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## $\Lambda_c^+$ production and baryon-to-meson ratios in pp and p–Pb collisions at $\sqrt{s_{\text{NN}}} = 5.02$ TeV at the LHC

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### Abstract

The prompt production of the charm baryon  $\Lambda_c^+$  and the  $\Lambda_c^+/D^0$  production ratios were measured at midrapidity with the ALICE detector in pp and p–Pb collisions at  $\sqrt{s_{\text{NN}}} = 5.02$  TeV. These new measurements show a clear decrease of the  $\Lambda_c^+/D^0$  ratio with increasing transverse momentum ( $p_T$ ) in both collision systems in the range  $2 < p_T < 12$  GeV/ $c$ , exhibiting similarities with the light-flavour baryon-to-meson ratios  $p/\pi$  and  $\Lambda/K_S^0$ . At low  $p_T$ , predictions that include additional colour-reconnection mechanisms beyond the leading-colour approximation; assume the existence of additional higher-mass charm-baryon states; or include hadronisation via coalescence can describe the data, while predictions driven by charm-quark fragmentation processes measured in  $e^+e^-$  and  $e^-p$  collisions significantly underestimate the data. The results presented in this letter provide significant evidence that the established assumption of universality (colliding-system independence) of parton-to-hadron fragmentation is not sufficient to describe charm-baryon production in hadronic collisions at LHC energies.

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\*See Appendix A for the list of collaboration members

Heavy-flavour hadron production in hadronic collisions occurs through the fragmentation of a charm or beauty quark, created in hard parton-parton scattering processes, into a given meson or baryon. Theoretical calculations of heavy-flavour production generally utilise the QCD factorisation theorem [1], which describes the hadron cross section as the convolution of three terms: the parton distribution functions, the parton hard-scattering cross sections, and the fragmentation functions. It is generally assumed that the fragmentation functions are universal between collision systems and energies, and the measurement of the relative production of different heavy-flavour hadron species is sensitive to fragmentation functions used in pQCD-based calculations. While perturbative calculations at next-to-leading order with next-to-leading-log resummation [2–5] generally describe the D- and B-meson cross-section measurements [6–10] and the ratios of strange and non-strange D mesons [6, 10] within uncertainties, heavy-flavour baryon production is less well understood.

The  $\Lambda_c^+$  production cross section in pp collisions at  $\sqrt{s} = 7$  TeV and p–Pb collisions at  $\sqrt{s_{\text{NN}}} = 5.02$  TeV was reported by ALICE [11]. It was shown that in both collision systems the  $p_T$ -differential  $\Lambda_c^+$  production cross section is higher than predictions from pQCD calculations with charm fragmentation tuned on previous  $e^+e^-$  and  $e^-p$  measurements [2, 3]. The  $\Lambda_c^+/D^0$  ratio in pp and p–Pb collisions is consistent in both collision systems and also significantly underestimated by several Monte Carlo (MC) generators implementing different charm-quark fragmentation processes [12–15], suggesting that the fragmentation fractions of charm quarks into different hadronic states are non-universal with respect to collision system and centre-of-mass energy. The production of charm baryons has recently been calculated within the  $k_T$ -factorisation approach using unintegrated gluon distribution functions and the Peterson fragmentation functions [16], and with the GM-VFNS scheme using updated fragmentation functions from OPAL and Belle [17]. These approaches are unable to simultaneously describe ALICE and LHCb data with the same set of parameters, suggesting that the independent parton fragmentation scheme is insufficient to fully describe the results. An alternative explanation has been offered by a statistical hadronisation model, taking into account an augmented list of charm-baryon states based on guidance from the Relativistic Quark Model (RQM) [18] and lattice QCD [19], which is able to reproduce the  $\Lambda_c^+/D^0$  ratio measured by ALICE. The magnitude of the relative yields of  $\Lambda_b^0$  baryons and beauty mesons in pp collisions measured by LHCb [20–22] and CMS [23] offers further evidence that the fragmentation fractions in the beauty sector also vary between collision systems.

The measurement of baryon production has also been important in heavy-ion collisions, where the high energy density and temperature create a colour-deconfined state of matter [24]. A measured enhancement of the light-flavour [25, 26] and charm [27–29] baryon-to-meson ratio at the LHC and RHIC can be explained via an additional mechanism of hadronisation known as coalescence (or recombination), where soft quarks from the medium recombine to form a meson or baryon [30], in addition to hydrodynamical radial flow. Measurements in p–Pb collisions are crucial to provide an ‘intermediate’ collision system where the generated particle multiplicities and energy densities are between those generated in pp and A–A collisions. ALICE and CMS reported an enhancement of the baryon-to-meson ratios in the light-flavour sector ( $p/\pi$  and  $\Lambda/K_S^0$ ) at intermediate  $p_T$  ( $2 < p_T < 10$  GeV/ $c$ ) in high-multiplicity pp and p–Pb collisions, similar to that observed in heavy-ion collisions [31, 32]. This adds to the evidence that small systems also exhibit collective behaviour, which may have similar physical origins in pp, p–A, and A–A collisions [33]. It has been suggested that hadronisation of charm quarks via coalescence may also occur in pp and p–Pb collisions [34–36].

In this letter, the measurements of the prompt production of the charm baryon  $\Lambda_c^+$  in pp collisions at  $\sqrt{s} = 5.02$  TeV in  $|y| < 0.5$  and in p–Pb collisions at  $\sqrt{s_{\text{NN}}} = 5.02$  TeV in  $-0.96 < y < 0.04$

are presented, with a focus on the  $\Lambda_c^+/D^0$  production ratios. The measurement is performed as an average of the  $\Lambda_c^+$  and its charge conjugate  $\bar{\Lambda}_c^-$ , collectively referred to as  $\Lambda_c^+$  in the following. Two hadronic decay channels were measured:  $\Lambda_c^+ \rightarrow pK^-\pi^+$  (branching ratio BR =  $6.28 \pm 0.33\%$ ), and  $\Lambda_c^+ \rightarrow pK_S^0$  (BR =  $1.59 \pm 0.08\%$ ) [37], which were reconstructed exploiting the topology of the weakly-decaying  $\Lambda_c^+$  ( $c\tau = 60.7\mu\text{m}$ ) [37]. The results from both decay channels were averaged to obtain more precise production cross sections. With respect to the results presented in [11], this work studies a different centre-of-mass energy for pp collisions, and the cross section is measured in finer  $p_T$  intervals and over a wider  $p_T$  range. The overall precision of the measurements is significantly improved, by a factor of 1.5–2, depending on  $p_T$ , for both pp and p–Pb collisions. For a detailed description of the analysis techniques, corrections, systematic uncertainty determination, and supplementary measurements, the reader is referred to [38].

A description of the ALICE detector and its performance are reported in [39, 40]. The pp data sample was collected in 2017 and the p–Pb data sample was collected in 2016 during the LHC Run 2. Both pp and p–Pb collisions were recorded using a Minimum Bias (MB) trigger, which required coincident signals in the two V0 scintillator detectors located on either side of the interaction vertex. Further offline selection was applied in order to remove background from beam–gas collisions and other machine-induced backgrounds. To reduce superposition of more than one interaction within the colliding bunches (pile-up), events with multiple reconstructed primary vertices were rejected. Only events with a  $z$ -coordinate of the reconstructed vertex position within 10 cm from the nominal interaction point were used. With these requirements, approximately one billion MB-triggered pp events were selected, corresponding to an integrated luminosity of  $\mathcal{L}_{\text{int}} = 19.5 \text{ nb}^{-1} (\pm 2.1\% [41])$ . Approximately 600 million MB-triggered p–Pb events were selected, corresponding to  $\mathcal{L}_{\text{int}} = 287 \mu\text{b}^{-1} (\pm 3.7\% [42])$ .

The analysis techniques used for the results presented here are described in detail in [38]. Charged-particle tracks and particle decay vertices are reconstructed in the central barrel using the Inner Tracking System (ITS) and the Time Projection Chamber (TPC), which are located inside a solenoid magnet of field strength 0.5 T. In order to reduce the large combinatorial background, selections on the  $\Lambda_c^+$  candidates were made based on the particle identification (PID) signals and the displacement of the decay tracks from the collision point. The PID was performed using information on the specific energy loss of charged particles as they pass through the gas of the TPC and, where available, with flight-time measurements given by the Time-Of-Flight detector (TOF).

For the  $\Lambda_c^+ \rightarrow pK^-\pi^+$  analysis, candidates were built by reconstructing triplets of tracks with the correct configuration of charges. For this analysis, the high-resolution tracking provided by the detectors meant that the decay vertex of the  $\Lambda_c^+$  candidates could be resolved from the interaction point. To identify each of the p, K, and  $\pi$  daughter tracks, information from the TPC and TOF was combined using the ‘maximum-probability’ Bayesian approach described in [43]. Kinematic selections were made on the  $p_T$  of the decay products of the  $\Lambda_c^+$ , and geometrical selections were made on topological properties related to the displaced vertex of the  $\Lambda_c^+$  decay.

The reconstruction of  $\Lambda_c^+ \rightarrow pK_S^0$  candidates relied on reconstructing the V-shaped decay of the  $K_S^0$  meson into two pions, which was then combined with a proton track (bachelor). In pp collisions, candidates were further selected using criteria related to PID and properties of the  $\Lambda_c^+ \rightarrow pK_S^0$  decay. The Bayesian probability of the combined TPC and TOF response for the bachelor track to be a proton was required to be above 80%. The selection criteria on kinematical and geometrical variables included the distance of closest approach between the decay daughters, the invariant mass, and the cosine of the pointing angle of the neutral decay vertex ( $K_S^0$ ) to the primary vertex.

For the  $\Lambda_c^+ \rightarrow pK_S^0$  decay channel in p-Pb collisions, the analysis was performed using a multivariate technique based on the Boosted Decision Tree (BDT) algorithm provided by the Toolkit for Multivariate Data Analysis (TMVA) [44]. The BDT algorithm was trained using signal and background  $\Lambda_c^+ \rightarrow pK_S^0$  decay candidates simulated using PYTHIA 6.4.25 [45] with the Perugia2011 tune [46], and the underlying p-Pb event simulated with HIJING 1.36 [47]. Candidates obtained with the same reconstruction strategy previously described were preselected using loose geometrical selections and PID selection on the bachelor proton track. The model was trained independently for each  $p_T$  interval analysed, with input variables comprising the  $p_T$  and Bayesian PID probability of the proton track, the  $c\tau$  and invariant mass of the  $K_S^0$ , and the impact parameters of the  $\Lambda_c^+$  decay tracks to the primary vertex. This model was then applied on data, and a selection on the output response was chosen based on the expected maximum significance determined from simulations.

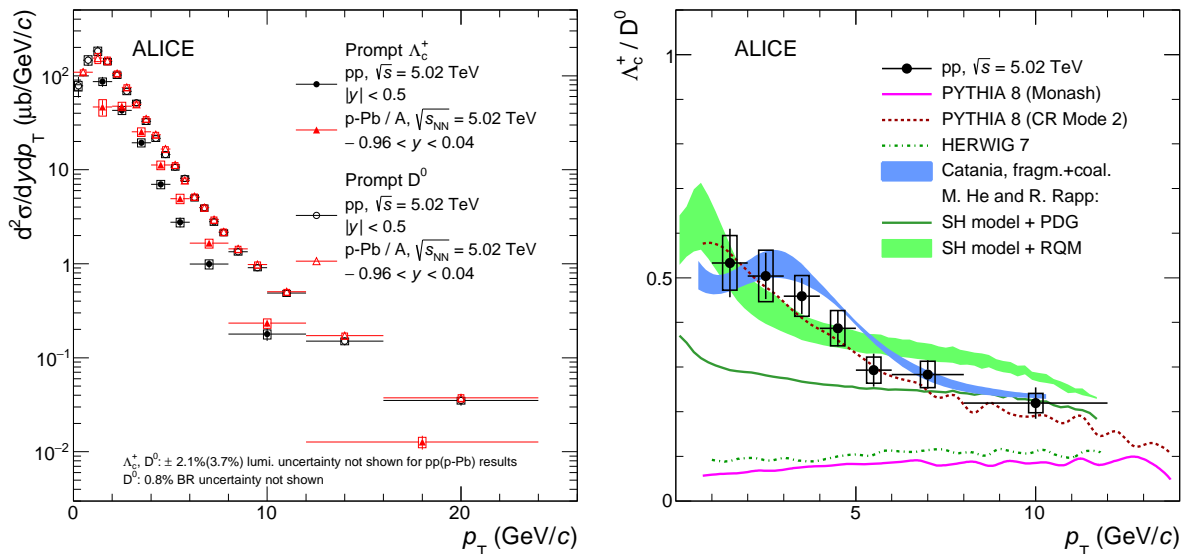
For both decay channels the yield of  $\Lambda_c^+$  baryons was extracted in each  $p_T$  interval via fits to the candidate invariant-mass distributions. The fitting function consisted of a Gaussian to estimate the signal and an exponential or polynomial function to estimate the background. The width of the Gaussian was fixed in each  $p_T$  interval to values obtained from Monte Carlo simulations, and the mean was treated as a free parameter. A statistical significance higher than 4 standard deviations was achieved in all  $p_T$  intervals.

Several corrections were applied to the measurement of the  $\Lambda_c^+$  cross section. The geometrical acceptance of the detector as well as the selection and reconstruction efficiencies for prompt  $\Lambda_c^+$  were taken into account. These correction factors were determined from pp collisions generated with PYTHIA 6 and PYTHIA 8.243 [48], with each event including either a  $c\bar{c}$  or a  $b\bar{b}$  pair. For p-Pb collisions, this was supplemented with an underlying event from the HIJING event generator. In p-Pb collisions the efficiency was calculated after reweighting the events based on their charged particle multiplicity. This accounts for the fact that the event multiplicity in simulation does not reproduce the one in data, and the efficiency depends on the multiplicity of the event as a consequence of the improvement of the resolution of the primary vertex and thus of the performance of the topological selections at higher multiplicities. The fraction of the  $\Lambda_c^+$  yield originating from beauty decays (feed-down) was obtained using the beauty-quark production cross section from FONLL [4, 5], the fraction of beauty quarks that fragment into beauty hadrons  $H_b$  from LHCb measurements [22], and  $H_b \rightarrow \Lambda_c^+ + X$  decay kinematics from PYTHIA 8, as well as the selection and reconstruction efficiency of  $\Lambda_c^+$  from beauty-hadron decays. The fraction of the  $\Lambda_c^+$  yield from beauty decays was found to be 2% at low  $p_T$  and up to 16% at high  $p_T$ , and was subtracted from the measured yield. As done in the D-meson analysis [49], the possible modification of beauty-hadron production in p-Pb collisions was included in the feed-down calculation by scaling the beauty-quark production by a nuclear modification factor  $R_{pPb}^{\text{feed-down}}$ , where it was assumed that  $R_{pPb}^{\text{feed-down}} = R_{pPb}^{\text{prompt}}$  with their ratio varied in the range  $0.9 < R_{pPb}^{\text{feed-down}}/R_{pPb}^{\text{prompt}} < 1.3$  to evaluate the systematic uncertainties.

Systematic uncertainties on the  $\Lambda_c^+$  cross sections were estimated considering the same sources as described in [11]. The contributions from the raw-yield extraction were evaluated by repeating the fits varying the fit interval and the functional form of the background fit function. For each of these variations the four combinations of free and fixed Gaussian mean and width parameters of the fit were considered. Overall, the relative uncertainty ranged from 4% to 11% depending on the  $p_T$  and analysis. The uncertainties on the track reconstruction efficiency were estimated by adding in quadrature the uncertainty due to track quality selection and the uncertainty due to the TPC-ITS matching efficiency (from 3% to 7%). The former is estimated by varying the track-quality selection criteria and the latter is estimated by comparing the probability to match the tracks from the TPC to the ITS hits in data and simulation. The uncertainty

on the  $\Lambda_c^+$  selection efficiency was estimated by varying the selection on the kinematical and topological properties of the  $\Lambda_c^+$  decays, or the selection on the BDT response (from 3% to 15%). The uncertainty on the PID efficiency was estimated by varying the selection on the Bayesian probability variables (from 2% to 5%). The systematic effect on the efficiencies due to the shape of the simulated  $\Lambda_c^+$   $p_T$  distribution was evaluated by reweighting the generated  $\Lambda_c^+$  from PYTHIA 6 to match the  $p_T$  distribution obtained from FONLL calculations for D mesons (maximum 1% uncertainty). The relative statistical uncertainty on the acceptance and efficiency correction was considered as an additional systematic uncertainty source (from 1–2% at low  $p_T$  to 3–5% at high  $p_T$ ). The uncertainties on  $f_{\text{prompt}}$  were estimated by varying the hypothesis on the production of  $\Lambda_c^+$  from B-hadron decays to account for the theoretical uncertainties of b-quark production within FONLL and experimental uncertainties on B-hadron fragmentation (around 2% at low  $p_T$ , and from 4% to 7% at high  $p_T$ , depending on the analysis). Global uncertainties of the measurement include those from the luminosity and  $\Lambda_c^+$  branching ratios. The raw-yield extraction uncertainty source are considered to be uncorrelated across  $p_T$  bins, while all other sources are considered to be correlated.

The results in each collision system from the two  $\Lambda_c^+$  decay channels were averaged to obtain the final results. A weighted average of the results was calculated, with weights defined as the inverse of the quadratic sum of the relative statistical and uncorrelated systematic uncertainties. The sources of systematic uncertainty assumed to be uncorrelated between different decay channels were those due to the raw-yield extraction, the statistical uncertainties on the efficiency and acceptance, and those related to the  $\Lambda_c^+$  selection. The remaining uncertainties were assumed to be correlated, except the branching ratio uncertainties, which were treated as partially correlated among the hadronic-decay modes as defined in [37].



**Fig. 1:** Left: Prompt  $\Lambda_c^+$  and  $D^0$   $p_T$ -differential cross section in pp collisions and in p-Pb collisions at  $\sqrt{s_{\text{NN}}} = 5.02$  TeV. The results in p-Pb collisions are scaled with the atomic mass number  $A$  of the Pb nucleus. Right: the  $\Lambda_c^+ / D^0$  ratio as a function of  $p_T$  measured in pp collisions at  $\sqrt{s} = 5.02$  TeV compared with theoretical predictions (see text for details). Statistical uncertainties are shown as vertical bars, while systematic uncertainties are shown as boxes, and the bin widths are shown as horizontal bars.

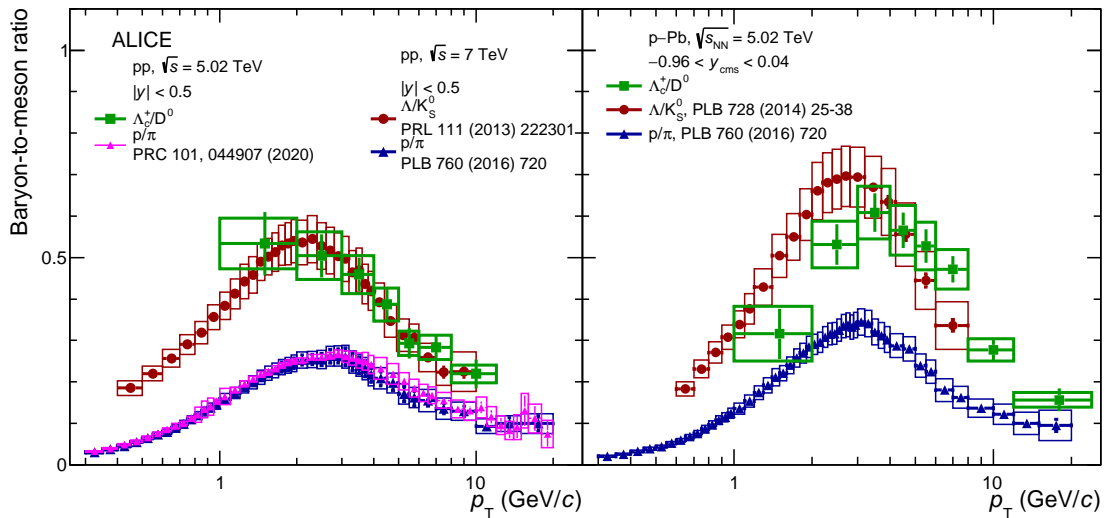
Figure 1 (left) shows a comparison of the  $\Lambda_c^+$   $p_T$ -differential cross sections in pp and in p-Pb collisions at  $\sqrt{s_{\text{NN}}} = 5.02$  TeV. The  $D^0$   $p_T$ -differential cross sections measured in the same collision systems and at the same centre-of-mass energy during the same data taking periods [10, 50] are also shown. In order to compare the spectral shapes in the two different collision systems

at the same energy, the results in p-Pb collisions are scaled by the atomic mass number of the lead nucleus. For  $\Lambda_c^+$  baryons the spectral shape in p-Pb collisions is slightly harder than in pp collisions, while for  $D^0$  mesons the spectral shapes are fully consistent within uncertainties.

Figure 1 (right) shows the baryon-to-meson ratio  $\Lambda_c^+/D^0$  measured in pp collisions at  $\sqrt{s} = 5.02$  TeV as a function of  $p_T$ , compared to theoretical predictions. The uncertainty on the luminosity cancels in the ratio. The  $\Lambda_c^+/D^0$  ratio is measured to be 0.4–0.5 at low  $p_T$ , and decreases to around 0.2 at high  $p_T$ . The previous results at  $\sqrt{s} = 7$  TeV hinted at a decrease of the  $\Lambda_c^+/D^0$  ratio with  $p_T$ , although the precision was not enough to confirm this [11]. The results in pp collisions at  $\sqrt{s} = 5.02$  TeV, with much higher precision than  $\sqrt{s} = 7$  TeV results, show a clear decrease with increasing  $p_T$ . The strong  $p_T$ -dependence of the  $\Lambda_c^+/D^0$  ratio is in contrast with the ratios of strange and non-strange D mesons in pp collisions at  $\sqrt{s} = 5.02$  TeV and  $\sqrt{s} = 7$  TeV [10, 51] and in p-Pb collisions at  $\sqrt{s_{\text{NN}}} = 5.02$  TeV [50], which do not show a significant  $p_T$  dependence within uncertainties and thus indicate that there are no large differences between fragmentation functions of charm quarks to charm mesons. The result presented here instead provides strong indications that the fragmentation functions of baryons and mesons differ significantly.

The measured  $\Lambda_c^+/D^0$  ratios in pp collisions are compared with predictions from several MC generators and models in which different hadronisation processes are implemented. The PYTHIA 8 predictions include the Monash tune [12] and a tune that implements colour reconnection beyond the leading-colour approximation, corresponding to CR Mode 2 as defined in [13]. Hadronisation in PYTHIA is built on the Lund string fragmentation model [52, 53], where quarks and gluons connected by colour strings fragment into hadrons, and colour reconnection allows for partons created in the collision to interact via colour strings. The latter tune introduces new colour reconnection topologies beyond the leading-colour approximation, including ‘junctions’ that fragment into baryons, leading to increased baryon production. As a technical point, the PYTHIA 8 simulations are generated with all soft QCD processes switched on [48]. The PYTHIA 8 Monash tune and HERWIG 7.2 [15] predictions are driven by the fragmentation fraction  $f(c \rightarrow \Lambda_c^+)$  implemented in these generators, which all suggest a relatively constant  $\Lambda_c^+/D^0$  ratio versus  $p_T$  of about 0.1, significantly underestimating the data at low  $p_T$ . At high  $p_T$ , the data approach the predictions from these generators, although the measurement in  $8 < p_T < 12$  GeV/c is still underestimated by about a factor of 2. A significant enhancement of the  $\Lambda_c^+/D^0$  ratio is seen with colour reconnection beyond the leading-colour approximation (PYTHIA 8 CR Mode 2). This prediction is consistent with the measured  $\Lambda_c^+/D^0$  ratio in pp collisions, also reproducing the downward  $p_T$  trend. The statistical hadronisation model (‘SH model’ in the legend) [19] uses either an underlying charm-baryon spectrum taken from the PDG, or includes additional excited charm baryons that have not yet been observed but are predicted by the RQM. These additional states decay strongly to  $\Lambda_c^+$  baryons, which contribute to the prompt  $\Lambda_c^+$  spectrum. The RQM predictions include a source of uncertainty related to the branching ratios of the excited baryon states into  $\Lambda_c^+$  final states, which is estimated by varying the branching ratios between 50% and 100%. With the PDG charm-baryon spectrum the model underpredicts the data. With the additional baryon states the model instead gives a good description of the pp data, both in the magnitude of the ratio, and the decreasing trend with  $p_T$ . The Catania model [36] assumes that a colour-deconfined state of matter is formed and hadronisation can occur via coalescence in addition to fragmentation. Coalescence is implemented through the Wigner formalism, where a blast wave model is used to determine the  $p_T$  spectrum of light quarks and FONLL pQCD calculations are used for heavy quarks. Hadronisation via coalescence is predicted to dominate at low  $p_T$ , while fragmentation dominates at high  $p_T$ . This model provides a good description of both the magnitude and shape of the data over the full  $p_T$  range.

Figure 2 shows the  $\Lambda_c^+/D^0$  baryon-to-meson ratio measured in pp collisions at  $\sqrt{s} = 5.02$  TeV



**Fig. 2:** The charm baryon-to-meson ratio  $\Lambda_c^+/D^0$  in pp collisions (left), and p-Pb collisions (right) at  $\sqrt{s_{\text{NN}}} = 5.02$  TeV, compared to the light-flavour baryon-to-meson ratios  $\Lambda/K_S^0$  and  $p/\pi$ . Statistical uncertainties are shown as vertical bars, while systematic uncertainties are shown as boxes, and the bin widths are shown as horizontal bars.

(left) and in p-Pb collisions at  $\sqrt{s_{\text{NN}}} = 5.02$  TeV (right) as a function of  $p_T$ , compared to baryon-to-meson ratios in the light-flavour sector,  $\Lambda/K_S^0$  [25, 54] and  $p/\pi$  [31, 55] (calculated as the sum of both charged particles and antiparticles,  $(p + \bar{p})/(\pi^+ + \pi^-)$ ). The  $p/\pi$  ratio in pp collisions is shown at both  $\sqrt{s} = 5.02$  TeV and  $\sqrt{s} = 7$  TeV, displaying consistent results at both centre-of-mass energies, while the  $\Lambda/K_S^0$  ratio in pp collisions is shown only at  $\sqrt{s} = 7$  TeV. Unlike heavy-flavour hadron production, which occurs primarily through the fragmentation of a charm quark produced in the initial hard scattering, light-flavour hadrons have a significant contribution from gluon fragmentation. Low- $p_T$  light-flavour hadrons also primarily originate from soft scattering processes involving small momentum transfers. All particle yields in these ratios were corrected for feed-down from weak decays, although the pion spectrum is expected to have significant feed-down contributions also from the strong decays of other particle species, primarily  $\rho$  and  $\omega$  mesons. Despite these differences, the three ratios,  $\Lambda_c^+/D^0$ ,  $\Lambda/K_S^0$ , and  $p/\pi$  demonstrate some remarkably similar characteristics in both collision systems. All ratios exhibit a decreasing trend after  $p_T \gtrsim 2-3$  GeV/c. The  $\Lambda_c^+/D^0$  and  $\Lambda/K_S^0$  ratios are consistent, in terms of both shape and magnitude, within uncertainties. The light-flavour ratios both peak at  $\sim 2-3$  GeV/c in both pp and p-Pb collisions, and there is an indication of a peak at  $2 < p_T < 4$  GeV/c in the  $\Lambda_c^+/D^0$  ratio in p-Pb collisions. These similarities between heavy-flavour and light-flavour measurements hint at a potential common mechanism for light- and charm-baryon formation in pp and p-Pb collisions at LHC energies. It is interesting to note that all baryon-to-meson ratios also indicate a shift toward higher momenta in p-Pb collisions, which for light-flavour particle production is often attributed to radial flow [54]. However, while flow effects in the charm sector ( $D^0$  and heavy-flavour decay leptons) have been observed in high-multiplicity p-Pb collisions [56, 57], these effects are expected to be smaller at lower multiplicities, and also smaller for charm than for light-flavour hadrons.

In summary,  $\Lambda_c^+$ -baryon production was measured in pp collisions at midrapidity ( $|y| < 0.5$ ) and in p-Pb collisions in the rapidity interval  $-0.96 < y < 0.04$  at  $\sqrt{s_{\text{NN}}} = 5.02$  TeV. A clear  $p_T$ -dependence of the  $\Lambda_c^+/D^0$  ratio is reported, with the ratio decreasing as the  $p_T$  increases. This trend is similar to that of baryon-to-meson ratios measured in the light-flavour sector in pp and p-Pb collisions, suggesting common mechanisms for light- and charm-baryon formation.

While models incorporating fragmentation parameters from  $e^+e^-$  and  $e^-p$  collisions significantly underestimate the  $\Lambda_c^+/D^0$  ratio, three models can reproduce the measurements. The first is a tune of PYTHIA 8 which considers that, in pp collisions at high energy, multi-parton interactions produce a rich hadronic environment that requires an extension of colour reconnection in hadronisation processes beyond the leading-colour approximation. The second method is the SH+RQM model, which relies on the presence of a large set of yet-unobserved higher-mass charm-baryon states with relative yields following the Statistical Hadronisation model. The third relies on hadronisation via coalescence and fragmentation after the formation of a colour-deconfined state of matter. All three models imply a substantially different description of the charm-baryon production in pp collisions with respect to  $e^+e^-$  and  $e^-p$  collisions, indicating that the assumption of universal parton-to-hadron fragmentation between collision systems is not sufficient to describe charm-baryon production.

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