# 77. Remarks on Hadamard's Variation of Eigenvalues of the Laplacian 

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§1. Introduction. The study in this note is a continuation of our previous paper [6]. Let $\Omega$ be a bounded domain in $\mathrm{R}^{n}(n \geq 2)$ with $\mathcal{C}^{\infty}$ boundary $\gamma$. Let $\rho(x)$ be a smooth function on and $\nu_{x}$ be the exterior unit normal vector at $x \in \gamma$. For any sufficiently small $\varepsilon \geq 0$, let $\Omega_{\varepsilon}$, be the bounded domain whose boundary $\gamma_{s}$ is defined by $\gamma_{\varepsilon}=\left\{x+\varepsilon \rho(x) \nu_{x}\right.$; $x \in \gamma\}$. Let $U_{s}(x, y, t)$ be the Green kernel of the heat equation in $\Omega_{\text {。 }}$ with the Dirichlet boundary condition on $\gamma_{s}$. Let $T_{r}(t ; \varepsilon)$ be the trace of $U_{s}$ on $\Omega_{s}$. When $t$ tends to zero, we have the asymptotic expansion $T_{r}(t ; \varepsilon) \sim \sum_{j=0}^{\infty} a_{n-j}(\varepsilon)(\sqrt{t})^{-n+j}$ which was given by MinakshisundarumPleijel [5]. In [6], the author gave the asymptotic expansion $\delta T_{r}(t)$ $\sim \sum_{j=0}^{\infty} b_{n-j}(\sqrt{t})^{-n+j}$ near $t=0$ of the variational term $\delta T_{r}(t)$ of the trace which was defined by $\delta T_{r}(t)=\lim _{\varepsilon \rightarrow 0} \varepsilon^{-1}\left(T_{r}(t ; \varepsilon)-T_{r}(t ; 0)\right)$. We proposed the following problem ( E$)_{k}^{n}$ in [6] and gave an affirmative answer for the case $k=0$.

Problem (E) $)_{k^{+}} \quad$ Can we say that the following is valid? $(\mathrm{E})_{k}^{n}$

$$
b_{n-k}=\lim _{\varepsilon \rightarrow 0} \varepsilon^{-1}\left(a_{n-k}(\varepsilon)-a_{n-k}(0)\right) .
$$

In this paper, we shall prove the following
Theorem 1. ( E$)_{1}^{n}$ is valid for any $n \geq 2$.
The aims of this note are verification of Theorem 1 and an application of Theorem 1 to some eigenvalue problem which will be stated in this section.

We now mention the following
Problem (Q). Characterize the bounded domain $\Omega$ with smooth boundary $\gamma$ having the following property.
(I) For any $\rho(z) \in \mathcal{C}^{\infty}(\gamma)$ such that $\int_{r} \rho(z) d \sigma_{z}=0$, we have $\delta \lambda_{1}=0$, where $\delta \lambda_{1}$ is the variational term of the first eigenvalue $\lambda_{1}<0$ of the Laplacian with the Dirichlet condition. Here d $\sigma_{z}$ denotes the surface element of $\gamma$ at $z$.

The condition $\int_{r} \rho(z) d \sigma_{z}=0$ means that the perturbation of domain we considered preserves the volume of domains infinitesimally. We
call the domain $\Omega$ satisfying the property (I) stationary domain. By the theorem of Hadamard-Garabedian-Schiffer we conclude that $\Omega$ is stationary if and only if

$$
\left\{\begin{align*}
\Delta \varphi_{1}(x)-\lambda_{1} \varphi_{1}(x)=0 & \text { in } \Omega  \tag{1}\\
\varphi_{1}(x)=0 & \text { on } \gamma \\
\frac{\partial \varphi_{1}}{\partial \nu_{x}}(x)=C & \text { on } \gamma,
\end{align*}\right.
$$

where $\varphi_{1}$ is the normalized eigenfunction of the Laplacian with the Dirichlet condition and $C$ is a constant. See [4] and [2].

The classical Faber-Krahn theorem states that the two dimensional domain $\Omega$ with the fixed area $A$ which maximize the first eigenvalue $\lambda_{1}(\Omega)<0$ of the Laplacian with the Dirichlet condition is the disk of radius $\left(A \pi^{-1}\right)^{1 / 2}$. In connection with this fact, we conjecture that any $n$-dimensional stationary domain is an open $n$-ball. For the $n$-ball with radius $p$, we have (1), so our problem may be restated as follows.
$(Q)^{\text {bis }}$ : Are the following two statements equivalent?
$(\mathrm{Q})_{1}: \Omega$ is an open $n$-ball.
$(\mathrm{Q})_{2}: \quad \Omega$ is an $n$-dimensional stationary domain.
Even if $n=2$, the implication from $(\mathrm{Q})_{2}$ to $(\mathrm{Q})_{1}$ is not trivial. It should be remarked that in general there is a gap between the concept of stationary value and that of maximum value in variational problem. In this paper, we give a partial answer to the problem ( $Q)^{\text {bis }}$ by using Theorem 1.

The domain satisfying the following property will be called to be $T$-stationary :
(II) For any $\rho(z) \in \mathcal{C}^{\infty}(\gamma)$ such that $\int_{\gamma} \rho(z) d \sigma_{z}=0$ and for any $t>0$, we have $\delta T_{r}(t)=0$.

By Theorem 3 in [6], we know that any $T$-stationary domain is stationary. We have the following

Theorem 2. Assume that $\Omega$ is T-stationary. Then for $n=2 \Omega$ is a disk, and for $n \geq 3$ every component of $\gamma$ is a hypersurface of constant mean curvature.

In §2, we consider some geometry of hypersurfaces and the asymptotics of heat equation. In $\S 3$, we give proofs of Theorems 1 and 2.
§2. Geometry of hypersurfaces and asymptotic expansion. We proved in [9] that

$$
\begin{equation*}
\delta T_{r}(t)=\left.t \int_{r} \frac{\partial^{2} U(y, w, t)}{\partial \nu_{y} \partial \nu_{w}}\right|_{y=w=z} \rho(z) d \sigma_{z} . \tag{2}
\end{equation*}
$$

Here we abbreviate $U_{0}(x, y, t)$ in $\S 1$ as $U(x, y, t)$. In [8], we proved the following proposition by using the hard calculus of pseudo-differential operators such as [10].

Proposition 1. If $t$ tends to zero, then

$$
\begin{equation*}
\left.t \cdot \frac{\partial^{2} U(y, w, t)}{\partial \nu_{y} \partial \nu_{w}}\right|_{y=w=z} \sim \sum_{k=0}^{\infty} B_{n-k}(z) t^{-(n-k) / 2} \tag{3}
\end{equation*}
$$

for any $z \in \gamma$. Here $B_{n-k}(z) \in \mathcal{C}^{\infty}(\gamma)$.
We fix an arbitrary point $z$ on $\gamma$. Without loss of generalities we can assume that $z$ is the origin of $\mathrm{R}^{n}$. We take the $z_{n}$-axis as being coincident with the direction of the interior normal and the hyperplane $z_{n}=0$ as being coincident with the tangent hyperplane of $\gamma$ at $z$. We take an orthonormal coordinate system $z^{\prime}=\left(z_{1}, \cdots, z_{n-1}\right)$ on this hyperplane. Then $\gamma$ can be locally written as $z_{n}=\theta\left(z_{1}, \cdots, z_{n-1}\right)$. Here $\theta \in C^{\infty}\left(\mathbf{R}^{n-1}\right)$ and it has the Taylor expansion $\theta\left(z_{1}, \cdots, z_{n-1}\right)=\sum_{|\alpha| \geq 2}^{\infty} \omega_{\alpha} z^{\prime \alpha}$. In [8], we also proved the following proposition by using the calculus of pseudo-differential operators. See also [7].

Proposition 2. For any $k$, there exists a constant $w(k)$ such that $B_{n-k}(z)$ is a polynomial of the variables $\left\{\omega_{\alpha}\right\}_{2 \leq|\alpha| \leq w(k)}$ whose coefficients depend only on $k$.

We write $B_{n-k}(z)=B_{n-k}\left(\left\{\omega_{\alpha}\right\}\right)$. Following the idea of [1] and [3], we use the notion of the weight of polynomials. We give the weight $|\alpha|-1$ to the variable $\omega_{\alpha}$, and we give the weight $\sum_{j=1}^{s} \beta_{j}\left(\left|\alpha_{j}\right|-1\right)$ to the monomial $\prod_{j=1}^{S}\left(\omega_{\alpha_{j}}\right)^{\beta_{j}}$. If $P$ is the sum of the monomials of the same weight, we say that $P$ is homogeneous. Now we prove the following

Theorem 3. $B_{n-k}(z) \equiv B_{n-k}\left(\left\{\omega_{\alpha}\right\}\right)$ is the homogeneous polynomial of the weight $k$.

Proof. Fix $q>0$. We take a new coordinate system $\tilde{z}$ as follows; $\tilde{z}_{j}=q z_{j}(j=1, \cdots, n)$. Then $\gamma$ can be locally written as $\tilde{z}_{n}=\tilde{\theta}\left(\tilde{z}_{1}, \cdots, \tilde{z}_{n-1}\right)$ $=\tilde{\theta}\left(\tilde{z}^{\prime}\right)=\sum_{|\alpha| \geq 2} \tilde{\omega}_{\alpha} \tilde{z}^{\prime \alpha}$. We have the relation

$$
\begin{equation*}
\tilde{\omega}_{\alpha}=q^{-(|\alpha|-1)} \omega_{\alpha} . \tag{4}
\end{equation*}
$$

On the other hand, the fundamental solution $\tilde{U}(\tilde{x}, \tilde{y}, t)$ of the heat equation $\frac{\partial}{\partial t}-\Delta_{x x}$ with the Dirichlet boundary condition in new coordinates is related to $U(x, y, t)$ by

$$
\begin{equation*}
\tilde{U}(\tilde{x}, \tilde{y}, t)=q^{-n} U\left(x, y, q^{-2} t\right) \tag{5}
\end{equation*}
$$

where $\Delta_{\tilde{x}}$ is the Laplacian in $x$-coordinates. It is easy to see

$$
\begin{equation*}
\left.\frac{\partial^{2} U(\tilde{y}, \tilde{w}, t)}{\partial \nu_{\tilde{y}} \partial \nu_{\tilde{w}}}\right|_{\tilde{y}=\tilde{w}=\tilde{z}}=\left.q^{-n-2} \frac{\partial^{2} U\left(y, w, q^{-2} t\right)}{\partial \nu_{y} \partial \nu_{w}}\right|_{y=w=z}, \tag{6}
\end{equation*}
$$

where $\frac{\partial}{\partial \nu_{z}}$ is the derivative along the exterior normal vector. We compare the asymptotic expansion of both sides of (6). Then we get

$$
\begin{equation*}
\boldsymbol{B}_{n-k}\left(\left\{\tilde{\omega}_{\alpha}\right\}\right)=q^{-k} \boldsymbol{B}_{n-k}\left(\left\{\omega_{\alpha}\right\}\right) . \tag{7}
\end{equation*}
$$

By (4) and (7), we obtain Theorem 3.
We transform $z^{\prime}=\left(z_{1}, \cdots, z_{n-1}\right)$ to another orthonormal basis $\hat{z}^{\prime}=\left(\hat{z}_{1}, \cdots, \hat{z}_{n-1}\right)=\left(z_{1}, \cdots, z_{n-1}\right) V$. Here $V \in 0(n-1, \mathbf{R})$. Then $\gamma$ can be locally written as $z_{n}=\hat{\theta}\left(\hat{z}_{1}, \cdots, \hat{z}_{n-1}\right)=\sum_{|\alpha| \geq 2} \hat{\omega}_{\alpha} \hat{z}^{\prime \alpha}$. We can represent $\left\{\hat{\omega}_{\alpha}\right\}$ explicitly by $\left\{\omega_{\alpha}\right\}$ and $V \in 0(n-1, \mathrm{R})$. It should be remarked that (8)

$$
B_{n-k}\left(\left\{\hat{\omega}_{\alpha}\right\}\right)=q^{-k} B_{n-k}\left(\left\{\omega_{\alpha}\right\}\right),
$$

since both sides are equal to $B_{n-k}(z)$ which does not depend on the choice of orthonormal basis on the tangent hyperplane of $\gamma$ at $z$.

We have the following
Proposition 3. $B_{n}(z)=e_{n}$ on $\gamma$, and $B_{n-1}(z)=\beta H_{1}(z)$ on $\gamma$. Here $H_{1}(z)$ denotes the first mean curvature of $\gamma$ at $z$, and $e_{n}, \beta$ are constants depending only on $n$.

Proof. Since $B_{n}(z)$ is homogeneous of weight 0 , it must be a constant. Next we study $B_{n-1}(z)$. If we transform $z^{\prime}$ to $z^{\prime}=z^{\prime} V$, where $V \in 0(n-1, \mathrm{R})$, then the following relation holds.

$$
\begin{equation*}
\hat{\omega}_{k h}=\sum_{i, j=1}^{n-1} V_{k i} \omega_{i j} V_{i n} . \tag{9}
\end{equation*}
$$

The above relation can be rewritten as

$$
\begin{equation*}
\hat{\omega}=V \omega V^{-1}, \tag{10}
\end{equation*}
$$

where $\omega$ and $\hat{\omega}$ is the matrix with the components $\left\{\omega_{i j}\right\}$ and $\left\{\hat{\omega}_{k h}\right\}$ respectively. Since $B_{n-1}(z)$ is homogeneous of weight 1, it must be written as $B_{n-1}(z)=\sum_{i, j=1}^{n-1} s_{i j} \omega_{i j}$, where $s_{i j}$ is a constant which depend only on $n$. So we have also $B_{n-1}(z)=\sum_{i, j=1}^{n-1} s_{i j} \hat{\omega}_{i j}$ for any $\hat{\omega}_{i j}$ with respect to any other orthonormal coordinates $\hat{z}^{\prime}$. Finally by the theory of orthogonal invariants such as [11], we conclude that $s_{i j}=\tau \delta_{i j}$ where $\tau$ is a constant and $\delta_{i j}$ is the Kronecker delta. It is well known that $H_{1}(z)$ $=2(n-1)^{-1} \sum_{i=1}^{n-1} \omega_{i i}$. The proof is over.
§3. Proofs of Theorems 1 and 2. The coefficients $a_{n}(\varepsilon)$ and $a_{n-1}(\varepsilon)$ in the asymptotic expansion of $T_{r}(t ; \varepsilon)$ are represented as $a_{n}(\varepsilon)$ $=C_{n}\left|\Omega_{\mathrm{s}}\right|, a_{n-1}(\varepsilon)=C_{n-1}\left|\gamma_{\mathrm{c}}\right|$. Here $C_{n}$ and $C_{n-1}$ are non zero constants which depend only on $n$. And here $\left|\Omega_{\varepsilon}\right|$ denotes the volume of $\Omega_{\varepsilon}$ and $\left|\gamma_{c}\right|$ denotes the area of $\gamma_{c}$. Therefore we have to show

$$
\begin{equation*}
b_{n}=C_{n} \int_{r} \rho(z) d \sigma_{z} \tag{11}
\end{equation*}
$$

and

$$
\begin{equation*}
b_{n-1}=(n-1) C_{n-1} \int_{r} H_{1}(z) \rho(z) d \sigma_{z} . \tag{12}
\end{equation*}
$$

By Theorems 1 and 2 in [6], we have

$$
b_{n-k}=\int_{r} B_{n-k}(z) \rho(z) d \sigma_{z} .
$$

So we can restate (E) $)_{k}^{n}(k=0,1)$ as follows.
(E) ${ }_{0}^{n}$ (E) ${ }_{1}^{n}$

$$
\begin{gathered}
e_{n}=C_{n}, \\
e_{n-1}(z)=(n-1) C_{n-1} H_{1}(z) .
\end{gathered}
$$

Let $B_{R}$ be the open ball with radius $R$ centered at the origin. Let $0 \geq \lambda_{1}(R) \geq \lambda_{2}(R) \geq \cdots$ be the eigenvalues of the Laplacian with the Dirichlet boundary condition at $\partial B_{R}$. We arrange them according to their multiplicities. We have $\lambda_{i}(R)=\lambda_{i}(1) R^{-2}$. Put $T_{r}\left(t \mid B_{R}\right)=\sum_{j=1}^{\infty} e^{\lambda_{j}(R) t}$. And put $\rho(z)=1$ on $\gamma=\partial B_{R}$. Then we have

$$
\delta T_{r}(t)=\frac{\partial}{\partial R} T_{r}\left(t \mid B_{R}\right)=-2 t R^{-1} \frac{\partial}{\partial t} T_{r}\left(t \mid B_{R}\right) .
$$

On the other hand, by the calculus of pseudo-differential operators we can get the following proposition. We need nontrivial calculations to prove it. See [8].

Proposition 4. When $t$ tends to zero, $\frac{\partial}{\partial t} T_{r}\left(t \mid B_{R}\right)$ has the asymptotic expansion

$$
\frac{\partial}{\partial t} T_{r}\left(t \mid B_{R}\right) \sim \sum_{k=0}^{\infty} M_{n-k}(R) t^{-(n+2-k) / 2} .
$$

By Proposition 4, we can differentiate with respect to $t$ the asymptotic expansion $T_{r}\left(t \mid B_{R}\right) \sim \sum_{k=0}^{\infty} a_{n-k}(\sqrt{t})^{-n+k}$ term by term. Therefore we have
(13) $\delta T_{r}(t) \sim n a_{n} R^{-1}(\sqrt{t})^{-n}+\cdots+(n-k) a_{n-k} R^{-1}(\sqrt{t})^{-n+k}+\cdots$. On the other hand, we have

$$
\begin{equation*}
\delta T_{r}(t) \sim e_{n}|\gamma|(\sqrt{t})^{-n}+\beta \int_{T} H_{1}(z) d \sigma_{z} \cdot(\sqrt{t})^{-n+1}+0\left(t^{-n / 2+1}\right) . \tag{14}
\end{equation*}
$$

By (13) and (14), we have $|\gamma| e_{n}=n a_{n} R^{-1}$ and $\beta|\gamma| R^{-1}=(n-1) a_{n-1} R^{-1}$. Since $a_{n}=C_{n}\left|B_{R}\right|, a_{n-1}=C_{n-1}\left|\partial B_{R}\right|$, we get Theorem 1 .

By Theorem 1 we have the following
Corollary. If tends to zero, then

$$
\begin{align*}
\delta T_{r}(t)= & C_{n} \int_{r} \rho(z) d \sigma_{z} \cdot t^{-(n+2) / 2}  \tag{15}\\
& +(n-1) C_{n-1} \cdot \int_{r} H_{1}(z) \rho(z) d \sigma_{z} \cdot t^{-(n+1) / 2}+0\left(t^{-n / 2}\right),
\end{align*}
$$

where $C_{n} \neq 0, C_{n-1} \neq 0$.
Proof of Theorem 2. The proof of Theorem 2 is obvious for $n \geq 3$, since we have (15). Now we study the case $n=2$. By (15) we can conclude that $\gamma$ consists of finite disjoint union of circles. Since we know that any $T$-stationary domain is stationary by Theorem 3 in [6], we have (1). Let $\gamma$ be the largest circle contained in $\gamma$. By the uniqueness theorem of Holmgren we conclude that the level set $\gamma(s)$ $=\left\{z ; \varphi_{1}(z)=s, z \in \Omega\right\}$ consists of the finite disjoint union of circles with the same center as $\gamma$. It is well known that $\varphi_{1}(z)$ does not take zero in
$\Omega$. Summing up these facts we conclude that $\Omega$ must be the disk or the annulus. For the annulus $\left\{z \in \mathrm{R}^{2} ; r_{1}<|z|<r_{2}\right\}, \varphi_{1}(z)$ can be calculated explicitly. If $r_{1}>0$, then $\varphi_{1}(z)$ with respect to this domain does not satisfy (1). The proof is over.

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