



## Saturation properties of four-wave mixing between short optical pulses in semiconductor optical amplifiers

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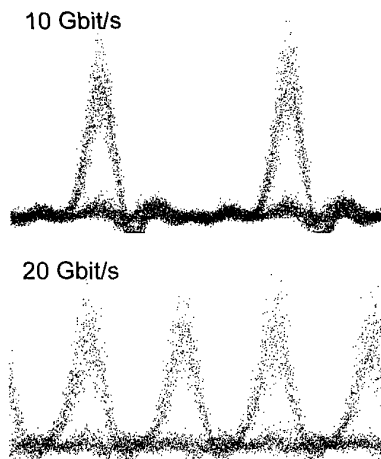
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CTuW3 Fig. 3. Experimental measured eye diagrams for a bi-directional transmitted signal at 10 and 20 Gbit/s, respectively.

directional MSSl swapping at 10 and 20 Gbit/s are shown in Fig. 3. The coupled CW pump power in the experiment was 6 to 7 dBm. The coupled signal power was around -6 to -7 dBm. The difference in CW pump power for equal FWM output power for up- and down-conversion was 2-3 dB. These values are in good agreement with the model predictions.

In conclusion, large signal simulations of a novel scheme of bi-directional FWM are shown to be in good agreement with experiments and identify critical aspects of the system performance.

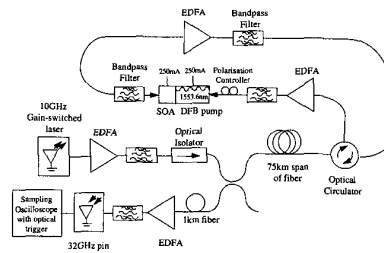
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#### CTuW4 5:15 pm

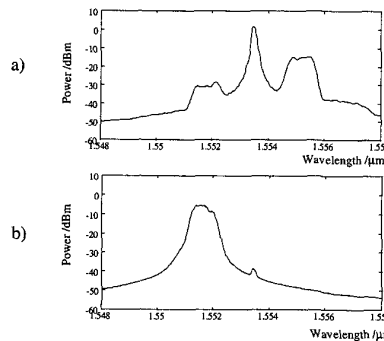
##### Demonstration of four wave mixing in an integrated pump laser and semiconductor optical amplifier for mid-span spectral inversion dispersion compensation

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Non-degenerate four wave mixing (NDFWM) has been studied extensively in recent years as both a method for compensating against the dispersion of standard single mode optical fibre via mid-span spectral inversion (MSSI)<sup>1</sup> and as a transparent wavelength conversion mechanism. NDFWM has been reported in many devices, including semiconductor optical amplifiers (SOAs), distributed feedback (DFB) lasers, nonlinear waveguides and optical fibre itself.<sup>2</sup> SOAs are promising devices for NDFWM due to their small size, large gain bandwidth and high speed (>100 Gb/s).<sup>3</sup> However, an external pump laser and EDFA



CTuW4 Fig. 1. Schematic of the experimental MSSl arrangement.



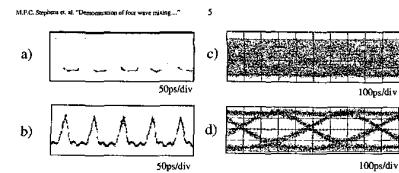
CTuW4 Fig. 2. (a) Spectrum taken directly after the integrated DFB pump/SOA, clearly showing the newly generated phase conjugate at 1552 nm. (b) The phase conjugate signal after filtering and amplification.

are required which adds both cost and complexity to the system. Recently, an integrated DFB laser and SOA has been reported for wavelength conversion under static conditions.<sup>3</sup> Here, for the first time we believe, the use of a similar device for dynamic MSSl dispersion compensation is demonstrated.

The experimental arrangement is depicted in Fig. 1.

A 10 GHz train of 18 ps pulses is generated by gain-switching a 1555.2 nm DFB laser. These pulses are amplified by an EDFA and injected into a 75 km span of standard fibre. The output of the span is passed via an optical circulator into another EDFA and filter before injection into the DFB laser section of the phase conjugator. A polarisation controller is used to match the polarisation of the incoming signal to that of the DFB pump beam. The DFB pump wavelength is at 1553.6 nm and has a linewidth of ~2 MHz when both it and the SOA are biased at 250 mA. The newly generated phase conjugate signal at 1552 nm is filtered to remove the pump and original signal before being amplified, filtered again and transmitted back over the span. The received signal is analysed after traversing an additional 1 km of standard fibre, which compensates for the dispersion induced by the wavelength difference on the two spans. The pulses are examined using a 32 GHz photodiode and 50 GHz sampling oscilloscope.

The optical spectra produced from the output of the phase conjugator are shown in Fig. 2 before and after filtering. The average injected signal power to the phase conjugator is +3.5



CTuW4 Fig. 3. (a) 10 GHz 18 ps input pulses at 1555.2 nm. (b) Received phase conjugate pulses at 1552 nm after 151 km of standard fibre. (c) 2.5 Gb/s eye with no MSSl at 1555.6 nm. (d) Phase conjugate 2.5 Gb/s eye after 151 km.

dBm. After filtering it can be seen that the residual peak pump level is over 30 dB lower than the phase conjugate.

The back to back and final received pulses are shown in Figs. 3(a), 3(b). It can be seen that there is some degradation of the signal after the 150 km span. In addition, a highly chirped laser is modulated at 2.5 Gb/s and transmitted over the system. The received eye is shown in [Fig. 3(d)] with no significant degradation. Without MSSl, the same eye diagram is completely closed [Fig. 3(c)].

In conclusion, for the first time, an integrated DFB/SOA laser has been shown to operate as an MSSl dispersion compensator via NDFWM with a 10 GHz pulse train and 2.5 Gb/s data. Although the former does show some degradation, it is believed that further optimisation of the device (such as reduced DFB linewidth) will lead to far superior performance. A full bit error rate analysis and results at 10 Gb/s will be presented at the conference.

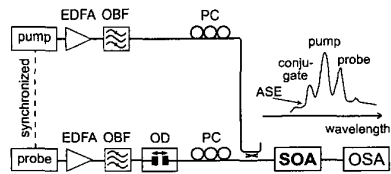
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#### CTuW5 5:30 pm

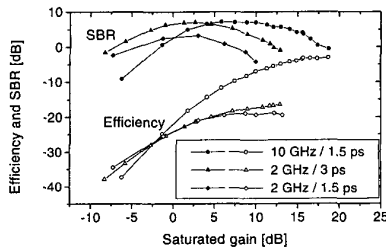
##### Saturation properties of four-wave mixing between short optical pulses in semiconductor optical amplifiers

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Four-wave mixing (FWM) between short optical pulses in a semiconductor optical amplifier (SOA) is of considerable interest for performing high-speed optical de-multiplexing or optical sampling.<sup>1-3</sup> The large gain of the SOA makes efficient FWM over detunings of several nanometers possible. Since FWM relies on physical processes (like carrier density modulation and ultrafast intraband carrier dynamics



**CTuW5 Fig. 1.** Experimental setup. Pump (1540 nm) and probe (1550 nm) pulses are amplified (EDFA) and band-pass filtered (OBF) before injection into the SOA. An optical delay line (OD) and polarization controllers (PC) are used to align the pulses in time and polarization. The schematic FWM spectrum shows the different wavelength components at the SOA output.

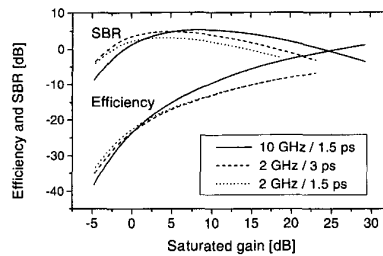


**CTuW5 Fig. 2.** Experimental results. The average input powers were 3 dBm for the pump and -7 dBm for the probe. The saturated gain is the fiber-to-fiber gain, including an estimated total coupling loss of -7 dB.

due to carrier heating and spectral holeburning) that all contribute to gain saturation and therefore counteract the FWM process, the characteristics depend strongly on the operation conditions, i.e., SOA gain, input power, bitrate and pulsewidth.

Figure 1 shows the experimental set-up. Pump and probe pulses are generated using two synchronized external cavity semiconductor modelocked lasers, which can be tuned in pulsewidth and repetition rate.<sup>4</sup> The pulses are injected, co-polarized and temporally overlapping, into a 980  $\mu\text{m}$  long bulk InGaAsP amplifier. The efficiency (conjugate power at output divided by probe input) and the signal-to-background ratio (SBR) are measured using an optical spectrum analyzer (1 nm resolution bandwidth) as a function of the saturated SOA gain. The SOA driving current is varied to adjust to different gain values, which are recorded using a weak CW beam at the probe wavelength. Measurements for a constant average pump power of 3 dBm (probe is 10 dB lower) and pulsewidths of 1.5 ps and 3 ps as well as repetition rates of 2 GHz and 10 GHz are shown in Fig. 2. Figure 3 depicts calculated results using the known parameter values from the experiment. The theoretical model incorporates a detailed treatment of propagation effects as well as ultrafast carrier dynamics, but the application of a simple mathematical trick transforms it into a tractable and numerically very effective approach that avoids the solution of partial differential equations.<sup>5</sup>

The calculated results capture all the qualitative effects seen in the experiment. We attribute the quantitative differences between experiment and theory to uncertainties in the parameter values as well as dynamical effects



**CTuW5 Fig. 3.** Calculated results.

due to saturation of the gain coefficients themselves, which is not included in the model.

The existence of a certain gain, or injection current, that optimizes the SBR is special to the interaction between trains of short optical pulses. For CW beams, in contrast, both the efficiency and the SBR increase monotonically with the gain.<sup>4</sup> The reason for this behavior is that the conjugate pulse experiences the maximum gain saturation since it overlaps with the pump pulse, while the amplified spontaneous emission (ASE) noise has time to recover in-between the pulses, and therefore experiences a gain which is closer to the unsaturated level. For increasing gain the ASE therefore grows faster than the conjugate energy and a maximum is obtained. Both the calculated and measured data show that the maximum moves to larger gain for the higher repetition rate, but changes only in amplitude for the investigated pulsewidths. We notice that the inclusion of saturation effects due to ultrafast carrier dynamics is absolutely necessary in order to obtain the qualitative dependencies seen in the experiment.

In conclusion, we report the first comparison between theory and experiment on the FWM between trains of short pulses in SOAs. The theory is able to explain all qualitative features seen in the experiment.

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#### CTuW6 5:45 pm

##### Cross-gain saturation reduction in a double-stage overpumped EDFA with pump re-use

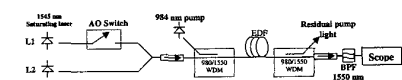
D.A. Carter, Jeff Korn, Lincoln Laboratory, Massachusetts Institute of Technology, 244 Wood Street, Lexington, Massachusetts 02420-9108 USA; E-mail: Dedric@mit.edu, korn@LL.mit.edu

Cross-gain saturation in EDFAs gives rise to deleterious effects for WDM, packet TDM net-

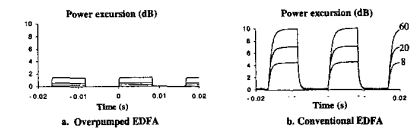
works, and for secure all-optical networks.<sup>1</sup> In this paper we make use of the fact that the signal saturation power depends upon the pump saturation level.<sup>2</sup> Using a short piece of Erbium-doped fiber and utilizing only a fraction of pump light (~10%) power excursions are greatly reduced while providing good noise figure and gain. A similar scheme has been analyzed numerically by Willems *et al.*<sup>7</sup> Also demonstrated is pump re-use in a double-stage configuration with inter-stage loss. This architecture could be used to prevent jamming of a frequency-selective switch.<sup>1</sup>

The single-stage experiment is shown in Fig. 1. Light from a 1545 nm laser, L1, is sent through an acousto-optic switch to simulate the addition/dropping of channels. A 1550 nm probe laser, L2, serves as the surviving channel, and is combined with L1 via a 50/50 coupler, followed by the overpumped EDFA, consisting of input and output isolators, 13 m of Lucent's HP-980 erbium-doped fiber, and two WDMs. The pump laser is a fiber-coupled SDL MOPA, providing 450 mW of co-propagating 984 nm power. The output of the EDFA is then passed through a 1-nm bandwidth filter and the 1550 nm signal is displayed on an oscilloscope. The small-signal gain and noise-figure are measured to be 20 dB and 3.8 dB, respectively, with ~40 mW of absorbed pump power.

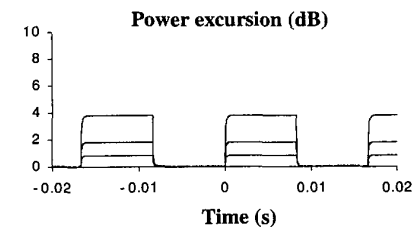
The power of L2 into the EDFA is -20 dBm and the peak power of L1 is either -11, -7, or -2.2 dBm. This simulates a 9, 21, or 61 (-20 dBm/channel) channel WDM system where 8, 20, or 60 channels are added/dropped, respectively, in a purely homogeneous gain medium. Figure 2(a) shows the power excursions were 0.27, 0.6, and 1.5 dB, respectively, for the 3 cases. By comparison, Fig. 2(b) shows the power excursions when a



**CTuW6 Fig. 1.** Experimental setup.



**CTuW6 Fig. 2.** Power excursion (dB) versus time (ms) for dropping 8, 20, or 60 channels (-20 dBm/channel).



**CTuW6 Fig. 3.** Power excursion (dB) versus time (ms) for dropping 8, 20, or 60 channels (-20 dBm/channel) for the double-stage EDFA with 15-dB mid-stage loss and pump re-use.