

# Realizability of a Model in Infinite Statistics

**Don Zagier**

Max-Planck-Institut für Mathematik, Bonn, FRG and University of Maryland, College Park, MD 20742, USA

Received August 26, 1991; in revised form November 22, 1991

**Abstract.** Following Greenberg and others, we study a space with a collection of operators  $a(k)$  satisfying the “ $q$ -mutator relations”  $a(l)a^\dagger(k) - qa^\dagger(k)a(l) = \delta_{k,l}$  (corresponding for  $q = \pm 1$  to classical Bose and Fermi statistics). We show that the  $n! \times n!$  matrix  $A_n(q)$  representing the scalar products of  $n$ -particle states is positive definite for all  $n$  if  $q$  lies between  $-1$  and  $+1$ , so that the commutator relations have a Hilbert space representation in this case (this has also been proved by Fivel and by Bożejko and Speicher). We also give an explicit factorization of  $A_n(q)$  as a product of matrices of the form  $(1 - q^j T)^{\pm 1}$  with  $1 \leq j \leq n$  and  $T$  a permutation matrix. In particular,  $A_n(q)$  is singular if and only if  $q^M = 1$  for some integer  $M$  of the form  $k^2 - k$ ,  $2 \leq k \leq n$ .

## 1. Introduction

In this paper we study the following object: a Hilbert space  $\mathbf{H}$  together with a non-zero distinguished vector  $|0\rangle$  (vacuum state) and a collection of operators  $a_k: \mathbf{H} \rightarrow \mathbf{H}$  satisfying the commutation relations (“ $q$ -mutator relations”)

$$a(l)a^\dagger(k) - qa^\dagger(k)a(l) = \delta_{k,l} \quad (\forall k, l) \quad (1)$$

and the relations

$$a(k)|0\rangle = 0 \quad (\forall k). \quad (2)$$

Here  $q$  is a fixed real number and  $a^\dagger(l)$  denotes the adjoint of  $a(l)$ . The statistics based on the commutation relation (1) generalizes classical Bose and Fermi statistics, corresponding to  $q = 1$  and  $q = -1$ , respectively, as well as the intermediate case  $q = 0$  suggested by Hegstrom and investigated by Greenberg [1]. The study of the general case was initiated by Polyakov and Biedenharn [2].

Our first main result is a realizability theorem saying that the object just described exists if  $-1 < q < 1$ . In view of (2), we can think of the  $a(k)$  as annihilation operators and the  $a^\dagger(k)$  as creation operators. As well as the 0-particle state  $|0\rangle$ , our space must contain the many-particle states obtained by applying combinations of  $a(k)$ 's and  $a^\dagger(k)$ 's to  $|0\rangle$ . To prove the realizability of our model it is obviously necessary and sufficient to consider the minimal space

containing these vectors. We therefore define for each  $q \in \mathbb{R}$  an inner product space  $\mathbf{H}(q)$  generated by  $|0\rangle$  and its images under polynomials in the operators  $a(k)$  and  $a^\dagger(k)$ , subject to the relations (1) and (2). It has a basis consisting of  $n$ -particle states

$$\mathbf{x}_k = a^\dagger(k_1) \cdots a^\dagger(k_n) |0\rangle$$

for each  $n \geq 0$  and each  $n$ -tuple of indices  $\mathbf{k} = (k_1, \dots, k_n)$ , since we can use (1) to write any monomial in the  $a(k)$ 's and  $a^\dagger(k)$ 's as a sum of monomials having all the  $a(k)$ 's on the right and all the  $a^\dagger(k)$ 's on the left, and the only ones of these which do not annihilate  $|0\rangle$  are those consisting of  $a^\dagger(k)$ 's only (the linear independence is clear). By the same argument, we can use (1) and (2) to calculate each scalar product  $(\mathbf{x}_l, \mathbf{x}_k)$  as a polynomial in  $q$ , for instance, for  $k \neq l$  we have

$$\begin{aligned} (\mathbf{x}_{kl}, \mathbf{x}_{lk}) &= \langle 0 | a(l) a(k) a^\dagger(l) a^\dagger(k) |0\rangle = q \langle 0 | a(l) a^\dagger(l) a(k) a^\dagger(k) |0\rangle \\ &= q \langle 0 | (1 + qa^\dagger(l) a(l)) (1 + qa^\dagger(k) a(k)) |0\rangle = q \langle 0 |0\rangle = q. \end{aligned}$$

(Here  $\langle 0|$  denotes the operator  $(|0\rangle, \cdot)$  and we have normalized by  $\langle 0|0\rangle = 1$ .) In particular, for each value of  $q$  the infinite matrix  $A(q) = \{(x_l, x_k)\}_{l,k}$  is well-defined. The condition for the Hilbert space realizability of the  $q$ -mutator relation (1) is then that  $A(q)$  be positive definite, i.e., that  $(\mathbf{x}, \mathbf{x}) > 0$  for every non-zero vector  $\mathbf{x} \in \mathbf{H}(q)$ .

**Theorem 1.** *The matrix  $A(q)$  is positive definite for  $-1 < q < 1$ , so that the  $q$ -mutator relation (1) has a Hilbert space realization for  $q$  in this range.*

It is easy to see that  $(x_k, x_l)$  vanishes unless  $\mathbf{k}$  is a permutation of  $\mathbf{l}$ . Thus the space  $\mathbf{H}(q)$  [respectively the matrix  $A(q)$ ] is the direct sum of infinitely many finite-dimensional spaces (respectively matrices) indexed by all *unordered*  $n$ -tuples  $\{k_1, \dots, k_n\}$ , and we only have to show the positive definiteness of these. We will show in Sect. 2 that the general case of this follows from the case when all of the indices  $k_i$  are distinct. It is not hard to see (Sect. 2) that

$$(\mathbf{x}_{\pi(1)\dots\pi(n)}, \mathbf{x}_{1\dots n}) = q^{I(\pi)} \quad (3)$$

for each permutation  $\pi$  in the  $n^{\text{th}}$  symmetric group  $\mathfrak{S}_n$ , where  $I(\pi)$  denotes the number of *inversions* of  $\pi$ , i.e., the number of  $i, j \in [1, n]$  for which  $i < j$  but  $\pi(i) > \pi(j)$ . Thus the problem reduces to showing that the  $n! \times n!$  matrix  $A_n = A_n(q)$  defined by

$$A_n(\pi, \sigma) = q^{I(\sigma^{-1}\pi)} \quad (\pi, \sigma \in \mathfrak{S}_n) \quad (4)$$

is positive definite for  $q$  between  $-1$  and  $1$ . For this, in turn, it is sufficient by continuity to show that  $A_n(q)$  is non-singular in this range, since  $A_n(0)$  is the identity matrix and the eigenvalues of  $A_n(q)$  vary continuously with  $q$  and are real for  $q$  real (because  $A_n(q)$  is real and symmetric). We will prove the following stronger statement.

**Theorem 2.** *The determinant of the matrix  $A_n(q)$  is given by*

$$\det A_n(q) = \prod_{k=1}^{n-1} (1 - q^{k^2+k})^{\frac{n!(n-k)}{k^2+k}}. \quad (5)$$

*In particular,  $A_n(q)$  is non-singular for all complex numbers  $q$  except the  $N^{\text{th}}$  roots of unity for  $N = 2, 6, 12, \dots, n^2 - n$ .*

We will also describe explicitly the inverse of  $A_n(q)$ . Based on calculations for  $n \leq 5$ , we conjecture that

$$A_n(q)^{-1} \stackrel{?}{\in} \frac{1}{\Delta_n} M_{n!}(\mathbb{Z}[q]), \quad \Delta_n := \prod_{k=1}^{n-1} (1 - q^{k^2+k}). \tag{6}$$

For instance, for  $n = 3$  we have

$$A_3(q) = \begin{pmatrix} 1 & q & q & q^2 & q^2 & q^3 \\ q & 1 & q^2 & q & q^3 & q^2 \\ q & q^2 & 1 & q^3 & q & q^2 \\ q^2 & q & q^3 & 1 & q^2 & q \\ q^2 & q^3 & q & q^2 & 1 & q \\ q^3 & q^2 & q^2 & q & q & 1 \end{pmatrix}, \tag{7}$$

where the rows and columns are indexed by the elements of  $\mathfrak{S}_3$  in the order [123], [213], [132], [231], [312], [321] (we use  $[j_1 \dots j_n]$  to denote the element  $\pi$  of  $\mathfrak{S}_n$  defined by  $\pi(i) = j_i$ ). The determinant of this matrix is  $(1 - q^2)^6 (1 - q^6)$  and its inverse is

$$A_3(q)^{-1} = \Delta_3^{-1} \begin{pmatrix} 1 + q^2 & -q & -q & -q^4 & -q^4 & q^3 + q^5 \\ -q & 1 + q^2 & -q^4 & -q & q^3 + q^5 & -q^4 \\ -q & -q^4 & 1 + q^2 & q^3 + q^5 & -q & -q^4 \\ -q^4 & -q & q^3 + q^5 & 1 + q^2 & -q^4 & -q \\ -q^4 & q^3 + q^5 & -q & -q^4 & 1 + q^2 & -q \\ q^3 + q^5 & -q^4 & -q^4 & -q & -q & 1 + q^2 \end{pmatrix}. \tag{8}$$

Finally, we remark that the matrix  $A_n(q)$  splits as a direct sum of pieces corresponding to the irreducible representations of  $\mathfrak{S}_n$ , the piece corresponding to a representation  $\Pi$  of dimension  $d$  being the direct sum of  $d$  copies of a  $d \times d$  matrix  $A_{n,\Pi}(q)$ . For the bosonic and fermionic cases  $q = 1$  and  $q = -1$  all of these matrices are identically zero except for the one corresponding to the one-dimensional trivial or alternating representation, respectively, but for  $-1 < q < 1$  Theorem 1 says that every representation of every symmetric group occurs in a non-trivial (indeed, non-degenerate) way. (This is the reason for the term ‘‘infinite statistics’’ used by the physicists.) It would be of interest to calculate the determinants of the matrices  $A_{n,\Pi}(q)$ , say in terms of the Young diagram corresponding to  $\Pi$ . By Theorem 2, each of these determinants is a product of cyclotomic polynomials  $\Phi_m(q)$  for integers  $m$  dividing some  $k^2 + k$ ,  $1 \leq k \leq n - 1$ .

The paper is organized as follows. In Sect. 2 we give some generalities on group determinants and show that Theorem 1 follows from Theorem 2, which is then proved in Sect. 3. In Sect. 4 we give an explicit description of the inverse matrix of  $A_n(q)$ , while Sect. 5 gives a conjectural formula for the ‘‘number operators’’ in the Hilbert space  $\mathbf{H}(q)$ .

The author would like to thank O. W. Greenberg who told him about the  $q$ -mutator relation and suggested the problem of proving the positive definiteness for  $-1 < q < 1$ . This positive definiteness has been proved independently by Fivel and by Bozejko and Speicher [3]. (However, Fivel apparently asserts that the zeros of  $A_n(q)$  are all roots of  $q^{2n} = 1$ , which contradicts Theorem 2 and is false for all

$n \geq 4$ .) Consequences and related results are discussed in several subsequent papers by Greenberg [4].

### 2. Group Determinants and the Reduction to $A_n(q)$

Let  $G$  be a finite group of order  $m$  and  $\rho: G \rightarrow GL(V)$  a representation of  $G$  on a (finite-dimensional) complex vector space  $V$ . We can extend  $\rho$  to an algebra homomorphism from the group algebra

$$\mathbb{C}[G] = \left\{ \sum_{g \in G} t_g g \mid t_g \in \mathbb{C} \text{ for } g \in G \right\}$$

to the matrix algebra  $\text{End}(V)$  by  $\rho(\sum t_g g) = \sum t_g \rho(g)$ . The determinant of  $\rho(\sum t_g g)$  is a polynomial  $F_\rho(\mathbf{t})$  of degree  $\dim(V)$  in the  $m$  variables  $\mathbf{t} = \{t_g\}_{g \in G}$  which is determined by and uniquely determines the isomorphism class of the representation  $\rho$ . Thus the entire representation theory of  $G$  can be expressed in terms of the ‘‘group determinants’’  $F_\rho(\mathbf{t})$ ; this is in fact the way that representation theory was developed in its early years (see for instance Weber’s *Lehrbuch der Algebra*, Vol. 2, Chap. 7).

If  $V$  is reducible, say  $V = V_1 \oplus V_2$ , then  $F_\rho(\mathbf{t})$  splits as  $F_{\rho_1}(\mathbf{t}) F_{\rho_2}(\mathbf{t})$ , so the study of group determinants can be reduced to the case of irreducible representations of  $G$ . At the other extreme, let  $(V, R)$  be the (right) regular representation of  $G$ , i.e.  $V = \mathbb{C}^G$  is the  $m$ -dimensional vector space of functions  $f: G \rightarrow \mathbb{C}$  and  $\rho = R$  is given by

$$(R(g)f)(g') = f(g'g) \quad (g, g' \in G).$$

The matrix representation of  $R$  with respect to the basis of  $\delta$ -functions on  $G$  is clearly given by

$$R(g)_{g_1, g_2} = \begin{cases} 1 & \text{if } g_1 g = g_2, \\ 0 & \text{otherwise,} \end{cases}$$

so that the group determinant  $F_R(\mathbf{t})$  is the determinant of the  $m \times m$  matrix  $(t_{g_1^{-1}g_2})_{g_1, g_2 \in G}$ . It is well known that  $R$  contains every irreducible representation  $\Pi$  of  $G$  with positive multiplicity (equal to  $\dim \Pi$ ). Hence if  $F_R(\mathbf{t}) \neq 0$  for some  $\mathbf{t} \in \mathbb{C}^m$  then  $F_\Pi(\mathbf{t}) \neq 0$  for every irreducible representation  $\pi$  and consequently  $F_\rho(\mathbf{t}) \neq 0$  for every representation  $\rho$  of  $G$ . —

Now apply this to  $G = \mathfrak{S}_n$ ,  $m = n!$ . Formula (4) and the discussion just given say that  $A_n = \alpha_n(q)$  is just the matrix representation  $R(\alpha_n)$  of the element

$$\alpha_n = \alpha_n(q) = \sum_{\pi \in \mathfrak{S}_n} q^{l(\pi)} \pi \in \mathbb{C}[\mathfrak{S}_n] \tag{9}$$

acting on the regular representation  $(\mathbb{C}^{\mathfrak{S}_n}, R)$ . Here we are thinking of  $q$  as being a complex number; if  $q$  is thought of as a variable, then  $\alpha_n(q)$  belongs to the group ring  $\mathbb{Z}[q][\mathfrak{S}_n]$ . We will usually consider  $q$  as fixed and omit it from the notation. To prove Theorems 1 and 2, we will forget that  $\alpha_n$  is acting on  $\mathbb{C}^{\mathfrak{S}_n}$  and simply show that it is invertible in the group algebra if  $\prod_{k=1}^{n-1} (1 - q^{k^2+k}) \neq 0$ , in which case the inverse of the matrix  $A_n$  is simply the matrix  $R(\alpha_n^{-1})$ .

We now use this point of view to show how the positive definiteness of  $A(q) = \{(x_l, x_k)\}$  follows from that of the  $n! \times n!$  matrices  $A_n(q)$  for  $n = 1, 2, 3, \dots$ . Equation (1) gives by induction the formula for any indices  $l, k_1, \dots, k_n$  (not

necessarily distinct)

$$a(l) a^\dagger(k_1) \cdots a^\dagger(k_n) = q^n a^\dagger(k_1) \cdots a^\dagger(k_n) a(l) + \sum_{\substack{1 \leq i \leq n \\ k_i = l}} q^{i-1} a^\dagger(k_1) \cdots \widehat{a^\dagger(k_i)} \cdots a^\dagger(k_n),$$

where the sum runs over those indices  $i$  for which  $k_i$  equals  $l$  and the hat over the  $i^{\text{th}}$  term of the product indicates that this term is to be omitted. Combining this with (2) gives

$$a(l) a^\dagger(k_1) \cdots a^\dagger(k_n) |0\rangle = \sum_{\substack{1 \leq i \leq n \\ k_i = l}} q^{i-1} a^\dagger(k_1) \cdots \widehat{a^\dagger(k_i)} \cdots a^\dagger(k_n) |0\rangle.$$

Now induction on  $m$  gives a formula for  $a(l_m) \cdots a(l_1) a^\dagger(k_1) \cdots a^\dagger(k_n) |0\rangle$  as a sum of terms  $q^N a^\dagger(k_1) \cdots \widehat{a^\dagger(k_{i_1})} \cdots \widehat{a^\dagger(k_{i_m})} \cdots a^\dagger(k_n) |0\rangle$ , where  $i_1, \dots, i_m$  are distinct indices with  $k_{i_1}, \dots, k_{i_m}$  equal to  $l_1, \dots, l_m$  in some order, the final result for  $m = n$  being

$$a(l_n) \cdots a(l_1) a^\dagger(k_1) \cdots a^\dagger(k_n) |0\rangle = \sum_{\substack{1 \leq i_1, \dots, i_n \leq n \\ i_1, \dots, i_n \text{ distinct} \\ k_{i_1} = l_1, \dots, k_{i_n} = l_n}} q^{\#\{1 \leq r < s \leq n, i_r > i_s\}} |0\rangle,$$

i.e., in the notation of Sect. 1,

$$(\mathbf{x}_1, \mathbf{x}_k) = \langle 0 | a(l_n) \cdots a(l_1) a^\dagger(k_1) \cdots a^\dagger(k_n) |0\rangle = \sum_{\substack{\pi \in \mathfrak{S}_n \\ l_i = k_{\pi(i)} (i=1, \dots, n)}} q^{I(\pi)}. \quad (10)$$

This formula includes (4) and also shows that  $(\mathbf{x}_1, \mathbf{x}_k) = 0$  unless  $\mathbf{l}$  and  $\mathbf{k}$  are permutations of one another, as already mentioned in Sect. 1, so that  $A(q)$  splits up into the matrices  $A_{\mathbf{k}_0}$  having as entries the numbers  $(\mathbf{x}_1, \mathbf{x}_k)$  for  $\mathbf{l}$  and  $\mathbf{k}$  ranging over all permutations of a given index set  $\mathbf{k}_0$ , e.g. for  $\mathbf{k}_0 = (k, k, l)$  with  $k \neq l$

$$A_{\mathbf{k}_0} = \begin{pmatrix} (\mathbf{x}_{kkl}, \mathbf{x}_{kkl}) & (\mathbf{x}_{kkl}, \mathbf{x}_{klk}) & (\mathbf{x}_{kkl}, \mathbf{x}_{lkk}) \\ (\mathbf{x}_{klk}, \mathbf{x}_{kkl}) & (\mathbf{x}_{klk}, \mathbf{x}_{klk}) & (\mathbf{x}_{klk}, \mathbf{x}_{lkk}) \\ (\mathbf{x}_{lkk}, \mathbf{x}_{kkl}) & (\mathbf{x}_{lkk}, \mathbf{x}_{klk}) & (\mathbf{x}_{lkk}, \mathbf{x}_{lkk}) \end{pmatrix} = \begin{pmatrix} 1+q & q+q^2 & q^2+q^3 \\ q+q^2 & 1+q^3 & q+q^2 \\ q^2+q^3 & q+q^2 & 1+q \end{pmatrix}.$$

In each such matrix, the rows and columns are indexed by the permutations  $\mathbf{k} = \pi \mathbf{k}_0$  of  $\mathbf{k}_0$  or equivalently by the left cosets  $G/H$ , where  $G = \mathfrak{S}_n$  and  $H$  is the subgroup of permutations of  $\mathfrak{S}_n$  fixing  $\mathbf{k}_0$ . Write  $\mathbf{k} = \sigma \mathbf{k}_0$ ,  $\mathbf{l} = \tau \mathbf{k}_0$  with  $\sigma, \tau \in \mathfrak{S}_n$ ; then (10) says that the  $(\mathbf{l}, \mathbf{k})$  matrix coefficient of  $A_{\mathbf{k}_0}$  is equal to

$$\sum_{\substack{\pi \in \mathfrak{S}_n \\ \pi \sigma H = \tau H}} q^{I(\pi)}.$$

But a moment's thought shows that this is simply the  $(\tau H, \sigma H)$ -matrix coefficient (with respect to the basis of  $\delta$ -functions) of the element

$$\alpha_n = \sum_{\pi} q^{I(\pi)} \pi$$

on the subspace  $V = \mathbb{C}^{G/H}$  of  $\mathbb{C}^G$  consisting of functions  $f: G \rightarrow \mathbb{C}$  which satisfy  $f(gh) = f(g)$  for all  $g \in G, h \in H$ . This subspace is invariant under the action  $R$  of  $G$  on  $\mathbb{C}^G$ , so that  $(V, R)$  is a representation of  $G$ . Hence if  $\alpha_n$  is invertible in the group algebra  $\mathbb{C}[G]$ , then the matrix  $A_{k_0}$  is invertible. This completes the reduction of Theorem 1 to Theorem 2.

### 3. Factorization of $\alpha_n$ ; Proof of Theorem 2

We first introduce some notations. As in Sect. 1 we denote by  $[i_1, i_2, \dots, i_n]$  the permutation in  $\mathfrak{S}_n$  which sends 1 to  $i_1, 2$  to  $i_2, \dots, n$  to  $i_n$ . We identify  $\mathfrak{S}_{n-1}$  with the subgroup of  $\mathfrak{S}_n$  consisting of permutations fixing  $n$ . For  $1 \leq k \leq n$  we denote by  $T_{k,n}$  the element  $[1, \dots, k-1, n, k, k+1, \dots, n-1]$  of  $\mathfrak{S}_n$ , i.e.

$$T_{k,n}(i) = \begin{cases} i & 1 \leq i < k, \\ n & i = k, \\ i-1 & k < i \leq n, \end{cases} \quad T_{k,n}^{-1}(i) = \begin{cases} i & 1 \leq i < k, \\ i+1 & k \leq i < n, \\ k & i = n. \end{cases}$$

Any element  $\pi \in \mathfrak{S}_n$  can be represented uniquely as  $\sigma T_{k,n}$  with  $\sigma \in \mathfrak{S}_{n-1}$  and  $1 \leq k \leq n$  (namely  $k = \pi^{-1}(n), \sigma = \pi T_{k,n}^{-1}$ ), and a short calculation shows that then  $I(\pi)$  equals  $I(\sigma) + n - k$ . Hence

$$\alpha_n = \sum_{\pi \in \mathfrak{S}_n} q^{I(\pi)} \pi = \sum_{\substack{\sigma \in \mathfrak{S}_{n-1} \\ 1 \leq k \leq n}} q^{I(\sigma T_{k,n})} \sigma T_{k,n} = \left( \sum_{\sigma \in \mathfrak{S}_{n-1}} q^{I(\sigma)} \sigma \right) \left( \sum_{k=1}^n q^{n-k} T_{k,n} \right).$$

In other words,

**Proposition 1.** Define  $\beta_n = \beta_n(q) = \sum_{k=1}^n q^{n-k} T_{k,n} \in \mathbb{C}[\mathfrak{S}_n]$ . Then  $\alpha_n = \alpha_{n-1} \beta_n$ .

Here  $\alpha_{n-1}$  is considered as an element of  $\mathbb{C}[\mathfrak{S}_n]$  via the inclusion  $\mathfrak{S}_{n-1} \subset \mathfrak{S}_n$ . In particular, the representation of  $\alpha_{n-1}$  in  $R_n$ , the  $n!$ -dimensional regular representation of  $\mathfrak{S}_n$ , consists of  $n$  copies of the representation of  $\alpha_{n-1}$  in  $R_{n-1}$ . Thus in terms of the matrices  $A_n$  we can rewrite Proposition 1 as  $A_n = (A_{n-1} \otimes 1_n) \cdot B_n$ , where  $A_{n-1} \otimes 1_n$  denotes the  $n! \times n!$  block matrix with  $n$  copies of  $A_{n-1}$  on the diagonal blocks and zeros elsewhere and  $B_n = B_n(q)$  has the matrix coefficients

$$B_n(\pi, \sigma) = \begin{cases} q^{n-k} & \text{if } \pi \sigma^{-1} = T_{k,n} \text{ for some } 1 \leq k \leq n, \\ 0 & \text{otherwise.} \end{cases}$$

In particular,  $\det(A_n(q)) = \det(A_{n-1}(q))^n \det(B_n(q))$ , so by induction on  $n$  we have reduced Theorem 2 to the simpler

**Theorem 2'.**  $\det(B_n(q)) = \prod_{i=1}^{n-1} (1 - q^{k^2+k})^{\frac{n!}{k^2+k}}$ .

We now make a second reduction by expressing  $B_n$  in turn as a product of yet simpler matrices.

**Proposition 2.** For each  $n$  define elements  $\gamma_n, \delta_n$  in the group algebra  $\mathbb{C}[\mathfrak{S}_n]$  by

$$\begin{aligned} \gamma_n &= (1 - q^{n-1} T_{1,n}) (1 - q^{n-2} T_{2,n}) \cdots (1 - q T_{n-1,n}), \\ \delta_n &= (1 - q^{n+1} T_{1,n}) (1 - q^n T_{2,n}) \cdots (1 - q^2 T_{n,n}). \end{aligned}$$

Then  $\beta_n \gamma_n = \delta_{n-1}$ .

*Proof.* Let  $\beta_{r,n} = \sum_{k=r}^n q^{n-k} T_{k,n}$ , so that  $\beta_{1,n} = \beta_n, \beta_{n,n} = 1$  (note that  $T_{n,n} = 1 \in \mathfrak{S}_n$ ).

Using the easily checked commutation relation

$$T_{k,n} T_{r,n} = T_{r,n-1} T_{k+1,n} \quad (r \leq k \leq n-1),$$

we find

$$\begin{aligned} \beta_{r,n} \cdot (1 - q^{n-r} T_{r,n}) &= \sum_{k=r+1}^n q^{n-k} T_{k,n} + q^{n-r} T_{r,n} \\ &\quad - q^{n-r} T_{r,n} - \sum_{k=r}^{n-1} q^{2n-k-r} T_{k,n} T_{r,n} \\ &= \sum_{k=r+1}^n q^{n-k} T_{k,n} - \sum_{k=r+1}^n q^{2n-k+1-r} T_{r,n-1} T_{k,n} \\ &= (1 - q^{n-r+1} T_{r,n-1}) \cdot \beta_{r+1,n} \end{aligned}$$

and hence by induction on  $r$  (starting with the trivial case  $r = 0$ )

$$\begin{aligned} &\beta_{1,n} (1 - q^{n-1} T_{1,n}) \cdots (1 - q^{n-r} T_{r,n}) \\ &= (1 - q^n T_{1,n-1}) \cdots (1 - q^{n+1-r} T_{r,n-1}) \beta_{r+1,n}. \end{aligned}$$

The case  $r = n - 1$  of this identity is the desired identity.  $\square$

To complete the proof of Theorem 2 we need to compute the determinants of the factors in  $\gamma_n$  and  $\delta_{n-1}$  under the regular representation  $R_n$  of  $\mathfrak{S}_n$ . We use the inclusions  $\mathfrak{S}_b \subset \mathfrak{S}_{b+1} \subset \cdots \subset \mathfrak{S}_n$  to define elements  $T_{a,b} \in \mathfrak{S}_n$  for all  $1 \leq a \leq b \leq n$  (we actually need only the cases  $b = n - 1$  and  $n$ ). Its characteristic polynomial is given by:

**Lemma.** For  $1 \leq a \leq b \leq n$  the determinant of  $R_n(1 - tT_{a,b})$  is  $(1 - t^{b-a+1})^{\frac{n!}{b-a+1}}$ .

*Proof.* The element  $T_{a,b} \in \mathfrak{S}_n$  is a cyclic permutation of the indices  $a, a + 1, \dots, b$  and hence has order  $b - a + 1$ . But if  $G$  is an arbitrary finite group of order  $m$  and  $g \in G$  an element of order  $d$ , then the characteristic polynomial  $\det(1 - tR(g))$  of  $g$  under the regular representation is  $(1 - t^d)^{m/d}$ , because the cycle structure of the permutation of  $G$  given by left multiplication by  $g^{-1}$  consists of  $m/d$  disjoint cycles of length  $d$ . The lemma follows.  $\square$

The proof of Theorem 2 is now immediate: we have

$$\begin{aligned} \det(R_n(\gamma_n)) &= \prod_{k=1}^{n-1} \det(R_n(1 - q^k T_{n-k,n})) = \prod_{k=1}^{n-1} (1 - q^{k(k+1)})^{\frac{n!}{k+1}}, \\ \det(R_n(\delta_n)) &= \prod_{k=1}^n \det(R_n(1 - q^{k+1} T_{n-k+1,n})) = \prod_{k=1}^n (1 - q^{k(k+1)})^{\frac{n!}{k}}, \end{aligned}$$

and hence

$$\begin{aligned} \det(B_n) &= \det(R_n(\beta_n)) = \frac{\det(R_n(\delta_{n-1}))}{\det(R_n(\gamma_n))} = \frac{\det(R_{n-1}(\delta_{n-1}))^n}{\det(R_n(\gamma_n))} \\ &= \prod_{k=1}^{n-1} (1 - q^{k(k+1)})^{\frac{n \cdot (n-1)!}{k} - \frac{n!}{k+1}} \\ &= \prod_{k=1}^{n-1} (1 - q^{k(k+1)})^{\frac{n!}{k(k+1)}}, \end{aligned}$$

which is Theorem 2'; Theorem 2 then follows by induction from this and Proposition 1.

#### 4. Formula for $A_n(q)^{-1}$

According to Propositions 1 and 2 we have

$$\begin{aligned} \beta_n &= \delta_{n-1} \gamma_n^{-1}, \\ \alpha_n &= \beta_2 \cdots \beta_n = \delta_1 \gamma_2^{-1} \delta_2 \gamma_3^{-1} \cdots \gamma_{n-1}^{-1} \delta_{n-1} \gamma_n^{-1} \end{aligned}$$

and hence

$$\alpha_n^{-1} = \gamma_n \delta_{n-1}^{-1} \gamma_{n-1} \cdots \gamma_2 \delta_1^{-1}.$$

To invert  $\alpha_n$ , therefore, the first step is to invert  $\delta_k$  for each  $k$ .

**Proposition 3.** For  $\pi \in \mathfrak{S}_n$  define  $W(\pi) \in \mathbb{Z}$  by

$$\begin{aligned} W(\pi) &= \sum_{\substack{1 \leq i < j \leq n \\ \pi(i) > \pi(j)}} (1 + (n+1-i)(n+1-j)) \delta_{j-1,i} \\ &= I(\pi) + \sum_{\substack{1 \leq i \leq n-1 \\ \pi(i+1) < \pi(i)}} (n+1-i)(n-i) \end{aligned}$$

and set  $\varepsilon_n = \sum_{\pi \in \mathfrak{S}_n} q^{W(\pi)} \pi^{-1} \in \mathfrak{S}_n$ . Then  $\delta_n^{-1} = \frac{1}{\Delta_{n+1}} \varepsilon_n$  with  $\Delta_{n+1}$  as in Eq. (6).

*Proof.* Denote by  $\sigma \mapsto \tilde{\sigma}$  the map  $\mathfrak{S}_{n-1} \rightarrow \mathfrak{S}_n$  defined by  $\tilde{\sigma}(1) = 1$ ,  $\tilde{\sigma}(i) = \sigma(i-1) + 1$  for  $i > 1$  (this is a homomorphism since  $\tilde{\sigma}$  is just  $T_{1,n}^{-1} \sigma T_{1,n}$ ). Then  $\tilde{T}_{a,b} = T_{a+1,b+1}$  for  $1 \leq a < b \leq n-1$ , so  $\delta_n = (1 - q^{n+1} T_{1,n}) \tilde{\delta}_{n-1}$ . Hence by induction it suffices to show that  $\varepsilon_n (1 - q^{n+1} T_{1,n}) = (1 - q^{n^2+n}) \tilde{\varepsilon}_{n-1}$ .

For  $\pi \in \mathfrak{S}_n$ , let  $k = \pi^{-1}(1)$  and denote by  $\pi'$  the element  $T_{1,n} \pi$  of  $\mathfrak{S}_n$ . Since  $\pi'(k) = n$  but  $\pi'(i) = \pi(i) - 1$  for all  $i \neq k$ , all the terms in the definition of  $W(\pi)$  and of  $W(\pi')$  are the same except those with  $i$  or  $j$  equal to  $k$ , so



$$\begin{aligned}
 W(\pi') - W(\pi) &= \sum_{k < j \leq n} (1 + (n + 1 - k)(n + 1 - j) \delta_{j, k+1}) \\
 &\quad - \sum_{1 \leq i < k} (1 + (n + 1 - i)(n + 1 - k) \delta_{i, k-1}) \\
 &= n - k + (n + 1 - k)(n - k) - (k - 1) \\
 &\quad - \begin{cases} (n + 1 - k)(n + 2 - k) & \text{if } k > 1 \\ 0 & \text{if } k = 1 \end{cases} \\
 &= \begin{cases} -n - 1 & \text{if } k \neq 1, \\ n^2 - 1 & \text{if } k = 1. \end{cases}
 \end{aligned}$$

Hence

$$\begin{aligned}
 \varepsilon_n(1 - q^{n+1} T_{1,n}) &= \sum_{\pi \in \mathfrak{S}_n} (q^{W(\pi)} - q^{W(T_{1,n}\pi) + n + 1}) \pi^{-1} \\
 &= \sum_{\substack{\pi \in \mathfrak{S}_n \\ \pi(1) = 1}} (q^{W(\pi)} - q^{W(\pi) + n^2 + n}) \pi^{-1} \\
 &= (1 - q^{n^2 + n}) \sum_{\sigma \in \mathfrak{S}_{n-1}} q^{W(\sigma)} \tilde{\sigma}^{-1}
 \end{aligned}$$

as desired.  $\square$

We next give a formula expressing  $\gamma_n$  as a sum rather than a product.

**Proposition 4.** *The element  $\gamma_n \in \mathfrak{S}_n$  defined in Proposition 2 is given by*

$$\gamma_n = \sum_{k=1}^n \gamma_{n,k}, \quad \gamma_{n,k} = (-1)^{n-k} \sum_{\pi \in \mathfrak{S}_{n,k}} q^{I(\pi)} \pi^{-1},$$

where  $\mathfrak{S}_{n,k}$  is the subset of  $\mathfrak{S}_n$  of cardinality  $\binom{n-1}{k-1}$  consisting of those permutations  $\pi$  for which  $\pi(1) < \dots < \pi(k) > \dots > \pi(n)$ .

*Proof.* Multiplying out the terms in the product defining  $\gamma_n$ , we find

$$\gamma_n = \sum_{s=0}^{n-1} (-1)^s \sum_{1 \leq i_1 < \dots < i_s \leq n-1} q^{(n-i_1) + \dots + (n-i_s)} T_{i_1,n} T_{i_2,n} \dots T_{i_s,n}.$$

The element  $\sigma = T_{i_1,n} T_{i_2,n} \dots T_{i_s,n}$  of  $\mathfrak{S}_n$  maps  $i_1$  to  $n$ ,  $i_2$  to  $n - 1$ ,  $\dots$ , and  $i_s$  to  $n - s + 1$ , and maps the rest of  $\{1, 2, \dots, n\}$  monotone increasingly to  $\{1, 2, \dots, n - s\}$ . Moreover, it is easy to check that  $(n - i_1) + \dots + (n - i_s)$  equals  $I(\sigma)$ . The proposition now follows on setting  $\pi = \sigma^{-1}$  and  $k = n - s$ .  $\square$

The explicit formulas for  $\delta_n^{-1}$  and  $\gamma_n$  just given together with the formula  $\alpha_n^{-1} = \gamma_n \delta_{n-1}^{-1} \alpha_{n-1}^{-1}$  give an inductive method to calculate  $\alpha_n$  for each  $n$ . To describe this a little more explicitly, we define another element of  $\mathbb{C}[\mathfrak{S}_n]$  by

$$\zeta_n = \varepsilon_n \alpha_n^{-1}$$

with  $\varepsilon_n$  as in Proposition 3. We conjecture that  $\zeta_n$  has coefficients which are polynomials in  $q$ . Propositions 1–3 give  $\alpha_n^{-1} = \mathcal{A}_n^{-1} \gamma_n \zeta_{n-1}$ , so this conjecture implies the conjecture in (6). In fact, the two propositions are equivalent. Indeed, for each  $\pi \in \mathfrak{S}_{n,k}$  we have  $\pi(k) = n$  and hence  $\pi = \sigma T_{n,k}$  for some  $\sigma \in \mathfrak{S}_{n-1}$ , so  $\gamma_{n,k}$

equals  $T_{n,k}^{-1} \gamma_{n-1,k}^*$  with  $\gamma_{n-1,k}^* \in \mathbb{C}[\mathfrak{S}_{n-1}]$  (in fact  $\gamma_{n-1,k}^* = \gamma_{n-1,k-1} - \gamma_{n-1,k}$ ). It follows that if  $\pi$  is any element of  $\mathfrak{S}_n$ , and  $\pi = \sigma T_{n,k}$  ( $1 \leq k \leq n$ ,  $\sigma \in \mathfrak{S}_{n-1}$ ) its canonical decomposition as at the beginning of Sect. 3, then the coefficient of  $\pi$  in  $\Delta_n \alpha_n^{-1}$  equals the coefficient of  $\sigma^{-1}$  in  $\gamma_{n-1,k}^* \zeta_{n-1}$ . In particular, taking  $k = n$  we find that the first  $(n - 1)!$  coefficients in  $\Delta_n \alpha_n^{-1}$  are exactly the coefficients of  $\zeta_{n-1}$ , so that the integrality of  $\Delta_n \alpha_n$  implies that of  $\zeta_{n-1}$  for each  $n$ .

We illustrate with numerical examples for  $n \leq 4$ . For  $n = 2$  we have

$$\alpha_2 = 1 + qT_{1,2}, \quad \alpha_2^{-1} = \frac{1}{1 - q^2} (1 - qT_{1,2}), \quad \varepsilon_2 = 1 + q^3 T_{1,2},$$

$$\zeta_2 = \varepsilon_2 \alpha_2^{-1} = (1 + q^2) - qT_{1,2}.$$

We see that  $\zeta_2$  is integral and that its coefficients are the first two coefficients of  $\Delta_3 \alpha_3^{-1}$ , i.e., the first two coefficients of the matrix in (8). The other coefficients of  $\alpha_3^{-1}$  are obtained by multiplying  $\zeta_2$  by the elements  $\gamma_{2,k}^*$  for  $k = 2$  and  $k = 3$ , and we find

$$(1 - q^2)(1 - q^6) \alpha_3^{-1} = (1 - T_{2,3}^{-1}(qT_{1,3} + q^2 T_{2,3}) + T_{1,3}^{-1}(q^3 T_{1,3} T_{2,3})) \zeta_2$$

$$= (1 + q^2)[123] - q[213] - q[132] - q^4[231]$$

$$- q^4[312] + (q^3 + q^5)[321],$$

giving the remaining coefficients in the first row of the matrix in (8) (the other rows are permutations of the first one). Write this as  $\Delta_3 \alpha_3^{-1} = \{1 + q^2, -q, -q, -q^4, -q^4, q^3 + q^5\}$  in the obvious notation. Using this value of  $\alpha_3^{-1}$  and the value  $\varepsilon_3 = \{1, q^7, q^3, q^4, q^8, q^{11}\}$  we find  $\zeta_3 = \{1 + 2q^2 + q^4 + 2q^6 + q^8, -q - q^3 - q^5 - q^7, -q - q^7, -q^4, -q^4, q^3 + q^5\}$ , which is integral as claimed. Now multiplying this by the various components  $\gamma_{4,k}$  ( $1 \leq k \leq 4$ ), we find

$$\alpha_4^{-1} = (1 - q^2)^{-1} (1 - q^6)^{-1} (1 - q^{12})^{-1}$$

$$\times \{1 + 2q^2 + q^4 + 2q^6 + q^8, -q - q^3 - q^5 - q^7, -q - q^7, -q^4, -q^4, q^3 + q^5;$$

$$-q - q^3 - q^5 - q^7, q^2 + q^4 + q^6, -q^4, -q^9 - q^{11}, 0, q^{10};$$

$$-q^4, 0, q^3 + q^5, q^{10}, -q^8 - q^{10} - q^{12}, q^7 + q^9 + q^{11} + q^{13};$$

$$-q^9 - q^{11}, q^{10}, q^{10}, q^7 + q^{13}, q^7 + q^9 + q^{11} + q^{13},$$

$$-q^6 - 2q^8 - q^{10} - 2q^{12} - q^{14}\},$$

where the 24 components have been listed in the obvious order (namely, the elements  $\sigma \in \mathfrak{S}_3$  in the order above, followed by the elements  $T_{3,4} \sigma$  with the same  $\sigma$ , then the  $T_{2,4} \sigma$ , then  $T_{1,4} \sigma$ ). This gives the 24 elements of the first row of the matrix  $A_4(q)^{-1}$ , the other rows of course being permutations of this one. We have also checked the  $\zeta_4$  has integral coefficients and thus that (6) holds for  $n = 5$ .

### 5. Number Operators

For each index  $k$ , the  $k^{\text{th}}$  number operator is the operator on  $\mathbf{H}$  having each vector  $\mathbf{x}_k = a^\dagger(k_1) \cdots a^\dagger(k_n) |0\rangle$  as an eigenvector with eigenvalue equal to the number of  $i$  with  $k_i = k$ , so that the eigenspace of  $N(k)$  with eigenvalue  $r$  is the space spanned

by the states containing exactly  $r$  particles of type  $k$ . It is easy to see that this definition is equivalent to the requirements

$$N^\dagger(k) = N(k), \quad N(k)|0\rangle = 0, \quad [N(k), a^\dagger(l)] = \delta_{kl}a^\dagger(l) \quad \text{for all } l \quad (11)$$

(and hence  $[N(k), a(l)] = -\delta_{kl}a(l)$  for all  $l$ ).

Consider first the case in which there is only one operator  $a(1)$  and its adjoint, i.e., only one kind of particle. In this case  $\mathbf{H}(q)$  can be realized explicitly as the space spanned by vectors  $e_0 = |0\rangle, e_1, e_2, \dots$  with

$$a(1)e_n = \sqrt{n \frac{1-q^n}{1-q}} e_{n-1} \quad (=0 \text{ for } n=0), \quad a^\dagger(1)e_n = \sqrt{\frac{1}{n+1} \frac{1-q^{n+1}}{1-q}} e_{n+1},$$

since then

$$a(1)a^\dagger(1)e_n - qa^\dagger(1)a(1)e_n = \frac{1-q^{n+1}}{1-q} e_n - q \frac{1-q^n}{1-q} e_n = e_n,$$

while the number operator  $N(1)$  is given by either of the two formulas [5]

$$N(1) = \sum_{n=1}^{\infty} \frac{(1-q)^n}{1-q^n} a^\dagger(1)^n a(1)^n = \sum_{n=1}^{\infty} \frac{(1-q)^n}{\log(1/q^n)} (a^\dagger(1)a(1))^n, \quad (12)$$

as one sees either by computing the action of the expressions on the right on the vectors  $e_n$  or else by verifying the relations (11) using (1) and (2). The first formula makes sense for all  $q$  between  $-1$  and  $1$ , the second (which can be rewritten  $\frac{1-q^{N(1)}}{1-q} = a^\dagger(1)a(1)$ ) only for  $0 < q < 1$ . Both reduce to  $N(1) = a^\dagger(1)a(1)$  in the limit as  $q$  tends to  $1$ . For  $q = 0$  the first formula reduces to

$$N(1) = \sum_{n=1}^{\infty} a^\dagger(1)^n a(1)^n \quad (q = 0), \quad (13)$$

which makes sense since only finitely many of the terms act non-trivially on any given state.

In [1], Greenberg showed that the generalization of (13) to the case when there are many indices  $k$  is

$$N(k) = \sum_{n=1}^{\infty} \sum_{k_2, \dots, k_n} a^\dagger(k_n) \cdots a^\dagger(k_2) a^\dagger(k) a(k) a(k_2) \cdots a(k_n) \quad (q = 0).$$

We now give a conjectural generalization of this formula to the case of arbitrary  $q$  between  $-1$  and  $1$ . It is convenient to express the formula for all  $N(k)$  simultaneously by giving a formula for the energy operator  $\mathcal{E} = \sum_k E_k N(k)$ , where the  $E_k$  (interpreted as the energy of particle  $k$ ) are scalar coefficients.

**Conjecture.** *The operator  $\mathcal{E}$  is given by*

$$\mathcal{E} = \sum_{n=1}^{\infty} \sum_{k_1, \dots, k_n} \sum_{\pi \in \mathfrak{S}_n} \sum_{i=1}^n c_i(q, \pi) E_{k_i} a^\dagger(k_{\pi(n)}) \cdots a^\dagger(k_{\pi(1)}) a(k_1) \cdots a(k_n),$$

where the coefficients  $c_i(q, \pi)$  are given by

$$\sum_{\substack{\pi \in \mathfrak{S}_n \\ 1 \leq i \leq n}} c_i(q, \pi) X^{i-1} \pi = \alpha_n^{-1} (1 - qXT_{1,2}) (1 - q^2XT_{1,3}) \cdots (1 - q^{n-1}XT_{1,n}) \in \mathbb{C}[X][\mathfrak{S}_n].$$

This formula gives the correct result up to terms annihilating all 1-, 2-, and 3-particle states, viz:

$$\begin{aligned} \mathcal{E} &= \sum_k E_k a^\dagger(k) a(k) \\ &+ \frac{1}{1-q^2} \sum_{k,l} \left\{ (E_k + q^2 E_l) a^\dagger(l) a^\dagger(k) - q(E_k + E_l) a^\dagger(k) a^\dagger(l) \right\} a(k) a(l) \\ &+ \frac{1}{(1-q^2)(1-q^6)} \times \\ &\sum_{k,l,m} \left\{ ((1+q^2)E_k + (q^2+q^6)E_l + (q^6+q^8)E_m) a^\dagger(m) a^\dagger(l) a^\dagger(k) \right. \\ &\quad - q(E_k + E_l + q^6 E_m) a^\dagger(m) a^\dagger(k) a^\dagger(l) \\ &\quad - q(E_k + q^6(E_l + E_m)) a^\dagger(l) a^\dagger(m) a^\dagger(k) \\ &\quad - q^4(E_k + E_l + E_m) a^\dagger(k) a^\dagger(m) a^\dagger(l) \\ &\quad - q^4(E_k + E_l + E_m) a^\dagger(l) a^\dagger(k) a^\dagger(m) \\ &\quad \left. + q^3(1+q^2)(E_k + E_l + E_m) a^\dagger(k) a^\dagger(l) a^\dagger(m) \right\} a(k) a(l) a(m) \\ &+ \dots \end{aligned}$$

**Note added in proof.** The conjecture stated in this section has now been proved by Sonia Stanciu (see paper following this one).

## References

- Greenberg, O.W.: Example of infinite statistics. *Phys. Rev. Lett.* **64**, 705–708 (1990)
- Biedenharn, L.C.: *J. Phys.* **A22**, L873–L878 (1989)
- Fivel, D.: Interpolation between Fermi and Bose statistics using generalized commutators. *Phys. Rev. Lett.* **65**, 3361–3364 (1990)  
Bożejko, M., Speicher, R.: *Commun. Math. Phys.* **137**, 519–531 (1991)
- Greenberg, O.W.: Q-mutators and violations of statistics. *Argonne Workshop on Quantum Groups*. T. Curtright, D. Fairlie, C. Zachos (eds.). Singapore: World Scientific, pp. 166–180 (1991)  
Greenberg, O.W.: Particles with small violations of Fermi or Bose statistics. *Phys. Rev. D* **43**, 4111–4120 (1991)  
Greenberg, O.W.: Interactions of particles having small violations of statistics. *Physica A* (to appear)
- Freund, P.G.O., Nambu, Y.: Private communication (1989)