

Review

## Recent Advances in Photoinduced Electron Transfer Processes of Fullerene-Based Molecular Assemblies and Nanocomposites

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Received: 29 March 2012; in revised form: 3 May 2012 / Accepted: 8 May 2012 /

Published: 16 May 2012

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**Abstract:** Photosensitized electron-transfer processes of fullerenes hybridized with electron donating or other electron accepting molecules have been surveyed in this review on the basis of the recent results reported mainly from our laboratories. Fullerenes act as photo-sensitizing electron acceptors with respect to a wide variety of electron donors; in addition, fullerenes in the ground state also act as good electron acceptors in the presence of light-absorbing electron donors such as porphyrins. With single-wall carbon nanotubes (SWCNTs), the photoexcited fullerenes act as electron acceptor. In the case of triple fullerene/porphyrin/SWCNT architectures, the photoexcited porphyrins act as electron donors toward the fullerene and SWCNT. These mechanisms are rationalized with the molecular orbital considerations performed for these huge supramolecules. For the confirmation of the electron transfer processes, transient absorption methods have been used, in addition to time-resolved fluorescence spectral measurements. The kinetic data obtained in solution are found to be quite useful to predict the efficiencies of photovoltaic cells.

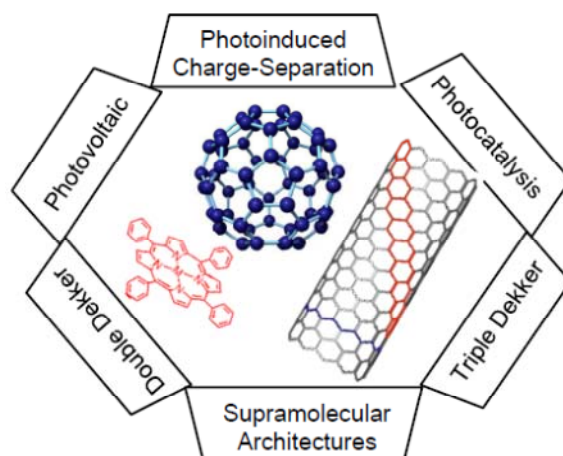
**Keywords:** fullerenes; porphyrins; phthalocyanine; single-wall carbon nanotubes; electron-transfer; photoelectrochemistry

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## 1. Introduction

The study of photoinduced electron transfer in donor-acceptor systems is one of the growing research areas driven primarily by solar energy conversion [1], ultimately used in the construction of molecular electronic and optoelectronic devices [2,3]. Among the donor-acceptor systems, porphyrin-fullerene systems are one of the most widely studied classes of compounds due to their rich photo- and redox chemical properties [4–25]. Both covalently linked systems and non-covalent systems assembled via metal-ligand coordination or  $\pi$ - $\pi$  stacking have been elegantly designed and studied. Fullerenes, owing to their spherical shape, possess high electron-affinity and require small reorganization energy in the electron-transfer processes [21]. Consequently, in donor-acceptor systems, fullerenes tend to accelerate forward electron transfer and slow down backward electron transfer, resulting in the formation of long-lived charge-separated (CS) states [9,19,20]. This is a key factor for utilizing fullerenes in building the solar energy conversion devices [21–23]. The electron-rich macrocyclic compounds such as porphyrins have been widely used as biomimetic photosensitizing electron donor in these studies, since they absorb lights over wide wavelengths in the visible region and exhibit favorable redox potentials [3,4]. In the present review, we include chemically functionalized single-wall carbon nanotubes (SWCNTs) as photoactive electron-conductive materials [26,27] for the development of photovoltaic cells. We expect that a combination of these three kinds of molecules, acting as light-harvesting donors-acceptor systems, will be useful materials for photocatalytic and light-energy conversion applications, as shown in Figure 1.

**Figure 1.** Functionalized fullerenes with porphyrins and SWCNTs, which induce charge separation by light illumination applicable for photo-voltaic and photosensing systems.



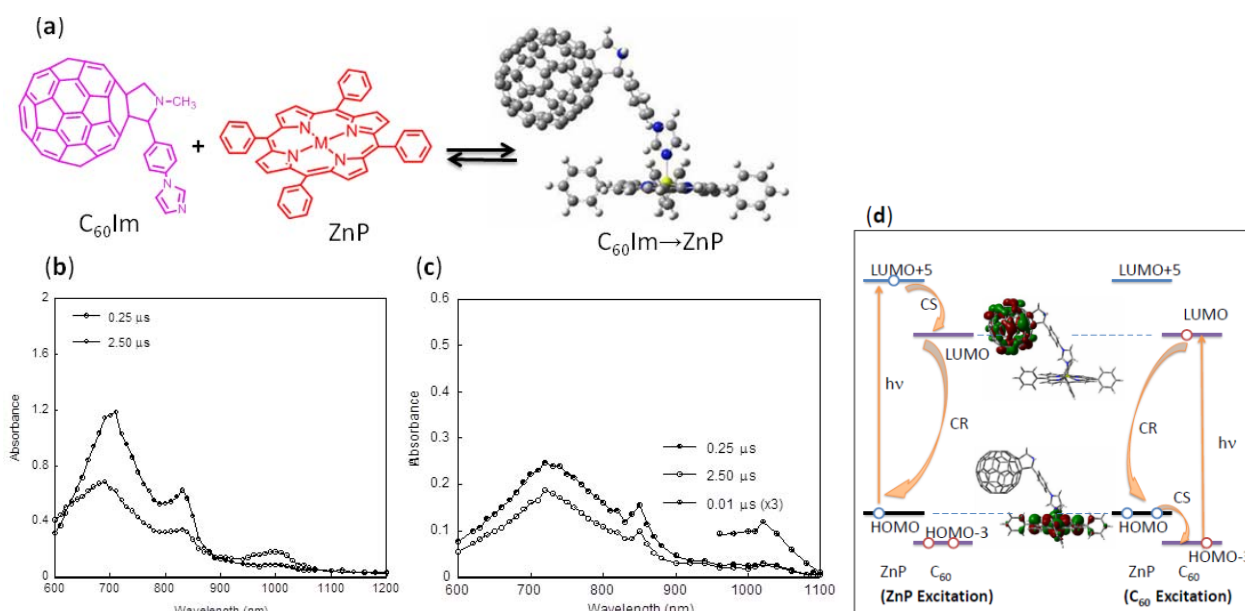
## 2. Fullerene-Porphyrin Systems

We describe here the recent developments in the construction of self-assembled supramolecular donor-acceptor conjugates with porphyrin as donor and fullerene (C<sub>60</sub>) as electron acceptor by adopting different self-assembly mechanisms. The typical examples are the functionalized fullerenes coordinated to the central metal atom of the porphyrins [28]. In addition, the photoinduced charge-separation of the chemically functionalized fullerenes with porphyrin via the covalent bond can be controlled by the additive effect involving coordination to the central metal of the porphyrin [29–33].

## 2.1. Fullerene-Porphyrin Coordination Systems

Usually in polar solvents, a mixture of pristine  $C_{60}$  and zinc or magnesium porphyrins (ZnP or MgP) shows intermolecular electron transfer by visible light illumination, giving their radical anion ( $C_{60}^{\bullet-}$ ) and radical cation ( $MP^{\bullet+}$ ), respectively, as revealed directly by the transient absorption spectral measurements [34,35]. This is also recognized when  $C_{60}$  is functionalized with imidazole ( $C_{60}Im$ ) or pyridine ( $C_{60}Py$ ) entities; that is, they undergo intermolecular electron transfer as revealed by the slow rises of the  $C_{60}^{\bullet-}Im$  at 1,000 nm and  $MP^{\bullet+}$  at 620 nm as shown in the transient absorption spectra and their time-profiles in polar solvents such as  $PhC\equiv N$  (Figure 2b) [36–38]. Under the usual concentrations of  $C_{60}Im$  and MP, the electron transfer takes place via their triplet excited states ( $^3C_{60}^*Im$  at 700 nm and  $^3MP^*$  at 820 nm), which also show a slow decay, confirming the intermolecular electron transfer. These findings are evidence for the destruction of the coordination complex via blocking of the metal atom of MP by the solvent  $PhC\equiv N$  molecules. From the time profiles for the decay of  $^3C_{60}^*Im$  and the rise of  $C_{60}^{\bullet-}Im$  (inset of Figure 2b), the pseudo-first order rate constant can be evaluated to be *ca.*  $3 \times 10^5 s^{-1}$ , giving the second-order rate constant of  $3 \times 10^9 M^{-1}\cdot s^{-1}$  for concentrations of 0.1 mM.

**Figure 2.** (a) Equilibrium between  $C_{60}Im$  and ZnP [for  $C_{60}Im \rightarrow ZnP$ , Zn (yellow), N (blue), C (grey), H (white)]; (b and c) Nanosecond transient absorption spectra and time profiles in Ar-saturated solution; (b) in  $PhC\equiv N$  and (c) in *o*-DCB (adopted from [36]); (d) Energy diagram and HOMO and LUMO of  $C_{60}Im \rightarrow ZnP$  complex (CS; charge separation and CR; charge recombination) constructed from the data in [36].



When both  $C_{60}$  and ZnP are in high concentrations ( $>100$  mM), it is notable that the intermolecular electron transfer takes place via the singlet excited states ( $^1C_{60}^*$  and  $^1ZnP^*$ ), as evidenced by their fluorescence quenching and the quick rise of  $C_{60}^{\bullet-}$  and  $ZnP^{\bullet+}$ ; furthermore, these  $C_{60}^{\bullet-}$  and  $ZnP^{\bullet+}$  decay quickly within 10 ns, even in  $PhC\equiv N$ , due to the radical ion-pairing (RIP) with the singlet-spin character, in which the back electron transfer usually occurs very fast [39].

In *o*-dichlorobenzene (*o*-DCB), however, the supramolecular complex ( $C_{60}Im \rightarrow MP$ ) exists as a major equilibrium species in solution ( $K_{\text{complex}} = 11 \times 10^4 \text{ M}^{-1}$ ); thus, quite different features are found in the transient absorption spectra, as shown in Figure 2c. The time-profile of the  $C_{60}^{\bullet-}Im$  peak at 1,020 nm reveals a sharp rise-and-decay curve within 100 ns immediately after the ns laser-light pulse, although the absorption peak of the  $ZnP^{\bullet+}$  in the 600–680 nm region is hidden by the huge absorptions of the  $C_{60}^*Im$  at 700 nm. The sharp rise of the  $C_{60}^{\bullet-}Im$  corresponds to the  $ZnP$ -fluorescence decay within 0.2 ns ( $k_{\text{CS}} = 1.7 \times 10^{10} \text{ s}^{-1}$ ), suggesting the charge-separation (CS) process takes place via  $^1ZnP^*$  entity within the coordination complex. From the decay of the  $C_{60}^{\bullet-}Im$  within ca. 20 ns, the rate of charge recombination,  $k_{\text{CR}}$ , in the studied dyads was obtained to be  $k_{\text{CR}} = \text{ca. } 5 \times 10^7 \text{ s}^{-1}$ , which is also reasonable due to the close distance between the electron and hole in the  $C_{60}^{\bullet-}Im \rightarrow ZnP^{\bullet+}$ .

In the picosecond time region, the charge-separation in such coordination complexes takes place very rapidly, within 200 ps, with the rise of the radical ions with concomitant decay of the  $S_n \rightarrow S_1$  transition. Then the radical ions decay with ca.  $5 \times 10^7 \text{ s}^{-1}$ , giving the lifetime of the radical ion pair ( $\tau_{\text{RIP}}$ ) to be 20 ns in non-polar solvents.

Figure 2d shows the molecular-orbital (MO) energy-diagram for photoinduced CS and CR processes of the  $C_{60}Im \rightarrow ZnP$  complex via  $^1ZnP^*$  entity with their HOMO and LUMO; the HOMO is mainly localized on the  $ZnP$  moiety, whereas the LUMO is localized on the  $C_{60}$  moiety, suggesting the most stable CS state is the RIP like as  $C_{60}^{\bullet-}Im \rightarrow ZnP^{\bullet+}$ . When the  $ZnP$  is photo-excited, the HOMO-electron of the  $ZnP$  is risen to the upper LUMO+5 localizing on the  $ZnP$ , from which the risen electron transferred to the LUMO localizing on the  $C_{60}$  moiety, leaving the hole on the  $ZnP$ -HOMO, which gives  $C_{60}^{\bullet-}Im \rightarrow ZnP^{\bullet+}$  as a stable RIP. When the  $C_{60}$  moiety is photo-excited, the electron of the HOMO-3 on the  $C_{60}$  moiety is risen to the LUMO of  $C_{60}$ , subsequently, the half-vacant HOMO-3 abstracts an electron from the HOMO level of the  $ZnP$  moiety, giving the  $C_{60}^{\bullet-}Im \rightarrow ZnP^{\bullet+}$ .

From the redox potentials, the free energy change for photoinduced intermolecular electron transfer for the mixture system ( $-\Delta G_{\text{PET}}$ ) can be calculated to be 0.3 eV via either the  $^3C_{60}^*Im$  (ca. 1.5 eV) or  $^3ZnP^*$  (ca. 1.5 eV), whereas the free energy for charge recombination ( $-\Delta G_{\text{CR}}$ ) is calculated to be 1.2 eV, which also represents the energy of the RIP. Since the reorganization energy for electron transfer of the spherical  $C_{60}$  molecule is reported to be as small as 0.5 eV, the CS process positions almost in the top region of the Marcus bell curve, whereas the CR process positions in the deep inverted region.

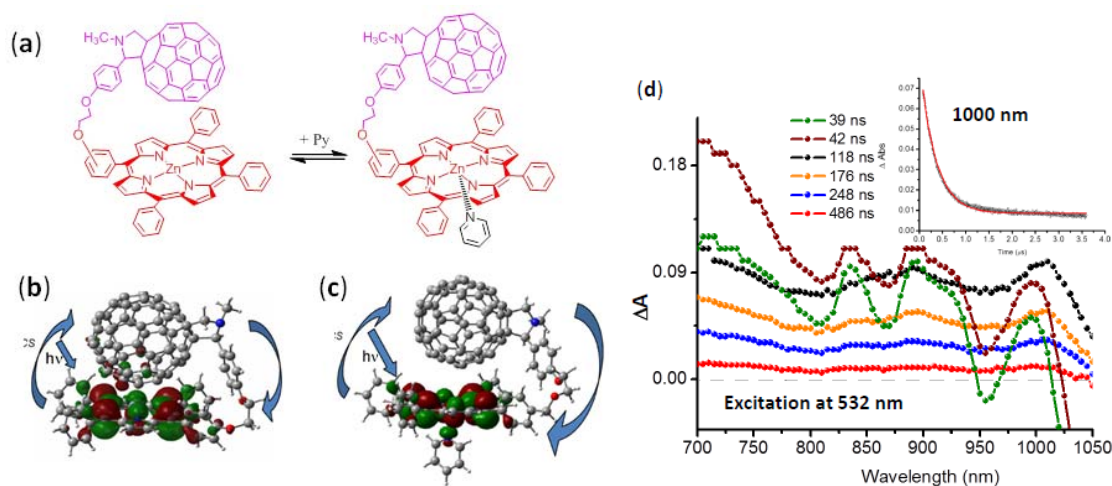
From the redox potentials, the free energy change for intramolecular charge separation ( $\Delta G_{\text{CS}}$ ) for the  $C_{60}Im \rightarrow ZnP$  complex can be calculated to be 0.5 eV via the  $^1C_{60}^*Im$  (ca. 1.8 eV) and to be 0.8 eV via  $^1ZnP^*$  (ca. 2.1 eV). Thus, the CS process positions almost at the top region of the Marcus bell curve, whereas the CR process is still in the deep inverted region. Therefore, the lifetime of the  $C_{60}^{\bullet-}Im \rightarrow ZnP^{\bullet+}$  is prolonged over 1 ns.

From the energies of the excited states, the energy transfer from  $^1ZnP^*$  to  $C_{60}Im$  is also possible; although the  $C_{60}$ -fluorescence is not observed after the decay of the  $ZnP$ -fluorescence, the absorption bands of  $^3C_{60}^*Im$  is observed as shown in Figure 2c, suggesting that the energy transfer might occur concurrently with the CS within  $C_{60}Im \rightarrow ZnP$ .

## 2.2. Fullerene-Porphyrin Charge-Transfer Systems

In another study, the pyridine coordination effect on the charge separation and charge recombination of covalently linked  $C_{60}$ ~ZnP dyads is investigated (Figure 3a) [30]. In *o*-DCB, photoexcitation gives the  $C_{60}^{\bullet-}$ ~ZnP $^{\bullet+}$  via  $C_{60}^{\bullet-}$ ~ZnP\* and/or  ${}^1C_{60}^*~ZnP$  in ca.  $5 \times 10^9$  s $^{-1}$  as revealed by the fluorescence lifetime and ps-transient absorption spectra. The  $C_{60}^{\bullet-}$ ~ZnP $^{\bullet+}$  displays relatively short lifetimes similar to the coordination complex (ca. 20 ns in Figure 2c), because of the vicinity of  $C_{60}^{\bullet-}$  and ZnP $^{\bullet+}$ . However, on addition of a small amount of pyridine to coordinate the Zn ion of the ZnP moiety in *o*-DCB solution, the transient absorption spectra changes to Figure 3d [40]. After 10 ns laser-light pulse, the 1,000 nm peak reaches a maximum at 100 ns. Before 100 ns, sharp peaks appear at 700 nm and 800–900 nm, which can be attributed to  ${}^3C_{60}^*$  and  ${}^3ZnP^*$ , respectively. A sharp peak is also observed at 1,000 nm, which can be ascribed to  $C_{60}^{\bullet-}$ . After 100 ns, the broad bands are found at 1,000 nm and 880 nm, which are also attributed to  $C_{60}^{\bullet-}$  and ZnP $^{\bullet+}$ , respectively, although the solvation of the RIP ( $C_{60}^{\bullet-}$  and ZnP $^{\bullet+}$ ) may change from that before 100 ns. The  $\tau_{RIP}$  of the  $C_{60}^{\bullet-}$  and ZnP $^{\bullet+}$  prolongs up to 200 ns (Figure 3d).

**Figure 3.** Effect of pyridine (Py) addition on the  $C_{60}$ ~ZnP dyad; (a) equilibrium; (b) HOMO of  $C_{60}$ ~ZnP; (c) HOMO of  $C_{60}$ ~ZnP coordinated with Py; (d) noanosecond transient absorption spectra of  $C_{60}$ ~ZnP in the presence of excess Py in Ar-saturated *o*-DCB. Inset: Time profile at 1,000 nm (adopted from [30] and [40]).



The MO calculations of the  $C_{60}$ ~ZnP are shown in Figure 3b, in which the optimized structure indicates the short distance between the  $C_{60}$  sphere and ZnP plane in  $C_{60}$ ~ZnP. Earlier, such calculations on fullerene-based supramolecular nano-architectures predicted geometries and the frontier molecular orbitals to a certain accuracy at the B3LYP/3-21G(\*) level [41]. In the HOMO, although the electron mainly localizes on the ZnP moiety, a small amount of the electron is distributed to the near-side of the  $C_{60}$  sphere, suggesting charge-transfer (CT) interactions exist in the ground state. Therefore, the observed transient species must be short-lived, because of the excited state of the CT state, which is considered to be delocalized RIP in the whole  $C_{60}$ ~ZnP, but not pure RIP like as the HOMO-LUMO in Figure 2c.

On the other hand, the coordination of the pyridine at the ZnP of the opposite side to the  $C_{60}$  sphere tends to elongate the distance between the  $C_{60}$  sphere and the ZnP plane as revealed in the optimized

structure by  $\sim 0.5$  Å as shown in Figure 3c; furthermore, such coordination makes the HOMO predominantly localized on the ZnP moiety (Figure 3c), destroying the CT interaction. Instead, the electron distributes to the coordinated pyridine, which improves the hole-delocalization in the RIP. These characters of the HOMO result in the prolonging the RIP, as observed in Figure 3d.

The coordination of the additional pyridine to the ZnP plane on the same side to the C<sub>60</sub> sphere can separate the distance between the both entities (right structure in Figure 3a); then, the photo-excitation produces pure RIP (C<sub>60</sub><sup>•-</sup>~ZnP<sup>•+</sup>) and its CR rate is slowed down, which may be one of the reasons for long-lived C<sub>60</sub><sup>•-</sup>~ZnP<sup>•+</sup> in Figure 3d.

In addition, other electronic factors can be considered to be responsible for this observation, because of the broadening of the transient bands after 100 ns suggests the characteristic interactions between the pyridine molecules and C<sub>60</sub> molecule, and sometimes induces aggregation of C<sub>60</sub> molecules resulting in fine particles.

When the C<sub>60</sub> moiety with a pyridine unit in C<sub>60</sub>~ZnP is employed, intramolecular coordination of the Py to ZnP forms a parachute-like structure between C<sub>60</sub> and ZnP in *o*-DCB; in such a parachute-like structure, strong interaction between two entities would be anticipated (Figure 2a). Indeed, similar transient spectra to Figure 2c are observed. On addition of excess pyridine molecules, intermolecular coordination bond is destroyed, giving loose structures similar to Figure 3c, and transient spectra similar to Figure 3d with long-living RIP [29].

Summarizing this section, for the photoinduced electron transfer processes for the intermolecular, supramolecular, and flexible covalent bonding systems between C<sub>60</sub> and porphyrins, the appearance of the transient absorption band of C<sub>60</sub><sup>•-</sup> at 1,000 nm is diagnostic evidence of the intermolecular electron transfer and intramolecular charge-separation. These mechanisms and the lifetimes of the C<sub>60</sub><sup>•-</sup> entities in these systems, which are sensitively changed by the delicate molecular structures and environments guided by the MO calculations, are important findings to utilize in the construction the photo-active energy harvesting or fuel production devices.

### 3. Fullerene-SWCNT Systems

Single walled carbon nanotubes (SWCNTs) are emerging new materials for diverse applications including electronics, composites, and biosensors [26,27]. The unique structural properties have made them novel materials in nanotechnology. On combination with light-sensitive molecules, SWCNTs can be modified as light-energy harvesting materials over a wide wavelength range [9,21,23]. Accordingly, various light-sensitive devices and photovoltaic applications using chemically modified SWCNTs have been reported [3–13]. When SWCNT samples including both metallic and semiconducting tubes with different chirality indices are employed as “a material”, no meaningful structure-reactivity relationship could be derived. However, we recently overcame this hurdle by employing single-diameter enriched SWCNT(*n,m*), that allowed us to consider each SWCNT as “a molecule”, which may open a new paradigm for chemically modified SWCNT systems [6]. Thus, by constructing self-assembled supramolecular C<sub>60</sub>/SWCNT(*n,m*) hybrids and C<sub>60</sub>→ZnP/SWCNT(*n,m*), we investigated the photoinduced charge separation processes in solution by various photochemical and photophysical techniques. We also constructed photovoltaic cells by adsorbing the donor-acceptor nanohybrids on an appropriate electrode surface decorated with semiconductor nanoparticles.

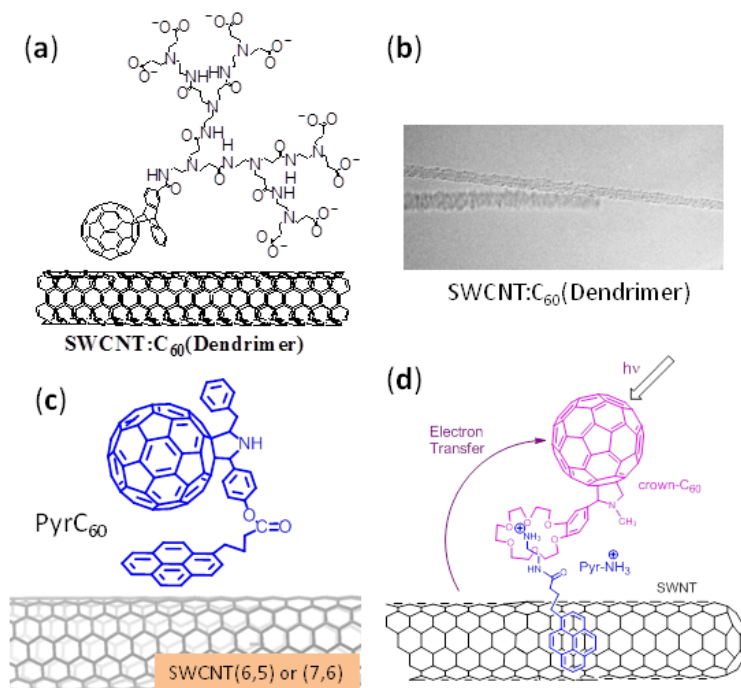


### 3.1. Fullerene-SWCNT Systems

A quick literature survey shows that only a few examples of studies involving SWCNTs with a combination of  $C_{60}$  for probing photoinduced electron transfer have been reported [42–46]. Moreover, no systematic study has ever been done on enriched semi-conducting SWCNT( $n,m$ ) for their ability to be electron donors or acceptors.

To accomplish the non-covalent supramolecular attachment,  $C_{60}$ -dendrimers are mixed with SWCNT in solution [43,44]. Figure 4a is an example of a water-soluble  $C_{60}$ -dendrimer, which tends to surround SWCNT with  $\pi$ - $\pi$  interactions and also with alkyl chain entanglements. Indeed, the TEM images in Figure 4b show the rough surface along SWCNT, suggesting the adsorption of the  $C_{60}$ -dendrimers [32].

**Figure 4.** (a) SWCNT/ $C_{60}$ -dendrimers; (b) TEM image; (c)  $C_{60}$ Pyr /SWCNT assembled via  $\pi$ - $\pi$  interaction (SWCNT(6,5) and SWCNT(7,6)); and (d) SWCNT/pyrNH $_3^+$ /Crown- $C_{60}$  assembled via cation-dipole inclusion complex formation (adopted from [43–46]).



We also newly functionalized  $C_{60}$  to possess a pyrene moiety ( $C_{60}$ Pyr) as shown in Figure 4c. To produce  $C_{60}$ Pyr/SWCNT( $n,m$ ) nanohybrids,  $C_{60}$ Pyr were mixed with SWCNT( $n,m$ ) in DMF [46]. After purification by repeated sonication and centrifugation, un-complexed  $C_{60}$ Pyr were eliminated, increasing the content of attached  $C_{60}$ Pyr molecules [46]. In  $C_{60}$ Pyr/SWCNT nanohybrids thus produced, the Pyr moiety attaches to the SWCNT surface via  $\pi$ - $\pi$  interactions (Figure 4c). In addition, the  $\pi$  orbitals of the  $C_{60}$  sphere may also assist the complex formation via the  $\pi$ - $\pi$  interaction.

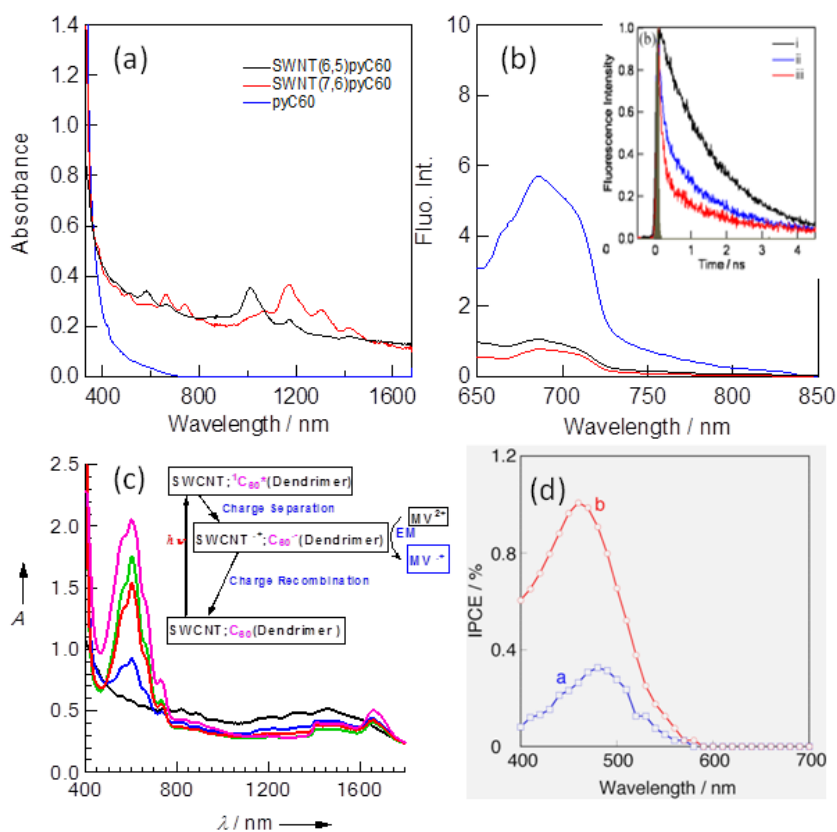
As a third approach, a pyrene functionalized with an alkylammonium cation, PyrNH $_3^+$ , is stacked to SWCNT via  $\pi$ - $\pi$  interaction, forming SWCNT/PyrNH $_3^+$ , to which a benzo-18-crown-6 functionalized  $C_{60}$  (Crown- $C_{60}$ ) is self-assembled via cation-dipole interactions, finally generating SWCNT/PyrNH $_3^+$ /Crown- $C_{60}$  as shown in Figure 4d [45]. The SWCNT/ $C_{60}$  nanohybrids generated by these approaches are characterized by TEM and thermal analysis methods.

### 3.2. Steady-State Spectral Studies of Fullerene-SWCNT Systems

Figure 5a shows the steady-state absorption spectra of  $C_{60}$ Pyr/SWCNT( $n,m$ ) in the wide region from UV-Visible to NIR; the blue spectrum (ii) for  $C_{60}$ Pyr/SWCNT(6,5) shows the absorption peak at ca. 1,000 nm in addition to the absorption of  $C_{60}$ Pyr [black spectrum (i)] in the UV and Visible region, whereas the red spectrum (iii) for  $C_{60}$ Pyr/SWCNT(7,6) shows a peak at 1,200 nm. These longest absorption peaks correspond to the band gaps of the semiconducting SWCNT; the larger diameter one has narrower band gap than that of the smaller one. The 500–400 nm band is attributed to the characteristic absorption of the  $C_{60}$  moiety.

Figure 5b shows the steady-state fluorescence spectra and time profiles of the  $C_{60}$  moiety in  $C_{60}$ Pyr/SWCNT( $n,m$ ) hybrids. Compared with the fluorescence intensity of free  $C_{60}$ Pyr at 680 nm, the intensities decrease down to 1/5, suggesting that photo-physical events originate from the  $^1C_{60}^*Pyr/SWCNT$  state. The fluorescence time profiles gave more reliable information than the steady-state fluorescence measurements, because the fluorescence time profiles do not depend on the concentration and light scattering. Quick decays of the  $C_{60}$  fluorescence give the short lifetimes of ca. 200 ps, which correspond to the fluorescence quenching rate constant of ca.  $5 \times 10^9 \text{ s}^{-1}$ .

**Figure 5.** (a) Steady-state absorption spectra; blue for of  $C_{60}$ Pyr/SWCNT(6,5) and red for  $C_{60}$ Pyr/SWCNT(7,6), black for black for Pyr $C_{60}$ ; (b) fluorescence spectra and time profiles; (c) methyl viologen radical cation ( $MV^{\bullet+}$ ) accumulation by the light excitation of  $C_{60}$  in SWCNT/ $C_{60}$ -dendrimer ( $MV^{2+}$  and BNAH) in Ar-saturated  $H_2O$ ; (d) IPCE for voltaic cell of FTO/ $SnO_2$ / $C_{60}$ Pyr/SWCNT( $n,m$ ) in the presence of  $LiI/I_2$  in acetonitrile solution ( $CH_3C\equiv N$ ) (adopted from [46]).





In order to estimate the photoinduced CS events, it is most convenient to observe the electron-pooling in solution on addition of methyl viologen ( $MV^{2+}$ ) in the presence of appropriate hole-transfer reagents such as 1-benzyl-1,4-dihydronicotinamide (BNAH). On selective photo-excitation of the  $C_{60}$  moiety in the SWCNT/ $C_{60}$ -dendrimers in the presence of  $MV^{2+}$  and BNAH, the  $MV^{2+}$  absorption band intensity at 620 nm is accumulated and preserved as shown in Figure 5c. This implies that  $MV^{2+}$  accepts an electron from the  $C_{60}^{\bullet-}$ , which can be formed via the CS process between the  ${}^1C_{60}^*$  and SWCNT under photo-excitation of the  $C_{60}$  moiety as shown in scheme in inset of Figure 5c. This suggests that the photocatalytic reduction of  $MV^{2+}$  takes place by the photo-excitation of  $C_{60}Pyr/SWCNT(n,m)$  with consuming BNAH.

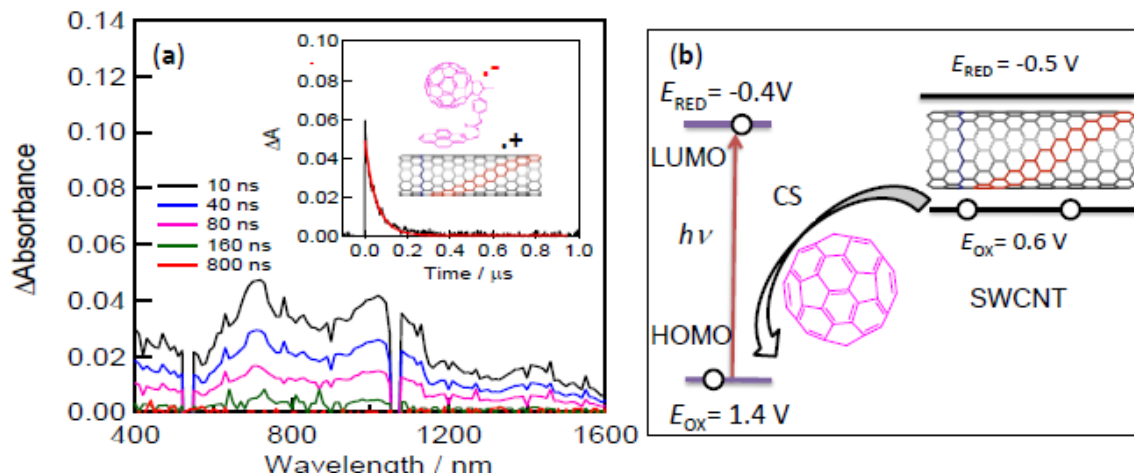
A photovoltaic cell is also constructed by attaching  $C_{60}Pyr/SWCNT(n,m)$  to the  $SnO_2/FTO$  surface and the incident photon to current conversion efficiency (IPCE%) spectra are measured. As shown in Figure 5d, an IPCE peak is observed at 460 nm, which is in agreement with the absorption peak of the  $C_{60}$  moiety, supporting the electron transfer starts from the light absorption of the  $C_{60}$  moiety, but not from SWCNT, although the NIR region data is lacking. Although the IPCE values are not high, a relatively higher value is obtained for  $C_{60}Pyr/SWCNT(7,6)$  compared to  $C_{60}Pyr/SWCNT(6,5)$ , which can be interpreted with the wider band-gap and higher conductivity of SWCNT(7,6) [46].

### 3.3. Direct Evidence for Charge Separation of Fullerene-SWCNT Systems

Nanosecond transient absorption spectra of  $C_{60}Pyr/SWCNT(7,6)$  are measured by the selective photo-excitation of the  $C_{60}$  moiety. Figure 6a shows the observed spectra for the  $C_{60}Pyr/SWCNT(7,6)$  hybrid; similar spectra are obtained for the  $C_{60}Pyr/SWCNT(6,5)$  hybrid. The absorption bands around 1,000 nm corresponding to  $C_{60}^{\bullet-}Pyr$  are clearly recognized with weak peaks in the 1,200–1,600 nm region possibly due to the one-electron oxidized SWCNTs, suggesting the formation of  $C_{60}^{\bullet-}Pyr/SWCNT(7,6)^{\bullet+}$ . Appearance of the 740-nm band corresponding to  ${}^3C_{60}^*$  indicates the loosely bound or free  $C_{60}Pyr$ . The CR rate ( $k_{CR}$ ) evaluated from the  $C_{60}^{\bullet-}Pyr$  decay (inset of Figure 6a) is  $1.80 \times 10^7 \text{ s}^{-1}$  for  $C_{60}^{\bullet-}Pyr/SWCNT(7,6)^{\bullet+}$ ; similarly, the  $k_{CR}$  value evaluated to be  $1.21 \times 10^7 \text{ s}^{-1}$  for  $C_{60}^{\bullet-}Pyr/SWCNT(6,5)^{\bullet+}$ . The larger  $k_{CR}$  value for  $C_{60}^{\bullet-}/SWCNT(6,5)^{\bullet+}$  may be related to the higher electron-conductivity of SWCNT(7,6). The ratios of  $k_{CS}/k_{CR}$ , which are measure of the exploitable electron-hole of the CS state, are found to be  $\sim 360$ , suggesting feasibility of these nanohybrids for the construction of devices using the CS states.

To further understand the excited state events originating from the  $C_{60}$  moiety, an energy level diagram is constructed from the reported electrochemical redox data [47]. As shown in Figure 6b, since the LUMO of  $C_{60}$  is lower than the conduction band of SWCNT, a photo-excited electron of  $C_{60}$  cannot transfer to the conduction band SWCNT, whereas an electron of the valence band of SWCNT( $n,m$ ) can be poured down to the half-vacant  $C_{60}$ -HOMO, generating the charge-separated RIP state,  $C_{60}^{\bullet-}Pyr/SWCNT(m,n)^{\bullet+}$ . The energies of  $C_{60}^{\bullet-}Pyr/SWCNT(m,n)^{\bullet+}$  can be evaluated to be 0.9–1.0 eV from the difference between  $E_{OX(SWCNT)}$  and  $E_{RED(C60)}$ . Thus, the free-energy changes for CS process ( $\Delta G_{CS}$ ) via  ${}^1C_{60}^*Pyr$  (1.80 eV) are negative ( $\Delta G_{CS} = -0.8 - -0.9 \text{ eV}$ ), predicting exothermic CS process.

**Figure 6.** (a) Nanosecond transient absorption spectra of SWCNT(7,6)/PyrC<sub>60</sub> in Ar-saturated *o*-DCB (negative spikes at 532 and 1024 nm are due to the scattering of YAG laser light). Inset: Time profile at 1,000 nm; (b) Energy diagram for charge separation of SWCNT(7,6)/PyrC<sub>60</sub> induced by the photo-excitation of the C<sub>60</sub> moiety (adopted from [46]).



Usually, the fluorescence quenching observed in Figure 5b suggests both the energy transfer and electron transfer. However, direct energy transfer from  $^1\text{C}_{60}^*$  (half-occupied LUMO) to the conduction-band of SWCNT( $n,m$ ) is not efficient, via electron-exchange mechanism since an endothermic process is anticipated since  $E_{\text{LUMO}(\text{C}_{60})} < E_{\text{conduction-band}(\text{SWCNT})}$  in Figure 6b. Resonance energy transfer from  $^1\text{C}_{60}^*$  to SWCNT( $n,m$ ) may also not be efficient, because the C<sub>60</sub>-fluorescence band (650–730 nm) is far from the SWCNT( $n,m$ )-absorption bands (1,000–1,200 nm); furthermore, the spectroscopic overlap is not large, because the C<sub>60</sub>-fluorescence intensity and absorption intensity of the semiconducting SWCNTs in the Visible region are both very low. Therefore, the fluorescence quenching of the C<sub>60</sub> moiety can be attributed to the CS process via  $^1\text{C}_{60}^*$  as confirmed by the transient absorption spectra.

From the energy level diagram (Figure 6b), an electron-transfer from the conduction band of SWCNT( $n,m$ ) to the LUMO of C<sub>60</sub> is plausible after the excitation of SWCNT( $n,m$ ). However, the lifetimes of the excited state of the SWCNT( $n,m$ ) are very short (<1 ps), relaxing to their ground states before an electron transfer to the vicinal C<sub>60</sub>.

To summarize this section, stable nanohybrids are produced with functionalized fullerenes with SWCNT( $n,m$ ), positioning the C<sub>60</sub> spheres in the close proximity to the surface of SWCNT, making the two-layer nanohybrids. Such nanohybrids are found to be stable by the TEM images and steady-state optical spectral studies. Photoinduced CS processes in the supramolecular C<sub>60</sub>/SWCNT( $n,m$ ) hybrids having (6,5)- and (7,6)-enriched SWCNTs was demonstrated with the help of time-resolved fluorescence studies. The nanosecond transient absorption technique confirms the CS products, C<sub>60</sub><sup>•-</sup>/SWCNT<sup>•+</sup>, which upon surface modification operated as photovoltaic cell under solar visible light irradiation.

#### 4. Fullerene-Porphyrin-SWCNT Systems

Since the absorption intensity of the C<sub>60</sub> moiety in the Visible region is not high enough to utilize them in solar light harvesting, we tried to introduce the photo-sensitizers such as porphyrins to the C<sub>60</sub>-SWCNT(*n,m*) systems [48]. Three-layer supramolecular hybrids, C<sub>60</sub>-porphyrin-SWCNT are constructed from semiconducting (7,6)- and (6,5)-enriched SWCNTs.

##### 4.1. Molecular Structure and Energy Level Diagram

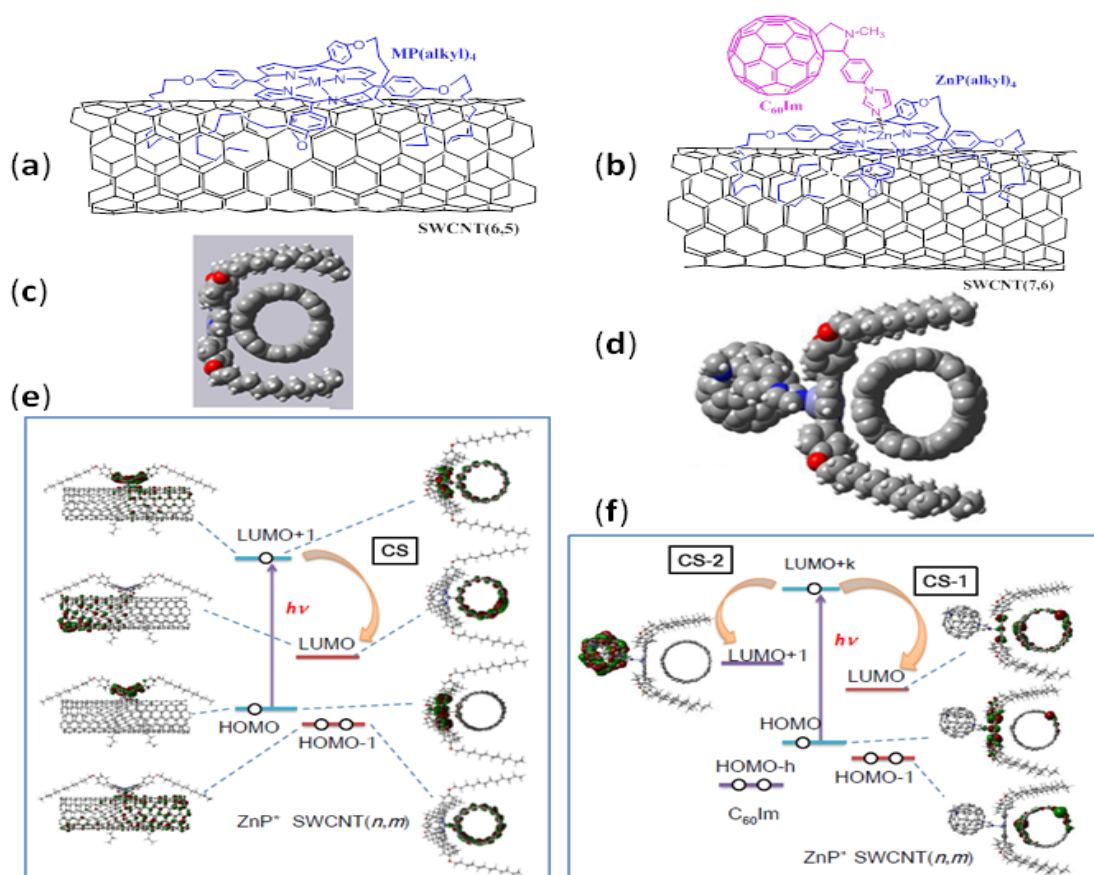
C<sub>60</sub>-porphyrin-SWCNT are constructed via a supramolecular method; first, long alkyl-chain substituted zinc porphyrins [tetrakis(4-dodecylphenyl) zinc porphyrins (ZnP(alkyl)<sub>4</sub>)] are mixed with SWCNT(*n,m*) in organic solvents. The ZnP(alkyl)<sub>4</sub>/SWCNT(*n,m*) hybrids are formed via the intermolecular alkyl- $\pi$  and  $\pi$ - $\pi$  interactions between the ZnP(alkyl)<sub>4</sub> and SWCNT(*n,m*) (Figure 7a). The formation of the ZnP(alkyl)<sub>4</sub>/SWCNT(*n,m*) hybrids are confirmed by the ZnP-fluorescence quenching in organic solution, suggesting occurrence of some photophysical events via <sup>1</sup>ZnP\*. On addition of C<sub>60</sub>Im to ZnP(alkyl)<sub>4</sub>/SWCNT(*n,m*), the C<sub>60</sub>Im $\rightarrow$ ZnP(alkyl)<sub>4</sub>/SWCNT(*n,m*) is formed via the coordination of C<sub>60</sub>Im to Zn ion (Figure 7b).

The intermolecular alkyl- $\pi$  and  $\pi$ - $\pi$  interactions and relative orientation of (ZnP(alkyl)<sub>4</sub>) on the surface of the semiconducting SWCNT can be visualized by performing DFT-MO calculations [41]. The edge-on view in Figure 7c shows an optimized structure of ZnP(alkyl)<sub>4</sub>/SWCNT(7,6), which clearly reveals that the ZnP  $\pi$ -plane is interacting with the surface of SWCNT(7,6); the distance between Zn and nearest carbon of the nanotube is found to be 2.1 Å, suggesting close association of the entities [48]. In addition, the alkyl chains also get their arms around the SWCNT(7,6).

The DFT-MO calculations of C<sub>60</sub>Im $\rightarrow$ ZnP(alkyl)<sub>4</sub>/SWCNT(7,6) show an optimized structure including C<sub>60</sub>Im-Zn axial coordination in addition to ZnP(alkyl)<sub>4</sub>/SWCNT(7,6) as shown in Figure 7d. The optimized distance Im $\rightarrow$ Zn  $\approx$  2.0 Å is similar to that reported earlier for C<sub>60</sub>Im $\rightarrow$ ZnP dyad computationally [41] and by X-ray structure [49]. On the other hand, the Zn-SWCNT distance increased to about 2.2 Å due to slight pulling of the Zn from the porphyrin macrocycle due to axial coordination.

The frontier orbitals for the two-layer ZnP(alkyl)<sub>4</sub>/SWCNT(7,6) system are also shown in Figure 7e. The majority of the HOMO is found on the ZnP entity, while the majority of the HOMO-1 is found on SWCNT(7,6). The LUMO was found on SWCNT(7,6), whereas the LUMO+1 is mainly located on the ZnP entity and partially on SWCNT(7,6). The pair of the HOMO and the LUMO corresponds to the most stable CS state as described to be ZnP<sup>•+</sup>(alkyl)<sub>4</sub>/SWCNT<sup>•-</sup>, whereas the pair of the HOMO and the LUMO+1 corresponds to the local excitation of the ZnP moiety. Consequently, excitation of the HOMO on the ZnP moiety rises an electron to the LUMO+1 within ZnP, from which the electron falls down to the LUMO of SWCNT(7,6), leaving a hole on the HOMO of ZnP, which results in ZnP<sup>•+</sup>(alkyl)<sub>4</sub>/SWCNT<sup>•-</sup> via <sup>1</sup>ZnP\*. If the excited state of SWCNT(7,6) is long enough, electron on the ZnP falls down to the vacant valence bond of SWCNT(7,6), giving ZnP<sup>•+</sup>(alkyl)<sub>4</sub>/SWCNT<sup>•-</sup> via <sup>1</sup>SWCNT\*, which will open the utilization of the NIR light for solar energy conversion.

**Figure 7.** Illustration of (a)  $\text{ZnP}(\text{alkyl})_4/\text{SWCNT}(n,m)$  and (b)  $\text{C}_{60}\text{Im}\rightarrow\text{ZnP}(\text{alkyl})_4/\text{SWCNT}(n,m)$  donor-acceptor nanohybrids. Edge-on view of the optimized structures at the B3LYP/3-21G(\*) level of (c)  $\text{ZnP}(\text{alkyl})_4$  interacting with  $\text{SWCNT}(7,6)$ ; and (d)  $\text{C}_{60}\text{Im}\rightarrow\text{ZnP}(\text{alkyl})_4/\text{SWCNT}$ ; C (gray), H (white), O (red), Zn (light blue), N (blue). For the calculation, 300 carbon atoms are used to build  $\text{SWCNT}(7,6)$ . Energy diagrams and electron distributions for (e)  $\text{ZnP}(\text{alkyl})_4/\text{SWCNT}(7,6)$  and (f)  $\text{C}_{60}\text{Im}\rightarrow\text{ZnP}(\text{alkyl})_4/\text{SWCNT}(7,6)$  calculated at the B3LYP/3-21G(\*) level (adopted from [48]).



The frontier orbitals for the three-layered  $\text{C}_{60}\text{Im}\rightarrow\text{ZnP}(\text{alkyl})_4/\text{SWCNT}(7,6)$  system are also shown in Figure 7f. The majority of the HOMO was found to be on ZnP entity, while the majority of the HOMO-1 was found to be on  $\text{SWCNT}(7,6)$ . The LUMO was found to be mainly localized on  $\text{SWCNT}(7,6)$  and partially on ZnP, whereas the LUMO+1 is fully located on the  $\text{C}_{60}$  entity. Pair of the HOMO and the LUMO corresponds to the CS state as described around  $\text{ZnP}^{*+}(\text{alkyl})_4/\text{SWCNT}^{*-}(7,6)$ , while having slight CT character. On the other hand, pair of the HOMO and the LUMO+1 corresponds to the CS state as  $\text{C}_{60}^{*-}\text{Im}\rightarrow\text{ZnP}^{*+}(\text{alkyl})_4$ . From the calculated MO energy levels,  $\text{ZnP}^{*+}(\text{alkyl})_4/\text{SWCNT}^{*-}$  is more stable than  $\text{C}_{60}^{*-}\text{Im}\rightarrow\text{ZnP}^{*+}(\text{alkyl})_4$ . When the electron on the LUMO+k ( $k>5$ ) of the ZnP moiety transfers to the LUMO of  $\text{SWCNT}$ ,  $\text{ZnP}^{*+}(\text{alkyl})_4/\text{SWCNT}^{*-}$  can be produced, which is denoted as CS-1 process in Figure 7f. On the other hand, when the electron on the LUMO+k on the ZnP moiety transfers to LUMO+1 of  $\text{C}_{60}$ ,  $\text{C}_{60}^{*-}\text{Im}\rightarrow\text{ZnP}^{*+}(\text{alkyl})_4$  can be produced as CS-2 process. The photo-excitation of  $\text{C}_{60}$  makes a hole on the HOMO-h ( $h > 10$ ), to which the HOMO-electron of ZnP flows down to the hole on the HOMO-h of  $\text{C}_{60}$ , giving  $\text{C}_{60}^{*-}\text{Im}\rightarrow\text{ZnP}^{*+}(\text{alkyl})_4$ , too. If the lifetime

of the excited state of SWCNT(7,6) is long enough, electron on the ZnP falls down to the vacant valence bond of SWCNT(7,6), giving  $\text{ZnP}^+(\text{alkyl})_4/\text{SWCNT}^{\bullet-}$  via  $^1\text{SWCNT}^*$ .

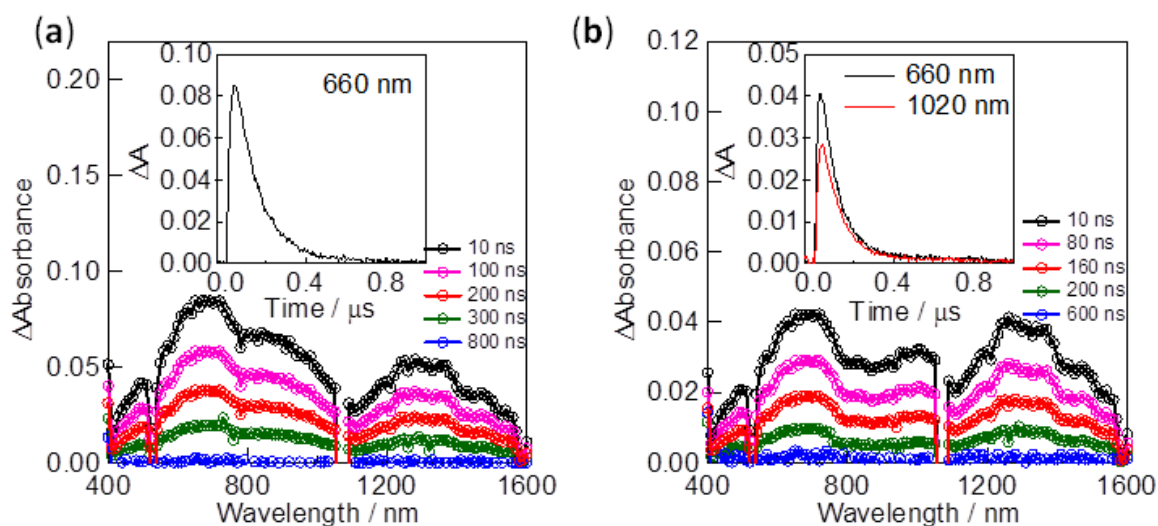
The energy diagrams can be constructed from the reported redox potentials for  $\text{C}_{60}\text{Im}$ ,  $\text{MP}(\text{alkyl})_4$ , SWCNT(6,5) and SWCNT(7,6) in DMF; they are almost the same as the HOMO and LUMO in Figures 7e and 7f, excepting the order of the HOMO of ZnP and the HOMO-1 of SWCNT, which makes electron transfer from the SWCNT to half-vacant ZnP, giving  $\text{ZnP}^+(\text{alkyl})_4/\text{SWCNT}^{\bullet-}$ .

Experimentally, these nanohybrids are characterized by TEM imaging, steady-state absorption and fluorescence spectra. The ZnP-fluorescence intensity is quenched on complexing with SWCNT( $n,m$ ) accompanying the shortening of the fluorescence lifetimes ( $\tau_F$ ) of the  $^1\text{ZnP}^*(\text{alkyl})_4$  in two-layered nanohybrids ( $\text{ZnP}(\text{alkyl})_4/\text{SWCNT}$ ), giving the  $k_{CS}$  values to produce  $\text{C}_{60}^{\bullet-}\text{Im} \rightarrow \text{ZnP}^+(\text{alkyl})_4$  in the range of  $(5-6) \times 10^9 \text{ s}^{-1}$ . On further addition of  $\text{C}_{60}\text{Im}$ , the ZnP-fluorescence intensity is additionally quenched. The time-resolved fluorescence studies of three-layered nanohybrids ( $\text{C}_{60}\text{Im} \rightarrow \text{ZnP}(\text{alkyl})_4/\text{SWCNT}$ ) gave further shortening of the fluorescence lifetimes ( $\tau_F$ ) of the  $^1\text{ZnP}^*(\text{alkyl})_4$ , giving the total  $k_{CS}$  values are  $(8-9) \times 10^9 \text{ s}^{-1}$ , from which the  $k_{CS}$  values from  $^1\text{ZnP}^*(\text{alkyl})_4$  to  $\text{C}_{60}\text{Im}$  to form  $\text{C}_{60}^{\bullet-}\text{Im} \rightarrow \text{ZnP}^+(\text{alkyl})_4$  are evaluated to be  $(2-3) \times 10^9 \text{ s}^{-1}$ .

#### 4.2. Evidence for Photoinduced Charge Separation

Information for the CS products and kinetics of CR are obtained from the transient absorption spectral studies. Figure 8a shows the nanosecond transient absorption spectra of the  $\text{ZnP}(\text{alkyl})_4/\text{SWCNT}(n,m)$  nanohybrids observed using a 532 nm laser light to excite selectively the  $\text{ZnP}(\text{alkyl})_4$  moiety in Ar-saturated DMF. Transient absorption bands appear in the wide region from the Visible region to the NIR region. The absorption bands at 650 nm can be attributed to the  $\text{ZnP}^+(\text{alkyl})_4$  moiety. Since the NIR-band intensities decrease in almost the same rates with the visible  $\text{ZnP}^+$  bands, they can be thought as a pair of the RIP. Thus, the absorption bands at 900, 1,250 and 1,380 nm can be ascribed to the  $\text{SWCNT}(6,5)^{\bullet-}$ .

**Figure 8.** Nanosecond transient absorption spectra observed by 532 nm laser-light pulse irradiation; (a)  $\text{ZnP}(\text{alkyl})_4/\text{SWCNT}(6,5)$  in Ar-saturated DMF and (b)  $\text{C}_{60}\text{Im} \rightarrow \text{ZnP}(\text{alkyl})_4/\text{SWCNT}(6,5)$  in Ar-saturated *o*-DCB: Inset, Absorption-time profiles (adopted from [48]).



In the case of  $\text{ZnP}(\text{alkyl})_4/\text{SWCNT}(7,6)$ , the absorption bands of the  $\text{ZnP}^{\bullet+}(\text{alkyl})_4$  moiety similarly appear at 650 nm in addition to the broad NIR bands in 1,200–1,500 nm, whereas the 900 nm band is not found, suggesting that the broad 900-nm band in Figure 8a is one of the absorption bands of  $\text{SWCNT}(6,5)^{\bullet-}$ . These assignments afford evidence for the occurrence of the CS process in the supramolecular nanohybrid to form  $\text{ZnP}^{\bullet+}(\text{alkyl})_4/\text{SWCNT}^{\bullet-}$ . The time profile of  $\text{ZnP}^{\bullet+}(\text{alkyl})_4$  at 650 nm is shown in the inset of Figure 8a. The rise of  $\text{ZnP}^{\bullet+}(\text{alkyl})_4$  is fast within the 6-ns laser light pulse, which corresponds to the fast CS rate within ca. 200 ps as estimated from the fluorescence lifetimes of  $\text{ZnP}(\text{alkyl})_4$ . The decay of the  $\text{ZnP}^{\bullet+}(\text{alkyl})_4$  mostly persists until about 300–400 ns; from the first-order fitting to the decay curve, the  $k_{\text{CR}}$  value is evaluated to be  $6 \times 10^6 \text{ s}^{-1}$ . From the  $k_{\text{CR}}$  value, the lifetimes of  $\text{ZnP}^{\bullet+}(\text{alkyl})_4/\text{SWCNT}^{\bullet-}$  ( $\tau_{\text{RIP}}$ ) can be calculated to be 140 ns in DMF. Similar values are obtained from the decay-time profiles at the NIR bands. The  $k_{\text{CS}}/k_{\text{CR}}$  ratios, which are proportional to the charge stabilization, are 600~700 for  $\text{ZnP}(\text{alkyl})_4/\text{SWCNT}(n,m)$ ; these values are considerably large enough to prove the charge stabilization in DMF.

Following is a comparison of kinetics of CS and CR with other kinds of  $\text{ZnP}/\text{SWCNTs}$  using  $\text{SWCNT}(6,5)$  and  $\text{SWCNT}(7,6)$ . In the case of nanohybrids formed using  $\text{ZnP}$  covalently bonded with four pyrene moieties [ $(\text{ZnP}(\text{Pyr})_4)$ ], which attaches on  $\text{SWCNT}$  surfaces,  $k_{\text{CS}} = (3\text{--}4) \times 10^9 \text{ s}^{-1}$ ,  $k_{\text{CR}} = (2\text{--}3) \times 10^7 \text{ s}^{-1}$  ( $\tau_{\text{RIP}} = 40\text{--}50 \text{ ns}$ ) and  $k_{\text{CS}}/k_{\text{CR}} \approx 150$  [50]. For nanohybrids formed by  $\text{ZnP}$  coordinately bonded to an axial imidazole ligand tethered to a pyrene moiety ( $\text{PyrIm} \rightarrow \text{ZnP}$ ), in which the Pyr moiety attaches to the  $\text{SWCNT}$  surfaces ( $\text{SWCNT}/\text{PyrIm} \rightarrow \text{ZnP}$ ),  $k_{\text{CS}} = (5\text{--}6) \times 10^9 \text{ s}^{-1}$ ,  $k_{\text{CR}} = (1\text{--}2) \times 10^7 \text{ s}^{-1}$  ( $\tau_{\text{RIP}} = 50\text{--}70 \text{ ns}$ ) and  $k_{\text{CS}}/k_{\text{CR}} \approx 500$  [51]. For nanohybrids formed via ion-pairing, *viz.*,  $\text{ZnPs}$  with charged groups at the macrocycle periphery attached to the opposite charged pyrene moiety on  $\text{SWCNT}$  surface,  $k_{\text{CS}} = (2\text{--}8) \times 10^9 \text{ s}^{-1}$ ,  $k_{\text{CR}} = (0.5\text{--}0.8) \times 10^7 \text{ s}^{-1}$  ( $\tau_{\text{RIP}} = 50\text{--}70 \text{ ns}$ ) and  $k_{\text{CS}}/k_{\text{CR}} \approx 500$  [52]. For  $\text{ZnP}$  with crown ethers complexed with ammonium ion-pyrene on  $\text{SWCNT}$  surface ( $\text{SWCNT}/\text{Pyr-NH}_3^+:\text{CrownZnP}$ ),  $k_{\text{CS}} = (1\text{--}5) \times 10^9 \text{ s}^{-1}$ ,  $k_{\text{CR}} = (0.9\text{--}1.0) \times 10^7 \text{ s}^{-1}$  ( $\tau_{\text{RIP}} = 90\text{--}110 \text{ ns}$ ) and  $k_{\text{CS}}/k_{\text{CR}} \approx 500$  [53]. Clearly, these high  $k_{\text{CS}}/k_{\text{CR}}$  ratios are favorable for further electron mediation and hole-shift, opening wider applications for solar photovoltaic cells and solar  $\text{H}_2$ -evolution.

On addition of  $\text{C}_{60}\text{Im}$  in *o*-DCB, new band appears at 1,020 nm in addition to the Visible and NIR bands due to  $\text{ZnP}^{\bullet+}(\text{alkyl})_4/\text{SWCNT}(6,5)^{\bullet-}$  as shown in Figure 8b. The 1,020-nm band is a diagnostic evidence of  $\text{C}_{60}^{\bullet-}\text{Im}$ , confirming the formation of  $\text{C}_{60}^{\bullet-}\text{Im} \rightarrow \text{ZnP}^{\bullet+}(\text{alkyl})_4/\text{SWCNT}(6,5)$  and  $\text{C}_{60}^{\bullet-}\text{Im} \rightarrow \text{ZnP}(\text{alkyl})_4/\text{SWCNT}(6,5)^{\bullet+}$  via  $^1\text{ZnP}^*(\text{alkyl})_4$ . The decay of  $\text{C}_{60}^{\bullet-}\text{Im}$  is almost the same as those of the visible band [ $\text{ZnP}^{\bullet+}(\text{alkyl})_4$ ] and the NIR bands [ $\text{SWCNT}(6,5)^{\bullet+}$ ].

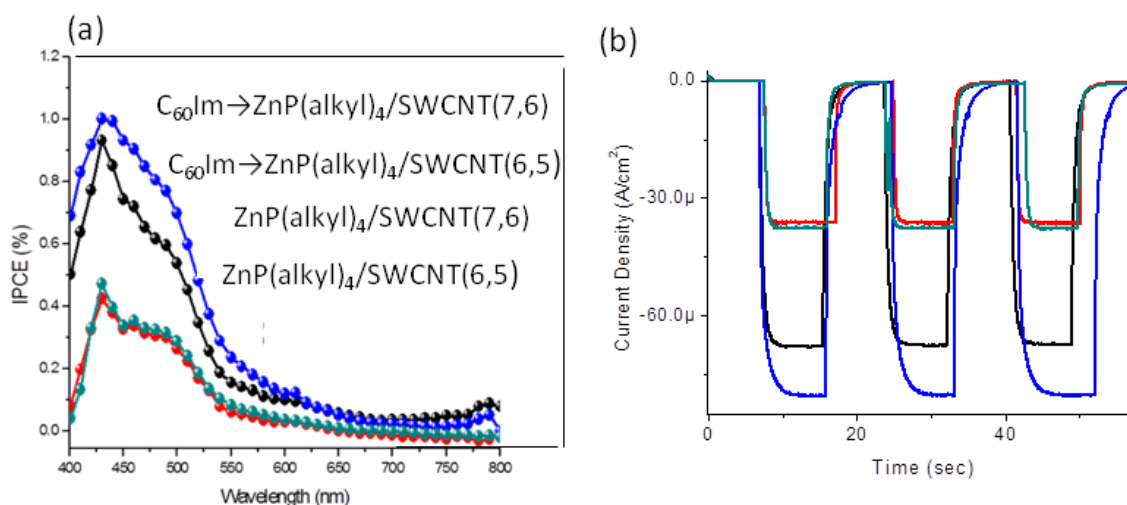
Similar transient absorption spectral changes are observed for  $\text{C}_{60}\text{Im} \rightarrow \text{ZnP}(\text{alkyl})_4/\text{SWCNT}(7,6)$ . The  $k_{\text{CR}}$  values evaluated from the 660 nm band and 1,020 nm bands become faster on coordination of  $\text{C}_{60}\text{Im}$ , giving the larger  $k_{\text{CR}}$  values and shorter  $\tau_{\text{RIP}}$  values; the faster CR process may occur between  $\text{C}_{60}^{\bullet-}\text{Im} \rightarrow \text{ZnP}^{\bullet+}(\text{alkyl})_4$ . A slight decrease in the  $k_{\text{CS}}/k_{\text{CR}}$  ratios are induced on coordination of the  $\text{C}_{60}\text{Im}$  due to structural and electronic reasons. It may be mentioned here that the decay curves of  $\text{ZnP}^{\bullet+}(\text{alkyl})_4/\text{SWCNT}^{\bullet-}$  and  $\text{C}_{60}^{\bullet-}\text{Im} \rightarrow \text{ZnP}^{\bullet+}(\text{alkyl})_4/\text{SWCNT}$  followed the trend expected for intramolecular process and not intermolecular bimolecular process occurring on several  $\mu\text{s}$  time scales, suggesting that both  $\text{ZnP}(\text{alkyl})_4/\text{SWCNT}$  and  $\text{C}_{60}\text{Im} \rightarrow \text{ZnP}(\text{alkyl})_4/\text{SWCNT}$  are intact in solution during photochemical investigations [48].



### 4.3. Photoelectrochemical Studies

The observed photo-induced CS processes in the investigated nano-hybrids prompted us to perform photoelectro-chemical studies to visualize their ability to convert the light energy into the electricity. For this, the  $\text{ZnP(alkyl)}_4/\text{SWCNT}(n,m)$  and  $\text{C}_{60}\text{Im}\rightarrow\text{ZnP(alkyl)}_4/\text{SWCNT}(n,m)$  nano-hybrids are drop-coated on  $\text{SnO}_2$ -modified fluorine doped tin oxide (FTO) electrode surface, since this  $\text{SnO}_2$ -modified FTO is known to yield high IPC with suitable porphyrins or phthalocyanines [54,55]. The photocurrent action spectra show IPCE values in the wide wavelength region covering up to 700 nm with maximum near 420 nm, which tracks the absorption spectra of the ZnP moiety (Figure 9a). The switching responses are shown in Figure 9b, in which quick response and light-stability are revealed. The higher IPCE values at 420 nm are obtained for  $\text{C}_{60}\text{Im}\rightarrow\text{ZnP(alkyl)}_4/\text{SWCNT}(n,m)$  than for  $\text{ZnP(alkyl)}_4/\text{SWCNT}(n,m)$  by a factor of ca. 2, suggesting that  $\text{C}_{60}\text{Im}$  contributes to the additional CS process via  $^1\text{C}_{60}\text{Im}^*$ . In  $\text{C}_{60}\text{Im}\rightarrow\text{ZnP(alkyl)}_4/\text{SWCNT}(n,m)$ ,  $\text{SWCNT}(7,6)$  is slightly higher than  $\text{SWCNT}(6,5)$ .

**Figure 9.** (a) IPCE spectra and (b) light-current switching curves for the photoelectrochemical solar cells of  $\text{ZnP(alkyl)}_4/\text{SWCNT}$  and  $\text{C}_{60}\text{Im}\rightarrow\text{ZnP(alkyl)}_4/\text{SWCNT}$  coated on  $\text{SnO}_2$ -modified FTO electrode surface in acetonitrile containing 0.5 M LiI and 0.1 M  $\text{I}_2$  ( $\text{I}^-/\text{I}_3^-$ ) mediator. Red curve;  $\text{ZnP(alkyl)}_4/\text{SWCNT}(6,5)$ , black curve;  $\text{C}_{60}\text{Im}\rightarrow\text{ZnP(alkyl)}_4/\text{SWCNT}(6,5)$ , green curve;  $\text{ZnP(alkyl)}_4/\text{SWCNT}(7,6)$ , and blue curve;  $\text{C}_{60}\text{Im}\rightarrow\text{ZnP(alkyl)}_4/\text{SWCNT}(7,6)$  (adopted from [47]).



In summary of this section, in the case of  $\text{C}_{60}\text{Im}\rightarrow\text{ZnP}/\text{SWCNT}(n,m)$  nanohybrids, further charge-separation from  $^1\text{ZnP}^*$  to the  $\text{C}_{60}$  moiety takes place, in addition to from  $^1\text{ZnP}^*$  to the  $\text{SWCNT}(n,m)$ , giving larger  $k_{\text{CS}}$  values than those of  $\text{ZnP(alkyl)}_4/\text{SWCNT}(n,m)$ . Photoelectrochemical studies of these nanohybrids modified on  $\text{SnO}_2/\text{FTO}$  electrodes revealed photo-current generation. Higher IPCE values are obtained for  $\text{C}_{60}\text{Im}\rightarrow\text{ZnP}/\text{SWCNT}(n,m)$  than those of  $\text{ZnP}/\text{SWCNT}(n,m)$ .

## 5. Conclusions

In the present article, we have shown that supramolecular methods are simple and versatile approaches to construct donor-acceptor hybrids involving fullerenes to induce photoinduced events.

Occurrence of photoinduced electron transfer is established by photophysical and photochemical methods such as fluorescence lifetime measurements and transient absorption spectral measurements. These phenomena are interpreted on the basis of the molecular orbital energies and electron-distributions calculated using DFT methods. The MO calculations of the supramolecular systems including the SWCNT( $n,m$ ) open a molecular paradigm of size-selected SWCNT from the material of SWCNT mixtures. Practical applications, especially in the area of photovoltaic cells and photocatalysis, including electron pooling and H<sub>2</sub> production, are promising. Such applications are currently being explored in our laboratories.

### Acknowledgments

We are thankful to K. Miyazawa and T. Wakahara of NIMS (Tsukuba, Japan) for their useful support during these studies. We also thank our collaborators at Tohoku University (Sendai, Japan), JAIST (Nomi, Japan), Wichita State University (Wichita, KS, USA) and University of North Texas (Denton, TX, USA). This work was financially supported by the National Science Foundation (Grant No. 1110942 to FD).

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