

Schrödinger Flow near Harmonic Maps

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Abstract

For the Schrödinger flow from $\mathbb{R}^2 \times \mathbb{R}^+$ to the 2-sphere \mathbb{S}^2 , it is not known if finite energy solutions can blow up in finite time. We study equivariant solutions whose energy is near the energy of the family of equivariant harmonic maps. We prove that such solutions remain close to the harmonic maps until the blowup time (if any), and that they blow up if and only if the length scale of the nearest harmonic map goes to 0. © 2006 Wiley Periodicals, Inc.

1 Introduction and Main Result

The Schrödinger flow for maps from \mathbb{R}^n to \mathbb{S}^2 (also known as the Schrödinger map, and, in ferromagnetism, as the Heisenberg model or Landau-Lifshitz equation) is given by the equation

$$(1.1) \quad u_t = u \times \Delta u, \quad u(x, 0) = u_0(x).$$

Here $u(x, t)$ is the unknown map from $\mathbb{R}^n \times \mathbb{R}^+$ to the 2-sphere

$$\mathbb{S}^2 := \{u \in \mathbb{R}^3 \mid |u| = 1\} \subset \mathbb{R}^3,$$

Δ denotes the Laplace operator in \mathbb{R}^n , and \times denotes the cross product in \mathbb{R}^3 . A more geometric way to write this equation is

$$(1.2) \quad u_t = JP\Delta u, \quad P\Delta u = \Delta u + |\nabla u|^2 u,$$

where $P = P_u$ denotes the orthogonal projection from \mathbb{R}^3 onto the tangent plane

$$T_u\mathbb{S}^2 := \{\xi \in \mathbb{R}^3 \mid \xi \cdot u = 0\}$$

to \mathbb{S}^2 at u , and $J = J^u := u \times$ is a rotation through $\pi/2$ on $T_u\mathbb{S}^2$.

On one hand, equation (1.1) is a borderline case of the Landau-Lifshitz-Gilbert equations that model isotropic ferromagnetic spin systems:

$$(1.3) \quad u_t = aP\Delta u + bJP\Delta u, \quad a \geq 0,$$

(see, e.g., [10, 11]). The Schrödinger flow corresponds to the case $a = 0$. The case $b = 0$ is the well-studied harmonic map heat flow, for which some finite-energy solutions do blow up in finite time [2].

On the other hand, equation (1.1) is a particular case of the Schrödinger flow from a Riemannian manifold into another one with a complex structure (see, e.g., [6, 7, 8, 9, 14, 18, 19]). We will limit ourselves to the case $u : \mathbb{R}^n \times \mathbb{R}^+ \rightarrow \mathbb{S}^2$ in this paper.

Equation (1.1) can be written in the divergence form $u_t = \sum_{j=1}^n \partial_j (u \times \partial_j u)$, which is useful in the construction of global weak solutions [16]. Its formal equivalence to a nonlinear Schrödinger equation (NLS) can be seen by applying the stereographic projection from \mathbb{S}^2 to \mathbb{C}_∞ , the extended complex plane,

$$(1.4) \quad w = \frac{u_1 + iu_2}{1 + u_3}, \quad iw_t = -\Delta w + \frac{2\bar{w}}{1 + |w|^2} \sum_j (\partial_j w)^2.$$

It is also known to be equivalent to an integrable cubic NLS in space dimension $n = 1$ (see, e.g., [5, 17]).

Equation (1.1) formally conserves the *energy*

$$(1.5) \quad \mathcal{E}(u) = \frac{1}{2} \int_{\mathbb{R}^n} |\nabla u|^2 dx = \frac{1}{2} \int_{\mathbb{R}^n} \sum_{j=1}^n \sum_{k=1}^3 |\partial_j u_k|^2 dx.$$

The space dimension $n = 2$ is critical in the sense that $\mathcal{E}(u)$ is invariant under scaling. In general,

$$(1.6) \quad \mathcal{E}(u) = s^{2-n} \mathcal{E}(u_s), \quad u_s(x) := u\left(\frac{x}{s}\right), \quad s > 0.$$

For our problem $u : \mathbb{R}^n \times \mathbb{R}^+ \rightarrow \mathbb{S}^2$, local-in-time well-posedness (LWP) is established in [16] in the class $|u| = 1$ and $\nabla u \in H^k(\mathbb{R}^n)$, where $k > n/2 + 1$ is an integer. The authors there also proved global-in-time well-posedness (GWP) in the same class when $n = 1$ and when $n \geq 2$ for data that is small in certain Sobolev norms. For $n = 2$, global existence is proved in [3] for small energy radial or equivariant data. Also, for $n = 2$, LWP for a closely related system of nonlinear Schrödinger equations is established in [12, 13] for data corresponding to $\nabla u \in H^{1+\epsilon}$. There are known to be self-similar blowup solutions for $n = 2$ [4]; however, these do *not* have finite energy.

Fix $m \in \mathbb{Z}$, a nonzero integer. By an m -equivariant map $u : \mathbb{R}^2 \rightarrow \mathbb{S}^2$, we mean a map of the form

$$(1.7) \quad u(r, \theta) = e^{m\theta R} v(r)$$

where (r, θ) are polar coordinates on \mathbb{R}^2 , $v : [0, \infty) \rightarrow \mathbb{S}^2$, and R is the matrix generating rotations around the u_3 -axis:

$$(1.8) \quad R = \begin{bmatrix} 0 & -1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix}, \quad e^{\alpha R} = \begin{bmatrix} \cos \alpha & -\sin \alpha & 0 \\ \sin \alpha & \cos \alpha & 0 \\ 0 & 0 & 1 \end{bmatrix}.$$

Radial maps arise as the case $m = 0$. The class of m -equivariant maps is formally preserved by the Schrödinger flow.

If u is m -equivariant, we have $|\nabla u|^2 = |u_r|^2 + r^{-2}|u_\theta|^2 = |v_r|^2 + (m^2/r^2)|Rv|^2$, and so

$$(1.9) \quad \mathcal{E}(u) = \pi \int_0^\infty \left(|v_r|^2 + \frac{m^2}{r^2}(v_1^2 + v_2^2) \right) r \, dr.$$

If $\mathcal{E}(u) < \infty$, the limits $\lim_{r \rightarrow 0} v(r)$ and $\lim_{r \rightarrow \infty} v(r)$ make sense (see (2.2) in the next section), and so we must have $v(0), v(\infty) = \pm \hat{k}$, where $\hat{k} = (0, 0, 1)^\top$. We may and will fix $v(0) = -\hat{k}$. The two cases $v(\infty) = \pm \hat{k}$ correspond to different topological classes of maps. We denote by Σ_m the class of m -equivariant maps with $v(\infty) = \hat{k}$:

$$(1.10) \quad \Sigma_m = \{u : \mathbb{R}^2 \rightarrow \mathbb{S}^2 \mid u = e^{m\theta R} v(r), \mathcal{E}(u) < \infty, \\ v(0) = -\hat{k}, v(\infty) = \hat{k}\}.$$

The energy $\mathcal{E}(u)$ can be rewritten as follows:

$$(1.11) \quad \begin{aligned} \mathcal{E}(u) &= \pi \int_0^\infty \left(|v_r|^2 + \frac{m^2}{r^2} |J^v Rv|^2 \right) r \, dr \\ &= \pi \int_0^\infty \left| v_r - \frac{|m|}{r} J^v Rv \right|^2 r \, dr + \mathcal{E}_{\min} \end{aligned}$$

where $J^v := v \times$, and

$$(1.12) \quad \begin{aligned} \mathcal{E}_{\min} &= 2\pi \int_0^\infty v_r \cdot \frac{|m|}{r} J^v Rv \, r \, dr \\ &= 2\pi |m| \int_0^\infty (v_3)_r \, dr = 2\pi |m| [v_3(\infty) + 1] \end{aligned}$$

(using $v_1^2 + v_2^2 + v_3^2 = 1$). The number \mathcal{E}_{\min} , which depends only on the boundary conditions, is in fact 4π times the absolute value of the *degree* of the map u , considered as a map from \mathbb{S}^2 to itself (defined, for example, by integrating the pullback by u of the volume form on \mathbb{S}^2). It provides a lower bound for the energy of an m -equivariant map, $\mathcal{E}(u) \geq \mathcal{E}_{\min}$, and this lower bound is attained if and only if

$$(1.13) \quad v_r = \frac{|m|}{r} J^v Rv.$$

If $v(\infty) = -\hat{k}$, the minimal energy is $\mathcal{E}_{\min} = 0$ and is attained by the constant map $u \equiv -\hat{k}$. On the other hand, if $v(\infty) = \hat{k}$, the minimal energy is

$$\mathcal{E}(u) \geq \mathcal{E}_{\min} = 4\pi|m|$$

and is attained by the 2-parameter family of harmonic maps

$$(1.14) \quad \mathcal{O}_m := \{e^{m\theta R} h^{s,\alpha}(r) \mid s > 0, \alpha \in [0, 2\pi)\}$$

where

$$h^{s,\alpha}(r) := e^{\alpha R} h\left(\frac{r}{s}\right),$$

and

$$(1.15) \quad h(r) = \begin{pmatrix} h_1(r) \\ 0 \\ h_3(r) \end{pmatrix}, \quad h_1(r) = \frac{2}{r^{|m|} + r^{-|m|}}, \quad h_3(r) = \frac{r^{|m|} - r^{-|m|}}{r^{|m|} + r^{-|m|}}.$$

The fact that $h(r)$ satisfies (1.13) means

$$(1.16) \quad (h_1)_r = -\frac{m}{r} h_1 h_3, \quad (h_3)_r = \frac{m}{r} h_1^2.$$

So \mathcal{O}_m is the orbit of the single harmonic map $e^{m\theta R} h(r)$ under the symmetries of the energy \mathcal{E} that preserve equivariance: scaling and rotation. Explicitly,

$$(1.17) \quad e^{m\theta R} h^{s,\alpha}(r) = \begin{pmatrix} \cos(m\theta + \alpha)h_1(r/s) \\ \sin(m\theta + \alpha)h_1(r/s) \\ h_3(r/s) \end{pmatrix}.$$

The solution (1.15) is easily found by solving system (1.13) of ODEs directly. Alternately, under the stereographic projection (1.4), equation (1.13) amounts to the Cauchy-Riemann equations, and these harmonic maps correspond to the antimeromorphic (if $m > 0$) functions

$$(1.18) \quad w = e^{i\alpha} \left(\frac{\bar{z}}{s}\right)^{-m}, \quad z = r e^{i\theta}.$$

We are now ready to state our main result. We denote $\|u\|_{\dot{H}^k} = \|\nabla^k u\|_{L^2}$.

THEOREM 1.1 *There exist $\delta > 0$ and $C_0, C_1 > 0$ such that if $u \in C([0, T]; \dot{H}^2 \cap \Sigma_m)$ is a solution of the Schrödinger flow (1.1) conserving energy and satisfying*

$$\delta_1^2 := \mathcal{E}(u_0) - 4\pi|m| < \delta^2,$$

then there exist $s(t) \in \mathcal{C}([0, T]; (0, \infty))$ and $\alpha(t) \in \mathcal{C}([0, T]; \mathbb{R})$ so that

$$(1.19) \quad \left\| u(x, t) - e^{(m\theta + \alpha(t))R} h\left(\frac{r}{s(t)}\right) \right\|_{\dot{H}^1(\mathbb{R}^2)} < C_0 \delta_1 \quad \forall t \in [0, T].$$

Moreover, $s(t) > C_1/\|u(t)\|_{\dot{H}^2}$. Furthermore, if $T < \infty$ is the maximal time of existence for u in \dot{H}^2 (i.e., $\lim_{t \rightarrow T^-} \|u(t)\|_{\dot{H}^2} = \infty$), then

$$(1.20) \quad \liminf_{t \rightarrow T^-} s(t) = 0.$$

Remark 1.2.

(1) This theorem can be viewed, on one hand, as an *orbital stability* result for the family of harmonic maps (at least up to the possible blowup time), and on the other hand as a characterization of blowup for energy near \mathcal{E}_{\min} : solutions blow up if and only if the \dot{H}^1 -nearest harmonic map “collapses” (i.e., its length scale goes to 0).

(2) The assumption $\mathcal{E}(u_0) - 4\pi|m| < \delta^2$ implies u_0 is close to \mathcal{O}_m in \dot{H}^1 (see (2.7)), but not necessarily in \dot{H}^2 .

(3) The existence of local (in time) $\dot{H}^2 \cap \dot{H}^1$ solutions of (1.1) is established in [16] for sufficiently regular initial data. In particular, this ensures Theorem 1.1 is nonempty (see also [12]–[13] for local well-posedness results). Local well-posedness with data in $\dot{H}^2 \cap \dot{H}^1$ still appears to be open. If we had this, (1.20) could be replaced by $\lim_{t \rightarrow T^-} s(t) = 0$.

(4) From now on we will assume $m > 0$. The cases $m < 0$ of Theorem 1.1 follow from the change of variables $(x_1, x_2, x_3) \rightarrow (x_1, -x_2, x_3)$.

The plan for the paper is as follows: In Section 2 we study maps whose energy is close to that of the family of harmonic maps. The analysis here is completely time independent. In Section 3 we apply Strichartz estimates to a certain nonlinear Schrödinger equation, obtained via the Hasimoto transformation introduced in [3], and present the proof of the main theorem. The proofs of some of the more technical lemmas are relegated to Section 4 in order to streamline the presentation. Without loss of generality, we assume $m > 0$ for the rest of the paper.

REMARK ON NOTATION. Throughout the paper, the letter C is used to denote a generic constant, the value of which may change from line to line.

2 Maps with Energy near the Harmonic Map Energy

This section is devoted entirely to static m -equivariant maps (i.e., there is no time dependence anywhere in this section). We establish some properties of maps with energy close to the harmonic map energy $4\pi m$. Roughly speaking, we prove that such maps are \dot{H}^1 -close to harmonic maps. Precise statements appear in Theorem 2.1 below.

We define the distance from any map u to the family \mathcal{O}_m of m -equivariant harmonic maps to be

$$\text{dist}(u, \mathcal{O}_m) := \inf_{\substack{s \in (0, \infty) \\ \alpha \in \mathbb{S}^1}} \|u - e^{m\theta R} h^{s, \alpha}\|_{\dot{H}^1}.$$

Here $\mathbb{S}^1 = \mathbb{R}/2\pi$.

The following theorem defines a (nonlinear) projection from the set Σ_m of m -equivariant maps with energy close to $4\pi m$ onto the family \mathcal{O}_m and establishes a key fact: for maps in this set, the squared distance $\text{dist}^2(u, \mathcal{O}_m)$ is bounded by the energy difference $\mathcal{E}(u) - 4\pi m$.

THEOREM 2.1 *There are constants $\delta > 0$ and $C_0, C_1 > 0$ such that if $u \in \Sigma_m$ satisfies*

$$\mathcal{E}(u) < 4\pi m + \delta^2,$$

then the following hold:

- (i) *There exist unique $s(u) \in (0, \infty)$ and $\alpha(u) \in \mathbb{S}^1$ such that*

$$\text{dist}(u, \mathcal{O}_m) = \|u - e^{m\theta R} h^{s(u), \alpha(u)}\|_{\dot{H}^1}.$$

Moreover, $s(u)$ and $\alpha(u)$ are continuous functions of $u \in \dot{H}^1$.

- (ii) $\text{dist}(u, \mathcal{O}_m) < C_0[\mathcal{E}(u) - 4\pi m]^{1/2} < C_0\delta$.
- (iii) *If $u \in \dot{H}^2(\mathbb{R}^2)$, then $s(u)\|u\|_{\dot{H}^2(\mathbb{R}^2)} > C_1$.*

PROOF: The proof is long, so we break it into a series of steps. At each step, we may need to take δ smaller than in the previous step.

Step 1. A Change of Variable. Recall that $u \in \Sigma_m$ implies that in polar coordinates (r, θ) , $u(x) = e^{m\theta R} v(r)$, with $v(0) = -\hat{k}$ and $v(\infty) = \hat{k}$. The change of variables

$$r \rightarrow y = m \log(r) \in (-\infty, \infty) \quad \text{or} \quad e^y = r^m$$

turns out to be very useful for our purposes.

Set

$$\tilde{v}(y) := v(e^{y/m}).$$

With this change of variables, the \dot{H}^1 inner product of m -equivariant maps changes as follows:

$$\begin{aligned} \langle e^{m\theta R} v(r), e^{m\theta R} w(r) \rangle_{\dot{H}^1} &= \int_{\mathbb{R}^2} \nabla[e^{m\theta R} v(r)] \cdot \nabla[e^{m\theta R} w(r)] dx \\ &= 2\pi \int_0^\infty \left(v_r \cdot w_r + \frac{m^2}{r^2} Rv \cdot R w \right) r dr \\ &= 2\pi m \int_{\mathbb{R}} (\tilde{v}' \cdot \tilde{w}' + R\tilde{v} \cdot R\tilde{w}) dy, \end{aligned}$$

where $'$ denotes d/dy . In particular,

$$(2.1) \quad \mathcal{E}(u) = 2\pi m \tilde{E}(\tilde{v}), \quad \tilde{E}(\tilde{v}) := \frac{1}{2} \int_{\mathbb{R}} (|\tilde{v}'(y)|^2 + \tilde{v}_1(y)^2 + \tilde{v}_2(y)^2) dy.$$

Note that this implies $\tilde{v}_j \in H^1(\mathbb{R})$ for $j = 1, 2$, and in particular $(\tilde{v}_j^2)' \in L^1(\mathbb{R})$ so that the limits $\lim_{y \rightarrow \pm\infty} \tilde{v}_j^2(y)$ exist and are equal 0. The function \tilde{v} is continuous, and $\tilde{v}_3^2 = 1 - \tilde{v}_1^2 - \tilde{v}_2^2$ has limit 1 as $y \rightarrow \pm\infty$. Thus the limits $\lim_{y \rightarrow \pm\infty} \tilde{v}$ exist, justifying our earlier claim

$$(2.2) \quad \lim_{r \rightarrow 0} v(r) \quad \text{and} \quad \lim_{r \rightarrow \infty} v(r) \quad \text{exist.}$$

Recall that for $v \in \Sigma_m$, we have chosen $\tilde{v}(-\infty) = -\hat{k}$, $\tilde{v}(\infty) = \hat{k}$.

$\tilde{E}(\tilde{v})$ inherits the “topological lower bound”

$$(2.3) \quad \tilde{E}(\tilde{v}) = 2 + \frac{1}{2} \int_{\mathbb{R}} |\tilde{v}' - J^{\tilde{v}} R \tilde{v}|^2 dy \geq 2.$$

In this new variable, scaling r corresponds to translating y . In particular, the family of harmonic maps is composed of translations and rotations of a simple explicit map:

$$h^{s,\alpha}(r) = e^{\alpha R} h\left(\frac{r}{s}\right) = e^{\alpha R} \tilde{h}(y - m \log(s)),$$

$$\tilde{h}(y) = \begin{pmatrix} \operatorname{sech} y \\ 0 \\ \tanh y \end{pmatrix}.$$

Step 2. Energy Close to $4\pi m$ Implies u Close to a Harmonic Map. Using the variable y , we would like to prove the following:

LEMMA 2.2 *For any $\epsilon > 0$, there exists $\mu > 0$ such that if a map $\tilde{v} : \mathbb{R} \rightarrow \mathbb{S}^2$ satisfies $\tilde{E}(\tilde{v}) < 2 + \mu$, and $\tilde{v}(-\infty) = -\hat{k}$, $\tilde{v}(\infty) = \hat{k}$, then*

$$\inf_{\alpha \in \mathbb{S}^1, a \in \mathbb{R}} \|\tilde{v} - e^{\alpha R} \tilde{h}(\cdot - a)\|_{H^1(\mathbb{R})} < \epsilon.$$

PROOF: Suppose not. Then there exist $v_j(r)$, $j = 1, 2, 3, \dots$, and $\epsilon_0 > 0$ such that

$$(2.4) \quad |v_j(r)| \equiv 1, \quad \tilde{E}(v_j) < 2 + \frac{1}{j},$$

$$v_j(\pm\infty) = \pm\hat{k}, \quad \|v_j - e^{\alpha R} \tilde{h}(\cdot - a)\|_{H^1(\mathbb{R})} \geq \epsilon_0,$$

for every $\alpha \in \mathbb{S}^1$ and $a \in \mathbb{R}$. Since $\tilde{E}(v_j) < \infty$, v_j is continuous and thus $v_{j3}(a_j) = 0$ for some $a_j \in \mathbb{R}$. We replace v_j by $w_j(y) := e^{-\alpha_j R} v_j(y + a_j)$, where $\alpha_j \in \mathbb{S}^1$ is chosen so that $w_j(0) = \hat{i} = (1, 0, 0)^T$. The properties (2.4) still hold for $\{w_j\}$:

$$(2.5) \quad |w_j(r)| \equiv 1, \quad \tilde{E}(w_j) < 2 + \frac{1}{j}, \quad w_j(\pm\infty) = \pm\hat{k}, \quad \|w_j - \tilde{h}\|_{H^1(\mathbb{R})} \geq \epsilon_0.$$

Since $\sup_j \tilde{E}(w_j) < \infty$, there is a subsequence (which we continue to denote by $\{w_j\}$) and a limit vector function $w^*(y)$ satisfying

$$(2.6) \quad \begin{cases} w_{j1} \longrightarrow w_1^*, & w_{j2} \longrightarrow w_2^* & \text{weakly in } H^1(\mathbb{R}) \\ w'_{j3} \longrightarrow w_3^* & & \text{weakly in } L^2(\mathbb{R}) \\ w_j \longrightarrow w^* & & \text{strongly in } L^2_{\text{loc}}(\mathbb{R}) \cap C^0_{\text{loc}}(\mathbb{R}). \end{cases}$$

Because of the local uniform convergence, we have $|w^*| \equiv 1$ and $w^*(0) = \hat{i}$. On the other hand, by the topological lower bound (2.3), we have

$$2 \leq \tilde{E}(w_j) = 2 + \frac{1}{2} \int_{\mathbb{R}} |w'_j - J^{w_j} R w_j|^2 < 2 + \frac{1}{j},$$

from which it is immediate that $\|w'_j - J^{w_j} R w_j\|_{L^2(\mathbb{R})} \rightarrow 0$ as $j \rightarrow \infty$. For any bounded interval I , using $J^w R w = \hat{k} - w_3 w$, we have

$$\|J^{w_j} R w_j - J^{w^*} R w^*\|_{L^2(I)} = \|w_{j3} w_j - w_3^* w^*\|_{L^2(I)} \rightarrow 0 \quad \text{as } j \rightarrow \infty.$$

Hence $w'_j \rightarrow J^{w^*} R w^*$ strongly in $L^2(I)$. On the other hand, $w^{*'}$ is the weak limit of w'_j , and so we obtain $w^{*'} = J^{w^*} R w^*$ almost everywhere, with $w^*(0) = \hat{i}$. By uniqueness of H^1_{loc} solutions of this system of ordinary differential equations, we must have

$$w^*(y) = \tilde{h}(y) = (\text{sech } y, 0, \tanh y)^\top.$$

Now we note that

$$\tilde{E}(w^*) = \tilde{E}(\tilde{h}) = 2 = \lim_{j \rightarrow \infty} \tilde{E}(w_j).$$

By (2.1),

$$\|R w^*\|_{H^1(\mathbb{R})}^2 + \|w_3^{*'}\|_{L^2(\mathbb{R})}^2 = \lim_{j \rightarrow \infty} (\|R w_j\|_{H^1(\mathbb{R})}^2 + \|w'_{j3}\|_{L^2(\mathbb{R})}^2).$$

By weak lower semicontinuity, i.e.,

$$\|R w^*\|_{H^1(\mathbb{R})}^2 \leq \liminf_{j \rightarrow \infty} \|R w_j\|_{H^1(\mathbb{R})}^2, \quad \|w_3^{*'}\|_{L^2(\mathbb{R})}^2 \leq \liminf_{j \rightarrow \infty} \|w'_{j3}\|_{L^2(\mathbb{R})}^2,$$

we have $\|R w_j\|_{H^1(\mathbb{R})}^2 \rightarrow \|R w^*\|_{H^1(\mathbb{R})}^2$ and $\|w'_{j3}\|_{L^2(\mathbb{R})}^2 \rightarrow \|w_3^{*'}\|_{L^2(\mathbb{R})}^2$, which implies $R w_j \rightarrow R w^*$ strongly in $H^1(\mathbb{R})$ and $w'_{j3} \rightarrow w_3^{*'}$ strongly in $L^2(\mathbb{R})$. Finally, we will show that $w_j - \tilde{h}$ converges to 0 strongly in $H^1(\mathbb{R})$, which will contradict assumption (2.5), and so complete the proof of the lemma. Indeed,

$$\|w_j - \tilde{h}\|_{H^1(\mathbb{R})}^2 = \|w'_j - \tilde{h}'\|_{L^2(\mathbb{R})}^2 + \|R w_j - R \tilde{h}\|_{L^2(\mathbb{R})}^2 + \|w_{j3} - \tilde{h}_3\|_{L^2(\mathbb{R})}^2.$$

We have already shown that the first two terms go to 0 in the limit, and so it remains to consider the last term. For this, we need another lemma. For $f : (a, b) \rightarrow \mathbb{R}$, denote by $T_{(a,b)}(f)$ the total variation of f on (a, b) . The following lemma shows that the total variation of v_3 is close to 2 if $\mathcal{E}(u)$ is close to $4\pi m$.

LEMMA 2.3 *If $u = e^{m\theta R} v(r) \in \Sigma_m$ and $\mathcal{E}(u) = 4\pi m + \epsilon_0$, then $T_{(0,\infty)}(v_3) \leq 2 + C\epsilon_0$.*

PROOF: Make the change of variable $\tilde{v}(y) = v(e^{y/m})$, and write

$$\tilde{v} = (\tilde{\rho} \cos(\tilde{\omega}), \tilde{\rho} \sin(\tilde{\omega}), \tilde{v}_3),$$

so that $\tilde{\rho}^2 = \tilde{v}_1^2 + \tilde{v}_2^2$. We have

$$\begin{aligned}
 4 + \frac{\epsilon_0}{\pi m} &\geq \int_{\mathbb{R}} (|\tilde{v}'|^2 + |R\tilde{v}|^2) dy = \int_{\mathbb{R}} \left(\tilde{\rho}^2 (\tilde{\omega}')^2 + \frac{(\tilde{\rho}')^2}{1 - \tilde{\rho}^2} + \tilde{\rho}^2 \right) dy \\
 &\geq 2 \int_{\mathbb{R}} \frac{\tilde{\rho} |\tilde{\rho}'|}{(1 - \tilde{\rho}^2)^{1/2}} dy = 2 \int_{\mathbb{R}} \left| \frac{d}{dy} (1 - \tilde{\rho}^2)^{1/2} \right| dy = 2T_{\mathbb{R}}(\tilde{v}_3).
 \end{aligned}$$

Dividing by 2 on both sides completes the proof. □

Applying this lemma to w_j , we have $T_{\mathbb{R}}(w_{j3}) \leq 2 + C/j$. Since $w_{j3}(-\infty) = -1$, $w_{j3}(0) = 0$, and $w_{j3}(\infty) = 1$, we have $w_{j3}(y) > -C/j$ for $y \geq 0$ (and similarly for $y \leq 0$). Fix $\epsilon_1 > 0$. For $|y| \geq \epsilon_1$ and j sufficiently large (depending on ϵ_1),

$$\begin{aligned}
 |w_{3j} - \tilde{h}_3| &= \frac{|w_{3j}^2 - \tilde{h}_3^2|}{|w_{3j} + \tilde{h}_3|} \leq \frac{|w_{3j}^2 - \tilde{h}_3^2|}{|\tilde{h}_3| - C/j} \\
 &\leq \frac{2}{\tanh \epsilon_1} |w_{3j}^2 - \tilde{h}_3^2| = \frac{2}{\tanh \epsilon_1} ||R\tilde{h}|^2 - |Rw_j|^2|,
 \end{aligned}$$

and so

$$\int_{|y| \geq \epsilon_1} |w_{3j} - \tilde{h}_3|^2 dy \rightarrow 0 \quad \text{as } j \rightarrow \infty.$$

Since $\int_{|y| < \epsilon_1} |w_{3j} - \tilde{h}_3|^2 dy \leq C\epsilon_1$ and ϵ_1 is arbitrary, we conclude

$$\|w_{3j} - \tilde{h}_3\|_{L^2(\mathbb{R})} \rightarrow 0.$$

This completes the proof of Lemma 2.2. □

Translating Lemma 2.2 back to the original variable $r = e^{y/m}$, we find

(2.7) given $\epsilon > 0$, there is $\mu > 0$ such that if $u \in \Sigma_m$, $\mathcal{E}(u) < 4\pi m + \mu$, then $\text{dist}(u, \mathcal{O}_m) < \epsilon$.

Step 3. Existence of $s(u)$ and $\alpha(u)$. Recall that since $u \in \Sigma_m$, $\mathcal{E}(u) \geq 4\pi m$, and set $\delta_1 := [\mathcal{E}(u) - 4\pi m]^{1/2} < \delta$. We observe first that

(2.8) $\liminf_{s \rightarrow \infty} \inf_{\alpha} \|u - e^{m\theta R} h^{s,\alpha}\|_{\dot{H}^1}^2 = \liminf_{s \rightarrow 0} \inf_{\alpha} \|u - e^{m\theta R} h^{s,\alpha}\|_{\dot{H}^1}^2 = 8\pi m + \delta_1^2$.

Indeed, we have

(2.9) $\|u - e^{m\theta R} h^{s,\alpha}\|_{\dot{H}^1}^2 = 8\pi m + \delta_1^2 - 2 \int_{\mathbb{R}^2} \nabla u \cdot \nabla (e^{m\theta R} h^{s,\alpha}) dx,$

and it suffices to show that for any $\alpha \in \mathbb{S}^1$

(2.10) $\int_{\mathbb{R}^2} \nabla u \cdot \nabla (e^{m\theta R} h^{s,\alpha}) dx \rightarrow 0$

as $s \rightarrow 0$ or $s \rightarrow \infty$. Since $h^{s,\alpha}(r) = h^{1,\alpha}(r/s)$, this latter fact follows from an easy lemma:

LEMMA 2.4 *If $f \in L^2(\mathbb{R}^2)$, then*

$$\frac{1}{s}f\left(\frac{x}{s}\right) \rightarrow 0 \quad \text{weakly in } L^2(\mathbb{R}^2) \text{ as } s \rightarrow \infty \text{ and } s \rightarrow 0.$$

PROOF: First suppose $f \in L^2 \cap L^\infty$ and fix $g \in L^2$. By Hölder's inequality,

$$\begin{aligned} \int_{\mathbb{R}^2} \frac{1}{s}f\left(\frac{x}{s}\right)g(x)dx &= \int_{|x| \leq \sqrt{s}} \frac{1}{s}f\left(\frac{x}{s}\right)g(x)dx + \int_{|x| > \sqrt{s}} \frac{1}{s}f\left(\frac{x}{s}\right)g(x)dx \\ &\leq \|f\|_{L^\infty} \|g\|_{L^2} \frac{\sqrt{\pi}}{\sqrt{s}} + \|f\|_{L^2} \left(\int_{|x| > \sqrt{s}} |g(x)|^2 dx \right)^{1/2} \rightarrow 0 \end{aligned}$$

as $s \rightarrow \infty$. A similar argument covers the $s \rightarrow 0$ case. For general $f \in L^2$, choose $f_\epsilon \in L^2 \cap L^\infty$ with $\|f_\epsilon - f\|_{L^2} < \epsilon$. Then by the above argument,

$$\left| \int_{\mathbb{R}^2} \frac{1}{s}f\left(\frac{x}{s}\right)g(x)dx \right| \leq \epsilon \|g\|_{L^2} + o(s).$$

Since ϵ is arbitrary, we are done. \square

To prove the claim (2.10), just take $g = \nabla u$ and $\frac{1}{s}f(\cdot/s) = \nabla(e^{m\theta R}h^{s,\alpha})$ in this lemma.

So on one hand, $\inf_\alpha \|u - e^{m\theta R}h^{s,\alpha}\|_{\dot{H}^1}^2$ approaches $8\pi m + \delta_1^2$ as $s \rightarrow 0$ and $s \rightarrow \infty$. On the other hand, (2.7) shows that if δ (and hence δ_1) is sufficiently small, then for some s and α , $\|u - e^{m\theta R}h^{s,\alpha}\|_{\dot{H}^1}^2 < 8\pi m$. Thus to minimize

$$F(s, \alpha) := \|u - e^{m\theta R}h^{s,\alpha}\|_{\dot{H}^1}^2$$

over $s \in (0, \infty)$, $\alpha \in \mathbb{S}^1$, it suffices to consider s in a compact subset of $(0, \infty)$. Since $F(s, \alpha)$ is continuous, there must exist $s(u) \in (0, \infty)$ and $\alpha(u) \in \mathbb{S}^1$ such that

$$\text{dist}(u, \mathcal{O}_m) = \|u - e^{m\theta R}h^{s(u), \alpha(u)}\|_{\dot{H}^1}.$$

Step 4. Uniqueness of $s(u)$ and $\alpha(u)$. Denote $\sigma = (s, \alpha)$. Suppose there exist σ_1 and σ_2 with $\sigma_1 \neq \sigma_2$ such that

$$\delta_0 := \text{dist}(u, \mathcal{O}_m) = \|u - e^{m\theta R}h^{\sigma_1}\|_{\dot{H}^1} = \|u - e^{m\theta R}h^{\sigma_2}\|_{\dot{H}^1}.$$

Let μ be half the distance between $e^{m\theta R}h^{\sigma_1}$ and $e^{m\theta R}h^{\sigma_2}$, $\mu := \frac{1}{2}\|e^{m\theta R}h^{\sigma_1} - e^{m\theta R}h^{\sigma_2}\|_{\dot{H}^1}$. It follows that $\mu \leq \delta_0$. Now set $\varphi(t) = e^{m\theta R}h^{\sigma(t)}$ with $\sigma(t) = \frac{1}{2}[(\sigma_1 + \sigma_2) + t(\sigma_2 - \sigma_1)]$, so that $\varphi(-1) = e^{m\theta R}h^{\sigma_1}$ and $\varphi(1) = e^{m\theta R}h^{\sigma_2}$. Set $\bar{\varphi} := \frac{1}{2}[\varphi(-1) + \varphi(1)]$. Lemma 4.6 (stated and proved in Section 4) then yields

$$\|\bar{\varphi} - \varphi(0)\|_{\dot{H}^1} < C\mu^2.$$

This estimate amounts to a bound on the curvature of the family \mathcal{O}_m .

Now $\nabla[u - \bar{\varphi}]$ is L^2 -orthogonal to $\nabla[e^{m\theta R}(h^{\sigma_1} - h^{\sigma_2})]$, since $\|u - e^{m\theta R}h^{\sigma_1}\|_{\dot{H}^1} = \|u - e^{m\theta R}h^{\sigma_2}\|_{\dot{H}^1}$. By the Pythagorean theorem,

$$\|u - \bar{\varphi}\|_{\dot{H}^1}^2 = \|u - \varphi(-1)\|_{\dot{H}^1}^2 - \frac{1}{4} \|\varphi(-1) - \varphi(1)\|_{\dot{H}^1}^2 = \delta_0^2 - \mu^2,$$

and so

$$\begin{aligned} \|u - \varphi(0)\|_{\dot{H}^1} &\leq \|u - \bar{\varphi}\|_{\dot{H}^1} + \|\bar{\varphi} - \varphi(0)\|_{\dot{H}^1} \leq (\delta_0^2 - \mu^2)^{1/2} + C\mu^2 \\ &= \delta_0 - \frac{\mu^2}{\delta_0 + \sqrt{\delta_0^2 - \mu^2}} + C\mu^2. \end{aligned}$$

By (2.7), we can ensure $\delta_0 < 1/(2C)$ by choosing δ sufficiently small. This in turn implies $\|u - \varphi(0)\|_{\dot{H}^1} < \delta_0$, which contradicts the assumption that δ_0 is the minimal distance. This establishes uniqueness of $s(u)$ and $\alpha(u)$.

Step 5. Continuity of $s(u)$ and $\alpha(u)$. We could invoke the implicit function theorem, but we prefer to give a simple direct proof of continuity. Suppose $u_j \rightarrow u$ in \dot{H}^1 with $\mathcal{E}(u_j) < 4\pi m + \delta^2$ and $\mathcal{E}(u) < 4\pi m + \delta^2$. We have

$$\text{dist}(u_j, \mathcal{O}_m) \leq \|u_j - e^{m\theta R}h^{s(u_j), \alpha(u_j)}\|_{\dot{H}^1} \leq \|u_j - u\|_{\dot{H}^1} + \text{dist}(u, \mathcal{O}_m)$$

and

$$\text{dist}(u, \mathcal{O}_m) \leq \|u - e^{m\theta R}h^{s(u_j), \alpha(u_j)}\|_{\dot{H}^1} \leq \|u_j - u\|_{\dot{H}^1} + \text{dist}(u_j, \mathcal{O}_m)$$

and so $\text{dist}(u_j, \mathcal{O}_m) \rightarrow \text{dist}(u, \mathcal{O}_m)$. Since

$$\|e^{m\theta R}(h^{s(u_j), \alpha(u_j)} - h^{s(u), \alpha(u)})\|_{\dot{H}^1} \leq \text{dist}(u_j, \mathcal{O}_m) + \text{dist}(u, \mathcal{O}_m) + \|u_j - u\|_{\dot{H}^1},$$

$\{s(u_j)\}$ is contained in a compact subinterval of $(0, \infty)$ by Lemma 4.5, and so, up to subsequence, $s(u_j) \rightarrow s_*$ and $\alpha(u_j) \rightarrow \alpha_*$ for some s_* and α_* . Along this subsequence

$$\begin{aligned} \|u - h^{s(u), \alpha(u)}\|_{\dot{H}^1} &= \text{dist}(u, \mathcal{O}_m) = \lim_{j \rightarrow \infty} \text{dist}(u_j, \mathcal{O}_m) \\ &= \lim_{j \rightarrow \infty} \|u_j - h^{s(u_j), \alpha(u_j)}\|_{\dot{H}^1} = \|u - h^{s_*, \alpha_*}\|_{\dot{H}^1}. \end{aligned}$$

By the uniqueness we have already proved, $s_* = s(u)$ and $\alpha_* = \alpha(u)$. We conclude that $s(u_j) \rightarrow s(u)$ and $\alpha(u_j) \rightarrow \alpha(u)$ (for the full sequence). Continuity is proved. This completes the proof of part (i) of Theorem 2.1.

Step 6. The “Linearized Operator.” We now proceed to the proof of part (ii) of Theorem 2.1. The main idea is this: the “global” result of Lemma 2.2 allows us now to work “locally” — i.e., nearby a harmonic map. Indeed, to prove (ii), we study the second variation of the energy functional around the nearby harmonic map. We begin by discussing this “linearized operator.”

Given an m -equivariant map $u \in \Sigma_m$ with $\mathcal{E}(u) - 4\pi m < \delta^2$, we fix $s = s(u)$ and $\alpha = \alpha(u)$ and write $u = e^{m\theta R}v(r)$ with

$$(2.11) \quad v(r) = e^{\alpha R} \left(h \left(\frac{r}{s} \right) + \xi \left(\frac{r}{s} \right) \right).$$

This defines $\xi = \xi(r) = e^{-\alpha R} v(sr) - h(r)$ (note that the variable r here is no longer the original polar coordinate). Using $h_r - (m/r)J^h Rh = 0$, expand

$$\begin{aligned} \mathcal{E}(u) &= \frac{1}{2} \int_{\mathbb{R}^2} |\nabla u|^2 dx = 4\pi m + \frac{1}{2} \int_{\mathbb{R}^2} \left| v_r - \frac{m}{r} J^v Rv \right|^2 dx \\ (2.12) \qquad &= 4\pi m + \frac{1}{2} \int_{\mathbb{R}^2} \left| L\xi + \frac{m}{r} \xi_3 \xi \right|^2 dx \end{aligned}$$

where $L\xi$ is the linear part,

$$(2.13) \qquad L\xi := \xi_r + \frac{m}{r}(\xi_3 h + h_3 \xi).$$

In fact, the operator L maps tangent vector fields (i.e., vector functions $\eta(r)$ tangent to \mathbb{S}^2 at $h(r)$) into tangent vector fields. To see this explicitly, we specify an orthonormal basis of $T_h \mathbb{S}^2$:

$$e = \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix} \quad \text{and} \quad J^h e = \begin{pmatrix} -h_3(r) \\ 0 \\ h_1(r) \end{pmatrix}.$$

Then for any map $\xi : [0, \infty) \rightarrow \mathbb{R}^3$, we have an orthogonal decomposition

$$(2.14) \qquad \xi(r) = z_1(r)e + z_2(r)J^h e + \gamma(r)h,$$

which defines a complex-valued function $z(r) := z_1(r) + iz_2(r)$. Note that

$$Le = \frac{m}{r}h_3 e, \quad L(J^h e) = \frac{m}{r}h_3 J^h e, \quad Lh = \frac{m}{r}h_3 h + \frac{m}{r}\hat{k}.$$

Hence the operator L restricted to $T_h \mathbb{S}^2$ is equivalent to

$$L_0 := \partial_r + \frac{m}{r}h_3 = h_1 \partial_r \frac{1}{h_1}$$

in the sense that

$$L(z_1 e + z_2 J^h e) = (L_0 z_1) e + (L_0 z_2) J^h e.$$

Roughly speaking, our strategy for proving part (ii) of Theorem 2.1 is to show that $\text{dist}(u, \mathcal{O}_m)$ is controlled by z , which is controlled by $L_0 z$, which in turn is controlled by $\mathcal{E}(u) - 4\pi m$.

Define for radial complex-valued functions $f(r)$ and $g(r)$ the following inner product:

$$\langle f, g \rangle_X := \int_0^\infty \left(\bar{f}_r(r) g_r(r) + \frac{m^2}{r^2} \bar{f}(r) g(r) \right) r dr.$$

If we set $\tilde{f}(y) := f(e^{y/m})$ and $\tilde{g}(y) := g(e^{y/m})$, we have

$$(2.15) \qquad \langle f, g \rangle_X = m \int_{\mathbb{R}} (\bar{\tilde{f}}' \tilde{g}' + \bar{\tilde{f}} \tilde{g}) dy = m \langle \tilde{f}, \tilde{g} \rangle_{H^1(\mathbb{R})}.$$

LEMMA 2.5 *There are $\varepsilon, C > 0$ such that if $f : [0, \infty) \rightarrow \mathbb{C}$ satisfies*

$$(2.16) \quad |\langle f, h_1 \rangle_X| \leq \varepsilon \|f\|_X$$

then

$$(2.17) \quad \|f\|_X^2 = \int_0^\infty \left(|f_r|^2 + \frac{m^2}{r^2} |f|^2 \right) r \, dr \leq C \int_0^\infty |L_0 f|^2 r \, dr.$$

PROOF: We may assume $f(r)$ is real-valued. Under our change of variables $\tilde{f}(y) := f(e^{y/m})$, we have

$$\int_0^\infty \left(f_r^2 + \frac{m^2}{r^2} f^2 \right) r \, dr = m \int_{\mathbb{R}} ((\tilde{f}')^2 + \tilde{f}^2) dy$$

and

$$\int_0^\infty [L_0 f]^2 r \, dr = m \int_{\mathbb{R}} [\tilde{L}_0 \tilde{f}]^2 dy$$

where

$$\tilde{L}_0 := \frac{d}{dy} + \tanh y = \operatorname{sech}(y) \frac{d}{dy} \frac{1}{\operatorname{sech} y}.$$

In the variable y , the *assumption* of the lemma becomes

$$(2.18) \quad \langle \tilde{f}, \operatorname{sech} y \rangle_{H^1(\mathbb{R})} = \int_{\mathbb{R}} (\tilde{f}'(\operatorname{sech} y)' + \tilde{f} \operatorname{sech} y) dy \leq \frac{\varepsilon}{\sqrt{m}} \|\tilde{f}\|_{H^1(\mathbb{R})},$$

and so it suffices to prove that (2.18) implies

$$(2.19) \quad \|\tilde{f}\|_{H^1(\mathbb{R})}^2 \leq C \|\tilde{L}_0 \tilde{f}\|_{L^2(\mathbb{R})}^2 = C \langle \tilde{f}, \tilde{L}_0^* \tilde{L}_0 \tilde{f} \rangle_{L^2(\mathbb{R})}.$$

The second-order differential operator

$$H := \tilde{L}_0^* \tilde{L}_0 = -\frac{d^2}{dy^2} + 1 - 2 \operatorname{sech}^2(y)$$

is nonnegative with unique zero eigenfunction (“ground state”) $\operatorname{sech}(y)$ (in fact, this operator is well studied; see, e.g., [15]). Set $\phi := (1/\sqrt{2}) \operatorname{sech} y$ so that $\|\phi\|_{L^2(\mathbb{R})} = 1$. Write

$$\tilde{f} = a\phi + \varphi \quad \text{with } \langle \phi, \varphi \rangle = 0.$$

Similarly, decompose

$$\phi = b(\phi'' - \phi) + \psi \quad \text{with } \langle (\phi'' - \phi), \psi \rangle = 0.$$

Since $\phi'' - \phi = -\sqrt{2} \operatorname{sech}^3 y$, ψ is nonzero, $b = -\|\phi^2\|_{L^2(\mathbb{R})}^2 / (2\|\phi^3\|_{L^2(\mathbb{R})}^2)$, and $\|\psi\|_{L^2} < 1$.

Now by (2.18),

$$\begin{aligned} |a| &= |\langle \phi, \tilde{f} \rangle| \leq |\langle b(\phi'' - \phi), \tilde{f} \rangle| + |\langle \psi, \tilde{f} \rangle| \\ &= |b| |\langle \phi, \tilde{f} \rangle_{H^1}| + |\langle \psi, \tilde{f} \rangle| \leq C\varepsilon \|\tilde{f}\|_{H^1(\mathbb{R})} + |\langle \psi, \tilde{f} \rangle|. \end{aligned}$$

Using the above estimate and $(A + B)^2 \leq q'A^2 + qB^2$ where q' is the Hölder conjugate of q with $1 < q < \infty$, we have

$$\begin{aligned} \langle \tilde{f}, \tilde{f} \rangle &= |a|^2 + \|\varphi\|_{L^2}^2 \leq (C\varepsilon\|\tilde{f}\|_{H^1(\mathbb{R})} + |\langle \psi, \tilde{f} \rangle|)^2 + \|\varphi\|_{L^2}^2 \\ &\leq Cq'\varepsilon^2\|\tilde{f}\|_{H^1(\mathbb{R})}^2 + q|\langle \psi, \tilde{f} \rangle|^2 + \|\varphi\|_{L^2}^2 \\ &\leq Cq'\varepsilon^2\|\tilde{f}\|_{H^1(\mathbb{R})}^2 + q\|\psi\|_{L^2}^2\|\tilde{f}\|_{L^2}^2 + \|\varphi\|_{L^2}^2. \end{aligned}$$

Choose q to be such that $q\|\psi\|_{L^2}^2 < 1$ to obtain

$$\|\tilde{f}\|_{L^2}^2 \leq C\varepsilon^2\|\tilde{f}\|_{H^1(\mathbb{R})}^2 + C\langle \tilde{f}, H\tilde{f} \rangle,$$

where we used $\|\varphi\|_{L^2}^2 \leq C\langle \varphi, H\varphi \rangle = C\langle \tilde{f}, H\tilde{f} \rangle$. On the other hand, we have

$$\langle \tilde{f}', \tilde{f}' \rangle \leq \langle \tilde{f}, H\tilde{f} \rangle + C\|\tilde{f}\|_{L^2}^2.$$

Combining the two estimates above, we get

$$\|\tilde{f}\|_{H^1(\mathbb{R})}^2 \leq C\langle \tilde{f}, H\tilde{f} \rangle = C \int [\tilde{L}_0\tilde{f}]^2 dy,$$

provided ε is sufficiently small. Transforming back to the variable r , we obtain estimate (2.17), completing the proof. \square

Step 7. Almost Orthogonality. To apply the previous lemma to $z(r)$, we need to verify condition (2.16), for which we use the following lemma:

LEMMA 2.6 *For $z(r)$ and $\gamma(r)$ defined by (2.11) and (2.14),*

$$(2.20) \quad \langle z_1, h_1 \rangle_X = \int_0^\infty \left(z_{1r}h_{1r} + \frac{m^2}{r^2} z_1h_1 \right) r dr = 0,$$

$$(2.21) \quad \langle z_2, h_1 \rangle_X = \int_0^\infty \left(z_{2r}h_{1r} + \frac{m^2}{r^2} z_2h_1 \right) r dr = \int_0^\infty \frac{4m^2}{r^2} h_1^2 h_3 \gamma r dr.$$

PROOF: The pair $(\alpha(u), s(u))$ is the minimizer of the differentiable function

$$F(s, \alpha) = \int_{\mathbb{R}^2} |\nabla(e^{m\theta R}(v - h^{\alpha,s}))|^2 dx.$$

The lemma follows from the equations $\nabla_{\alpha,s} F(\alpha(u), s(u)) = 0$. \square

Step 8. Proof of (ii). We will use the following abbreviation:

$$ze := z_1e + z_2J^h e, \quad z = z_1 + iz_2.$$

By (2.7), we may choose δ sufficiently small so that

$$\|u - e^{m\theta R}h^{s,\alpha}\|_{\dot{H}^1} < \delta_0$$

for any given $\delta_0 > 0$ (which will be specified later). So we have

$$(2.22) \quad \|\xi_r\|_{L^2(\mathbb{R}^2)}^2 + \left\| \frac{m}{r} R\xi \right\|_{L^2(\mathbb{R}^2)}^2 = \|u - e^{m\theta R}h^{s,\alpha}\|_{\dot{H}^1}^2 < \delta_0^2.$$

It is proved in Lemma 4.7 that (2.22) implies the L^∞ -smallness of $\xi(r)$ for δ_0 sufficiently small: $\|\xi\|_{L^\infty} \leq C\delta_0$. This immediately implies

$$(2.23) \quad \|z\|_\infty + \|\gamma\|_\infty \leq C\delta_0.$$

Since $1 = |v|^2 = |z|^2 + (1 + \gamma)^2$ and $|\gamma|$ is small,

$$(2.24) \quad \gamma = \sqrt{1 - |z|^2} - 1 \leq 0, \quad |\gamma(r)| \leq C|z(r)|^2, \quad |\gamma_r(r)| \leq C|z(r)z_r(r)|.$$

By Lemma 2.6, equation (2.23) through (2.24), and $|h(r)| \leq 1$,

$$|\langle z, h_1 \rangle_X| = \left| \int_0^\infty \frac{4m^2}{r^2} h_1^2 h_3 \gamma r \, dr \right| \leq C\|z\|_\infty \left\| \frac{h_1}{r} \right\|_2 \left\| \frac{z}{r} \right\|_2 \leq C\delta_0 \|z\|_X.$$

Taking δ_0 small enough so that $C\delta_0$ is less than ε in Lemma 2.5, we have

$$(2.25) \quad \begin{aligned} \text{dist}^2(u, \mathcal{O}_m) &= 2\pi \int_0^\infty \left(|\xi_r|^2 + \frac{m^2}{r^2} |R\xi|^2 \right) r \, dr \\ &\leq C \int_0^\infty \left(|z_r|^2 + \frac{|z|^2}{r^2} \right) r \, dr \leq C \int_0^\infty |L_0 z|^2 r \, dr. \end{aligned}$$

On the other hand, by (2.12),

$$\delta_1^2 = \mathcal{E}(u) - 4\pi m = \frac{1}{2} \int_{\mathbb{R}^2} \left| L\xi + \frac{m}{r} \xi_3 \xi \right|^2 dx = \frac{1}{2} \int_{\mathbb{R}^2} \left| L(ze) + L(\gamma h) + \frac{m}{r} \xi_3 \xi \right|^2 dx.$$

Hence

$$\int_{\mathbb{R}^2} |L_0 z|^2 dx = \int_{\mathbb{R}^2} |L(ze)|^2 dx \leq 3\delta_1^2 + C \int_{\mathbb{R}^2} \left(|L(\gamma h)|^2 + \left| \frac{m}{r} \xi_3 \xi \right|^2 \right) dx.$$

Using (2.23) we have, for δ_0 sufficiently small,

$$\int_{\mathbb{R}^2} \left| \frac{m}{r} \xi_3 \xi \right|^2 dx \leq C\|z\|_\infty^2 \int_{\mathbb{R}^2} \left| \frac{z}{r} \right|^2 dx.$$

For the term $\|L(\gamma h)\|_{L^2}^2$, using (2.24) we find

$$\int_{\mathbb{R}^2} |L(\gamma h)|^2 \leq C \int_{\mathbb{R}^2} \left(|\gamma_r h|^2 + \left| \frac{\gamma h}{r} \right|^2 + \left| \frac{\gamma}{r} h_3 h \right|^2 \right) \leq C\|z\|_\infty^2 \int_{\mathbb{R}^2} |L_0 z|^2.$$

Thus we get $(1 - C\delta_0^2)\|L_0 z\|_{L^2}^2 \leq 4\delta_1^2$. Now choose $\delta_0 > 0$ sufficiently small (by choosing δ small) so that $\|L_0 z\|_{L^2}^2 \leq 5\delta_1^2$ and therefore $\text{dist}(u, \mathcal{O}_m) < C\delta_1$. This completes the proof of part (ii) of Theorem 2.1.

Step 9. Proof of (iii). By (2.9) with $s = s(u)$ and $\alpha = \alpha(u)$, we have $\int_{\mathbb{R}^2} \nabla u \cdot \nabla(e^{m\theta R} h^{s,\alpha}) dx > 2\pi m$, since δ_1 is small. On the other hand, inequality (4.2) from Lemma 4.1 in Section 4 gives an upper bound: for any $\sigma \in (0, 1)$,

$$\begin{aligned} \int_{\mathbb{R}^2} \nabla u \cdot \nabla(e^{m\theta R} h^{s,\alpha}) dx &\leq \left(\int_{\mathbb{R}^2} \frac{|\nabla u|^2}{|x|^{2\sigma}} dx \right)^{1/2} \left(\int_{\mathbb{R}^2} |x|^{2\sigma} |\nabla(e^{m\theta R} h^{s,\alpha})|^2 dx \right)^{1/2} \\ &\leq C \|\nabla u\|_{L^2(\mathbb{R}^2)}^{1-\sigma} \|\nabla^2 u\|_{L^2(\mathbb{R}^2)}^\sigma \left\| |x|^\sigma |\nabla(e^{m\theta R} h^{s,\alpha})| \right\|_{L^2(\mathbb{R}^2)} \\ &\leq C \|\nabla^2 u\|_{L^2(\mathbb{R}^2)}^\sigma s^\sigma. \end{aligned}$$

Thus we obtain $\|\nabla^2 u\|_{L^2(\mathbb{R}^2)} \geq C/s$, completing the proof of Theorem 2.1. □

3 Global Well-Posedness vs. Blowup

In this section we complete the proof of our main result, Theorem 1.1.

Let $u \in C([0, T]; \dot{H}^2 \cap \Sigma_m)$ be a solution of the Schrödinger flow equation (1.1) that conserves energy. We are assuming

$$\delta_1^2 := \mathcal{E}(u_0) - 4\pi m = \mathcal{E}(u(t)) - 4\pi m < \delta^2$$

where δ is taken to be sufficiently small. In particular, we choose δ small enough so that Theorem 2.1 applies for each $t \in [0, T)$ and so furnishes us with continuous functions $s(t) \in (0, \infty)$ and $\alpha(t) \in \mathbb{R}$ such that

$$(3.1) \quad \left\| u(x, t) - e^{(m\theta + \alpha(t))R} h\left(\frac{r}{s(t)}\right) \right\|_{\dot{H}^1} < C\delta_1$$

and

$$s(t)\|u(t)\|_{\dot{H}^2} > C > 0.$$

Thus the first part of Theorem 1.1 is proved. It remains to show that

$$(3.2) \quad T < \infty \quad \text{and} \quad \lim_{t \rightarrow T^-} \|u(t)\|_{\dot{H}^2} = \infty \implies \liminf_{t \rightarrow T^-} s(t) = 0.$$

Recall that we are writing

$$u(x, t) = e^{m\theta R} v(r), \quad v(r) = e^{\alpha(t)R} \left[h\left(\frac{r}{s(t)}\right) + \xi\left(\frac{r}{s(t)}\right) \right],$$

and (3.1) is equivalent to

$$\|e^{m\theta R} \xi(r)\|_{\dot{H}^1}^2 = \|\xi_r(r)\|_{L^2}^2 + m^2 \left\| \frac{R\xi(r)}{r} \right\|_{L^2}^2 < C\delta_1^2.$$

To prove (3.2), we need estimates showing that $\|u(t)\|_{\dot{H}^2}$ is controlled as long as $s(t)$ is bounded away from zero. These \dot{H}^2 -estimates are obtained using the fact that the coordinates of the tangent vector field $v_r - (m/r)J^v Rv$ with respect to a

certain orthonormal frame satisfy a nonlinear Schrödinger-type equation and can be estimated using Strichartz estimates.

This construction, which was introduced in [3], begins with a unit tangent vector field $\hat{e}(r) \in T_{v(r)}\mathbb{S}^2$ satisfying the parallel transport condition

$$D_r^v \hat{e} \equiv 0.$$

Recall that D_r^v is the covariant derivative acting on vector fields $\eta(r) \in T_{v(r)}\mathbb{S}^2$:

$$D_r^v \eta := P^v \eta_r = \eta_r - (v \cdot \eta_r)v = \eta_r + (v_r \cdot \eta)v.$$

So $\hat{e}(r)$ and $J^v \hat{e}(r)$ form an orthonormal frame on $T_v \mathbb{S}^2$. Then $q(r) = q_1(r) + iq_2(r)$ is defined to be the coordinates of $v_r - (m/r)J^v Rv \in T_v \mathbb{S}^2$ in the basis

$$v_r - \frac{m}{r} J^v Rv = q_1 \hat{e} + q_2 J^v \hat{e}.$$

We will sometimes write $q\hat{e} := q_1 \hat{e} + q_2 J^v \hat{e}$ for convenience. Note that by (1.11),

$$\|q\|_{L^2(r dr)}^2 = \left\| v_r - \frac{m}{r} J^v Rv \right\|_{L^2(r dr)}^2 = \frac{1}{\pi} (\mathcal{E}(u) - 4\pi m) = \frac{\delta_1^2}{\pi}$$

is constant in time and can be taken small.

Define $v = v_1 + iv_2$ as follows:

$$J^v Rv = v_1 \hat{e} + v_2 J^v \hat{e}.$$

Again, we will sometimes denote $v\hat{e} := v_1 \hat{e} + v_2 J^v \hat{e}$. It is now straightforward, if somewhat involved, to show that if $u(x, t)$ solves the Schrödinger map equation (1.1), the complex function $q(r, t)$ solves the following nonlinear Schrödinger equation with nonlocal nonlinearity (see [3] for more details):

$$(3.3) \quad iq_t = -\Delta_r q + \frac{(1 - mv_3)^2}{r^2} q + \frac{m(v_3)_r}{r} q + qN(q),$$

where

$$N(q) = \operatorname{Re} \int_r^\infty \left(\bar{q} + \frac{m}{r} \bar{v} \right) \left(q_r + \frac{1 - mv_3}{r} q \right) dr.$$

By changing variables to $\tilde{q} := e^{i(m+1)\theta} q$, we obtain

$$(3.4) \quad i\tilde{q}_t + \Delta \tilde{q} - \frac{m(1 + v_3)(mv_3 - m - 2)}{r^2} \tilde{q} - \frac{mv_{3,r}}{r} \tilde{q} - \tilde{q}N(q) = 0.$$

We will use this equation to obtain H^1 -estimates on q .

For these estimates of q to be useful, we need to bound the original map $u(x, t)$ — or, equivalently, $v(r, t)$ or $\xi(r, t)$ or $z(r, t)$ — by q . Since $v(r, t) = e^{\alpha R}[h(r/s) + \xi(r/s, t)]$, we have

$$(3.5) \quad q\hat{e} = \frac{1}{s} e^{\alpha R} \left[L\xi + \frac{m}{r} \xi_3 \xi \right] \left(\frac{r}{s} \right),$$

where, recall,

$$L\xi = \xi_r + \frac{m}{r} (\xi_3 h + h_3 \xi) \quad \text{and} \quad \xi = z_1 e + z_2 J^h e + \gamma h.$$

Since ξ is small, we have, very roughly speaking,

$$\xi \approx ze \quad \text{and} \quad L\xi \approx (L_0z)e,$$

and z can be controlled by L_0z . More precisely, Lemma 4.8, proved in the next section, gives the following bounds: for $2 \leq p < \infty$, and provided δ is sufficiently small,

$$(3.6) \quad \|z_r\|_{L^p(\mathbb{R}^2)} + \left\| \frac{z}{r} \right\|_{L^p(\mathbb{R}^2)} \leq C(s^{1-2/p}\|q\|_{L^p(\mathbb{R}^2)} + \|q\|_{L^2(\mathbb{R}^2)}),$$

$$(3.7) \quad \|z_{rr}\|_{L^2(\mathbb{R}^2)} \leq C\left(s\|q_r\|_{L^2(\mathbb{R}^2)} + s\left\| \frac{q}{r} \right\|_{L^2(\mathbb{R}^2)} + s\|q\|_{L^4(\mathbb{R}^2)}^2 + \|q\|_{L^2(\mathbb{R}^2)}\right),$$

$$(3.8) \quad \|u\|_{\dot{H}^2(\mathbb{R}^2)} \leq C\left(\frac{1}{s} + \|q_r\|_{L^2(\mathbb{R}^2)} + \left\| \frac{q}{r} \right\|_{L^2(\mathbb{R}^2)} + \|q\|_{L^4(\mathbb{R}^2)}^2 + \frac{1}{s}\|q\|_{L^2(\mathbb{R}^2)}\right).$$

We will use the following notation to denote space-time Lebesgue norms: for an interval $I \subset \mathbb{R}$,

$$\|f\|_{L_t^r L_x^p(\mathbb{R}^2 \times I)} := \int_I \left(\int_{\mathbb{R}^2} |f(x, t)|^p dx \right)^{r/p} dt.$$

We use (3.6)–(3.7) together with Strichartz estimates and equation (3.4) to prove the following estimates for q :

LEMMA 3.1 *For $\tau \geq 0$ and $\sigma > 0$, set $I := (\tau, \tau + \sigma)$, $Q := \mathbb{R}^2 \times I$, and $X(Q) := L_t^4 L_x^4(Q) \cap L_t^\infty L_x^2(Q) \cap L_t^{8/3} L_x^8(Q)$. Define $\underline{s} := \inf_{t \in I} s(t)$. If δ is sufficiently small, we have*

$$(3.9) \quad \|q\|_{X(Q)} \leq C(\|q_0\|_{L_x^2(\mathbb{R}^2)} + (\underline{s}^{-1}\sigma^{1/2} + \|q\|_{L_x^4 L_t^4(Q)}^2)\|q\|_{L_x^4 L_t^4(Q)}),$$

$$(3.10) \quad \|\nabla \tilde{q}\|_{X(Q)} \leq C(\|\nabla \tilde{q}(\tau)\|_{L_x^2} + (\underline{s}^{-1}\sigma^{1/2} + \underline{s}^{-3/2}\sigma^{3/4} + \|q\|_{X(Q)}^2)\|\nabla \tilde{q}\|_{L_t^\infty L_x^2(Q) \cap L_t^4 L_x^4(Q)}).$$

Before proving Lemma 3.1, we show how it completes the proof of our main theorem.

COMPLETION OF PROOF OF THEOREM 1.1: We need to prove (3.2), so suppose that

$$(3.11) \quad \lim_{t \rightarrow T^-} \|u(t)\|_{\dot{H}^2} = \infty$$

and

$$(3.12) \quad \liminf_{t \rightarrow T^-} s(t) = s_0 > 0.$$

Our goal is to derive a contradiction. By (3.12), we have $s(t) \geq s^*$ for all $0 \leq t < T$ for some $s^* > 0$. So we may take $\underline{s} = s^*$ in the estimates (3.9)–(3.10) for

any time interval $I \subset [0, T)$. If σ is sufficiently small depending on s^* ($\sigma^{1/2} < s^*/(2C)$), and $\|q_0\|_{L^2} \leq C\delta_1$ is taken sufficiently small, estimate (3.9) implies

$$\|q\|_{X(Q)} \leq C\delta_1.$$

Using this estimate in (3.10), for δ_1 and σ sufficiently small we obtain

$$\|\nabla\tilde{q}\|_{L_t^\infty L_x^2(Q)} \leq C\|\nabla\tilde{q}(\tau)\|_{L^2}.$$

In particular, taking τ close to T , we see

$$\limsup_{t \rightarrow T^-} \|\nabla\tilde{q}(t)\|_{L^2} < \infty.$$

Then using (3.8) and $|q_r| + |q/r| \leq C|\nabla\tilde{q}|$, we find $\limsup_{t \rightarrow T^-} \|u(t)\|_{\dot{H}^2} < \infty$, contradicting (3.11). This completes the proof of Theorem 1.1. \square

PROOF OF LEMMA 3.1: By applying Strichartz estimates for the inhomogeneous Schrödinger equation (see, e.g., [1]) to (3.4), we get

$$(3.13) \quad \|q\|_{X(Q)} \leq C(\|q(\tau)\|_{L^2} + \|F\|_{L_t^{4/3} L_x^{4/3}(Q)})$$

where

$$F := \frac{m(1+v_3)(mv_3-m-2)}{r^2} \tilde{q} + \frac{m(v_3)_r}{r} \tilde{q} + \tilde{q}N(q).$$

Conservation of the L^2 -norm of $q(t)$ (equivalent to conservation of energy $\mathcal{E}(u)$) means we can replace $\|q(\tau)\|_{L^2}$ by $\|q_0\|_{L^2}$ in (3.13). Using $v_3(r) = h_3(r/s) + \xi_3(r/s)$ and $\xi_3 = z_2h_1 + \gamma h_3$, we find

$$(3.14) \quad \begin{aligned} \left\| \frac{1+v_3}{r^2} \right\|_{L_x^2} &= \frac{1}{s} \left\| \frac{1+h_3+z_2h_1+\gamma h_3}{r^2} \right\|_{L_x^2} \\ &\leq \frac{C}{s} \left(\left\| \frac{1+h_3}{r^2} \right\|_{L_x^2} + \left\| \frac{z_2h_1}{r^2} \right\|_{L_x^2} + \left\| \frac{\gamma h_3}{r^2} \right\|_{L_x^2} \right) \\ &\leq \frac{C}{s} \left(1 + \left\| \frac{z}{r} \right\|_{L_x^2} + \left\| \frac{z}{r} \right\|_{L_x^4}^2 \right) \\ &\leq C(s^{-1}(1 + \|q\|_{L^2} + \|q\|_{L^2}^2) + \|q\|_{L^4}^2). \end{aligned}$$

Using (3.14) and the uniform boundedness of $\|q\|_{L^2}$, we get

$$\begin{aligned} \left\| \frac{m(1+v_3)(mv_3-m-2)}{r^2} \tilde{q} \right\|_{L_t^{4/3} L_x^{4/3}} &\leq C \left\| \frac{1+v_3}{r^2} \tilde{q} \right\|_{L_t^{4/3} L_x^{4/3}} \\ &\leq C \left\| \frac{1+v_3}{r^2} \right\|_{L_x^2} \|q\|_{L_x^4} \Big|_{L_t^{4/3}} \\ &\leq C(\underline{s}^{-1}\sigma^{1/2} + \|q\|_{L_t^4 L_x^4}^2) \|q\|_{L_t^4 L_x^4}. \end{aligned}$$

Next we estimate $\|(m(v_3)_r/r)q\|_{L_t^{4/3}L_x^{4/3}}$. Compute

$$\frac{1}{r}(v_3)_r = \frac{1}{sr} \left[\frac{mh_1^2}{r} + h_1(z_2)_r - \frac{mh_1h_3z_2}{r} + \gamma_r h_3 + \frac{m\gamma h_1^2}{r} \right] \left(\frac{r}{s}\right).$$

Remark 3.2. The notation means evaluate the quantity in the square brackets with r replaced by r/s . We use this notation frequently in what follows.

Therefore, using boundedness of z , h_3 , h_1/r , and $\gamma = O(|z|^2)$, $\gamma_r = O(|zz_r|)$, we have

$$(3.15) \quad \left\| \frac{(v_3)_r}{r} \right\|_{L^2} \leq \frac{C}{s} \left(\left\| \frac{h_1^2}{r^2} \right\|_{L^2} + \|z_r\|_{L^2} + \left\| \frac{z}{r} \right\|_{L^2} \right) \leq Cs^{-1}(1 + \|q\|_{L^2}).$$

So

$$\left\| \frac{m(v_3)_r}{r} q \right\|_{L_t^{4/3}L_x^{4/3}} \leq C \left\| \frac{v_{3,r}}{r} \right\|_{L_x^2} \|q\|_{L_x^4} \Big|_{L_t^{4/3}} \leq C(1 + \|q_0\|_{L^2}) \underline{s}^{-1} \sigma^{1/2} \|q\|_{L_t^4L_x^4}.$$

Next we need to estimate the nonlocal term

$$\begin{aligned} N(q) &= \operatorname{Re} \int_r^\infty \left(\bar{q} + \frac{m\bar{v}}{r} \right) \left(q_r + \frac{1-mv_3}{r} q \right) dr \\ &= -\frac{|q|^2}{2} + \operatorname{Re} \int_r^\infty \left(\frac{1-mv_3}{r} |q|^2 + \frac{m\bar{v}}{r} q_r + \frac{m\bar{v}(1-mv_3)}{r^2} q \right) dr \\ &:= -\frac{|q|^2}{2} + (R_1 + R_2 + R_3). \end{aligned}$$

First note

$$\| |q|^2 q \|_{L_t^{4/3}L_x^{4/3}} = \|q\|_{L_t^4L_x^4}^3.$$

Now consider $\tilde{q}R_1$. Using the estimate

$$\|f(r)\|_{L^2} \leq C \|rf_r(r)\|_{L^2}$$

(Hardy’s inequality in \mathbb{R}^4 for radial functions), we have

$$\|\tilde{q}R_1\|_{L_t^{4/3},L_x^{4/3}} \leq \|q\|_{L^4L^4} \|R_1\|_{L^2L^2} \leq C \|q\|_{L_t^4L_x^4}^3.$$

Next we note that $|\bar{v}| = |J^v Rv| = \sqrt{1-v_3^2} = |Rv|$. We consider next $\tilde{q}R_3$. By Hardy again,

$$\begin{aligned} (3.16) \quad \|\tilde{q}R_3\|_{L_x^{4/3}} &\leq C \left\| \tilde{q} \int_r^\infty \frac{m\bar{v}(1-mv_3)}{r^2} q dr \right\|_{L^{4/3}} \\ &\leq C \left\| \frac{\bar{v}}{r} q \right\|_{L^2} \|q\|_{L^4} \leq C \left\| \frac{\bar{v}}{r} \right\|_{L^4} \|q\|_{L^4}^2 \leq C \left\| \frac{1-v_3^2}{r^2} \right\|_{L^2}^{1/2} \|q\|_{L^4}^2 \\ &\leq C \left\| \frac{1+v_3}{r^2} \right\|_{L^2}^{1/2} \|q\|_{L^4}^2 \leq C(s^{-1/2} + \|q\|_{L^4}) \|q\|_{L^4}^2, \end{aligned}$$

where we used (3.14) again. Thus

$$(3.17) \quad \begin{aligned} \|\tilde{q} R_3\|_{L_t^{4/3}, L_x^{4/3}} &\leq C \left\| (s^{-1/2} + \|q\|_{L_x^4}) \|q\|_{L_x^4}^2 \right\|_{L_t^{4/3}} \\ &\leq C (\underline{s}^{-1/2} \sigma^{1/4} + \|q\|_{L_t^4 L_x^4}) \|q\|_{L_t^4 L_x^4}^2. \end{aligned}$$

It remains to estimate $\tilde{q} R_2$. We rewrite R_2 , using integration by parts,

$$\begin{aligned} R_2 &= -\operatorname{Re} \frac{m}{r} \bar{v} q - \operatorname{Re} \int_r^\infty \left(\frac{m}{r} \bar{v}_r - \frac{m}{r^2} \bar{v} \right) q \, dr \\ &= -\operatorname{Re} \frac{m}{r} \bar{v} q + \operatorname{Re} \int_r^\infty \left(\frac{m v_3}{r} \bar{q} + \frac{m(1 + m v_3)}{r^2} \bar{v} \right) q \, dr, \end{aligned}$$

where we used $\bar{v}_r = -v_3 \bar{q} - \frac{m}{r} v_3 \bar{v}$. So

$$\begin{aligned} \|\tilde{q} R_2\|_{L_t^{4/3}, L_x^{4/3}} &\leq C \left(\left\| q \frac{\bar{v}}{r} \bar{q} \right\|_{L_t^{4/3}, L_x^{4/3}} + \left\| q \int_r^\infty \frac{|q|^2}{r} \, dr \right\|_{L_t^{4/3}, L_x^{4/3}} \right. \\ &\quad \left. + \left\| q \int_r^\infty \frac{\bar{v}}{r^2} q \, dr \right\|_{L_t^{4/3}, L_x^{4/3}} \right). \end{aligned}$$

The last term is estimated as follows:

$$(3.18) \quad \begin{aligned} \left\| q \int_r^\infty \frac{\bar{v}}{r^2} q \, dr \right\|_{L_t^{4/3}, L_x^{4/3}} &\leq C \left\| \frac{\bar{v}}{r} q \right\|_{L_x^2} \|q\|_{L_x^4} \Big\|_{L_t^{4/3}} \\ &\leq C \left\| \frac{\bar{v}}{r} \right\|_{L_x^4} \|q\|_{L_x^4}^2 \Big\|_{L_t^{4/3}} \\ &\leq C \left\| (s^{-1/2} + \|q\|_{L_x^4}) \|q\|_{L_x^4}^2 \right\|_{L_t^{4/3}} \\ &\leq C (\underline{s}^{-1/2} \sigma^{1/4} + \|q\|_{L_t^4 L_x^4}) \|q\|_{L_t^4 L_x^4}^2, \end{aligned}$$

where we used the same computations as in (3.16) and (3.17). The first term can be treated in a similar manner, leading to the same estimate as in (3.18). The estimate for the second term has been done already.

By returning now to (3.13) and using the above estimates, we establish (3.9).

Next we need to estimate the derivative of \tilde{q} in order to establish (3.10). Denote $w := \partial_{x_i} \tilde{q}$ for $i = 1, 2$. Then w solves

$$(3.19) \quad \begin{aligned} i w_t + \Delta w &= \frac{m(1 + v_3)(m v_3 - m - 2)}{r^2} w \\ &\quad + \left(\frac{m(1 + v_3)(m v_3 - m - 2)}{r^2} \right)_{x_i} \tilde{q} \\ &\quad + \frac{m(v_3)_r}{r} w + \left(\frac{m(v_3)_r}{r} \right)_{x_i} \tilde{q} + N(q)w + N(q)_{x_i} \tilde{q}. \end{aligned}$$

Using the previous estimates, we can estimate the various terms involving w in the right-hand side:

- $$\left\| \frac{m(1+v_3)(mv_3-m-2)}{r^2} w \right\|_{L_t^{4/3} L_x^{4/3}} \leq C \|(s^{-1} + \|q\|_{L_x^4}^2) \|w\|_{L_x^4} \|L_t^{4/3}\|$$

$$\leq C(\underline{s}^{-1} \sigma^{1/2} + \|q\|_{L_t^4 L_x^4}^2) \|\nabla \tilde{q}\|_{L_t^4 L_x^4}$$
- $$\left\| \frac{m(v_3)_r}{r} w \right\|_{L_t^{4/3} L_x^{4/3}} \leq C \|s^{-1}\| \|w\|_{L_x^4} \|L_t^{4/3}\| \leq C \underline{s}^{-1} \sigma^{1/2} \|\nabla \tilde{q}\|_{L_t^4 L_x^4}$$
- $$\|N(q)w\|_{L_t^{4/3} L_x^{4/3}} \leq C(\underline{s}^{-1/2} \sigma^{1/4} + \|q\|_{L_t^4 L_x^4}) \|q\|_{L_t^4 L_x^4} \|\nabla \tilde{q}\|_{L_t^4 L_x^4}$$

Now the other terms. First note that $|\nabla \tilde{q}|^2 \sim |\tilde{q}|^2 + |\tilde{q}/r|^2$, and thus

$$\left\| \frac{\tilde{q}}{r} \right\|_{L^p(\mathbb{R}^2)} \leq C \|\nabla \tilde{q}\|_{L^p(\mathbb{R}^2)}$$

for any $1 \leq p \leq \infty$. Due to (3.14) and (3.15), we have

$$\begin{aligned} & \left\| \left(\frac{(1+v_3)(mv_3-m-2)}{r^2} \right)_{x_i} \tilde{q} \right\|_{L_t^{4/3} L_x^{4/3}} \\ & \leq C \left(\left\| \frac{(v_3)_r}{r^2} \tilde{q} \right\|_{L_t^{4/3} L_x^{4/3}} + \left\| \frac{(1+v_3)}{r^3} \tilde{q} \right\|_{L_t^{4/3} L_x^{4/3}} \right) \\ & \leq C \left(\left\| \frac{(v_3)_r}{r} \right\|_{L_x^2} + \left\| \frac{1+v_3}{r^2} \right\|_{L_x^2} \right) \left\| \frac{\tilde{q}}{r} \right\|_{L_x^4} \|L_t^{4/3}\| \\ & \leq C \|(s^{-1} + \|q\|_{L_x^4}^2) \|\nabla \tilde{q}\|_{L_x^4} \|L_t^{4/3}\| \\ & \leq C(\underline{s}^{-1} \sigma^{1/2} + \|q\|_{L_t^4 L_x^4}^2) \|\nabla \tilde{q}\|_{L_t^4 L_x^4}. \end{aligned}$$

Next we consider

$$\left\| \left(\frac{m(v_3)_r}{r} \right)_{x_i} \tilde{q} \right\|_{L_t^{4/3} L_x^{4/3}} \leq C \left(\left\| \frac{(v_3)_r}{r^2} \tilde{q} \right\|_{L_t^{4/3} L_x^{4/3}} + \left\| \frac{(v_3)_{rr}}{r} \tilde{q} \right\|_{L_t^{4/3} L_x^{4/3}} \right).$$

The first term in the right side can be estimated as before:

$$\left\| \left(\frac{(v_3)_r}{r^2} \right) \tilde{q} \right\|_{L_t^{4/3} L_x^{4/3}} \leq C \underline{s}^{-1} \sigma^{1/2} \|\nabla \tilde{q}\|_{L_t^4 L_x^4}.$$

Recalling $(h_1)_r = -(m/r)h_1h_3$, $(h_3)_r = (m/r)h_1^2$, and $h_1/r, h_3$ bounded, we find

$$\begin{aligned} |(v_3)_{rr}| & \leq \frac{1}{s^2} \left| m \left(\frac{h_1^2}{r} \right)_r + (h_1 z_{2,r})_r - m \left(\frac{h_1 h_3 z_2}{r} \right)_r + (\gamma_r h_3)_r + m \left(\frac{\gamma h_1^2}{r} \right)_r \right| \left(\frac{r}{s} \right) \\ & \leq \frac{C}{s^2} \left[\left| \left(\frac{h_1^2}{r} \right)_r \right| + \left| \frac{h_1 z_2}{r^2} \right| + \left| \frac{h_1 (z_2)_r}{r} \right| + |h_1 z_{rr}| + |z z_{rr}| + |z_r|^2 \right] \left(\frac{r}{s} \right). \end{aligned}$$

We estimate term by term:

- $$\begin{aligned} & \left\| \frac{1}{s^2} \left[\left(\frac{h_1^2}{r} \right)_r \left(\frac{r}{s} \right) \right] \frac{q}{r} \right\|_{L_t^{4/3} L_x^{4/3}} \\ & \leq C \left\| \frac{1}{s^2} \left(\left| \frac{h_1^2}{r^2} + \frac{h_1^2 |h_3|}{r^2} \right| \left(\frac{r}{s} \right) \right) \frac{q}{r} \right\|_{L_t^{4/3} L_x^{4/3}} \leq C \underline{s}^{-1} \sigma^{1/2} \|\nabla \tilde{q}\|_{L_t^4 L_x^4}, \end{aligned}$$
- $$\begin{aligned} & \left\| \frac{1}{s^2} \left(\left| \frac{h_1 |z_2|}{r^2} + \frac{h_1 |(z_2)_r|}{r} \right| \left(\frac{r}{s} \right) \right) \frac{q}{r} \right\|_{L_t^{4/3} L_x^{4/3}} \\ & \leq C \left\| s^{-1} \left(\left\| \frac{z}{r} \right\|_{L_x^2} + \|z_r\|_{L_x^2} \right) \frac{q}{r} \right\|_{L_x^4} \Big\|_{L_t^{4/3}} \leq C \underline{s}^{-1} \sigma^{1/2} \|\nabla \tilde{q}\|_{L_t^4 L_x^4}, \end{aligned}$$
- $$\begin{aligned} & \left\| \frac{1}{s^2} \left(|z_r|^2 \left(\frac{r}{s} \right) \right) \frac{q}{r} \right\|_{L_t^{4/3} L_x^{4/3}} \\ & \leq C \left\| \frac{1}{s} \|z_r\|_{L_x^4}^2 \frac{q}{r} \right\|_{L_x^4} \Big\|_{L_t^{4/3}} \leq C (\|q\|_{L_t^4 L_x^4}^2 + \underline{s}^{-1} \sigma^{1/2}) \|\nabla \tilde{q}\|_{L_t^4 L_x^4}. \end{aligned}$$

For the remaining term, using (3.7),

- $$\begin{aligned} & \left\| \frac{1}{s^2} \left((|h_1| + |z|) z_{rr} \right) \left(\frac{r}{s} \right) \frac{q}{r} \right\|_{L_t^{4/3} L_x^{4/3}} \\ & \leq C \left\| \frac{1}{s} \|z_{rr}\|_{L_x^2} \right\| \left\| (|h_1| + |z|) \left(\frac{r}{s} \right) \frac{q}{r} \right\|_{L_x^4} \Big\|_{L_t^{4/3}} \\ & \leq C \left\| \left(\|q_r\|_{L_x^2} + \left\| \frac{q}{r} \right\|_{L_x^2} \right) \right\| \left\| (|h_1| + |z|) \left(\frac{r}{s} \right) \frac{q}{r} \right\|_{L_x^4} \Big\|_{L_t^{4/3}} \\ & \quad + C \left\| \left(\|q\|_{L_x^4}^2 + s^{-1} \|q\|_{L_x^2}^2 \right) \right\| \left\| (|h_1| + |z|) \left(\frac{r}{s} \right) \frac{q}{r} \right\|_{L_x^4} \Big\|_{L_t^{4/3}} \\ & \leq C (\|q\|_{L_t^{8/3} L_x^8}^2 + \underline{s}^{-3/2} \sigma^{3/4}) \|\nabla \tilde{q}\|_{L_t^\infty L_x^2} \\ & \quad + C (\|q\|_{L_t^4 L_x^4}^2 + \underline{s}^{-1} \sigma^{1/2}) \|\nabla \tilde{q}\|_{L_t^4 L_x^4}, \end{aligned}$$

where we used, first by (3.6) with $p = 8$,

$$\begin{aligned} \left\| (|h_1| + |z|) \left(\frac{r}{s} \right) \frac{q}{r} \right\|_{L_x^4} & \leq C \left\| \frac{(|h_1| + |z|)(r/s)}{r} \right\|_{L_x^8} \|q\|_{L_x^8} \\ & \leq C \|q\|_{L_x^8}^2 + C s^{-3/2} (1 + \|q\|_{L_x^2}^2), \end{aligned}$$

and then $\|(|h_1| + |z|)(r/s)(q/r)\|_{L_x^4} \leq C \|q/r\|_{L_x^4}$.

It remains to estimate $N(q)_{x_i} q$:

$$\begin{aligned}
\|N(q)_{x_i} \tilde{q}\|_{L_t^{4/3} L_x^{4/3}} &\leq C \left\| \left(|q|^2 + \left| \frac{\bar{v}}{r} q \right| \right) \left(|q_r| + \left| \frac{q}{r} \right| \right) \right\|_{L_t^{4/3} L_x^{4/3}} \\
&\leq C \left\| \left(|q|^2 + \left| \frac{\bar{v}}{r} q \right| \right) \right\|_{L_x^2} \left\| \left(|q_r| + \left| \frac{q}{r} \right| \right) \right\|_{L_x^4} \left\| \right\|_{L_t^{4/3}} \\
&\leq C \left\| \left(\|q\|_{L_x^4}^2 + \left\| \frac{v}{r} \right\|_{L_x^4}^2 \right) \|\nabla \tilde{q}\|_{L_x^4} \right\|_{L_t^{4/3}} \\
&\leq C (\underline{s}^{-1} \sigma^{1/2} + \|q\|_{L_t^4 L_x^4}^2) \|\nabla \tilde{q}\|_{L_t^4 L_x^4}.
\end{aligned}$$

Now applying Strichartz estimates to (3.19) and using the estimates established above, we obtain

$$\|\nabla \tilde{q}\|_{L_{t,x}^{\infty,2} \cap L_t^4 L_x^4} \leq \|\nabla \tilde{q}(\tau)\|_{L_x^2} + C((\underline{s}^{-1} + \underline{s}^{-2})\sigma^{1/2} + \|q\|_{X(\mathcal{Q})}^2) \|\nabla \tilde{q}\|_{L_{t,x}^{\infty,2} \cap L_t^4 L_x^4}.$$

This completes the proof of Lemma 3.1. \square

4 Technical Lemmas

In this section we collect some of the technical lemmas used in the proof of the main theorem in the previous sections.

4.1 Some Inequalities for Radial Functions

We begin with some inequalities for radial functions.

LEMMA 4.1

(i) Let $0 \leq \sigma < 1$, and suppose $f \in H^1(\mathbb{R}^2)$ is radial. Then

$$(4.1) \quad \int_0^\infty \frac{f^2}{r^{2\sigma}} r dr \leq C_\sigma \left(\int_0^\infty f^2 r dr \right)^{1-\sigma} \left(\int_0^\infty |f_r|^2 r dr \right)^\sigma.$$

We have $\lim_{\sigma \rightarrow 1^-} C_\sigma = \infty$ and the estimate is false if $\sigma = 1$.

(ii) Let $0 \leq \sigma < 1$, and suppose $f \in H^1(\mathbb{R}^2)$. Then

$$(4.2) \quad \int_{\mathbb{R}^2} \frac{f^2}{|x|^{2\sigma}} dx \leq C_\sigma \left(\int_{\mathbb{R}^2} |f|^2 dx \right)^{1-\sigma} \left(\int_{\mathbb{R}^2} |\nabla f|^2 dx \right)^\sigma.$$

(iii) Suppose $f \in H_{\text{loc}}^1(\mathbb{R}^2)$ is radial with $f_r, f/r \in L^2(\mathbb{R}^2)$. Then

$$(4.3) \quad \|f\|_{L^\infty(\mathbb{R}^2)} \leq C \left(\int_0^\infty \left(|f_r|^2 + \frac{f^2}{r^2} \right) r dr \right).$$

PROOF: We first show that

$$(4.4) \quad \int_0^\infty \frac{f^2(r)}{r^{2\sigma}} r dr \leq \frac{1}{1-\sigma} \left(\int_0^\infty \frac{f^2(r)}{r^{4\sigma-2}} r dr \right)^{1/2} \left(\int_0^\infty (f_r(r))^2 r dr \right)^{1/2}.$$

Indeed, by changing the order of integration, we get

$$\begin{aligned} \int_0^\infty \frac{f^2(r)}{r^{2\sigma}} r dr &= - \int_0^\infty \frac{r}{r^{2\sigma}} \int_r^\infty [f^2(s)]_s ds dr \\ &= -2 \int_0^\infty f(s) f_s(s) ds \int_0^s \frac{r}{r^{2\sigma}} dr \\ &= -\frac{1}{1-\sigma} \int_0^\infty s^{2-2\sigma} f(s) f_s(s) ds \\ &\leq \frac{1}{1-\sigma} \left(\int_0^\infty s^{2-4\sigma} f^2(s) s ds \right)^{1/2} \left(\int_0^\infty (f_s(s))^2 s ds \right)^{1/2}. \end{aligned}$$

In particular, if $\sigma = \frac{1}{2}$, (4.4) immediately implies (4.1). We also note that estimate (4.1) is immediate in the case $\sigma = 0$.

Let $\sigma_i = 1 - (1/2^i)$ where $i \geq 0$ is an integer. From estimate (4.4) with $\sigma = \sigma_{i+1}$, we have

$$\begin{aligned} \int_0^\infty \frac{f^2(r)}{r^{2\sigma_{i+1}}} r dr &\leq 2^{i+1} \left(\int_0^\infty \frac{f^2(r)}{r^{2\sigma_i}} r dr \right)^{1/2} \left(\int_0^\infty (f_r(r))^2 r dr \right)^{1/2}, \quad i = 0, 1, 2, \dots \end{aligned}$$

Iterating this estimate, we obtain

$$(4.5) \quad \int_0^\infty \frac{f^2(r)}{r^{2\sigma_{n+1}}} r dr \leq C_{n+1} \left(\int_0^\infty f^2(r) r dr \right)^{1/2^{n+1}} \left(\int_0^\infty (f_r(r))^2 r dr \right)^{1-1/2^{n+1}}$$

for some constant C_{n+1} . (One can solve $C_{n+1} = 2^{n+1} \sqrt{C_n}$ and $C_0 = 1$ to get $C_{n+1} = 2^{2n+2-n}$, which is certainly not the best constant.)

It remains to consider the general case $\sigma \in [0, 1)$. Let $k \geq 0$ be an integer with $\sigma_k \leq \sigma < \sigma_{k+1}$. There exists $0 \leq \theta < 1$ such that $\sigma = \theta\sigma_{k+1} + (1-\theta)\sigma_k = 1 - (2-\theta)/2^{k+1}$. Using the Hölder inequality and (4.5), we get

$$\begin{aligned} \int_0^\infty \frac{f^2}{r^{2\sigma}} r dr &\leq \left(\int_0^\infty \frac{f^2}{r^{2\sigma_{k+1}}} r dr \right)^\theta \left(\int_0^\infty \frac{f^2}{r^{2\sigma_k}} r dr \right)^{1-\theta} \\ &\leq C_\sigma \left(\int_0^\infty f^2(r) r dr \right)^{1-\sigma} \left(\int_0^\infty (f')^2(r) r dr \right)^\sigma, \end{aligned}$$

where $C_\sigma = C_{k+1}^\theta C_k^{1-\theta}$. This completes the proof of the first estimate (4.1).

To see that this estimate fails at the endpoint $\sigma = 1$, fix a smooth, nonnegative, nondecreasing function $\eta(r)$, supported in $(\frac{1}{2}, \infty)$, and with $1 - \eta(r)$ supported in $[0, \frac{3}{2})$. Then it is easy to check that $f_\delta(r) := \eta(r/\delta) - \eta(r)$ provides a counterexample to the endpoint estimate as $\delta \rightarrow 0$. Note that $f_\delta(0) = 0$ for all δ .

The second estimate (4.2) is an immediate consequence of (4.1). Using polar coordinates, we obtain

$$\begin{aligned} & \int_0^{2\pi} \int_0^\infty \frac{|f(r, \theta)|^2}{r^{2\sigma}} r \, dr \, d\theta \\ & \leq C_\sigma \int_0^{2\pi} \left(\int_0^\infty |f(r, \theta)|^2 r \, dr \right)^{1-\sigma} \left(\int_0^\infty |\partial_r f(r, \theta)|^2 r \, dr \right)^\sigma d\theta. \end{aligned}$$

Using $|\partial_r f|^2 \leq |\nabla f|^2$, we obtain estimate (4.2) from Hölder’s inequality.

For the third estimate (4.3), we introduce the new variable y defined by $r = e^y$ and denote $g(y) = f(e^y) = f(r)$. Then it is immediate that

$$\int_0^\infty \left(|f_r(r)|^2 + \frac{f^2(r)}{r^2} \right) r \, dr = \int_{-\infty}^\infty (|g'(y)|^2 + |g(y)|^2) dy.$$

By Sobolev embedding, we have $\|g\|_{L^\infty(\mathbb{R})} \leq C \|g\|_{H^1(\mathbb{R})}$. Transforming back to the original variable completes our proof. \square

LEMMA 4.2 *Let $g : \mathbb{R}^2 \rightarrow \mathbb{C}$ be radial and bounded with $g, g' \in L^p_{\text{loc}}$, $2 < p < \infty$. Assume $(\partial_r - m/r)g(r) \in L^p(\mathbb{R}^2)$ for some $m \geq 1$. Then $g(r)/r \in L^p(\mathbb{R}^2)$ and*

$$\left\| \frac{g(r)}{r} \right\|_{L^p(\mathbb{R}^2)} \leq C \left\| \left(\partial_r - \frac{m}{r} \right) g(r) \right\|_{L^p(\mathbb{R}^2)}.$$

PROOF: Let $0 < r_1 < r_2 < \infty$ and denote $A = \{x \in \mathbb{R}^2 : r_1 < |x| < r_2\}$. Consider

$$I := -2\pi \operatorname{Re} \int_{r_1}^{r_2} \left(g_r - \frac{m}{r} g \right) \left| \frac{g}{r} \right|^{p-2} \frac{\bar{g}}{r} r \, dr.$$

On one hand, $I \leq C \|g/r\|_{L^p(A)}^{p-1} \|g_r - (m/r)g\|_{L^p(A)}$ by the Hölder inequality. On the other hand,

$$\begin{aligned} I &= m \left\| \frac{g}{r} \right\|_{L^p(A)}^p - \frac{2\pi}{p} \frac{|g(r)|^p}{r^{p-2}} \Big|_{r_1}^{r_2} - \frac{p-2}{p} \left\| \frac{g}{r} \right\|_{L^p(A)}^p \\ &\geq \left(m - 1 + \frac{2}{p} \right) \left\| \frac{g}{r} \right\|_{L^p(A)}^p - \frac{2\pi}{p} \frac{|g(r_2)|^p}{r_2^{p-2}}. \end{aligned}$$

Thus

$$\left(m - 1 + \frac{2}{p} \right) \left\| \frac{g}{r} \right\|_{L^p(A)}^p \leq \frac{2\pi}{p} \frac{|g(r_2)|^p}{r_2^{p-2}} + \left\| \frac{g}{r} \right\|_{L^p(A)}^{p-1} \left\| g_r - \frac{m}{r} g \right\|_{L^p(A)}.$$

This gives a bound for $\|g(r)/r\|_{L^p(A)}$ uniformly in r_1 and r_2 . Hence $g(r)/r \in L^p(\mathbb{R}^2)$. As $r_2 \rightarrow \infty$ and $r_1 \rightarrow 0$, we get

$$\left(m - 1 + \frac{2}{p} \right) \left\| \frac{g}{r} \right\|_{L^p(\mathbb{R}^2)}^p \leq \left\| \frac{g}{r} \right\|_{L^p(\mathbb{R}^2)}^{p-1} \left\| g_r - \frac{m}{r} g \right\|_{L^p(\mathbb{R}^2)},$$

where we used $p > 2$ and the boundedness of g . This completes the proof. \square

Remark 4.3. It is essential to assume g is bounded, as can be seen by the example $g(r) = r^m$. If we assume in the above lemma that $g(r) = o(1)$ as $r \rightarrow \infty$, then we also have

$$\left\| \frac{g(r)}{r} \right\|_{L^2(\mathbb{R}^2)} \leq \frac{1}{m} \left\| \left(\partial_r - \frac{m}{r} \right) g(r) \right\|_{L^2(\mathbb{R}^2)}.$$

Using Lemma 4.2, we prove an L^p -version of Lemma 2.5.

LEMMA 4.4 *Let $2 < p < \infty$. There exists $\epsilon > 0$ such that if $f(r)$ is a radial function satisfying $|\langle f, h_1 \rangle_X| \leq \epsilon \|f\|_X$, then*

$$\|f_r\|_{L^p(\mathbb{R}^2)} + \left\| \frac{f}{r} \right\|_{L^p(\mathbb{R}^2)} \leq C(\|L_0 f\|_{L^p(\mathbb{R}^2)} + \|L_0 f\|_{L^2(\mathbb{R}^2)}).$$

Recall $L_0 f := f_r + (m/r)h_3 f$.

PROOF: We note first that it suffices to prove

$$\left\| \frac{f}{r} \right\|_{L^p} \leq C(\|L_0 f\|_{L^p} + \|L_0 f\|_{L^2}),$$

since

$$\|f_r\|_{L^p} \leq C \left(\left\| f_r + \frac{m}{r} h_3 f \right\|_{L^p} + \left\| \frac{m}{r} h_3 f \right\|_{L^p} \right) \leq C \left(\|L_0 f\|_{L^p} + \left\| \frac{f}{r} \right\|_{L^p} \right).$$

Let $\varphi : [0, \infty) \rightarrow \mathbb{R}$ be a standard cutoff function with

$$0 \leq \varphi \leq 1, \quad \varphi(r) \equiv 1 \quad \text{for } r < 1, \quad \varphi(r) \equiv 0 \quad \text{for } r > 2,$$

$$\int_0^\infty \left| \frac{f}{r} \right|^p r \, dr \leq C \left(\int_0^\infty \left| \frac{f\varphi}{r} \right|^p r \, dr + \int_0^\infty \left| \frac{f(1-\varphi)}{r} \right|^p r \, dr \right) \equiv C(\text{I} + \text{II}).$$

We consider the second term II. Since $1 - \varphi \equiv 0$ if $r < 1$, we have

$$\begin{aligned} \text{II} &= \int_0^\infty \left| \frac{f(1-\varphi)}{r} \right|^p r \, dr \leq \|f\|_{L^\infty(\mathbb{R}^2)}^{p-2} \int_0^\infty \left| \frac{f}{r} \right|^2 (1-\varphi) r \, dr \\ &\leq C \|f\|_{L^\infty(\mathbb{R}^2)}^{p-2} \left\| \frac{f}{r} \right\|_{L^2(\mathbb{R}^2)}^2 \leq C \|f\|_{L^\infty(\mathbb{R}^2)}^{p-2} \|L_0 f\|_{L^2(\mathbb{R}^2)}^2, \end{aligned}$$

where we used Lemma 2.5.

Next we consider the term I. Using Lemma 4.2, we have

$$\begin{aligned} \left\| \frac{f\varphi}{r} \right\|_{L^p}^p &\leq C \left\| (f\varphi)_r - \frac{m}{r} f\varphi \right\|_{L^p}^p \\ &\leq C \left(\|\varphi_r f\|_{L^p}^p + \left\| \left(f_r + \frac{mh_3}{r} f \right) \varphi \right\|_{L^p}^p + \left\| \frac{m}{r} (1+h_3) f \varphi \right\|_{L^p}^p \right). \end{aligned}$$

Since φ_r is supported only on $(1, 2)$, and $(1 + h_3)/r$ is bounded, we have

$$\left\| \frac{f\varphi}{r} \right\|_{L^p}^p \leq C(\|f\|_{L^\infty}^p + \|L_0 f\|_{L^p}^p).$$

Since, by Lemma 2.5, $\|f\|_{L^\infty} \leq C\|L_0 f\|_{L^2}$, we obtain

$$\left\| \frac{f}{r} \right\|_{L^p}^p \leq C(\|L_0 f\|_{L^p}^p + \|L_0 f\|_{L^2}^p),$$

completing the proof. \square

4.2 Some Harmonic Map Estimates

Here we prove some facts about the family \mathcal{O}_m of m -equivariant harmonic maps.

The first lemma shows that if m -equivariant harmonic maps $h = h^{0,1}$ and $h^{\alpha,s}$ are close in the sense of energy, then α and s are also close to 0 and 1, respectively.

LEMMA 4.5 *Let $0 < s < \infty$ and $-\pi \leq \alpha < \pi$. There exists $\epsilon > 0$ and $C > 0$ such that if $\int_{\mathbb{R}^2} |\nabla(e^{m\theta R}(h - h^{s,\alpha}))|^2 dx < \delta^2$ for any $\delta < \epsilon$, then $|\alpha| + |s - 1| \leq C\delta$.*

PROOF: Consider the case $s > 1$ (the case $s < 1$ can be treated in the same way). We first note that

$$h - h^{s,\alpha} = \begin{pmatrix} h_1(r) - h_1(r/s) \cos \alpha \\ -h_1(r/s) \sin \alpha \\ h_3(r) - h_3(r/s) \end{pmatrix}.$$

Our assumption is

$$\begin{aligned} \int_{\mathbb{R}^2} |\nabla(e^{m\theta R}(h - h^{s,\alpha}))|^2 dx \\ = 2\pi \int_0^\infty \left(|\partial_r(h - h^{s,\alpha})|^2 + \frac{m^2}{r^2} |R(h - h^{s,\alpha})|^2 \right) r dr < \delta^2. \end{aligned}$$

For $0 \leq r \leq 1$, we have $(s^{2m} - 1)(1 - r^{2m}) \geq 0$, which, when rearranged, yields

$$h_1\left(\frac{r}{s}\right) \leq \frac{2s^m}{s^{2m} + 1} h_1(r) \leq h_1(r), \quad 0 \leq r \leq 1.$$

Using this inequality, we find

$$\begin{aligned} \delta^2 &\geq \int_{\mathbb{R}^2} |\nabla(e^{m\theta R}(h - h^{s,\alpha}))|^2 dx \geq 2\pi \int_0^1 \left| \partial_r \left(h_3(r) - h_3\left(\frac{r}{s}\right) \right) \right|^2 r dr \\ &= 2\pi \int_0^1 \frac{m^2}{r^2} \left(h_1^2(r) - \frac{1}{s^2} h_1^2\left(\frac{r}{s}\right) \right)^2 r dr \end{aligned}$$

$$\begin{aligned} &\geq 2\pi \int_0^1 \frac{m^2}{r^2} \left(h_1^2(r) - h_1^2\left(\frac{r}{s}\right) \right)^2 r \, dr \\ &\geq 2\pi m^2 \frac{(s^{2m} - 1)^4}{(s^{2m} + 1)^4} \int_0^1 \frac{h_1^4(r)}{r^2} r \, dr \\ &= C \frac{(s^{2m} - 1)^4}{(s^{2m} + 1)^4}. \end{aligned}$$

It follows that $|s - 1| \leq C\delta^{1/2}$ if δ is sufficiently small. Now $g(s) := \|\partial_r(h_3(r) - h_3(r/s))\|_{L^2}^2$ is a smooth function of s with $g(1) = g'(1) = 0$ and $g''(1) > 0$, and so by Taylor's theorem, we have $g(s) \geq C(s - 1)^2$ for some $C > 0$, and for $|s - 1| \leq C\delta^{1/2}$, δ sufficiently small. Thus $|s - 1| \leq C\delta$, as required.

Next consider the second component of $h - h^{s,\alpha}$:

$$\delta^2 > 2\pi \sin^2(\alpha) \int_0^\infty \left| \partial_r h_1\left(\frac{r}{s}\right) \right|^2 r \, dr \geq C \sin^2(\alpha)$$

and so $|\sin(\alpha)| < C\delta$ for sufficiently small δ .

Finally, use

$$(1 - \cos(\alpha))h_1(r) = (h - h^{s,\alpha})_1 + \cos(\alpha) \left(h_1\left(\frac{r}{s}\right) - h_1(r) \right)$$

together with the previous results to arrive at $1 - \cos(\alpha) \leq C\delta$, from which (for $\alpha \in [-\pi, \pi)$) $|\alpha| \leq C\delta$ follows. \square

The next lemma is a bound on the curvature of the family \mathcal{O}_m of m -equivariant harmonic maps.

LEMMA 4.6 *There are $\varepsilon > 0$ and $C > 0$ such that if*

$$\|e^{m\theta R}(h^{s_1,\alpha_1}(r) - h^{s_2,\alpha_2}(r))\|_{\dot{H}^1} < \varepsilon,$$

then setting $\bar{s} := \frac{1}{2}[s_1 + s_2]$, $\bar{\alpha} := \frac{1}{2}[\alpha_1 + \alpha_2]$, and $\bar{h} := \frac{1}{2}[h^{s_1,\alpha_1} + h^{s_2,\alpha_2}]$, we have

$$(4.6) \quad \|e^{m\theta R}(\bar{h}(r) - h^{\bar{s},\bar{\alpha}}(r))\|_{\dot{H}^1} < C \|e^{m\theta R}(h^{s_1,\alpha_1}(r) - h^{s_2,\alpha_2}(r))\|_{\dot{H}^1}^2.$$

PROOF: By rotating and rescaling, we may assume $(s_1, \alpha_1) = (1, 0)$. If ε is sufficiently small, Lemma 4.5 gives (taking $\alpha_2 \in [-\pi, \pi)$)

$$(4.7) \quad |s_2 - 1| + |\alpha_2| \leq C \|e^{m\theta R}(h^{1,0}(r) - h^{s_2,\alpha_2}(r))\|_{\dot{H}^1} \leq C\varepsilon.$$

Now set $s(t) := \bar{s} + (t/2)(s_2 - 1)$, $\alpha(t) := \bar{\alpha} + (t/2)\alpha_2$, and $\phi(t) := h^{s(t),\alpha(t)}$. Then

$$\begin{aligned} \bar{h} - h^{\bar{s},\bar{\alpha}} &= \frac{1}{2}[\phi(-1) - \phi(0) + \phi(1) - \phi(0)] \\ &= \frac{1}{2} \left[\int_0^1 \phi'(t) dt - \int_{-1}^0 \phi'(t) dt \right] \\ &= \frac{1}{2} \int_0^1 [\phi'(t) - \phi'(-t)] dt = \frac{1}{2} \int_0^1 \int_{-t}^t \phi''(\tau) d\tau. \end{aligned}$$

Using (4.7), we have

$$\|e^{m\theta R}\phi''(\tau)\|_{\dot{H}^1} \leq C[(s_2 - 1)^2 + \alpha_2^2] \leq C\|e^{m\theta R}(h^{1,0}(r) - h^{s_2,\alpha_2}(r))\|_{\dot{H}^1}^2$$

and (4.6) follows. □

Our next lemma gives L^∞ -smallness for \dot{H}^1 -small perturbations of harmonic maps.

LEMMA 4.7 *For $u \in \Sigma_m$, set $s = s(u)$ and $\alpha = \alpha(u)$, and write*

$$u(r, \theta) = e^{m\theta R}v(r), \quad v(r) = e^{\alpha R}\left[h\left(\frac{r}{s}\right) + \xi\left(\frac{r}{s}\right)\right].$$

There exists $\epsilon > 0$ and $C > 0$ such that if $\delta_0 < \epsilon$ and

$$(4.8) \quad \|\xi_r\|_{L^2(\mathbb{R}^2)}^2 + \left\|\frac{m}{r}R\xi\right\|_{L^2(\mathbb{R}^2)}^2 = \|u - e^{m\theta R}h^{s,\alpha}\|_{\dot{H}^1}^2 < \delta_0^2,$$

then

$$\|\xi\|_{L^\infty} \leq C\delta_0.$$

PROOF: Without loss of generality, we may assume $s = 1$ and $\alpha = 0$. It follows immediately from (4.8) and Lemma 4.1 that

$$\|\xi_i\|_\infty \leq C\delta_0, \quad i = 1, 2.$$

For ξ_3 , we have, as yet, only $\|(\xi_3)_r\|_{L^2} < \delta_0$, and so our aim is to show that $\|\xi_3\|_{L^\infty} \leq C\delta_0$. Under our change of variable $\tilde{\xi}(y) := \xi(m \log(r))$, it suffices to prove that $\|\tilde{\xi}_3\|_{L^2(\mathbb{R})} \leq C\delta_0$, since $\|\tilde{\xi}_3\|_{L^\infty(\mathbb{R})} \leq C\|\tilde{\xi}_3\|_{H^1(\mathbb{R})}$. By the continuity and boundary conditions of $v(r)$ (for $u \in \Sigma_m$), there must exist $y_0 \in \mathbb{R}$ such that $\tilde{v}_3(y_0) = 0$. Note that since $\tilde{h}_3(y) = \tanh(y)$, we have

$$(4.9) \quad |\tilde{v}_3^2(y) - \tanh^2(y)| = \left| \sum_{j=1}^2 (\tilde{v}_j^2 - \tilde{h}_j^2) \right| \leq C(|\xi_1| + |\xi_2|),$$

and in particular,

$$\tanh^2(y_0) \leq C(\|\xi_1\|_{L^\infty} + \|\xi_2\|_{L^\infty}) < C\delta_0.$$

So for $-1 \leq y \leq 1$,

$$(4.10) \quad |\tilde{v}_3(y) - \tanh(y)| = \left| \int_{y_0}^y (\tilde{v}_3' - \tanh'(y))dy - \tanh(y_0) \right| \leq C(\|\tilde{v}_3' - \tanh'\|_{L^2} + \delta_0^{1/2}) \leq C\delta_0^{1/2},$$

and in particular, for δ_0 sufficiently small, $|\tilde{v}_3(\pm 1)| > (1/2)\tanh(1)$. Then with the aid of Lemma 2.3, for δ_0 sufficiently small, we have $|\tilde{v}_3(y)| > (1/4)\tanh(1)$ for $|y| \geq 1$. Estimate (4.9) then yields

$$\int_{|y| \geq 1} (\tilde{v}_3(y) - \tanh(y))^2 dy \leq C\delta_0^2,$$

which also gives us $\sup_{|y| \geq 1} |\tilde{v}_3(y) - \tanh(y)| \leq C\delta_0$ (Sobolev embedding), and in particular $|\tilde{v}_3(1) - \tanh(1)| \leq C\delta_0$. Finally, we get the same result for $|y| < 1$ by integrating the derivative:

$$|\tilde{v}_3(y) - \tanh(y)| \leq \left| \int_y^1 (\tilde{v}'_3 - \tanh') dy \right| + C\delta_0 < C\delta_0.$$

This completes the proof. □

4.3 Perturbation Is Bounded by q

Here we prove estimates used in Section 3. We show that $z(r)$ is controlled by $q(r)$, where, recall,

$$v_r - \frac{m}{r} J^v Rv = q_1 \hat{e} + q_2 J^v \hat{e}, \quad D_r^v \hat{e} \equiv 0,$$

and

$$(4.11) \quad v(r) = e^{\alpha R} \left[h \left(\frac{r}{s} \right) + \xi \left(\frac{r}{s} \right) \right], \quad \xi = z_1 e + z_2 J^h e + \gamma h.$$

LEMMA 4.8 *Let $2 \leq p < \infty$. For δ sufficiently small,*

$$(4.12) \quad \|z_r\|_{L^p(\mathbb{R}^2)} + \left\| \frac{z}{r} \right\|_{L^p(\mathbb{R}^2)} \leq C(s^{1-2/p} \|q\|_{L^p(\mathbb{R}^2)} + \|q\|_{L^2(\mathbb{R}^2)}),$$

$$(4.13) \quad \|z_{rr}\|_{L^2(\mathbb{R}^2)} \leq C \left(s \|q_r\|_{L^2(\mathbb{R}^2)} + s \left\| \frac{q}{r} \right\|_{L^2(\mathbb{R}^2)} + s \|q\|_{L^4(\mathbb{R}^2)}^2 + \|q\|_{L^2(\mathbb{R}^2)} \right),$$

and

$$(4.14) \quad \|u\|_{\dot{H}^2(\mathbb{R}^2)} \leq C \left(\frac{1}{s} + \|q_r\|_{L^2(\mathbb{R}^2)} + \left\| \frac{q}{r} \right\|_{L^2(\mathbb{R}^2)} + \|q\|_{L^4(\mathbb{R}^2)}^2 + \frac{1}{s} \|q\|_{L^2(\mathbb{R}^2)} \right).$$

PROOF: We will first show the following:

$$(4.15) \quad \|z_r\|_{L^2(\mathbb{R}^2)} + \left\| \frac{z}{r} \right\|_{L^2(\mathbb{R}^2)} \leq C \|q\|_{L^2(\mathbb{R}^2)}.$$

By (4.11), we have

$$(4.16) \quad e^{-\alpha R} s(q\hat{e})(sr) = (L_0 z)e + (\gamma h)_r + \frac{2m}{r} h_3 \gamma h + \frac{m}{r} \xi_3 \xi.$$

Since $\|z\|_X \leq C \|L_0 z\|_{L^2}$ by (2.25), it suffices to prove that $\|L_0 z\| \leq C \|q\|_{L^2}$. We first show that $\|\xi_r\|_{L^2} + \|\xi/r\|_{L^2} \leq C \|z\|_X$. Indeed, since $(J^h e)_r = -(m/r)h_1 h$ and $h_r = (m/r)h_1 J^h e$, we find

$$\xi_r = z_r e - z_2 \frac{m}{r} h_1 h + \gamma_r h + \frac{m}{r} \gamma h_1 J^h e.$$

Therefore, since $\gamma = O(|z|^2)$ and $\gamma_r = O(|z||z_r|)$, we obtain

$$\|\xi_r\|_{L^2} \leq \|z_r\|_{L^2} + \left\| \frac{z}{r} \right\|_{L^2} + \|z\|_{L^\infty} \|z_r\|_{L^2} + \|z\|_{L^\infty} \left\| \frac{z}{r} \right\|_{L^2},$$

where we used the boundedness of h . By (2.23), we have $\|z\|_{L^\infty} \leq C\delta$, which can be chosen sufficiently small to yield $\|\xi_r\|_{L^2} \leq C\|z\|_X$. In a similar manner, we can show $\|\xi/r\|_{L^2} \leq C\|z/r\|_{L^2} \leq C\|z\|_X$. Combining, we obtain

$$(4.17) \quad \|\xi_r\|_{L^2} + \left\| \frac{\xi}{r} \right\|_{L^2} \leq C\|z\|_X.$$

Now we are ready to prove $\|L_0 z\|_{L^2} \leq C\|q\|_{L^2}$. Using again $\gamma = O(|z|^2)$, $\gamma_r = O(|z||z_r|)$, and the boundedness of h , we find

$$\begin{aligned} \left\| (\gamma h)_r + \frac{2m}{r} h_3 \gamma h + \frac{m}{r} \xi_3 \xi \right\|_{L^2} &\leq C\|z\|_{L^\infty} \left(\|z_r\|_{L^2} + \left\| \frac{z}{r} \right\|_{L^2} \right) + C\|\xi\|_{L^\infty} \left\| \frac{\xi}{r} \right\|_{L^2} \\ &\leq C\|\xi\|_{L^\infty} \left(\|z\|_X + \left\| \frac{\xi}{r} \right\|_{L^2} \right) \\ &\leq C\|\xi\|_{L^\infty} \|z\|_X \leq C\|\xi\|_{L^\infty} \|L_0 z\|_{L^2}, \end{aligned}$$

where we used (4.17) and $\|z\|_{L^\infty} \leq \|\xi\|_{L^\infty}$. Thus we have

$$\begin{aligned} \|L_0 z\|_{L^2} &\leq \|sq(s \cdot)\|_{L^2} + \left\| (\gamma h)_r + \frac{2m}{r} h_3 \gamma h + \frac{m}{r} \xi_3 \xi \right\|_{L^2} \\ &\leq \|q\|_{L^2} + C\|\xi\|_{L^\infty} \|L_0 z\|_{L^2}. \end{aligned}$$

Since $\|\xi\|_{L^\infty} \leq C\|\xi\|_X \leq C\|z\|_X \leq C\delta$ can be taken sufficiently small, the above inequality implies

$$(4.18) \quad \|L_0 z\|_{L^2} \leq C\|q\|_{L^2},$$

which completes the proof of (4.15). Similarly, for any p with $2 \leq p < \infty$, we have

$$\begin{aligned} \|L_0 z\|_{L^p} &\leq C \left(\|sq(s \cdot)\|_{L^p} + \|z\|_{L^\infty} \left(\|z_r\|_{L^p} + \left\| \frac{z}{r} \right\|_{L^p} \right) \right) \\ &\leq C(s^{1-2/p} \|q\|_{L^p} + \|z\|_{L^\infty} (\|L_0 z\|_{L^p} + \|L_0 z\|_{L^2})), \end{aligned}$$

where we used Lemma 4.4. Since $\|z\|_{L^\infty}$ can be taken sufficiently small, by using (4.18), we finally have $\|L_0 z\|_{L^p} \leq C(s^{1-2/p} \|q\|_{L^p} + \|q\|_{L^2})$, completing the proof of (4.12).

Next we prove (4.13). We first show

$$(4.19) \quad \|z_{rr}\|_{L^2} \leq C \left(\|\partial_r L(z e)\|_{L^2} + \left\| \frac{L(z e)}{r} \right\|_{L^2} + \|L_0 z\|_{L^2} \right).$$

Indeed, recalling $L(ze) = z_re + (m/r)h_3ze$, we have

$$\begin{aligned} \left(\partial_r + \frac{1}{r} + \frac{m}{r}h_3\right)z_re &= \left(\partial_r + \frac{1}{r}\right)z_re + \frac{m}{r}h_3z_re \\ &= \left(\partial_r + \frac{1}{r}\right)\left(L(ze) - \frac{m}{r}h_3ze\right) + \frac{m}{r}h_3z_re \\ &= \left(\partial_r + \frac{1}{r}\right)L(ze) - \frac{m^2}{r^2}h_1^2ze + \frac{m^2}{r^2}z_2h_3h_1h. \end{aligned}$$

Set

$$\eta := e^{m\theta R}ze \quad \text{and} \quad H := e^{m\theta R}h.$$

Since $\partial_j\eta = D_j^H\eta - (\partial_jH \cdot \eta)H$ and $(h_3)_r = (m/r)h_1^2$, we have

$$\begin{aligned} e^{-m\theta R}D_j^H D_j^H \eta &= \left(\left(\partial_r + \frac{1}{r}\right)z_re - \frac{m^2}{r^2}h_3^2ze\right) \\ &= \left(\partial_r + \frac{1}{r}\right)(z_re) + \frac{m}{r}z_2'h_1h - \frac{m^2}{r^2}h_3^2ze \\ &= \left(\partial_r + \frac{1}{r} + \frac{m}{r}h_3\right)(z_re) - \frac{m}{r}h_3L(ze) + \frac{m}{r}z_2'h_1h, \end{aligned}$$

and therefore we obtain

$$\begin{aligned} \|D_j^H D_j^H \eta\|_{L^2} &= \left\| \left(\left(\partial_r + \frac{1}{r}\right)z_re - \frac{m^2}{r^2}h_3^2ze\right) \right\|_{L^2} \\ &= \left\| \left(\partial_r + \frac{1}{r}\right)L(ze) - \frac{m}{r}h_3L(ze) \right. \\ &\quad \left. - \frac{m^2}{r^2}h_1^2ze + \frac{m^2}{r^2}z_2h_3h_1h + \frac{m}{r}z_2'h_1h \right\|_{L^2} \\ &\leq C \left(\left\| \left(\partial_r + \frac{1}{r}\right)L(ze) \right\|_{L^2} + \left\| \frac{L(ze)}{r} \right\|_{L^2} + \|z'\|_{L^2} + \left\| \frac{z}{r} \right\|_{L^2} \right) \\ &\leq C \left(\|\partial_r L(ze)\|_{L^2} + \left\| \frac{L(ze)}{r} \right\|_{L^2} + \|L_0z\|_{L^2} \right), \end{aligned}$$

where we used $|h_1/r|$ bounded. Since $\partial_j\eta = D_j^H\eta - (\partial_jH \cdot \eta)H$, we have

$$\begin{aligned} \Delta\eta &= \partial_j\partial_j\eta = \partial_j(D_j^H\eta - (\partial_jH \cdot \eta)H) \\ &= \partial_jD_j^H\eta - \partial_j((\partial_jH \cdot \eta)H) \\ &= D_j^H D_j^H \eta - (\partial_jH \cdot D_j^H\eta)H - (\partial_j\partial_jH \cdot \eta)H - (\partial_jH \cdot \partial_j\eta)H \\ &\quad - (\partial_jH \cdot \eta)\partial_jH. \end{aligned}$$

Therefore, we obtain

$$\begin{aligned}
\|\Delta\eta\|_{L^2} &\leq \|D_j^H D_j^H \eta\|_{L^2} + \|(\partial_j H \cdot D_j^H \eta)H\|_{L^2} + \|(\partial_j H \cdot \partial_j \eta)H\|_{L^2} \\
&\quad + \|(\partial_j \partial_j H \cdot \eta)H\|_{L^2} + \|(\partial_j H \cdot \eta)\partial_j H\|_{L^2} \\
&\leq \|D_j^H D_j^H \eta\|_{L^2} + C\left(\|\nabla\eta\|_{L^2} + \left\|\frac{\eta}{r}\right\|_{L^2}\right) \\
&\leq \|D_j^H D_j^H \eta\|_{L^2} + C\left(\|z_r\|_{L^2} + \left\|\frac{z}{r}\right\|_{L^2}\right).
\end{aligned}$$

Thus

$$(4.20) \quad \|\eta\|_{\dot{H}^2} \leq C\|\Delta\eta\|_{L^2} \leq C\left(\|\partial_r L(ze)\|_{L^2} + \left\|\frac{L(ze)}{r}\right\|_{L^2} + \|L_0 z\|_{L^2}\right).$$

Let

$$A = (ze)_{rr} - \frac{(ze)_r}{r} + \frac{m^2 ze}{r^2}, \quad B = (ze)_{rr}, \quad E = \frac{mR(ze)_r}{r} - \frac{mRze}{r^2}.$$

Then direct calculations show

$$\begin{aligned}
(e^{m\theta R} ze)_{xx} &= e^{m\theta R} \left(-\frac{y^2}{r^2} A + B - 2\frac{xy}{r^2} E\right), \\
(e^{m\theta R} ze)_{yy} &= e^{m\theta R} \left(-\frac{x^2}{r^2} A + B + 2\frac{xy}{r^2} E\right), \\
(e^{m\theta R} ze)_{xy} &= e^{m\theta R} \left(\frac{xy}{r^2} A + \left(\frac{x^2}{r^2} - \frac{y^2}{r^2}\right) E\right).
\end{aligned}$$

Each of A , B , and E can be expressed in terms of combinations of second derivatives of $\eta = e^{m\theta R} ze$. This implies, in particular, estimate (4.19).

It remains to control $\|\partial_r L(ze)\|_{L^2} + \|L(ze)/r\|_{L^2}$. Noting that $\hat{e}_r = -(v_r \cdot \hat{e})v$ and $(J^v \hat{e})_r = (v_r \cdot \hat{e})v$, so

$$(q\hat{e})_r = q_r \hat{e} - q_1 (v_r \cdot \hat{e})v + q_2 (v_r \cdot \hat{e})v,$$

and recalling $v_r(r) = e^{\alpha R}(1/s)(h_r(r/s) + \xi_r(r/s))$, we have

$$\begin{aligned}
\|\partial_r(q\hat{e})\|_{L^2} &\leq C\left(\|q_r\|_{L^2} + \left\|\frac{q}{r}\right\|_{L^2} + \left\|q \frac{1}{s} z_r \left(\frac{\cdot}{s}\right)\right\|_{L^2}\right) \\
&\leq C\left(\|q_r\|_{L^2} + \left\|\frac{q}{r}\right\|_{L^2} + \|q\|_{L^4}^2 + s^{-1}\|q\|_{L^2}^2\right),
\end{aligned}$$

where we used (4.12). Taking the derivative of Equation (4.16), we get

$$\begin{aligned} \|\partial_r L(ze)\|_{L^2} &\leq C \left(s \|q_r\|_{L^2} + s \left\| \frac{q}{r} \right\|_{L^2} + s \|q\|_{L^4}^2 + \|q\|_{L^2}^2 \right. \\ &\quad \left. + \left\| (\gamma h)_{rr} + \left(2 \frac{m}{r} h_3 \gamma h \right)_r + \left(\frac{m}{r} \xi_3 \xi \right)_r \right\|_{L^2} \right). \end{aligned}$$

We consider first $(\gamma h)_{rr}$. Using $|h_r| + |rh_{rr}| < C$, we have

$$\begin{aligned} \|(\gamma h)_{rr}\|_{L^2} &\leq C \| |\gamma_{rr}| + |\gamma_r h_r| + |\gamma h_{rr}| \|_{L^2} \\ &\leq C \left(\|z\|_{L^\infty} \|z_{rr}\|_{L^2} + \|z_r\|_{L^4}^2 + \left\| |z_r| + \left| \frac{z}{r} \right| \right\|_{L^2} \right) \\ &\leq C (\|z\|_{L^\infty} \|z_{rr}\|_{L^2} + s \|q\|_{L^4}^2 + \|q\|_{L^2}^2 + \|q\|_{L^2}). \end{aligned}$$

Next we consider $((m/r)h_3\gamma h)_r$. In a similar manner, we find

$$\begin{aligned} \left\| \left(\frac{m}{r} h_3 \gamma h \right)_r \right\|_{L^2} &\leq C \left\| \left| \frac{\gamma}{r^2} \right| + \left| \frac{\gamma_r}{r} \right| \right\|_{L^2} \leq C \left(\left\| \frac{z}{r} \right\|_{L^4}^2 + \|z_r\|_{L^4} \left\| \frac{z}{r} \right\|_{L^4} \right) \\ &\leq C (s \|q\|_{L^4}^2 + \|q\|_{L^2}^2). \end{aligned}$$

For the term $(\xi_3 \xi / r)_r$, we have the estimate

$$\left\| \left(\frac{m}{r} \xi_3 \xi \right)_r \right\|_{L^2} \leq C (s \|q\|_{L^4}^2 + \|q\|_{L^2}^2).$$

Following a similar procedure for $\|L(ze)/r\|_{L^2}$, we obtain

$$\begin{aligned} \|\partial_r L(ze)\|_{L^2} + \left\| \frac{L(ze)}{r} \right\|_{L^2} &\leq C \left(s \|q_r\|_{L^2} + s \left\| \frac{q}{r} \right\|_{L^2} + \|z\|_{L^\infty} \|z_{rr}\|_{L^2} + s \|q\|_{L^4}^2 + \|q\|_{L^2}^2 \right) \\ &\leq C \left(s \|q_r\|_{L^2} + s \left\| \frac{q}{r} \right\|_{L^2} + \|z\|_{L^\infty} \left(\|\partial_r L(ze)\|_{L^2} + \left\| \frac{L(ze)}{r} \right\|_{L^2} \right) \right. \\ &\quad \left. + s \|q\|_{L^4}^2 + \|q\|_{L^2}^2 \right), \end{aligned}$$

using (4.19). Since $\|z\|_{L^\infty}$ can be taken sufficiently small, we conclude

$$(4.21) \quad \|\partial_r L(ze)\|_{L^2} + \left\| \frac{L(ze)}{r} \right\|_{L^2} \leq C \left(s \|q_r\|_{L^2} + s \left\| \frac{q}{r} \right\|_{L^2} + s \|q\|_{L^4}^2 + \|q\|_{L^2}^2 \right),$$

having used the smallness of $\|q\|_{L^2}$. Combining (4.21) with (4.19) completes the proof of (4.13).

It remains to prove (4.14). Since $u(x) = e^{(m\theta+\alpha)R}(h(r/s) + \xi(r/s))$, it is straightforward to check that for δ sufficiently small,

$$\|u\|_{\dot{H}^2} \leq \frac{C}{s}(1 + \|\eta\|_{\dot{H}^2})$$

and so (4.14) follows from (4.20) and (4.21). The proof of Lemma 4.8 is complete. \square

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