



Annealed central limit theorems for the Ising model on random graphs

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Abstract. The aim of this paper is to prove central limit theorems with respect to the annealed measure for the magnetization rescaled by \sqrt{N} of Ising models on random graphs. More precisely, we consider the general rank-1 inhomogeneous random graph (or generalized random graph), the 2-regular configuration model and the configuration model with degrees 1 and 2. For the generalized random graph, we first show the existence of a finite annealed inverse critical temperature $0 \leq \beta_c^{\text{an}} < \infty$ and then prove our results in the uniqueness regime, i.e., the values of inverse temperature β and external magnetic field B for which either $\beta < \beta_c^{\text{an}}$ and $B = 0$, or $\beta > 0$ and $B \neq 0$.

In the case of the configuration model, the central limit theorem holds in the whole region of the parameters β and B , because phase transitions do not exist for these systems as they are closely related to one-dimensional Ising models. Our proofs are based on explicit computations that are possible since the Ising model on the generalized random graph in the annealed setting is reduced to an inhomogeneous Curie-Weiss model, while the analysis of the configuration model with degrees only taking values 1 and 2 relies on that of the classical one-dimensional Ising model.

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1. Introduction and main results

1.1. *Motivation.*

1.1.1. *Ising models on random graphs.* The ferromagnetic *Ising model* is the most well-known example of statistical mechanics system describing cooperative behavior. Its probabilistic formulation Ellis (1985) amounts to an infinite family of random variables taking values in $\{-1, 1\}$ (so-called spins) whose joint law is given by the Boltzmann-Gibbs distribution. The properties of such families of random variables are crucially determined by the spatial structure where the spin variables are sitting. For instance, for the Ising model on \mathbb{Z}^d with nearest-neighbor interactions, the model displays a second-order phase transition for $d \geq 2$. Furthermore, the universality prediction states that the precise details of the interactions are not relevant for the near-critical behavior, so that around the critical temperature each universality class is described by a single set of critical exponents.

Besides regular lattices, in recent years much attention has been devoted to the setting in which the spin variables are placed on the vertices of *random graphs* Albert and Barabási (2002); Dembo and Montanari (2010); Dembo et al. (2014, 2013); De Sanctis and Guerra (2008); Dommers et al. (2010, 2014); Dorogovtsev et al. (2002, 2008); Leone et al. (2002); Montanari et al. (2012). Such random graphs aim to model emergent properties of complex systems consisting of many interacting agents described by a network. Several studies on empirical networks have found that two random elements of the network are typically within relatively short graph distance (the so-called *small-world* paradigm), whereas there does not exist a typical scale for the number of neighbors that a random element has (the so-called *scale-free* paradigm where degrees in the network are proposed to have a power-law distribution) Newman (2003, 2010); van der Hofstad (2014a,b). As a consequence, there are vertices with very high degree that often play an important role in the functionality of the network.

Thus, the combination of the ferromagnetic Ising model on a random graph describes situations in which single units establish macroscopic cooperative behavior in the presence of the random and complex connectivity structure described by a network. Here the playing field has two levels of randomness: firstly, the probabilistic law of the spins and, secondly, the probability distribution of the graph. So far, most of the studies have focused on the so-called (*random*) *quenched* state, in which the random graph is considered to be fixed once and for all. In this paper, we instead consider the *annealed* state: the Ising model at every time sees an average of the possible random graphs Bianconi (2012); Krasnytska et al. (2014), rather than one realization of the graph. The annealed measure is particularly relevant for applications in socio-economic systems, in which the graph dynamic models the evolution of social acquaintances, or the brain, in which graph edge rearrangements represent the evolution of synaptic connections. We explain the role of the annealed and the quenched laws in more detail in the following section.

1.1.2. *Annealing.* To understand the difference between the quenched and annealed settings, it is convenient to think of a microscopic dynamics yielding the equilibrium state. For instance, one could imagine that the spins are subject to a Glauber dynamics with a reversible Boltzmann-Gibbs distribution and the graph also has its own dynamical evolution approaching the graph's stationary distribution. In

general, these two dynamics are intertwined and both concur to determine the equilibrium state, i.e., the asymptotic value of an ergodic dynamical time average. The quenched and annealed state arise as follows:

- (a) In the *quenched state*, the changes of the graph happen on a time-scale that is infinitely longer than the time-scale over which the changes of the spin variables occur. Thus in the quenched state the graph viewed by the evolving spins is frozen. One distinguishes between the *random quenched measure*, i.e., the random Boltzmann-Gibbs distribution of a given realization of the graph, and the *averaged quenched measure*, i.e., the average of the Boltzmann-Gibbs distribution over the graph ensemble. Several thermodynamic observables (e.g., the free energy per particle, the internal energy per particle, etc.) are self-averaging, and therefore the random quenched values and their averaged quenched expectations do coincide in the thermodynamic limit. In the study of the *fluctuations* of the properly rescaled magnetization one finds a Gaussian limiting law. Interestingly, the asymptotic variances of the random quenched and averaged quenched state might be different [Giardinà et al. \(2015\)](#) due to local Gaussian fluctuations of graph properties.
- (b) In the *annealed state*, the environment seen by the spins includes all possible arrangements of the random graph. The annealed measure (defined later in (1.7)) is given by the stationary reversible measure of a Glauber spin dynamics in which the transition from a configuration σ to another configuration σ' occurs with probability

$$\frac{\mathbb{E}[e^{-\beta H(\sigma')}]}{\mathbb{E}[e^{-\beta H(\sigma)}]} \wedge 1, \quad (1.1)$$

where H is the Hamiltonian and $\mathbb{E}[\cdot]$ represents the average over the graph ensemble. The above dynamics corresponds to an extremely fast random graph dynamics in which we do not even *observe* the graph at any time, but merely see it averaged over the random graph distribution. This is equivalent to an *effective* Glauber dynamics with (annealed) Hamiltonian equal to

$$H^{\text{an}}(\sigma) = -\frac{1}{\beta} \log(\mathbb{E}[e^{-\beta H(\sigma)}]). \quad (1.2)$$

Thus, by construction, the annealed measure is necessarily non-random. We will be interested in the properties of the Gibbs measure corresponding to the dynamics in (1.1), which corresponds to the stationary or infinite-time distribution of the spins under the dynamics. While the Glauber dynamics (1.1) corresponds to infinitely fast graph dynamics compared to the spin dynamics, the stationary distribution can equally well be viewed as a dynamics where the graph and the spin evolve at equal speeds, as is the more usual viewpoint in statistical mechanics. Note that, in the definition of the annealed pressure (see (1.8)), the averages taken w.r.t. the spins and the graph are completely symmetric, which can be seen as another argument in favor of the view that the corresponding dynamics run equally fast and that the limiting measure corresponds to the average w.r.t. graph and spins alike. In this paper, we will study annealed central limit theorems for the

ferromagnetic Ising model on random graphs, in order to deduce what the effect of annealing on the macroscopic properties of the Ising model is.

The definition of the annealed measure in the context of Ising models on random graphs is thus different than in other class of problems with disorder, such as random walks in random environment [Comets et al. \(2000\)](#). In that context, annealing is rather similar to what here we have called the averaged quenched measure.

In disordered systems (such as spin glasses [Mézard et al. \(1987\)](#); [Contucci and Giardinà \(2013\)](#)), annealed disorder is usually considered to be easier to deal with mathematically, since the average on the disorder and the thermal average are treated on the same footing. This is true whenever the edges of the graph are independent, due to the form of the Hamiltonian that allows a factorization of expectations w.r.t. the bond variables. If instead the edge distribution in the graph does not have a product structure, the annealed case can actually be more difficult than the quenched case. Indeed, whereas the random-quenched case is dominated by the typical realization of the graph (often having the local structure of a random tree), in the annealed case (as in the averaged-quenched case) the rare graph samples actually give a contribution that can not be ignored. This is due to the fact that the Ising model gives rise to *exponential functionals* on the random graph, and expectations of exponential functionals tend to be dominated by rare events in which the exponential functional is larger than it would be under the quenched law. Deriving such statement rigorously requires a deep understanding of the large deviation properties of random graphs, a highly interesting but also challenging topic.

In this paper, we consider graph ensembles of both types, i.e., random graphs with independent edges (these are generalized random graphs) or dependent edges (in this case, we study the configuration model). These are described in the following section.

1.2. Random graph models. We denote by $G_N = (V_N, E_N)$ a random graph with vertex set $V_N = [N]$ and edge set $E_N \subset V_N \times V_N$. Here and in the rest of this paper, we write $[N] = \{1, \dots, N\}$ for the vertex set of G_N . For any $N \in \mathbb{N}$, we denote by Q_N the probability law of the random graph G_N . In this work, we consider two classes of random graphs: the configuration model and the generalized random graph. We next introduce these models.

1.2.1. The Generalized Random Graph. In the generalized random graph, each vertex $i \in [N]$ receives a weight $w_i > 0$. Given the weights, edges are present independently, but the occupation probabilities for different edges are *not* identical, instead, they are moderated by the vertex weights. For a given sequence of weights $\mathbf{w} = (w_i)_{i \in [N]}$, the graph is denoted by $\text{GRG}_N(\mathbf{w})$. We call I_{ij} the Bernoulli indicator that the edge between vertex i and vertex j is present and $p_{ij} = \mathbb{P}(I_{ij} = 1)$ is equal to

$$p_{ij} = \frac{w_i w_j}{\ell_N + w_i w_j}, \quad (1.3)$$

where ℓ_N is the total vertex weight given by

$$\ell_N = \sum_{i=1}^N w_i. \quad (1.4)$$

Denote by $W_N = w_{I_N}$ the weight of a uniformly chosen vertex $I_N \in [N]$. The weight sequence of the generalized random graph $\text{GRG}_N(\mathbf{w})$ is often assumed to satisfy a *regularity condition*, which is expressed as follows:

Condition 1.1 (Weight regularity). *There exists a random variable W such that, as $N \rightarrow \infty$,*

- (a) $W_N \xrightarrow{\mathcal{D}} W$,
- (b) $\mathbb{E}[W_N] = \frac{1}{N} \sum_{i \in [N]} w_i \rightarrow \mathbb{E}[W] < \infty$,
- (c) $\mathbb{E}[W_N^2] = \frac{1}{N} \sum_{i \in [N]} w_i^2 \rightarrow \mathbb{E}[W^2] < \infty$,

where $\xrightarrow{\mathcal{D}}$ denotes convergence in distribution. Further, we assume that $\mathbb{E}[W] > 0$.

In the following, we will consider deterministic sequences of weights that satisfy Condition 1.1. In many cases, one could also work with weights $\mathbf{w} = (w_i)_{i \in [N]}$ that are i.i.d. random variables. For the annealed setting, however, one has to be careful, as we will argue in more detail in Section 1.5.1 below. Indeed, when the weights are themselves random variables, they introduce a double randomness in the random graphs: firstly there is the randomness introduced by the weights, and secondly there is the randomness introduced by the edge occupation statuses, which are conditionally independent given the weights. Whereas the thermodynamic properties (pressure, magnetization, etc.) in the quenched measures of the Ising model on $\text{GRG}_N(\mathbf{w})$ are not affected by the choice of deterministic or random weight sequences, the pressure of the annealed Ising model becomes *infinite* when the weights have sufficiently heavy tails.

1.2.2. *The Configuration Model.* The configuration model is a *multigraph*, that is, a graph possibly having self-loops and multiple edges between pairs of vertices. Fix an integer N and consider a sequence of integers $\mathbf{d} = (d_i)_{i \in [N]}$. The aim is to construct an undirected multigraph with N vertices, where vertex j has degree d_j . We assume that $d_j \geq 1$ for all $j \in [N]$ and we denote the total degree in the graph ℓ_N by

$$\ell_N := \sum_{i \in [N]} d_i. \quad (1.5)$$

We assume ℓ_N to be even in order to be able to construct the graph.

Assuming that initially d_j half-edges are attached to each vertex $j \in [N]$, one way of obtaining a uniform multigraph with the given degree sequence is to pair the half-edges belonging to the different vertices in a uniform way. Two half-edges together form an edge, thus creating the edges in the graph. To construct the multigraph with degree sequence \mathbf{d} , the half-edges are numbered in an arbitrary order from 1 to ℓ_N . Then we start by randomly connecting the first half-edge with one of the $\ell_N - 1$ remaining half-edges. Once paired, two half-edges form a single edge of the multigraph. We continue the procedure of randomly choosing and pairing the half-edges until all half-edges are connected, and call the resulting graph the *configuration model with degree sequence \mathbf{d}* , abbreviated as $\text{CM}_N(\mathbf{d})$.

We will consider, in particular, the following models:

- (1) The *2-regular random graph*, i.e., the configuration model with $d_i = 2$ for all $i \in [N]$, which we denote by $\text{CM}_N(\mathbf{2})$.

- (2) The configuration model with $d_i \in \{1, 2\}$ for all $i \in [N]$, which we denote by $\text{CM}_N(\mathbf{1}, \mathbf{2})$. In $\text{CM}_N(\mathbf{1}, \mathbf{2})$, for a given $p \in [0, 1]$, we have $N - \lfloor pN \rfloor$ vertices of degree 1 and $\lfloor pN \rfloor$ vertices of degree 2.

1.2.3. *Properties of $\text{GRG}_N(\mathbf{w})$ and $\text{CM}_N(\mathbf{d})$.* The existence of a phase transition in the structural properties of the graph depends on the asymptotic degree D , i.e. the weak limit, provided it exists, of the sequence $(D_N)_{N \geq 1}$ where D_N is the degree of a uniformly chosen vertex $I_N \in [N]$ in the graph. In order to state this result, we introduce some notation that we will frequently rely upon. Let the integer-valued random variable D have distribution $P = (p_k)_{k \geq 1}$, i.e., $\mathbb{P}(D = k) = p_k$, for $k \geq 1$. We define the *size-biased law* $\rho = (\rho_k)_{k \geq 0}$ of D by

$$\rho_k = \frac{(k+1)p_{k+1}}{\mathbb{E}[D]},$$

where the expected value of D is supposed to be finite, and introduce the average value of ρ by

$$\nu := \sum_{k \geq 0} k \rho_k = \frac{\mathbb{E}[D(D-1)]}{\mathbb{E}[D]}. \quad (1.6)$$

For $\text{CM}_N(\mathbf{2})$, the asymptotic degree distribution equals $\mathbb{P}(D = 2) = 1$, while for $\text{CM}_N(\mathbf{1}, \mathbf{2})$, the asymptotic degree distribution equals $\mathbb{P}(D = 2) = p, \mathbb{P}(D = 1) = 1 - p$. For $\text{GRG}_N(\mathbf{w})$ with asymptotic weight distribution W , the asymptotic degree D is a mixed Poisson random variable $\text{Poi}(W)$ where W appears in Condition 1.1, i.e.,

$$\mathbb{P}(D = k) = \mathbb{E} \left[e^{-W} \frac{W^k}{k!} \right],$$

see e.g., [van der Hofstad \(2014a, Chapter 6\)](#).

It is well known [Bollobás et al. \(2007\)](#); [Janson and Luczak \(2009\)](#) that the above random graphs have a phase transition in their maximal component. Indeed, when $\nu > 1$ a giant component exists, while for $\nu \leq 1$ the maximal component has $o(N)$ vertices. Here, since the degree distribution for $\text{GRG}_N(\mathbf{w})$ is $D = \text{Poi}(W)$, we have that $\nu = \frac{\mathbb{E}[W^2]}{\mathbb{E}[W]}$, because $\mathbb{E}[D] = \mathbb{E}[W]$ and $\mathbb{E}[D(D-1)] = \mathbb{E}[W^2]$. Thus, depending on W , a giant component for $\text{GRG}_N(\mathbf{w})$ may exist, while it does not exist for $\text{CM}_N(\mathbf{2})$ and $\text{CM}_N(\mathbf{1}, \mathbf{2})$ since, in these cases, $\nu \leq 1$. In fact, for $\text{CM}_N(\mathbf{2})$, the connectivity structure is quite interesting and explained in more detail in [Janson and Luczak \(2009\)](#).

1.3. *Annealed measure and thermodynamic quantities.* We continue by introducing the ferromagnetic Ising model and the *annealed measure*. We define them on finite graphs with N vertices and then study asymptotic results in the limit $N \rightarrow \infty$. We denote a configuration of N spins by σ , where σ is defined on the vertices of the random graph G_N whose law is Q_N .

In our previous work [Giardinà et al. \(2015\)](#), we have considered two Ising models. The *random-quenched measure* $\mu_{G_N}(\sigma)$ coincides with the random Boltzmann–Gibbs distribution, where the randomness is given by the graph G_N . The *averaged-quenched measure* $P_N(\sigma)$ is obtained by averaging the random Boltzmann–Gibbs distribution over all possible random graphs, i.e., $P_N(\sigma) = Q_N(\mu_{G_N}(\sigma))$.

In defining the *annealed measure*, the numerator and denominator of the Boltzmann–Gibbs distribution μ_{G_N} are averaged separately with respect to Q_N , as formalized in the following definition:

Definition 1.1 (Annealed measure). *For spin variables $\sigma = (\sigma_1, \dots, \sigma_N)$ taking values on the space of spin configurations $\Omega_N = \{-1, 1\}^N$, we define the annealed measure by*

$$\tilde{P}_N(\sigma) = \frac{Q_N \left(\exp \left[\beta \sum_{(i,j) \in E_N} \sigma_i \sigma_j + B \sum_{i \in [N]} \sigma_i \right] \right)}{Q_N(Z_N(\beta, B))}, \quad (1.7)$$

where

$$Z_N(\beta, B) = \sum_{\sigma \in \Omega_N} \exp \left[\beta \sum_{(i,j) \in E_N} \sigma_i \sigma_j + B \sum_{i \in [N]} \sigma_i \right]$$

is the partition function. Here $\beta \geq 0$ is the inverse temperature and $B \in \mathbb{R}$ is the uniform external magnetic field.

In this paper, with a slight abuse of notation, we use the same symbols to denote both a measure and the corresponding expectation. Moreover, we remark that the measure defined above depends sensitively on the two parameters (β, B) . However, for the sake of notation, we will drop the dependence of the measure on these parameters. Sometimes we will use $\text{Var}_\mu(X)$ to denote the variance of a random variable X with law μ .

We now define the thermodynamic quantities with respect to the annealed measure:

Definition 1.2 (Thermodynamic quantities [Dembo and Montanari \(2010\)](#); [Dommers et al. \(2010\)](#)). *For a given $N \in \mathbb{N}$, we introduce the following thermodynamics quantities in finite volume:*

(i) *The annealed pressure is given by*

$$\tilde{\psi}_N(\beta, B) = \frac{1}{N} \log(Q_N(Z_N(\beta, B))). \quad (1.8)$$

(ii) *The annealed magnetization is given by*

$$\tilde{M}_N(\beta, B) = \tilde{P}_N \left(\frac{S_N}{N} \right),$$

where the total spin is defined as

$$S_N = \sum_{i \in [N]} \sigma_i.$$

(iii) *The annealed susceptibility equals*

$$\tilde{\chi}_N(\beta, B) := \frac{\partial}{\partial B} \tilde{M}_N(\beta, B) = \text{Var}_{\tilde{P}_N} \left(\frac{S_N}{\sqrt{N}} \right).$$

We are interested in the thermodynamic limit of these quantities, i.e., their limits as $N \rightarrow \infty$. In this limit, critical phenomena may appear. If $\mathcal{M}(\beta, B) := \lim_{N \rightarrow \infty} \mathcal{M}_N(\beta, B)$, where $\mathcal{M}_N(\beta, B)$ is the average of S_N/N with respect to $\mu_{G_N}(\cdot)$,

$P_N(\cdot)$ or $\tilde{P}_N(\cdot)$ and provided this limit exists, criticality manifests itself in the behavior of the *spontaneous magnetization* defined as $\mathcal{M}(\beta, 0^+) = \lim_{B \rightarrow 0^+} \mathcal{M}(\beta, B)$. In more detail, the *critical inverse temperature* is defined as

$$\beta_c := \inf\{\beta > 0: \mathcal{M}(\beta, 0^+) > 0\}. \quad (1.9)$$

and thus, depending on the setting, we can obtain the *quenched* and *annealed critical points* denoted by β_c^{qu} and β_c^{an} , respectively. When $0 < \beta_c < \infty$, we say that the system undergoes a *phase transition* at $\beta = \beta_c$.

From [Dommers et al. \(2014\)](#), we recall that, in the general setting of tree-like random graphs to which our models belong, the quenched critical inverse temperature is given by

$$\beta_c^{\text{qu}} = \text{atanh}(1/\nu), \quad (1.10)$$

where ν is defined in (1.6). Let us remark that, in the quenched setting, since $\nu \leq 1$ for both $\text{CM}_N(\mathbf{2})$ and $\text{CM}_N(\mathbf{1}, \mathbf{2})$, from (1.10) it follows immediately that $\beta_c^{\text{qu}} = \infty$, which means that there is no quenched phase transition in these models. In the annealed setting, we will prove the absence of phase transition for $\text{CM}_N(\mathbf{2})$ and $\text{CM}_N(\mathbf{1}, \mathbf{2})$ below. On the contrary, we will see that a critical inverse temperature appears for $\text{GRG}_N(\mathbf{w})$.

1.4. Results. We focus first on the study of the generalized random graph under the annealed measure, obtaining the Strong Law of Large Numbers (SLLN) and the Central Limit Theorem (CLT) for the total spin S_N . Then we present the results in the annealed setting for the configuration models $\text{CM}_N(\mathbf{2})$ and $\text{CM}_N(\mathbf{1}, \mathbf{2})$.

1.4.1. Results for $\text{GRG}_N(\mathbf{w})$. The proofs of the SLLN and CLT for $\text{GRG}_N(\mathbf{w})$ require to investigate the uniqueness regime for $\text{GRG}_N(\mathbf{w})$. For this, we first investigate the existence of the thermodynamic quantities in the infinite volume limit with respect to the annealed law. These results will be obtained in the next theorem. They show, in particular, that annealing changes the critical inverse temperature. Indeed, the annealed critical inverse temperature β_c^{an} is *strictly smaller* than the quenched critical inverse temperature β_c^{qu} , when the latter exists. In the statement of the theorem below, we will use the notation \mathcal{U}^{an} for the annealed uniqueness regime, i.e.,

$$\mathcal{U}^{\text{an}} := \{(\beta, B) : \beta \geq 0, B \neq 0 \text{ or } 0 < \beta < \beta_c^{\text{an}}, B = 0\}.$$

Theorem 1.1 (Thermodynamic limits for the annealed $\text{GRG}_N(\mathbf{w})$). *Let $(G_N)_{N \geq 1}$ be a sequence of $\text{GRG}_N(\mathbf{w})$ satisfying Condition 1.1. Then the following conclusions hold:*

- (i) *For all $0 \leq \beta < \infty$ and for all $B \in \mathbb{R}$, the annealed pressure exists in the thermodynamic limit $N \rightarrow \infty$ and is given by*

$$\tilde{\psi}(\beta, B) := \lim_{N \rightarrow \infty} \tilde{\psi}_N(\beta, B), \quad (1.11)$$

its value is given in (2.8).

- (ii) *For all $(\beta, B) \in \mathcal{U}^{\text{an}}$, the magnetization per vertex exists in the limit $N \rightarrow \infty$, i.e.,*

$$\tilde{M}(\beta, B) := \lim_{N \rightarrow \infty} \tilde{M}_N(\beta, B). \quad (1.12)$$

For $B \neq 0$ the limit value $\widetilde{M}(\beta, B)$ equals $\widetilde{M}(\beta, B) = \frac{\partial}{\partial B} \widetilde{\psi}(\beta, B)$ and is given by

$$\widetilde{M}(\beta, B) = \mathbb{E} \left[\tanh \left(\sqrt{\frac{\sinh(\beta)}{\mathbb{E}[W]}} W z^* + B \right) \right],$$

where $z^* = z^*(\beta, B)$ is the solution of the fixed-point equation

$$z = \mathbb{E} \left[\tanh \left(\sqrt{\frac{\sinh(\beta)}{\mathbb{E}[W]}} W z + B \right) \sqrt{\frac{\sinh(\beta)}{\mathbb{E}[W]}} W \right]$$

and W is the limiting random variable defined in Condition 1.1.

(iii) The spontaneous magnetization is given by

$$\widetilde{\mathcal{M}}(\beta) := \lim_{B \rightarrow 0^+} \widetilde{M}(\beta, B) = \begin{cases} 0 & \text{if } \beta \in \mathcal{U}^{\text{an}} \\ \neq 0 & \text{if } \beta \notin \mathcal{U}^{\text{an}} \end{cases}$$

and the annealed critical inverse temperature is

$$\beta_c^{\text{an}} = \text{asinh}(1/\nu),$$

where ν , defined in (1.6), is given by $\nu = \mathbb{E}[W^2]/\mathbb{E}[W]$ and W is the limiting random variable introduced in Condition 1.1. In particular, if $\nu > 1$, then $\beta_c^{\text{an}} < \beta_c^{\text{qu}}$.

(iv) For all $(\beta, B) \in \mathcal{U}^{\text{an}}$, the thermodynamic limit of the susceptibility exists and is given by

$$\widetilde{\chi}(\beta, B) := \lim_{N \rightarrow \infty} \widetilde{\chi}_N(\beta, B) = \frac{\partial^2}{\partial B^2} \widetilde{\psi}(\beta, B). \quad (1.13)$$

Having investigated the phase diagram of the annealed Ising model on the $\text{GRG}_N(\mathbf{w})$, we next state the SLLN and CLT for the total spin in the following two theorems:

Theorem 1.2 (Annealed SLLN). *Let $(G_N)_{N \geq 1}$ be a sequence of $\text{GRG}_N(\mathbf{w})$ graphs satisfying Condition 1.1 then, for all $(\beta, B) \in \mathcal{U}^{\text{an}}$, for any $\varepsilon > 0$ there exists a number $L = L(\varepsilon) > 0$ such that the total spin is exponentially concentrated in the form*

$$\widetilde{P}_N \left(\left| \frac{S_N}{N} - \widetilde{M} \right| \geq \varepsilon \right) \leq e^{-NL} \quad \text{for all sufficiently large } N,$$

where $\widetilde{M} = \widetilde{M}(\beta, B)$ is the annealed magnetization defined in (1.12).

Theorem 1.3 (Annealed CLT). *Let $(G_N)_{N \geq 1}$ be a sequence of $\text{GRG}_N(\mathbf{w})$ graphs satisfying Condition 1.1. Then, for all $(\beta, B) \in \mathcal{U}^{\text{an}}$, the total spin satisfies a CLT of the form*

$$\frac{S_N - \widetilde{P}_N(S_N)}{\sqrt{N}} \xrightarrow{\mathcal{D}} \mathcal{N}(0, \widetilde{\chi}), \quad \text{w.r.t. } \widetilde{P}_N, \quad \text{as } N \rightarrow \infty,$$

where $\widetilde{\chi} = \widetilde{\chi}(\beta, B)$ is the thermodynamic limit of the annealed susceptibility defined in (1.13) and $\mathcal{N}(0, \sigma^2)$ denotes a centered normal random variable with variance σ^2 .

The proofs of Theorems 1.1, 1.2 and 1.3 all heavily rely on the fact that the annealed $\text{GRG}_N(\mathbf{w})$ gives rise to an inhomogeneous Curie-Weiss model, which is interesting in its own right. We continue by studying the annealed measure on $\text{CM}_N(\mathbf{2})$.

1.4.2. *Results for $\text{CM}_N(\mathbf{2})$.* Our main result for $\text{CM}_N(\mathbf{2})$ concerns its thermodynamic limits, a SLLN and a CLT for its total spin, as formulated in the following theorems:

Theorem 1.4 (Thermodynamic limits for the annealed $\text{CM}_N(\mathbf{2})$). *Let $(G_N)_{N \geq 1}$ be a sequence of $\text{CM}_N(\mathbf{2})$ graphs. Then, for all $\beta > 0$, $B \in \mathbb{R}$, the following hold:*

(i) *The annealed pressure exists in the thermodynamic limit $N \rightarrow \infty$ and is given by*

$$\tilde{\psi}(\beta, B) := \lim_{N \rightarrow \infty} \tilde{\psi}_N(\beta, B) = \log \lambda_+(\beta, B),$$

where

$$\lambda_+(\beta, B) = e^\beta \left[\cosh(B) + \sqrt{\sinh^2(B) + e^{-4\beta}} \right].$$

(ii) *The magnetization per vertex exists in the limit $N \rightarrow \infty$ and is given by*

$$\tilde{M}(\beta, B) := \lim_{N \rightarrow \infty} \tilde{M}_N(\beta, B) = \frac{\sinh(B)}{\sqrt{\sinh^2(B) + e^{-4\beta}}}. \quad (1.14)$$

Remark 1.1. *Since $\lim_{B \rightarrow 0^+} \tilde{M}(\beta, B) = 0$ for all $\beta > 0$, by definition (1.9) we conclude that there is no annealed phase transition for $\text{CM}_N(\mathbf{2})$. This is not surprising, since $\text{CM}_N(\mathbf{2})$ consists of a collection of disjoint cycles, and the Ising model does not have a phase transition in dimension one.*

Next we state the SLLN for the total spin in $\text{CM}_N(\mathbf{2})$:

Theorem 1.5 (Annealed SLLN for $\text{CM}_N(\mathbf{2})$). *Let $(G_N)_{N \geq 1}$ be a sequence of $\text{CM}_N(\mathbf{2})$ graphs. Then, for all $\beta \geq 0$, $B \in \mathbb{R}$, for any $\varepsilon > 0$ there exists a number $L = L(\varepsilon) > 0$ such that the total spin is exponentially concentrated in the form*

$$\tilde{P}_N \left(\left| \frac{S_N}{N} - \tilde{M} \right| \geq \varepsilon \right) \leq e^{-NL} \quad \text{for all sufficiently large } N,$$

where $\tilde{M} = \tilde{M}(\beta, B)$ is the annealed magnetization defined in (1.14).

Finally, we investigate the CLT for $\text{CM}_N(\mathbf{2})$:

Theorem 1.6 (Annealed CLT for $\text{CM}_N(\mathbf{2})$). *Let $(G_N)_{N \geq 1}$ be a sequence of $\text{CM}_N(\mathbf{2})$ graphs. Then, for all $\beta \geq 0$, $B \in \mathbb{R}$, the total spin satisfies a CLT of the form*

$$\frac{S_N - \tilde{P}_N(S_N)}{\sqrt{N}} \xrightarrow{\mathcal{D}} \mathcal{N}(0, \chi), \quad \text{w.r.t. } \tilde{P}_N, \quad \text{as } N \rightarrow \infty,$$

where $\chi = \chi(\beta, B)$ is the thermodynamic limit of the quenched susceptibility (see [Giardinà et al. \(2015, Theorem 1.1\)](#)) of the Ising model on $\text{CM}_N(\mathbf{2})$. Moreover, $\chi(\beta, B)$ is also equal to the susceptibility of the one-dimensional Ising model, i.e.,

$$\chi(\beta, B) = \chi^{d=1}(\beta, B) = \frac{\cosh(B)e^{-4\beta}}{(\sinh(B) + e^{-4\beta})^{3/2}}.$$

Theorems 1.4, 1.5 and 1.6 are proved in Section 3. Their proofs heavily rely on the fact that $\text{CM}_N(\mathbf{2})$ consists of a collection of cycles, and the partition function on a cycle can be computed explicitly.

1.4.3. *Results for $\text{CM}_N(\mathbf{1}, \mathbf{2})$.* Our main result for $\text{CM}_N(\mathbf{1}, \mathbf{2})$ again concerns its thermodynamic limits, SLLN and CLT for its total spin. Some of the quantities involved in the statement of these results are defined in Section 4.

Theorem 1.7 (Thermodynamic limits for the annealed $\text{CM}_N(\mathbf{1}, \mathbf{2})$). *Let $(G_N)_{N \geq 1}$ be a sequence of $\text{CM}_N(\mathbf{1}, \mathbf{2})$ graphs for a given $p \in (0, 1)$. Then, for all $\beta > 0$, $B \in \mathbb{R}$, the following hold:*

(i) *The annealed pressure exists in the thermodynamic limit $N \rightarrow \infty$ and is given by*

$$\begin{aligned} \tilde{\psi}(\beta, B) &:= \lim_{N \rightarrow \infty} \tilde{\psi}_N(\beta, B) \\ &= \log \lambda_+(\beta, B) + \frac{1-p}{2} \log A_+(\beta, B) + H(s^*, t^*), \end{aligned} \quad (1.15)$$

where $A_+(\beta, B)$ is defined in (4.3) below,

the function $H: [0, \frac{1-p}{2}] \times [0, p] \rightarrow \mathbb{R}$ is defined in (4.25) below, and

(s^*, t^*) is the unique maximum point of $(s, t) \mapsto H(s, t)$ on $[0, \frac{1-p}{2}] \times [0, p]$.

(ii) *The magnetization per vertex exists in the limit $N \rightarrow \infty$, i.e.,*

$$\tilde{M}(\beta, B) := \lim_{N \rightarrow \infty} \tilde{M}_N(\beta, B) = \frac{\partial}{\partial B} \tilde{\psi}(\beta, B), \quad (1.16)$$

and is given in (4.46) below.

Remark 1.2. (a) *Since $\lim_{B \rightarrow 0^+} \tilde{M}(\beta, B) = 0$ for all $\beta > 0$ (the explicit expression of the magnetization is given in (4.46)), we again conclude that there is not an annealed phase transition also for $\text{CM}_N(\mathbf{1}, \mathbf{2})$. Again this is not surprising, since $\text{CM}_N(\mathbf{1}, \mathbf{2})$ consists of a collection of one-dimensional lines and cycles, and the one-dimensional Ising model does not have a phase transition.*

(b) *Broutin and de Panafieu (2014) proved a CLT for the number of lines of given lengths in $\text{CM}_N(\mathbf{1}, \mathbf{2})$. Leveraging on this result, we proved in Giardinà et al. (2015) the averaged quenched CLT for the total spin of the Ising model on $\text{CM}_N(\mathbf{1}, \mathbf{2})$. We have applied the result of Broutin and de Panafieu (2014) to compute also the annealed pressure (1.8), but obtaining a result different form (1.15). While we are able to see numerically that the two formulas agree, we have no analytic proof that they coincide.*

We next state the SLLN for the total spin in $\text{CM}_N(\mathbf{1}, \mathbf{2})$:

Theorem 1.8 (Annealed SLLN for $\text{CM}_N(\mathbf{1}, \mathbf{2})$). *Let $(G_N)_{N \geq 1}$ be a sequence of $\text{CM}_N(\mathbf{1}, \mathbf{2})$ graphs. Then, for all $\beta \geq 0, B \in \mathbb{R}$, for any $\varepsilon > 0$ there exists a number $L = L(\varepsilon) > 0$ such that the total spin is exponentially concentrated in the form*

$$\tilde{P}_N \left(\left| \frac{S_N}{N} - \tilde{M} \right| \geq \varepsilon \right) \leq e^{-NL} \quad \text{for all sufficiently large } N,$$

where $\tilde{M} = \tilde{M}(\beta, B)$ is the annealed magnetization defined in (1.16).

We finish with the annealed CLT in $\text{CM}_N(\mathbf{1}, \mathbf{2})$:

Theorem 1.9 (Annealed CLT for $\text{CM}_N(\mathbf{1}, \mathbf{2})$). *Let $(G_N)_{N \geq 1}$ be a sequence of $\text{CM}_N(\mathbf{1}, \mathbf{2})$ graphs. Then, for all $\beta > 0, B \in \mathbb{R}$,*

$$\frac{S_N - \tilde{P}_N(S_N)}{\sqrt{N}} \xrightarrow{\mathcal{D}} \mathcal{N}(0, \sigma_2^2), \quad \text{w.r.t. } \tilde{P}_N, \quad \text{as } N \rightarrow \infty,$$

where σ_2^2 is defined in (4.15) below.

Theorems 1.7, 1.8 and 1.9 are proved in Section 4.

1.5. Discussion.

1.5.1. *Properties of annealing.* From the results described above, the following general picture emerges on the effect of annealing:

- (i) First of all, in the presence of a ferromagnetic phase transition, annealing can change the critical temperature, meaning that $\beta_c^{\text{an}} < \beta_c^{\text{qu}}$. We proved this for the rank-1 inhomogeneous graph. For the configuration models with vertex degrees at most two that we have analyzed, it holds $\beta_c^{\text{an}} = \beta_c^{\text{qu}} = \infty$. We conjecture that in the general case when there is a positive proportion of vertices of degree at least 3 and $\nu > 1$ (so that there exists a giant component), an annealed positive critical temperature exists. We believe that this annealed critical temperature is strictly larger than the quenched critical temperature whenever the vertex degrees fluctuate and a positive proportion of the vertices have at least degree three.
- (ii) Furthermore, the annealed state satisfies a central limit theorem for the rescaled magnetization, as the quenched state does as proved in our previous paper [Giardinà et al. \(2015\)](#). Unfortunately, we can only prove this for certain random graph sequences, but we believe this to be true in general. The variance of the annealed CLT and the variance of the quenched CLT are different whenever the degrees are allowed to fluctuate. We showed this in the case of the generalized random graph, where they can not be ordered because the quenched and annealed critical temperatures are different and the quenched and annealed susceptibilities diverge at the critical point. For $\text{CM}_N(\mathbf{1}, \mathbf{2})$ having zero critical temperature and fluctuating degrees the variances are also different, and we believe the annealed variance to be *larger* than the quenched variance. Unfortunately, we have not been able to prove this.
- (iii) From the analysis of the $\text{CM}_N(\mathbf{2})$, we see that both the annealed critical temperature and the annealed variance are the same of their quenched counterparts. We conjecture this behavior to occur for all random regular graphs.
- (iv) In the $\text{GRG}_N(\mathbf{w})$, when the weights $(w_i)_{i \in [N]}$ are i.i.d and such such $\mathbb{P}(w_1 > w) = cw^{-(\tau-1)}(1 + o(1))$ for some $\tau > 1$, the annealed partition function satisfies

$$Q_N(Z_N(\beta, B)) = e^{\frac{\beta N^2}{2}(1+o(1))}. \quad (1.17)$$

Thus, the effect of annealing of the weights is dramatic, as the pressure becomes infinite for every $\beta > 0$. To see (1.17), we first note that the upper bound is trivial, as $H(\sigma) \leq N(N-1)/2$. Thus, it suffices to prove a matching lower bound. With K_N the complete graph on N vertices and

for $a > 0$,

$$\begin{aligned} \mathbb{P}(\text{GRG}_N(\mathbf{w}) = K_N) &= \mathbb{E}\left[\prod_{ij} p_{ij}\right] \geq \mathbb{P}(w_i \in [N^a, 2N^a] \forall i \in [N]) \\ &\quad \times \mathbb{E}\left[\prod_{ij} p_{ij} \mid w_i \in [N^a, 2N^a] \forall i \in [N]\right]. \end{aligned}$$

We analyze both terms separately. Firstly, since the weights are i.i.d.,

$$\mathbb{P}(w_i \in [N^a, 2N^a] \forall i \in [N]) = \mathbb{P}(w_1 \in [N^a, 2N^a])^N \geq \left(cN^{-a(\tau-1)}\right)^N = e^{o(N^2)}.$$

Secondly, when $w_i \in [N^a, 2N^a]$ for every i , there exists $b > 0$ such that

$$p_{ij} \geq 1 - bN^{1-a}.$$

Therefore,

$$\mathbb{E}\left[\prod_{ij} p_{ij} \mid w_i \in [N^a, 2N^a] \forall i \in [N]\right] \geq \left(1 - bN^{1-a}\right)^{N(N-1)/2} = 1 - o(1),$$

when $a > 3$. Thus,

$$Q_N(Z_N(\beta, B)) \geq Z_N^{K_N}(\beta, B)e^{o(N^2)} \geq e^{\frac{\beta N^2}{2}(1+o(1))},$$

where $Z_N^{K_N}(\beta, B)$ is the partition function on the complete graph. This proves the claim.

We will expand on the analysis of the annealed critical behavior of Ising models on generalized random graphs in a forthcoming paper [Dommers et al. \(2016\)](#), where we study critical exponents around the annealed critical temperature and we derive non-classical asymptotic laws *at* criticality.

1.5.2. CLT proof strategy. By applying a commonly used strategy [Ellis \(1985\)](#), we can prove CLTs for $(S_N)_{N \geq 1}$ by showing that the moment generating function of the rescaled total spin $V_N = \frac{S_N - \mathbb{E}(S_N)}{\sqrt{N}}$, converges in a neighborhood of $t = 0$ to the moment generating function of a centered Gaussian random variable. The convergence can be achieved by considering the so-called scaled *cumulant generating functions* of S_N , defined as

$$c_N(t) = \frac{1}{N} \log \mathbb{E}[\exp(tS_N)], \quad (1.18)$$

and by proving the convergence of the sequence $(c_N''(t_N))_{N \geq 1}$ for $t_N = o(1)$ to a finite value χ , which turns out to be the variance of the normal limit. This strategy has been followed in the quenched setting in [Giardinà et al. \(2015\)](#) where, specializing \mathbb{E} to the relevant measures, the CLT was proved for the Ising model on the whole class of locally tree-like random graphs in the random quenched setting, and for the $\text{CM}_N(\mathbf{2})$ and $\text{CM}_N(\mathbf{1}, \mathbf{2})$ models in the averaged quenched setting. In the former case, the limit $c(t) := \lim_{N \rightarrow \infty} c_N(t)$ can be established as a simple consequence of the existence of the random quenched pressure on locally tree-like graphs, while the convergence of $(c_N''(t_N))_{N \geq 1}$ follows from the concavity of the first derivatives of the cumulant generating functions. In the random quenched setting, this in turn is a consequence of the GHS inequality, which holds for the ferromagnetic Boltzmann-Gibbs measure μ_{G_N} . On the other hand, under the averaged quenched measure this derivative can not be expressed in terms of the averaged quenched magnetization

to exploit the GHS inequality. Because of that, only the $\text{CM}_N(\mathbf{2})$ and $\text{CM}_N(\mathbf{1}, \mathbf{2})$ setting have been treated in [Giardinà et al. \(2015\)](#) explicitly, by exploiting the structure of the graphs and connecting these systems to the one-dimensional Ising model.

A similar scenario is found in this paper, where the approach to the proof of the CLT described above is applied to the annealed setting, i.e., with (1.18) replaced by the *annealed* cumulant generating function

$$\tilde{c}_N(t) = \frac{1}{N} \log \tilde{P}_N [\exp(tS_N)]$$

that can be connected to the annealed pressure, since $\tilde{c}_N(t) = \tilde{\psi}_N(\beta, B + t) - \tilde{\psi}_N(\beta, B)$, see (1.8).

We will show by an explicit computation that the annealed pressure of $\text{GRG}_N(\mathbf{w})$ coincides with that of an inhomogeneous Curie-Weiss model. From this fact, the thermodynamic limit of the annealed pressure, magnetization and susceptibility can be obtained. This again relies on the GHS inequality that is valid also for this inhomogeneous ferromagnetic system. Thus, for the generalized random graph, the annealed CLT can be proven in a similar way as for the random quenched measure.

On the other hand, the proofs of the CLT for the configuration models do *not* follow from the abstract argument based on the GHS inequality, since GHS is not available in the general annealed context. Because of that, we have to explicitly control the limit $(\tilde{c}_N''(t_N))_{N \geq 1}$ throughout the computation of the annealed pressure. It is relatively simple to accomplish this task in the case of the regular $\text{CM}_N(\mathbf{2})$ graph consisting of cycles only. The fluctuating degree of $\text{CM}_N(\mathbf{1}, \mathbf{2})$ makes the computation of the pressure and of the limit $(\tilde{c}_N''(t_N))_{N \geq 1}$ much more involved. $\text{CM}_N(\mathbf{1}, \mathbf{2})$ consists of both lines and cycles. While the cycles give a vanishing contribution to the thermodynamic limit, the distribution of the length of the lines has to be carefully analyzed and its Gaussian fluctuations appear in the CLT for the total spin.

1.5.3. Paper organization. The rest of the paper is organized as follows. In Section 2 we deal with $\text{GRG}_N(\mathbf{w})$ for which we compute the pressure and magnetization in the thermodynamic limit, identify the critical temperature and then prove the SLLN and CLT. All of these results rely on the fact that the Ising model on $\text{GRG}_N(\mathbf{w})$ in the annealed setting turns into an inhomogeneous Curie-Weiss model. The pressures and CLTs for the 2-regular configuration model are considered in Section 3 and for the configuration model with vertex degrees 1 and 2 in Section 4. In the former case, we show that the variance of the limiting normal variable is the susceptibility of the one-dimensional Ising model. In the latter case, which is much more difficult, the varying degrees of the vertices affect the pressure and the limiting distribution. In fact, the limiting variance is the sum of that of the one-dimensional Ising model and of an extra term emerging from the fluctuations of the connected structures of the graph.

2. Proofs for $\text{GRG}_N(\mathbf{w})$

In this section, we derive our results for the generalized random graph $\text{GRG}_N(\mathbf{w})$ stated in Theorems 1.1, 1.2 and 1.3.

2.1. *Annealed thermodynamic limits: Proof of Theorem 1.1.* The proof is divided into several steps.

Annealed partition function. We start by analyzing the average of the partition function for $\text{GRG}_N(\mathbf{w})$. By remembering that in this random graph the edges are independent and denoting by I_{ij} the Bernoulli indicator that the edge between vertex i and vertex j is present, we compute

$$\begin{aligned} Q_N(Z_N(\beta, B)) &= Q_N\left(\sum_{\sigma \in \Omega_N} \exp\left[\beta \sum_{i < j} I_{ij} \sigma_i \sigma_j + B \sum_{i \in [N]} \sigma_i\right]\right) \\ &= \sum_{\sigma \in \Omega_N} e^{B \sum_{i \in [N]} \sigma_i} Q_N\left(e^{\beta \sum_{i < j} I_{ij} \sigma_i \sigma_j}\right) \\ &= \sum_{\sigma \in \Omega_N} e^{B \sum_{i \in [N]} \sigma_i} \prod_{i < j} Q_N(e^{\beta I_{ij} \sigma_i \sigma_j}) \\ &= \sum_{\sigma \in \Omega_N} e^{B \sum_{i \in [N]} \sigma_i} \prod_{i < j} (e^{\beta \sigma_i \sigma_j} p_{ij} + (1 - p_{ij})). \end{aligned}$$

We rewrite

$$e^{\beta \sigma_i \sigma_j} p_{ij} + (1 - p_{ij}) = C_{ij} e^{\beta_{ij} \sigma_i \sigma_j},$$

where β_{ij} and C_{ij} are chosen such that

$$e^{-\beta} p_{ij} + (1 - p_{ij}) = C_{ij} e^{-\beta_{ij}}, \quad \text{and} \quad e^{\beta} p_{ij} + (1 - p_{ij}) = C_{ij} e^{\beta_{ij}}.$$

Now, by adding and dividing the two equations of the system above, we get

$$C_{ij} \cosh(\beta_{ij}) = p_{ij} \cosh(\beta) + (1 - p_{ij}), \quad \beta_{ij} = \frac{1}{2} \log \frac{e^{\beta} p_{ij} + (1 - p_{ij})}{e^{-\beta} p_{ij} + (1 - p_{ij})}.$$

Then, using the symmetry $\beta_{ij} = \beta_{ji}$ we arrive at

$$\begin{aligned} Q_N(Z_N(\beta, B)) &= \left(\prod_{i < j} C_{ij}\right) \sum_{\sigma \in \Omega_N} e^{B \sum_{i \in [N]} \sigma_i} e^{\sum_{i < j} \beta_{ij} \sigma_i \sigma_j} \\ &= G(\beta) G_1(\beta) \sum_{\sigma \in \Omega_N} e^{B \sum_{i \in [N]} \sigma_i} e^{\frac{1}{2} \sum_{i, j \in [N]} \beta_{ij} \sigma_i \sigma_j}, \end{aligned} \quad (2.1)$$

where $G(\beta) = \prod_{i < j} C_{ij}$ and $G_1(\beta) = \prod_{i \in [N]} e^{-\beta_{ii}/2}$ and we write $p_{ii} = w_i^2 / (\ell_N + w_i^2)$. This is the starting point of our analysis. We can recognize the r.h.s. as the partition function of an inhomogeneous Ising model on the complete graph, where the coupling constant between vertices i and j is equal to β_{ij} . In the next step, we analyze this partition function in detail.

Towards an inhomogeneous Curie-Weiss model. We continue by showing that β_{ij} is close to factorizing into a contribution due to i and to j . For this, by a Taylor expansion of $x \mapsto \log(1 + x)$,

$$\begin{aligned} \beta_{ij} &= \frac{1}{2} \log(1 + p_{ij}(e^{\beta} - 1)) - \frac{1}{2} \log(1 + p_{ij}(e^{-\beta} - 1)) \\ &= \frac{1}{2} p_{ij}(e^{\beta} - 1) - \frac{1}{2} p_{ij}(e^{-\beta} - 1) + O(p_{ij}^2) = \sinh(\beta) p_{ij} + O(p_{ij}^2). \end{aligned}$$

Then,

$$\begin{aligned} Q_N(Z_N(\beta, B)) & \quad (2.2) \\ &= G_2(\beta) \sum_{\sigma \in \Omega_N} e^{B \sum_{i \in [N]} \sigma_i} e^{\frac{1}{2} \sinh(\beta) \sum_{i, j \in [N]} p_{ij} \sigma_i \sigma_j + O(\sum_{i, j \in [N]} p_{ij}^2 \sigma_i \sigma_j)}. \end{aligned}$$

where $G_2(\beta) = G(\beta)G_1(\beta)$. To control the error in the exponent, we use $p_{ij} \leq w_i w_j / \ell_N$ and the assumptions in Condition 1.1, to obtain

$$\left| \sum_{i, j \in [N]} p_{ij}^2 \sigma_i \sigma_j \right| \leq \sum_{i, j \in [N]} \left(\frac{w_i w_j}{\ell_N} \right)^2 = \left(\frac{\sum_{i \in [N]} w_i^2}{\ell_N} \right)^2 = o(N).$$

Then,

$$\begin{aligned} Q_N(Z_N(\beta, B)) &= G_2(\beta) e^{o(N)} \sum_{\sigma \in \Omega_N} e^{B \sum_{i \in [N]} \sigma_i} e^{\frac{1}{2} \sinh(\beta) \sum_{i, j \in [N]} \frac{w_i w_j}{\ell_N} \sigma_i \sigma_j} \\ &= G_2(\beta) e^{o(N)} \sum_{\sigma \in \Omega_N} e^{B \sum_{i \in [N]} \sigma_i} e^{\frac{1}{2} \frac{\sinh(\beta)}{\ell_N} (\sum_{i \in [N]} w_i \sigma_i)^2}. \end{aligned}$$

When $w_i \equiv w$ for all i , so that $\text{GRG}_N(\mathbf{w})$ is the Erdős-Rényi random graph, we retrieve the Curie-Weiss model at inverse temperature $\beta' = \sinh(\beta)w$. In our inhomogeneous setting, we obtain an inhomogeneous Curie-Weiss model that we will analyze next.

Analysis of the inhomogeneous Curie-Weiss model. We use the Hubbard-Stratonovich identity, i.e., we write $e^{t^2/2} = \mathbb{E}[e^{tZ}]$, with Z standard Gaussian. Then, we find

$$\begin{aligned} Q_N(Z_N(\beta, B)) &= G_2(\beta) e^{o(N)} \sum_{\sigma \in \Omega_N} e^{B \sum_{i \in [N]} \sigma_i} \mathbb{E} \left[e^{\sqrt{\frac{\sinh(\beta)}{\ell_N}} (\sum_{i \in [N]} w_i \sigma_i) Z} \right] \\ &= G_2(\beta) e^{o(N)} 2^N \mathbb{E} \left[\prod_{i=1}^N \cosh \left(\sqrt{\frac{\sinh(\beta)}{\ell_N}} w_i Z + B \right) \right] \\ &= G_2(\beta) e^{o(N)} 2^N \mathbb{E} \left[\exp \left\{ \sum_{i=1}^N \log \cosh \left(\sqrt{\frac{\sinh(\beta)}{\ell_N}} w_i Z + B \right) \right\} \right]. \end{aligned}$$

We rewrite the sum in the exponential, using the fact that $W_N = w_{I_N}$, where we recall that I_N is a uniform vertex in $[N]$, to obtain

$$\begin{aligned} Q_N(Z_N(\beta, B)) &= G_2(\beta) e^{o(N)} 2^N \mathbb{E} \left[\exp \left\{ N \mathbb{E} \left[\log \cosh \left(\sqrt{\frac{\sinh(\beta)}{N \mathbb{E}[W_N]}} W_N Z + B \right) \middle| Z \right] \right\} \right] \\ &= G_2(\beta) e^{o(N)} 2^N \mathbb{E} \left[e^{N F_N \left(\frac{Z}{\sqrt{N}} \right)} \right], \end{aligned}$$

where

$$F_N(z) = \mathbb{E} \left[\log \cosh \left(\sqrt{\frac{\sinh(\beta)}{\mathbb{E}[W_N]}} W_N z + B \right) \right]. \quad (2.3)$$

Here we emphasize the fact that in (2.3), the expectation is w.r.t. W_N only. We continue by analyzing $F_N(z)$. We claim that, uniformly for $|z| \leq a$ and any $a < \infty$,

$$\sup_{|z| \leq a} |F_N(z) - F(z)| = o(1), \quad (2.4)$$

where

$$F(z) = \mathbb{E} \left[\log \cosh \left(\sqrt{\frac{\sinh(\beta)}{\mathbb{E}[W]}} Wz + B \right) \right].$$

To see (2.4), we note that $F_N(z) \rightarrow F(z)$ for every z fixed by Condition 1.1(a)-(b), and the fact that $\log \cosh(x) \leq |x|$. Further,

$$|F'_N(z)| \leq \frac{\sinh(\beta)}{\mathbb{E}[W_N]} \mathbb{E} \left[\tanh \left(\sqrt{\frac{\sinh(\beta)}{\mathbb{E}[W_N]}} W_N z + B \right) W_N \right] \leq \sinh(\beta),$$

since $\tanh(x) \leq 1$ for all x , so that $|F'_N(z)|$ is uniformly bounded in N and z . Therefore, $(F_N)_{N \geq 1}$ forms a uniformly equicontinuous family of functions, so that (2.4) follows from Arzelà-Ascoli. Since $F_N(z) \leq \sinh(\beta)|z|$, it further follows that, for $a > 4 \sinh(\beta)$,

$$\begin{aligned} \mathbb{E} \left[e^{NF_N\left(\frac{z}{\sqrt{N}}\right)} \mathbb{1}_{\{|Z| > a\sqrt{N}\}} \right] &\leq \mathbb{E} \left[e^{\sqrt{N} \sinh(\beta) |Z|} \mathbb{1}_{\{|Z| > a\sqrt{N}\}} \right] \\ &= 2 \mathbb{E} \left[e^{\sqrt{N} \sinh(\beta) Z} \mathbb{1}_{\{Z > a\sqrt{N}\}} \right] \\ &= \frac{2}{\sqrt{2\pi}} \int_{a\sqrt{N}}^{\infty} e^{\sqrt{N} \sinh(\beta) z} e^{-z^2/2} dz \\ &\leq e^{a \sinh(\beta) N - a^2 N/2} \int_0^{\infty} e^{\sqrt{N}(\sinh(\beta) - a)x} dx \leq e^{-a^2 N/4}, \end{aligned}$$

which, for a sufficiently large, is negligible compared to $\mathbb{E} \left[e^{NF_N\left(\frac{z}{\sqrt{N}}\right)} \mathbb{1}_{\{|Z| \leq a\sqrt{N}\}} \right]$.

We conclude that

$$Q_N(Z_N(\beta, B)) = G_2(\beta) e^{o(N)} 2^N \mathbb{E} \left[e^{NF\left(\frac{Z}{\sqrt{N}}\right)} \right] (1 + o(1)). \quad (2.5)$$

A large deviation analysis. The expectation in (2.5) is an expectation of an exponential functional, to which we apply large deviation machinery. The Gaussian variable Z/\sqrt{N} satisfies a large deviation principle with rate function $I(z) = z^2/2$ and speed N , because $Z/\sqrt{N} \stackrel{d}{=} \frac{1}{N} (Z_1 + \dots + Z_N)$, where $(Z_i)_{i \in [N]}$ are i.i.d. standard Gaussian variables. Using Varadhan's Lemma and the fact that $z \mapsto F(z)$ is continuous, we calculate the thermodynamic limit of the pressure as

$$\begin{aligned} \lim_{N \rightarrow \infty} \frac{1}{N} \log Q_N(Z_N(\beta, B)) &= \log 2 + \lim_{N \rightarrow \infty} \frac{1}{N} \log G_2(\beta) + \sup_z [F(z) - I(z)] \\ &= \log 2 + \alpha(\beta) \\ &\quad + \sup_z \left[\mathbb{E} \left[\log \cosh \left(\sqrt{\frac{\sinh(\beta)}{\mathbb{E}[W]}} Wz + B \right) \right] - \frac{z^2}{2} \right]. \end{aligned} \quad (2.6)$$

where $\alpha(\beta) = \lim_{N \rightarrow \infty} \frac{1}{N} \log G_2(\beta)$. The equation that defines the supremum is

$$z^* = z^*(\beta, B) = \mathbb{E} \left[\tanh \left(\sqrt{\frac{\sinh(\beta)}{\mathbb{E}[W]}} Wz^* + B \right) \sqrt{\frac{\sinh(\beta)}{\mathbb{E}[W]}} W \right], \quad (2.7)$$

and the annealed pressure is obtained by substituting the supremum point z^* in the right hand side of (2.6) as

$$\tilde{\psi}(\beta, B) = \log 2 + \alpha(\beta) + \mathbb{E} \left[\log \cosh \left(\sqrt{\frac{\sinh(\beta)}{\mathbb{E}[W]}} W z^*(\beta, B) + B \right) \right] - z^*(\beta, B)^2 / 2. \quad (2.8)$$

This completes the proof of Theorem 1.1(i).

The critical inverse temperature. To identify β_c^{an} as stated in Theorem 1.1(ii), we evaluate (2.7) when $B \searrow 0$ to obtain

$$z^* = H(z^*) \quad \text{where} \quad H(z) = \mathbb{E} \left[\tanh \left(\sqrt{\frac{\sinh(\beta)}{\mathbb{E}[W]}} W z \right) \sqrt{\frac{\sinh(\beta)}{\mathbb{E}[W]}} W \right]. \quad (2.9)$$

We investigate the solutions of $z^* = H(z^*)$ in (2.9). We note that $z \mapsto H(z)$ is an increasing and concave function in $[0, \infty)$. When $H'(0) > 1$, we have three solutions of (2.9), i.e., $\pm z^*$ and 0, where $z^* = z^*(\beta, 0^+) > 0$. When $H'(0) \leq 1$, instead, $z^* = 0$ is the only solution. This leads us to compute that

$$H'(0) = \sinh(\beta) \frac{\mathbb{E}[W^2]}{\mathbb{E}[W]} = \sinh(\beta) \nu.$$

Thus, the annealed critical temperature β_c^{an} satisfies $\sinh(\beta_c^{\text{an}}) = 1/\nu$. Since $\tanh(\beta_c^{\text{qu}}) = 1/\nu$, and $\tanh(x) < \sinh(x) \forall x > 0$, we obtain $\beta_c^{\text{qu}} > \beta_c^{\text{an}}$, unless when $\nu = \infty$, in which case $\beta_c^{\text{an}} = \beta_c^{\text{qu}} = 0$.

Thermodynamic limit of the magnetization. To prove the existence of the magnetization in the thermodynamic limit stated in Theorem 1.1(ii), we follow the strategy used in Dommers et al. (2010). We use the following lemma:

Lemma 2.1. *Let $(f_n)_{n \geq 1}$ be a sequence of functions that are twice differentiable in x . Assume that*

- (a) $\lim_{n \rightarrow \infty} f_n(x) = f(x)$ for some function $y \mapsto f(y)$ that is differentiable in x ;
- (b) $\frac{d}{dx} f_n(x)$ is monotone in $[x-h, x+h]$ for all $n \geq 1$ and some $h > 0$.

Then,

$$\lim_{n \rightarrow \infty} \frac{d}{dx} f_n(x) = \frac{d}{dx} f(x).$$

We apply Lemma 2.1 with $n = N$ and f_n equal to $B \mapsto \tilde{\psi}_N(\beta, B)$. We verify the conditions in Lemma 2.1 and start by noting that

$$\tilde{M}_N(\beta, B) = \tilde{P}_N(S_N/N) = \frac{\partial}{\partial B} \tilde{\psi}_N(\beta, B),$$

and $\lim_{N \rightarrow \infty} \tilde{\psi}_N(\beta, B) = \tilde{\psi}(\beta, B)$ by Theorem 1.1(i) with $B \mapsto \tilde{M}_N(\beta, B)$ non-decreasing:

$$\frac{\partial}{\partial B} \tilde{M}_N(\beta, B) = \frac{1}{N} \left[\tilde{P}_N(S_N^2) - \tilde{P}_N(S_N)^2 \right] \geq 0.$$

Thus, we can indeed conclude that

$$\tilde{M}(\beta, B) = \lim_{N \rightarrow \infty} \tilde{M}_N(\beta, B) = \lim_{N \rightarrow \infty} \frac{\partial}{\partial B} \tilde{\psi}_N(\beta, B) = \frac{\partial}{\partial B} \tilde{\psi}(\beta, B).$$

The limit magnetization $\widetilde{M}(\beta, B)$ can be explicitly computed by taking the derivative of $\widetilde{\psi}(\beta, B)$, (2.8) and using the fixed point equation (2.7), to obtain

$$\widetilde{M}(\beta, B) = \mathbb{E} \left[\tanh \left(\sqrt{\frac{\sinh(\beta)}{\mathbb{E}[W]}} W z^* + B \right) \right].$$

Thermodynamic limit of the susceptibility. Finally, the thermodynamic limit of the susceptibility in Theorem 1.1(iv) is proved using Lemma 2.1 by combining Theorem 1.1(ii) and the fact that $B \mapsto \frac{\partial}{\partial B} \widetilde{M}_N(\beta, B)$ is non-increasing by the GHS inequality. Indeed, by the explicit computation in (2.1), we see that the annealed partition function can be viewed as the partition function of an inhomogeneous Curie-Weiss model, where the field is homogeneous and the coupling constants depend on the edges. Since such an inhomogeneous Ising model also satisfies the GHS inequality, the same follows for the annealed partition function for $\text{GRG}_N(\mathbf{w})$. Therefore,

$$\frac{\partial^2}{\partial B^2} \widetilde{M}_N(\beta, B) = \frac{1}{N} \sum_{i \in [N]} \frac{\partial^2}{\partial B^2} \widetilde{P}_N(\sigma_i) \leq 0.$$

□

2.2. Annealed SLLN and CLT: Proofs of Theorems 1.2 and 1.3. With Theorem 1.1 in hand, we now have all the hypotheses to prove Theorems 1.2 and 1.3 following the strategy used for the random quenched setting in Sections 2.2 and 2.3 of Giardinà et al. (2015) verbatim. Indeed, for the proof of the annealed SLLN, referring to Giardinà et al. (2015, Section 2.2), we obtain the existence of the thermodynamic limit of the annealed cumulant generating function

$$\widetilde{c}_N(t) = \frac{1}{N} \log \widetilde{P}_N[\exp(tS_N)] = \widetilde{\psi}_N(\beta, B + t) - \widetilde{\psi}_N(\beta, B)$$

by Theorem 1.1(i). Then, from Ellis (1985, Theorem II.6.3) and Theorem 1.1(ii) we conclude the proof.

To prove the annealed CLT (see Giardinà et al. (2015, Section 2.3) for the proof of the random quenched CLT) we need the existence in the thermodynamic limit of pressure, magnetization and susceptibility given by Theorem 1.1 together with the GHS inequality that is still true in the annealed setting thanks to the mapping to the inhomogeneous Curie-Weiss model.

3. Proofs for $\text{CM}_N(\mathbf{2})$

In this section we prove the CLT with respect to the annealed measure for the 2-regular random graph. We start by computing the annealed pressure using the partition functions for the one-dimensional Ising model with periodic boundary conditions.

3.1. Annealed thermodynamic limits and SLLN: Proofs of Theorems 1.4 and 1.5. From our previous paper Giardinà et al. (2015), we remember that any 2-regular random graph is formed by cycles only. Thus, as in Giardinà et al. (2015), denoting the random number of cycles in the graph by K_N^t , we can enumerate them in an

arbitrary order from 1 to K_N^t and call $L_N(i)$ the length (i.e., the number of vertices) of the i th cycle. The random variable K_N^t has distribution given by

$$K_N^t = \sum_{j=1}^N I_j, \quad (3.1)$$

where I_j are independent Bernoulli variables given by

$$I_j = \text{Bern} \left(\frac{1}{2N - 2j + 1} \right). \quad (3.2)$$

See [Giardinà et al. \(2015\)](#) for a proof of this fact. Since the random graph splits into (disjoint) cycles, its quenched partition function factorizes into the product of the partition functions of each cycle. Therefore,

$$Z_N(\beta, B) = \prod_{i=1}^{K_N^t} Z_{L_N(i)}^{(t)}(\beta, B). \quad (3.3)$$

By [Giardinà et al. \(2015, Section 3.1\)](#), we have that the partition function of the one-dimensional Ising model with periodic boundary conditions $Z_N^{(t)}$ is given by

$$Z_N^{(t)}(\beta, B) = \lambda_+^N(\beta, B) + \lambda_-^N(\beta, B), \quad (3.4)$$

where

$$\lambda_{\pm}(\beta, B) = e^{\beta} \left[\cosh(B) \pm \sqrt{\sinh^2(B) + e^{-4\beta}} \right], \quad (3.5)$$

so we can write

$$Z_{L_N(i)}^{(t)}(\beta, B) = \lambda_+^{L_N(i)}(\beta, B) + \lambda_-^{L_N(i)}(\beta, B).$$

Because $\beta > 0$, we have $0 < \lambda_-(\beta, B) < \lambda_+(\beta, B)$, so that, for every i ,

$$\lambda_+^{L_N(i)}(\beta, B) \leq Z_{L_N(i)}^{(t)}(\beta, B) \leq 2\lambda_+^{L_N(i)}(\beta, B).$$

As a result, we can bound the the pressure as follows:

$$\prod_{i=1}^{K_N^t} \lambda_+^{L_N(i)}(\beta, B) \leq \prod_{i=1}^{K_N^t} Z_{L_N(i)}^{(t)}(\beta, B) \leq \prod_{i=1}^{K_N^t} 2\lambda_+^{L_N(i)}(\beta, B),$$

and, since $\sum_{i=1}^{K_N^t} L_N(i) = N$, we finally obtain

$$\lambda_+^N(\beta, B) \leq Z_N(\beta, B) \leq 2^{K_N^t} \lambda_+^N(\beta, B). \quad (3.6)$$

The thermodynamic limit of the annealed pressure $\tilde{\psi}_N(\beta, B)$, defined in (1.8), can be computed along the same lines of the averaged quenched one in [Giardinà et al. \(2015\)](#). Indeed, by applying the monotone operator $N^{-1} \log(Q_N(\cdot))$ to (3.6) and using the fact that $\lambda_+(\beta, B)$ is non random, we obtain

$$\log \lambda_+(\beta, B) \leq \tilde{\psi}_N(\beta, B) \leq \frac{1}{N} \log \left(Q_N \left(2^{K_N^t} \right) \right) + \log \lambda_+(\beta, B).$$

Now using the fact

$$\begin{aligned}
 \frac{1}{N} \log \left(Q_N \left(2^{K_N^t} \right) \right) &= \frac{1}{N} \log \prod_{i=1}^N Q_N \left(2^{I_i} \right) \\
 &= \frac{1}{N} \log \prod_{i=1}^N \left[\frac{2}{2N - 2i + 1} + \left(1 - \frac{1}{2N - 2i + 1} \right) \right] \\
 &= \frac{1}{N} \sum_{i=1}^N \log \left(1 + \frac{1}{2N - 2i + 1} \right) \xrightarrow{N \rightarrow \infty} 0,
 \end{aligned} \tag{3.7}$$

we conclude that the annealed pressure of $\text{CM}_N(\mathbf{2})$ coincides with the pressure of the one-dimensional Ising model $\psi^{d=1}(\beta, B)$, i.e.,

$$\tilde{\psi}(\beta, B) = \psi^{d=1}(\beta, B) \equiv \log \lambda_+(\beta, B). \tag{3.8}$$

Moreover, it also agrees with the averaged and random quenched pressures [Giardinà et al. \(2015\)](#), i.e.,

$$\tilde{\psi}(\beta, B) = \bar{\psi}(\beta, B) = \psi(\beta, B),$$

where

$$\bar{\psi}(\beta, B) := \lim_{N \rightarrow \infty} \frac{1}{N} Q_N(\ln Z_N(\beta, B)) \quad \text{and} \quad \psi(\beta, B) := \lim_{N \rightarrow \infty} \frac{1}{N} \ln Z_N(\beta, B).$$

It also straightforwardly follows that the annealed cumulant generating function of $\text{CM}_N(\mathbf{2})$ coincides with the random and averaged quenched ones [Giardinà et al. \(2015\)](#) i.e.,

$$\tilde{c}(t) = \bar{c}(t) = c(t) = \log \lambda_+(\beta, B + t) - \log \lambda_+(\beta, B). \tag{3.9}$$

The existence of the magnetization in the thermodynamic limit (Theorem 1.4(ii)) can be proved, as in the previous section, using Lemma 2.1 and the existence of the thermodynamic limit of the pressure (3.8), so we obtain

$$\tilde{M}(\beta, B) = \frac{\partial}{\partial B} \tilde{\psi}(\beta, B) = \frac{\sinh(B)}{\sqrt{\sinh^2(B) + e^{-4\beta}}},$$

as required. \square

Proof of Theorem 1.5: The proof, as for Theorem 1.2, follows immediately from the existence of the annealed pressure in the thermodynamic limit and its differentiability with respect to B . See also [Giardinà et al. \(2015, Section 2.2\)](#). \square

3.2. Annealed CLT: Proof of Theorem 1.6. To prove the CLT in the annealed setting, we follow the strategy used in [Giardinà et al. \(2015\)](#) for the averaged quenched CLT.

Rewrite in terms of cumulant generating functions. Using the annealed cumulant generating function and using a Taylor expansion, we write

$$\log \tilde{P}_N \left[\exp \left(\frac{tS_N - t\tilde{P}_N(S_N)}{\sqrt{N}} \right) \right] = \frac{t^2}{2} \tilde{c}_N''(t_N), \tag{3.10}$$

where $t_N \in [0, t/\sqrt{N}]$. Then the aim is to prove that $\lim_{N \rightarrow \infty} \tilde{c}_N''(t_N)$ exists as a finite limit.

By expressing $\tilde{c}_N(t)$ in terms of $Z_N^{(t)} = \lambda_+^N + \lambda_-^N$ and using (3.9), we can compute the difference as

$$\tilde{c}_N(t) - \tilde{c}(t) = \frac{1}{N} \log \left[\frac{Q_N \left(\frac{Z_N^{(t)}(\beta, B+t)}{(\lambda_+(\beta, B+t))^N} \right)}{Q_N \left(\frac{Z_N^{(t)}(\beta, B)}{(\lambda_+(\beta, B))^N} \right)} \right] = \frac{1}{N} \log \left[\frac{Q_N \left(\prod_{i=1}^{K_N^t} (1 + (r_{B+t})^{L_N(i)}) \right)}{Q_N \left(\prod_{i=1}^{K_N^t} (1 + (r_B)^{L_N(i)}) \right)} \right],$$

where, as in Giardinà et al. (2015), we have defined

$$r_B = r(\beta, B) = \frac{\lambda_-(\beta, B)}{\lambda_+(\beta, B)}. \quad (3.11)$$

Then

$$\tilde{c}_N(t) = \log \lambda_+(\beta, B+t) - \log \lambda_+(\beta, B) + \frac{1}{N} \log \left[\frac{Q_N \left(\prod_{i=1}^{K_N^t} (1 + (r_{B+t})^{L_N(i)}) \right)}{Q_N \left(\prod_{i=1}^{K_N^t} (1 + (r_B)^{L_N(i)}) \right)} \right]. \quad (3.12)$$

Our aim is to show that the double derivative arises from the first term only, the second derivative of the last term vanishes.

Computation of the second derivative of the cumulant generating function. The second derivative of (3.12) is

$$\tilde{c}_N''(t) = \frac{\partial^2}{\partial t^2} \log \lambda_+(\beta, B+t) + \frac{1}{N \tilde{D}_N(t)} \left[\tilde{I}_N(t) + \tilde{II}_N(t) + \frac{\tilde{III}_N(t)}{\tilde{D}_N(t)} \right], \quad (3.13)$$

where

$$\begin{aligned} \tilde{I}_N(t) &= Q_N \left[\sum_{i=1}^{K_N^t} L_N(i) (L_N(i) - 1) (r_{B+t})^{L_N(i)-2} (r'_{B+t})^2 \right. \\ &\quad \left. + L_N(i) (r_{B+t})^{L_N(i)-1} r''_{B+t} \prod_{\substack{j=1 \\ j \neq i}}^{K_N^t} (1 + (r_{B+t})^{L_N(j)}) \right], \\ \tilde{II}_N(t) &= Q_N \left[\sum_{i=1}^{K_N^t} \sum_{\substack{j=1 \\ j \neq i}}^{K_N^t} L_N(i) L_N(j) (r_{B+t})^{L_N(i)+L_N(j)-2} (r'_{B+t})^2 \prod_{\substack{l=1 \\ l \neq i, j}}^{K_N^t} (1 + (r_{B+t})^{L_N(l)}) \right], \\ \tilde{III}_N(t) &= \left[Q_N \left(\sum_{i=1}^{K_N^t} L_N(i) (r_{B+t})^{L_N(i)-1} r'_{B+t} \prod_{\substack{j=1 \\ j \neq i}}^{K_N^t} (1 + (r_{B+t})^{L_N(j)}) \right) \right]^2, \\ \tilde{D}_N(t) &= Q_N \left[\prod_{i=1}^{K_N^t} (1 + (r_{B+t})^{L_N(i)}) \right]. \end{aligned}$$

Uniform bound of the averaged normalized partition function. To analyze the contributions above we show that the averaged normalized partition function of $\text{CM}_N(\mathbf{2})$ is uniformly bounded:

Lemma 3.1 (The partition function on tori). *For every $\gamma < 1$ and $\alpha \in (0, \infty)$, there exists a constant $A = A(\alpha, \gamma)$ such that, uniformly in N ,*

$$Q_N \left[\prod_{i=1}^{K_N^t} (1 + \alpha \gamma^{L_N^{(i)}}) \right] \leq A.$$

Proof: Denote $\mathcal{Z}_N = Q_N \left[\prod_{i=1}^{K_N^t} (1 + \alpha \gamma^{L_N^{(i)}}) \right]$. For the proof we use induction in N . The induction hypothesis is that there exists an $A > 1$ such that

$$\mathcal{Z}_N \leq A \left(1 - \frac{1}{2\sqrt[3]{N+1}} \right). \quad (3.14)$$

Fix $M \geq 1$ large. We note that we can fix A so large that the inequality is trivially satisfied for $N \leq M$. To advance the induction hypothesis we first derive a recursion relation for \mathcal{Z}_N . We have

$$\begin{aligned} \mathcal{Z}_N &= \sum_{l=1}^N Q_N(L_N(1) = l) Q_N \left(\prod_{i=1}^{K_N^t} (1 + \alpha \gamma^{L_N^{(i)}}) \mid L_N(1) = l \right) \\ &= \sum_{l=1}^N Q_N(L_N(1) = l) (1 + \alpha \gamma^l) \mathcal{Z}_{N-l}. \end{aligned} \quad (3.15)$$

Indeed, the average of $\prod_{i=1}^{K_N^t} (1 + \alpha \gamma^{L_N^{(i)}})$ conditioned on $L_N(1) = l$, reduces to the average on a $\text{CM}_N(\mathbf{2})$ graph with $N - l$ vertices of a similar product. This average gives rise to the factor \mathcal{Z}_{N-l} in (3.15), while the term corresponding to the first cycle is factorized, being $(1 + \alpha \gamma^l)$. Substituting the induction hypothesis into (3.15) leads to

$$\begin{aligned} \mathcal{Z}_N &\leq A \sum_{l=1}^N Q_N(L_N(1) = l) (1 + \alpha \gamma^l) \left(1 - \frac{1}{2\sqrt[3]{N-l+1}} \right) \\ &\leq A \sum_{l=1}^N Q_N(L_N(1) = l) (1 + \alpha \gamma^l) - A \sum_{l=1}^N Q_N(L_N(1) = l) \frac{1}{2\sqrt[3]{N-l+1}}. \end{aligned}$$

It is not hard to see that

$$\sum_{l=1}^N Q_N(L_N(1) = l) \gamma^l \leq c/(N+1),$$

while there exists a constant $\theta > 1$ such that

$$\sum_{l=1}^N Q_N(L_N(1) = l) \frac{1}{\sqrt[3]{N-l+1}} \geq \frac{\theta}{\sqrt[3]{N+1}}. \quad (3.16)$$

Indeed, by [van der Hofstad \(2014b, Exercise 4.1\)](#), or an explicit computation, $L_N(1)/N \xrightarrow{D} T$, where T has density $f_T(x)$ given by

$$f_T(x) = \frac{1}{2\sqrt{1-x}}.$$

Therefore, rewriting the sum in (3.16) we have:

$$\sum_{l=1}^N Q_N(L_N(1) = l) \frac{1}{\sqrt[3]{N-l+1}} = \frac{1}{\sqrt[3]{N+1}} Q_N \left[\frac{1}{\sqrt[3]{1 - L_N(1)/(N+1)}} \right],$$

and by Fatou's Lemma and weak convergence, we obtain

$$\liminf_{N \rightarrow \infty} Q_N \left[\frac{1}{\sqrt[3]{1 - L_N(1)/(N+1)}} \right] \geq \mathbb{E} \left[\frac{1}{\sqrt[3]{1 - T}} \right] > 1.$$

Since we can assume that $N \geq M$, which is sufficiently large, we thus obtain (3.16). Thus,

$$\mathcal{Z}_N \leq A \left(1 + \frac{c}{N+1} - \frac{\theta}{2\sqrt[3]{N+1}} \right) \leq A \left(1 - \frac{1}{2\sqrt[3]{N+1}} \right),$$

when N is sufficiently large. This advances the induction hypothesis and completes the proof of the lemma. \square

Analysis of the second derivative of the cumulant generating function.

Armed with Lemma 3.1, it is now easy to show that all the contributions in the second term of the r.h.s. of (3.13) indeed vanish on a sequence $t_N = o(1)$. To see this, let $t > 0$ and $(t_N)_{N \geq 1}$ a sequence of real numbers such that $t_N \in [0, t/\sqrt{N}]$. We consider first the term $\tilde{I}_N(t_N)$. As in Giardinà et al. (2015, Lemma 3.1), there exists a constant $C > 0$ such that

$$\sum_{i=1}^{K_N^t} |L_N(i)(L_N(i)-1)(r_{B+t_N})^{L_N(i)-2}(r'_{B+t_N})^2 + L_N(i)(r_{B+t_N})^{L_N(i)-1}r''_{B+t_N}| \leq C \cdot K_N^t,$$

since $r_{B+t_N} < 1$. Then, using the Cauchy-Schwarz inequality and recalling that $0 < r < 1$, we obtain

$$\begin{aligned} |\tilde{I}_N(t_N)| &\leq C \cdot Q_N \left(K_N^t \prod_{i=1}^{K_N^t} (1 + (r_{B+t_N})^{L_N(i)}) \right) \\ &\leq C \cdot Q_N \left((K_N^t)^2 \right)^{1/2} \cdot Q_N \left(\prod_{i=1}^{K_N^t} (1 + (r_{B+t_N})^{L_N(i)})^2 \right)^{1/2} \\ &\leq C Q_N \left((K_N^t)^2 \right)^{1/2} \cdot Q_N \left(\prod_{i=1}^{K_N^t} (1 + 3(r_{B+t_N})^{L_N(i)}) \right)^{1/2}. \end{aligned}$$

Using Lemma 3.1 with $\alpha = 3$ and $\gamma = r_{B+t_N}$ we conclude that

$$Q_N \left(\prod_{i=1}^{K_N^t} (1 + 3(r_{B+t_N})^{L_N(i)}) \right)^{1/2} \leq A^{\frac{1}{2}}.$$

Finally, since $\tilde{D}_N(t_N) \geq 1$,

$$\frac{|\tilde{I}_N(t_N)|}{N\tilde{D}_N(t_N)} \leq \frac{C \cdot A^{\frac{1}{2}} \cdot \log N}{N} \xrightarrow{N \rightarrow \infty} 0.$$

Similar computations allow us to estimate $\tilde{II}_N(t_N)$ and $\tilde{III}_N(t_N)$ to obtain

$$\lim_{N \rightarrow \infty} \frac{\tilde{II}_N(t_N)}{N\tilde{D}_N(t_N)} = 0, \quad \lim_{N \rightarrow \infty} \frac{\tilde{III}_N(t_N)}{N \left(\tilde{D}_N(t_N) \right)^2} = 0.$$

Completion of the proof of Theorem 1.6. Having proved that

$$\lim_{N \rightarrow \infty} \frac{1}{N\tilde{D}_N(t_N)} \left[\tilde{I}_N(t_N) + \tilde{II}_N(t_N) + \frac{\tilde{III}_N(t_N)}{\tilde{D}_N(t_N)} \right] = 0,$$

the combination of (3.10) and (3.13) yields the proof of the annealed CLT, i.e.,

$$\begin{aligned} \lim_{N \rightarrow \infty} \log \tilde{P}_N \left[\exp \left(t \frac{S_N - \tilde{P}_N(S_N)}{\sqrt{N}} \right) \right] &= \frac{t^2}{2} \frac{\partial^2}{\partial t^2} \log \lambda_+(\beta, B + t) \Big|_{t=0} \\ &= \frac{t^2}{2} \frac{\cosh(B) e^{-4\beta}}{(\sinh(B) + e^{-4\beta})^{3/2}}. \end{aligned}$$

Therefore, we conclude that the annealed CLT has the same variance as in averaged quenched case [Giardinà et al. \(2015\)](#), i.e., the variance in both cases is the susceptibility of the one-dimensional Ising model. \square

4. Proofs for $\text{CM}_N(\mathbf{1}, \mathbf{2})$

In this section, we consider the Configuration Model $\text{CM}_N(\mathbf{1}, \mathbf{2})$, introduced in Section 1.2. In this graph, the connected components are either cycles or tori (which we indicate by a superscript (t)) connecting vertices of degree 2, or lines (indicated by a superscript (l)) having vertices of degree 2 between two vertices of degree 1. In order to state some properties of the number of lines and tori, we need to introduce some notation. By taking $p \in (0, 1)$, let us define the number of vertex of degree 1 and 2 by

$$n_1 := \# \{i \in [N] : d_i = 1\} = N - \lfloor pN \rfloor, \quad n_2 := \# \{i \in [N] : d_i = 2\} = \lfloor pN \rfloor,$$

and the total degree of the graph by

$$\ell_N = \sum_{i \in [N]} d_i = 2n_2 + n_1 = N + \lfloor pN \rfloor. \quad (4.1)$$

Then, the number of edges is given by $\ell_N/2$. Let us also denote by K_N the number of connected components in the graph and by $K_N^{(l)}$ and $K_N^{(t)}$ the number lines and tori. Obviously,

$$K_N = K_N^{(l)} + K_N^{(t)}.$$

Because every line uses up two vertices of degree 1, the number of lines is given by $n_1/2$, i.e., $K_N^{(l)} = (N - \lfloor pN \rfloor)/2$ a.s.. Regarding the number of cycles, we have that $K_N^{(t)}$ has the same distribution of $K_{\bar{N}}^t$, where \bar{N} is the (random) number of vertices with degree 2 that do not belong to any line and $K_{\bar{N}}^t$ is the number of tori on this set of vertices. Then, since this subset forms a $\text{CM}_{\bar{N}}(\mathbf{2})$ graph, we can apply [Giardinà et al. \(2015, \(3.16\)\)](#) in Section 3.2), obtaining that $K_N^{(t)}/N \xrightarrow{\mathbb{P}} 0$, so that also

$$K_N/N \xrightarrow{\mathbb{P}} (1 - p)/2.$$

Denoting the length (i.e. the number of vertices) in the i th line and j th torus (for an arbitrary labeling) by $L_N^{(l)}(i)$ and $L_N^{(t)}(j)$, the partition function can be computed as

$$Z_N(\beta, B) = \prod_{i=1}^{K_N^{(l)}} Z_{L_N^{(l)}(i)}(\beta, B) \cdot \prod_{i=1}^{K_N^{(t)}} Z_{L_N^{(t)}(i)}(\beta, B), \quad (4.2)$$

where, by (3.4),

$$Z_{L_N^{(t)}(i)}(\beta, B) = \lambda_+^{L_N^{(t)}(i)} + \lambda_-^{L_N^{(t)}(i)},$$

while the partition function on each line is obtained using the partition function on one-dimensional Ising model with free boundary condition [Giardinà et al. \(2015\)](#), Section 3.1) as

$$Z_N^{(l)} = A_+ \lambda_+^N + A_- \lambda_-^N,$$

where

$$A_{\pm} = A_{\pm}(\beta, B) = \frac{e^{-2\beta} e^{\pm B} + (\lambda_+ - e^{\beta+B})^2 e^{\mp B} \pm 2e^{-\beta} (\lambda_+ - e^{\beta+B})}{[e^{-2\beta} + (\lambda_+ - e^{\beta+B})^2] \lambda_{\pm}}. \quad (4.3)$$

This is the starting point of our analysis of the annealed Ising measure on $\text{CM}_N(\mathbf{1}, \mathbf{2})$.

4.1. *Annealed CLT: proof of Theorem 1.9.* In order to prove the CLT in the annealed setting, we will show that

$$\lim_{N \rightarrow \infty} \tilde{P}_N \left[\exp \left(\frac{t}{\sqrt{N}} (S_N - \tilde{P}_N(S_N)) \right) \right] = \exp(\sigma_2^2 t^2 / 2), \quad t \in \mathbb{R}. \quad (4.4)$$

From now on, to alleviate notation we will omit the dependence on β and abbreviate $B_N = B + \frac{t}{\sqrt{N}}$. We start by writing

$$\tilde{P}_N \left[\exp \left(\frac{t}{\sqrt{N}} S_N \right) \right] = \frac{Q_N[Z_N(B_N)]}{Q_N[Z_N(B)]} = \frac{Q_N[e^{NF_{B_N}(p^{(N)}) + NE_N(B_N)}]}{Q_N[e^{NF_B(p^{(N)}) + NE_N(B)}]},$$

where, by [Giardinà et al. \(2015\)](#),

$$F_B(p^{(N)}) = \log \lambda_+(B) + \sum_{l \geq 2} p_l^{(N)} \log \left(A_+(B) + A_-(B) (r(B))^l \right), \quad (4.5)$$

$$E_N(B) = \frac{1}{N} \sum_{i=1}^{K_N^{(t)}} \log \left(1 + r(B)^{L_N^{(t)}(i)} \right), \quad (4.6)$$

with $r(B) = r_B$ defined in (3.11) and $p^{(N)} = (p_l^{(N)})_{l \geq 2}$ the empirical distribution of the lines lengths given by

$$p_l^{(N)} := \frac{1}{N} \sum_{i=1}^{K_N^{(l)}} \mathbb{1}_{\{L_N^{(l)}(i)=l\}}. \quad (4.7)$$

Analysis of the annealed partition function. We have

$$\begin{aligned} e^{NF_B(p^{(N)})} &= (\lambda_+(B))^N \prod_{l=2}^{\infty} \left(A_+(B) + A_-(B) (r(B))^l \right)^{N_l} \\ &= (\lambda_+(B))^N \prod_{l=2}^{\infty} (A_+(B))^{N_l} \prod_{l=2}^{\infty} \left(1 + a(B) (r(B))^l \right)^{N_l}. \end{aligned} \quad (4.8)$$

where

$$a(B) = \frac{A_-(B)}{A_+(B)} \quad (4.9)$$

and $N_l = N p_l^{(N)}$ is the number of lines of length l . We rewrite the second factor in (4.8) as

$$\prod_{l=2}^{\infty} (A_+(B))^{N_l} = (A_+(B))^{n_1/2},$$

since $\sum_{l \geq 2} N_l = n_1/2$. Therefore, we arrive at

$$e^{NF_B(p^{(N)})} = \lambda_+^N(B) A_+^{n_1/2}(B) \prod_{l=2}^{\infty} (1 + a(B)r^l(B))^{N_l} = \lambda_+^N(B) A_+^{n_1/2}(B) \prod_{l=2}^{\infty} c_l(B)^{N_l},$$

where we define

$$c_l(B) := 1 + a(B)r^l(B).$$

Next, define

$$M_N = N - \sum_{l \geq 2} l N_l$$

for the number of vertices that are *not* part of a line. Then, denoting by $Z_N^{(2)}(B)$ the partition function of $\text{CM}_N(\mathbf{2})$,

$$\begin{aligned} \frac{Q_N[Z_N(B_N)]}{Q_N[Z_N(B)]} &= \frac{Q_N[e^{NF_{B_N}(p^{(N)})} \bar{Z}_{M_N}^{(2)}(B_N)]}{Q_N[e^{NF_B(p^{(N)})} \bar{Z}_{M_N}^{(2)}(B)]} \\ &= \frac{\lambda_+^N(B_N) A_+^{n_1/2}(B_N) Q_N[\bar{Z}_{M_N}^{(2)}(B_N) \prod_{l=2}^{\infty} c_l(B_N)^{N_l}]}{\lambda_+^N(B) A_+^{n_1/2}(B) Q_N[\bar{Z}_{M_N}^{(2)}(B) \prod_{l=2}^{\infty} c_l(B)^{N_l}]}, \end{aligned} \quad (4.10)$$

where we write

$$\bar{Z}_N^{(2)}(B) = \lambda_+^{-N} Z_N^{(2)}(B).$$

Asymptotic behavior of the annealed partition function. The key result for the proof of Theorem 1.9 is the following proposition that establishes the exponential growth of the annealed partition function with polynomial corrections:

Proposition 4.1. *The following holds true:*

(a) For $B \neq 0$, there exist $I = I(B)$ and $J = J(B)$ such that, as $N \rightarrow \infty$,

$$Q_N[\bar{Z}_{M_N}^{(2)}(B) \prod_{l=2}^{\infty} c_l(B)^{N_l}] = J(B) e^{I(B)N} (1 + o(1)). \quad (4.11)$$

The function $B \mapsto J(B)$ is continuous, while $B \mapsto I(B)$ is infinitely differentiable.

(b) Given $t \in \mathbb{R}$ there exist $\bar{I} = \bar{I}(t)$ and \bar{J} such that, as $N \rightarrow \infty$,

$$Q_N[\bar{Z}_{M_N}^{(2)}\left(\frac{t}{\sqrt{N}}\right) \prod_{l=2}^{\infty} c_l\left(\frac{t}{\sqrt{N}}\right)^{N_l}] = \bar{J} e^{\bar{I}(t)/\sqrt{N}N} (1 + o(1)). \quad (4.12)$$

The function $t \mapsto \bar{I}(t)$ is infinitely differentiable.

Proof of Theorem 1.9 subject to Proposition 4.1. We start proving the theorem for $B \neq 0$. We substitute (4.11) into (4.10) to arrive at

$$\begin{aligned} \frac{Q_N[Z_N(B_N)]}{Q_N[Z_N(B)]} &= (1 + o(1)) \frac{\lambda_+^N(B_N) A_+^{n_1/2}(B_N) J(B_N) e^{I(B_N)N}}{\lambda_+^N(B) A_+^{n_1/2}(B) J(B) e^{I(B)N}} \\ &= (1 + o(1)) \left(\frac{\lambda_+(B_N)}{\lambda_+(B)}\right)^N \left(\frac{A_+(B_N)}{A_+(B)}\right)^{n_1/2} e^{N(I(B_N) - I(B))}, \end{aligned} \quad (4.13)$$

where we use the fact that $B \mapsto J(B)$ is continuous to obtain that $J(B_N) = (1 + o(1))J(B)$. We can next use the differentiability of $B \mapsto I(B)$ and the fact that $B_N = B + t/\sqrt{N}$ to expand out

$$\frac{Q_N[Z_N(B_N)]}{Q_N[Z_N(B)]} = (1 + o(1))e^{t\sqrt{N}\left[\frac{\partial}{\partial t} \log \lambda_+(B+t)|_{t=0} + \frac{n_1}{2N} \frac{\partial}{\partial t} \log A_+(B+t)|_{t=0} + \frac{\partial}{\partial t} I(B+t)|_{t=0}\right]} \times e^{\sigma_2^2 t^2/2}, \quad (4.14)$$

where

$$\sigma_2^2 = \frac{\partial^2}{\partial t^2} \log \lambda_+(B+t)|_{t=0} + \frac{(1-p)}{2} \frac{\partial^2}{\partial t^2} \log A_+(B+t)|_{t=0} + \frac{\partial^2}{\partial t^2} I(B+t)|_{t=0}. \quad (4.15)$$

Since

$$\begin{aligned} \tilde{P}_N(S_N) = & N \left[\frac{\partial}{\partial t} \log \lambda_+(B+t)|_{t=0} + \frac{n_1}{2N} \frac{\partial}{\partial t} \log A_+(B+t)|_{t=0} + \frac{\partial}{\partial t} I(B+t)|_{t=0} \right] \\ & + o(\sqrt{N}), \end{aligned}$$

then (4.14) implies (4.4), thus proving the theorem in the case $B \neq 0$.

For $B = 0$, in a similar way now using (4.12), we get

$$\frac{Q_N[Z_N(t/\sqrt{N})]}{Q_N[Z_N(0)]} = (1 + o(1))e^{t\sqrt{N}\left[\frac{\partial}{\partial t} \log \lambda_+(t)|_{t=0} + \frac{n_1}{2N} \frac{\partial}{\partial t} \log A_+(t)|_{t=0} + \frac{\partial}{\partial t} \bar{I}(t)|_{t=0}\right]} e^{\bar{\sigma}_2^2 t^2/2},$$

where

$$\bar{\sigma}_2^2 = \frac{\partial^2}{\partial t^2} \log \lambda_+(t)|_{t=0} + \frac{(1-p)}{2} \frac{\partial^2}{\partial t^2} \log A_+(t)|_{t=0} + \frac{\partial^2}{\partial t^2} \bar{I}(t)|_{t=0}. \quad \square$$

Strategy to prove asymptotic behavior. The remainder of this section is devoted to the proof of Proposition 4.1. We use the law of total probability to write

$$\begin{aligned} Q_N[\bar{Z}_{M_N}^{(2)}(B) \prod_{l=2}^{\infty} c_l(B)^{N_l}] \\ = \sum_{m=0}^{n_2} \mathbb{E}_m[\bar{Z}_m^{(2)}(B)] Q_N\left[\prod_{l=2}^{\infty} c_l(B)^{N_l} \mid M_N = m\right] Q_N(M_N = m), \end{aligned} \quad (4.16)$$

where we denote by the symbol \mathbb{E}_m the expectation with respect to an independent $\text{CM}_m(\mathbf{2})$.

Our aim is to prove that the asymptotic behavior of (4.16) is essentially dominated by the term with $m = 0$, which gives the exponential growth $J(B)e^{NI(B)}$ stated in Proposition 4.1. To achieve a full control we analyze in the following the three contributions whose product gives rise to the summand of (4.16):

- i) $\mathbb{E}_m[\bar{Z}_m^{(2)}(B)]$: this is subdominant in the limit $N \rightarrow \infty$ since, by Lemma 3.1, $\sup_m \mathbb{E}_m[\bar{Z}_m^{(2)}(B)]$ is bounded. Therefore it will appear only in the prefactor $J(B)$.
- ii) $Q_N(M_N = m)$: we study the distribution of the number of vertices in tori M_N in Lemma 4.1; in particular we prove the existence of a limiting distribution function in the limit $N \rightarrow \infty$.
- iii) $Q_N\left[\prod_{l=2}^{\infty} c_l(B)^{N_l} \mid M_N = m\right]$: this is rewritten explicitly in Lemma 4.2 and its asymptotics is computed in Lemmata 4.3 and 4.4.

The number of vertices in tori. We start by analyzing the random variable M_N representing the number of vertices belonging to tori.

Lemma 4.1 (The number of vertices in tori). *When $N \rightarrow \infty$, there exists a random variable M such that*

$$M_N \xrightarrow{D} M.$$

Further,

$$\begin{aligned} Q_N(M_N = m) &= \frac{1}{(n_1 + 2n_2 - 1)!!} \binom{n_2}{m} 2^{n_2 - m} (n_2 - m)! (n_1 - 1)! (2m - 1)! \binom{n_1/2 + n_2 - m - 1}{n_2 - m} \\ &= 2^{n_2} \frac{(n_1 - 1)! n_2!}{(n_1 + 2n_2 - 1)!!} 2^{-2m} \binom{2m}{m} \binom{n_1/2 + n_2 - m - 1}{n_2 - m}. \end{aligned} \quad (4.17)$$

Proof: Number the vertices of degree 2 in an arbitrary way. We write

$$M_N = \sum_{l=1}^{\infty} l N_l^{(t)}, \quad \text{where} \quad N_l^{(t)} = \sum_{i=1}^{n_2} J_i(l),$$

and $J_i(l)$ is the indicator that vertex i is in a cycle of length l of which vertex i has the smallest label. We compute that

$$Q_N[N_l^{(t)}] = \frac{n_2}{l} Q_N(\text{vertex 1 is in cycle of length } l) \rightarrow \frac{1}{2l} (2p/(1+p))^l \equiv \lambda_l.$$

It is not hard to see, along the lines of [van der Hofstad \(2014a, Proposition 7.12\)](#), that $(N_l^{(t)})_{l \geq 1}$ converges in distribution to a collection of *independent* Poisson random variables $(P_l)_{l \geq 1}$ with parameters $(\lambda_l)_{l \geq 1}$. Further, since $Q_N[N_l^{(t)}] \leq \frac{1}{2l} (n_2/\ell_N)^l$, which decays exponentially, the contribution from large l equals zero whp, i.e., $Q_N(\exists l > T: N_l^{(t)} > 0)$ is small uniformly in N for T large. This shows that

$$M_N \xrightarrow{D} \sum_{l \geq 1} l P_l \equiv M.$$

Note that

$$Q_N[b^M] = \prod_{l=1}^{\infty} Q_N[b^{lP_l}] = \prod_{l=1}^{\infty} e^{(b^l - 1)\lambda_l} = e^{\sum_{l \geq 1} (b^l - 1)\lambda_l}, \quad (4.18)$$

which is finite only when $b < (1+p)/(2p)$.

To prove (4.17), we note that

$$Q_N(M_N = m) = \frac{1}{(n_1 + 2n_2 - 1)!!} N(n_1, n_2, m), \quad (4.19)$$

where $N(n_1, n_2, m)$ is the number of ways in which the half-edges can be paired such that there are precisely m degree 2 vertices in cycles. We claim that

$$N(n_1, n_2, m) = \binom{n_2}{m} 2^{n_2 - m} (n_2 - m)! (n_1 - 1)! (2m - 1)! \binom{n_1/2 + n_2 - m - 1}{n_2 - m}. \quad (4.20)$$

For this, note that

- (1) there are $\binom{n_2}{m}$ ways to choose the m vertices of degree 2 that are in cycles;
- (2) there are $(2m - 1)!!$ ways to pair the half-edges that are incident to vertices in cycles;

- (3) there are $(n_1-1)!!$ ways to pair the vertices of degree 1 (and this corresponds to the pairing of degree 1 vertices in lines);
- (4) there are $(n_2-m)!$ ways to order the vertices that are in lines;
- (5) there are 2 ways to attach the half-edges of a degree 2 vertex inside a line, and there are in total n_2-m degree 2 vertices in lines, giving 2^{n_2-m} ways to attach their half-edges; and
- (6) finally, there are $\binom{n_1/2+n_2-m-1}{n_2-m}$ ways to create $n_1/2$ lines with n_2-m vertices of degree 2.

Multiplying these numbers out gives (4.20). This completes the proof of Lemma 4.1. \square

Combinatorial expression of the partition function. To perform the asymptotic analysis of the partition function $Q_N[\prod_{l=2}^{\infty} c_l(B)^{N_l} \mid M_N = m]$, we rewrite it as double sum in Lemma 4.2 and then we investigate the asymptotics of the summand in Lemma 4.3 by Stirling's formula. Finally, in Lemma 4.4, we use the Laplace method to estimate the asymptotics of the double sum.

Lemma 4.2 (Generating function of number of lines in $\text{CM}_N(\mathbf{1}, \mathbf{2})$). *For every a, r , for $c_l = 1 + ar^l$ for every $l \geq 2$,*

$$Q_N\left[\prod_{l=2}^{\infty} c_l(B)^{N_l} \mid M_N = m\right] = \sum_{\ell=0}^{n_1/2} \sum_{k=0}^{n_2-m} B_{\ell,k}^{(N)}(n_2-m), \quad (4.21)$$

where

$$B_{\ell,k}^{(N)}(n_2-m) = \binom{n_1/2}{\ell} (ar^2)^\ell r^k \frac{\binom{\ell+k-1}{k} \binom{n_1/2-\ell+n_2-m-k-1}{n_2-m-k}}{\binom{n_1/2+n_2-m-1}{n_2-m}}. \quad (4.22)$$

Proof: When $M_N = m$, we have that n_2-m vertices of degree 2 have to be divided over $n_1/2$ lines. Number the lines as $1, \dots, n_1/2$ in an arbitrary way. Denote the number of degree 2 vertices in line j by Y_j and rewrite

$$Q_N\left[\prod_{l=2}^{\infty} c_l(B)^{N_l} \mid M_N = m\right] = \sum_{(i_1, \dots, i_{n_1/2})} Q_N(Y_1 = i_1, \dots, Y_{n_1/2} = i_{n_1/2}) \prod_{j=1}^{n_1/2} (1 + ar^{i_j+2})$$

where $(i_1, \dots, i_{n_1/2})$ is such that $i_1 + \dots + i_{n_1/2} = n_2-m$. Let $[n_1/2] = \{1, \dots, n_1/2\}$, and expand out $\prod_{j=1}^{n_1/2} (1 + ar^{i_j+2})$ to obtain

$$\sum_{(i_1, \dots, i_{n_1/2})} Q_N(Y_1 = i_1, \dots, Y_{n_1/2} = i_{n_1/2}) \sum_{\Gamma \subseteq [n_1/2]} (ar^2)^{|\Gamma|} \prod_{j \in \Gamma} r^{i_j}. \quad (4.23)$$

where the sum over Γ is over all subsets of $[n_1/2]$. We denote

$$\begin{aligned} N_{n_1, n_2-m} &= \#\{(i_1, \dots, i_{n_1/2}) : i_j \geq 0 \forall j \text{ and } \sum_{j=1}^{n_1/2} i_j = n_2-m\} \\ &= \binom{n_1/2 + n_2 - m - 1}{n_2 - m}, \end{aligned}$$

so that (4.23) is equal to

$$\begin{aligned}
& \sum_{(i_1, \dots, i_{n_1/2})} \frac{1}{\binom{n_1/2+n_2-m-1}{n_2-m}} \sum_{\Gamma \subseteq [n_1/2]} (ar^2)^{|\Gamma|} \sum_{k=0}^{n_2-m} r^k \mathbb{1}_{\{\sum_{j \in \Gamma} i_j = k\}} \\
&= \sum_{(i_1, \dots, i_{n_1/2})} \frac{1}{\binom{n_1/2+n_2-m-1}{n_2-m}} \sum_{\ell=0}^{n_1/2} \binom{n_1/2}{\ell} (ar^2)^\ell \sum_{k=0}^{n_2-m} r^k \mathbb{1}_{\{(i_1 + \dots + i_\ell = k)\}} \\
&= \frac{1}{\binom{n_1/2+n_2-m-1}{n_2-m}} \\
&\quad \times \sum_{\ell=0}^{n_1/2} \binom{n_1/2}{\ell} (ar^2)^\ell \sum_{k=0}^{n_2-m} r^k \binom{\ell+k-1}{k} \binom{n_1/2-\ell+n_2-m-k-1}{n_2-m-k} \\
&= \sum_{\ell=0}^{n_1/2} \sum_{k=0}^{n_2-m} \binom{n_1/2}{\ell} (ar^2)^\ell r^k \frac{\binom{\ell+k-1}{k} \binom{n_1/2-\ell+n_2-m-k-1}{n_2-m-k}}{\binom{n_1/2+n_2-m-1}{n_2-m}}.
\end{aligned}$$

□

Asymptotics by Stirling's formula. We continue the analysis by investigating the asymptotics of $B_{\ell,k}^{(N)}(n_2)$ in (4.22) when ℓ, k and n_2 are of the same asymptotic order. To alleviate the notation we write $B_{a,b}^{(N)}(n_2) := B_{[a],[b]}^{(N)}(n_2)$ when a, b are not necessarily integers.

Lemma 4.3 (Asymptotics of $B_{\ell,k}^{(N)}(n_2)$). *Let $D_p = [0, (1-p)/2] \times [0, p]$. For external fields $B \neq 0$, there exists a function $H(s, t)$ continuous in D_p and smooth in D_p° (the interior of D_p) and a function $C(s, t)$ smooth in D_p° , such that*

$$B_{sN,tN}^{(N)}(n_2) = \frac{C(s, t)}{N} \exp\{NH(s, t)\} (1 + o(1)), \quad \text{as } N \rightarrow \infty, \quad (4.24)$$

Moreover, $H(s, t)$ is strictly concave on its domain D_p and its (unique) maximum point (s^*, t^*) lies in the interior D_p° . In D_p° , the functions are defined as follows:

$$\begin{aligned}
H(s, t) &= (1-p) \log\left(\frac{1-p}{2}\right) - 2s \log(s) - 2\left(\frac{1-p}{2} - s\right) \log\left(\frac{1-p}{2} - s\right) \\
&\quad + s \log(ar^2) + t \log(r) + (s+t) \log(s+t) - t \log(t) \\
&\quad + \left(\frac{1+p}{2} - s - t\right) \log\left(\frac{1+p}{2} - s - t\right) \\
&\quad - (p-t) \log(p-t) - \left(\frac{1+p}{2}\right) \log\left(\frac{1+p}{2}\right) + p \log(p),
\end{aligned} \quad (4.25)$$

and

$$C(s, t) = \frac{1}{2\pi} \frac{\frac{1-p}{2}}{s \left(\frac{1-p}{2} - s\right)} \frac{\sqrt{\left(\frac{1+p}{2} - s - t\right)(s+t)p}}{\sqrt{\left(\frac{1+p}{2}\right)(p-t)t}}. \quad (4.26)$$

Finally, uniformly in $(s, t) \in D_p$,

$$B_{sN,tN}^{(N)}(n_2) \leq CN^{1/2} \exp\{NH(s, t)\}. \quad (4.27)$$

Proof: Using Stirling's approximation in the form $n! = e^{-n}n^n\sqrt{2\pi n}(1+o(1))$ for n large, taking $a, b \in \mathbb{N}$, we can rewrite the binomial coefficients as

$$\binom{b}{a} = e^{[b \log b - a \log a - (b-a) \log (b-a)]n} \cdot \frac{\sqrt{b}(1+o(1))}{\sqrt{a}\sqrt{b-a}\sqrt{2\pi n}}.$$

Plugging the previous formula into (4.22), then (4.24) follows. By inspection, $H(s, t)$ and $C(s, t)$ are smooth functions in D_p° for $B \neq 0$. The function $H(s, t)$ can be further extended by continuity to the boundary ∂D_p of D_p , while $C(s, t)$ cannot be defined in $D_p \setminus D_p^\circ$ since it is unbounded there. In order to prove concavity of $H(s, t)$, we check that its Hessian matrix $\mathbf{Q}(s, t)$ is negative definite on each point of D_p° . For this, we compute the Hessian $\mathbf{Q}(s, t)$ as

$$\begin{pmatrix} \frac{1}{\frac{1-p}{2} - s + p - t} + \frac{1}{s+t} - \frac{2}{\frac{1-p}{2} - s} - \frac{2}{s} & \frac{1}{\frac{1-p}{2} - s + p - t} + \frac{1}{s+t} \\ \frac{1}{\frac{1-p}{2} - s + p - t} + \frac{1}{s+t} & \frac{1}{\frac{1-p}{2} - s + p - t} + \frac{1}{s+t} - \frac{1}{p-t} - \frac{1}{t} \end{pmatrix}.$$

The eigenvalues μ^+ and μ^- of $\mathbf{Q}(s, t)$ are

$$\begin{aligned} \mu^\pm &= \frac{1}{2} \left[\frac{2}{\frac{1+p}{2} - s - t} + \frac{2}{s+t} - \frac{2}{\frac{1-p}{2} - s} - \frac{2}{s} - \frac{1}{p-t} - \frac{1}{t} \right. \\ &\quad \left. \pm \sqrt{\left(\frac{2}{\frac{1-p}{2} - s} + \frac{2}{s} - \frac{1}{p-t} - \frac{1}{t} \right)^2 + 4 \left(\frac{1}{\frac{1+p}{2} - s - t} + \frac{1}{s+t} \right)^2} \right]. \end{aligned}$$

We can easily see that $\mu^- < 0$, and in order to show that also μ^+ is negative, we observe that the determinant of the Hessian matrix is positive. Therefore, $H(s, t)$ is strictly concave in D_p° and, by continuity, concave in D_p . Concavity implies that $(s, t) \mapsto H(s, t)$ has a unique global maximum in D_p , the uniqueness follows by strict concavity and the fact that the maximizer is not on the boundary. In order to find $(s^*, t^*) := \operatorname{argmax}_{(s,t) \in D_p} H(s, t)$, and to prove that it lies in D_p° , we calculate

$$\begin{aligned} \frac{\partial H(s, t)}{\partial s} &= 2 \log \left(\frac{1-p}{2} - s \right) - \log \left(\frac{1+p}{2} - s - t \right) + \log(s+t) - 2 \log(s) \\ &\quad + \log(ar^2), \end{aligned} \tag{4.28}$$

$$\frac{\partial H(s, t)}{\partial t} = \log(s+t) - \log(t) - \log \left(\frac{1+p}{2} - s - t \right) + \log(p-t) + \log(r), \tag{4.29}$$

so that (s^*, t^*) is a solution of the system

$$\begin{cases} \frac{\left(\frac{1-p}{2} - s\right)^2}{s^2} \frac{(s+t)}{\left(\frac{1+p}{2} - s - t\right)} = \frac{1}{ar^2}, \\ \frac{(p-t)}{t} \frac{(s+t)}{\left(\frac{1+p}{2} - s - t\right)} = \frac{1}{r}. \end{cases} \tag{4.30}$$

Since $B \neq 0$, and then both ar^2 and r are finite and larger than zero, it is easy to see from (4.30) that the maximum point cannot be attained on the boundary.

The proof of (4.27) follows similarly, now using that $e^{-n}n^n\sqrt{2\pi n} \leq n! \leq e^{-n}n^n\sqrt{2\pi n}(1+\frac{1}{12n})$ for every $n \geq 1$. The power of N is needed to make the estimate uniform, e.g., by bounding $s((1-p)/2 - s) \geq c/N$ uniformly for $s \geq 1/N$. \square

Asymptotics by Laplace method. In the next Lemma we compute the asymptotic behavior of (4.21) by using the discrete analogue of Laplace method.

Lemma 4.4 (Asymptotics of $Q_N[\prod_{l=2}^{\infty} c_l(B)^{N_l} \mid M_N = m]$). *For every $m \geq 0$ fixed,*

$$\begin{aligned} & Q_N\left[\prod_{l=2}^{\infty} c_l(B)^{N_l} \mid M_N = m\right] \\ &= 2\pi C(s^*, t^*) (\det \mathbf{Q}(s^*, t^*))^{-\frac{1}{2}} (b^*)^m \exp\{NH(s^*, t^*)\} (1 + o(1)), \end{aligned} \quad (4.31)$$

where $(s, t) \mapsto H(s, t)$ and $(s, t) \mapsto C(s, t)$ are defined in Lemma 4.3, (s^*, t^*) is the maximum point of $H(s, t)$, $\mathbf{Q}(s, t)$ is the Hessian matrix of H and

$$b^* = \left(\frac{1+p}{2p}\right) \left(\frac{p-t^*}{\frac{1+p}{2} - s^* - t^*}\right).$$

Proof: We start by proving (4.31) for $m = 0$. Due to Lemma 4.3, we may estimate the asymptotic behavior of the double sum

$$K(N) := \sum_{\ell=0}^{n_1/2} \sum_{k=0}^{n_2} B_{\ell,k}^{(N)}(n_2), \quad (4.32)$$

by making use of the function $f_N(s, t) = N^{-1}C(s, t) \exp\{NH(s, t)\}$ that appeared in (4.24). The correspondence between the two sets of variables (ℓ, k) and (s, t) is given by the simple transformation $s = \ell/N$ and $t = k/N$. We denote this transformation by T_N .

Let us define $\ell_N^* := \lfloor s^* N \rfloor$, $k_N^* := \lfloor t^* N \rfloor$ and introduce $0 < \delta < \min\{p, (1-p)/2\}$. The precise value of δ will be chosen later on. We partition the domain of the summation appearing in the sum of $B_{\ell,k}^{(N)}(n_2)$

$$\Lambda_N = \{(\ell, k) : \ell = 1, \dots, n_1/2, k = 1, \dots, n_2\},$$

into two subsets

$$U_{\delta, N} = \{(\ell, k) \in \Lambda_N : |\ell - \ell_N^*| \leq \delta N + 1, |k - k_N^*| \leq \delta N + 1\}, \quad U_{\delta, N}^c = \Lambda_N \setminus U_{\delta, N}.$$

The set $U_{\delta, N}$ is to be considered as a neighborhood of (ℓ_N^*, k_N^*) , the ‘‘maximum’’ point of $B_{\ell,k}^{(N)}(n_2)$. We observe that $T_N(U_{\delta, N})$ is contained in the neighborhood of (s^*, t^*) in D_p , i.e.,

$$W_{\delta + \frac{1}{N}} = \{(s, t) \in D_p : |s - s^*| \leq \delta + \frac{1}{N}, |t - t^*| \leq \delta + \frac{1}{N}\},$$

while $T_N(U_{\delta, N}^c)$ is contained in its complement $W_{\delta + \frac{1}{N}}^c := D_p \setminus W_{\delta + \frac{1}{N}}$. We rewrite (4.32) as $K(N) = K_1(\delta, N) + K_2(\delta, N)$ where

$$K_1(\delta, N) := \sum_{(\ell, k) \in U_{\delta, N}} B_{\ell,k}^{(N)}(n_2), \quad K_2(\delta, N) := \sum_{(\ell, k) \in U_{\delta, N}^c} B_{\ell,k}^{(N)}(n_2). \quad (4.33)$$

We aim to prove that the asymptotic behavior of $K(N)$ is given by $K_1(\delta, N)$, while $K_2(\delta, N)$ gives a sub-dominant contribution. We start by proving the latter statement.

Bound on $K_2(\delta, N)$. Making use of (4.27), we upper bound

$$K_2(\delta, N) \leq CN^{1/2} \sum_{(\ell, k) \in U_{\delta, N}^c} \exp\{NH(\ell/N, k/N)\}. \quad (4.34)$$

Defining

$$M(\delta) := \sup_{\substack{|s-s^*|>\delta \\ |t-t^*|>\delta}} H(s,t) \geq \sup_{(s,t) \in U_{\delta,N}^c} H(s,t),$$

since the values $(\ell/N, k/N)$ in (4.34) belong to $W_{\delta+\frac{1}{N}}^c$ we can bound $H(\ell/N, k/N) \leq M(\delta)$. We conclude that

$$K_2(\delta, N) \leq CN^{1/2} \exp\{NM(\delta)\} |U_{\delta,N}^c| \leq CN^{5/2} \exp\{NM(\delta)\}, \quad (4.35)$$

which, together with $M(\delta) < H(s^*, t^*)$, implies that

$$\exp\{-NH(s^*, t^*)\} K_2(\delta, N) \rightarrow 0, \quad \text{as } N \rightarrow \infty. \quad (4.36)$$

Let us remark that, besides the condition $0 < \delta < \min\{p, (1-p)/2\}$ (which guarantees that (ℓ, k) and (s, t) are contained in the domains of $B_{\ell,k}^{(N)}(n_2)$ and $f_N(s, t)$), in the previous argument no further condition has been imposed on δ .

Asymptotics of $K_1(\delta, N)$. Here we consider the sum $K_1(\delta, N)$ defined in (4.33). Choose $\varepsilon > 0$ arbitrary and small. By continuity of $C(s, t)$, we can choose $\delta > 0$ small enough so that, for large N ,

$$C(s^*, t^*) - \varepsilon \leq C(s, t) \leq C(s^*, t^*) + \varepsilon, \quad \text{for all } (s, t) \in W_{\delta+\frac{1}{N}}$$

Then using (4.24), we obtain

$$K_1(\delta, N) \leq \frac{C(s^*, t^*) + \varepsilon}{N} \sum_{(\ell,k) \in U_{\delta,N}} \exp\left\{NH\left(\frac{\ell}{N}, \frac{k}{N}\right)\right\} (1 + o(1)), \quad (4.37)$$

and a similar lower bound with $C(s^*, t^*) + \varepsilon$ replaced by $C(s^*, t^*) - \varepsilon$. Recalling that $\mathbf{Q}(s, t)$ is the Hessian matrix of $H(s, t)$, by Taylor expanding up to second order and using that (s^*, t^*) is the maximum, we can write

$$H(s, t) - H(s^*, t^*) \leq \frac{1}{2} \mathbf{x} \cdot \mathbf{Q}(s^*, t^*) \mathbf{x} + c\delta \|\mathbf{x}\|^2/2, \quad \text{for all } (s, t) \in W_{\delta+\frac{1}{N}}$$

where $\mathbf{x} = (s - s^*, t - t^*)$, and a similar lower bound with $c\delta$ replaced by $-c\delta$.

By multiplying (4.37) by $\exp[-NH(s^*, t^*)]$ and applying the previous inequality, we obtain

$$\begin{aligned} & \exp[-NH(s^*, t^*)] K_1(\delta, N) \\ & \leq \frac{C(s^*, t^*) + \varepsilon}{N} \sum_{(\ell,k) \in U_{\delta,N}} \exp\left\{\frac{N}{2} \mathbf{x}^T \cdot \mathbf{Q}(s^*, t^*) \mathbf{x} + c\delta N \|\mathbf{x}\|^2/2\right\}, \end{aligned} \quad (4.38)$$

(where \mathbf{x} is computed with $s = \ell/N$ and $t = k/N$) and a similar lower bound with $C(s^*, t^*) + \varepsilon$ replaced with $C(s^*, t^*) - \varepsilon$ and $+c\delta N \|\mathbf{x}\|^2$ replaced with $-c\delta N \|\mathbf{x}\|^2$. The last step is bounding the sum

$$\tilde{K}_1(\delta, N) := \sum_{(\ell,k) \in U_{\delta,N}} \exp\left\{\frac{N}{2} \mathbf{x}^T \cdot (\mathbf{Q}(s^*, t^*) \pm c\delta \mathbf{I}) \mathbf{x}\right\},$$

Now we can substitute the finite sum in the previous display with the infinite one, since the difference is exponentially small [de Bruijn \(1961\)](#). It is known that (as can be seen by extending [de Bruijn \(1961, \(3.9.4\)\)](#) to two-dimensional sums) that

$$\sum_{\mathbf{j} \in \mathbb{Z}^2} e^{-\mathbf{j}^T \mathbf{A} \mathbf{j} / (2N)} = \frac{2\pi N}{\det(\mathbf{A})^{1/2}} (1 + o(1)),$$

Therefore,

$$\tilde{K}_1(\delta, N) = \frac{2\pi}{\det(\mathbf{Q}(s^*, t^*) \pm c\delta\mathbf{I})^{1/2}} N(1 + o(1)).$$

From the previous equation, recalling (4.38), we obtain

$$\exp[-NH(s^*, t^*)]K_1(\delta, N) \leq \frac{2\pi(C(s^*, t^*) + \varepsilon)}{\det(\mathbf{Q}(s^*, t^*) - c\delta\mathbf{I})^{1/2}} (1 + o(1)),$$

and a similar lower bound with $C(s^*, t^*) + \varepsilon$ replaced with $C(s^*, t^*) - \varepsilon$ and $\mathbf{Q}(s^*, t^*) - c\delta\mathbf{I}$ with $\mathbf{Q}(s^*, t^*) + c\delta\mathbf{I}$. Since ε is arbitrary, the previous inequality implies

$$\lim_{N \rightarrow \infty} \exp[-NH(s^*, t^*)]K(N) = \frac{2\pi C(s^*, t^*)}{\det(\mathbf{Q}(s^*, t^*))^{1/2}},$$

which proves the claim.

Next, we want to generalize the previous result by computing the asymptotic of (4.21) in the case $m \neq 0$. We start by rewriting (4.21) in the following fashion:

$$Q_N\left[\prod_{l=2}^{\infty} c_l(B)^{N_l} \mid M_N = m\right] = \sum_{\ell=0}^{n_1/2} \sum_{k=0}^{n_2-m} G_{\ell,k}^{(N)}(m; n_2) B_{\ell,k}^{(N)}(n_2),$$

where

$$G_{\ell,k}^{(N)}(m; n_2) := \frac{B_{\ell,k}^{(N)}(n_2 - m)}{B_{\ell,k}^{(N)}(n_2)} = \frac{\prod_{j=1}^m \left(\frac{n_1}{2} + n_2 - j\right) \prod_{j=0}^{m-1} (n_2 - k - j)}{\prod_{j=0}^{m-1} (n_2 - j) \prod_{j=1}^m \left(\frac{n_1}{2} + n_2 - \ell - k - j\right)},$$

for $m = 0, 1, \dots, n_2$. By defining the function $F(s, t; m)$ on D_p° given by

$$F(s, t; m) = \left(\frac{1+p}{2p}\right)^m \left(\frac{p-t}{\frac{1+p}{2} - s - t}\right)^m,$$

we obtain that

$$G_{\ell,k}^{(N)}(m; n_2) = F\left(\frac{\ell}{N}, \frac{k}{N}; m\right) (1 + o(1))$$

as $N \rightarrow \infty$. Then the proof is obtained from that for $m = 0$ by replacing $C(s, t)$ by $C(s, t)F(s, t)$. \square

Remark 4.1 (Bound on b^*). *Since $(s^*, t^*) \in D_p^\circ$,*

$$\frac{p-t^*}{\frac{1+p}{2} - s^* - t^*} = \frac{p-t^*}{\left(\frac{1-p}{2} - s^*\right) + (p-t^*)} < 1, \quad \text{so that} \quad b^* < \frac{1+p}{2p}.$$

This will allow us to use in the following the moment generating function $Q_N[(b^)^M]$ defined in (4.18).*

Boundary contribution. Lemma 4.4 proves, for any fixed $0 \leq m < \infty$, the asymptotic exponential growth of $Q_N[\prod_{l=2}^{\infty} c_l(B)^{N_l} \mid M_N = m]$ as $N \rightarrow \infty$. However in formula (4.16) we need to sum over a range of values of m that increases with the volume N . In order to overcome this problem, in the proof of Proposition 4.1 we introduce a cut-off in the sum over m (and then send the cut-off to infinity at the end). In doing so we need to exclude the contribution arising from $B_{\ell,k}^{(N)}$ for ℓ close to the boundary $n_1/2$. This is achieved in the following Lemma.

Lemma 4.5 (Boundary contribution). *For every $\varepsilon > 0$ sufficiently small, as $N \rightarrow \infty$*

$$\sum_{m=0}^{n_2} \mathbb{E}_m[\bar{Z}_m^{(2)}(B)] \sum_{\ell > (1-\varepsilon)\frac{n_1}{2}}^{n_1/2} \sum_{k=0}^{n_2-m} B_{\ell,k}^{(N)}(n_2-m)Q_N(M_N = m) = o(e^{NH(s^*, t^*)}).$$

Proof: By Lemma 4.3, we know that $s^* < \frac{1-p}{2}$. Defining

$$\underline{H}(s^*, t^*) := \sup_{\substack{(s,t) \\ s > (1-\varepsilon)\frac{1-p}{2}}} H(s, t),$$

it follows that $\underline{H}(s^*, t^*) < H(s^*, t^*)$. Further, we define

$$D_{\ell,k}^{(N)}(n_1, n_2, m) = B_{\ell,k}^{(N)}(n_2-m)Q_N(M_N = m),$$

and using (4.17) and the following bound

$$2^{-2(m+1)} \binom{2(m+1)}{m+1} \leq 2^{-2m} \binom{2m}{m},$$

we obtain

$$\begin{aligned} \frac{D_{\ell,k}^{(N)}(n_1, n_2, m+1)}{D_{\ell,k}^{(N)}(n_1, n_2, m)} &\leq \frac{n_2 - m - k}{n_1/2 - \ell + n_2 - m - k - 1} \\ &\leq \frac{n_2}{n_1/2 - \ell + n_2 - 1} \leq \frac{n_2}{n_2 - 1}. \end{aligned} \quad (4.39)$$

As a consequence, using (4.27),

$$\begin{aligned} B_{\ell,k}^{(N)}(n_2-m)Q_N(M_N = m) &\leq B_{\ell,k}^{(N)}(n_2)Q_N(M_N = 0) \left(\frac{n_2}{n_2-1} \right)^m \\ &\leq aN^{1/2} \exp\{H(\ell/N, k/N)\}, \end{aligned}$$

since, $\left(\frac{n_2}{n_2-1}\right)^m \leq \left(\frac{n_2}{n_2-1}\right)^{n_2} \leq a$ for $m \leq n_2$ and some $a > e$. Therefore, using this inequality together with Lemma 3.1, we obtain that

$$\begin{aligned} e^{-NH(s^*, t^*)} \sum_{m=0}^{n_2} \mathbb{E}_m[\bar{Z}_m^{(2)}(B)] \sum_{\ell > (1-\varepsilon)\frac{n_1}{2}}^{n_1/2} \sum_{k=0}^{n_2-m} B_{\ell,k}^{(N)}(n_2-m)Q_N(M_N = m) \\ \leq aAn_2N^{1/2} e^{-NH(s^*, t^*)} \sum_{\ell > (1-\varepsilon)\frac{n_1}{2}}^{n_1/2} \sum_{k=0}^{n_2} \exp\{H(\ell/N, k/N)\} \\ \leq aAN^{7/2} e^{N(\underline{H}(s^*, t^*) - H(s^*, t^*))} \xrightarrow{N \rightarrow \infty} 0. \end{aligned}$$

□

Now we are finally ready for the proof of Proposition 4.1. We treat first the case in the presence of an external field B and then the case without field.

Proof of Proposition 4.1 (a). We fix $\mu \in \{0, \dots, n_2\}$ and $\varepsilon > 0$ sufficiently small. Using (4.16) and Lemma 4.2 we write

$$Q_N[\bar{Z}_{M_N}^{(2)}(B)] \prod_{l=2}^{\infty} c_l(B)^{N_l} = X_{N, \ell \leq (1-\varepsilon)\frac{n_1}{2}}^{(1)}(\mu) + X_{N, \ell \leq (1-\varepsilon)\frac{n_1}{2}}^{(2)}(\mu) + X_{N, \ell > (1-\varepsilon)\frac{n_1}{2}}^{(3)}$$

where

$$X_{N,\ell \leq (1-\varepsilon)\frac{n_1}{2}}^{(1)}(\mu) = \sum_{m=0}^{\mu} \mathbb{E}_m[\bar{Z}_m^{(2)}(B)] \sum_{\ell=0}^{(1-\varepsilon)\frac{n_1}{2}} \sum_{k=0}^{n_2-m} B_{\ell,k}^{(N)}(n_2-m) Q_N(M_N = m), \quad (4.40)$$

$$X_{N,\ell \leq (1-\varepsilon)\frac{n_1}{2}}^{(2)}(\mu) = \sum_{m=\mu+1}^{n_2} \mathbb{E}_m[\bar{Z}_m^{(2)}(B)] \sum_{\ell=0}^{(1-\varepsilon)\frac{n_1}{2}} \sum_{k=0}^{n_2-m} B_{\ell,k}^{(N)}(n_2-m) Q_N(M_N = m), \quad (4.41)$$

$$X_{N,\ell > (1-\varepsilon)\frac{n_1}{2}}^{(3)} = \sum_{m=0}^{n_2} \mathbb{E}_m[\bar{Z}_m^{(2)}(B)] \sum_{\ell > (1-\varepsilon)\frac{n_1}{2}}^{n_1/2} \sum_{k=0}^{n_2-m} B_{\ell,k}^{(N)}(n_2-m) Q_N(M_N = m). \quad (4.42)$$

We analyze the three pieces separately, showing that only the first of them contributes to the exponential growth of $Q_N[\bar{Z}_{M_N}^{(2)}(B) \prod_{i=2}^{\infty} c_i(B)^{N_i}]$. By Lemma 4.1 and Lemma 4.4

$$\begin{aligned} X_{N,\ell \leq (1-\varepsilon)\frac{n_1}{2}}^{(1)}(\mu) &= 2\pi C(s^*, t^*) (\det \mathbf{Q}(s^*, t^*))^{-\frac{1}{2}} \exp\{NH(s^*, t^*)\} \\ &\quad \times \sum_{m=0}^{\mu} \mathbb{E}_m[\bar{Z}_m^{(2)}(B)] Q_N(M = m) (1 + o(1)). \end{aligned} \quad (4.43)$$

The expression in (4.41) can be rewritten as

$$\begin{aligned} X_{N,\ell \leq (1-\varepsilon)\frac{n_1}{2}}^{(2)}(\mu) &= \sum_{m=\mu+1}^{n_2} \mathbb{E}_m[\bar{Z}_m^{(2)}(B)] \sum_{\ell=0}^{(1-\varepsilon)\frac{n_1}{2}} \sum_{k=0}^{n_2-m} \frac{B_{\ell,k}^{(N)}(n_2-m) Q_N(M_N = m)}{B_{\ell,k}^{(N)}(n_2) Q_N(M_N = 0)} \\ &\quad \times B_{\ell,k}^{(N)}(n_2) Q_N(M_N = 0). \end{aligned}$$

Now, by (4.39), $\frac{B_{\ell,k}^{(N)}(n_2-m) Q_N(M_N = m)}{B_{\ell,k}^{(N)}(n_2-(m-1)) Q_N(M_N = m-1)}$ is uniformly bounded by $1 - \delta$ for $\delta > 0$ sufficiently small, because $\ell \leq (1-\varepsilon)n_1/2$. Using this bound together with Lemma 3.1 yields

$$X_{N,\ell \leq (1-\varepsilon)\frac{n_1}{2}}^{(2)}(\mu) \leq A \sum_{\ell=0}^{n_1/2} \sum_{k=0}^{n_2} B_{\ell,k}^{(N)}(n_2) \sum_{m=\mu+1}^{n_2} (1-\delta)^m.$$

Thus there exist $\varepsilon(\mu)$ (with $\varepsilon(\mu) \rightarrow 0$ as $\mu \rightarrow \infty$) such that, uniformly in N ,

$$X_{N,\ell \leq (1-\varepsilon)\frac{n_1}{2}}^{(2)}(\mu) \leq \varepsilon(\mu) \exp\{NH(s^*, t^*)\} (1 + o(1)).$$

Finally, from Lemma 4.5 it results that $X_{N,\ell > (1-\varepsilon)\frac{n_1}{2}}^{(3)} = o(\exp\{NH(s^*, t^*)\})$. Thus, from (4.43) we can identify $I = H(s^*, t^*)$ and

$$J = 2\pi C(s^*, t^*) (\det \mathbf{Q}(s^*, t^*))^{-\frac{1}{2}} \sum_{m=0}^{\infty} \mathbb{E}_m[\bar{Z}_m^{(2)}(B)] (b^*)^m \mathbb{P}(M = m).$$

We remark the previous expression is well-defined since, from Lemma 3.1,

$$\sum_{m=0}^{\infty} \mathbb{E}_m[\bar{Z}_m^{(2)}(B)] (b^*)^m \mathbb{P}(M = m) \leq A \mathbb{E}[(b^*)^M],$$

which is finite because $b^* < \frac{1+p}{2p}$ (see (4.18) in Lemma 4.1 and Remark 4.1).

Proof of Proposition 4.1 (b). In this case we work with a vanishing external field. Defining $t_N := \frac{t}{\sqrt{N}}$ and by Taylor expanding (4.9) around 0, we have

$$a(t_N) := \frac{A_-(t_N)}{A_+(t_N)} = C \frac{t^2}{N} (1 + o(1)), \quad \text{as } N \rightarrow \infty,$$

where C is a constant (whose value actually depends on β). Thus,

$$\begin{aligned} Q_N \left[\prod_{l=2}^{\infty} c_l(t_N)^{N_l} \mid M_N = m \right] &= Q_N \left[\prod_{l=2}^{\infty} (1 + a(t_N)r^l(t_N))^{N_l} \mid M_N = m \right] \\ &= Q_N \left[e^{\sum_{l=2}^{\infty} a(t_N)r^l(t_N)N_l} (1 + o(1)) \mid M_N = m \right] \\ &= Q_N \left[e^{Ct^2 \sum_{l=2}^{\infty} r^l(t_N)p_l^{(N)}} \mid M_N = m \right] (1 + o(1)). \end{aligned}$$

By writing

$$\begin{aligned} Q_N \left[e^{Ct^2 \sum_{l=2}^{\infty} r^l(t_N)p_l^{(N)}} \mid M_N = m \right] & \tag{4.44} \\ &= e^{Ct^2 \sum_{l=2}^{\infty} r^l(t_N)Q_N(p_l^{(N)})} \\ &+ e^{Ct^2 \sum_{l=2}^{\infty} r^l(t_N)Q_N(p_l^{(N)})} Q_N \left[e^{Ct^2 \sum_{l=2}^{\infty} r^l(t_N)(p_l^{(N)} - Q_N(p_l^{(N)}))} - 1 \mid M_N = m \right], \end{aligned}$$

and using formula (4.16), we can rewrite $Q_N [\bar{Z}_{M_N}^{(2)}(t_N) \prod_{l=2}^{\infty} c_l(t_N)^{N_l}] = S_1(N) + S_2(N)$ as the sum of two contributions, due to the two terms in (4.44). Now we analyze $S_1(N)$ and $S_2(N)$ as $N \rightarrow \infty$. First, we remark that the sum in the exponential factor converges in this limit. This can be shown by observing that $r(B) < 1$ (see (3.11)). Therefore, calling $r^* = r(0)$ and given any $\epsilon > 0$ such that $r^* + \epsilon < 1$, thanks to the convergence of $r(t_N)$ to r^* , we have that for all N sufficiently large,

$$\sum_{l=2}^{\infty} r^l(t_N)Q_N(p_l^{(N)}) \equiv \sum_{l=2}^N r^l(t_N)Q_N(p_l^{(N)}) \leq \sum_{l=2}^N (r^* + \epsilon)^l$$

where we used the fact that $p_l^{(N)} \leq 1$ for $l \leq N$ and $p_l^{(N)} = 0$ for all $l > N$. Since the geometric sum in the r.h.s. of the previous display is convergent, the positive series in the l.h.s. is also convergent to some positive value \bar{I}_0 . Thus, by inserting the first term of the r.h.s. of (4.44) in (4.16) and applying bounded convergence, we obtain (4.12) with

$$\bar{I}(t) = \bar{I}_0 C t^2.$$

and

$$\bar{J} = \sum_{m=0}^{\infty} \mathbb{E}_m [\bar{Z}_m^{(2)}(0)] \mathbb{P}(M = m).$$

Further, by Lemma 3.1 and the law of total expectation,

$$S_2(N) \leq A e^{Ct^2 \sum_{l=2}^{\infty} r^l(t_N)Q_N(p_l^{(N)})} Q_N \left[\left[e^{Ct^2 \sum_{l=2}^{\infty} r^l(t_N)(p_l^{(N)} - Q_N(p_l^{(N)}))} - 1 \right] \right].$$

We use the Cauchy-Schwarz inequality to bound

$$\begin{aligned} Q_N \left[\left[e^{Ct^2 \sum_{l=2}^{\infty} r^l(t_N)(p_l^{(N)} - Q_N(p_l^{(N)}))} - 1 \right] \right] & \tag{4.45} \\ \leq Q_N \left[\left(e^{Ct^2 \sum_{l=2}^{\infty} r^l(t_N)(p_l^{(N)} - Q_N(p_l^{(N)}))} - 1 \right)^2 \right]^{1/2}, \end{aligned}$$

and, by Jensen's inequality and Hölder inequality,

$$\begin{aligned} 1 &\leq Q_N \left[e^{Ct^2 \sum_{i=2}^{\infty} r^i(t_N) (p_i^{(N)} - Q_N(p_i^{(N)}))} \right] \\ &\leq \left(Q_N \left[e^{Ct^2 \sqrt{N} \sum_{i=2}^{\infty} r^i(t_N) (p_i^{(N)} - Q_N(p_i^{(N)}))} \right] \right)^{\frac{1}{\sqrt{N}}} \leq 1 + o(1), \end{aligned}$$

due to the existence of the finite limit of $Q_N \left[e^{Ct^2 \sqrt{N} \sum_{i=2}^{\infty} r^i(t_N) (p_i^{(N)} - Q_N(p_i^{(N)}))} \right]$ by [Giardinà et al. \(2015, Lemma 4.3\)](#). Therefore, also the term in (4.45) converges to 0, showing that $S_2(N)$ gives a vanishing contribution. This completes the proof. \square

4.2. *Annealed thermodynamic limits and SLLN: Proof of Theorems 1.7 and 1.8.* Finally, we prove the existence of the thermodynamic limits.

Proof of Theorem 1.7. The thermodynamic limit of the annealed pressure is given by

$$\tilde{\psi}(\beta, B) = \lim_{N \rightarrow \infty} \tilde{\psi}_N(\beta, B) = \lim_{N \rightarrow \infty} \frac{1}{N} \log(Q_N(Z_N(\beta, B))).$$

From (4.10) we can rewrite

$$Q_N(Z_N(\beta, B)) = \lambda_+^N(B) A_+^{n_1/2}(B) Q_N \left[\bar{Z}_{M_N}^{(2)}(B) \prod_{l=2}^{\infty} c_l(B)^{N_l} \right],$$

and then

$$\tilde{\psi}(\beta, B) = \log \lambda_+ + \frac{1-p}{2} \log A_+ + \lim_{N \rightarrow \infty} \frac{1}{N} \log \left\{ Q_N \left[\bar{Z}_{M_N}^{(2)}(B) \prod_{l=2}^{\infty} c_l(B)^{N_l} \right] \right\}.$$

Using Proposition 4.1 we find

$$\tilde{\psi}(\beta, B) = \log \lambda_+(\beta, B) + \frac{1-p}{2} \log A_+(\beta, B) + H(s^*, t^*).$$

To prove the existence of the thermodynamic limit of the magnetization, we use Lemma 2.1 and the existence of the pressure in the thermodynamic limit. Then, remembering that (s^*, t^*) is the maximum point of the function $H(s, t)$, we compute

$$\begin{aligned} \widetilde{M}(\beta, B) &= \frac{\partial}{\partial B} \tilde{\psi}(\beta, B) = \frac{\partial}{\partial B} \log \lambda_+(\beta, B) + \frac{1-p}{2} \frac{\partial}{\partial B} \log A_+(\beta, B) \\ &\quad + s^*(\beta, B) \frac{\partial}{\partial B} \log [a(\beta, B) r^2(\beta, B)] \\ &\quad + t^*(\beta, B) \frac{\partial}{\partial B} \log r(\beta, B) \end{aligned} \tag{4.46}$$

In the limit of small external field B by Taylor expanding (4.9) one has $a(\beta, B) = O(B^2)$. Also, from the fixed point equations (4.30) one can check that $s^*(\beta, B) = O(B^2)$. As a consequence $\lim_{B \rightarrow 0^+} \widetilde{M}(\beta, B) = 0$ for all $\beta > 0$, and therefore, by the definition in (1.9), we conclude that there is no phase transition for $\text{CM}_N(\mathbf{1}, \mathbf{2})$. \square

Proof of Theorem 1.8: Again, the SLLN follows immediately from the existence of the annealed pressure in the thermodynamic limit and its differentiability with respect to B . See the proof of Theorem 1.2 and [Giardinà et al. \(2015, Section 2.2\)](#). \square

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