

Search for CP Violation in $D_s^+ \rightarrow K_S^0 \pi^+$, $D^+ \rightarrow K_S^0 K^+$, and $D^+ \rightarrow \phi \pi^+$ Decays

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A search for charge-parity (CP) violation in Cabibbo-suppressed $D_s^+ \rightarrow K_S^0 \pi^+$, $D^+ \rightarrow K_S^0 K^+$, and $D^+ \rightarrow \phi \pi^+$ decays is reported using proton-proton collision data, corresponding to an integrated luminosity of 3.8 fb^{-1} , collected at a center-of-mass energy of 13 TeV with the LHCb detector. High-yield samples of kinematically and topologically similar Cabibbo-favored $D_{(s)}^+$ decays are analyzed to subtract nuisance asymmetries due to production and detection effects, including those induced by CP violation in the neutral kaon system. The results are

$$\begin{aligned} \mathcal{A}_{CP}(D_s^+ \rightarrow K_S^0 \pi^+) &= (1.3 \pm 1.9 \pm 0.5) \times 10^{-3}, \\ \mathcal{A}_{CP}(D^+ \rightarrow K_S^0 K^+) &= (-0.09 \pm 0.65 \pm 0.48) \times 10^{-3}, \\ \mathcal{A}_{CP}(D^+ \rightarrow \phi \pi^+) &= (0.05 \pm 0.42 \pm 0.29) \times 10^{-3}, \end{aligned}$$

where the first uncertainties are statistical and the second systematic. They are the most precise measurements of these quantities to date, and are consistent with CP symmetry. A combination with previous LHCb measurements, based on data collected at 7 and 8 TeV, is also reported.

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Violation of charge-parity (CP) symmetry arises in the standard model (SM) of particle physics through the complex phase of the Cabibbo-Kobayashi-Maskawa (CKM) quark-mixing matrix [1,2]. CP violation is well established in K - and B -meson systems [3–7], and has been observed only recently in charm decays [8]. CP violation in charm decays can arise from the interference between tree- and loop-level diagrams through Cabibbo-suppressed $c \rightarrow d\bar{d}u$ and $c \rightarrow s\bar{s}u$ transition amplitudes. In the loop-level processes, contributions from physics beyond the SM may arise that can lead to additional sources of CP violation [9]. However, the expected SM contribution is difficult to compute due to the presence of low-energy strong-interaction effects, with current predictions spanning several orders of magnitude [9–13]. A promising handle to determine the origin of possible CP -violation signals are correlations between CP asymmetries in flavor- $SU(3)$ related decays [14–22]. Particularly interesting in this respect are D_s^+ and D^+ decays to two-body (or quasi-two-body) final states, such as $D_s^+ \rightarrow K_S^0 \pi^+$, $D^+ \rightarrow K_S^0 K^+$, and $D^+ \rightarrow \phi \pi^+$. (The inclusion of charge-conjugate processes is implied throughout this Letter, unless stated

otherwise.) Searches for CP violation in these modes have been performed by the CLEO [23], BABAR [24,25], Belle [26–28], and LHCb [29,30] collaborations. No evidence for CP violation has been found within a precision of a few per mille.

This Letter presents measurements of CP asymmetries in $D_s^+ \rightarrow K_S^0 \pi^+$, $D^+ \rightarrow K_S^0 K^+$, and $D^+ \rightarrow \phi \pi^+$ decays performed using proton-proton collision data collected with the LHCb detector between 2015 and 2017 at a center-of-mass energy of 13 TeV, and corresponding to an integrated luminosity of 3.8 fb^{-1} . In the presence of a K_S^0 meson in the final state, a CP asymmetry is expected to be induced by $K^0-\bar{K}^0$ mixing [31]. This effect is well known and predictable, allowing for a precise measurement of CP violation in the charm-quark transition. The $D^+ \rightarrow \phi \pi^+$ decay is reconstructed with the $\phi \rightarrow K^+ K^-$ mode. Several intermediate states contribute to the $D^+ \rightarrow K^+ K^- \pi^+$ decay amplitude [32]. In this Letter, no attempt is made to separate them through an amplitude analysis, and the measurement is performed by simply restricting the $K^+ K^-$ pair to the mass region around the $\phi(1020)$ resonance.

The CP asymmetry of a $D_{(s)}^+$ meson decaying to the final state f^+ is defined as

$$\mathcal{A}_{CP}(D_{(s)}^+ \rightarrow f^+) \equiv \frac{\Gamma(D_{(s)}^+ \rightarrow f^+) - \Gamma(D_{(s)}^- \rightarrow f^-)}{\Gamma(D_{(s)}^+ \rightarrow f^+) + \Gamma(D_{(s)}^- \rightarrow f^-)}, \quad (1)$$

where Γ is the partial decay rate. If CP symmetry is violated in the decay, $\mathcal{A}_{CP} \neq 0$. An experimentally

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convenient quantity to measure is the “raw” asymmetry of the observed yields N ,

$$A(D_{(s)}^+ \rightarrow f^+) \equiv \frac{N(D_{(s)}^+ \rightarrow f^+) - N(D_{(s)}^- \rightarrow f^-)}{N(D_{(s)}^+ \rightarrow f^+) + N(D_{(s)}^- \rightarrow f^-)}. \quad (2)$$

The raw asymmetry can be approximated as

$$A(D_{(s)}^+ \rightarrow f^+) \approx \mathcal{A}_{CP}(D_{(s)}^+ \rightarrow f^+) + A_P(D_{(s)}^+) + A_D(f^+), \quad (3)$$

where $A_P(D_{(s)}^+)$ is the asymmetry of the $D_{(s)}^+$ -meson production cross section [33,34] and $A_D(f^+)$ is the asymmetry of the reconstruction efficiency for the final state f^+ . When $f^+ = K_S^0 h^+$ (with $h = K, \pi$), the detection asymmetry receives contributions from the h^+ hadron (indicated as companion hadron in the following), $A_D(h^+)$, and from the neutral kaon, $A_D(\bar{K}^0)$. Relevant instrumental effects contributing to $A_D(h^+)$ may include differences in interaction cross sections with matter between positive and negative hadrons and the slightly charge-asymmetric performance of the reconstruction algorithms. The contribution to $A_D(\bar{K}^0)$ arises from K^0 and \bar{K}^0 mesons having different interaction cross sections with matter and from their propagation in the detector being affected by the presence of CP violation in the K^0 - \bar{K}^0 system. When $f^+ = \phi(\rightarrow K^+ K^-)\pi^+$, the detection asymmetry is mostly due to the charged pion, as the contributions from the oppositely charged kaons cancel to a good precision.

The detection and production asymmetries are canceled by using the decays $D^+ \rightarrow K_S^0 \pi^+$, $D_s^+ \rightarrow K_S^0 K^+$, and $D_s^+ \rightarrow \phi \pi^+$, which proceed through the Cabibbo-favored $c \rightarrow s \bar{d} u$ transition. In the SM, these decays are expected to have CP asymmetries that are negligibly small compared to the Cabibbo-suppressed modes, when effects induced by the neutral kaons are excluded [31,35]. Hence, their raw asymmetries can be approximated as in Eq. (3), but with $\mathcal{A}_{CP} = 0$. The CP asymmetries of the decay modes of interest are determined by combining the raw asymmetries as follows:

$$\mathcal{A}_{CP}(D_s^+ \rightarrow K_S^0 \pi^+) \approx A(D_s^+ \rightarrow K_S^0 \pi^+) - A(D_s^+ \rightarrow \phi \pi^+), \quad (4)$$

$$\begin{aligned} \mathcal{A}_{CP}(D^+ \rightarrow K_S^0 K^+) &\approx A(D^+ \rightarrow K_S^0 K^+) - A(D^+ \rightarrow K_S^0 \pi^+) \\ &\quad - A(D_s^+ \rightarrow K_S^0 K^+) + A(D_s^+ \rightarrow \phi \pi^+), \end{aligned} \quad (5)$$

$$\mathcal{A}_{CP}(D^+ \rightarrow \phi \pi^+) \approx A(D^+ \rightarrow \phi \pi^+) - A(D^+ \rightarrow K_S^0 \pi^+), \quad (6)$$

where the contribution from $A_D(\bar{K}^0)$ is omitted and should be subtracted from any of the measured asymmetries where it is present.

The LHCb detector [36,37] is a single-arm forward spectrometer designed for the study of particles containing b or c quarks. The detector elements that are particularly relevant to this analysis are a silicon-strip vertex detector that allows for a precise measurement of the impact parameter, i.e., the minimum distance of a charged-particle trajectory to a pp interaction point (primary vertex), a tracking system that provides a measurement of the momentum of charged particles, two ring-imaging Cherenkov detectors that are able to discriminate between different species of charged hadrons, and a calorimeter system that is used for the identification of photons, electrons and hadrons. The polarity of the magnetic field is periodically reversed during data-taking to mitigate the differences between reconstruction efficiencies of oppositely charged particles.

The online event selection is performed by a trigger, which consists of a hardware stage followed by a two-level software stage. In between the two software stages, an alignment and calibration of the detector is performed in near real-time and their results are used in the trigger [38]. Events with candidate $D_{(s)}^+$ decays are selected by the hardware trigger by imposing either that one or more $D_{(s)}^+$ decay products are associated with large transverse energy deposits in the calorimeter or that the accept decision is independent of the $D_{(s)}^+$ decay products (i.e., it is caused by other particles in the event). In the first level of the software trigger, one or more $D_{(s)}^+$ decay products must have large transverse momentum and be inconsistent with originating from any primary vertex. In the second level, the candidate decays are fully reconstructed using kinematic, topological and particle-identification criteria. The $D_{(s)}^+ \rightarrow K_S^0 h^+$ candidates are made by combining charged hadrons with $K_S^0 \rightarrow \pi^+ \pi^-$ candidates that decay early enough for the final-state pions to be reconstructed in the vertex detector. This requirement suppresses to a negligible level possible CP -violation effects due to interference between Cabibbo-favored and doubly Cabibbo-suppressed amplitudes with neutral-kaon mixing in the control-sample decays $D^+ \rightarrow K_S^0 \pi^+$ and $D_s^+ \rightarrow K_S^0 K^+$ [35].

The $D_{(s)}^+$ candidates reconstructed in the trigger are used directly in the offline analysis [39,40]. The candidates with a K_S^0 meson in the final state are further selected offline using an artificial neural network (NN), based on the multilayer perceptron algorithm [41], to suppress background due to random combinations of K_S^0 mesons and hadrons not originating from a $D_{(s)}^+ \rightarrow K_S^0 h^+$ decay. The quantities used in the NN to discriminate signal from combinatorial background are the K_S^0 candidate momentum, the transverse momenta of the $D_{(s)}^+$ candidate and of the companion hadron, the angle between the $D_{(s)}^+$ candidate momentum

and the vector connecting the primary and secondary vertices, the quality of the secondary vertex, and the track quality of the companion hadron. The NN is trained using signal and background data samples, obtained with the *sPlot* method [42], from a $\mathcal{O}(1\%)$ fraction of candidates randomly sampled. In the $D_s^+ \rightarrow K_S^0 \pi^+$ case, thanks to similar kinematics, background-subtracted $D^+ \rightarrow K_S^0 \pi^+$ decays are exploited as a signal proxy to profit from larger yields. The thresholds on the NN response are optimized for the $D_s^+ \rightarrow K_S^0 \pi^+$ and $D^+ \rightarrow K_S^0 K^+$ decays by maximizing the value of $S/\sqrt{S+B}$, where S and B stands for the signal and background yield observed in the mass ranges $1.93 < m(K_S^0 \pi^+) < 2.01$ GeV/ c^2 and $1.83 < m(K_S^0 K^+) < 1.91$ GeV/ c^2 , respectively. Candidate $D_{(s)}^+ \rightarrow \phi(\rightarrow K^+ K^-) \pi^+$ decays are selected offline with requirements on the transverse momenta of the $D_{(s)}^+$ candidate and of the companion hadron, on the quality of the secondary vertex, and on the $K^+ K^-$ mass to be within 10 MeV/ c^2 of the nominal $\phi(1020)$ -meson mass [32]. The mass window is chosen considering that the observed width is dominated by the $\phi(1020)$ -meson natural width of 4.2 MeV/ c^2 [32] and is only marginally affected by the experimental resolution of 1.3 MeV/ c^2 .

The contribution of $D_{(s)}^+$ mesons produced through decays of b hadrons, referred to as secondaries throughout, is suppressed by requiring that the $D_{(s)}^+$ impact parameter in the plane transverse to the beam (TIP) is smaller than 40 μm . The remaining percent-level contribution is evaluated by means of a fit to the TIP distribution when such requirement is released, as shown in Fig. 1 for the $D_s^+ \rightarrow K_S^0 \pi^+$ decay. The impact of the secondary background on the results is accounted for in the systematic uncertainties.

Typical sources of background from $D_{(s)}^+$ meson and Λ_c^+ baryon decays are the $D_s^+ \rightarrow K_S^0 K^+$ and $\Lambda_c^+ \rightarrow K_S^0 p$ decays, where the kaon and the proton are misidentified as a pion, when the signal is the $D_s^+ \rightarrow K_S^0 \pi^+$ decay, the

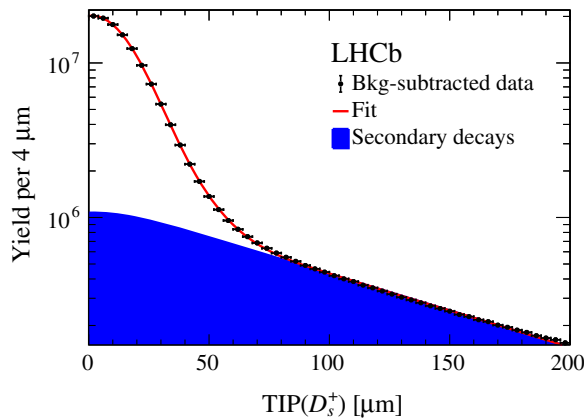


FIG. 1. Distribution of the transverse impact parameter (TIP) for background-subtracted $D_s^+ \rightarrow K_S^0 \pi^+$ candidates with fit projections overlaid.

$D^+ \rightarrow K_S^0 \pi^+$ and $\Lambda_c^+ \rightarrow K_S^0 p$ decays, where the pion and the proton are misidentified as a kaon, in the $D^+ \rightarrow K_S^0 K^+$ case, and the $\Lambda_c^+ \rightarrow \phi p$ decay, where the proton is misidentified as a pion, when the signal is the $D^+ \rightarrow \phi \pi^+$ decay. These are all reduced to a negligible level using particle-identification requirements and kinematic vetos.

Fiducial requirements are imposed to exclude kinematic regions that induce a large asymmetry in the companion-hadron reconstruction efficiency. These regions occur because low momentum particles of one charge at large (small) angles in the bending plane may be deflected out of the detector acceptance (into the noninstrumented beam pipe region), whereas particles with the other charge are more likely to remain within the acceptance. About 78%, 93%, and 94% of the selected candidates are retained by these fiducial requirements for $D_{(s)}^+ \rightarrow K_S^0 \pi^+$, $D_{(s)}^+ \rightarrow K_S^0 K^+$, and $D_{(s)}^+ \rightarrow \phi \pi^+$ decays, respectively.

Detection and production asymmetries may depend on the kinematics of the involved particles. Therefore, the cancellation provided by the control decays is accurate only if the kinematic distributions agree between any pair of signal and control modes, or pair of control modes entering Eqs. (4)–(6). Differences are observed, and the ratio between background-subtracted [42] signal and control sample distributions of transverse momentum, azimuthal angle and pseudorapidity are used to define candidate-by-candidate weights. The background-subtracted candidates of the control decays are weighted such that their distributions agree with those of the signal using an iterative procedure. The process consists of calculating the weights in each one-dimensional distribution of the weighting variables and repeating the procedure until good agreement is achieved among all the distributions. For the measurements of the $D_s^+ \rightarrow K_S^0 \pi^+$ and $D^+ \rightarrow \phi \pi^+$ CP asymmetries, the $D_s^+ \rightarrow \phi \pi^+$ and $D^+ \rightarrow K_S^0 \pi^+$ control samples are weighted so that the $D_{(s)}^+$ meson and companion-pion kinematic distributions agree with their respective signal samples to cancel the $D_{(s)}^+$ production and companion-pion detection asymmetries. In the case of the $\mathcal{A}_{CP}(D^+ \rightarrow K_S^0 K^+)$ measurement, the D^+ kinematic distributions of the $D^+ \rightarrow K_S^0 \pi^+$ sample are weighted to those of the $D^+ \rightarrow K_S^0 K^+$ signal to cancel the D^+ production asymmetry, and the K^+ distributions of the $D_s^+ \rightarrow K_S^0 K^+$ decays are weighted to those of the $D^+ \rightarrow K_S^0 K^+$ signal to cancel the kaon detection asymmetry. The $D^+ \rightarrow K_S^0 \pi^+$ and $D_s^+ \rightarrow K_S^0 K^+$ control decays then introduce their own additional nuisance asymmetries, which need to be corrected for using the $D_s^+ \rightarrow \phi \pi^+$ control decay. Hence, the D_s^+ and companion-pion kinematic distributions of the $D_s^+ \rightarrow \phi \pi^+$ sample are made to agree with those of the $D_s^+ \rightarrow K_S^0 K^+$ and $D^+ \rightarrow K_S^0 \pi^+$ samples, respectively, to cancel the D_s^+ production and companion-pion detection asymmetries.

Simultaneous least-squares fits to the mass distributions of weighted $D_{(s)}^+$ and $D_{(s)}^-$ candidates determine the raw asymmetries for each decay mode considered. To avoid experimenter bias, the raw asymmetries of the Cabibbo-suppressed signals were shifted by unknown offsets sampled uniformly between -1% and 1% , such that the results remained blind until the analysis procedure was finalized. In the fits, the signal and control decays are modeled as the

sum of a Gaussian function to describe the core of the peaks, and a Johnson S_U distribution [43], which accounts for the asymmetric tails. The combinatorial background is described by the sum of two exponential functions. All shape parameters are determined from the data. In each fit, signal and control decays share the same shape parameters apart from a mass shift, which accounts for the known difference between the D_s^+ and D^+ masses [32], and a relative scale factor between the peak widths, which is also determined from the data. The means and widths of the peaks, as well as all background shape parameters, are allowed to differ between $D_{(s)}^+$ and $D_{(s)}^-$ decays. The projections of the fits to the combined $D_{(s)}^+$ and $D_{(s)}^-$ data are shown in Fig. 2. The samples contain approximately 600 thousand $D_s^+ \rightarrow K_S^0 \pi^+$, 5.1 million $D^+ \rightarrow K_S^0 K^+$, and 53.3 million $D^+ \rightarrow \phi \pi^+$ signal candidates, together with approximately 30.5 million $D^+ \rightarrow K_S^0 \pi^+$, 6.5 million $D_s^+ \rightarrow K_S^0 K^+$, and 107 million $D_s^+ \rightarrow \phi \pi^+$ control decays.

The raw asymmetries are, where relevant, corrected for the neutral-kaon detection asymmetry. The net correction is estimated following Ref. [44] to be $(+0.084 \pm 0.005)\%$ for $\mathcal{A}_{CP}(D_s^+ \rightarrow K_S^0 \pi^+)$, $(-0.086 \pm 0.005)\%$ for $\mathcal{A}_{CP}(D^+ \rightarrow K_S^0 K^+)$, and $(-0.068 \pm 0.004)\%$ for $\mathcal{A}_{CP}(D^+ \rightarrow \phi \pi^+)$, where the uncertainty is dominated by the accuracy of the detector modeling in the simulation. The asymmetries are combined following Eqs. (4)–(6) to obtain $\mathcal{A}_{CP}(D_s^+ \rightarrow K_S^0 \pi^+) = (1.3 \pm 1.9) \times 10^{-3}$, $\mathcal{A}_{CP}(D^+ \rightarrow K_S^0 K^+) = (-0.09 \pm 0.65) \times 10^{-3}$, $\mathcal{A}_{CP}(D^+ \rightarrow \phi \pi^+) = (0.05 \pm 0.42) \times 10^{-3}$, where the uncertainties are only statistical.

Several sources of systematic uncertainty affecting the measurement are considered as reported in Table I. The dominant contribution is due to the assumed shapes in the mass fits. This is evaluated by fitting with the default model large sets of pseudoexperiments where alternative models that describe data equally well are used in generation. For $\mathcal{A}_{CP}(D_s^+ \rightarrow K_S^0 \pi^+)$ and $\mathcal{A}_{CP}(D^+ \rightarrow K_S^0 K^+)$, the second leading contribution is due to the residual

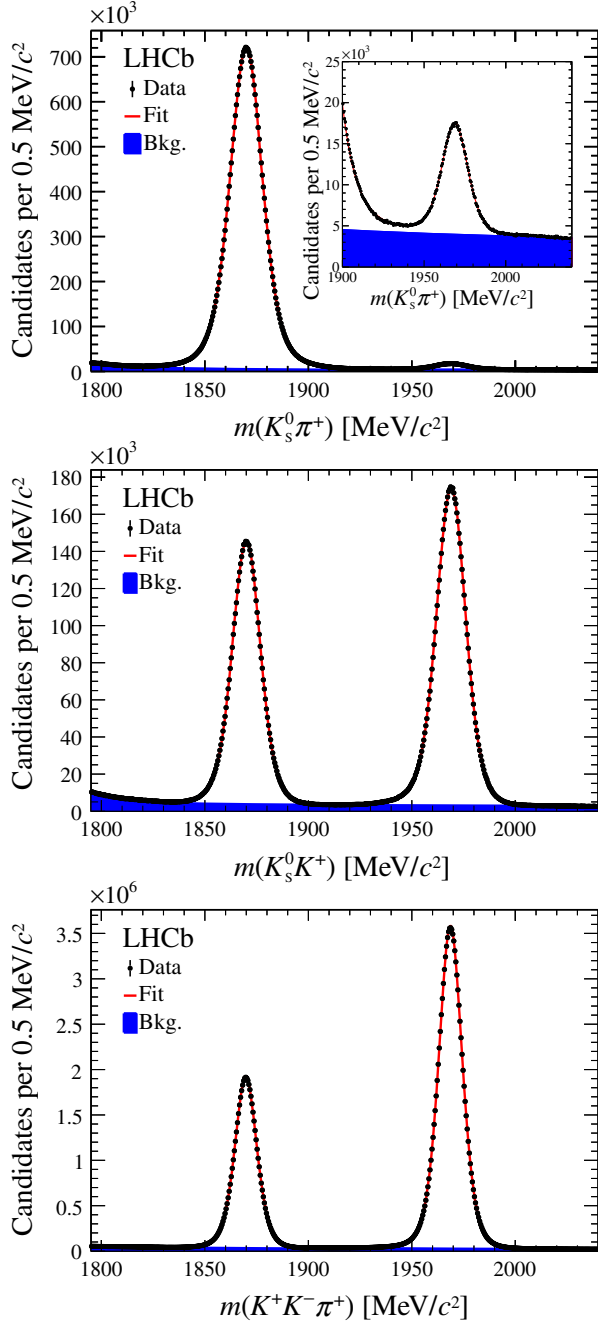


FIG. 2. Mass distributions of the selected (top) $D_{(s)}^+ \rightarrow K_S^0 \pi^+$, (middle) $D_{(s)}^+ \rightarrow K_S^0 K^+$, and (bottom) $D_{(s)}^+ \rightarrow \phi \pi^+$ candidates with fit projections overlaid. The inset in the top plot shows the mass distribution around the $D_s^+ \rightarrow K_S^0 \pi^+$ signal region.

TABLE I. Summary of the systematic uncertainties (in units of 10^{-3}) on the measured quantities. The total is the sum in quadrature of the different contributions.

Source	$\mathcal{A}_{CP}(D_s^+ \rightarrow K_S^0 \pi^+)$	$\mathcal{A}_{CP}(D^+ \rightarrow K_S^0 K^+)$	$\mathcal{A}_{CP}(D^+ \rightarrow \phi \pi^+)$
Fit model	0.39	0.44	0.24
Secondary decays	0.30	0.12	0.03
Kinematic differences	0.09	0.09	0.04
Neutral kaon asymmetry	0.05	0.05	0.04
Charged kaon asymmetry	0.08	0.09	0.15
Total	0.51	0.48	0.29

contamination from secondary $D_{(s)}^+$ decays, which introduces a small difference between the asymmetry of $D_{(s)}^+$ -meson production cross sections of the signal and control modes. For $\mathcal{A}_{CP}(D^+ \rightarrow \phi\pi^+)$, instead, the second leading systematic uncertainty arises from neglected kinematic differences between the ϕ -meson decay products. These differences, mainly caused by the interference between the S wave and $\phi\pi^+$ decay amplitudes in the K^+K^- -mass region under study, result in an imperfect cancelation of the charged-kaon detection asymmetry. Other subleading contributions are due to the inaccuracy in the equalization of the kinematic distributions between signal and control samples, and to the uncertainty in the neutral-kaon detection asymmetry.

In addition, several consistency checks are performed to investigate possible unexpected biases by comparing results obtained in subsamples of the data defined according to the data-taking year and magnetic-field polarity, the per-event track multiplicity, the configurations of the hardware- and software-level triggers, and the $D_{(s)}^+$ momentum. A χ^2 test has been performed for each cross-check and the corresponding p values are consistent with being uniformly distributed; the lowest (largest) p value is 4% (86%). Therefore, the observed variations in results are consistent with statistical fluctuations and no additional sources of systematic uncertainties are considered.

In summary, using proton-proton collision data collected with the LHCb detector at a center-of-mass energy of 13 TeV, and corresponding to 3.8 fb^{-1} of integrated luminosity, the following CP asymmetries are measured:

$$\begin{aligned}\mathcal{A}_{CP}(D_s^+ \rightarrow K_S^0\pi^+) &= (1.3 \pm 1.9 \pm 0.5) \times 10^{-3}, \\ \mathcal{A}_{CP}(D^+ \rightarrow K_S^0K^+) &= (-0.09 \pm 0.65 \pm 0.48) \times 10^{-3}, \\ \mathcal{A}_{CP}(D^+ \rightarrow \phi\pi^+) &= (0.05 \pm 0.42 \pm 0.29) \times 10^{-3},\end{aligned}$$

where the first uncertainties are statistical and the second systematic. Effects induced by CP violation in the neutral kaon system are subtracted from the measured asymmetries. The results represent the most precise determination of these quantities to date and are consistent with CP symmetry. They are in agreement with previous LHCb determinations based on independent data samples collected at center-of-mass energies of 7 and 8 TeV [29,30], as well as with measurements from other experiments [23–28]. The results are combined with previous LHCb measurements using the BLUE method [45]. The systematic uncertainties are considered uncorrelated, apart from those due to the neutral- and charged-kaon detection asymmetries that are fully correlated. The combination yields

$$\begin{aligned}\mathcal{A}_{CP}(D_s^+ \rightarrow K_S^0\pi^+) &= (1.6 \pm 1.7 \pm 0.5) \times 10^{-3}, \\ \mathcal{A}_{CP}(D^+ \rightarrow K_S^0K^+) &= (-0.04 \pm 0.61 \pm 0.45) \times 10^{-3}, \\ \mathcal{A}_{CP}(D^+ \rightarrow \phi\pi^+) &= (0.03 \pm 0.40 \pm 0.29) \times 10^{-3},\end{aligned}$$

where the first uncertainties are statistical and the second systematic. No evidence for CP violation in these decays is found. More precise measurements of these asymmetries can be expected when the data already collected by LHCb in 2018 are included in a future analysis, and when much larger samples will become available at the upgraded LHCb detector [46].

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