowleft}=$ $\mathcal{A}_{\circlearrowleft}^{*} \mathcal{C}_{\circlearrowleft}^{*} \mathcal{C}_{\circlearrowleft} \mathcal{A}_{\circlearrowleft}$. Finally, the mixed Lyapunov equation is also deduced in the same way, starting from the fact that $\mathcal{A}_{\circlearrowleft}^{*} \Pi \mathcal{A}=\mathcal{A}_{\circlearrowleft}^{*} \mathcal{C}_{\circlearrowleft}^{*} \mathcal{C} \mathcal{A}$.

To prove that the open loop Lyapunov equation determines $\Pi$ uniquely we argue as follows. Take $x_{0}$ and $x_{1}$ in $W$. Then $\left\langle\mathcal{A} x_{1}, \Pi \mathcal{A} x_{0}\right\rangle$ is continuously differentiable on $\mathbf{R}^{+}$, with the derivate $\left\langle A \mathcal{A} x_{1}, \Pi A \mathcal{A} x_{0}\right\rangle$. Thus, if $\Pi$ satisfies the open loop

Lyapunov equation, then for each $t>0$,

$$
\begin{aligned}
\left\langle x_{1}, \Pi x_{0}\right\rangle_{H}- & \left\langle\mathcal{A}(t) x_{1}, \Pi \mathcal{A}(t) x_{0}\right\rangle_{H} \\
= & \int_{0}^{t}\left\langle C \mathcal{A}(s) x_{1}, C \mathcal{A}(s) x_{0}\right\rangle_{Y} d s \\
& -\int_{0}^{t}\left\langle E K \mathcal{A}(s) x_{1}, E K \mathcal{A}(s) x_{0}\right\rangle_{U} d s \\
= & \int_{0}^{t}\left\langle\left(\mathcal{C} x_{1}\right)(s),\left(\mathcal{C} x_{1}\right)(s) x_{0}\right\rangle_{Y} d s \\
& -\int_{0}^{t}\left\langle E\left(\mathcal{K} x_{1}\right)(s), E\left(\mathcal{K} x_{1}\right)(s)\right\rangle_{U} d s
\end{aligned}
$$

Let $t \rightarrow \infty$. Because of the strong stability of $\Psi, \mathcal{A}(t) x_{0} \rightarrow 0$ and $\mathcal{A}(t) x_{1} \rightarrow 0$ in $H$, and we conclude that

$$
\left\langle x_{1}, \Pi x_{0}\right\rangle_{H}=\left\langle\mathcal{C} x_{1}, \mathcal{C} x_{0}\right\rangle_{L^{2}\left(\mathbf{R}^{+} ; Y\right)}-\left\langle E \mathcal{K} x_{1}, E \mathcal{K} x_{0}\right\rangle_{L^{2}\left(\mathbf{R}^{+} ; U\right)} .
$$

This being true for all $x_{0}$ and $x_{1}$ in $W$, we must have $\Pi=\mathcal{C}^{*} \mathcal{C}-\mathcal{K}^{*} E^{*} E \mathcal{K}$.
The proofs that also the closed loop and the mixed Lyapunov equations determine $\Pi$ uniquely are similar (cf. Lemma 21).

The uniqueness proofs given above have been adapted from [19, Theorems 3 and 4].

Remark 42. As is well known, in the classical case the different Lyapunov equations given in Theorem 41 are equivalent; we can pass from one to another by using the appropriate version of (24) and the connection between the Riccati operator $\Pi$ and the feedback operator $K$ displayed in Theorem 44. A crucial role in these formulas is played by the feed-through operator $D$ of the original system. Since a general well-posed linear system does not have a well-defined feed-through operator, it is difficult if not impossible to pass from one of the Lyapunov equations given above to another in the general case. Note that the closed loop and mixed Lyapunov equations are quite elementary, since they are based on the identities $\Pi=\mathcal{C}_{\circlearrowleft}^{*} \mathcal{C}_{\circlearrowleft}$ and $\Pi=\mathcal{C}_{\circlearrowleft}^{*} \mathcal{C}$ proved in Lemma 15 , but the validity of the open loop Lyapunov equation is a deeper result.

Corollary 43. Make the same same assumption and introduce the same notation as in Theorem 41.
(i) The Riccati operator $\Pi$ satisfies the equations

$$
\begin{aligned}
\Pi A x & =-\left(A^{*} \Pi+C^{*} C-K^{*} E^{*} E K\right) x=-\left(A_{\circlearrowleft}^{*} \Pi+C_{\circlearrowleft}^{*} C\right) x, \quad x \in \operatorname{dom}(A), \\
\Pi A_{\circlearrowleft} x & =-\left(A^{*} \Pi+C^{*} C_{\circlearrowleft}\right) x=-\left(A_{\circlearrowleft}^{*} \Pi+C_{\circlearrowleft}^{*} C_{\circlearrowleft}\right) x, \quad x \in \operatorname{dom}\left(A_{\circlearrowleft}\right) .
\end{aligned}
$$

In particular, the operators $A^{*} \Pi+C^{*} C-K^{*} E^{*} E K$ and $A_{\circlearrowleft}^{*} \Pi+C_{\circlearrowleft}^{*} C$ map $\operatorname{dom}(A)$ continuously into $H$, and the operators $A^{*} \Pi+C^{*} C_{\circlearrowleft}$ and $A_{\circlearrowleft}^{*} \Pi+$ $C_{\circlearrowleft}^{*} C_{\circlearrowleft} \operatorname{map} \operatorname{dom}\left(A_{\circlearrowleft}\right)$ continuously into $H$.
(ii) Construct the range space $V_{(C, K)}^{*}$ in the same way as the space $V_{C}^{*}$ was constructed, except that $\Psi$ is replaced by $\Psi_{\text {ext }}$. Then $\Pi \in \mathcal{L}\left(\operatorname{dom}(A) ; V_{(C, K)}^{*}\right) \cap$ $\mathcal{L}\left(\operatorname{dom}\left(A_{\circlearrowleft}\right) ; V_{C}^{*}\right)$.
(iii)

$$
\begin{gathered}
\Pi \in \mathcal{L}\left(\operatorname{dom}(A) ; \operatorname{dom}\left(A^{*}\right)\right) \quad \text { iff } \quad C^{*} C-K^{*} E^{*} E K \in \mathcal{L}(\operatorname{dom}(A) ; H), \\
\Pi \in \mathcal{L}\left(\operatorname{dom}(A) ; \operatorname{dom}\left(A_{\circlearrowleft}^{*}\right)\right) \quad \text { iff } \quad C_{\circlearrowleft}^{*} C \in \mathcal{L}(\operatorname{dom}(A) ; H), \\
\Pi \in \mathcal{L}\left(\operatorname{dom}\left(A_{\circlearrowleft}\right) ; \operatorname{dom}\left(A^{*}\right)\right) \quad \text { iff } \quad C^{*} C_{\circlearrowleft} \in \mathcal{L}\left(\operatorname{dom}\left(A_{\circlearrowleft}\right) ; H\right), \\
\Pi \in \mathcal{L}\left(\operatorname{dom}\left(A_{\circlearrowleft}\right) ; \operatorname{dom}\left(A_{\circlearrowleft}^{*}\right)\right) \quad \text { iff } \quad C_{\circlearrowleft}^{*} C_{\circlearrowleft} \in \mathcal{L}\left(\operatorname{dom}\left(A_{\circlearrowleft}\right) ; H\right) .
\end{gathered}
$$

(iv) $\Pi \in \mathcal{L}\left(W_{B} ; V_{C, K}^{*}\right)$.

Proof. The first part follows immediately from Theorem 41. To prove parts (ii) and (iii) it suffices to choose some $\alpha$ in the resolvent set of $A^{*}$ or $A_{\circlearrowleft}^{*}$, to add $\alpha \Pi x_{0}$ to both the equations, and to solve the resulting equations for $\Pi x_{0}$. Claim (iv) is proved in [40].

## 10. The Algebraic Riccati Equation

In order to turn the Lyapunov equations developed in the previous section into algebraic Riccati equations we still need to show that in some sense $B^{*} \Pi=$ $-D^{*}(C+D K)$. Actually, as we shall see, this equation is not quite correct in the sense that in some cases we shall have to replace the feed-through operator $D$ of $\mathcal{D}$ by the feed-through operator $X$ of the outer factor $\mathcal{X}$ of $\mathcal{D}$, a phenomenon that was first discovered in [37].

To compute $B^{*} \Pi$ we study some coupled systems, where the output of the original open loop system $\Psi$ or the closed loop state feedback system $\Psi_{\circlearrowleft}$ is used as the input of the adjoint closed loop system $\Psi_{\circlearrowleft}^{*}$ or open loop system $\Psi^{*}$, respectively. The systems $\Psi$ and $\Psi_{\circlearrowleft}$ are considered in the initial value setting with initial time zero, and the adjoint systems $\Psi_{\circlearrowleft}^{*}$ and $\Psi^{*}$ are considered in the time-invariant setting on $\mathbf{R}^{+}$. We start with the connection drawn in Figure 5, which will tell us how $B^{*} \Pi$ behaves on $\operatorname{dom}(A)$.
Theorem 44. Make the same assumption and introduce the same notation as in Proposition 39. Then

$$
E K x=-\left(X^{*}\right)^{-1}\left(\bar{B}^{*} \Pi+D^{*} C\right) x, \quad x \in \operatorname{dom}(A)
$$

Proof. Let us first remark that $\bar{B}^{*} \Pi$ is well-defined, since, by Proposition 39 and Corollary $43, \bar{B}^{*} \in \mathcal{L}\left(V_{(C, K)}^{*} ; U\right)$ and $\Pi \in \mathcal{L}\left(W ; V_{(C, K)}^{*}\right)$.

Consider the connection in Figure 5. We take $u=0$ and $x_{0} \in W$, and define

$$
x(t)=\mathcal{A}(t) x_{0}, \quad y(t)=\left(\mathcal{C} x_{0}\right)(t)=C \mathcal{A}(t) x_{0}, \quad t \in \mathbf{R}^{+}
$$



Figure 5. Primal-Dual Connection with Dual Feedback

Then, by Proposition 29, $x \in C^{1}\left(\mathbf{R}^{+} ; H\right)$ and $y \in W^{1,2}\left(\mathbf{R}^{+} ; Y\right)$. We extend $y$ to an arbitrary function in $W^{1,2}(\mathbf{R} ; Y)$ (for example, define $y(t)=y(-t)$ for $t<0$ ), and define

$$
x^{*}=\mathcal{C}_{\circlearrowleft}^{*} \tau y=\mathcal{C}^{*}\left(I-\mathcal{Y} \pi_{+} \mathcal{Y}^{*}\right) \tau y, \quad u^{*}=\mathcal{D}_{\circlearrowleft}^{*} y=E^{*} \mathcal{Y}^{*} y
$$

By the same proposition, $x^{*} \in B C^{1}(\mathbf{R} ; H)$ and $u^{*} \in W^{1,2}(\mathbf{R} ; U)$. For nonnegative $t$ we have (use the anti-causality of $\mathcal{C}_{\circlearrowleft}^{*}$ and $\mathcal{Y}^{*}$ )

$$
\begin{aligned}
y(t) & =C x(t) \\
x^{*}(t) & =\mathcal{C}_{\circlearrowleft}^{*} \tau(t) \mathcal{C} x_{0}=\mathcal{C}_{\circlearrowleft}^{*} \mathcal{C} \mathcal{A}(t) x_{0}=\Pi \mathcal{A}(t) x_{0}=\Pi x(t) \\
u^{*}(t) & =E^{*}\left(\mathcal{Y}^{*} \mathcal{C} y\right)(t)=-E^{*} E(\mathcal{K} y)(t)=-E^{*} E K \mathcal{A}(t) x_{0}=-E^{*} E K x(t)
\end{aligned}
$$

By Propositions 36 and 39 , for all $t \in \mathbf{R}$,

$$
u^{*}(t)=\bar{B}_{\circlearrowleft}^{*} x^{*}(t)+D_{\circlearrowleft}^{*} y(t)=E^{*}\left(X^{*}\right)^{-1}\left(\bar{B}^{*} x^{*}(t)+D^{*} y(t)\right)
$$

For nonnegative $t$ we can combine this with the preceding equations and use the invertibility of $E^{*}$ to get

$$
E K x(t)=-\left(X^{*}\right)^{-1}\left(\bar{B}^{*} \Pi x(t)+D^{*} C\right) x(t)
$$

In particular, taking $t=0$ we find that $E K x_{0}=-\left(X^{*}\right)^{-1}\left(\bar{B}^{*} \Pi+D^{*} C\right) x_{0}$.
Corollary 45. In the regular case treated in Theorem 44, the Riccati operator satisfies the open loop algebraic Riccati Equation

$$
\begin{aligned}
& \left\langle A x_{1}, \Pi x_{0}\right\rangle_{H}+\left\langle x_{1}, \Pi A x_{0}\right\rangle_{H} \\
& \quad=-\left\langle C x_{1}, C x_{0}\right\rangle_{Y}+\left\langle\left(\bar{B}^{*} \Pi+D^{*} C\right) x_{1},\left(X^{*} X\right)^{-1}\left(\bar{B}^{*} \Pi+D^{*} C\right) x_{0}\right\rangle_{U}, \\
& \quad x_{0}, x_{1} \in \operatorname{dom}(A)
\end{aligned}
$$

Proof. Combine Theorems 41 and 44.
Theorem 44 tells us how $\bar{B}^{*} \Pi$ maps $W=\operatorname{dom}(A)$. It is an interesting fact that it is possible to define $\bar{B}^{*} \Pi$ also on $W_{\circlearrowleft}=\operatorname{dom}\left(A_{\circlearrowleft}\right)$, but that the resulting formula differs significantly from the one in Theorem 44:

Theorem 46. With the assumption and notation of Proposition 39, the Riccati operator $\Pi$ satisfies

$$
\bar{B}^{*} \Pi x+D^{*} C_{\circlearrowleft} x=\bar{B}^{*} \Pi x+D^{*}\left(\bar{C}+D X^{-1} E \bar{K}\right) x=0, \quad x \in \operatorname{dom}\left(A_{\circlearrowleft}\right)
$$

In particular, if $D^{*} D$ is invertible, then

$$
E \bar{K} x_{1}=-X\left(D^{*} D\right)^{-1}\left(\bar{B}^{*} \Pi+D^{*} \bar{C}\right) x_{1}, \quad x_{1} \in \operatorname{dom}\left(A_{\circlearrowleft}\right)
$$

In the classical case, $\operatorname{dom}(A)=\operatorname{dom}\left(A_{\circlearrowleft}\right)$ and $X^{*} X=D^{*} D$, so in that case this theorem is identical to Theorem 44. However, they differ from each other in the general case. In particular, although $X^{*} X$ is always invertible, $D^{*} D$ need not be so. See the examples in [38, Section 8] and [41].

Proof. Let us consider the connection in Figure 6. We take $u=0$ and $x_{0} \in$ $\operatorname{dom}\left(A_{\circlearrowleft}\right)$, and define

$$
x(t)=\mathcal{A}_{\circlearrowleft}(t) x_{0}, \quad y(t)=\left(\mathcal{C}_{\circlearrowleft} x_{0}\right)(t)=C_{\circlearrowleft} \mathcal{A}_{\circlearrowleft}(t) x_{0}, \quad t \in \mathbf{R}^{+} .
$$



Figure 6. Primal-Dual Connection with Primal Feedback
Then, by Proposition $29, x \in C^{1}(\mathbf{R} ; H)$ and $y \in W^{1,2}(\mathbf{R} ; Y)$, and we can extend $y$ to a function in $W^{1,2}(\mathbf{R} ; Y)$. Define

$$
x^{*}=\mathcal{C}^{*} \tau y, \quad u^{*}=\mathcal{D}^{*} y .
$$

By the same proposition, $x^{*} \in B C^{1}(\mathbf{R} ; H)$ and $u^{*} \in W^{1,2}(\mathbf{R} ; U)$. For $t \in \mathbf{R}^{+}$we have, by Proposition 39 , by the anti-causality of $\mathcal{C}^{*}$ and $\mathcal{D}^{*}=\left(\mathcal{X}^{*}\right)^{-1} \mathcal{Y}^{*}$, and by the fact that $\mathcal{Y}^{*} \mathcal{Y}=I$,

$$
\begin{aligned}
y(t) & =C_{\circlearrowleft} x(t)=\left(\bar{C}+D X^{-1} E \bar{K}\right) x(t) \\
x^{*}(t) & =\mathcal{C}^{*} \tau(t) \mathcal{C}_{\circlearrowleft} x_{0}=\mathcal{C}^{*} \mathcal{C}_{\circlearrowleft} \mathcal{A}_{\circlearrowleft}(t) x_{0}=\Pi \mathcal{A}_{\circlearrowleft}(t) x_{0}=\Pi x(t), \\
\pi_{+} u^{*} & =\pi_{+}\left(\mathcal{X}^{*}\right)^{-1} \mathcal{Y}^{*}\left(I-\mathcal{Y} \pi_{+} \mathcal{Y}^{*}\right) \mathcal{C} x_{0}=\pi_{+}\left(\mathcal{X}^{*}\right)^{-1}\left(\pi_{+} \mathcal{Y}^{*}-\left(\mathcal{Y}^{*} \mathcal{Y}\right) \pi_{+} \mathcal{Y}^{*}\right)=0
\end{aligned}
$$

By Propositions 36 and 39 , for all $t \in \mathbf{R}$,

$$
u^{*}(t)=\bar{B}^{*} x^{*}(t)+D^{*} y(t),
$$

which for nonnegative $t$ we may combine with the equations above to get

$$
0=\bar{B}^{*} \Pi x(t)+D^{*}\left(\bar{C}+D X^{-1} E \bar{K}\right) x(t)
$$

Take $t=0$ to get the claim of Theorem 46 .
The diagrams in Figures 5 and 6 are simplified versions of the more complete diagrams in [40].
Remark 47. Most of the results in this paper remain true if we throughout replace the algebra of time-invariant bounded linear operators from $L^{2}(\mathbf{R} ; U)$ into $L^{2}(\mathbf{R} ; Y)$ by some subalgebra, for example, by the algebra studied in [37]. The main exception is that spectral factorizations and inner-outer factorizations need not exist. In particular, Theorem 27 remains valid in that setting, too.

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