

SCALING LIMIT FOR A CLASS OF GRADIENT FIELDS WITH NON-CONVEX POTENTIALS

BY MAREK BISKUP AND HERBERT SPOHN

UCLA and University of South Bohemia, and TU Munich

We consider gradient fields $(\phi_x: x \in \mathbb{Z}^d)$ whose law takes the Gibbs-Boltzmann form $Z^{-1} \exp\{-\sum_{\langle x,y \rangle} V(\phi_y - \phi_x)\}$, where the sum runs over nearest neighbors. We assume that the potential V admits the representation

$$V(\eta) := -\log \int \varrho(d\kappa) \exp[-\frac{1}{2}\kappa\eta^2]$$

where ϱ is a positive measure with compact support in $(0, \infty)$. Hence the potential V is symmetric but non-convex in general. While for strictly convex V 's the translation-invariant, ergodic gradient Gibbs measures are completely characterized by their tilt, a non-convex potential as above may lead to several ergodic gradient Gibbs measures with zero tilt. Still, every ergodic, zero-tilt gradient Gibbs measure for the potential V from above scales to a Gaussian free field.

1. Introduction. Gradient fields belong to a class of models that arise in equilibrium statistical mechanics, e.g., as approximations of critical systems and as effective interface models. Although their definition is rather simple, and in fact quite a lot is known (see the reviews by Funaki [15], Velenik [28] or Sheffield [25]), much is still there to be learned. In this note we study gradient fields on a lattice. Here the field is a collection of real-valued random variables $\phi := \{\phi_x: x \in \mathbb{Z}^d\}$ and the distribution of ϕ on $\mathbb{R}^{\mathbb{Z}^d}$ is given by the formal expression

$$\frac{1}{Z} \exp\left\{-\sum_{\langle x,y \rangle} V(\phi_y - \phi_x)\right\} \prod_{x \in \mathbb{Z}^d} d\phi_x, \quad (1.1)$$

where $d\phi_x$ is the Lebesgue measure, $\langle x, y \rangle$ refers to an unordered nearest-neighbor pair on \mathbb{Z}^d , and V is an even, measurable function — called the *potential* — which is bounded from below and grows superlinearly at $\pm\infty$.

Of course, to define the measure (1.1) precisely we have to restrict the above expression to a finite subset of \mathbb{Z}^d and fix the ϕ 's on its boundary; Z is then the normalizing constant. Another way to regularize the expression (1.1) is to consider directly measures on all of $\mathbb{R}^{\mathbb{Z}^d}$ whose conditional

AMS 2000 subject classifications: Primary 60K35, 60K35; secondary 60K35

Keywords and phrases: gradient fields, scaling limit, Gaussian free field

probabilities in finite sets take the above form. In our context this standard definition is hampered by the fact that, due to the unbounded nature of the fields ϕ_x , no such infinite-volume measures may exist at all. However, if one restricts attention to (the σ -algebra generated by) the gradient variables,

$$\eta_{xy} := \phi_y - \phi_x, \quad |x - y| = 1, \quad (1.2)$$

then infinite volume measures exist under the above conditions on V . Since the measure depends only on gradients, we refer to such measures as *gradient Gibbs measures* (GGM), in accord with Funaki [15] and Sheffield [25].

Throughout we will focus on translation-invariant GGMs. An important characteristic is their *tilt*. For a translation-invariant GGM μ there exists a unique tilt vector $t \in \mathbb{R}^d$ such that

$$E_\mu(\eta_b) = t \cdot b \quad (1.3)$$

for every edge b of \mathbb{Z}^d — which we regard as a vector in this formula. Of course, this definition is really meaningful only for the GGMs that are ergodic — i.e., trivial on the σ -algebra of events invariant under all lattice translations. Indeed, in the ergodic case t represents the average incline of typical configurations.

For the case of quadratic V — the *massless free field* — the measure (1.1) is Gaussian and so many desired characteristics are amenable to explicit computations. The challenge for mathematicians has been to develop an equivalent level of understanding for non-quadratic V 's. A good amount of progress in this direction has been made in the last 10 to 15 years: Brydges and Yau [7] (and also earlier works, e.g., by Gawędzki and Kupiainen [12] and Magnen and Sénéor [19]) studied the effect of analytic perturbations of the quadratic potentials and concluded that the large scale behavior is that of the massless free field. Naddaf and Spencer [23] proved the same non-perturbatively for strictly convex potentials V and GGMs with zero tilt. The corresponding extension to non-zero tilt was obtained by Giacomin, Olla, and Spohn [13]. For the same class of potentials, Funaki and Spohn [16] proved a bijection between the ergodic GGMs and their tilt. Sheffield [25] characterized translation-invariant GGMs by means of a Gibbs variational principle and extended Funaki-Spohn's results to fields taking only a discrete set of values. We refer to the reviews by Funaki [15], Velenik [28] and Sheffield [25] for further results and references.

As a unifying feature, all the (non-perturbative) results mentioned are based on the strict convexity of the potential V — be it for the avail of the Brascamp-Lieb inequality [23, 13, 16], Helffer-Sjöstrand random walk

representation [23, 13], coupling to Langevin dynamics [16], and the cluster swapping-algorithm [25]. One would naturally like to have a non-perturbative approach that works even for non-convex potentials. With this motivation, Biskup and Kotecký [5] recently studied the GGMs for non-convex V that are a log-mixture of centered Gaussians,

$$V(\eta) := -\log \int \varrho(d\kappa) e^{-\frac{1}{2}\kappa\eta^2}, \quad (1.4)$$

where ϱ is a positive measure with compact support in $(0, \infty)$. Surprisingly, already for the simplest non-trivial case,

$$\varrho := p\delta_{\kappa_1} + (1-p)\delta_{\kappa_2} \quad (1.5)$$

with $\kappa_1 \gg \kappa_2 > 0$, it was shown that, in $d = 2$, there is a value $p \in (0, 1)$ at which one can construct *two* distinct, translation-invariant, gradient Gibbs measures of zero tilt.

The relevant conclusion from [5] for the general theory is that the one-to-one correspondence between ergodic GGMs and their tilt breaks down once V is sufficiently non-convex. The next natural question is to understand what happens to the scaling limit. The purpose of this note is to show that, regardless of the occurrence of phase transition, for potentials of the form (1.4), every translation-invariant, ergodic GGM with zero tilt scales to a Gaussian free field.

The proof is based on the fact — utilized already in [5] — that (1.4) allows us to represent every GGM as a mixture over Gaussian gradient measures with a random coupling constant κ_{xy} for each edge $\langle x, y \rangle$. Its covariance is simply the inverse of the operator

$$(\mathcal{L}_\kappa f)(x) := \sum_{y: |y-x|=1} \kappa_{xy} [f(y) - f(x)], \quad (1.6)$$

where we take, once for all, $\kappa_{xy} = \kappa_{yx}$. The *fluctuations* in the Gaussian measure can be analyzed by invoking a random walk representation; \mathcal{L}_κ is the generator of a random walk with symmetric random jump rates known, equivalently, under the name *random conductance model*. The name arises naturally from the electrostatic interpretation of this problem; cf. Doyle and Snell [9], in which one views \mathbb{Z}^d as a resistor network with *conductance* κ_{xy} — or resistivity $1/\kappa_{xy}$ — assigned to an edge $\langle x, y \rangle$. As it turns out (see Lemma 3.2), if the initial GGM is ergodic then so is the law of the conductances. This makes homogenization a possible tool.

Much work has been done in the past two decades on the problem of random walks with random conductances. For our purposes it suffices to

invoke two known results: Kipnis and Varadhan's [18] invariance principle (i.e., scaling of the random walk to Brownian motion) and Delmotte and Deuschel's [10] annealed derivative heat-kernel bounds. (Note that in the Helffer-Sjöstrand random walk representation, as used in [23, 13], one also has to study a random walk in a random environment. However, this random environment fluctuates in time, while in our case it is static.) This handles the fluctuations of the field; an important technical issue is thus the control of the mean. This is where the zero tilt restriction comes in (see Lemma 3.4, Corollary 5.8 and discussion in Sect. 6).

Note: While this paper has been in the process of revision, further developments occurred in the study of gradient models with non-convex potentials. Cotar, Deuschel and Müller [8] have shown that for non-convex perturbations of potentials V where the size of non-convex region is small compared to typical fluctuations of the field, the conclusions are as in the convex case. (Their precise condition is a bound on the L^1 -norm of the negative part of second derivative.) This is a high-temperature result; work in progress by Adams, Kotecký and Müller [1] addresses the low-temperature case when non-convexities are allowed only sufficiently far away from the absolute minimum of V . Our contribution remains valuable despite these advances as it applies to all potentials of the kind (1.4), including those for which phase coexistence occurs.

The plan of this paper is as follows: In Sect. 2 we define precisely the concept of the gradient Gibbs measure and state our main theorem. In Sect. 3 we introduce the extended gradient Gibbs measures and characterize their conditional marginals. This will naturally lead to the aforementioned connections with random walks in reversible random environments. To keep the main line of the argument intact, we first finish proving our main result in Sect. 4 and only then expound on the random-walk connections in Sect. 5. Sect. 6 is devoted to discussion of the limitations to zero tilt and some open questions concerning gradient Gibbs measures.

2. Main results.

2.1. *Gradient Gibbs measures.* As mentioned above, infinite volume measures on the field variables (ϕ_x) may not always exist — particularly, in sufficiently low dimensions. To make our statements uniform in dimension, we will focus attention on the gradient variables. However, not even that will be entirely straightforward because the gradient variables satisfy a host of “hard-core” constraints which, in a sense, encapsulate most of the

interaction. Since η is gradient, one has

$$\eta_{x_1, x_2} + \eta_{x_2, x_3} + \eta_{x_3, x_4} + \eta_{x_4, x_1} = 0, \quad (2.1)$$

whenever (x_1, \dots, x_4) are the vertices of a cycle in \mathbb{Z}^d of length four. We will often write η_b for the *positively-oriented* edge b in \mathbb{Z}^d . Throughout we will only work with positively oriented edges and will use $\mathbb{B}(\Lambda)$ to denote the set of such edges with both endpoints in the set $\Lambda \subset \mathbb{Z}^d$.

The constraints (2.1) are implemented at the level of the *a priori* measure which is defined as follows: Fix a configuration $\eta \in \mathbb{R}^{\mathbb{B}(\mathbb{Z}^d)}$ that obeys (2.1) and, for $\Lambda \subset \mathbb{Z}^d$ finite, let $\nu_\Lambda^{(\eta_{\mathbb{B}(\Lambda)^c})}$ be the Lebesgue measure on the linear subspace of configurations (η'_b) such that $\eta'_b = \eta_b$ for all $b \notin \mathbb{B}(\Lambda)$ and that η' obeys the constraints (2.1). Note that, if $\bar{\phi}$ is a configuration such that $\eta_{xy} = \bar{\phi}_y - \bar{\phi}_x$ for every nearest-neighbor pair $\langle x, y \rangle$, then $\nu_\Lambda^{(\eta_{\mathbb{B}(\Lambda)^c})}$ is to within a normalization constant the projection to gradient variables of the Lebesgue measure on $\{\phi_x : x \in \Lambda\}$ subject to the boundary condition $\bar{\phi}$.

Next we will give a precise definition of the notion of gradient Gibbs measure. For a finite $\Lambda \subset \mathbb{Z}^d$, consider the *specification* γ_Λ , which is a measure in the first coordinate and a function of the boundary condition in the second coordinate, that is defined by

$$\begin{aligned} \gamma_\Lambda(d\eta_{\mathbb{B}(\Lambda)} | \eta_{\mathbb{B}(\Lambda)^c}) \\ := \frac{1}{Z_\Lambda(\eta_{\mathbb{B}(\Lambda)^c})} \exp\left\{ - \sum_{\substack{\langle x, y \rangle \\ x \in \Lambda, y \in \Lambda \cup \partial\Lambda}} V(\eta_{xy}) \right\} \nu_\Lambda^{(\eta_{\mathbb{B}(\Lambda)^c})}(d\eta_{\mathbb{B}(\Lambda)}). \end{aligned} \quad (2.2)$$

Here $Z_\Lambda(\eta_{\mathbb{B}(\Lambda)^c})$ is the normalizing constant.

DEFINITION 2.1. *Let $\mathcal{E}_{\mathbb{B}(\Lambda)} := \sigma(\{\eta_b : b \in \mathbb{B}(\Lambda)\})$. We say that a measure μ on $\mathbb{R}^{\mathbb{B}(\mathbb{Z}^d)}$ is a gradient Gibbs measure if the regular conditional probability $\mu(- | \mathcal{E}_{\mathbb{B}(\Lambda)^c})$ in any finite $\Lambda \subset \mathbb{Z}^d$ satisfies*

$$\mu(- | \mathcal{E}_{\mathbb{B}(\Lambda)^c})(\eta) = \gamma_\Lambda(- | \eta_{\mathbb{B}(\Lambda)^c}) \quad (2.3)$$

for μ -a.e. η .

Most of this paper is restricted to translation invariant gradient Gibbs measures. To define the required notation, for each $x \in \mathbb{Z}^d$ let $\tau_x : \mathbb{R}^{\mathbb{B}(\mathbb{Z}^d)} \rightarrow \mathbb{R}^{\mathbb{B}(\mathbb{Z}^d)}$ be the “translation by x ” which acts on configurations η by shifting the origin to position x ,

$$(\tau_x \eta)_{yz} := \eta_{y+x, z+x}, \quad (y, z) \in \mathbb{B}(\mathbb{Z}^d). \quad (2.4)$$

We say that μ is translation invariant if $\mu \circ \tau_x^{-1} = \mu$ for all $x \in \mathbb{Z}^d$, and that it is ergodic if $\mu(A) \in \{0, 1\}$ for every event A such that $\tau_x^{-1}(A) = A$ for all $x \in \mathbb{Z}^d$.

2.2. Scaling limit. As is usual for problems involving random fields, we will interpret samples from gradient Gibbs measures as random linear functionals on an appropriate space of functions. Let $C_0^\infty(\mathbb{R}^d)$ denote the set of all infinitely differentiable functions $f: \mathbb{R}^d \rightarrow \mathbb{R}$ with compact support. Given a configuration $\eta = (\eta_b)$ of gradients satisfying the conditions (2.1), we can find a configuration of the field $\phi = (\phi_x)$ such that (1.2) holds for every nearest-neighbor pair of sites. The configuration ϕ is determined uniquely once we fix the value at one site, e.g., ϕ_0 . For any function $f \in C_0^\infty(\mathbb{R}^d)$ we introduce the random linear functional

$$\phi(f) := \int dx f(x) \phi_{[x]} \quad (2.5)$$

which, under the condition

$$\int dx f(x) = 0, \quad (2.6)$$

does not depend on the choice of the special value ϕ_0 .

The functional $\phi(f)$ can be naturally extended to a somewhat larger space defined as follows. Let Δ denote the Laplace differential operator in \mathbb{R}^d and consider the set

$$\mathcal{H}_0 := \{\Delta g: g \in C_0^\infty(\mathbb{R}^d)\}. \quad (2.7)$$

Note that each $f \in \mathcal{H}_0$ automatically obeys (2.6). The set \mathcal{H}_0 is endowed with a natural quadratic form $f \mapsto (f, f) + (f, -\Delta^{-1}f)$ defined as

$$(\Delta g, \Delta g) + (\Delta g, -\Delta^{-1}\Delta g) = \int_{\mathbb{R}^d} dx (|\Delta g(x)|^2 + |\nabla g(x)|^2). \quad (2.8)$$

We thus define the norm

$$\|f\|_{\mathcal{H}} := [(f, f) + (f, -\Delta^{-1}f)]^{1/2} \quad (2.9)$$

and let \mathcal{H} be the completion of \mathcal{H}_0 in this norm. Note that \mathcal{H} is the case of $k = -1/2$ in the family of Sobolev spaces $W^{k,2}(\mathbb{R}^d)$. The condition that $(f, -\Delta^{-1}f) < \infty$ is natural once we realize that this quantity will represent the variance of the limiting Gaussian field.

The extension of ϕ to \mathcal{H} is implied by the following lemma:

LEMMA 2.2. *Suppose ϱ in (1.4) has support bounded away from zero and let μ be a translation-invariant, ergodic zero-tilt gradient Gibbs measure for the potential V . Then there exists a constant $c < \infty$ such that for each $f \in \mathcal{H}_0$,*

$$\|\phi(f)\|_{L^2(\mu)} \leq c\|f\|_{\mathcal{H}}. \quad (2.10)$$

In particular, ϕ extends to a linear functional $\phi: \mathcal{H} \rightarrow \mathbb{R}$.

Note that (2.10) means that the map $f \mapsto \phi(f)$ is continuous in L^2 -norm. If we want to avoid questions about accumulations of null sets, this permits us to work with only a countable number of f 's at each time. (In particular, we do not claim that $f \mapsto \phi(f)$ is continuous in any pointwise sense.) This will not pose any problem because we will content ourselves with the (weaker) definition of a Gaussian free field based on the standard approach via Gaussian Hilbert spaces (cf. Sheffield [26, Sect. 2.4]):

DEFINITION 2.3. *We say that a family $\{\psi(f): f \in \mathcal{H}\}$ of random variables on a probability space (Ω, \mathcal{F}, P) is a Gaussian free field if the map $f \mapsto \psi(f)$ is linear a.s. and each $\psi(f)$ is Gaussian with mean zero and variance*

$$E(\psi(f)^2) = (f, -\Delta^{-1}f). \quad (2.11)$$

Our goal is to show that the family of random variables $\{\phi(f): f \in \mathcal{H}\}$ has asymptotically, in the scaling limit, the law of a linear transformation of a Gaussian free field. To pass to this limit, we have to impose that the test functions are slowly varying, which we take to be on the scale ϵ^{-1} . For $\epsilon > 0$ and a function $f: \mathbb{R}^d \rightarrow \mathbb{R}$, let

$$f_\epsilon(x) := \epsilon^{(d/2+1)}f(\epsilon x) \quad (2.12)$$

and note that the normalization ensures

$$\begin{aligned} \|f_\epsilon\|_{\mathcal{H}}^2 &= (f_\epsilon, f_\epsilon) + (f_\epsilon, (-\Delta)^{-1}f_\epsilon) \\ &= \epsilon^2(f, f) + (f, (-\Delta)^{-1}f) \underset{\epsilon \leq 1}{\leq} \|f\|_{\mathcal{H}}^2. \end{aligned} \quad (2.13)$$

Let ϕ_ϵ denote the linear functional acting on $f \in C_0^\infty(\mathbb{R}^d)$ via

$$\phi_\epsilon(f) := \phi(f_\epsilon) = \int dx f(x) (\epsilon^{(-d/2+1)}\phi_{\lfloor x/\epsilon \rfloor}). \quad (2.14)$$

The main theorem is the Gaussian scaling limit for $\phi_\epsilon(f)$:

THEOREM 2.4 (Scaling to GFF). *Suppose V is as in (1.4) with q compactly supported in $(0, \infty)$. Let μ be a gradient Gibbs measure for the potential V which we assume to be ergodic with respect to the translations of \mathbb{Z}^d and to have zero tilt. Then for every $f \in \mathcal{H}$,*

$$\lim_{\epsilon \downarrow 0} E_\mu(e^{i\phi_\epsilon(f)}) = \exp\left\{\frac{1}{2} \int dx f(x) (Q^{-1}f)(x)\right\}, \quad (2.15)$$

where Q^{-1} is the inverse of the operator

$$Qf := \sum_{i,j=1}^d q_{ij} \frac{\partial^2}{\partial x_j \partial x_i} f, \quad (2.16)$$

with (q_{ij}) denoting some positive semi-definite, non-degenerate, $d \times d$ matrix. In other words, the law of ϕ_ϵ on the linear dual \mathcal{E}' of any finite-dimensional linear subspace $\mathcal{E} \subset \mathcal{H}$ converges weakly to that of a Gaussian field with mean zero and covariance $(-Q)^{-1}$.

REMARKS 2.5. Here are some additional observations and remarks concerning the model under consideration and the results above:

- (1) Since $(-Q)$ is dominated by a multiple of $(-\Delta)$ from below, the integral in (2.15) — interpreted as the quadratic form $(f, Q^{-1}f)$ — is well defined for all $f \in \mathcal{H}$.
- (2) Note that, in (2.14), the individual ϕ 's get scaled by $\epsilon^{(-d/2+1)}$ and not $\epsilon^{-d/2}$ as one might expect from the conventional central limiting reasoning. This has to do with the fact that the variables (ϕ_x) are strongly correlated. These correlations are weaker for the gradients $\eta_{xy} := \phi_y - \phi_x$ which adhere to the “usual” central limit scaling. In $d = 1$ and for general potentials V , the increments η_b are in fact i.i.d. and the scaling limit follows from the standard CLT.
- (3) In $d > 1$ the matrix (q_{ij}) is not necessarily a multiple of unity, since, in general, μ is not guaranteed to be invariant under reflections and rotations of \mathbb{Z}^d . (Nevertheless, we expect that every zero-tilt, translation-invariant, ergodic measure for the isotropic interaction (2.2) will inherit these symmetries.) To get convergence of ϕ_ϵ to GFF in the sense of Definition 2.3, one thus has to scale the argument of ϕ by the root of the corresponding eigenvalue of q in each of its principal directions.
- (4) The absence of strict convexity does not permit us to use the general argument of Funaki and Spohn [16] for the *existence* of an ergodic GGM with zero (or any other prescribed) tilt. To show that such GGMs do exist — and that our Theorem 2.4 is not vacuous — we note that, by

Lemma 4.8 of Biskup and Kotecký [5], every weak limit of torus measures exhibits exponential concentration of the empirical tilt; one then just needs to pick any ergodic component. Note that this lemma applies only to zero tilt (cf. [5, Remark 4.9]).

- (5) The restriction to zero tilt is actually a significant drawback of our analysis. The main reason is our inability to characterize the scaling limit of the so called corrector for the corresponding random walk problem. See Sect. 6 for more details.
- (6) In the example studied by Biskup and Kotecký [5], cf. (1.5), the two GGMs coexisting at the transitional value p_t of p were proved to exhibit different characteristic fluctuations. It follows that the corresponding scaling limits will be distinguished by their stiffness coefficients q_{ij} . Moreover, by Theorem 2.5 of [5], for $\kappa_1 \gg \kappa_2$ the transition in $d = 2$ model with (1.5) lies on a self-dual line, i.e.,

$$\frac{p_t}{1 - p_t} = \left(\frac{\kappa_2}{\kappa_1}\right)^{1/4}. \quad (2.17)$$

The transition presumably stays on this line even as one slides the ratio κ_1/κ_2 towards one. However, it disappears before κ_1/κ_2 hits one because for $\kappa_1 \approx \kappa_2$ the potential V is convex and so there is only one GGM with zero tilt [16]. At such point of disappearance physicists often expect non-trivial critical fluctuations. Notwithstanding, our results show this is not the case.

- (7) We avoid the context of the “stronger” definition of GFF as a random element in an appropriate Banach space (cf. Gross [14] or Sheffield [26, Sect. 2.2]). This definition is appealing in $d = 1$, where the limiting functional $f \mapsto \psi(f)$ actually admits the integral representation

$$\psi(f) = \int_{\mathbb{R}} f(t) \psi_t dt \quad (2.18)$$

with $t \mapsto \psi_t$ denoting a continuous diffusion with generator Q , but in $d > 1$ the corresponding field becomes less and less regular with increasing dimension and the appeal is lost. However, this context would be in demand if one wishes to discuss the notion of tightness and convergence in law for the limit in Theorem 2.4.

Both Lemma 2.2 and Theorem 2.4 are proved in Sect. 4.

3. Extended gradient Gibbs measures.

3.1. *Coupling to random conductance model.* The key idea underlying the representation (1.4) is that the auxiliary variable κ in the expression for V may be elevated to a genuine degree of freedom associated with the corresponding edge. Specifically, given a gradient Gibbs measure μ with potential (1.4), for each finite $\Lambda \subset \mathbb{B}(\mathbb{Z}^d)$ consider the measure $\tilde{\mu}_\Lambda$ on $\mathbb{R}^{\mathbb{B}(\mathbb{Z}^d)} \times \mathbb{R}^\Lambda$ defined by

$$\tilde{\mu}_\Lambda(\mathcal{A} \times \mathcal{B}) := \int_{\mathcal{B}} \prod_{b \in \Lambda} \varrho(d\kappa_b) E_\mu \left(\mathbf{1}_{\mathcal{A}} \prod_{b \in \Lambda} e^{V(\eta_b) - \frac{1}{2}\kappa_b \eta_b^2} \right), \quad (3.1)$$

where $\mathcal{A} \subset \mathbb{R}^{\mathbb{B}(\mathbb{Z}^d)}$ and $\mathcal{B} \subset \mathbb{R}^\Lambda$ are Borel sets. The representation (1.4) ensures that $(\tilde{\mu}_\Lambda)$ is a consistent family of measures; by Kolmogorov's Extension Theorem these are projections from a unique measure $\tilde{\mu}$ on configurations $(\eta_b, \kappa_b) \in \mathbb{R}^{\mathbb{B}(\mathbb{Z}^d)} \times \mathbb{R}^{\mathbb{B}(\mathbb{Z}^d)}$. The restriction of $\tilde{\mu}$ on the η 's gives us back μ ; we call $\tilde{\mu}$ an *extension* of μ . The measure $\tilde{\mu}$ is Gibbs for the Hamiltonian $\sum_{\langle x, y \rangle} \frac{1}{2} \kappa_{xy} \eta_{xy}^2$, so we will refer to it as an *extended gradient Gibbs measure* (see Biskup and Kotecký [5] for further facts on the extended GGMs).

To ease the notation, whenever b is an edge between x and y , we may interchangeably write κ_b and κ_{xy} for the same quantity. Furthermore, for the same reasons it will even be convenient to assume

$$\kappa_{xy} = \kappa_{yx}, \quad |x - y| = 1, \quad (3.2)$$

and work with the κ 's as symmetric objects.

We proceed by a series of lemmas that characterize the properties of $\tilde{\mu}$.

LEMMA 3.1. *Let μ be a gradient Gibbs measure for the potential V and let $\tilde{\mu}$ be its extension to $\mathbb{R}^{\mathbb{B}(\mathbb{Z}^d)} \times \mathbb{R}^{\mathbb{B}(\mathbb{Z}^d)}$. Consider the σ -field $\mathcal{E} := \sigma(\{\eta_b : b \in \mathbb{B}(\mathbb{Z}^d)\})$. For $\tilde{\mu}$ -a.e. η , the regular conditional distribution $\tilde{\mu}(-|\mathcal{E})(\eta)$, regarded as a measure on the κ 's, takes the product form*

$$\tilde{\mu}(d\kappa|\mathcal{E})(\eta) = \bigotimes_{b \in \mathbb{B}(\mathbb{Z}^d)} [e^{V(\eta_b) - \frac{1}{2}\kappa_b \eta_b^2} \varrho(d\kappa_b)]. \quad (3.3)$$

PROOF. Recall that $\mathcal{E}_{\mathbb{B}(\Lambda)} := \sigma(\{\eta_b : b \in \mathbb{B}(\Lambda)\})$. The identity (3.1) implies that $\tilde{\mu}_\Lambda$ coincides with $\tilde{\mu}$ on $\mathbb{R}^{\mathbb{B}(\Lambda)} \times \mathbb{R}^{\mathbb{B}(\Lambda)}$. But $\mu_\Lambda(-|\mathcal{E}_\Lambda)$ has the desired product form by definition and so the claim follows by standard approximation arguments. \square

LEMMA 3.2. *Let μ be a gradient Gibbs measure and let $\tilde{\mu}$ be its extension to $\mathbb{R}^{\mathbb{B}(\mathbb{Z}^d)} \times \mathbb{R}^{\mathbb{B}(\mathbb{Z}^d)}$. If μ is translation-invariant and ergodic then so is $\tilde{\mu}$.*

PROOF. The uniqueness of the extension of measures (3.1) implies that $\tilde{\mu}$ is translation-invariant if μ is translation invariant, and so it remains to prove that ergodicity is inherited as well. Let $A \subset \mathbb{R}^{\mathbb{B}(\mathbb{Z}^d)} \times \mathbb{R}^{\mathbb{B}(\mathbb{Z}^d)}$ be a translation-invariant event, i.e., $(\eta, \kappa) \in A$ iff $(\tau_x \eta, \tau_x \kappa) \in A$ for all x . Our task is to show that $\tilde{\mu}(A) \in \{0, 1\}$.

First we invoke the ergodicity of μ . Consider the function

$$f(\eta) := E_{\tilde{\mu}}(1_A | \mathcal{E})(\eta). \quad (3.4)$$

Since A and $\tilde{\mu}$ are translation invariant, we have

$$f(\tau_x \eta) = E_{\tilde{\mu}}(1_A | \mathcal{E})(\tau_x \eta) = E_{\tilde{\mu}}(1_A \circ \tau_x^{-1} | \mathcal{E})(\eta) = f(\eta), \quad \tilde{\mu}\text{-a.s.} \quad (3.5)$$

But f is \mathcal{E} -measurable and the restriction of $\tilde{\mu}$ to \mathcal{E} is μ , which we assumed to be ergodic. Hence f is constant almost surely. Let c denote this constant.

We will use a standard approximation argument to show that $c \in \{0, 1\}$. Since A is an event from the product σ -algebra, there exists a sequence of events

$$A_n \in \sigma(\{\eta_{xy}, \kappa_{xy} : |x - y| = 1, |x| \leq n\}) \quad (3.6)$$

such that

$$\tilde{\mu}(A \Delta A_n) \xrightarrow{n \rightarrow \infty} 0. \quad (3.7)$$

The bound

$$\|1_{A_n} - 1_A\|_{L^1(\tilde{\mu})} \leq \tilde{\mu}(A \Delta A_n) \quad (3.8)$$

shows that then $1_{A_n} \rightarrow 1_A$ in $L^1(\tilde{\mu})$. Since A is translation invariant, we have $1_A = 1_A 1_{\tau_x^{-1}(A)}$. Each indicator can be approximated by the indicator of the event A_n ; a simple bound gives

$$\|1_{A_n} 1_{\tau_x^{-1}(A_n)} - 1_A 1_{\tau_x^{-1}(A)}\|_{L^1(\tilde{\mu})} \leq 2\tilde{\mu}(A \Delta A_n). \quad (3.9)$$

For x with $|x| > 2n + 1$, the fact that $\tilde{\mu}(-| \mathcal{E})$ is a product measure (cf Lemma 3.1) implies that A_n and $\tau_x^{-1}(A_n)$ are independent. Hence,

$$E_{\tilde{\mu}}(1_{A_n} 1_{\tau_x^{-1}(A_n)} | \mathcal{E}) = E_{\tilde{\mu}}(1_{A_n} | \mathcal{E}) E_{\tilde{\mu}}(1_{\tau_x^{-1}(A_n)} | \mathcal{E}). \quad (3.10)$$

Rolling the approximations backwards, we then conclude that the left-hand side converges to c in $L^1(\tilde{\mu})$, while the right-hand side converges to c^2 (note that all expectations are bounded). It follows that $c = c^2$, i.e., $c \in \{0, 1\}$. As

$$\tilde{\mu}(A) = E_{\tilde{\mu}}(f) = c, \quad (3.11)$$

the proof is finished. \square

3.2. *Random walk connections.* Our next goal will be to characterize also the conditional measure given the κ 's. This will in turn require some facts from the theory of random walks with random conductances. We will frequently borrow facts from an associated potential theory which will be expounded in Sect. 5.

Let us pick a configuration $\kappa = (\kappa_b)$ with $\kappa_b \in (0, \infty)$, and recall the formula (1.6) for the generator \mathcal{L}_κ of the random walk among conductances κ . We will focus on the action of \mathcal{L}_κ on functions of both the environment κ and the position x that satisfy the following *shift-covariance* property

$$g(\kappa, x + b) - g(\kappa, x) = g(\tau_x \kappa, b), \quad (3.12)$$

with $x \in \mathbb{Z}^d$ and b a coordinate unit vector in \mathbb{R}^d , subject to the condition

$$g(\kappa, 0) = 0. \quad (3.13)$$

This makes the function completely determined by its values at the neighbors of the origin. A function of this kind is said to be *harmonic* for the above random walk if

$$\mathcal{L}_\kappa g(\kappa, \cdot) = 0 \quad (3.14)$$

for (almost) every κ . As it turns out, harmonic, shift-covariant functions are uniquely determined (a.s.) by their mean with respect to ergodic measures on the conductances:

LEMMA 3.3. *Let ν be a translation-invariant, ergodic probability measure on configurations $\kappa = (\kappa_b) \in \mathbb{R}^{\mathbb{B}(\mathbb{Z}^d)}$ such that $\nu(\epsilon \leq \kappa_b \leq 1/\epsilon) = 1$ for some $\epsilon > 0$. Let $g: \mathbb{R}^{\mathbb{B}(\mathbb{Z}^d)} \times \mathbb{Z}^d \rightarrow \mathbb{R}$ be a measurable function which is*

- (1) *harmonic in the sense of (3.14), ν -a.s.,*
- (2) *shift-covariant in the sense of (3.12–3.13), ν -a.s.,*
- (3) *square integrable in the sense that $E_\nu |g(\cdot, x)|^2 < \infty$ for all x with $|x| = 1$.*

If $E_\nu(g(\cdot, x)) = 0$ for all x with $|x| = 1$, then $g(\cdot, x) = 0$ a.s. for all $x \in \mathbb{Z}^d$.

We defer the proof, and further discussion of the consequences of shift-covariance and harmonicity, to Sect. 5. Returning to the gradient fields, we now characterize the conditional law given the κ 's:

LEMMA 3.4. *Let μ be a translation-invariant, ergodic gradient Gibbs measure with zero tilt and let $\tilde{\mu}$ be its extension to $\mathbb{R}^{\mathbb{B}(\mathbb{Z}^d)} \times \mathbb{R}^{\mathbb{B}(\mathbb{Z}^d)}$. Consider the σ -field*

$\mathcal{F} := \sigma(\{\kappa_b : b \in \mathbb{B}(\mathbb{Z}^d)\})$. For $\tilde{\mu}$ -a.e. κ , the conditional law $\tilde{\mu}(-|\mathcal{F})(\kappa)$, regarded as a measure on the set of configurations $\{(\phi_x) \in \mathbb{R}^{\mathbb{Z}^d} : \phi_0 = 0\}$ with the ϕ 's defined from the η 's via (1.2), is Gaussian with mean zero

$$E_{\tilde{\mu}}(\phi_x|\mathcal{F})(\kappa) = 0, \quad x \in \mathbb{Z}^d, \quad (3.15)$$

and the covariance given by $(-\mathcal{L}_\kappa)^{-1}$. Explicitly, for each $f : \mathbb{Z}^d \rightarrow \mathbb{R}$ with finite support and $\sum_x f(x) = 0$,

$$\text{Var}_{\tilde{\mu}}\left(\sum_x f(x)\phi_x \middle| \mathcal{F}\right)(\kappa) = \sum_x f(x)(-\mathcal{L}_\kappa^{-1}f)(x). \quad (3.16)$$

PROOF. The fact that the conditional measure is multivariate Gaussian law with covariance \mathcal{L}_κ^{-1} is checked by direct inspection of (3.1). The only non-trivial task is to identify the mean. First we note that the loop conditions (2.1) ensure that there exists a function $u : \mathbb{R}^{\mathbb{B}(\mathbb{Z}^d)} \times \mathbb{Z}^d \rightarrow \mathbb{R}$ such that

$$u(\kappa, 0) = 0 \quad (3.17)$$

and

$$u(\kappa, x+b) - u(\kappa, x) = E_{\tilde{\mu}}(\eta_{x,x+b}|\mathcal{F})(\kappa) \quad (3.18)$$

for all unit vectors b in the coordinate directions. We claim that u is harmonic in the sense of (3.14). Indeed,

$$\mathcal{L}_\kappa u(\kappa, x) = E_{\tilde{\mu}}\left(\sum_{y:|y-x|=1} \kappa_{xy}(\phi_y - \phi_x) \middle| \mathcal{F}\right)(\kappa), \quad (3.19)$$

where we write, thanks to the loop conditions, $\eta_{xy} = \phi_y - \phi_x$. Using that $\tilde{\mu}$ is Gibbs, we can now also condition on $\sigma(\phi_y : y \neq x)$; the conditional measure $\mu_{\{x\}}$ is Gaussian with the explicit form

$$\mu_{\{x\}}(d\phi_x) = \frac{1}{Z} \exp\left\{-\frac{1}{2}\phi_x^2 \sum_{y:|y-x|=1} \kappa_{xy} + \phi_x \sum_{y:|x-y|=1} \kappa_{xy}\phi_y\right\} d\phi_x, \quad (3.20)$$

where Z is an appropriate normalization constant. It is easy to check that the mean of $\phi_x \sum_{y:|y-x|=1} \kappa_{xy}$ under $\mu_{\{x\}}$ is exactly $\sum_{y:|y-x|=1} \kappa_{xy}\phi_y$, proving $\mathcal{L}_\kappa u(\kappa, x) = 0$.

Next we observe that the translation invariance of $\tilde{\mu}$ implies

$$\begin{aligned} u(\tau_x \kappa, b) - u(\tau_x \kappa, 0) &= E_{\tilde{\mu}}(\eta_{0,b}|\mathcal{F})(\tau_x \kappa) \\ &= E_{\tilde{\mu}}(\eta_{x,x+b}|\mathcal{F})(\kappa) = u(\kappa, x+b) - u(\kappa, x) \end{aligned} \quad (3.21)$$

and so u is shift-covariant as defined in (3.12–3.13). Finally, the definition of u and the fact that $\tilde{\mu}$ has zero tilt imply

$$E_{\tilde{\mu}}(u(\cdot, x)) = E_{\tilde{\mu}}(\eta_{0,x}) = 0, \quad |x| = 1. \quad (3.22)$$

As u obeys all conditions of Lemma 3.3, we have $E_{\tilde{\mu}}(\phi_x | \mathcal{F}) = u(\cdot, x) = 0$ $\tilde{\mu}$ -a.s. \square

Our reference to the random walk with generator \mathcal{L}_κ is not limited to Lemma 3.3; we will also need to know some specific properties of this random walk. First, we will need to know that the position of the walk satisfies a Central Limit Theorem. Let $X = (X_t)$ denote the continuous-time random walk with the generator \mathcal{L}_κ and let P_κ^x denote the law of X subject to the initial condition $P_\kappa^x(X_0 = x) = 1$. The following lemma goes back to Kipnis and Varadhan [18].

LEMMA 3.5 (Annealed CLT). *Let μ be a measure on $\mathbb{R}^{\mathbb{B}(\mathbb{Z}^d)}$ which is translation-invariant and ergodic, and obeys $\mu(\epsilon \leq \omega_b \leq 1/\epsilon) = 1$ for some $\epsilon > 0$. Then there exists a positive semidefinite, non-degenerate, $d \times d$ matrix \mathfrak{q} such that, for every $t > 0$, the annealed distribution $E_\mu P_\kappa^0(\epsilon X_{te^{-2}} \in \cdot)$ converges weakly to the law of the multivariate normal $\mathcal{N}(0, t\mathfrak{q})$.*

The main result of [18] actually shows that the annealed law of the entire path $t \mapsto \epsilon X_{te^{-2}}$ converges to that of (a linear transform of) Brownian motion. However the above is all what will be needed for the purposes of the present paper.

Apart from a central-limit asymptotics, we will also need to have an estimate on the heat kernel of the above random walk. The following lemma is a consequence of the main result of Delmotte and Deuschel [10].

LEMMA 3.6 (Heat-kernel upper bound). *Let μ be a law of the conductances satisfying the ellipticity condition $\mu(\epsilon < \kappa_b < 1/\epsilon) = 1$ for some $\epsilon > 0$. Then there is a $c_1 < \infty$ such that*

$$E_\mu |\nabla_i \nabla_j P_\kappa^0(X_t = \cdot)| \leq \frac{c_1}{t^{d/2+1}}, \quad 1 \leq i, j \leq d, t > 0. \quad (3.23)$$

Here ∇_i is the discrete spatial derivative in the i -th coordinate direction, i.e., $\nabla_i f(x) := f(x + \hat{e}_i) - f(x)$.

PROOF. By formula (1.5b) in [10, Theorem 1.1],

$$E_\mu |\nabla_i \nabla_j P_\kappa^0(X_t = x)| \leq c_1' \frac{P_\kappa^{c_2 t}(0, x)}{t}, \quad (3.24)$$

where $P^{c_2^t}(0, x)$ is the probability of the continuous-time simple random walk to be at x at time c_2^t . This probability is bounded from above by a constant times $t^{-d/2}$. \square

4. Proof of main result.

4.1. *Regularity estimates.* The goal of this section is to prove Theorem 2.4 concerning the scaling limit of $\phi(f_\epsilon)$. We begin by proving L^2 -continuity of the random functional $f \mapsto \phi(f)$ on \mathcal{H} as stated in Lemma 2.2. For convenience of notation, whenever \mathcal{R} is an operator on $\ell^2(\mathbb{Z}^d)$ we will extend it to an operator on $L^2(\mathbb{R}^d)$ via the formula

$$(f, \mathcal{R}f) := \int dx dy f(x)f(y) \mathcal{R}(\lfloor x \rfloor, \lfloor y \rfloor), \quad (4.1)$$

where $\mathcal{R}(x, y)$ is the kernel of \mathcal{R} in the canonical basis in $\ell^2(\mathbb{Z}^d)$.

PROOF OF LEMMA 2.2. Let $\tilde{\mu}$ be the extended gradient Gibbs measure corresponding to μ . Recall the notation \mathcal{L}_κ for the generator of the random walk among conductances $\kappa = (\kappa_b)$ and let \mathcal{L} denote the generator of the simple-random walk (i.e., the special case of \mathcal{L}_κ when all $\kappa_b = 1$). Pick $f \in \{\Delta g: g \in C_0^\infty(\mathbb{R}^d)\}$. Lemma 3.4 and the fact that $\phi(f)$ is linear in the η 's imply

$$\|\phi(f)\|_{L^2(\tilde{\mu})}^2 = E_{\tilde{\mu}}((f, -\mathcal{L}_\kappa^{-1}f)) \quad (4.2)$$

where $(f, \mathcal{L}_\kappa^{-1}f)$ is as defined above. By assumption on the support of ϱ we know that $\kappa_b \geq a$ $\tilde{\mu}$ -a.s. by which we immediately have the operator inequalities

$$(-\mathcal{L}_\kappa) \geq a(-\mathcal{L}) \quad \text{and} \quad (-\mathcal{L}_\kappa)^{-1} \leq a^{-1}(-\mathcal{L})^{-1}. \quad (4.3)$$

Therefore, it suffices to bound the quadratic form associated with the (homogeneous) discrete Laplacian \mathcal{L} by the quadratic form defining the \mathcal{H} -space:

$$(f, (-\mathcal{L})^{-1}f) \leq c\|f\|_{\mathcal{H}}^2 \quad (4.4)$$

for some constant $c < \infty$ and all f in a dense subset of \mathcal{H} .

To this end we pick $f \in \mathcal{H}$ in the Schwartz class and let

$$\hat{f}(k) := (2\pi)^{-d/2} \int f(x)e^{ik \cdot x} dx \quad (4.5)$$

be its (L^2 -norm preserving) Fourier transform. A direct calculation now yields

$$(f, (-\mathcal{L})^{-1}f) = \int_{[-\pi, \pi]^d} dk \frac{1}{(-\hat{\mathcal{L}})(k)} \left| \sum_{\substack{k': \exists \ell \in \mathbb{Z}^d \\ k-k'=2\pi\ell}} \hat{f}(k') \prod_{j=1}^d \frac{1 - e^{-ik'_j}}{ik'_j} \right|^2, \quad (4.6)$$

where

$$(-\hat{\mathcal{L}})(k) := 4 \sum_{j=1}^d \sin^2(k_j/2) \quad (4.7)$$

is the generalized eigenvalue of the lattice Laplacian. Introducing $-\hat{\Delta}(k) = |k|^2$ to denote the corresponding quantity for the continuum Laplacian, we invoke the Cauchy-Schwarz inequality for the sum over k' to get

$$(f, (-\mathcal{L})^{-1}f) \leq \int_{[-\pi, \pi]^d} dk \left\{ \sum_{\substack{k': \exists \ell \in \mathbb{Z}^d \\ k-k'=2\pi\ell}} |\hat{f}(k')|^2 (1 - \hat{\Delta}(k')^{-1}) \right\} c(k) \quad (4.8)$$

where

$$c(k) := \sum_{\substack{k': \exists \ell \in \mathbb{Z}^d \\ k-k'=2\pi\ell}} \frac{1}{(-\hat{\mathcal{L}})(k)} \frac{-\hat{\Delta}(k')}{1 - \hat{\Delta}(k')} \prod_{j=1}^d \left(\frac{2 \sin(k'_j/2)}{k'_j} \right)^2 \quad (4.9)$$

is well defined on the set $\mathcal{B} := \{k \in [-\pi, \pi]^d : k_j \neq 0, j = 1, \dots, d\}$ of full Lebesgue measure in $[-\pi, \pi]^d$. We claim that

$$c := \sup_{k \in \mathcal{B}} c(k) < \infty. \quad (4.10)$$

We will show this by proving that the summand in (4.9) is bounded by a constant times the product $\prod_j (|k'_j| + 1)^{-2}$. Indeed, for the $k' = k$ term we use that the ratio $\hat{\Delta}(k)/\hat{\mathcal{L}}(k)$ is bounded throughout \mathcal{B} , and same for the ratios $k \mapsto \sin(k_j/2)/k_j$. When $k' \neq k$, we set i to be the first index j where $k'_j \neq k_j$ and bound the $4 \sin(k'_i/2)^2$ term by $-\hat{\mathcal{L}}(k)$. Then we bound the ratio of $\hat{\Delta}(k)$ -terms by unity and the j -th term in the product by a constant times $(|k'_j|^2 + 1)^{-2}$. (The $\sin(k'_j/2)$ -term is not needed because $|k'_i| \geq 2\pi$.)

The product $\prod_j (|k'_j| + 1)^{-2}$ is summable over $k' \in k + (2\pi\mathbb{Z})^d$ uniformly in $k \in \mathcal{B}$ and so (4.10) is proved. Bounding $c(k)$ by its supremum in (4.8), we can merge the sum and the integral to get $\|f\|_{\mathcal{H}}^2$. The desired bound (4.4) follows. \square

REMARK 4.1. The inclusion of L^2 -norm of f into $\|f\|_{\mathcal{H}}$ is crucial for the bound (4.4). Indeed, on the basis of (4.6) it is not hard to construct functions for which the ratio $(f, -\mathcal{L}^{-1}f)/(f, -\Delta^{-1}f)$ is arbitrarily large. This is caused by the fact that the spectrum of $-\Delta$ extends all the way to infinity while that of $-\mathcal{L}$ is bounded.

The continuity established in Lemma 2.2 allows us to work only with smooth and compactly-supported test functions. We will nevertheless need one more regularity bound before we can delve into the proof of our main result:

LEMMA 4.2. *Let μ be a translation-invariant law on the conductances subject to the ellipticity condition, $\mu(\epsilon < \kappa_b < 1/\epsilon) = 1$ for some $\epsilon > 0$. Then there exists $c < \infty$ such that whenever $f = \Delta g$ for some $g \in C_0^\infty(\mathbb{R}^d)$,*

$$E_\mu(f, e^{t\mathcal{L}_\kappa} f) \leq c \|\nabla g\|_\infty^2 \lambda(\text{supp } g)^2 \frac{1}{t^{d/2+1}}, \quad (4.11)$$

where $\lambda(A)$ is the set function on Borel subsets of \mathbb{R}^d defined by

$$\lambda(A) := \sum_{x \in \mathbb{Z}^d} \sum_{i=1}^d \int_{R_i} dz \mathbf{1}_{\{x+z \in A\}} \quad (4.12)$$

with R_i denoting the set of points in $[0, 1]^d$ with vanishing i -th coordinate.

PROOF. Translation invariance of μ and a simple integration by parts tells us

$$\begin{aligned} & E_\mu(\Delta g, e^{t\mathcal{L}_\kappa} \Delta g) \\ &= \sum_{x, y \in \mathbb{Z}^d} \int_{[0, 1]^d} dz \int_{[0, 1]^d} dz' \Delta g(x+z) \Delta g(y+z') E_\mu P_\kappa^x(X_t = y) \\ &= \sum_{x, y \in \mathbb{Z}^d} \sum_{i, j=1}^d \int_{R_i} dz \int_{R_j} dz' \partial_i g(x+z) \partial_j g(y+z') \nabla_i \nabla_j E_\mu P_\kappa^x(X_t = y), \end{aligned} \quad (4.13)$$

where ∂_i stands for the partial derivative with respect to the i -th coordinate. Restricting the integrations and sums so that the arguments $x+z$ and $y+z'$ are in the support of g , bounding the partial derivatives by $\|\nabla g\|_\infty$, and applying the estimate (3.23), we obtain the desired bound. \square

The consequence of Lemma 4.2 that will concern us is as follows:

COROLLARY 4.3. For μ as in Lemma 4.2 and any $f \in \{\Delta g : g \in C_0^\infty(\mathbb{R}^d)\}$,

$$\lim_{M \rightarrow \infty} \sup_{0 < \epsilon < 1} \int_M^\infty dt E_\mu \epsilon^{-2} (f_\epsilon, e^{t\epsilon^{-2}\mathcal{L}_\kappa} f_\epsilon) = 0. \quad (4.14)$$

PROOF. Pick f of the form $f = \Delta g$ and note that $f_\epsilon = \Delta g^{(\epsilon)}$, where $g^{(\epsilon)}(x) := \epsilon^{d/2+1}g(x\epsilon)$. First we observe

$$\|\nabla g^{(\epsilon)}\|_\infty = \epsilon^{d/2} \|\nabla g\|_\infty. \quad (4.15)$$

Next we note that, since the support of g is the closure of a non-empty bounded open set, a simple covering argument tells us that

$$\epsilon^d \lambda(\text{supp } g^{(\epsilon)}) \xrightarrow{\epsilon \downarrow 0} d |\text{supp } g|, \quad (4.16)$$

where $|\text{supp } g|$ is the Lebesgue measure of $\text{supp } g$. As a consequence, there exists a constant $C(g) < \infty$ such that

$$\lambda(\text{supp } g^{(\epsilon)}) \leq C(g) \epsilon^{-d}, \quad 0 < \epsilon < 1. \quad (4.17)$$

Plugging (4.15) and (4.17) into (4.11), we get for $\epsilon \in (0, 1)$,

$$E_\mu \epsilon^{-2} (f_\epsilon, e^{t\epsilon^{-2}\mathcal{L}_\kappa} f_\epsilon) \leq c \|\nabla g\|_\infty^2 C(g)^2 \frac{1}{t^{d/2+1}}. \quad (4.18)$$

The functions on the left (indexed by ϵ) are uniformly integrable in t in all $d \geq 1$. \square

4.2. *Scaling limit.* Having dispensed with regularity considerations, we can now proceed to establish the principal fact underlying the proof of Theorem 2.4:

PROPOSITION 4.4. Let μ be a translation-invariant, ergodic measure on $\kappa = (\kappa_b) \in \mathbb{R}^{\mathbb{B}(\mathbb{Z}^d)}$ such that $\mu(\delta \leq \kappa_b \leq 1/\delta) = 1$ for some $\delta > 0$. Then there exists a positive semidefinite, non-degenerate, $d \times d$ matrix $\mathfrak{q} = (q_{ij})$ such that

$$\lim_{\epsilon \downarrow 0} (f_\epsilon, (-\mathcal{L}_\kappa)^{-1} f_\epsilon) = (f, (-Q)^{-1} f) \quad (4.19)$$

in μ -probability for each $f \in C_0^\infty(\mathbb{R}^d) \cap \mathcal{H}$, where Q is defined from \mathfrak{q} by (2.16).

Key to the proof is the following lemma:

LEMMA 4.5. For every $t > 0$ and any $f \in C_0^\infty(\mathbb{R}^d) \cap \mathcal{H}$

$$\Theta_\epsilon(t) := \epsilon^{-2}(f_\epsilon, e^{t\epsilon^{-2}\mathcal{L}_\kappa} f_\epsilon) - (f, e^{tQ} f) \xrightarrow[\epsilon \downarrow 0]{} 0 \quad \text{in } L^2(\mu). \quad (4.20)$$

PROOF. Let $(X_t)_{t \geq 0}$ be the continuous-time random walk with the generator \mathcal{L}_κ and let P_κ^x denote the law of the walk started from x . By Lemma 3.5, the annealed law of $\epsilon X_{t\epsilon^{-2}}$ tends weakly to that of multivariate normal

$$\mathcal{N}_t := \mathcal{N}(0, t\mathfrak{q}) \quad (4.21)$$

for some positive semidefinite, non-degenerate, $d \times d$ matrix $\mathfrak{q} = (q_{ij})$. As a consequence, if $\mathcal{G} \subset C_0^\infty(\mathbb{R}^d)$ is a family of functions that are uniformly equicontinuous and bounded, we have

$$E_\mu \left(\sup_{g \in \mathcal{G}} \left| E_\kappa^0(g(\epsilon X_{t\epsilon^{-2}})) - E g(\mathcal{N}_t) \right|^2 \right) \xrightarrow[\epsilon \downarrow 0]{} 0. \quad (4.22)$$

Now let Q be the generator of the Brownian motion with mean zero and covariance \mathfrak{q} ; i.e., Q is the operator in (2.16). Then we have

$$(f, e^{tQ} f) = \int dy f(y) E(f(y + \mathcal{N}_t)). \quad (4.23)$$

As \mathcal{L}_κ is the generator of the random walk (X_t) , we derive similarly

$$\epsilon^{-2}(f_\epsilon, e^{t\epsilon^{-2}\mathcal{L}_\kappa} f_\epsilon) = e^d \int_{[0,1]^d \times [0,1]^d} dz_1 dz_2 \sum_{x \in \mathbb{Z}^d} f(\epsilon x + \epsilon z_1) E_\kappa^x(f(\epsilon z_2 + \epsilon X_{t\epsilon^{-2}})). \quad (4.24)$$

We thus need to show that the right-hand side of (4.24) tends to that of (4.23). Note that if f is supported in $[-M, M]^d$, the integral in (4.23) can be restricted to $y \in [-M, M]^d$ and the sum over x in (4.24) to, say, $|x| \leq 2M/\epsilon$ (once $\epsilon \ll 1$).

Substituting $y = \epsilon x + \epsilon z$ with $x \in \mathbb{Z}^d$, $|x| \leq 2M/\epsilon$, and $z \in [0, 1]^d$ in (4.23) allows us to put both terms on the same footing. Subtracting (4.23) from (4.24), taking expectation with respect to μ and applying the Cauchy-Schwarz inequality we thus get

$$\begin{aligned} & E_\mu \Theta_\epsilon(t)^2 \\ & \leq \|f\|_2^2 e^d \int_{[0,1]^d} dz \sum_{x: |x| \leq 2M/\epsilon} E_\mu \left| E_\kappa^x f(\epsilon z + \epsilon X_{t\epsilon^{-2}}) - E f(\epsilon x + \epsilon z + \mathcal{N}_t) \right|^2. \end{aligned} \quad (4.25)$$

Using translation invariance of μ , we may replace $E_\kappa^x f(\epsilon z + \epsilon X_{t\epsilon^{-2}})$ by the expression $E_\kappa^0 f(\epsilon x + \epsilon z + \epsilon X_{t\epsilon^{-2}})$ inside the expectation. For $f \in C_0^\infty(\mathbb{R}^d)$,

$$\mathcal{G} := \{f(\epsilon x + \epsilon z + \cdot) : |x| \leq 2M/\epsilon, z \in [0, 1]^d\} \quad (4.26)$$

is an equicontinuous family of uniformly bounded functions. Then (4.22) tells us that the right-hand side of (4.25) tends to zero as $\epsilon \downarrow 0$, which proves the desired claim (4.20). \square

PROOF OF PROPOSITION 4.4. To extract (4.19) from (4.20), we note that for any $f \in C_0^\infty(\mathbb{R}^d) \cap \mathcal{H}$,

$$(f, (-\mathcal{L}_\kappa)^{-1}f) = \int_0^\infty dt (f, e^{t\mathcal{L}_\kappa} f). \quad (4.27)$$

Replacing f by f_ϵ and scaling t by ϵ^2 , we find

$$(f_\epsilon, (-\mathcal{L}_\kappa)^{-1}f_\epsilon) = \int_0^\infty dt \epsilon^{-2} (f_\epsilon, e^{t\epsilon^{-2}\mathcal{L}_\kappa} f_\epsilon). \quad (4.28)$$

By (4.20) the function under integration tends to $(f, e^{tQ} f)$ in probability for each t ; the monotonicity in t (and continuity of the limit) ensures that the convergence is actually uniform (in probability) on compact intervals. By Corollary 4.3, the integral can be truncated to a finite interval in L^1 -norm, and similarly for the integral of the limit which is finite as f is in the domain of Q^{-1} . It follows that

$$(f_\epsilon, (-\mathcal{L}_\kappa)^{-1}f_\epsilon) \xrightarrow{\epsilon \downarrow 0} \int_0^\infty dt (f, e^{tQ} f) = (f, -Q^{-1}f), \quad (4.29)$$

in μ -probability (and $L^1(\mu)$). This is the desired conclusion (4.19). \square

REMARK 4.6. We note that, to control the tail of the integral in (4.27) in $d \geq 3$ it suffices to invoke the diagonal heat-kernel estimate,

$$E_\mu P_\kappa^0(X_t = x) \leq \frac{c_1}{t^{d/2}}, \quad x \in \mathbb{Z}^d, \quad (4.30)$$

which in the elliptic case is an immediate consequence of the mixing theory for Markov chains based on isoperimetric inequalities. This is sufficient because the finiteness of the Green function in $d \geq 3$ permits us to define $(f, (-\mathcal{L}_\kappa)^{-1}f)$ even for $f \geq 0$. This enables us to reduce the general case to positive f by decomposing the test function into a positive and negative part and applying

$$E_\nu (f, e^{t\mathcal{L}_\kappa} f) \leq \|f\|_1^2 \frac{c_1}{t^{d/2}}, \quad (4.31)$$

which is uniformly integrable when $d \geq 3$. However, to include $d = 2$, we cannot disregard the cancellations due to the vanishing of $\int f(x)dx$ and thus the stronger derivative bound (3.23) is necessary. A similar situation occurred in Giacomini, Olla and Spohn [13] where a stronger Nash continuity estimate was required to include $d = 2$.

Finally, we are ready to establish the main result of this paper:

PROOF OF THEOREM 2.4. Let μ be a translation-invariant, ergodic gradient Gibbs measure with zero tilt and let $\tilde{\mu}$ be its extension to $\mathbb{R}^{\mathbb{B}(\mathbb{Z}^d)} \times \mathbb{R}^{\mathbb{B}(\mathbb{Z}^d)}$. We want to prove that $\phi(f_\epsilon)$ tends weakly to a normal random variable with mean zero and variance $(f, (-Q)^{-1}f)$. By Lemma 2.2 it suffices to prove this for $f \in C_0^\infty(\mathbb{R}^d) \cap \mathcal{H}$.

By Lemma 3.4 we know that $\phi(f)$ is Gaussian conditional on κ . The standard formula for any Gaussian random variable X ,

$$E(e^{iX}) = e^{iE(X) - \frac{1}{2}\text{Var}(X)}, \quad (4.32)$$

implies via (3.15–3.16)

$$E_{\tilde{\mu}}(e^{i\phi(f)} | \mathcal{F})(\kappa) = e^{-\frac{1}{2}(f, (-\mathcal{L}_\kappa)^{-1}f)}. \quad (4.33)$$

By Proposition 4.4 we have $(f_\epsilon, (-\mathcal{L}_\kappa)^{-1}f_\epsilon) \rightarrow (f, (-Q)^{-1}f)$ in $\tilde{\mu}$ -probability. Since the right-hand side of (4.33) is a bounded continuous function of this inner product, (2.15) follows by means of the Bounded Convergence Theorem. \square

5. Potential theory for random conductance model. The proof of the key Lemma 3.3 leads us to the study of potential theory for operators depending on a random environment that falls into the class of *random conductance models*. A good deal of what is to follow exists explicitly, or implicitly, in the literature. We borrowed some of the notation from the paper of Mathieu and Piatnitski [21] although the formalism draws on earlier works in homogenization theory; see, e.g., the book by Jikov, Kozlov and Oleinik [17]. Notwithstanding, the content of Sect. 5.2 is presumably new.

5.1. Basic notions. Consider a translation-invariant ν probability measure on $\Omega := \mathbb{R}_+^{\mathbb{B}(\mathbb{Z}^d)}$ (endowed with the product σ -algebra) satisfying the ellipticity condition

$$\exists \epsilon > 0: \quad \nu(\epsilon \leq \kappa_b \leq 1/\epsilon) = 1, \quad b \in \mathbb{B}(\mathbb{Z}^d). \quad (5.1)$$

Let $L^2(\nu)$ denote the closure of the set of all local functions in the topology induced by the inner product

$$\langle h, g \rangle := E_\nu(h(\kappa)g(\kappa)). \quad (5.2)$$

Let $B := \{\hat{e}_1, \dots, \hat{e}_d\}$ denote the set of coordinate vectors in \mathbb{Z}^d . The translation by the vectors in B induce natural unitary maps T_1, \dots, T_d on $L^2(\nu)$ defined via

$$(T_j h) := h \circ \tau_{\hat{e}_j}, \quad j = 1, \dots, d. \quad (5.3)$$

Apart from the square integrable functions we will also need to work with vector fields by which we will generally mean measurable functions $u: \Omega \times B \rightarrow \mathbb{R}$ or $\Omega \times B \rightarrow \mathbb{R}^d$ depending on the context. We will sometimes write u_1, \dots, u_d for $u(\cdot, \hat{e}_1), \dots, u(\cdot, \hat{e}_d)$ — note that these may still be vector valued.

REMARK 5.1. While we index vector fields only by the *positive* coordinate vectors, in certain situations it is convenient to have them defined also for the negative coordinate directions via

$$u(\kappa, -b) := -u(\tau_{-b}\kappa, b), \quad b \in B. \quad (5.4)$$

As we will see this definition will automatically ensure that the cycle condition (see Lemma 5.2 below) holds for the trivial cycles crossing only a single edge.

Let $L^2_{\text{vec}}(\nu)$ be the set of all vector fields with $(u, u) < \infty$, where (\cdot, \cdot) denotes the inner product

$$(u, v) := E_\nu\left(\sum_{b \in B} \kappa_b u(\kappa, b) \cdot v(\kappa, b)\right). \quad (5.5)$$

Examples of such functions are the gradients ∇h of local functions $h \in L^2(\nu)$ defined componentwise via the formula

$$(\nabla h)_j := T_j h - h, \quad j = 1, \dots, d. \quad (5.6)$$

We denote by $L^2_{\nabla}(\nu)$ the closure of the set of gradients of local functions in the topology induced by the above inner product.

LEMMA 5.2. *Let $u \in L^2_{\nabla}$. Then u satisfies the cycle condition*

$$\sum_{j=0}^n u(\tau_{x_j}\kappa, x_{j+1} - x_j) = 0 \quad (5.7)$$

for any finite (nearest-neighbor) cycle $(x_0, x_1, \dots, x_n = x_0)$ on \mathbb{Z}^d . In particular, there exists a shift-covariant function $\bar{u}: \Omega \times \mathbb{Z}^d \rightarrow \mathbb{R}^d$ such that $u(\kappa, b) = \bar{u}(\kappa, b)$ for every $b \in B$.

PROOF. The cycle condition (5.7) holds trivially for all gradients of local functions. Indeed, if $u = \nabla h$ then, in light of (5.4) we have

$$u(\tau_{x_j}\kappa, x_{j+1} - x_j) = (\nabla h)_{x_{j+1}-x_j}(\tau_{x_j}\kappa) = h \circ \tau_{x_{j+1}}(\kappa) - h \circ \tau_{x_j}(\kappa). \quad (5.8)$$

A corresponding limit extends this to all of L^2_{∇} . To define $\bar{u}(\cdot, x)$, we integrate properly shifted values of u along a path from zero to x ; the cycle condition guarantees that the result is independent of the choice of the path and that \bar{u} is shift covariant. \square

We will henceforth use the convention to write \bar{u} for the extension of a shift-covariant vector field $u \in L^2_{\text{vec}}$ to a function on \mathbb{Z}^d . Notice that the shift T_j extends naturally via

$$T_j \bar{u}(\kappa, x) := \bar{u}(\tau_{\hat{e}_j}\kappa, x) = \overline{(T_j u)}(\kappa, x). \quad (5.9)$$

Next let us characterize the functions in $(L^2_{\nabla})^{\perp}$:

LEMMA 5.3. For $u \in L^2_{\text{vec}}(\nu)$, let $\mathcal{L}u$ be the function in $L^2(\nu)$ defined by

$$(\mathcal{L}u)(\kappa) := \sum_{b \in B} [\kappa_b u(\kappa, b) - (\tau_{-b}\kappa)_b u(\tau_{-b}\kappa, b)], \quad (5.10)$$

where $-b$ is the coordinate vector opposite to b . Then we have

$$u \in (L^2_{\nabla})^{\perp} \iff \mathcal{L}u = 0, \quad \nu\text{-a.s.} \quad (5.11)$$

If u satisfies the cycle condition and \bar{u} is its extension, then $\mathcal{L}u(\tau_x\kappa) = \mathcal{L}_{\kappa}\bar{u}(\kappa, x)$.

PROOF. These are direct consequences of the definitions, the translation invariance of ν and a simple calculation. \square

Note that $\mathcal{L}u$ plays the role of *divergence* — i.e., the total flow out of a given vertex — of vector field u . However, we prefer to denote it by \mathcal{L} to emphasize its connection with the operator \mathcal{L}_{κ} .

5.2. *Uniqueness of harmonic embedding.* Clearly, all $u \in L^2_{\nabla}$ are shift covariant and have zero mean. A natural question is whether every shift-covariant zero-mean u is in L^2_{∇} . (Note that this is analogous to asking whether every closed differential form is exact.) Our answer to this is in the affirmative:

THEOREM 5.4. *Suppose ν is ergodic. Then every $u \in L^2_{\text{vec}}(\nu)$ which obeys the cycle condition (5.7) and $E_{\nu}u = 0$ satisfies $u \in L^2_{\nabla}$.*

Recall again that (5.7) and zero expectation are necessary for $u \in L^2_{\nabla}$. The above says that these conditions are also sufficient. To prove the theorem, we will need the following lemma:

LEMMA 5.5. *Let P_j denote the orthogonal projection onto $\text{Ker}(1 - T_j)$ in $L^2_{\text{vec}}(\nu)$. If ν is ergodic and $u \in L^2_{\text{vec}}(\nu)$ satisfies (5.7), then $P_j u = E_{\nu}(P_j u)$, ν -almost surely.*

PROOF. Fix $u \in L^2_{\text{vec}}(\nu)$ that obeys (5.7) and let \bar{u} be the corresponding shift-covariant function. We will prove the claim only for the component $u_1 = u(\cdot, \hat{e}_1)$; the other cases follow analogously. By translation covariance and the L^2 -Ergodic Theorem,

$$\frac{\bar{u}(\cdot, n\hat{e}_1)}{n} = \frac{1}{n} \sum_{k=0}^{n-1} T_1^k u_1 \xrightarrow{n \rightarrow \infty} P_1 u_1, \quad \text{in } L^2(\nu). \quad (5.12)$$

If ν were separately ergodic (i.e., ergodic with respect to T_1 alone) then the claim would immediately follow by the fact that every T_1 -invariant function must be constant. To make up for the potential lack of separate ergodicity, we note that translation covariance of u and the fact that \bar{u} obeys the cycle conditions yield

$$\begin{aligned} T_j \bar{u}(\cdot, n\hat{e}_1) &= \bar{u}(\cdot, n\hat{e}_1 + \hat{e}_j) - \bar{u}(\cdot, \hat{e}_j) \\ &= \bar{u}(\cdot, n\hat{e}_1) - u(\cdot, \hat{e}_j) + T_{n\hat{e}_1} u(\cdot, \hat{e}_j) \end{aligned} \quad (5.13)$$

It follows that $T_j \frac{1}{n} \bar{u}(\cdot, n\hat{e}_1)$ also converges to $P_1 u_1$ (in L^2) and so, by the continuity of T_j ,

$$T_j P_1 u_1 = P_1 u_1. \quad (5.14)$$

Hence $P_1 u$ is invariant under all shifts and so it is constant ν -a.s. \square

PROOF OF THEOREM 5.4. Suppose ν is ergodic and let $u \in L^2_{\text{vec}}(\nu)$ obey (5.7) and $E_{\nu}u = 0$. The boundedness of the κ_b 's away from zero and infinity

ensures that $u \in L^2_{\text{vec}}(\nu)$ iff all of its components are in $L^2(\nu)$. Our task is to construct functions $h_\epsilon \in L^2(\nu)$ such that $\nabla h_\epsilon \rightarrow u$ in $L^2_{\text{vec}}(\nu)$. We define

$$h_\epsilon := - \sum_{n \geq 0} \frac{T_1^n u_1}{(1 + \epsilon)^{n+1}} \quad (5.15)$$

and note that this is the unique solution of the equation $(1 + \epsilon - T_1)h_\epsilon = -u_1$. This observation implies

$$(1 - T_1)h_\epsilon = -u_1 - \epsilon h_\epsilon \quad (5.16)$$

and so the first component of ∇h_ϵ converges to that of u provided $\epsilon h_\epsilon \rightarrow 0$ in $L^2(\nu)$. To see what happens with the other components of ∇h_ϵ , we note that the cycle condition (5.7) translates into

$$(1 - T_j)u_1 = (1 - T_1)u_j. \quad (5.17)$$

Applying this to the definition of h_ϵ , we conclude

$$(1 - T_j)h_\epsilon = -u_j - \epsilon \tilde{h}_\epsilon, \quad (5.18)$$

where \tilde{h}_ϵ is defined as h_ϵ but with u_1 replaced by u_j . Again, it suffices to show that $\epsilon \tilde{h}_\epsilon \rightarrow 0$ in $L^2(\nu)$ which will boil down to the same argument as for $j = 1$.

To prove that $\epsilon h_\epsilon \rightarrow 0$ we note that, for the inner product in (5.2),

$$\langle h_\epsilon, h_\epsilon \rangle = \sum_{n \geq 0} \frac{n+1}{(1 + \epsilon)^{n+2}} \langle u_1, T_1^n u_1 \rangle. \quad (5.19)$$

Introducing the notation $A_n u$ for the average

$$A_n u := \frac{1}{n} \sum_{k=0}^{n-1} T_1^k u, \quad (5.20)$$

inserting this into the above sum and reordering the terms, we get

$$\langle h_\epsilon, h_\epsilon \rangle = \frac{\langle u_1, u_1 \rangle}{(1 + \epsilon)^2} + \sum_{n \geq 1} \frac{n(\epsilon n - 1)}{(1 + \epsilon)^{n+2}} \langle u_1, A_n u_1 \rangle. \quad (5.21)$$

By Lemma 5.5, the L^2 -Ergodic Theorem and the fact that u has zero expectation in ν ,

$$A_n u_1 \xrightarrow{n \rightarrow \infty} P_1 u_1 = E_\nu u_1 = 0, \quad \text{in } L^2(\nu), \quad (5.22)$$

and so $\langle u_1, A_n u_1 \rangle \rightarrow 0$ as $n \rightarrow \infty$. A straightforward estimate now shows that the sum in (5.21) is $o(\epsilon^{-2})$ and so $\epsilon^2 \langle h_\epsilon, h_\epsilon \rangle \rightarrow 0$ as desired. \square

Not every function in $L_{\text{vec}}^2(\nu)$ necessarily belongs to $L_{\nabla}^2(\nu)$. A prime example is the *position* vector field $x(\kappa, b) = b$. Indeed, let $\chi: \mathbb{R}^{\mathbb{B}(\mathbb{Z}^d)} \times B \rightarrow \mathbb{R}^d$ be the projection

$$\chi := -\text{proj}_{L_{\nabla}^2(\nu)} x. \quad (5.23)$$

Since $\chi \in L_{\nabla}^2(\nu)$ it satisfies (5.7) and we may extend it to a function $\bar{\chi}$ mapping $\mathbb{R}^{\mathbb{B}(\mathbb{Z}^d)} \times \mathbb{Z}^d \rightarrow \mathbb{R}^d$ by setting $\bar{\chi}(\cdot, 0) = 0$ and integrating the gradients along oriented paths. Lemma 5.3 implies that $\bar{x} + \bar{\chi}$ is harmonic, $\mathcal{L}_{\kappa}(\bar{x} + \bar{\chi}) = 0$. Moreover, Lemma 5.5 and (5.12) show that $\bar{\chi}(\cdot, n\hat{e}_j)/n \rightarrow 0$ and so $x + \chi \neq 0$. It follows that $x \notin L_{\nabla}^2$ and so $L_{\nabla}^2 \neq L_{\text{vec}}^2$.

The function $\bar{\chi}$ is generally referred to as the *corrector* because it corrects for the non-harmonicity of the position function. The corrector can be defined by appealing to spectral theory (Kipnis and Varadhan [18]; see also Berger and Biskup [4]); the above “projection” definition is inspired by those in Giacomini, Olla and Spohn [13] or Mathieu and Piatnitski [21]. As a side remark we note that the function $\bar{x} + \bar{\chi}$ actually allows us to characterize the space of all square-integrable shift covariant functions:

COROLLARY 5.6. *Suppose ν is ergodic. Then every shift-covariant \mathbb{R} -valued $u \in (L_{\nabla}^2)^{\perp}$ can be obtained from $x + \chi$, where x is the position function and $-\chi$ is its orthogonal projection onto L_{∇}^2 , by means of the linear transformation*

$$u_j(\kappa) = \sum_{k=1}^d [\hat{e}_k \cdot (x_j + \chi_j)] E_{\nu} u_k, \quad j = 1, \dots, d. \quad (5.24)$$

Here u_1, \dots, u_d stand for $u(\cdot, \hat{e}_1), \dots, u(\cdot, \hat{e}_d)$ and similarly for (\mathbb{R}^d -valued objects) x_1, \dots, x_d and χ_1, \dots, χ_d . In particular, for \mathbb{R} -valued vector fields we have

$$\{u \in L_{\text{vec}}^2(\nu) : \text{shift covariant}\} = L_{\nabla}^2 \oplus \{\lambda \cdot (x + \chi) : \lambda \in \mathbb{R}^d\}. \quad (5.25)$$

REMARK 5.7. The observation (5.25) — which in the language borrowed from differential geometry implies that linear transforms of $x + \chi$ are the only closed forms that are not exact — has previously been discovered in the context of multicolored exclusion processes (Quastel [24, Theorem 9.1]). However, while our proof is based only on soft, Poisson-equation based argument (5.15–5.22), that of [24] requires an explicit bound on the spectral gap for the corresponding dynamics. The reason may be the reliance of [24] on the *spatial* ergodic theorem which, naturally, leads to bounds involving the Poincaré inequality and/or spectral gap.

PROOF OF COROLLARY 5.6. This is a simple consequence of part (1) of Theorem 5.4: Let $w = (w_1, \dots, w_d)$ be the d -component vector field whose (\mathbb{R} -valued) components are defined by

$$w_j := \sum_{k=1}^d [\hat{e}_k \cdot (x_j + \chi_j)] E_v u_k, \quad j = 1, \dots, d. \quad (5.26)$$

We will show that $w = u$. First, both u and w obey the cycle condition and so does $u - w$. As $\chi \in L_{\nabla}^2(v)$ implies $E_v \chi = 0$ the fact $E_v x_k = x_k = \hat{e}_k$ shows $E_v w_j = E_v u_j$, i.e., $E_v(u - w) = 0$. Theorem 5.4(1) implies $u - w \in L_{\nabla}^2$. On the other hand, $\mathcal{L}(x + \chi) = 0$ implies $\mathcal{L}w = 0$ and so, by Lemma 5.3, $w \in (L_{\nabla}^2)^\perp$. Thus $u - w \in (L_{\nabla}^2)^\perp$ holds as well. It follows that $u = w$ as claimed.

As to (5.25), the argument we just used ensures that a shift-covariant field $u \in L_{\text{vec}}^2(v)$ can be written as $\lambda \cdot (x + \chi)$, where λ is a vector with components $\lambda_j := E_v u_j$, plus a shift-covariant vector field with zero expectation. By Theorem 5.4(1), the latter is in L_{∇}^2 . \square

We still owe the reader the proof of Lemma 3.3:

PROOF. Proof of Lemma 3.3 Since g has square-integrable components, we have $g \in L_{\text{vec}}^2(v)$. As g is shift covariant and has zero expectation, Theorem 5.4 implies that $g \in L_{\nabla}^2$. But g is also harmonic and so Lemma 5.3 in turn forces $g \in (L_{\nabla}^2)^\perp$. Thus $g = 0$ as desired. \square

Recall that $\bar{\chi}$ is the extension of χ subject to the condition $\bar{\chi}(\cdot, 0) = 0$. With this object at hand, we may even remove the restriction to zero slope in Lemma 3.4:

COROLLARY 5.8. *Let μ be a translation-invariant, ergodic gradient Gibbs measure for the potential (1.4) and let $\tilde{\mu}$ be its extension to $\mathbb{R}^{\mathbb{B}(\mathbb{Z}^d)} \times \mathbb{R}^{\mathbb{B}(\mathbb{Z}^d)}$. Let $t \in \mathbb{R}^d$ be the tilt of μ and let $\mathcal{F} := \sigma(\{\kappa_b : b \in \mathbb{B}(\mathbb{Z}^d)\})$. Then $\tilde{\mu}(-|\mathcal{F})(\kappa)$ is Gaussian with mean*

$$E_{\tilde{\mu}}(\phi_x - \phi_0 | \mathcal{F})(\kappa) = t \cdot [x + \bar{\chi}(\kappa, x)] \quad (5.27)$$

and the covariance given by $(-\mathcal{L}_\kappa)^{-1}$.

PROOF. Let \bar{u} denote the extension of the shift-covariant vector field $u(\kappa, b) := E_{\tilde{\mu}}(\eta_{0,b} | \mathcal{F})(\kappa)$, $b \in B$. Clearly, we have

$$\bar{u}(\kappa, x) = E_{\tilde{\mu}}(\phi_x - \phi_0 | \mathcal{F})(\kappa), \quad x \in \mathbb{Z}^d. \quad (5.28)$$

Inspecting the proof of Lemma 3.4, all formulas carry over except (3.22) which becomes

$$E_{\bar{\mu}}(\bar{u}(\cdot, x)) = t \cdot x, \quad x \in \mathbb{Z}^d. \quad (5.29)$$

Thus $\bar{u} - t \cdot (\bar{x} + \bar{\chi})$ is harmonic, shift-covariant, and of zero mean and so $\bar{u} = t \cdot (\bar{x} + \bar{\chi})$ by Lemma 5.3 and Theorem 5.4. \square

6. Discussion and open problems. The proofs in the present note often rely on the fact that the random walk with generator \mathcal{L}_κ is uniformly elliptic. This enters via the assumption that ϱ , defining the potential V , has support bounded away from zero. While we believe that the general picture carries over even if we let the support extend all the way to zero, a number of steps in the proof become quite subtle. For instance, the pointwise heat-kernel asymptotic for this walk may take a radically different form (Berger, Biskup, Hoffman and Kozma [3]) and it is not known under what conditions on the environment the walk scales to Brownian motion. Progress in this direction for i.i.d. (or i.i.d. dominated) environments has been made only recently (Mathieu [20], Biskup and Prescott [6], Barlow and Deuschel [2]).

Another interesting open question concerns the annoying restriction to *zero-tilt* gradient Gibbs measures. As is seen from Corollary 5.8, once the tilt is nonzero, (4.33) has to be modified into

$$E_{\bar{\mu}}(e^{i[\phi(f) - E_{\bar{\mu}}(\phi(f))]} | \mathcal{F})(\kappa) = e^{-i(t \cdot \bar{\chi})(f) - \frac{1}{2}(f, (-\mathcal{L}_\kappa)^{-1}f)}. \quad (6.1)$$

When we plug in f_ϵ for f , the second term in the exponent still converges to $(f, (-Q)^{-1}f)$ in probability as $\epsilon \downarrow 0$. To show convergence of $\phi_\epsilon - E_{\bar{\mu}}\phi_\epsilon$ to Gaussian free field, we thus have to show that $(t \cdot \bar{\chi})(f_\epsilon)$ converges in law to a normal random variable. If V is strictly convex and of the form (1.4), we in fact know this to be true by the main result of Giacomin, Olla and Spohn [13]. However, the general case seems to be a hard open problem (cf. [4, Conjecture 5]). Already in $d = 1$, where the corrector can be written explicitly as the sum

$$\bar{\chi}(\kappa, x) = \frac{1}{C} \sum_{n=0}^{x-1} \left(\frac{1}{\kappa_{n,n+1}} - C \right), \quad (6.2)$$

with $C := E_{\bar{\mu}}(1/\kappa_b)$, the desired result requires that the resistivities $1/\kappa_b$ exhibit a specific correlation decay [11, Theorem 7.6]. Very little information concerning this is known in $d \geq 2$; see Mourrat [22] for some recent progress in this direction.

Our final remark concerns a generalization to potentials V where (1.4) has been modified into

$$V(\eta) := -\log \int \varrho(d\kappa) e^{-W_\kappa(\eta)}, \quad (6.3)$$

with (W_κ) denoting a family of strictly convex, even, measurable functions with uniformly superlinear growth at $\pm\infty$ and a uniform lower bound. In this case we may still consider the extended gradient Gibbs measures; however, the conditional law given the κ 's is no longer Gaussian. Notwithstanding, given κ , the Helffer-Sjöstrand representation still applies and leads to a random walk in a dynamic random environment. Its annealed CLT would imply the Gaussian scaling limit for the ϕ -field.

Acknowledgments. The research of M.B. was supported by the NSF grant DMS-0505356. This project was started at the Kavli Institute for Theoretical Physics during the program “Stochastic Geometry and Field Theory: From Growth Phenomena to Disordered Systems.” The KITP program was supported by the NSF under the Grant No. PHY99-07949. We wish to thank Jean-Dominique Deuschel for a number of interesting suggestions, and to Scott Sheffield and Jeremy Quastel for discussions on related subjects. We are also grateful to anonymous referees for careful reading of the first version of this paper and for a number of constructive suggestions.

REFERENCES

- [1] S. Adams, R. Kotecký and S. Müller, in preparation.
- [2] M.T. Barlow and J.-D. Deuschel (2008). *Invariance principle for the random conductance model with unbounded conductances*. Ann. Probab. (to appear), www.math.ubc.ca/~barlow/preprints/index.html.
- [3] N. Berger, M. Biskup, C.E. Hoffman and G. Kozma (2008). *Anomalous heat-kernel decay for random walk among bounded random conductances*. Ann. Inst. Henri Poincaré **44**, no. 2, 374–392.
- [4] N. Berger and M. Biskup (2007). *Quenched invariance principle for simple random walk on percolation clusters*. Probab. Theory Rel. Fields **137**, no. 1-2, 83–120.
- [5] M. Biskup and R. Kotecký (2007). *Phase coexistence of gradient Gibbs states*. Probab. Theory Rel. Fields **139**, no. 1–2, 1–39.
- [6] M. Biskup and T. Prescott (2007). *Functional CLT for random walk among bounded random conductances*. Electron. J. Probab. **12**, Paper no. 49, 1323–1348.
- [7] D. Brydges and H.-T. Yau (1990). *Grad ϕ perturbations of massless Gaussian fields*. Commun. Math. Phys. **129**, no. 2, 351–392.
- [8] C. Cotar, J.-D. Deuschel and S. Müller (2009). *Strict convexity of the free energy for a class of non-convex gradient models*, Commun. Math. Phys. **286**, no. 1, 359–376.
- [9] P.G. Doyle and J.L. Snell (1984). *Random Walks and Electric Networks*. Carus Mathematical Monographs, vol. 22, Mathematical Association of America, Washington, DC.

- [10] T. Delmotte and J.-D. Deuschel (2005). *On estimating the derivatives of symmetric diffusions in stationary random environment, with applications to $\nabla\phi$ interface model*, Probab. Theory Rel. Fields **133**, no. 3, 358–390.
- [11] R. Durrett (2005). *Probability: Theory and Examples* (Third edition), Brooks/Cole — Thomson Learning, Belmont, CA.
- [12] K. Gawędzki and A. Kupiainen, *Block spin renormalization group for dipole gas and $(\nabla\phi)^4$* , Ann. Physics **147** (1983), no. 1, 198–243.
- [13] G. Giacomin, S. Olla and H. Spohn (2001). *Equilibrium fluctuations for $\nabla\phi$ interface model*. Ann. Prob. **29**, no. 3, 1138–1172.
- [14] L. Gross (1967). *Abstract Wiener spaces*. Proc. Fifth Berkeley Sympos. Math. Statist. and Probability (Berkeley, Calif., 1965/66), Vol. II: Contributions to Probability Theory, Part 1 pp. 31–42 Univ. California Press, Berkeley.
- [15] T. Funaki (2005). *Stochastic Interface Models*. In.: J. Picard (ed.), *Lectures on probability theory and statistics*, Lecture Notes in Mathematics, vol. 1869, Springer-Verlag, Berlin.
- [16] T. Funaki and H. Spohn (1997). *Motion by mean curvature from the Ginzburg-Landau $\nabla\phi$ interface model*. Commun. Math. Phys. **185**, no. 1, 1–36.
- [17] V.V. Jikov, S.M. Kozlov and O.A. Oleinik, *Homogenization of differential operators and integral functionals*, Springer-Verlag, Berlin, 1994, pp. 570.
- [18] C. Kipnis, and S.R.S Varadhan (1986). *A central limit theorem for additive functionals of reversible Markov processes and applications to simple exclusions*. Commun. Math. Phys. **104**, no. 1, 1–19.
- [19] J. Magnen and R. Sénéor, *The infrared behaviour of $(\nabla\Phi)^4_3$* , Ann. Physics **152** (1984), no. 1, 130–202.
- [20] P. Mathieu (2008). *Quenched invariance principles for random walks with random conductances*. J. Stat. Phys. **130**, no. 5, 1025–1046.
- [21] P. Mathieu and A.L. Piatnitski (2007). *Quenched invariance principles for random walks on percolation clusters*. Proc. Roy. Soc. A **463**, 2287–2307.
- [22] J.-C. Mourrat (2009). *Variance decay for functionals of the environment viewed from the particle*. arxiv.org/pdf/0902.0204v1
- [23] A. Naddaf and T. Spencer (1997). *On homogenization and scaling limit of some gradient perturbations of a massless free field*. Commun. Math. Phys. **183**, no. 1, 55–84.
- [24] J. Quastel (1992). *Diffusion of color in the simple exclusion process*. Commun. Pure Appl. Math. **45**, no. 6, 623–679.
- [25] S. Sheffield (2005). *Random Surfaces*. Astérisque No. 304, 175 pp.
- [26] S. Sheffield (2007). *Gaussian free field for mathematicians*. Probab. Theory Related Fields **139**, no. 3–4, 521–541.
- [27] V. Sidoravicius and A.-S. Sznitman (2004). *Quenched invariance principles for walks on clusters of percolation or among random conductances*. Probab. Theory Rel. Fields **129**, no. 2, 219–244.
- [28] Y. Velenik (2006). *Localization and delocalization of random interfaces*. Probab. Surveys **3**, 112–169.

DEPARTMENT OF MATHEMATICS
UNIVERSITY OF CALIFORNIA AT LOS ANGELES
LOS ANGELES
CALIFORNIA 90095-1555
U.S.A.
AND
SCHOOL OF ECONOMICS
UNIVERSITY OF SOUTH BOHEMIA
STUDENTSKÁ 13
CZ-37005 ČESKÉ BUDĚJOVICE
CZECH REPUBLIC
E-MAIL: biskup@math.ucla.edu
URL: <http://www.math.ucla.edu/~biskup/>

ZENTRUM MATHEMATIK
TECHNISCHE UNIVERSITÄT MÜNCHEN
BOLTZMANNSTRASSE 3
D-85748 GARCHING BEI MÜNCHEN
GERMANY
E-MAIL: spohn@ma.tum.de
URL: <http://www-m5.ma.tum.de/pers/spohn/>