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Search for diboson resonances in hadronic final states in 139 fb^{-1} of pp collisions at $\sqrt{s} = 13 \text{ TeV}$ with the ATLAS detector

The ATLAS Collaboration

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Search for diboson resonances in hadronic final states in 139 fb^{-1} of pp collisions at $\sqrt{s} = 13 \text{ TeV}$ with the ATLAS detector



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ABSTRACT: Narrow resonances decaying into WW , WZ or ZZ boson pairs are searched for in 139 fb^{-1} of proton-proton collision data at a centre-of-mass energy of $\sqrt{s} = 13 \text{ TeV}$ recorded with the ATLAS detector at the Large Hadron Collider from 2015 to 2018. The diboson system is reconstructed using pairs of high transverse momentum, large-radius jets. These jets are built from a combination of calorimeter- and tracker-inputs compatible with the hadronic decay of a boosted W or Z boson, using jet mass and substructure properties. The search is performed for diboson resonances with masses greater than 1.3 TeV . No significant deviations from the background expectations are observed. Exclusion limits at the 95% confidence level are set on the production cross-section times branching ratio into dibosons for resonances in a range of theories beyond the Standard Model, with the highest excluded mass of a new gauge boson at 3.8 TeV in the context of mass-degenerate resonances that couple predominantly to gauge bosons.

KEYWORDS: Hadron-Hadron scattering (experiments)

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1 Introduction

The discovery of new phenomena in high-energy proton-proton (pp) collisions is one of the main goals of the Large Hadron Collider (LHC). New heavy, TeV-scale, resonances of vector bosons VV (where V represents a W or a Z boson) are a possible signature of such new physics and are predicted in several extensions to the Standard Model (SM). These include extended gauge-symmetry models [1–3], Grand Unified theories [4–7], theories with warped extra dimensions [8–12], two-Higgs-doublet models [13], little-Higgs models [14], theories with new strong dynamics [15], including technicolour [16–18], and more generic composite Higgs models [19]. The data sample of 36.7 fb^{-1} of pp collisions collected in 2015 and 2016 at the LHC at $\sqrt{s} = 13 \text{ TeV}$ offered improved sensitivity to heavy diboson resonances compared with earlier results. The ATLAS and CMS collaborations performed searches in the fully hadronic final states using this data [20–22] but no significant deviation

from a smooth background consistent with the SM expectation was observed. Searches by ATLAS [23, 24] and CMS [25, 26] for semileptonic decay modes of the boson pair, as well as statistical combinations of various decay channels [27], on the same data also did not reveal any hint of new physics.

This paper presents a search for narrow diboson resonances decaying into fully hadronic final states in 139 fb^{-1} of pp collision data collected by the ATLAS experiment between 2015 and 2018. The W and Z bosons produced in the decay of TeV-scale resonances are highly boosted, and are therefore reconstructed in ATLAS as a single large-radius-parameter jet. The signature of such heavy resonance decays is thus a resonant structure in the dijet invariant mass spectrum. Although the hadronic decays of vector bosons have the largest branching ratio (67% for W and 70% for Z bosons), they suffer from background contamination from the production of multijet events. This background is larger by several orders of magnitude, and to suppress it, the characteristic jet substructure of W/Z boson decays is used. Contributions to the background from SM processes containing bosons, $V + \text{jets}$, SM VV , $t\bar{t}$ and single top production, are significantly smaller.

To improve the sensitivity of this search, new techniques are used. Novel inputs are used for jet finding, which improve the jet substructure resolution of ATLAS in highly boosted topologies [28]. To further benefit from these developments, a new approach for identifying boosted boson candidates is introduced. The identification of the boosted boson candidates is validated using the known SM $V + \text{jets}$ production.

To avoid limitations caused by poor modelling or limited numbers of Monte Carlo (MC) generated background events, the observed background is characterised by a parametric function fit to the smoothly falling dijet invariant mass distribution. To assess the sensitivity of the search, to optimise the event selection and for comparison with the observed data, three specific benchmark models are used: a spin-0 radion [29] decaying into WW or ZZ ; a spin-1 Heavy Vector Triplet (HVT) Model [30] that provides signals such as $W' \rightarrow WZ$ and $Z' \rightarrow WW$; and a spin-2 graviton $G_{\text{KK}} \rightarrow WW$ or ZZ , corresponding to Kaluza-Klein (KK) modes [8, 9] of the Randall-Sundrum (RS) graviton [10–12]. These models assume production mechanisms either via gluon-gluon fusion or quark-antiquark annihilation.

2 ATLAS detector

The ATLAS detector [31] surrounds nearly the entire solid angle around the ATLAS collision point. It has an approximately cylindrical geometry¹ and consists of an inner tracking detector surrounded by electromagnetic and hadronic calorimeters and a muon spectrometer. The tracking detector is placed within a 2 T axial magnetic field provided by a superconducting solenoid and measures charged-particle trajectories with silicon pixel and

¹ATLAS uses a right-handed coordinate system with its origin at the nominal interaction point (IP) in the centre of the detector and the z -axis along the beam pipe. The x -axis points from the IP to the centre of the LHC ring, and the y -axis points upward. Cylindrical coordinates (r, ϕ) are used in the transverse plane, ϕ being the azimuthal angle around the beam pipe. The pseudorapidity is defined in terms of the polar angle θ as $\eta = -\ln \tan(\theta/2)$. Angular distance is measured in units of $\Delta R = \sqrt{(\Delta\eta)^2 + (\Delta\phi)^2}$.

silicon microstrip detectors that cover the pseudorapidity range $|\eta| < 2.5$, and with a straw-tube transition radiation tracker covering $|\eta| < 2.0$. A new innermost pixel layer [32, 33] inserted at a radius of 3.3 cm has been used since 2015.

Electromagnetic and hadronic calorimeter systems provide energy measurements with high granularity. The electromagnetic calorimeter is a liquid-argon (LAr) sampling calorimeter with lead absorbers, spanning $|\eta| < 3.2$ with barrel and endcap sections. The three-layer central hadronic calorimeter comprises scintillator tiles with steel absorbers and extends to $|\eta| = 1.7$. The hadronic endcap calorimeters measure particles in the region $1.5 < |\eta| < 3.2$ using liquid argon with copper absorbers. The forward calorimeters cover $3.1 < |\eta| < 4.9$, using LAr/copper modules for electromagnetic energy measurements and LAr/tungsten modules to measure hadronic energy.

The muon spectrometer surrounds the calorimetry system and provides precision muon tracking and triggering. It includes three large superconducting air-core toroids providing a magnetic field for accurate momentum measurements in tracking drift chambers arranged in a barrel, covering $|\eta| < 1.0$, and endcaps, extending to $|\eta| = 2.7$.

Events are recorded in ATLAS if they satisfy a two-level trigger requirement [34]. The level-1 trigger detects jet and particle signatures in the calorimeter and muon systems with a fixed latency of $2.5 \mu\text{s}$, and is designed to reduce the event rate to about 100 kHz. Jets are identified at level-1 with a sliding-window algorithm, searching for local maxima in square regions with size $\Delta\eta \times \Delta\phi = 0.8 \times 0.8$. The subsequent high-level trigger consists of software-based trigger filters that reduce the event rate to one kHz.

3 Data

The search is performed using data collected by the ATLAS experiment from 2015 to 2018 from $\sqrt{s} = 13 \text{ TeV}$ LHC pp collisions. Events used in this search satisfied a single-jet trigger requiring at least one jet reconstructed at each trigger level. The final filter in the high-level trigger required a jet to satisfy a high transverse momentum (p_T) threshold, $p_T \geq 360 \text{ GeV}$ (2015), $p_T \geq 420 \text{ GeV}$ (2016), $p_T \geq 440 \text{ GeV}$ (2017 and 2018), reconstructed with the anti- k_t algorithm [35] and a large radius parameter ($R = 1.0$). Calorimeter-cell energy clusters calibrated to the hadronic scale utilising the local cell signal weighting method [36] were used as inputs. After requiring that the data were collected during stable beam conditions and the detector components relevant to the analysis were functional, the integrated luminosity was 3.2 fb^{-1} in 2015, 33.0 fb^{-1} in 2016, 44.3 fb^{-1} in 2017 and 58.5 fb^{-1} in 2018.

4 Simulation

4.1 Signal models

MC simulation of signal events is used to optimise the sensitivity of the search and to interpret its results. Signals are simulated in three benchmark scenarios.

In the first scenario, the gravitational fluctuations in the extra dimension of the Randall-Sundrum framework correspond to scalar fields, known as the radion, which are massless in the simplest scenario. A fundamental problem in the original Randall-Sundrum

framework is that it lacks a mechanism to stabilise the radius of the compactified extra dimension, r_c . One possible mechanism to achieve this is to introduce an additional bulk scalar radion, produced via gluon-gluon fusion, which has its interactions localised on the two ends of the extra dimension [37, 38]. This causes the radion field to acquire a mass term, which is typically much smaller than the first KK excitation mass. The coupling of the radion field to SM fields scales inversely proportional to the model parameter $\Lambda_R = \sqrt{g} \times k \times e^{-k\pi r_c} \sqrt{M_5^3/k^3}$ where M_5 is the 5-dimensional Planck mass, which has been extensively studied in the literature [29, 39, 40, 40], k the curvature factor, and g is the 5-dimensional metric. The size of the extra dimension, defined as $k\pi r_c$, is another parameter of the model. In this analysis, the curvature factor is set to $k\pi r_c = 35$, and $\Lambda_R = 3 \text{ TeV}$ is used.

The couplings of the radion to fermions are proportional to the masses of the fermions, while the couplings are proportional to the square of the masses for bosons. For radion mass above $\sim 1 \text{ TeV}$, the dominant decay mode is into pairs of bosons. The decay width of the radion is approximately 10% of its pole mass, resulting in observable mass peaks with a width comparable to the experimental resolution (see section 6.3). The calculated production cross-section times branching ratio ($\sigma \times \mathcal{B}$) for a radion decaying into WW , with the W decaying hadronically, is 2.75 fb and 0.26 fb for radion masses of 2 TeV and 3 TeV, respectively. Corresponding values for a radion decaying into ZZ are 1.53 fb and 0.15 fb.

The second scenario is based on two benchmark models, A and B, of the HVT phenomenological Lagrangian [30]. The Lagrangian introduces a new heavy vector triplet V' , where V' refers to W' and Z' , produced via quark-antiquark annihilation, whose members are degenerate in mass, and parameterises its couplings to SM fields in a generic manner, such that a large class of extensions to the SM can be described.

Model A with the strength of the vector-boson interaction $g_V = 1$ [30] describes scenarios where the new triplet field couples weakly to the SM fields and arises from an extension of the SM gauge group, with the heavy vector bosons having comparable branching ratios into fermions and gauge bosons. For W' and Z' masses of interest, the width of the new heavy bosons is approximately 2.5%, which results in observable mass peaks with a width dominated by the experimental resolution. The branching fraction of the new heavy boson W' (Z') to each of the WZ and WH (WW and ZH) final states, where H represents the Higgs boson, is approximately 2%. The calculated $\sigma \times \mathcal{B}$ values for $W' \rightarrow WZ$, with W and Z bosons decaying hadronically, are 8.3 fb and 0.75 fb for W' masses of 2 TeV and 3 TeV, respectively. Corresponding values for $Z' \rightarrow WW$ are 3.8 fb and 0.34 fb.

Model B with $g_V = 3$ is representative of composite Higgs models, where the fermionic couplings to V' are suppressed. For the W' and Z' masses of interest, the branching fraction of the new heavy boson W' (Z') to each of the WZ and WH (WW and ZH) final states is close to 50%. Resonance widths and experimental signatures are similar to those obtained for model A and the predicted $\sigma \times \mathcal{B}$ values for $W' \rightarrow WZ$, with hadronic W and Z decays, are 13 fb and 1.3 fb for W' masses of 2 TeV and 3 TeV, respectively. Corresponding values for $Z' \rightarrow WW$ are 6.0 fb and 0.55 fb.

The third scenario considered is the bulk RS model [10] that extends the original RS model [8, 41] with a warped extra dimension, by allowing the SM fields to propagate in

the bulk of the extra dimension. This model is characterised by a dimensionless coupling constant $\kappa/\overline{M}_{\text{Pl}} \sim 1$, where κ is the curvature of the warped extra dimension, and \overline{M}_{Pl} is the reduced Planck mass. In this model, a Kaluza-Klein graviton, G_{KK} , predominately produced via gluon-gluon fusion, decays into pairs of top quarks, pairs of Higgs bosons, WW and ZZ with significant branching fractions. The branching fraction of the G_{KK} to WW (ZZ) ranges from 24% to 20% (12% to 10%) as the mass increases. The decay width of the G_{KK} is approximately 6% of its pole mass, resulting in observable mass peaks with a width comparable to the experimental resolution, and $\sigma \times \mathcal{B}$ for $G_{\text{KK}} \rightarrow WW$, with the W decaying hadronically, is 1.29 fb and 0.06 fb for G_{KK} masses of 2 TeV and 3 TeV, respectively. Corresponding values for $G_{\text{KK}} \rightarrow ZZ$ are 0.65 fb and 0.03 fb.

4.2 Simulated event samples

For all MC samples, all hadronic decays were imposed at the generator level. MC samples for the radion, HVT, and RS models, were generated using MADGRAPH 2.2.2 [42] interfaced to PYTHIA 8.186 [43] for hadronisation using the leading-order (LO) NNPDF 2.3 parton distribution function (PDF) set [44] and the ATLAS A14 set of tuned parameters for the underlying event [45]. In all signal samples, the W and Z bosons are primarily longitudinally polarised. The procedure to derive the optimal boson identification criteria (section 6.1) uses a dedicated sample of W' decaying only into W/Z bosons that in turn decay hadronically. PYTHIA 8.186 was used to generate this sample with the A14 set of tuned parameters for the underlying event and the NNPDF 2.3 LO PDF. The cross-section of the hard-scattering process was modified by applying an event-by-event weighting factor to broaden the width of the resonance and widen the p_{T} distribution of the electroweak bosons produced in its hadronic decays.

PYTHIA 8.186 with the NNPDF 2.3 LO PDF set and the A14 set of tuned parameters was used to generate and shower multijet background events. Samples of W +jets and Z +jets events were generated with SHERPA 2.2.5 [46–49] interfaced with the NNPDF 3.0 next-to-next-to-leading-order (NNLO) PDF set [50]. A $t\bar{t}$ sample generated with POWHEG-BOX v2 [51–53] with the NNPDF 3.0 next-to-leading-order (NLO) PDF [54], interfaced with PYTHIA 8.186 with the NNPDF 2.3 LO PDF and the A14 set of tuned parameters for parton showering is used for the V +jets study. EvtGen v1.2.0 [55] was used for properties of bottom and charm hadron decays, except for samples generated by SHERPA.

For all MC samples, the final-state particles produced by the generators were propagated through a detailed detector simulation based on GEANT4 [56, 57]. The mean number of pp interactions per bunch crossing, ‘pile-up’, was approximately 33 in the collision data being used for the analysis. The expected contribution from these minimum-bias pp interactions was accounted for by overlaying additional minimum-bias events generated with PYTHIA 8.186 using the ATLAS A3 [58] set of tuned parameter. The MC simulation events were weighted to match the distribution of the average number of interactions per bunch crossing observed in collision data. Simulated events were then reconstructed with the same algorithms as run on the collision data.

5 Reconstruction

The experimental signatures central to this analysis are hadronic jets. Since the decay products of TeV-scale resonances are highly boosted, their decay products become increasingly collimated and they are therefore reconstructed as a single large-radius jet. It is important that they can still be differentiated from multijet events where a jet is initiated by a single quark or gluon. This relies on both the energy and angular resolution of the detector used to reconstruct the jet. Although the analysis primarily relies on jets, reconstruction of lepton candidates is necessary to reject events that could bias the SM V +jets studies presented in section 6.2.

5.1 Track-CaloClusters

In previous analyses, ATLAS has mainly focused on the use of calorimeter-based jet substructure, which exploits the exceptional energy resolution of the ATLAS calorimetry [36]. However, as the event becomes even more energetic, jets become so collimated that the calorimeter lacks the angular resolution to resolve the desired structure inside the jet. For boson jets with high transverse momentum, p_T , only a handful of calorimeter-cell clusters are created, each with limited angular resolution, but excellent energy resolution. On the other hand, the tracking detector has excellent angular resolution and good reconstruction efficiency at very high energy [59], while its momentum resolution deteriorates. By combining information from the ATLAS calorimeter and tracking detectors, the precision of jet substructure techniques can be improved for a wide range of energies. This analysis uses a new unified object built from both the tracking and the calorimeter information, referred to as Track-CaloClusters (TCCs) [28]. This procedure is a type of particle flow, complementary to the energy subtraction algorithm described in the ATLAS particle flow paper [60] which improves the energy resolution of low-energy jets. The two algorithms are designed to improve the jet reconstruction performance in very different energy regimes, reflected in their distinct four-momentum construction and energy sharing procedures. Energy sharing in the TCC approach is based solely on a weighting scheme where only the relative track momenta are used to spatially redistribute the energy measured in the calorimeter. In practice, this means that the TCC algorithm uses the spatial coordinates of the tracker and the energy scale of the calorimeter. A more detailed description of TCCs can be found in ref. [28].

5.2 Jet reconstruction

This analysis uses anti- k_t , $R = 1.0$ jets reconstructed from both the combined and the neutral TCCs. Combined TCCs are four-momenta created by combining the angular information of tracks with the energy information of the calorimeters. Neutral TCCs are calorimeter topo-clusters that could not be matched to any track, most likely representing energy deposits from neutral particles. The use of combined and neutral TCCs captures most of the hard-scatter energy and provides the best representation of the total energy flow in the event, as there are both charged and neutral contributions. The combined TCC component is robust against effects from pile-up since only tracks consistent with coming

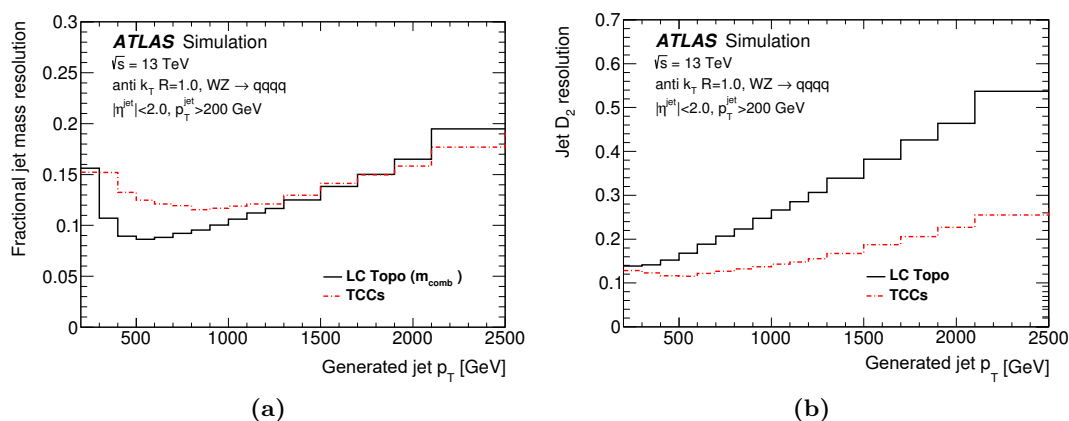


Figure 1. A comparison of (a) the fractional jet mass resolution for jets built from a linear combination of the calorimeter and track-only mass (LC Topo m_{comb} , solid line), and jets built using combined and neutral Track-CaloClusters objects (dashed lines) as a function of Monte Carlo generator-level jet p_T . The fractional jet resolution of the D_2 variable (b) is compared between Track-CaloClusters and pure calorimeter jets. Only the two jets with the highest p_T per event matched to a generated jet from a W or Z boson are shown.

from the primary vertex² are used. However, by including the neutral TCCs, these jets have a pile-up dependence similar to that of standard topo-cluster jets. Jets are therefore trimmed [61] to remove contributions from pile-up by removing any $R = 0.2$ subjet with less than 5% of the p_T of the associated $R = 1.0$ jet. The clustering and trimming algorithms use the FastJet package [62]. The combination of pile-up suppression through track-to-primary-vertex matching and trimming makes these jets very robust against pile-up [28]. A MC-based particle-level energy and mass calibration is applied to the jets, as described in ref. [63]. MC generator-level jets are built using the same algorithm and trimming procedure, with inputs of stable generator-level particles ($c\tau > 10$ mm) excluding muons and neutrinos, and excluding particles from pile-up. These serve as the reference in figure 1. The energy and mass of MC generator-level jets also serves as reference to which reconstructed detector-level jets are corrected to in the above mentioned calibration procedure. Consequently, the mass of jets from V bosons is not expected to match the pole mass of the bosons.

Several jet properties can be used to discriminate hadronic decays of W and Z bosons from background jets. Two providing strong discrimination are the jet mass and D_2 ,³ where the latter is defined as a ratio of two-point to three-point energy correlation functions that are based on the energies of the jets' constituents and their pairwise angular separations [64]. Signal jets are expected to peak at D_2 values below one, while jets from multijet background have significantly larger values. The radiation of a hard gluon can allow background jets to mimic a two-pronged structure and satisfy the tagging require-

²If more than one vertex is reconstructed, the one with the highest sum of p_T^2 of the associated tracks is regarded as the primary vertex.

³The angular exponent β , defined in ref. [64], is set to unity.

ments described above. Discrimination between boson jets and multijet background from such gluon-initiated jets can be attained by selecting on the charged hadron multiplicity, in form of the track multiplicity (n_{trk}) of the untrimmed $R = 1.0$ jet, considering tracks with $p_T > 0.5$ GeV consistent with coming from the primary vertex.

Figure 1 shows the striking improvement in D_2 resolutions⁴ achieved with TCC jets. The mass resolution is superior to previously used jet mass variables (m_{comb} [65]) starting around a jet p_T of 2 TeV. Below 1 TeV, the mass resolution in TCC jets is degraded. For identifying hadronically decaying V bosons, the improvement in D_2 resolution far outweighs the slight degradation in mass resolution.

5.3 Leptons

Electron identification is based on matching tracks to energy clusters in the electromagnetic calorimeter and calculating a likelihood based on several properties of the electron candidate. Electrons are required to have $p_T > 25$ GeV and $|\eta| < 2.5$, and to satisfy the ‘medium’ identification criterion [66] and the ‘loose’ track-based isolation [66].

Muon identification relies on matching tracks in the inner detector to muon spectrometer tracks or track segments. Muons are required to have $p_T > 25$ GeV and $|\eta| < 2.5$, and to satisfy the ‘loose’ selection criterion [67] and the ‘loose’ track isolation [67].

6 Event selection

To avoid contamination from non-collision backgrounds such as from calorimeter noise, beam halo, and cosmic rays, events containing an anti- k_t jet built from calorimeter-cell clusters with $R = 0.4$ and $p_T > 20$ GeV failing to meet the loose quality criteria for consistency with production in pp collisions are rejected [68]. In addition, events with at least one lepton meeting the requirements defined in section 5.3 are rejected. There are no further requirements on leptons that are aligned with jets.

Events are required to have at least two anti- k_t , $R = 1.0$, jets originating from the primary vertex, one with $p_T > 500$ GeV and the second with $p_T > 200$ GeV. The leading (highest p_T) and subleading of these jets must satisfy $|\eta| < 2.0$ (to guarantee a good overlap with the tracking acceptance), have masses $m_J > 50$ GeV, and their invariant mass, m_{JJ} , must be larger than 1.3 TeV. The last requirement ensures that the triggers in use are fully efficient for the backgrounds and the benchmark signals. These selections are referred to as pre-selections.

The pair of jets is then required to have a small separation in rapidity, $|\Delta y_{12}| < 1.2$. This requirement reduces the multijet background, which is mainly produced in t -channel processes with large rapidity differences, in contrast to signal events that are expected to be produced in s -channel processes with small rapidity differences. Additionally, to reject events with potentially badly reconstructed jets, a criterion is applied to the p_T asymmetry,

⁴The resolution is defined as $R^r = [Q_{84}(\mathcal{R}^r) - Q_{16}(\mathcal{R}^r)] / [2 \times Q_{50}(\mathcal{R}^r)]$ and $R^d = 1/2 [Q_{75}(\mathcal{R}^d) - Q_{25}(\mathcal{R}^d)]$ for the mass and D_2 , respectively, where Q_x is the $x\%$ quantile boundary, meaning that Q_{50} is the median. The mass response is defined as $\mathcal{R}^r = m_{\text{reco}}/m_{\text{gen}}$, while the residual of D_2 is $\mathcal{R}^d = D_{2,\text{reco}} - D_{2,\text{gen}}$, where ‘gen’ and ‘reco’ refer to the generated and reconstructed properties of the jets.

$A = (p_{T1} - p_{T2}) / (p_{T1} + p_{T2}) < 0.15$, where p_{T1} and p_{T2} are the transverse momenta of the leading and subleading jets, respectively.

6.1 Vector-boson identification

Jet substructure can be exploited to enhance the separation between signal boson jets and jets from multijet background. Several promising variables have been studied [63], with the largest sensitivity gain coming from the use of the three variables introduced in section 5.2: jet mass, D_2 , and n_{trk} .

A three-dimensional (jet mass, D_2 , n_{trk}) tagger using TCC jets is optimised to provide maximum significance for boosted vector-boson jets relative to background jets. A measure of significance, independent of the cross-sections of the new processes being searched for, is selected: $\epsilon / (a/2 + \sqrt{B})$, where ϵ is the per-signal-jet selection efficiency for masses in the range of 0.5 TeV to 10 TeV in the W' model described in section 4.1, a is the number of standard deviations corresponding to a one-sided Gaussian distribution, and B is the number of background jets after the selection [69] taken from MC simulation. This number does not rely on a specific signal, but is valid for all signals with similar experimental features. Compared with the often used S/\sqrt{B} , that breaks down for small values of B , as is the case here, this measure is more appropriate. A value of $a = 3$ is used, where the result of the optimisation is not very sensitive to the exact value. For each jet p_T bin, the optimal selection on the three variables is defined by the combination of selections that leads to the highest significance. This simultaneous treatment properly accounts for the correlations between the variables. The result of this optimisation does not depend on the pre-selections described at the beginning of this section. Next, the applied selection criteria on jet mass, D_2 , and n_{trk} are parameterised with jet- p_T -dependent functions. For the jet mass, the function follows its approximate experimental resolution. For the latter two a simple higher-order polynomial is used. The resulting smooth selections for the W and Z boson taggers as a function of jet p_T are shown in figure 2. Jets selected in the analysis are required to have jet masses inside the jet mass window and D_2 and n_{trk} values below the shown values. It should be noted that the W and Z boson mass windows overlap. This search is not sensitive to signatures containing massive particles with masses different from W/Z boson masses (for example top quarks or Higgs bosons).

Unlike previous boson taggers [20], the optimisation described above does not enforce a fixed signal efficiency nor a fixed background rejection, but rather creates a smooth behaviour that maximises the analysis sensitivity. Figure 3 shows the resulting W and Z boson efficiencies and multijet background rejections (defined as $1/\text{efficiency}$) as a function of jet p_T . The selection criteria retain about 20% efficiency for W and Z boson jets with $p_T = 0.5$ TeV. At higher jet p_T , the signal efficiency increases, reaching an efficiency close to 60% for jets with a p_T of 4 TeV. This is due to the behaviour of the multijet background, which decreases rapidly for high dijet masses, and thus higher jet p_T . In the regime of $m_{JJ} > 3.0$ TeV, where the number of background events is small, the tagger maintains a reasonable acceptance for signals with small cross-sections. This is also reflected in the background rejection as a function of jet p_T .

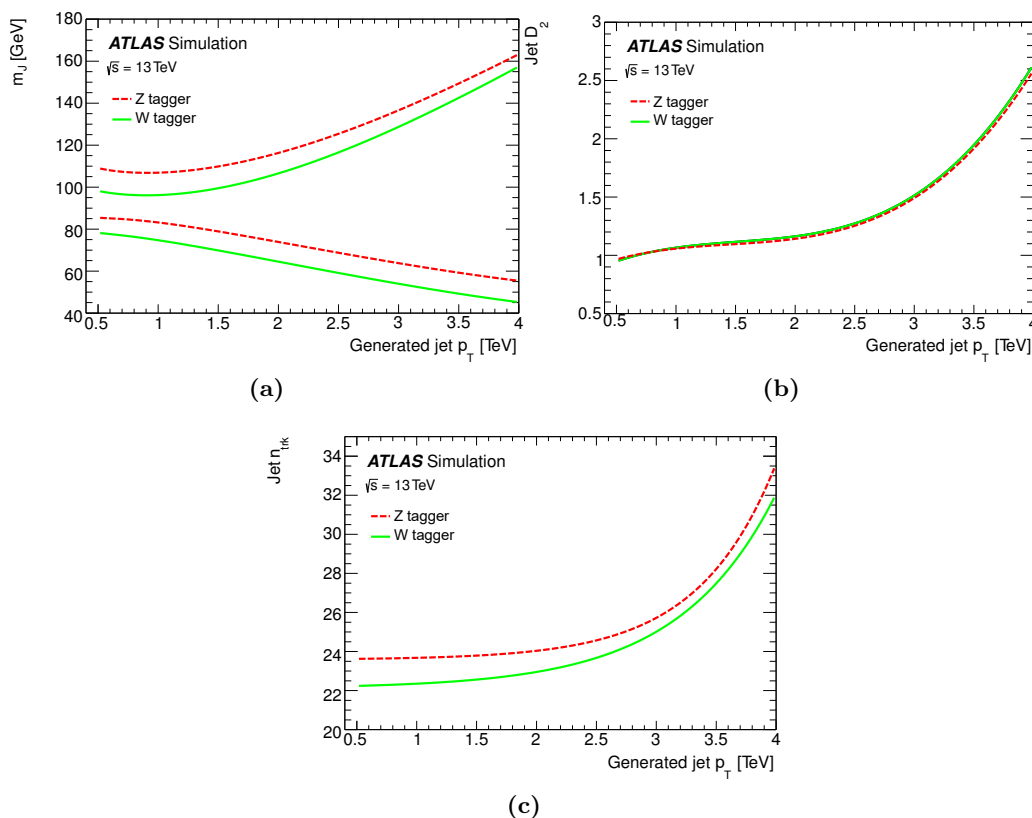


Figure 2. (a) Jet mass window, (b) D_2 selection and (c) n_{trk} selection of the W and Z taggers as a function of jet p_T . Jets selected in the analysis are required to have jet masses inside the jet mass window and D_2 and n_{trk} values below the shown values. The tagger is only valid for jets with a p_T between 0.5 TeV and 4.0 TeV and with $|\eta^{\text{jet}}| < 2.0$.

6.2 Measurement of boson-tagging efficiency

The modelling of the boson-tagging efficiency is evaluated in a data sample enriched in final states with a vector boson plus jets. This sample is obtained by requiring two large-radius jets with $|\eta| < 2.0$ and then requiring that the leading jet has $p_T > 600$ GeV. A higher minimum p_T requirement is imposed on the leading jet than in the nominal event selection to obtain a sample with higher average leading jet p_T that better corresponds to the jet p_T values probed in the search. Events with identified leptons are vetoed. Both jets are independently analysed for the presence of a vector boson by requiring them to satisfy the D_2 and n_{trk} selection for either a W or a Z boson. The opposite jet is required to not satisfy the same D_2 selection to guarantee independence of this control region from the main analysis signal region.

The mass distribution of the selected jets between 50 GeV and 200 GeV is fit by a signal-plus-background function, allowing the inclusive rate of $V + \text{jets}$ events to be measured. The contribution originating from $V + \text{jets}$ processes is modelled using a double-Gaussian distribution with the shape parameters determined from simulation, while the background contribution is fit to data using a fourth-order exponentiated polynomial. The ability of

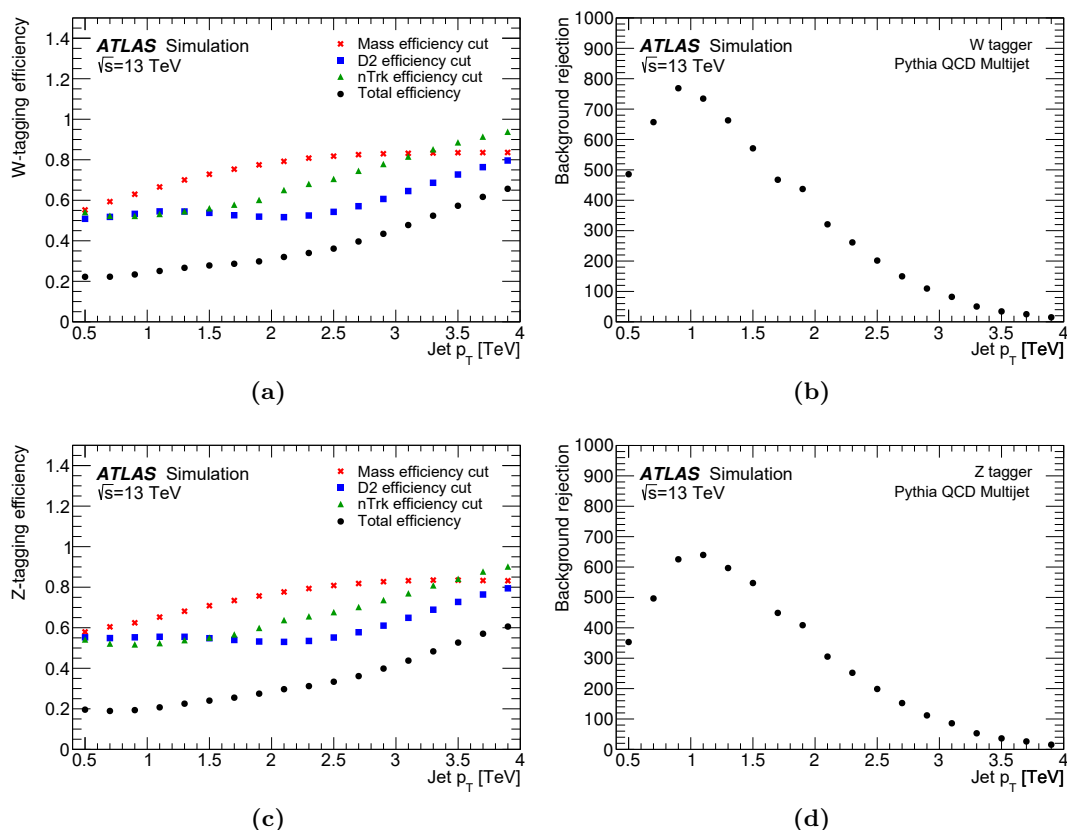


Figure 3. The (a) per-boson signal efficiency for the jet mass, D_2 , and n_{trk} selections, as well as the combined efficiency and (b) background rejection (1/efficiency) of the W tagger for HVT $W' \rightarrow WZ \rightarrow qq\bar{q}\bar{q}$ and MC simulated multijets as a function of the jet p_T . Corresponding values for the Z tagger are shown in (c) and (d).

the fit to extract the correct $V + \text{jets}$ yield (also called MC closure) is tested in simulation by injecting signals of various strengths. Good linearity is found and the method is deemed reliable. By comparing the measured event yield in data and MC simulation, potential differences in the selection efficiency (s_{Tag}) can be probed. Expected contributions of about 5% from $t\bar{t}$ events are subtracted based on MC simulation. The cross-section of $V + \text{jets}$ at a $V p_T$ of about 600 GeV is modelled with about 10% accuracy by the simulation [70]. Additional systematic uncertainties in the fitted $V + \text{jets}$ event yield from closure, from the uncertainty in the $t\bar{t}$ contribution, as well as from the fit parameterisation are considered. The relative efficiency of the D_2 and n_{trk} selections is extracted for V bosons with p_T starting from 600 GeV, while the analysis extends to $p_T = 3.5$ TeV. To estimate the dependence of the modelling on the jet p_T , the distribution of the D_2 and n_{trk} variables is compared in an inclusive sample in data and MC simulation as a function of jet p_T . The observed residual mismodelling as a function of jet p_T is taken into account as an additional 5% uncertainty in the relative efficiency.

The fit to data is shown in figure 4. This fit only extracts the overall yield, while the width and mean of the W/Z peaks are fixed from similar fits performed on MC simulation.

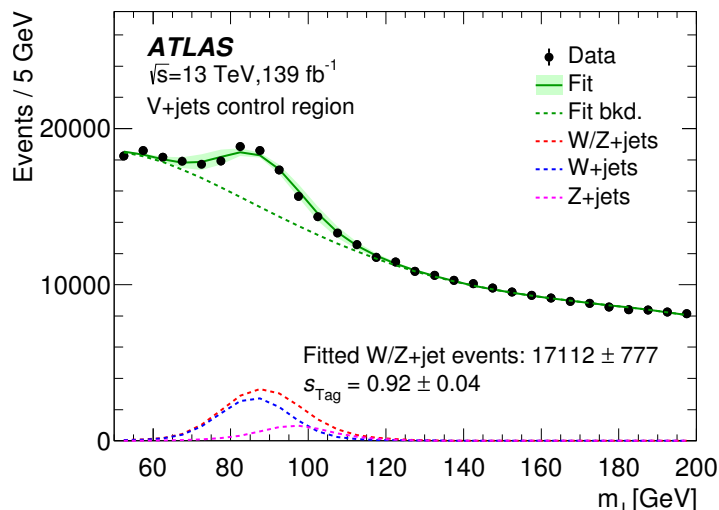


Figure 4. Jet mass distribution for data in the region enhanced in $V + \text{jets}$ events after boson tagging based only on the D_2 and n_{trk} variables. The result of fitting to the sum of functions for the $V + \text{jets}$ and background events is also shown. The shown fit uncertainty reflects the uncertainty in shape and positions of the W and Z peaks. At the bottom, the fitted contribution to the observed jet mass spectra from the $V + \text{jets}$ signal is shown. The fitted relative efficiency of the D_2 and n_{trk} selections is $s_{\text{Tag}} = 0.92 \pm 0.04$, where the uncertainty is purely statistical.

The fitted relative efficiency of the D_2 and n_{trk} selections in data compared with MC simulation is $s_{\text{Tag}} = 0.92 \pm 0.04$ (stat) ± 0.02 (closure) ± 0.03 ($t\bar{t}$) ± 0.02 (fit) ± 0.05 (p_{T} range) ± 0.10 (theory), or $s_{\text{Tag}} = 0.92 \pm 0.13$. This is applied as a scale-factor to the signal MC events, where the uncertainty in it reflects the uncertainty in the W/Z -tagging efficiency in the simulation. Additional fits allowing both the width and the mean of the W/Z peaks to float are used to compare the efficiency of the jet mass window of the boson taggers in data and simulation. Excellent agreement is found, and no additional uncertainty is assigned. The polarisation of the vector bosons in this control region will be different from those in different signal models, but as no other physics process has sufficient sensitivity to probe the data-to-simulation agreement it is assumed that the simulation models the polarisation effects sufficiently well that the scale-factors can be applied globally.

6.3 Signal and background selection efficiency

After boson tagging, the data is categorised into five non-exclusive signal regions (SRs): events with two jets identified as WW , ZZ , or WZ form three SRs, and events with two jets identified as either WZ or WW , and either WW or ZZ form two. The latter provides the highest sensitivity to the benchmark signals described in section 4.1 while the others probe the individual decay channels. Only the boson-tagging requirement differs between the regions. The highest mass jet of the two highest p_{T} jets is considered as the candidate for the higher boson mass requirement. For the WZ selection this means that the highest mass jet must satisfy the Z boson selections and the second highest one the W boson selections. The selection requirements are summarised in table 1.

Signal region	Veto events with leptons: No e or μ with $p_T > 25$ GeV and $ \eta < 2.5$ Event pre-selection: ≥ 2 large- R jets with $ \eta < 2.0$ and mass > 50 GeV $p_{T1} > 500$ GeV and $p_{T2} > 200$ GeV $m_{JJ} > 1.3$ TeV Topology and boson tag: $ \Delta y = y_1 - y_2 < 1.2$ $A = (p_{T1} - p_{T2}) / (p_{T1} + p_{T2}) < 0.15$ Boson tag with D_2 variable, n_{trk} variable, and W or Z mass window
V +jets control region	Veto events with leptons: No e or μ with $p_T > 25$ GeV and $ \eta < 2.5$ V +jets selection: ≥ 2 large- R jets with $ \eta < 2.0$ $p_{T1} > 600$ GeV and $p_{T2} > 200$ GeV Boson tag with D_2 and n_{trk} variables on either jet Anti-boson tag with D_2 variable on other jet

Table 1. Event selection requirements and definition of the different regions used in the analysis. Different requirements are indicated for the highest- p_T (leading) jet with index 1 and the second highest- p_T (subleading) jet with index 2.

The selection efficiency, defined as the number of selected events at different stages of the selection divided by the number of generated events, as a function of the resonance mass, is shown in figure 5 for the HVT Z' decaying into WW and for the bulk G_{KK} decaying into ZZ . Similar efficiencies are obtained in the WZ final state for the HVT model and in the WW final state for the bulk RS models. The figure shows that, among the different selection criteria described above, the boson tagging reduces the signal efficiency the most. However, this particular selection also provides the most significant suppression of the dominant multijet background. The resulting width of the m_{JJ} distributions in the signal region for a HVT model A $W' \rightarrow WZ$ (Bulk RS graviton $\rightarrow ZZ$) is about 6% (10%) of its mean value across the mass range studied, corresponding to about 120 GeV (200 GeV) at 2 TeV. Multijet background events are suppressed with a rejection factor greater than 10^5 across the entire m_{JJ} search range, as determined from simulation.

7 Background parameterisation

The search for diboson resonances is performed by looking for narrow peaks above the smoothly falling m_{JJ} distribution expected from the SM. The background in the search is estimated empirically from the observed m_{JJ} spectrum in the signal region. The background estimation procedure is based on a binned maximum-likelihood fit of the following parameterised form to the observed m_{JJ} spectrum:

$$\frac{dn}{dx} = p_1(1-x)^{p_2-\xi p_3} x^{-p_3} \tag{7.1}$$

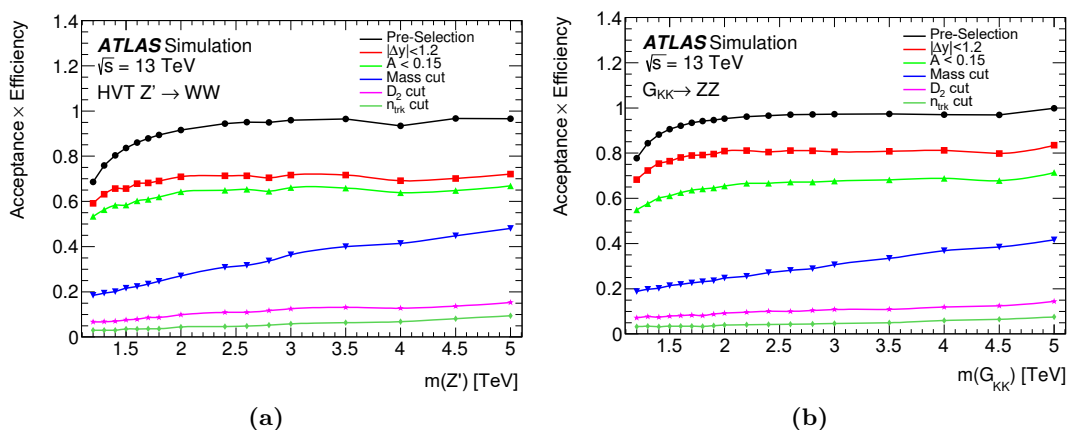


Figure 5. The acceptance \times efficiency for the selection, defined as the number of selected events at different stages of the selection divided by the number of generated events, for (a) $HVT Z' \rightarrow WW$ and (b) $G_{KK} \rightarrow ZZ$ as a function of mass. The selections are applied in sequence and include pre-selections, topological selections on $|\Delta y_{12}|$, p_T asymmetry A , and boson tagging using jet mass, D_2 , and n_{trk} .

where $x = m_{JJ}/\sqrt{s}$, p_1 is a normalisation factor, p_2 and p_3 are dimensionless shape parameters, and ξ is a constant. The value of ξ is derived in an iterative way, minimising the correlation between p_2 and p_3 in the fit, for each m_{JJ} distribution. It is confirmed that the complexity of this fit function is sufficient for the expected number of events in the signal regions by performing Wilks likelihood-ratio tests [71]. The fit is performed to the m_{JJ} distribution in each signal region in data with a constant bin size of 100 GeV. This choice is motivated by the experimental resolution.

The modelling of the parametric shape in eq. (7.1) is tested in dedicated fit control regions (CRs) in data. These CRs are designed to resemble the expected background in the SRs in both their shape and number of events, assuming that no signal contribution is present. Four regions are defined as shown in figure 6, where A and B differ for each tested SR. A possible contamination in region A, C, or D from a potential beyond-the-SM signal is negligible. Region B corresponds to the nominal signal regions.

The probability of misidentifying either the highest or second highest mass jet in an event as a W or Z boson in a data sample dominated by multijets is parameterised as a function of jet p_T using regions C and A. It is validated on data that such a probability is independent of $|\Delta y_{12}|$. Since a misidentification correlation between the two leading jets of the multijet background is observed after the pre-selections, the probability of the second highest mass jet is derived by requiring the highest mass jet to be in the mass window of the boson tagger. By applying per-jet weights, for the inverted selections, depending on the jet p_T , events of the region D are transformed to resemble region B — the fit CRs. To correctly take into account the expected statistical fluctuations and uncertainties, the CR distributions are assigned the correct Poisson errors, and fluctuated accordingly. The last step is repeated multiple times, fitting each distribution with the background fit function, and evaluating the goodness-of-fit χ^2/NDF . Bins with fewer than five events are

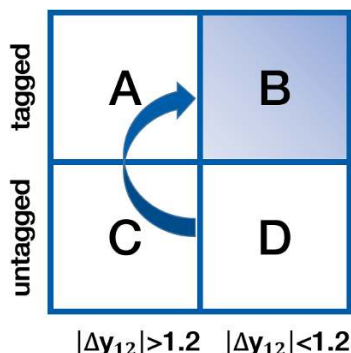


Figure 6. Four orthogonal regions used to build the fit control region for each signal region. A: $|\Delta y_{12}| > 1.2$ with both the jets boson-tagged, B: $|\Delta y_{12}| < 1.2$ with both the jets boson-tagged (this is the nominal signal region), C: $|\Delta y_{12}| > 1.2$ with the event not boson-tagged, D: $|\Delta y_{12}| < 1.2$ with the event not boson-tagged. Regions A and C are used to derive a per-event transfer factor from region D to the fit control region, which is representative of region B. A and C are also signal-depleted due to the $|\Delta y_{12}| > 1.2$ requirement.

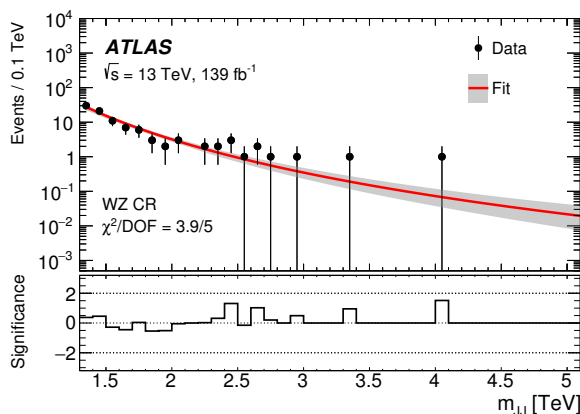


Figure 7. Comparison between fitted background shape and the m_{JJ} spectra in an example WZ fit control region in data. The fitted background distribution is normalised to the data shown in the displayed mass range. The shaded bands represent the uncertainty in the background expectation calculated from the maximum-likelihood function. The lower panel shows the significance, defined as the z -value as described in ref. [72].

combined with bins that contain at least five events to compute the number of degrees of freedom (NDF). On average, the χ^2/NDF is equal to unity with no cases for which the fit fails. Figure 7 shows the fit result performed in an example WZ fit CR. Similar results are obtained for the other CRs, confirming the ability of the chosen background fit function (eq. (7.1)) to describe the expected background dijet mass spectra in the SRs. It is validated on both the data and the simulation that this parametric background description is valid up to 8.0 TeV, which is also the mass up to which the observed m_{JJ} spectra are fit.

The statistical uncertainty in the background expectation comes directly from the uncertainty in the fitted parameters of the background function, which assumes a smoothly falling m_{JJ} distribution. Possible additional uncertainties due to the background model are

assessed by considering signal-plus-background fits (also called spurious signal tests) of the chosen function to the fit control regions of data in which a signal contribution is expected to be negligible. The background is modelled with eq. (7.1) and the signal is modelled using resonance mass distributions from simulation. These procedures were estimated to introduce a bias smaller than 25% of the statistical uncertainty in the background estimate at any mass in the search region, and no additional uncertainty is assigned.

8 Systematic uncertainties

The uncertainties affecting the background modelling are taken directly from the errors in the fit parameters of the background estimation procedure described in section 7. The systematic uncertainties in the expected signal yield and shapes arise from detector effects and MC modelling and are assessed and expressed in terms of nuisance parameters in the statistical analysis as described in section 9.2. The dominant sources of uncertainty in the signal modelling arise from uncertainties in the large- R jet tagging efficiency and the jet p_T calibration. These two uncertainties, and the uncertainty in the fitted background, are also the only ones significantly affecting the statistical results.

The uncertainty in the jet p_T scale (Jp_T S) is evaluated using track-to-calorimeter double ratios between data and simulation [73]. The ratio of the calorimeter and track measures of jet p_T is expected to be the same in data and simulation and any observed differences are assigned as baseline systematic uncertainties. Uncertainties obtained from this procedure assume no correlation between the two p_T measures, while any residual correlation would modify them by a certain factor. An upper limit to the correlation between the two p_T measures is found to be at the percent level by comparing the results of this double-ratio procedure between jets built from TCC inputs and jets built from calorimeter-only inputs. Additional uncertainties due to the track reconstruction efficiency, track impact parameter resolution, and track fake rate are taken into account. The size of the total Jp_T S uncertainty varies with jet p_T and is between 2.5% and 5% for the full mass range.

The impact of the jet p_T resolution uncertainty is evaluated event-by-event by re-running the analysis with an additional Gaussian smearing applied to the input jets' p_T to degrade the nominal resolution by the systematic uncertainty value. The systematic uncertainty in the width of the Gaussian distribution is an absolute 2% per jet, and is symmetrised.

Uncertainty in the jet mass scale and resolution influences the observed jet mass, affecting the boson-tagging efficiency. Any uncertainty in the value of the boson-tagging discriminant D_2 or n_{trk} , would also affect the selection efficiency of the analysis. A scale-factor for the W/Z -tagging efficiency is derived as described in section 6.2. The changes to the overall yield is hence corrected by $0.85^{+0.23}_{-0.21}$ per event with the boson-tagging efficiency scale-factor, assuming full correlation between the two jets. The uncertainty in the scale-factor is assigned as a two-sided variation in the yield. Additional studies comparing jet properties in data and simulation confirm this uncertainty to be valid up to 7.0 TeV.

The uncertainty in the combined 2015–2018 integrated luminosity is 1.7% [74], obtained using the LUCID-2 detector [75] for the primary luminosity measurements. The uncertainty from the trigger selection is found to be negligible, as the minimum requirement on the dijet invariant mass of 1.3 TeV guarantees that the trigger is fully efficient.

Uncertainties in the behaviour of the PDFs at high Q^2 values can potentially have a large effect on the signal acceptance. This systematic uncertainty is estimated by taking the envelope formed by the largest deviations produced by the errors of three PDF sets, as set out by the PDF4LHC group [76]. A constant 1% uncertainty is applied in the case of the RS and radion models, and a pole-mass-dependent uncertainty ranging from 1%–12% is applied in the case of the HVT model. Systematic variations are used to cover uncertainties in the A14 tuned parameter values describing initial-state radiation, final-state radiation, and multi-parton interactions. The uncertainty in the signal acceptance is evaluated at the generator level, before boson-tagging requirements. Following the same procedure as for the PDFs, constant uncertainties of 3% (5%) are applied for the HVT (RS and radion) models.

9 Results

9.1 Background fit

Figure 8 shows the comparison of the dijet mass distributions of the selected events in the combined $WW + WZ$ and $WW + ZZ$ signal regions with the expected background distribution from the background-only fits to the data. The fitted background functions shown, labelled ‘Fit’, are evaluated in bins between 1.3 TeV and 8.0 TeV. No events are observed beyond 5.0 TeV. A total of 119 and 113 events are observed above 1.3 TeV in the $WW + WZ$ and $WW + ZZ$ signal regions, respectively. Due to the non-exclusive selections of the boson taggers, about 50% of events satisfying the WW selection also satisfy the ZZ selection. The highest mass event at 4.4 TeV is the same for both signal regions, and it is compatible with the background expectation in the high mass region.

9.2 Statistical analysis

In the statistical analysis, the parameter of interest is the *signal strength*, which is defined as a scale-factor to the predicted signal normalisation of the model being tested. The analysis follows the *frequentist* approach with a test statistic based on the profile-likelihood ratio [77]. The test statistic extracts information about the signal strength from the binned maximum-likelihood fit of the signal-plus-background model to the data. The likelihood model is defined as,

$$\mathcal{L} = \prod_i P_{\text{pois}}(n_{\text{obs}}^i | n_{\text{exp}}^i) \times G(\alpha) \times \mathcal{N}(\theta)$$

where $P_{\text{pois}}(n_{\text{obs}}^i | n_{\text{exp}}^i)$ is the Poisson probability to observe n_{obs}^i events if n_{exp}^i events are expected, $G(\alpha)$ are a series of Gaussian probability density functions modelling the systematic uncertainties, α , related to the shape of the signal, and $\mathcal{N}(\theta)$ is a log-normal distribution for the nuisance parameters, θ , modelling the systematic uncertainty in the

signal normalisation. The expected number of events is the bin-wise sum of those expected for the signal and background: $\mathbf{n}_{\text{exp}} = \mathbf{n}_{\text{sig}} + \mathbf{n}_{\text{bg}}$. The expected number of background events in dijet mass bin i , n_{bg}^i , is obtained by integrating dn/dx obtained from eq. (7.1) over that bin. Thus \mathbf{n}_{bg} is a function of the dijet background parameters p_1 , p_2 and p_3 . The expected number of signal events, \mathbf{n}_{sig} , is evaluated from MC simulation assuming the cross-section of the model under test multiplied by the signal strength, including the effects of the systematic uncertainties described in section 8.

The significance of observed excesses over the background-only prediction is quantified using the local p_0 -value, defined as the probability of the background-only model to produce a signal-like fluctuation at least as large as that observed in the data. The most extreme p_0 has a local significance of 1.8 standard deviations, and is found when testing the HVT $W' \rightarrow WW$ hypothesis at a resonance mass of 1.8 TeV. This is within the expected fluctuation of the background.

Limits at 95% confidence level (CL) on the production cross-section times branching fraction to diboson final states for the benchmark signals are set with sampling distributions generated using pseudo-experiments. All systematic uncertainties are considered. The uncertainty in the W/Z -tagging efficiency is dominant at lower masses, while the uncertainty in the background modelling has largest impact at high masses. Uncertainties in the jet p_T scale are at the percent level but are subordinate across the full mass range. The cross-section limits extracted for the different benchmark scenarios in the $WW + WZ$ and $WW + ZZ$ signal regions are shown in figure 9 and table 2. Table 3 presents the resonance mass ranges excluded at the 95% CL in the various signal regions and signal models considered in the search.

10 Conclusion

A search for narrow heavy resonances decaying into dibosons in the all hadronic channel is performed using 139 fb^{-1} of proton-proton collisions at $\sqrt{s} = 13 \text{ TeV}$ collected by the ATLAS experiment at the LHC from 2015 to 2018. The results of the search are shown for the $WW + WZ$, $WW + ZZ$ channels, and are interpreted in terms of a radion model, two HVT benchmark models, and a bulk G_{KK} model. The data are in agreement with the background expectations in all channels. Upper limits on the production cross-section times branching ratio to diboson final states for new resonances with masses greater than 1.3 TeV are set at the 95% CL. These results exclude at the 95% CL the production of $WW + WZ$ from the HVT model A (model B) with $g_V = 1$ ($g_V = 3$) with masses in the range of 1.3 TeV–3.5 TeV (1.3 TeV–3.8 TeV). Production of a G_{KK} in the bulk RS model with $k/\overline{M}_{\text{Pl}} = 1$ is excluded in the range 1.3 TeV–1.8 TeV, at the 95% CL. Upper limits on the production cross section times branching ratio for a scalar-like radion are set with values of 5.72 fb and 1.86 fb at scalar masses of 2 TeV and 3 TeV, respectively.

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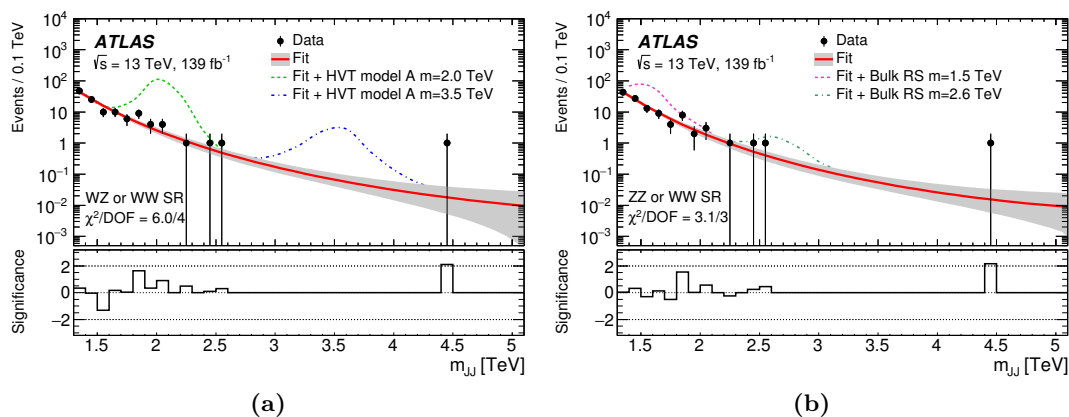


Figure 8. Background-only fits to the dijet mass (m_{JJ}) distributions in data after tagging in the combined (a) $WW + WZ$, and (b) $WW + ZZ$ signal region. The shaded bands represent the uncertainty in the background expectation calculated from the maximum-likelihood function. The lower panels show the significance, defined as the z -value as described in ref. [68]. Selected theoretical signal distributions are overlaid on top of the background.

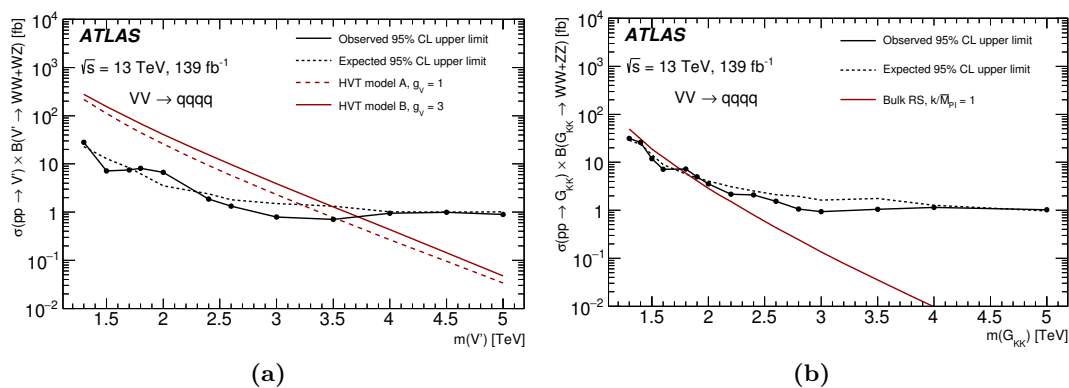


Figure 9. Observed and expected limits at 95% CL on the cross-section times branching ratio for $WW + WZ$ production as a function of (a) $m_{V'}$, and for $WW + ZZ$ production as a function of (b) the Bulk RS graviton $m_{G_{KK}}$. The predicted cross-section times branching ratio is shown (a) as dashed and solid lines for the HVT models A with $g_V = 1$ and B with $g_V = 3$, respectively, and (b) as a solid line for the bulk RS model with $k/\overline{M}_{Pl} = 1$.

Mass [TeV]	Observed Limit [fb]	Expected Limit [fb]	Prediction [fb]
2.0	5.72	5.75	4.286
3.0	1.86	2.85	0.415
4.0	1.98	2.34	0.040
5.0	1.98	2.02	0.006

Table 2. Observed and expected limits at 95% CL on cross-section times branching ratio for $WW + ZZ$ production for different radion masses m_{radion} , as well as the predicted cross-section times branching ratio.

Model	Signal Region	Excluded mass range [TeV]
Radion	WW	none
	ZZ	none
	$WW + ZZ$	none
HVT model A, $g_V = 1$	WW	1.3–2.9
	WZ	1.3–3.4
	$WW + WZ$	1.3–3.5
HVT model B, $g_V = 3$	WW	1.3–3.1
	WZ	1.3–3.6
	$WW + WZ$	1.3–3.8
Bulk RS, $k/\overline{M}_{\text{Pl}} = 1$	WW	1.3–1.6
	ZZ	none
	$WW + ZZ$	1.3–1.8

Table 3. Observed excluded resonance masses (at 95% CL) in the individual and combined signal regions for the HVT, bulk RS and radion models.

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