# Second order partial differential equations (SOPDEs) and nonlinear connections on the tangent bundle of $k^{1}$-velocities of a manifold 

Florian Munteanu


#### Abstract

In this paper I will describe two distinct ways by which it obtain a sequence of SOPDEs and a sequences of nonlinear connections on the tangent bundle of $k^{1}$-velocties $T_{k}^{1} M$, starting from a given SOPDE, follows the ideas of papers [25], [26], [27], [28]. Some properties about this sequences is also presented. Interesting cases appear in the presence of a regular Lagrangian on $T_{k}^{1} M$.


M.S.C. 2000: 53C05, 53C15, 53C60, 57R15, 58F05.

Key words: tangent bundle of $k^{1}$-velocities, almost $k$-tangent structure, $k$-vector field, SOPDE, nonlinear connection, regular Lagrangian on $T_{k}^{1} M$, generalized EulerLagrange equations.

## 1 Introduction

In this section we present the main notions used in the paper, with a strong accent on the necessary tools for obtain the results from the second section.

### 1.1 The tangent bundle of $k^{1}$-velocities of a manifold $M$

An almost tangent structure $J$ on a $2 n$-dimensional manifold $M$ is tensor field of type $(1,1)$ of constant rank $n$ such that $J^{2}=0$. The manifold $M$ is then called an almost tangent manifold. Almost tangent structures were introduced by Clark and Bruckheimer [5] and Eliopoulos [6] around 1960 and have been studied by numerous authors (see [10, 11, 12, 13, 14, 15, 16, 17, 18]).

The canonical model of these structures is the tangent bundle $\tau_{M}: T M \rightarrow M$ of an arbitrary manifold $M$. Recall that for a vector $X_{x}$ at a point $x \in M$ its vertical lift is the vector on $T M$ given by

$$
X_{x}^{V}\left(v_{x}\right)=\frac{d}{d t}\left(v_{x}+t X_{x}\right)_{\mid t=0} \in T_{v_{x}}(T M)
$$

for all points $v_{x} \in T M$.
Differential Geometry - Dynamical Systems, Vol.8, 2006, pp. 166-180. © Balkan Society of Geometers, Geometry Balkan Press 2006.

The canonical tangent structure $J$ on $T M$ is defined by

$$
J_{v_{x}}\left(Z_{v_{x}}\right)=\left(\left(\tau_{M}\right)_{*}\left(v_{x}\right) Z_{v_{x}}\right)_{v_{x}}^{V}
$$

for all vectors $Z_{v_{x}} \in T_{v_{x}}(T M)$, and it is locally given by

$$
\begin{equation*}
J=\frac{\partial}{\partial v^{i}} \otimes d x^{i} \tag{1.1}
\end{equation*}
$$

with respect the bundle coordinates on $T M$. This tensor $J$ can be regarded as the vertical lift of the identity tensor on $M$ to $T M$ [19].

The almost $k$-tangent structures were introduced as generalization of the almost tangent structures [7, 8].

Definition 1.1 An almost $k$-tangent structure $J$ on a manifold $M$ of dimension $n+$ $k n$ is a family $\left(J^{1}, \ldots, J^{k}\right)$ of tensor fields of type $(1,1)$ such that

$$
\begin{equation*}
J^{A} \circ J^{B}=J^{B} \circ J^{A}=0, \quad \operatorname{rank} J^{A}=n, \quad \operatorname{Im} J^{A} \cap\left(\oplus_{B \neq A} \operatorname{Im} J^{B}\right)=0 \tag{1.2}
\end{equation*}
$$

for $1 \leq A, B \leq k$. In this case the manifold $M$ is then called an almost $k$-tangent manifold.

The canonical model of these structures is the $k$-tangent vector bundle $T_{k}^{1} M=$ $J_{0}^{1}\left(R^{k}, M\right)$ of an arbitrary manifold $M$, that is the vector bundle with total space the manifold of 1-jets of maps with source at $0 \in R^{k}$ and with projection map $\tau: T_{k}^{1} M \rightarrow$ $M, \tau\left(j_{0}^{1} \sigma\right)=\sigma(0)$. This bundle is also known as the tangent bundle of $k^{1}$-velocities of $M$ [19].

The manifold $T_{k}^{1} M$ can be canonically identified with the Whitney sum of $k$ copies of $T M$, that is

$$
\begin{array}{rlc}
T_{k}^{1} M & \equiv & T M \oplus \cdots \oplus T M \\
j_{0}^{1} \sigma & \equiv & \left(j_{0}^{1} \sigma_{1}=v_{1}, \ldots, j_{0}^{1} \sigma_{k}=v_{k}\right)
\end{array}
$$

where $\sigma_{A}=\sigma(0, \ldots, t, \ldots, 0)$ with $t \in R$ at position $A$ and $v_{A}=\left(\sigma_{A}\right)_{*}(0)\left(\frac{d}{d t} 0\right)$.
If $\left(x^{i}\right)$ are local coordinates on $U \subseteq M$ then the induced local coordinates $\left(x^{i}, v_{A}^{i}\right)$, $1 \leq i \leq n, 1 \leq A \leq k$, on $\tau^{-1}(U) \equiv T_{k}^{1} U$ are given by

$$
x^{i}\left(j_{0}^{1} \sigma\right)=x^{i}(\sigma(0)), \quad v_{A}^{i}\left(j_{0}^{1} \sigma\right)=\left.\frac{d}{d t}\left(x^{i} \circ \sigma_{A}\right)\right|_{t=0}=v_{A}\left(x^{i}\right)
$$

A local change of coordinates on $T_{k}^{1} M,\left(x^{i}, v_{A}^{i}\right) \rightarrow\left(\tilde{x}^{i}, \tilde{v}_{A}^{i}\right)$, is given by the rule:

$$
\begin{cases}\tilde{x}^{i}=\tilde{x}^{i}\left(x^{1}, \ldots, x^{n}\right) & , \operatorname{rank}\left(\frac{\partial \tilde{x}^{i}}{\partial x^{j}}\right)=n  \tag{1.3}\\ \tilde{v}_{A}^{i}=\frac{\partial \tilde{x}^{i}}{\partial x^{j}} v_{A}^{j} & , 1 \leq A \leq k, 1 \leq i \leq n\end{cases}
$$

If $T_{u} T_{k}^{1} M$ is the tangent space to $T_{k}^{1} M$ in $u$, then $\left\{\frac{\partial}{\partial x^{i}}(u), \frac{\partial}{\partial v_{A}^{i}}(u)\right\}$ is its local basis and $\left\{d x^{i}(u), d v_{A}^{i}(u)\right\}$ is its dual local basis (local cobasis). With respect to (1.3)
the basis and the cobasis have the following rules of changing:

$$
\left\{\begin{align*}
\frac{\partial}{\partial x^{i}} & =\frac{\partial \tilde{x}^{j}}{\partial x^{i}} \frac{\partial}{\partial \tilde{x}^{j}}+\frac{\partial \tilde{v}_{A}^{j}}{\partial x^{i}} \frac{\partial}{\partial \tilde{v}_{A}^{j}}  \tag{1.4}\\
\frac{\partial}{\partial v_{A}^{i}} & =\frac{\partial \tilde{x}^{j}}{\partial x^{i}} \frac{\partial}{\partial \tilde{v}_{A}^{j}}
\end{align*}\right.
$$

and

$$
\left\{\begin{align*}
d x^{i} & =\frac{\partial \tilde{x}^{i}}{\partial x^{j}} d \tilde{x}^{j}  \tag{1.5}\\
d v_{A}^{i} & =\frac{\partial v_{A}^{i}}{\partial \tilde{x}^{j}} d \tilde{x}^{j}+\frac{\partial x^{i}}{\partial \tilde{x}^{j}} d \tilde{v}_{A}^{j}
\end{align*}\right.
$$

It is observe that $\frac{\partial \tilde{x}^{j}}{\partial x^{i}}=\frac{\partial \tilde{v}_{A}^{j}}{\partial v_{A}^{i}}$.
Definition 1.2 For a vector $X_{x}$ at $M$ we define its vertical $A$-lift $\left(X_{x}\right)^{A}$ as the vector on $T_{k}^{1} M$ given by
$\left(X_{x}\right)^{A}\left(j_{0}^{1} \sigma\right)=\left.\frac{d}{d t}\left(\left(v_{1}\right)_{x}, \ldots,\left(v_{A-1}\right)_{x},\left(v_{A}\right)_{x}+t X_{x},\left(v_{A+1}\right)_{x} \ldots,\left(v_{k}\right)_{x}\right)\right|_{t=0} \in T_{j_{0}^{1} \sigma}\left(T_{k}^{1} M\right)$
for all points $\left.j_{0}^{1} \sigma \equiv\left(\left(v_{1}\right)_{x}, \ldots,\left(v_{k}\right)_{x}\right)\right) \in T_{k}^{1} M$.
In local coordinates we have

$$
\begin{equation*}
\left(X_{x}\right)^{A}=\sum_{i=1}^{n} a^{i} \frac{\partial}{\partial v_{A}^{i}} \tag{1.6}
\end{equation*}
$$

for a vector $X_{x}=a^{i} \partial / \partial x^{i}$.
The canonical vertical vector fields $C_{B}^{A}$ on $T_{k}^{1} M$ are defined by

$$
\begin{equation*}
C_{B}^{A}\left(x, X_{1}, X_{2}, \ldots, X_{k}\right)=\left(X_{B}\right)^{A} \tag{1.7}
\end{equation*}
$$

and are locally given by $C_{B}^{A}=v_{B}^{i} \frac{\partial}{\partial v_{A}^{i}}$.
The canonical $k$-tangent structure $\left(J^{1}, \ldots, J^{k}\right)$ on $T_{k}^{1} M$ is defined by

$$
J^{A}\left(Z_{j_{0}^{1} \sigma}\right)=\left(\tau_{*}\left(Z_{j_{0}^{1} \sigma}\right)\right)^{A}
$$

for all vectors $Z_{j_{0}^{1} \sigma} \in T_{j_{0}^{1} \sigma}\left(T_{k}^{1} M\right)$. In local coordinates we have

$$
\begin{equation*}
J^{A}=\frac{\partial}{\partial v_{A}^{i}} \otimes d x^{i} \tag{1.8}
\end{equation*}
$$

The tensors $J^{A}$ can be regarded as the $\left(0, \ldots, 1_{A}, \ldots, 0\right)$-lift of the identity tensor on $M$ to $T_{k}^{1} M$ defined in [19].

### 1.2 The cotangent bundle of $k^{1}$-covelocities of $M$ and $\left(T_{k}^{1}\right)^{*} M$

Almost cotangent structures were introduced by Bruckheimer [4]. An almost cotangent structure on a $2 m$-dimensional manifold $M$ consists of a pair $(\omega, V)$ where $\omega$ is a symplectic form and $V$ is a distribution such that

$$
\text { (i) } \quad \omega\rfloor_{V \times V}=0, \quad \text { (ii) } \quad \operatorname{ker} \omega=\{0\}
$$

The canonical model of this structure is the cotangent bundle $\tau_{M}^{*}: T^{*} M \rightarrow M$ of an arbitrary manifold $M$, where $\omega$ is the canonical symplectic form $\omega_{0}=-d \theta_{0}$ on $T^{*} M$ and $V$ is the vertical distribution. Let us recall the definition of the Liouville form $\theta_{0}$ in $T^{*} M$ :

$$
\theta_{0}(\alpha)\left(\tilde{X}_{\alpha}\right)=\alpha\left(\left(\tau_{M}^{*}\right)_{*}(\alpha)\left(\tilde{X}_{\alpha}\right)\right)
$$

for all vectors $\tilde{X}_{\alpha} \in T_{\alpha}\left(T^{*} M\right)$. In local coordinates $\left(x^{i}, p_{i}\right)$ on $T^{*} M$

$$
\begin{equation*}
\theta_{0}=p_{i} d x^{i}, \quad \omega_{0}=d x^{i} \wedge d p_{i}, \quad V=\left\langle\frac{\partial}{\partial p_{1}}, \ldots, \frac{\partial}{\partial p_{k}}\right\rangle \tag{1.9}
\end{equation*}
$$

Definition 1.3 [1, 2] A $k$-symplectic structure on a manifold $M$ of dimension $N=$ $n+k n$ is a family $\left(\omega_{A}, V ; 1 \leq A \leq k\right)$, where each $\omega_{A}$ is a closed 2 -form and $V$ is an $n k$-dimensional distribution on $M$ such that

$$
\text { (i) } \quad \omega_{A_{\lrcorner V \times V}}=0, \quad \text { (ii) } \quad \cap_{A=1}^{k} \operatorname{ker} \omega_{A}=\{0\}
$$

In this case $\left(M, \omega_{A}, V\right)$ is called a $k$-symplectic manifold.
The canonical model of this structure is the $k$-cotangent bundle $\left(T_{k}^{1}\right)^{*} M=J^{1}\left(M, \mathbf{R}^{k}\right)_{0}$ of an arbitrary manifold $M$, that is the vector bundle with total space the manifold of 1-jets of maps with target at $0 \in R^{k}$, and projection $\tau^{*}\left(j_{x, o}^{1} \sigma\right)=x$.

The manifold $\left(T_{k}^{1}\right)^{*} M$ can be canonically identified with the Whitney sum of $k$ copies of $T^{*} M$, say

$$
\begin{aligned}
\left(T_{k}^{1}\right)^{*} M & \equiv T^{*} M \oplus \cdots \oplus T^{*} M \\
j_{x, 0} \sigma & \equiv\left(j_{x, 0}^{1} \sigma^{1}, \ldots, j_{x, 0}^{k} \sigma^{k}\right)
\end{aligned}
$$

where $\sigma^{A}=\pi_{A} \circ \sigma: M \longrightarrow R$ is the $A$-th component of $\sigma$.
The canonical $k$-symplectic structure $\left.\left(\omega_{0}\right)_{A}, V ; 1 \leq A \leq k\right)$, on $\left(T_{k}^{1}\right)^{*} M$ is defined by

$$
\begin{gathered}
\left(\omega_{0}\right)_{A}=\left(\tau_{A}^{*}\right)^{*}\left(\omega_{0}\right) \\
V\left(j_{x, 0}^{1} \sigma\right)=\operatorname{ker}\left(\tau^{*}\right)_{*}\left(j_{x, 0}^{1} \sigma\right)
\end{gathered}
$$

where $\tau_{A}^{*}=\left(T_{k}^{1}\right)^{*} M \rightarrow T^{*} M$ is the projection on the $A^{t h}$-copy $T^{*} M$ of $\left(T_{k}^{1}\right)^{*} M$, and $\omega_{0}$ is the canonical symplectic structure of $T^{*} M$.

One can also define the 2 -forms $\omega_{A}$ by $\omega_{A}=-d \theta_{A}$ where $\left(\theta_{0}\right)^{A}=\left(\tau_{A}^{*}\right)^{*} \theta_{0}$
If $\left(x^{i}\right)$ are local coordinates on $U \subseteq M$ then the induced local coordinates $\left(x^{i}, p_{i}^{A}\right), 1 \leq$ $i \leq n, 1 \leq A \leq k$ on $\left(T_{k}^{1}\right)^{*} U=\left(\tau^{*}\right)^{-1}(U)$ are given by

$$
x^{i}\left(j_{x, 0}^{1} \sigma\right)=x^{i}(x), \quad p_{i}^{A}\left(j_{x, 0}^{1} \sigma\right)=d_{x} \sigma^{A}\left(\left.\frac{\partial}{\partial x^{i}}\right|_{x}\right)
$$

Then the canonical $k$-symplectic structure is locally given by

$$
\begin{equation*}
\left(\omega_{0}\right)^{A}=\sum_{i=1}^{n} d x^{i} \wedge d p_{i}^{A}, \quad V=\left\langle\frac{\partial}{\partial p_{i}^{1}}, \ldots, \frac{\partial}{\partial p_{i}^{k}}\right\rangle \quad 1 \leq A \leq k \tag{1.10}
\end{equation*}
$$

and $\left(\theta_{0}\right)^{A}=p_{i}^{A} d x^{i}$.

### 1.3 Second Order Partial Differential Equations on $T_{k}^{1} M$

Let $M$ be an arbitrary manifold and $\tau: T_{k}^{1} M \longrightarrow M$ its tangent bundle of $k^{1}$ velocities.

Definition 1.4 $A$ section $X: M \longrightarrow T_{k}^{1} M$ of the projection $\tau$ will be called a $k$ vector field on $M$.

Since $T_{k}^{1} M$ can be canonically identified with the Whitney sum $T_{k}^{1} M \equiv T M \oplus$ $\cdots \oplus T M$ of $k$ copies of $T M$, we deduce that a $k$-vector field $X$ defines a family of vector fields $X_{1}, \ldots, X_{k}$ on $M$. Günther in [3] introduce the following definition.

Definition 1.5 $A n$ integral section of the $k$-vector field $X=\left(X_{1}, \ldots, X_{k}\right)$ passing through a point $x \in M$ on $M$ is a map $\phi: U_{0} \subset \mathbf{R}^{k} \rightarrow M$, defined on some neighborhood $U_{0}$ of $0 \in \mathbf{R}^{k}$, such that

$$
\phi(0)=x, \quad \phi_{*}(t)\left(\frac{\partial}{\partial t^{A}}\right)=X_{A}(\phi(t)) \quad \forall t \in U, \quad 1 \leq A \leq k
$$

or equivalently, $\phi$ satisfies

$$
\begin{equation*}
X \circ \phi=\phi^{(1)} \tag{1.11}
\end{equation*}
$$

where $\phi^{(1)}$ is the first prolongation of $\phi$ defined by

$$
\begin{aligned}
\phi^{(1)}: U_{0} \subset \mathbf{R}^{k} & \longrightarrow T_{k}^{1} M \\
t & \longrightarrow \phi^{(1)}(t)=j_{0}^{1} \phi_{t}
\end{aligned}
$$

where $\phi_{t}(s)=\phi(s+t)$ for all $t, s \in \mathbf{R}^{k}$ such that $s+t \in U_{0}$.
In local coordinates,

$$
\begin{equation*}
\phi^{(1)}\left(t^{1}, \ldots, t^{k}\right)=\left(\phi^{i}\left(t^{1}, \ldots, t^{k}\right), \frac{\partial \phi^{i}}{\partial t^{A}}\left(t^{1}, \ldots, t^{k}\right)\right), \quad 1 \leq A \leq k, 1 \leq i \leq n \tag{1.12}
\end{equation*}
$$

We say that a $k$-vector field $X=\left(X_{1}, \ldots, X_{k}\right)$ on $M$ is integrable if there is an integral section passing through each point of $M$.

Remark If $\phi$ is an integral section of a $k$-vector field $\left(X_{1}, \ldots, X_{k}\right)$ then each curve on $M$ defined by $\phi_{A}=\phi \circ h_{A}$, where $h_{A}: \mathbf{R} \rightarrow \mathbf{R}^{k}$ is the natural inclusion $h_{A}(t)=(0, \ldots, t, \ldots, 0)$, is an integral curve of the vector field $X_{A}$ on $M$, with $1 \leq A \leq k$. We refer to [20,21] for a discussion on the existence of integral sections.

Definition 1.6 $A$ k-vector field on $T_{k}^{1} M$, that is, a section $\xi \equiv\left(\xi_{1}, \ldots, \xi_{k}\right): T_{k}^{1} M \rightarrow$ $T_{k}^{1}\left(T_{k}^{1} M\right)$ of the projection $\tau_{T_{k}^{1} M}: T_{k}^{1}\left(T_{k}^{1} M\right) \rightarrow T_{k}^{1} M$, is a Second Order Partial Differential Equation (SOPDE) if and only if it is also a section of the vector bundle $T_{k}^{1}(\tau): T_{k}^{1}\left(T_{k}^{1} M\right) \rightarrow T_{k}^{1} M$, where $T_{k}^{1}(\tau)$ is defined by $T_{k}^{1}(\tau)\left(j_{0}^{1} \sigma\right)=j_{0}^{1}(\tau \circ \sigma)$.

Let $\left(x^{i}\right)$ be a coordinate system on $M$ and $\left(x^{i}, v_{A}^{i}\right)$ the induced coordinate system on $T_{k}^{1} M$. From the definition we deduce that the local expression of a SOPDE $\xi$ is

$$
\begin{equation*}
\xi_{A}\left(x^{i}, v_{A}^{i}\right)=v_{A}^{i} \frac{\partial}{\partial x^{i}}+\left(\xi_{A}\right)_{B}^{i} \frac{\partial}{\partial v_{B}^{i}}, \quad 1 \leq A \leq k \tag{1.13}
\end{equation*}
$$

We recall that the first prolongation $\phi^{(1)}$ of $\phi: U \subset \mathbf{R}^{k} \rightarrow M$ is defined by

$$
\begin{aligned}
\phi^{(1)}: U \subset \mathbf{R}^{k} & \left.\longrightarrow T_{k}^{1} M\right) \\
t & \longrightarrow \phi^{(1)}(t)=j_{0}^{1} \phi_{t}
\end{aligned}
$$

where $\phi_{t}(s)=\phi(s+t)$ for all $t, s \in R$.
Proposition 1.7 Let $\xi$ an integrable $k$-vector field on $T_{k}^{1} M$. The necessary and sufficient condition for $\xi$ to be a Second Order Partial Differential Equation (SOPDE) is that its integral sections are first prolongations $\phi^{(1)}$ of maps $\phi: \mathbf{R}^{k} \rightarrow M$. That is

$$
\xi_{A}\left(\phi^{(1)}(t)\right)=\left(\phi^{(1)}\right)_{*}(t)\left(\frac{\partial}{\partial t_{A}}\right)(t)
$$

for all $A=1, \ldots, k$. These maps $\phi$ will be called solutions of the SOPDE $\xi$.
From (1.12) and (1.13) we have
Proposition $1.8 \phi: \mathbf{R}^{k} \rightarrow M$ is a solution of the $\operatorname{SOPDE} \xi=\left(\xi_{1}, \ldots, \xi_{k}\right)$, locally given by (1.13), if and only if

$$
\frac{\partial \phi^{i}}{\partial t^{A}}(t)=v_{A}^{i}\left(\phi^{(1)}(t)\right), \quad \frac{\partial^{2} \phi^{i}}{\partial t^{A} \partial t^{B}}(t)=\left(\xi_{A}\right)_{B}^{i}\left(\phi^{(1)}(t)\right)
$$

If $\xi: T_{k}^{1} M \rightarrow T_{k}^{1} T_{k}^{1} M$ is an integrable SOPDE then for all integral sections $\sigma: U \subset \mathbf{R}^{k} \rightarrow T_{k}^{1} M$ we have $\left(\tau_{M} \circ \sigma\right)^{(1)}=\sigma$, where $\tau: T_{k}^{1} M \rightarrow M$ is the canonical projection.

Now we show how to characterize the SOPDEs using the canonical $k$-tangent structure of $T_{k}^{1} M$.
Definition 1.9 The Liouville vector field $C$ on $T_{k}^{1} M$ is the infinitesimal generator of the one parameter group

$$
\begin{array}{rlc}
\mathbf{R} \times\left(T_{k}^{1} M\right) & \longrightarrow & T_{k}^{1} M \\
\left(s,\left(x^{i}, v_{B}^{i}\right)\right) & \longrightarrow & \left(x^{i}, e^{s} v_{B}^{i}\right) .
\end{array}
$$

Thus $C$ is locally expressed as follows:

$$
\begin{equation*}
C=\sum_{B} C_{B}=\sum_{i, B} v_{B}^{i} \frac{\partial}{\partial v_{B}^{i}} \tag{1.14}
\end{equation*}
$$

where each $C_{B}$ corresponds with the canonical vector field on the $B$-th copy of $T M$ on $T_{k}^{1} M$.

Let us remark that each vector field $C_{A}$ on $T_{k}^{1} M$ can also be defined using the $A$-lifts of vectors as follows: $\left.C_{A}\left(\left(v_{1}\right)_{q}, \ldots,\left(v_{k}\right)_{q}\right)=\left(\left(v_{A}\right)_{q}\right)^{A}(v)\right)$.

From (1.8), (1.13) and (1.14) we deduce the following
Proposition 1.10 $A$-vector field $\xi=\left(\xi_{1}, \ldots, \xi_{k}\right)$ on $T_{k}^{1} M$ is a SOPDE if and only if

$$
J^{A}\left(\xi_{A}\right)=C_{A}, \quad \forall 1 \leq A \leq k
$$

where $\left(J^{1}, \ldots, J^{k}\right)$ is the canonical $k$-tangent structure on $T_{k}^{1} M$.
Proposition 1.11 a) With respect to a local changing of coordinates (1.3) the components $\left(\xi_{A}\right)_{B}^{i}$ of the SOPDE $\xi$ change as follows

$$
\begin{equation*}
\left(\widetilde{\xi}_{A}\right)_{B}^{i}=\left(\xi_{A}\right)_{B}^{j} \frac{\partial x^{i}}{\partial x^{j}}+\frac{\partial^{2} x^{i}}{\partial x^{j} \partial x^{k}} v_{A}^{j} v_{B}^{k} \tag{1.15}
\end{equation*}
$$

b) Conversely, if the function $\left(\xi_{A}\right)_{B}^{i}\left(x^{j}, v_{C}^{k}\right)$ are given on every domain of local chart on $T_{k}^{1} M$ so that (1.15) holds for a local change (1.3), then the $k$-vector field $\xi$ given by (1.13) is a SOPDE.

Next we will denote $\widetilde{T_{k}^{1} M}=T_{k}^{1} M \backslash\{0\}$, where 0 is the null section of the projection $\tau: T_{k}^{1} M \rightarrow M$.
Definition 1.12 A function $f: T_{k}^{1} M \rightarrow R$ differentiable of class $C^{\infty}$ on $\widetilde{T_{k}^{1} M}$ and continuous on the null section of the projection $\tau$ is called homogeneous of degree $r \in \mathbf{Z}$ on the fibres on $T_{k}^{1} M$ (briefly $r$-homogeneous) if

$$
\begin{equation*}
f\left(x^{i}, \lambda v_{A}^{i}\right)=\lambda^{r} f\left(x^{i}, v_{A}^{i}\right) \quad, \forall \lambda>0 \tag{1.16}
\end{equation*}
$$

Equivalently, a function $f: T_{k}^{1} M \rightarrow R$, differentiable on $\widetilde{T^{k} M}$ and continuous on the null section of $\tau$ is $r$-homogeneous if and only if

$$
\begin{equation*}
\mathcal{L}_{C} f=r f \tag{1.17}
\end{equation*}
$$

$\mathcal{L}_{C} f$ being the Lie operator of derivation with respect to the canonical vector field $C$.
Definition 1.13 A SOPDE $\xi$ is called homogeneous $S O P D E$, or spray on $T_{k}^{1} M$ if the components $\left(\xi_{A}\right)_{B}^{i}$ are 2-homogeneous functions on the fibres of $T_{k}^{1} M$, that is

$$
\left(\xi_{A}\right)_{B}^{i}\left(x^{j}, \lambda v_{C}^{j}\right)=\lambda^{2}\left(\xi_{A}\right)_{B}^{i}\left(x^{j}, v_{C}^{j}\right) \quad \text { for all } \quad \lambda>0
$$

From (1.16) we deduce the following proposition.
Proposition 1.14 $A S O P D E \xi \equiv\left(\xi_{1}, \ldots, \xi_{k}\right)$ is a homogeneous SOPDE, or spray, if and only if we have

$$
\begin{equation*}
v_{C}^{h} \frac{\partial\left(\xi_{A}\right)_{B}^{i}}{\partial v_{C}^{h}}=2\left(\xi_{A}\right)_{B}^{i} \tag{1.18}
\end{equation*}
$$

which is equivalent to $\left[C, \xi_{A}\right]=\xi_{A}, A=1, \ldots, k$.

### 1.4 Hamiltonian and Lagrangian formalisms

The role played by symplectic manifolds in classical mechanics is here played by the $k$-symplectic manifolds (see Günther, ([3])). Let $\left(M, \omega_{A}, V ; 1 \leq A \leq k\right)$ be a $k-$ symplectic manifold. Let us consider the vector bundle morphism defined by Günther ([3]):

$$
\left.\begin{array}{rl}
\Omega^{\sharp}: & T_{k}^{1} M
\end{array}\right] T^{*} M .
$$

Definition 1.15 Let $H: M \longrightarrow \mathbf{R}$ be a function on $M$. Any $k$-vector field $\left(X_{1}, \ldots, X_{k}\right)$ on $M$ such that

$$
\Omega^{\sharp}\left(X_{1}, \ldots, X_{k}\right)=d H
$$

will be called an evolution $k$-vector field on $M$ associated with the Hamiltonian function $H$.

It should be noticed that in general the solution to the above equation is not unique. Nevertheless, it can be proved [9] that there always exists an evolution $k$ vector field associated with a Hamiltonian function $H$.

Let $\left(x^{i}, p_{i}^{A}\right)$ be a local coordinate system on $M$. Then we have
Proposition 1.16 If $\left(X_{1}, \ldots, X_{k}\right)$ is an integrable evolution $k$-vector field associated to $H$ then its integral sections

$$
\begin{aligned}
\sigma: \quad \mathbf{R}^{k} & \longrightarrow M \\
\left(t^{B}\right) & \longrightarrow\left(\sigma^{i}\left(t^{B}\right), \sigma_{i}^{A}\left(t^{B}\right)\right),
\end{aligned}
$$

are solutions of the classical local Hamilton equations associated with a regular multiple integral variational problem [22]:

$$
\begin{equation*}
\frac{\partial H}{\partial x^{i}}=-\sum_{A=1}^{k} \frac{\partial \sigma_{i}^{A}}{\partial t^{A}}, \quad \frac{\partial H}{\partial p_{i}^{A}}=\frac{\partial \sigma^{i}}{\partial t^{A}}, \quad 1 \leq i \leq n, 1 \leq A \leq k \tag{1.20}
\end{equation*}
$$

Given a Lagrangian function of the form $L=L\left(x^{i}, v_{A}^{i}\right)$ one obtains, by using a variational principle, the generalized Euler-Lagrange equations for $L$ :

$$
\begin{equation*}
\sum_{A=1}^{k} \frac{d}{d t^{A}}\left(\frac{\partial L}{\partial v_{A}^{i}}\right)-\frac{\partial L}{\partial x^{i}}=0, \quad v_{A}^{i}=\frac{\partial x^{i}}{\partial t^{A}} \tag{1.21}
\end{equation*}
$$

Following the ideas of Günther [3], we will describe the above equations (1.21) in terms of the geometry of $k$-tangent structures. In classical mechanics the symplectic structure of Hamiltonian theory and the tangent structure of Lagrangian theory play complementary roles $[13,14,15,16,17]$. Also, that the $k$-symplectic structures and the $k$-tangent structures play similarly complementary roles.

First of all, we note that such a $L$ can be considered as a function $L: T_{k}^{1} M \rightarrow R$ with $M$ a manifold with local coordinates $\left(x^{i}\right)$. Next, we construct a $k$-symplectic structure on the manifold $T_{k}^{1} M$, using its canonical $k$-tangent structure for each $1 \leq A \leq k$.

Let us consider the 1 -forms $\left(\beta_{L}\right)_{A}=d L \circ J^{A}, 1 \leq A \leq k$. In a local coordinate $\operatorname{system}\left(x^{i}, v_{A}^{i}\right)$ we have

$$
\begin{equation*}
\left(\theta_{L}\right)_{A}=\frac{\partial L}{\partial v_{A}^{i}} d x^{i}, 1 \leq A \leq k \tag{1.22}
\end{equation*}
$$

Definition 1.17 A Lagrangian $L$ is called regular if and only if

$$
\begin{equation*}
\operatorname{det}\left(\frac{\partial^{2} L}{\partial v_{A}^{i} \partial v_{B}^{j}}\right) \neq 0, \quad 1 \leq i, j, \leq n, \quad 1 \leq A, B \leq k \tag{1.23}
\end{equation*}
$$

By introducing the following 2-forms $\left(\omega_{L}\right)_{A}=-d\left(\theta_{L}\right)_{A}, 1 \leq A \leq k$, one can easily prove the following.
Proposition $1.18 L: T_{k}^{1} M \longrightarrow \mathbf{R}$ is a regular Lagrangian if and only if $\left(\left(\omega_{L}\right)_{1}, \ldots,\left(\omega_{L}\right)_{k}, V\right)$ is a $k$-symplectic structure on $T_{k}^{1} M$, where $V$ denotes the vertical distribution of $\tau: T_{k}^{1} M \rightarrow M$.

Let $L: T_{k}^{1} M \longrightarrow \mathbf{R}$ be a regular Lagrangian and let us consider the $k$-symplectic structure
$\left(\left(\omega_{L}\right)_{1}, \ldots,\left(\omega_{L}\right)_{k}, V\right)$ on $T_{k}^{1} M$ defined by $L$. Let $\Omega_{L}^{\sharp}$ be the morphism defined by this $k$-symplectic structure

$$
\Omega_{L}^{\sharp}: T_{k}^{1}\left(T_{k}^{1} M\right) \longrightarrow T^{*}\left(T_{k}^{1} M\right)
$$

Thus, we can set the following equation:

$$
\begin{equation*}
\Omega_{L}^{\sharp}\left(X_{1}, \ldots, X_{k}\right)=d E_{L}, \tag{1.24}
\end{equation*}
$$

where $E_{L}=C(L)-L$, and where $C$ is the canonical vector field of the vector bundle $\tau: T_{k}^{1} M \rightarrow M$.

Proposition 1.19 Let $L$ be a regular Lagrangian. If $\xi=\left(\xi_{1}, \cdots, \xi_{k}\right)$ is a solution of (1.24) then it is a SOPDE. In addition if $\xi$ is integrable then the solutions of $\xi$ are solutions of the Euler-Lagrange equations (1.21).

Proof It is a direct computation in local coordinates using (1.13), (1.14), (1.22) and (1.23).

## 2 SOPDEs and nonlinear connections on $T_{k}^{1} M$

A nonlinear connection on the vector fiber bundle $\tau: T_{k}^{1} M \rightarrow M$ is a distribution $H: u \rightarrow H_{u} E$ on $E=T_{k}^{1} M$ which is supplementary to the vertical distribution $V: u \rightarrow V_{u} E$, where $V E=\operatorname{Ker}(\tau)_{*}$. The horizontal distribution $u \rightarrow H_{u} E$ may be given by $n$ local vector fields

$$
\begin{equation*}
\frac{\delta}{\delta x^{i}}=\frac{\partial}{\partial x^{i}}-N_{A i}^{j} \frac{\partial}{\partial v_{A}^{j}}, \quad i=1, \ldots, n \tag{2.1}
\end{equation*}
$$

Functions $N_{A i}^{j}$ are called the local coefficients (or components) of the nonlinear connection, which, for this reason, will be denoted by $N .\left\{\frac{\delta}{\delta x^{i}}(u), \frac{\partial}{\delta v_{A}^{i}}(u)\right\}$ is an adapted (local) basis to the decomposition $T_{u} E=H_{u} E \oplus V_{u} E$, called the adapted basis of the nonlinear connection. Its dual basis is

$$
\begin{equation*}
\left\{d x^{i}, \delta v_{A}^{i}=d v_{A}^{i}+N_{A, j}^{i} d x^{j}\right\} \tag{2.2}
\end{equation*}
$$

Follow [24] we have $\left[\delta / \delta x^{j}, \delta / \delta x^{k}\right]=R_{A j k}^{i} \partial / \partial v_{A}^{i}$ and so, the distribution horizontal $H$ is integrable if and only if $R_{A j k}^{i}=0$.

From the general theory we have ([23], [24]):
Theorem 2.1 a) The transformation rule, with respect to a change of coordinates, of the components $N_{A j}^{i}\left(x^{k}, v_{B}^{k}\right)$ of a nonlinear connection $N$ are

$$
\begin{equation*}
\frac{\partial \tilde{x}^{k}}{\partial x^{i}} \tilde{N}_{A k}^{j}=N_{A i}^{h} \frac{\partial \tilde{x}^{j}}{\partial x^{h}}-\frac{\partial^{2} \tilde{x}^{j}}{\partial x^{i} \partial x^{h}} v_{A}^{h} \tag{2.3}
\end{equation*}
$$

b) Conversely, if the functions $N_{A j}^{i}\left(x^{k}, v_{B}^{k}\right)$ are given on every domain of local chart on $T_{k}^{1} M$, so that (2.3) holds with respect to (1.3), then there exists a unique nonlinear connection $N$ on whose coefficients are the given functions $N_{A j}^{i}\left(x^{k}, v_{B}^{k}\right)$.
Using this theorem, we obtain the following results:
Theorem 2.2 Let $\xi \equiv\left(\xi_{1}, \ldots, \xi_{k}\right)$ be a SOPDE. Then the local functions on $T_{k}^{1} M$ defined by

$$
\begin{equation*}
N_{A j}^{i}=-\frac{1}{k+1} \sum_{B=1}^{k} \frac{\partial\left(\xi_{A}\right)_{B}^{i}}{\partial v_{B}^{j}} \tag{2.4}
\end{equation*}
$$

are the local components of a nonlinear connection.
Next, the nonlinear connection $N$ given by (2.4) is called the nonlinear connection associated to the SOPDE $\xi$.

Let us denote $\stackrel{1}{\xi}=\xi$ and

$$
\begin{equation*}
\stackrel{1}{N}_{A j}^{i}=-\frac{1}{k+1} \sum_{B=1}^{k} \frac{\partial\left(\stackrel{1}{( }_{A}\right)_{B}^{i}}{\partial v_{B}^{j}} \tag{2.5}
\end{equation*}
$$

Then $\stackrel{1}{N}$ is the nonlinear connection associated to the SOPDE $\stackrel{1}{\xi}$.
Let us consider the functions

$$
\begin{equation*}
\left(\stackrel{2}{\xi}_{A}\right)_{B}^{i}=\frac{1}{2} C\left(\left(\stackrel{1}{\xi}_{A}\right)_{B}^{i}=\frac{1}{2} \sum_{C=1}^{k} \sum_{j=1}^{n} \frac{\partial\left(\stackrel{1}{\xi}_{A}\right)_{B}^{i}}{\partial v_{C}^{j}} v_{C}^{j}\right. \tag{2.6}
\end{equation*}
$$

on every domain of a local chart. These functions change following (1.15) and therefore $\stackrel{2}{\xi}$ is a SOPDE. Using (2.4) it obtain the nonlinear connection $\stackrel{2}{N}$ associated to the SOPDE $\stackrel{2}{\xi}$.

So, starting with a SOPDE on $T_{k}^{1} M$ we obtain a sequence of $\operatorname{SOPDE}\binom{m}{\xi}_{m \geq 1}$ and a sequence of nonlinear connection $(\stackrel{m}{N})_{m \geq 1}$, for which we obtain the next results.

Theorem 2.3 The following assertions are equivalent:
i) the SOPDE $\stackrel{1}{\xi}$ is homogeneous (spray).
ii) $\stackrel{1}{\xi}=\stackrel{2}{\xi}$.

Theorem 2.4 If $\stackrel{1}{\xi}$ is a homogeneous SOPDE then the components $N_{A j}^{i}$ of the associated nonlinear connection $N$ are 1-homogeneous functions on the fibres of $T_{k}^{1} M$.

We remark that the converse of this theorem is not true, in generally (see [25] for the case $k=1$ ).

Theorem $2.5 \stackrel{1}{N}=\stackrel{2}{N}$ if and only if the components $N_{A j}^{i}$ are 1-homogeneous functions on the fibres of $T_{k}^{1} M$.

Corollary 2.6 If $\stackrel{1}{N}_{\text {Aj }}^{i}$ are 1-homogeneous functions on the fibres of $T_{k}^{1} M$, then the sequence $(\stackrel{m}{N})_{m \geq 1}$ is constant, but we can not deduce that $\binom{m}{\xi}_{m \geq 1}$ is constant.

Corollary 2.7 If $\stackrel{1}{\xi}$ is a homogeneous SOPDE, then the sequences $\binom{m}{\xi}_{m \geq 1}$ and $(\stackrel{m}{N})_{m \geq 1}$ are constants.

### 2.1 SOPDEs and nonlinear connections associated to a regular Lagrangian on $T_{k}^{1} M$

Using the previously notations, if we have a regular Lagrangian on $T_{k}^{1} M$, following the proposition 1.19 and the generalized Euler-Lagrange equations (1.21), we deduce the next results.

Proposition 2.8 Let $L$ be a regular Lagrangian. The $k$-vector field $\left(\left(\xi_{L}\right)_{1}, \ldots,\left(\xi_{L}\right)_{k}\right)$ locally given by

$$
\begin{equation*}
\left(\left(\xi_{L}\right)_{A}\right)_{B}^{j}=-\frac{1}{2} g_{A B}^{j i}\left(\frac{\partial^{2} L}{\partial v_{C}^{i} \partial x^{h}} v_{C}^{h}-\frac{\partial L}{\partial x^{i}}\right), \quad g_{i j}^{A B}=\frac{1}{2} \frac{\partial^{2} L}{\partial v_{A}^{i} \partial v_{B}^{j}}, \tag{2.7}
\end{equation*}
$$

where $g_{A B}^{j i}$ is the inverse of the matrix $g_{j i}^{A B}$, is a SOPDE and it is a solution of (1.24). Moreover, if $\xi_{L}$ is integrable, then its solutions are solutions of the generalized EulerLagrange equations (1.21).

Definition 2.9 The $\operatorname{SOPDE}\left(\left(\xi_{L}\right)_{1}, \ldots,\left(\xi_{L}\right)_{k}\right)$ from the above proposition will be called the canonical SOPDE associated to the regular Lagrangian $L$.

From (2.4) it results that there exists a nonlinear connection $N_{L}$ on $T_{k}^{1} M$ associated to $\left(\left(\xi_{L}\right)_{1}, \ldots,\left(\xi_{L}\right)_{k}\right)$ which depends only by $L . N_{L}$ will be called the canonical nonlinear connection of regular Lagrangian $L$.

Now we consider the particular case of a regular Lagrangian, $L=F^{2}$, where $F$ is a Finslerian function on $T_{k}^{1} M$ which is defined below.

Definition 2.10 A function $F: T_{k}^{1} M \rightarrow \mathbf{R}, F=F\left(x^{i}, v_{A}^{i}\right)$, which is differentiable on $\widetilde{T_{k}^{1} M}$ and continuous on null section of the projection $\tau$ is called Finslerian on $T_{k}^{1} M$ if
i) $F>0$ on $T_{k}^{1} M$,
ii) $F$ is 1-homogeneous on the fibres of $T_{k}^{1} M$,
iii) the matrix with the elements

$$
\begin{equation*}
g_{i j}^{A B}=\frac{1}{2} \frac{\partial^{2} F^{2}}{\partial v_{A}^{i} \partial v_{B}^{j}} \tag{2.8}
\end{equation*}
$$

is positively defined on $\widetilde{T_{k}^{1} M}$.
From Proposition $1.14,(1.17),(2.7)$ and ii), iii) we deduce that if $L=F^{2}$, then the canonical SOPDE $\xi_{F^{2}}$ is homogeneous (spray).

Theorem 2.11 If the Lagrangian is the square of a Finslerian on $T_{k}^{1} M$ then
a) the sequence $\binom{m}{\xi}_{m \geq 1}$ is constant, $\stackrel{1}{\xi}=\xi_{F^{2}}$,
b) the sequence $(\stackrel{m}{N})_{m \geq 1}$ is constant, $\stackrel{1}{N}=N_{F^{2}}$.

Example: Let $L\left(x^{h}, v_{D}^{h}\right)=\frac{1}{3} \alpha_{i j k}^{A B C}\left(x^{h}\right) v_{A}^{i} v_{B}^{j} v_{C}^{k}$, where $\alpha_{i j k}^{A B C}\left(x^{1}, \ldots, x^{n}\right)$ is totally symmetric and the matrix with the elements $g_{i j}^{A B}=\frac{1}{2} \frac{\partial^{2} L}{\partial v_{A}^{i} \partial v_{B}^{j}}=\alpha_{i j k}^{A B C} v_{C}^{k}$ is of rank $n k$ on $T_{k}^{1} M$. Then $L$ is a regular Lagrangian with $g_{i j}^{A B} 1$-homogeneous functions, $g_{A B}^{i j}$ homogeneous functions of degree -1 , according with (1.17). Using the relation (2.5) and Theorem 2.4, the components of the canonical nonlinear connection $N_{L}$ are homogeneous of degree 1. So, the sequence $(\stackrel{m}{N})_{m \geq 1}$ is constant, but we dont tell that the sequence $\binom{m}{\xi}_{m \geq 1}$ has the same property. This is an example of a regular Lagrangian on $T_{k}^{1} M$ which in not equal with the square of a Finslerian.

### 2.2 Other type of relationship between SOPDEs and nonlinear connections on $T_{k}^{1} M$

Let $\xi \equiv\left(\xi_{1}, \ldots, \xi_{k}\right)$ be a SOPDE on $T_{k}^{1} M$. Then the functions given by

$$
\begin{equation*}
N_{A l}^{j}=-\frac{1}{2} \frac{\partial\left(\xi_{A}\right)_{A}^{j}}{\partial v_{A}^{l}} \tag{2.9}
\end{equation*}
$$

are the components of a nonlinear connection $N$ on $T_{k}^{1} M$, determined by $\xi$. Indeed $N_{A l}^{j}$ verify (2.3), under (1.3).

From (2.9) and (1.15), we deduce that the functions

$$
\begin{equation*}
\left(\xi_{B}^{*}\right)_{A}^{j}=-N_{A l}^{j} v_{B}^{l} \tag{2.10}
\end{equation*}
$$

are the components of a SOPDE $\stackrel{*}{\xi}$ on $T_{k}^{1} M$.
We can continue and we will obtain a nonlinear connection $\stackrel{*}{N}$ from $\stackrel{*}{\xi}$ and so on.
By this procedure one can obtain again a sequence of SOPDEs and a sequence of nonlinear connection on $T_{k}^{1} M$, starting with a given SOPDE (or starting with a given nonlinear connection). But, now we have other results about this sequences.

Proposition 2.12 i) $\stackrel{*}{\xi}=\xi$ if and only if $C_{B}^{A}\left(\left(\xi_{A}\right)_{A}^{j}\right)=2\left(\xi_{B}\right)_{A}^{j}$ for all $j=1, \ldots, n$, $A, B=1, \ldots, k$;
ii) $\stackrel{*}{N}=N$ if and only if $C_{A}\left(N_{A l}^{j}\right)=N_{A l}^{j}$ for all $j, l=1, \ldots, n, A=1, \ldots, k$.

Remark $\quad \stackrel{*}{\xi}=\xi$ implies $\stackrel{*}{N}=N$, but the converse is not true. Only $\stackrel{*}{N}=N$ implies $\stackrel{*}{\xi}=\stackrel{*}{\xi}$

Now, if we done the next definition, then we get an interesting result.
Definition 2.13 $A \operatorname{map} \Phi: \mathbf{R}^{k} \mapsto M$ is said to be horizontal for a nonlinear connection $N$ on $T_{k}^{1} M$ if and only if

$$
\begin{equation*}
\left(\Phi^{(1)}\right)_{*}(t)\left(\frac{\partial}{\partial t^{A}}(t)\right) \in H\left(T_{\Phi^{(1)}(t)}\left(T_{k}^{1} M\right)\right), \forall A=1, \ldots, k \tag{2.11}
\end{equation*}
$$

Theorem 2.14 If the SOPDE $\xi$ associated to the nonlinear connection $N,\left(\xi_{A}\right)_{B}^{j}=$ $-N_{B l}^{j} v_{A}^{l}$ is integrable, then the solutions of $\xi$ are horizontals with respect to $N$.

Proof Since, for all $A=1, \ldots, k$,

$$
\left(\Phi^{(1)}\right)_{*}(t)\left(\frac{\partial}{\partial t^{A}}(t)\right)=\frac{\partial \Phi^{i}}{\partial t^{A}}(t) \frac{\delta}{\delta x^{i}}+\left(\frac{\partial^{2} \Phi^{j}}{\partial t^{A} \partial t^{B}}(t)+N_{B i}^{j} \frac{\partial \Phi^{i}}{\partial t^{A}}(t)\right) \frac{\partial}{\partial v_{B}^{j}}
$$

and taking into account the relationship between $\xi$ and $N$, we obtain that $\Phi$ is a solution of $\xi$ iff (2.11) holds, that is $\Phi$ is horizontal of $N$.

Remark Beginning with a SOPDE $\xi$ we can construct directly a sequence of SOPDEs:

$$
\begin{equation*}
\left(\stackrel{*}{*}_{A}\right)_{B}^{j}=\frac{1}{2} C_{A}^{B}\left(\left(\xi_{B}\right)_{B}^{j}\right)=\frac{1}{2} \frac{\partial\left(\xi_{B}\right)_{B}^{j}}{\partial v_{B}^{h}} v_{A}^{h} \tag{2.12}
\end{equation*}
$$

After this, we can obtain for each SOPDE of sequence a nonlinear connection by (2.9). It is the same construction as that presented above in (2.9) and (2.10).

Finally, starting with a SOPDE $\xi$ we can construct two sequence of SOPDEs. The first sequence by the method from (2.6) and the second by the method presented above (2.12). For each sequence of SOPDE we can obtain two sequences of nonlinear connections. The first sequence is obtained by the method from (2.4) and the second by the above method (2.9).

Acknowledgement. This work has been partially supported by a Grant CNCSIS MEN 17C/27661 (81/2005).

## References

[1] A. Awane, $k$-symplectic structures, J. Math. Phys. 33 (1992), 4046-4052.
[2] A. Awane, G-spaces $k$-symplectic homogènes, J. Geom. Phys. 13 (1994), 139-157.
[3] Ch. Günther, The polysymplectic Hamiltonian formalism in field theory and calculus of variations I: The local case, J. Diff. Geom. 25 (1987), 23-53.
[4] M. Bruckheimer, Ph.D. dissertation, University of Southampton, 1960
[5] R.S. Clark, M. Bruckheimer, Sur les structures presque tangents, C. R. Acad. Sci. Paris Sér . I Math. 251 (1960), 627-629.
[6] H.A. Eliopoulos, Structures presque tangents sur les variétés différentiables, C. R. Acad. Sci. Paris Sér . I Math. 255 (1962), 1563-1565.
[7] M. de León, I. Méndez, M. Salgado, p-almost tangent structures, Rend. Circ. Mat. Palermo Serie II XXXVII (1988), 282-294.
[8] M. de León, I. Méndez, M. Salgado, Integrable p-almost tangent structures and tangent bundles of $p^{1}$-velocities, Acta Math. Hungar. Vol. 58 (1-2) (1991), 45-54.
[9] M. de León, E. Merino, J. Oubina, P. Rodriguez, M. Salgado, Hamiltonian systems on $k$-cosymplectic manifolds, J. Math. Phys. 39 (1998), 876-893.
[10] F. Brickell, R.S. Clark, Integrable almost tangent structures, J. Diff. Geom. 9 (1974), 557-563.
[11] R. S. Clark, D.S. Goel, On the geometry of an almost tangent structure, Tensor (N. S.) 24 (1972), 243-252.
[12] M. Crampin, G. Thompson, Affine bundles and integrable almost tangent structures, Math. Proc. Cambridge Philos. Soc. 101 (1987), 61-67.
[13] M. Crampin, Tangent bundle geometry for Lagrangian dynamics, J. Phys. A: Math. Gen. 16 (1983), 3755-3772.
[14] M. Crampin, Defining Euler-Lagrange fields in terms of almost tangent structures, Phys. Lett. A 95 (1983), 466-468.
[15] J. Grifone, Structure presque-tangente et connexions, I, Ann. Inst. Fourier 22 (1972), 287-334.
[16] J. Grifone, Structure presque-tangente et connexions, II, Ann. Inst. Fourier 22 (1972), 291-338.
[17] J. Klein, Espaces variationelles et mécanique, Ann. Inst. Fourier 12 (1962),1-124.
[18] G. Thompson, U. Schwardmann, Almost tangent and cotangent structures in the large, Trans. Amer. Math. Soc. 327 (1991), 313-328.
[19] A. Morimoto, Liftings of some types of tensor fields and connections to tangent $p^{r}$-velocities, Nagoya Math. J. 40 (1970), 13-31.
[20] A. Echevarría-Enríquez, M.C. Muñoz-Lecanda, N. Román-Roy, Multivector fields and connections: Setting Lagrangian equations in field theories, J. Math. Phys. 39 (1998), 4578-4603.
[21] A. Echevarría-Enríquez, M.C. Muñoz-Lecanda, N. Román-Roy, Multivector field formulation of Hamiltonian field theories: equations and symmetries, J. Phys. A: Math. Gen. 32 (1999), 8461-8484.
[22] H. Rund, The Hamilton-Jacobi Theory in the Calculus of Variations, Reprinted Edition, Robert E. Krieger Publishing Co. Inc., Huntington, New York, 1973.
[23] R. Miron, M.S. Kirkovits, M. Anastasiei, A Geometrical Model for Variational Problems of Multiple Integrals, Proceedings of the Conference on Differential Geometry and Applications, June 26 - July 3, 1988, Dubrovnik, Yugoslavia, pp. 209-216.
[24] R. Miron, M. Anastasiei, The Geometry of Lagrange Spaces: Theory and Applications, FTPH, no. 59, Kluwer Academic Publishers, 1994.
[25] G. Munteanu, G. Pitiş, On the Nonlinear Connection of a Spray, Report to Math. Physics, vol. 50, no. 1, 2002, pp. 41-47.
[26] F. Munteanu, On the Semispray of Nonlinear Connections in Rheonomic Lagrange Geometry, Proc. Conf. Finsler-Lagrange Geometry, Iassy, Romania, August 26-31, 2001, "Al.I. Cuza" University of Iassy, ed. M. Anastasiei, Kluwer Academic Publishers, 2003, pp. 129-137.
[27] F. Munteanu, On the Nonlinear Connections of a $k$-Semispray in $k$-Tangent Bundle Geometry, Preprint.
[28] F. Munteanu, Spaţii Lagrange şi spaţii Hamilton. Aplicaţii la sisteme dinamice Hamiltoniene, Ph. D. Thesis, Craiova, Romania, 2002 (in Romanian).
[29] F. Munteanu, M. Salgado, A.M. Rey, The Günther's formalism in classical field theory: momentum map and reduction, J. Math. Phys. 45 (2004), 1730-1751.

Author's address:
Florian Munteanu
Department of Applied Mathematics, University of Craiova,
Al. I. Cuza st., no. 13, 200585-Craiova, Romania,
email: munteanufm@central.ucv.ro

